

*ELW. AND ALTERNATIVES WG
SUMMARY REPORT*

Michael Spira (PSI)

other convenors: A. Denner, K. Mönig, T. Ohl, G. Pasztor

- Signatures of new ρ -resonances from Ivan Melo
strong EWSB in $e^+e^- \rightarrow \nu\bar{\nu}t\bar{t}$
- Strong electroweak symmetry breaking Predrag Krstonosic
from $ee \rightarrow VV\nu\nu$ (presented by K. Mönig)
- Triple gauge couplings in $\gamma\gamma$ collisions Jadranka Sekaric
- Four-fermion production at future photon colliders Axel Bredenstein
- Precision measurements of beam polarization at the LC (The case of single- W production) Filip Franco-Sollova
- Two-loop Sudakov logarithms in electro- Bernd Feucht
weak processes
- Electroweak precision observables in the Sven Heinemeyer
MSSM with non minimal flavor violation

ECFA - Durham

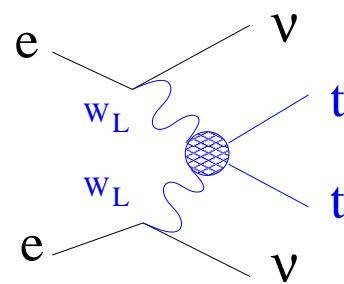
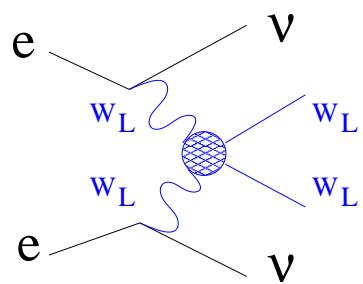
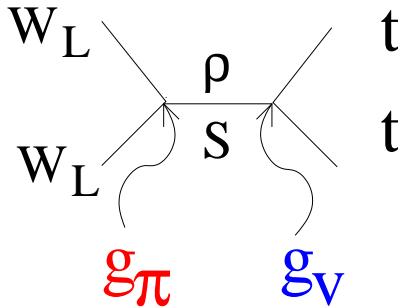
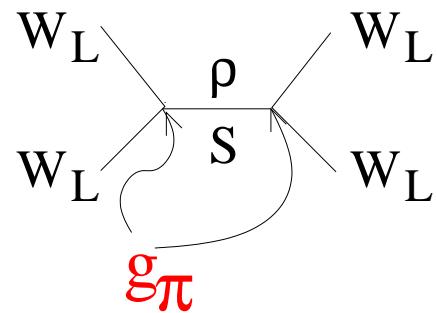
Sep 2004

Signatures of new vector resonances from strong EWSB in
 $e^+e^- \rightarrow \nu\bar{\nu}t\bar{t}$

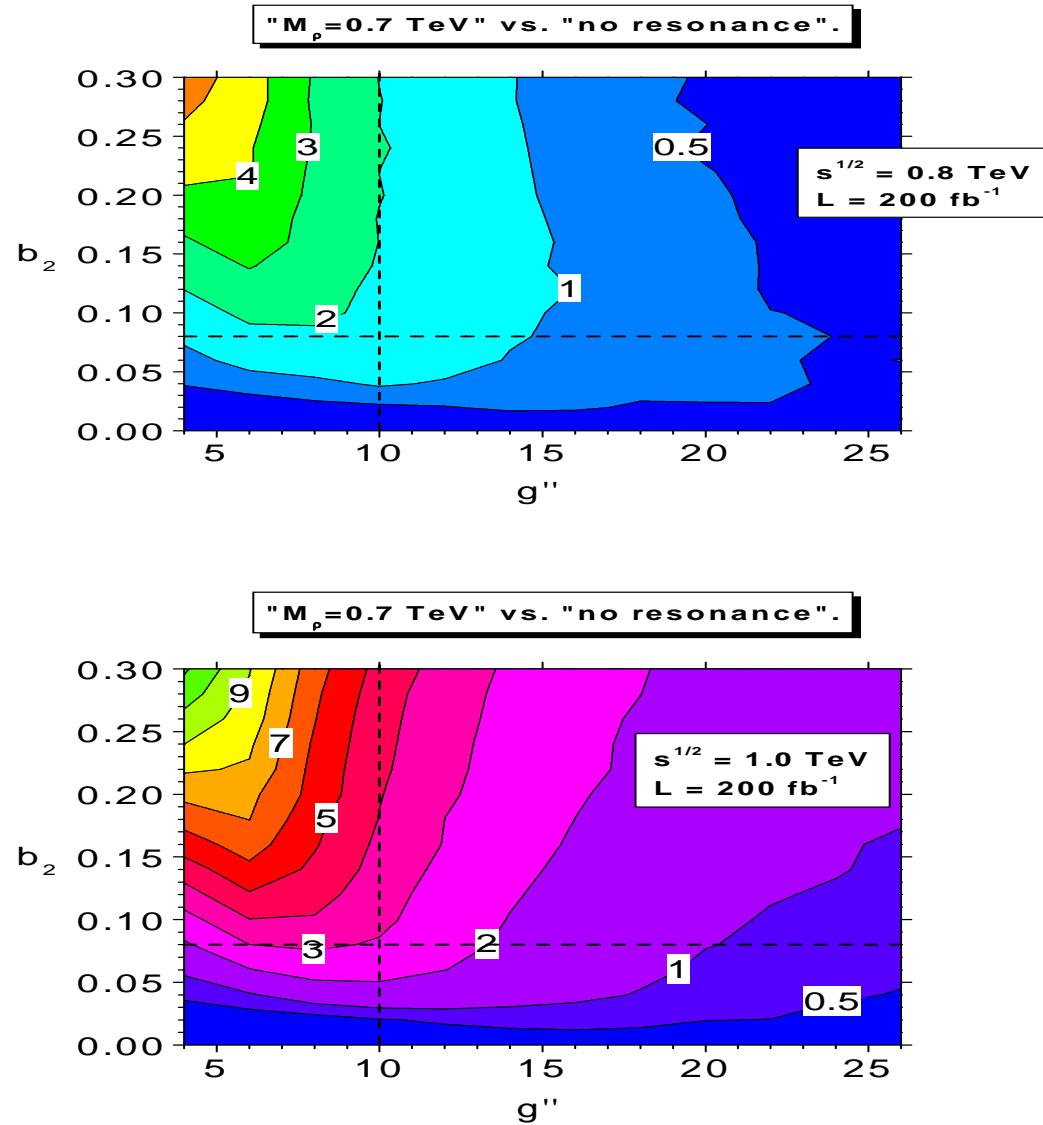
Ivan Melo

D. Bruncko (IEP SAS Kosice)
M. Gintner (University of Zilina)
I.Melo (University of Zilina)

$W_L W_L \rightarrow t\bar{t}$ scattering



$$R = \frac{|N(\rho) - N(\text{no resonance})|}{\sqrt{N(\text{Background}) + N(\text{no resonance})}}$$



Conclusions

ρ in $e^+e^- \rightarrow \nu\bar{\nu}t\bar{t}$ (Pythia and CompHEP)

- agreement within 10 %
- no t decays

$\sigma(0.8 \text{ TeV}) = 0.20 \text{ (0.13) fb}$

$\sigma(1.0 \text{ TeV}) = 0.16 \text{ (0.035) fb}$

R (ρ vs no resonance) values up to 8

- optimize cuts
- finalize analysis for all models considered
- $e^+e^- \rightarrow \nu\bar{\nu}W^+W^-$

Strong Electroweak Symmetry Breaking from $e^- e^- \rightarrow \nu \nu V$

K. Mönig , P. Krstonosic
DESY - Zeuthen



Introduction

- In absence of light Higgs interaction among gauge bosons becomes strong at high energy
- Effective Lagrangian contains five CP conserving operators that can be directly probed in weak boson scattering

$$J_{\text{scattering}} = \frac{\alpha_4}{6 \pi^2} r(V_\mu V_\nu) r(V^\mu V^\nu) \quad L_5 = \frac{\alpha_5}{6 \pi^2} r(V_\mu V^\mu) r(V_\nu V^\nu) \quad \left. \right\} \begin{matrix} SU(2)_C \\ \text{L} \end{matrix}$$

$$L_6 = \frac{\alpha_6}{6 \pi^2} r(V_\mu V_\nu) r(Y^{-\mu}) r(Y^{-\nu}) \quad L_7 = \frac{\alpha_7}{6 \pi^2} r(V_\mu V^\mu) r(Y_\nu) r(Y^\nu) \quad L_0 = \frac{\alpha_0}{6 \pi^2} (r(Y_\mu) r(Y_\nu))^2 \quad \left. \right\} \begin{matrix} SU(2)_C \\ \text{R} \end{matrix}$$

- Couplings are related to the scale of “new” physics $\frac{\alpha_i}{6 \pi^2} = \left(\frac{v}{b \lambda_i^*} \right)^2$

,

and it's necessary to do the multidimensional analysis in order to obtain the information on scale and dynamics of

Fit

- Binned Maximum Likelihood fit to the expected distributions for a

$$L_{\text{SM sample}} = \sum_{ijkl} N_{ijkl}^M - \sum_{ijkl} N_{ijkl}^B \quad (\text{all quartic couplings equal zero})$$

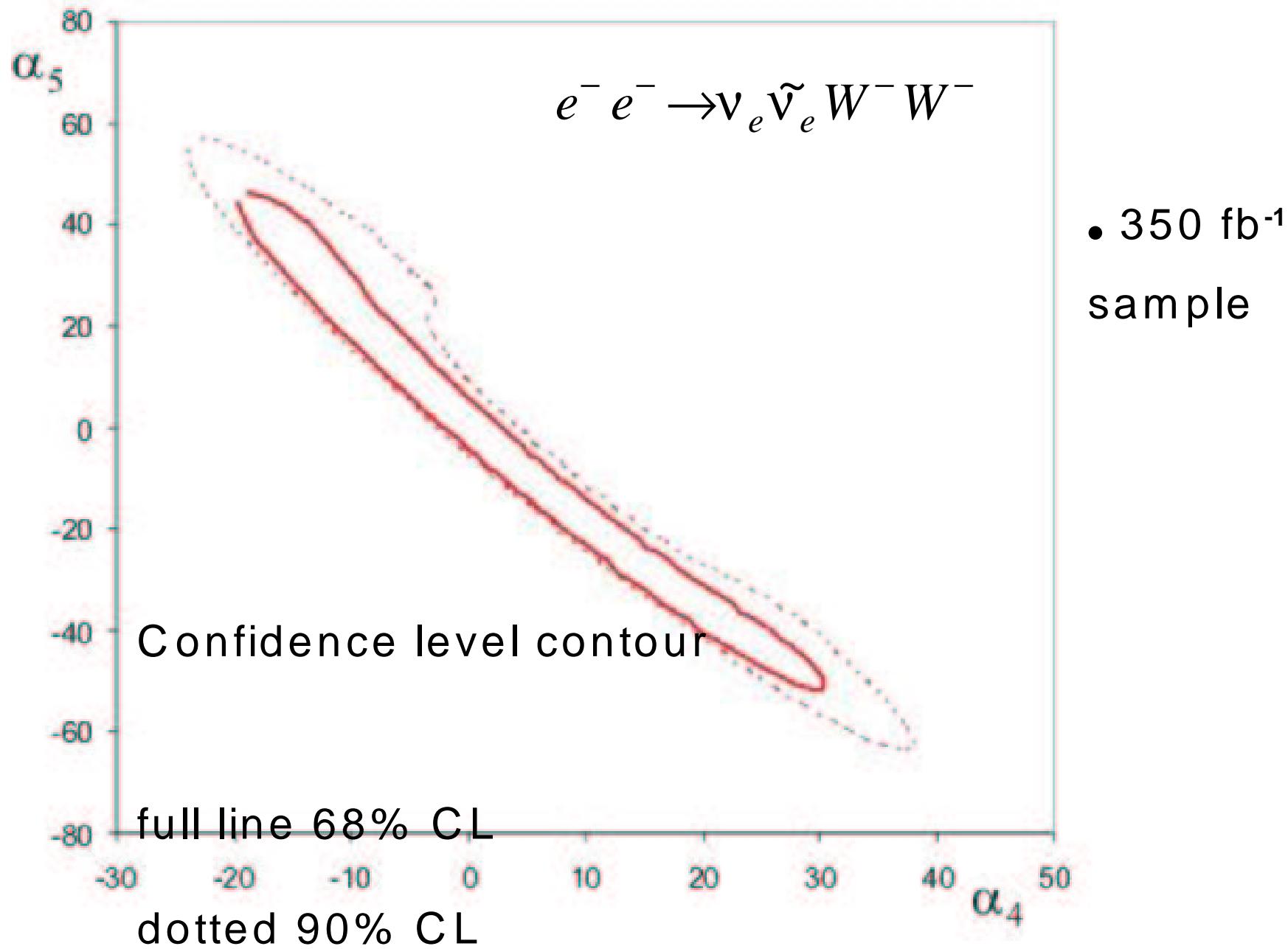
$$N_{ijkl}^M = \sum_B 1 + A \cdot \alpha_4 + B \cdot \alpha_4^2 + C \cdot \alpha_5 + D \cdot \alpha_5^2 + E \cdot \alpha_4 \cdot \alpha_5$$

- Coefficients obtained by recalculating matrix elements in 5 points in α_4 α_5 space for each event and solving system of equations:
- $$R_i = 1 + A \cdot \alpha_{4i} + B \cdot \alpha_{4i}^2 + C \cdot \alpha_{5i} + D \cdot \alpha_{5i}^2 + E \cdot \alpha_{4i} \cdot \alpha_{5i} \quad i = 1, 5$$

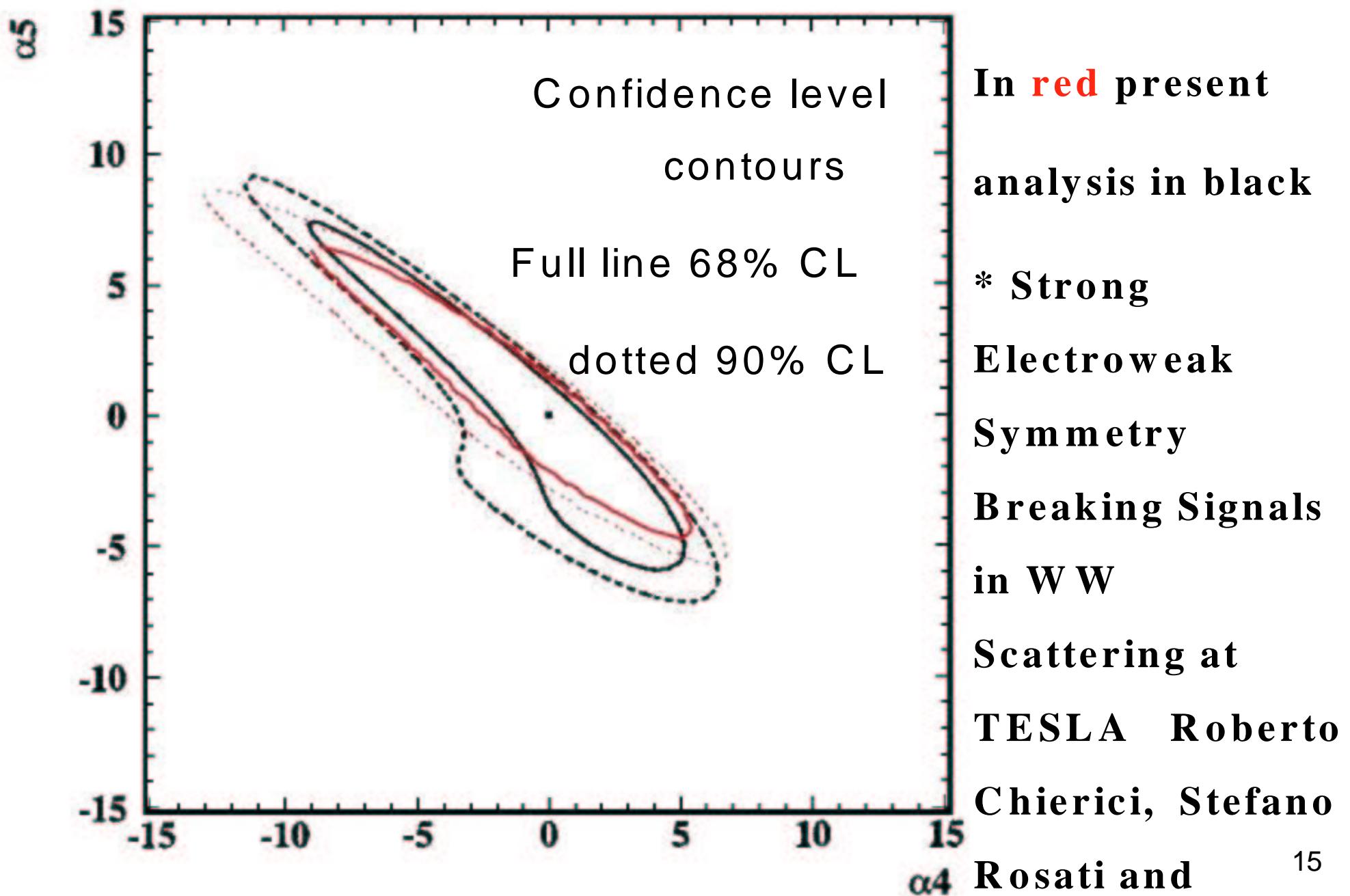
R is the ratio of the matrix element to the SM one.

- After separate analysis for signal processes double counted events were assigned to one or another sample according to the distance from $M(V) + M(V)$ mass ($V = W$ or Z)

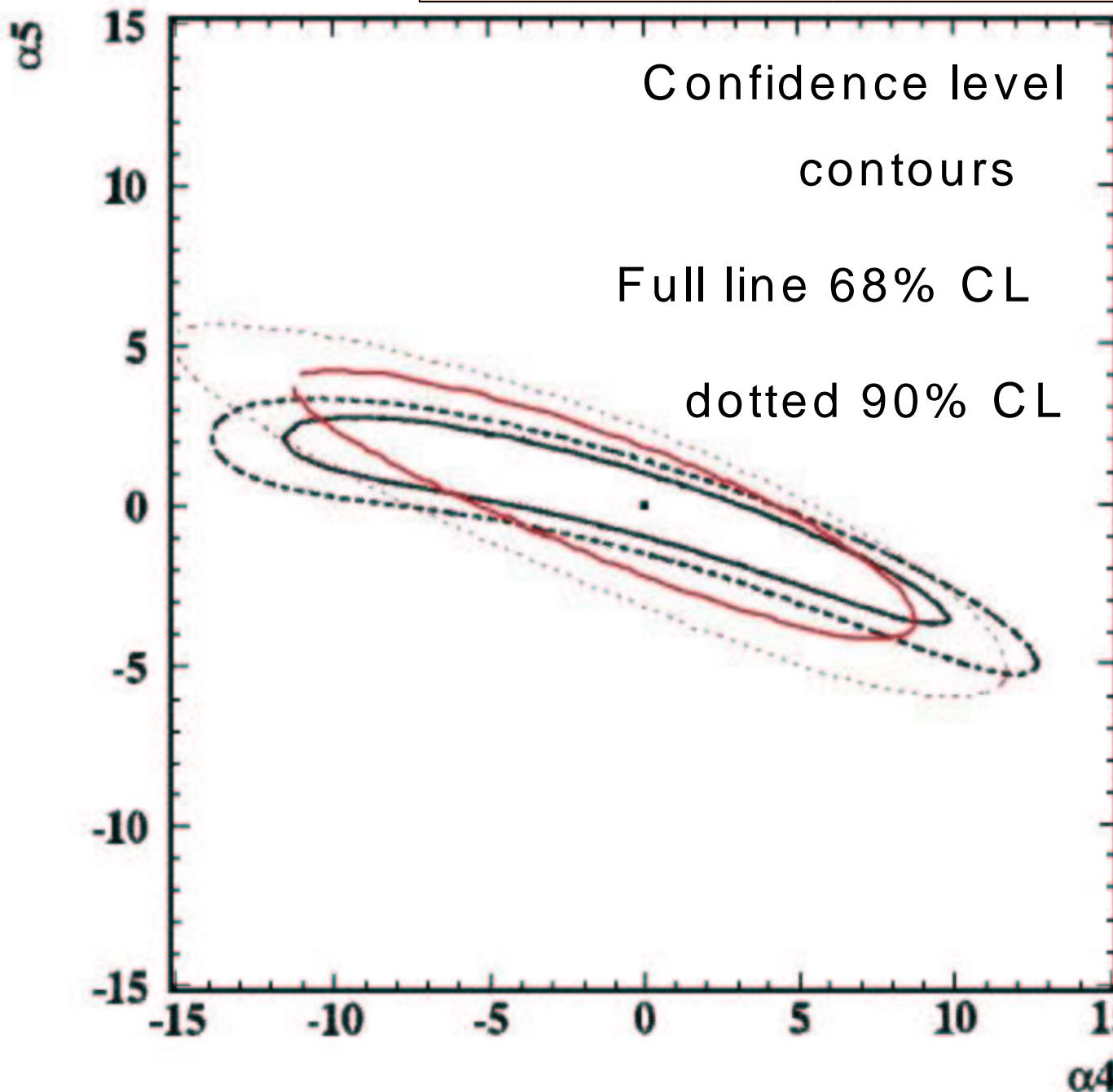
W W result



$W^+ W^-$ result



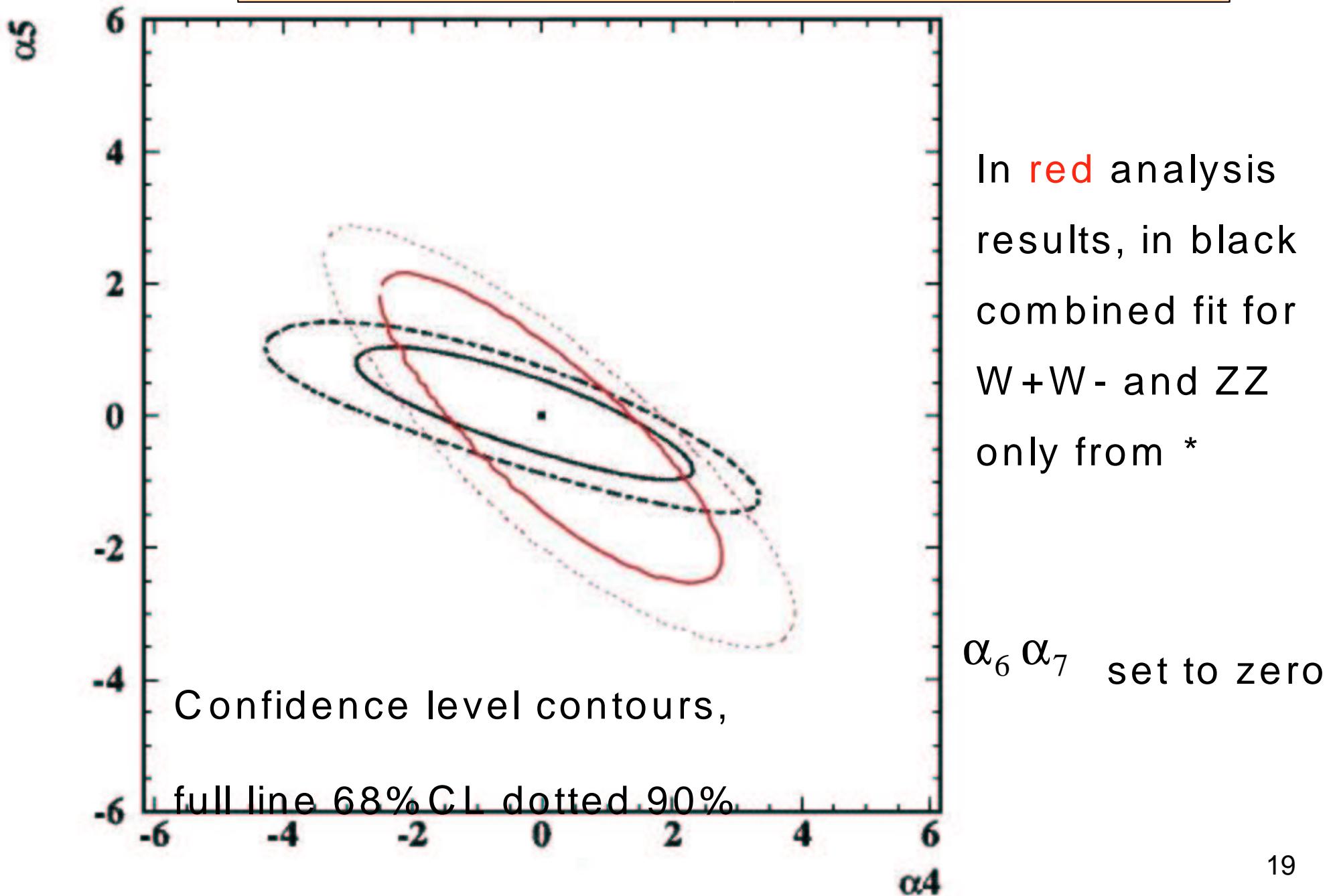
ZZ result



In red present
analysis in black

* Strong
Electroweak
Symmetry
Breaking Signals in
WW Scattering at
TESLA Roberto
Chierici, Stefano
Rosati and Michael
Kobel

e+e- & e-e- processes combined



Future

- improvement in $Z Z$ analysis still possible
- results in easy to use and combine form
- $W Z$ analysis in progress
- combined limits on full parameter set at 800GeV to be obtained soon
- going for 1TeV

WICsta
??-cicaESA

K.Mönig, J.Sekaric

DESY-Zeuthen



Sekaric Jadranka

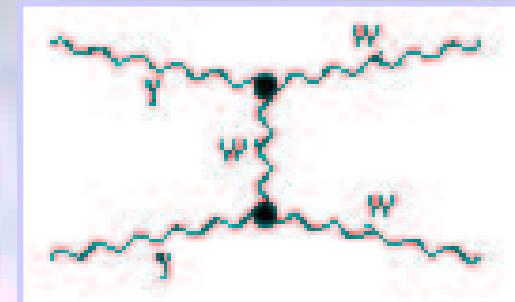
DESY-Zeuthen

Introduction

$$? \rightarrow W^+ W^- \rightarrow q\bar{q}q\bar{q} \quad \sqrt{S} = 7$$

**Dominating diagram for ?? $\rightarrow W^+ W^-$
TGC**

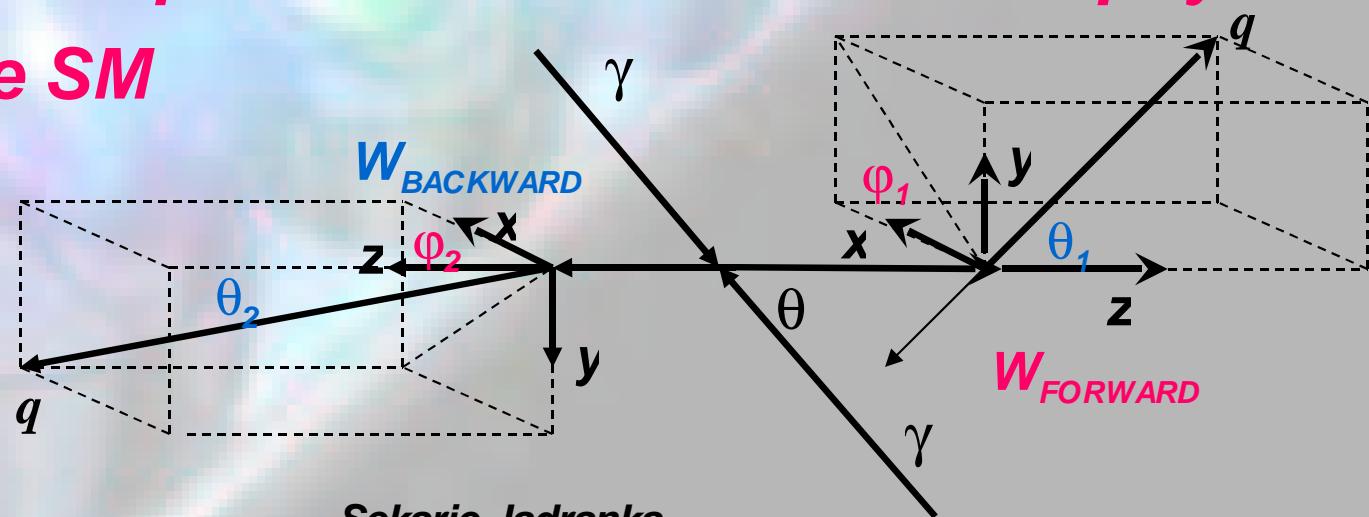
Two couplings, ?, in dim-4 and ?,



**in dim-6 operators
Deviations from SM TGC values \rightarrow test of EW theory,
probe of some possible extensions \rightarrow new physics
beyond the SM**

Ambiguities :

- ($\cos \theta_{1,2}, \sin \theta_{1,2}$) ?
- ($-\cos \theta_{1,2}, \sin \theta_{1,2} + ?$)



Sekaric Jadranka

2 Parameter Fit Results – $J_z=0$

$J_z=0$ <i>3D/5D</i>	$E_{CM} = 400 \text{ GeV}, L = 110 \text{ fb}^{-1}, FBE$	<i>Likelihood (normalized)</i>	
ΔL	1%	0.1%	<i>accurate</i>
$\Delta \kappa_\gamma \cdot 10^{-4}/10^{-4}$	26.0/14.4	6.2/5.4	3.8/2.6
$\Delta \lambda_\gamma \cdot 10^{-4}/10^{-4}$	14.4/3.0	13.7/3.0	13.7/3.0
$\Delta \kappa_\gamma \cdot 10^{-4} + \text{pileup}$	15.4	5.4	2.6
$\Delta \lambda_\gamma \cdot 10^{-4} + \text{pileup}$	3.8	3.8	3.8

3D → 5D decreases the error in κ_γ ; in λ_γ significantly!

Pile-up distorts the ?-distribution → increase of error

in λ_γ for ~20%

2 Parameter Fit Results – $J_z=2$

$J_z=2$ <i>3D/5D</i>	$E_{CM} = 400 \text{ GeV}, L = 110 \text{ fb}^{-1}, FBE$ <i>Likelihood (normalized)</i>		
ΔL	1%	0.1%	<i>accurate</i>
$\Delta \kappa_\gamma \cdot 10^{-4}/10^{-4}$	29.0/19.9	6.2/6.2	3.9/3.8
$\Delta \lambda_\gamma \cdot 10^{-4}/10^{-4}$	1.8/1.5	1.6/1.5	1.6/1.5
$\Delta \kappa_\gamma \cdot 10^{-4} + \text{pileup}$	23.2	6.3	3.9
$\Delta \lambda_\gamma \cdot 10^{-4} + \text{pileup}$	2.0	2.0	2.0

3D → 5D no influences in the error estimations

Pile-up distorts the χ^2 -distribution → increase of error

in λ_γ for ~25%

Conclusions

- $J_z=0,2$ are sensitive to the anomalous couplings → to the possible scenarios of EWSB
- ? distribution important for $J_z=0$
- Pile-up influences on ?, measurement (~25%)
- Promising channel for ?, - ?, measurements
??, ??, $\sim 10^4$

FUTURE

to include the possible backgrounds, variable beam energy
rejection of bad tracks from pileup

Sekaric Jadranka

Four-fermion production at the $\gamma\gamma$ Collider

Axel Bredenstein

in collaboration with Stefan Dittmaier und Markus Roth

Max-Planck-Institut für Physik, Munich

September 1, 2004

based on Eur. Phys. J. C **36** (2004) 341 [arXiv:hep-ph/0405169]



Contents

Calculation of $\gamma\gamma \rightarrow 4f$ und $\gamma\gamma \rightarrow 4f\gamma$ in lowest order

Construction of a Monte Carlo generator

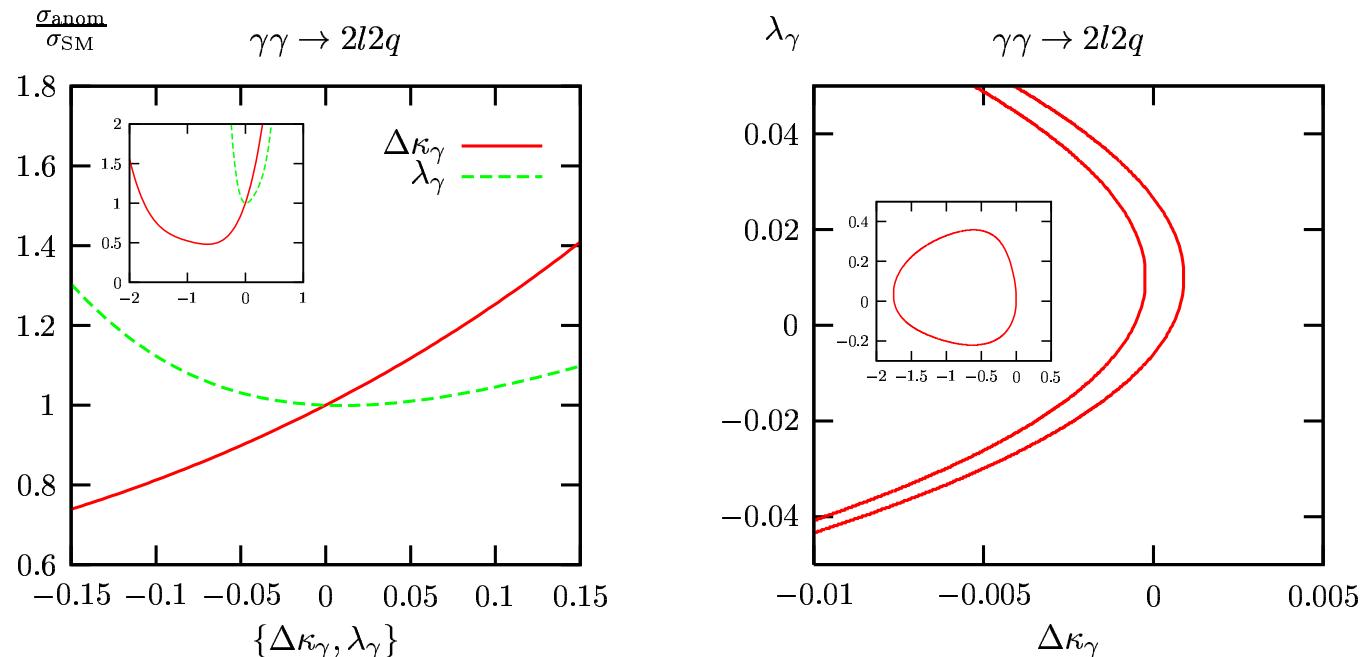
- Motivation
- Calculation of helicity amplitudes
- Phase space and photon spectrum integration
- Anomalous couplings
- Effective Higgs coupling
- Finite gauge-boson width
- Double-pole approximation



Anomalous triple couplings

$\gamma\gamma \rightarrow 4f$ all semi-leptonic final states photon spectrum included

$\sqrt{s_{ee}} = 500 \text{ GeV}$ $\int L dt = 100 \text{ fb}^{-1}$ $\chi^2 = 1$ $\chi^2 \equiv \frac{(N(a_i) - N_{\text{SM}})^2}{N_{\text{SM}}}$



→ large interference with SM amplitude

expected limits comparable to e^+e^- -mode (see also

Baillargeon et al. '97;
Bozovic-Jelisavcic et al. '02)

full study requires consideration of distributions



Summary

- Relevance of $\gamma\gamma \rightarrow WW$ due to its high cross section
- Calculation of Born amplitudes for $\gamma\gamma \rightarrow 4f$ and $\gamma\gamma \rightarrow 4f\gamma$
- Monte Carlo generator with multi-channel Monte Carlo integration
- Inclusion of a realistic photon spectrum
- Anomalous couplings, Higgs resonance
- Double-pole approximation is a promising approach for radiative corrections (cf. RacoonWW), work in progress



PRECISION MEASUREMENTS OF BEAM POLARIZATION AT THE LC (THE CASE OF SINGLE-W PRODUCTION)

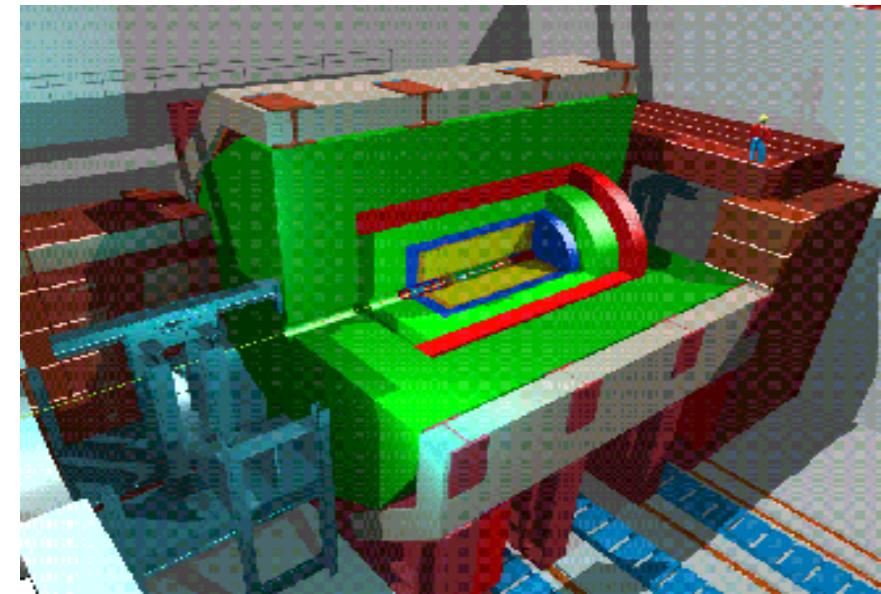
FILIP FRANCO-SOLLOVA



CONTENTS

- BEAM POLARIZATION AT THE LC
- MEASURING THE POLARIZATION
- SINGLE-W PRODUCTION
- BACKGROUND
- RESULTS
- CONCLUSIONS AND OUTLOOK

BEAM POLARIZATION AT THE LC



High Lum. Linear Collider

+

Longit. Beam Polarization

=

Precision Measurements
to test the SM

For some studies this is possible
ONLY if Beam Polarization is accurately determined

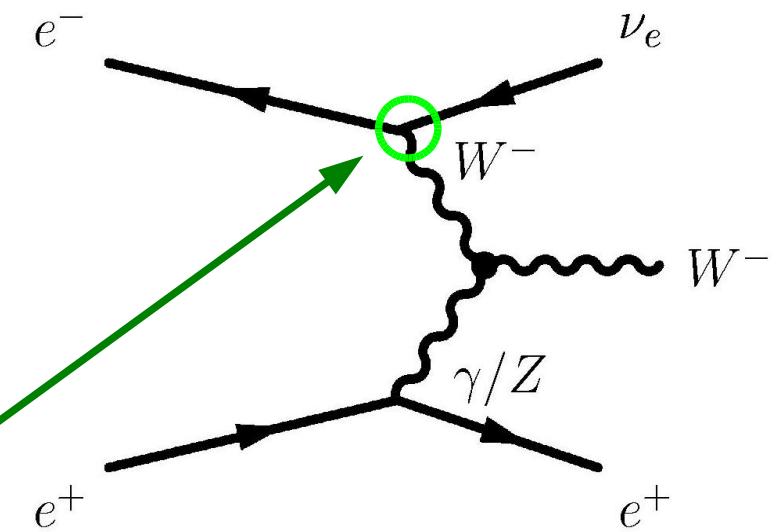
MEASURING THE POLARIZATION

A Solution: Measure the polarization directly using physical processes.
(It is not a replacement of the polarimeters !)

Advantage: No problem with depolarization effects !
(Measures the polarization directly at the interaction point)

SINGLE-W

- $\sigma_{\text{single-}W} \approx \sigma_{W\text{-pair}}$ (At 500 GeV)
- Single- W^- only sensitive to e^- polarization
- Single- W^+ only sensitive to e^+ polarization



How to measure the polarization ?

MEASURING THE POLARIZATION

a) Disadvantages of measuring the asymmetry:

$$|P| = \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-}$$

- The luminosity must be split for measuring σ_+ and σ_-

b) Disadvantages of measuring one cross section σ :

$$P_- = 1 - \frac{4\sigma}{(\sigma_{LL} + \sigma_{LR}) - P_+(\sigma_{LL} - \sigma_{LR})}$$

- The theoretical values of σ_{LL} and σ_{LR} can not be accurately determined
- Anomalous TGCs are problematic

CONCLUSIONS AND OUTLOOK

- From the conditions considered in this stage of the study, the best value for the polarization measurement error is $\approx 0.73\%$ (for $\cos \theta_{e+} > 0.90$)
- The best value of luminosity sharing is:
Left electrons $\approx 25\%$ Right electrons $\approx 75\%$
(considering only the $e^+ \nu_e \mu^- \bar{\nu}_\mu$ final state)
- Optimization of the signal selection
- Untagged events ($\mu\mu$ background)
- Solve technical issues:
 - Detector forward region
 - Problem of ISR and beamstrahlung in Whizard (not mentioned in the talk)
- Simultaneous measurement of electron and positron polarization

2nd workshop of the ECFA “Physics and Detectors for a Linear Collider” study series
Durham, 1–4 September 2004

Electroweak Sudakov logarithms

**The form factor in a massive $U(1)$ model
and in a $U(1) \times U(1)$ model with mass gap**

Bernd Feucht

Institut für Theoretische Teilchenphysik, Universität Karlsruhe

In collaboration with Johann H. Kühn, Alexander A. Penin and Vladimir A. Smirnov

- I Why logarithmic 2-loop results in EW theory?
- II Massive $U(1)$ form factor: evolution equation & 2-loop results
- III $U(1) \times U(1)$ model with mass gap: factorization of IR singularities
- IV Applications: how to treat the EW mass gaps $Z - W -$ photon
- V Summary & outlook

I Why logarithmic 2-loop results in EW theory?

Electroweak (EW) precision physics

- experimentally measured by now at energy scales up to $\sim M_{W,Z}$
- future generation of accelerators (LHC, LC) \rightarrow TeV region

Electroweak radiative corrections

at high energies $\sqrt{s} \sim \text{TeV} \gg M_{W,Z}$

Kühn et al. '00, '01; Fadin et al. '00;
Denner et al. '01, '03; B.F. et al. '03;
Pozzorini '04; ...

large negative corrections in *exclusive* cross sections

- EW corrections dominated by Sudakov logarithms $\alpha^n \ln^{2n}(s/M_{W,Z}^2)$
- 1-loop corrections $\gtrsim 10\%$
- 2-loop corrections $\sim 1\%$, need to be under control for LC

Massive U(1) form factor in 2-loop approximation

Known from resummation & full calculation of n_f contribution: $(n_f = \# \text{ fermions})$

$$\begin{aligned} \alpha^2 F_2 = \left(\frac{\alpha}{4\pi}\right)^2 & \left[+\frac{1}{2} \ln^4 \left(\frac{Q^2}{M^2}\right) - \left(\frac{4}{9} n_f + 3\right) \ln^3 \left(\frac{Q^2}{M^2}\right) \right. \\ & + \left(\frac{38}{9} n_f + \frac{2}{3} \pi^2 + 8\right) \ln^2 \left(\frac{Q^2}{M^2}\right) \\ & - \left(\frac{34}{3} n_f + \dots\right) \ln \left(\frac{Q^2}{M^2}\right) + \left(\frac{16}{27} \pi^2 + \frac{115}{9}\right) n_f + \dots \left. \right] \end{aligned}$$

Kühn, Moch, Penin, Smirnov '01
B.F., Kühn, Moch '03

- growing coefficients with alternating sign:

$$\begin{aligned} & -0.4 n_f \ln^3 + 4.2 n_f \ln^2 - 11.3 n_f \ln + 18.6 n_f \\ & + 0.5 \ln^4 - 3 \ln^3 + 14.6 \ln^2 - \dots \ln + \dots \end{aligned}$$

- $Q \sim 1 \text{ TeV} \rightarrow +\ln^4 \sim -\ln^3 \sim +\ln^2$
 \rightarrow large cancellations between logarithmic terms

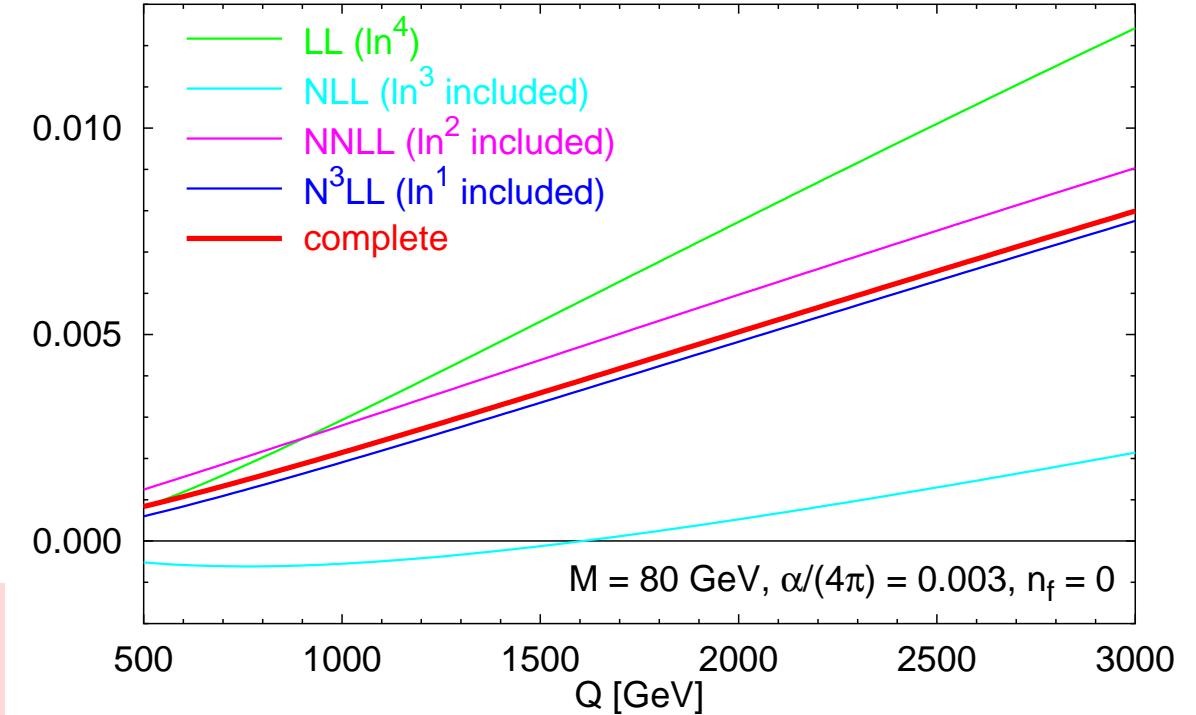
Complete 2-loop corrections in logarithmic approximation necessary.

Massive U(1) form factor in 2-loop approximation: result ($n_f = 0$)

B.F., Kühn, Penin, Smirnov, hep-ph/0404082

$$\alpha^2 F_2 = \left(\frac{\alpha}{4\pi}\right)^2 \left[\begin{aligned} & + \frac{1}{2} \ln^4 \left(\frac{Q^2}{M^2} \right) && \text{agreement ✓} \\ & - 3 \ln^3 \left(\frac{Q^2}{M^2} \right) \\ & + \left(\frac{2}{3} \pi^2 + 8 \right) \ln^2 \left(\frac{Q^2}{M^2} \right) \\ & - \left(-24\zeta_3 + 4\pi^2 + 9 \right) \ln \left(\frac{Q^2}{M^2} \right) \\ & + 256 \text{Li}_4 \left(\frac{1}{2} \right) + \frac{32}{3} \ln^4 2 - \frac{32}{3} \pi^2 \ln^2 2 - \frac{52}{15} \pi^4 + 80\zeta_3 + \frac{52}{3} \pi^2 + \frac{25}{2} \end{aligned} \right]$$

new!



size of coefficients: $+0.5 \ln^4 - 3 \ln^3 + 14.6 \ln^2 - 19.6 \ln + 26.4$
at $Q = 1 \text{ TeV}$: $+326 - 387 + 372 - 99.2 + 26.4$

⇒ alternating signs! small constant ($N^4 LL$) contribution

V Summary & outlook

Massive U(1) form factor

- complete 2-loop result in logarithmic approximation ✓
⇒ precise control of radiative corrections

U(1) \times U(1) model with mass gap

- factorization of IR singularities shown explicitly ✓

Applications

- calculation with mass gap reduced to the 1-mass case $M_W = M_Z = M_{\text{photon}}$
- $M_Z \neq M_W$ taken into account by expanding around the equal mass approximation

Outlook

- extend to non-Abelian models: SU(2), SU(N), SU(2) \times U(1)
- consider Higgs contributions
- 4-fermion scattering amplitude
- predictions for EW corrections to $f\bar{f} \rightarrow f'\bar{f}'$ cross sections

Electroweak precision observables in the MSSM with NMHV

Sven Heinemeyer, CERN

Durham, 09/2004

based on collaboration with
W. Hollik, F. Merz and S. Peñaranda

- 1.** Introduction
- 2.** Results for M_W and $\sin^2 \theta_{\text{eff}}$
- 3.** Results for m_h
- 4.** Conclusions

NMFV in the MSSM

NMFV: Non Minimal Flavor Violation

→ Mixing of scalar quark families (beyond CKM)

Mixing of stop/scharme

$$(\tilde{t}_L, \tilde{t}_R, \tilde{c}_L, \tilde{c}_R) \begin{pmatrix} \tilde{T} & 0 \\ 0 & \tilde{C} \end{pmatrix} \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \\ \tilde{c}_L \\ \tilde{c}_R \end{pmatrix}$$



add NMFV

$$(\tilde{t}_L, \tilde{t}_R, \tilde{c}_L, \tilde{c}_R) \begin{pmatrix} \tilde{T} & \neq 0 \\ \neq 0 & \tilde{C} \end{pmatrix} \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \\ \tilde{c}_L \\ \tilde{c}_R \end{pmatrix}$$

and of sbottom/sstrange:

$$(\tilde{b}_L, \tilde{b}_R, \tilde{s}_L, \tilde{s}_R) \begin{pmatrix} \tilde{B} & 0 \\ 0 & \tilde{S} \end{pmatrix} \begin{pmatrix} \tilde{b}_L \\ \tilde{b}_R \\ \tilde{s}_L \\ \tilde{s}_R \end{pmatrix}$$



$$(\tilde{b}_L, \tilde{b}_R, \tilde{s}_L, \tilde{s}_R) \begin{pmatrix} \tilde{B} & \neq 0 \\ \neq 0 & \tilde{S} \end{pmatrix} \begin{pmatrix} \tilde{b}_L \\ \tilde{b}_R \\ \tilde{s}_L \\ \tilde{s}_R \end{pmatrix}$$

$\neq 0$:

- experimentally only partially restricted
- can e.g. be induced by RGE running in mSUGRA
- changes Higgs-squark couplings
- changes Gauge boson-squark couplings

Analytical result:

evaluation with arbitrary NMHV couplings

Numerical result:

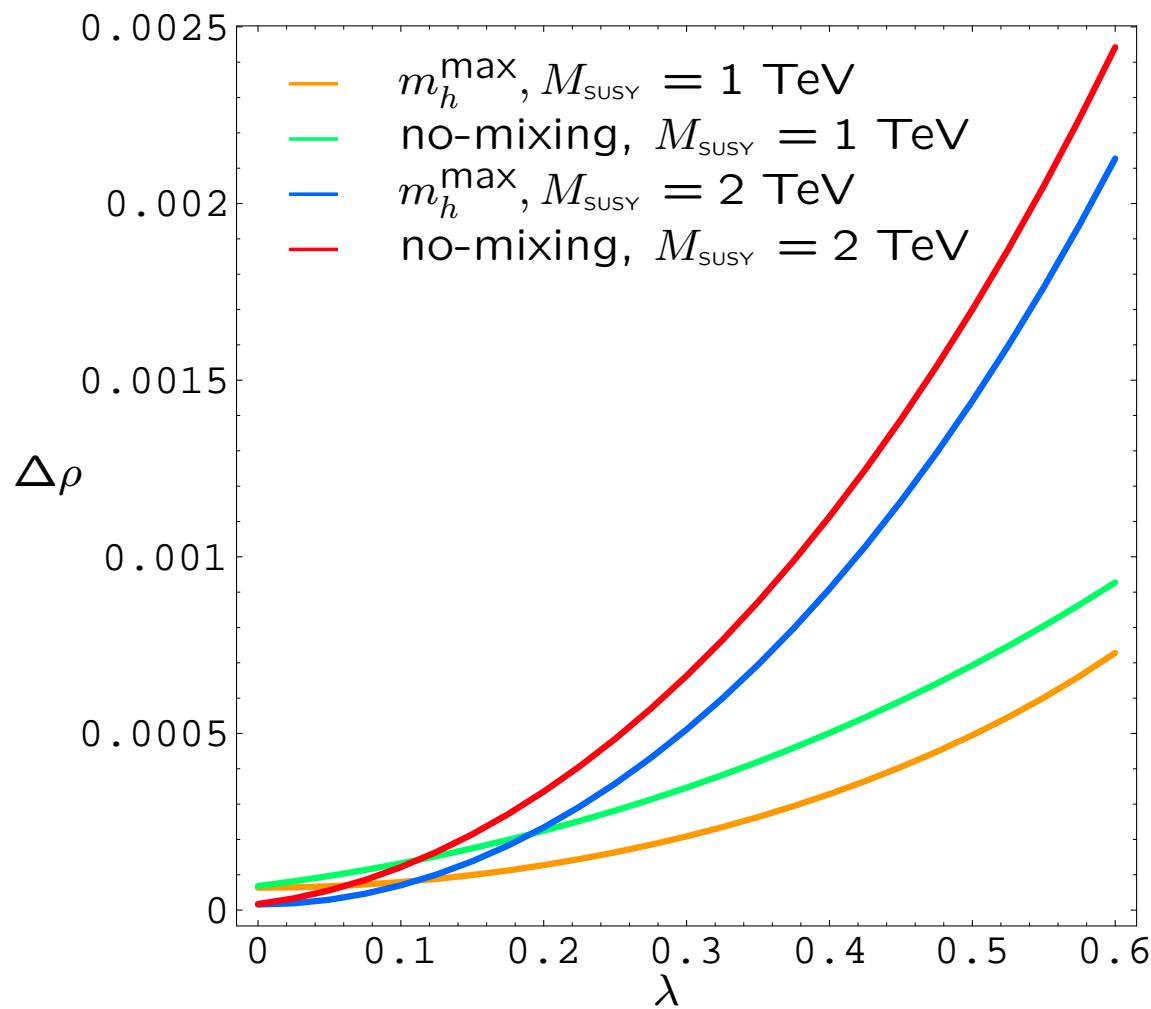
$$\tilde{t}/\tilde{c} : \begin{pmatrix} \lambda \sqrt{\tilde{T}_{LL} \tilde{C}_{LL}} & 0 \\ 0 & 0 \end{pmatrix} \quad \tilde{b}/\tilde{s} : \begin{pmatrix} \lambda \sqrt{\tilde{B}_{LL} \tilde{S}_{LL}} & 0 \\ 0 & 0 \end{pmatrix}$$

$SU(2)$: $\tilde{T}_{LL} \approx \tilde{B}_{LL}$, $\tilde{C}_{LL} \approx \tilde{S}_{LL}$

→ suggested by RGE analysis

→ no relevant experimental bounds on λ

$\Delta\rho$ as a function of λ :



increasing λ

\Rightarrow increasing mixing

\Rightarrow increasing $\Delta\rho$

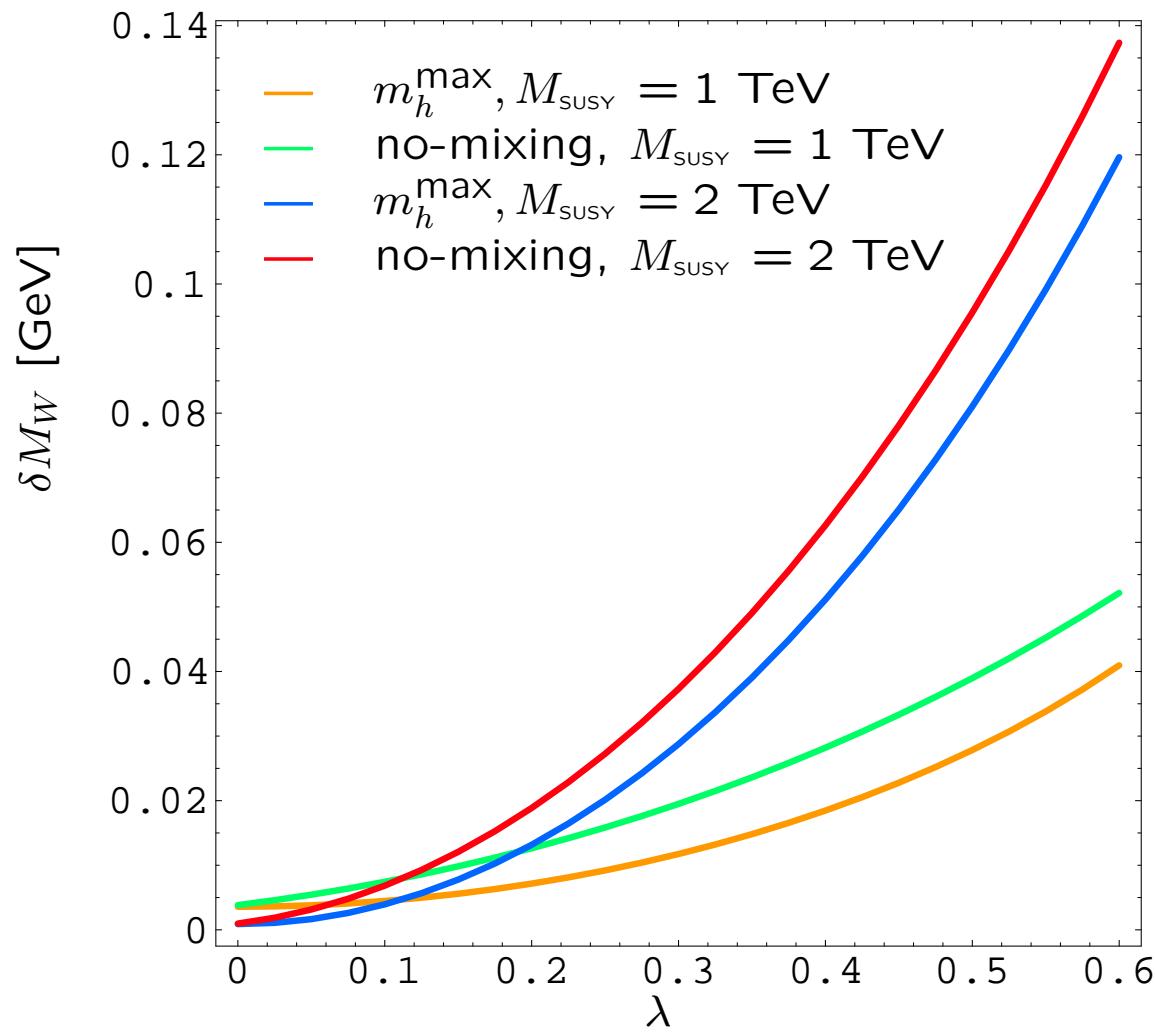
increasing M_{SUSY}

\Rightarrow increasing mixing

\Rightarrow increasing $\Delta\rho$

$\Delta\rho \lesssim 2 \times 10^{-3}$ can be saturated

δM_W as a function of λ :



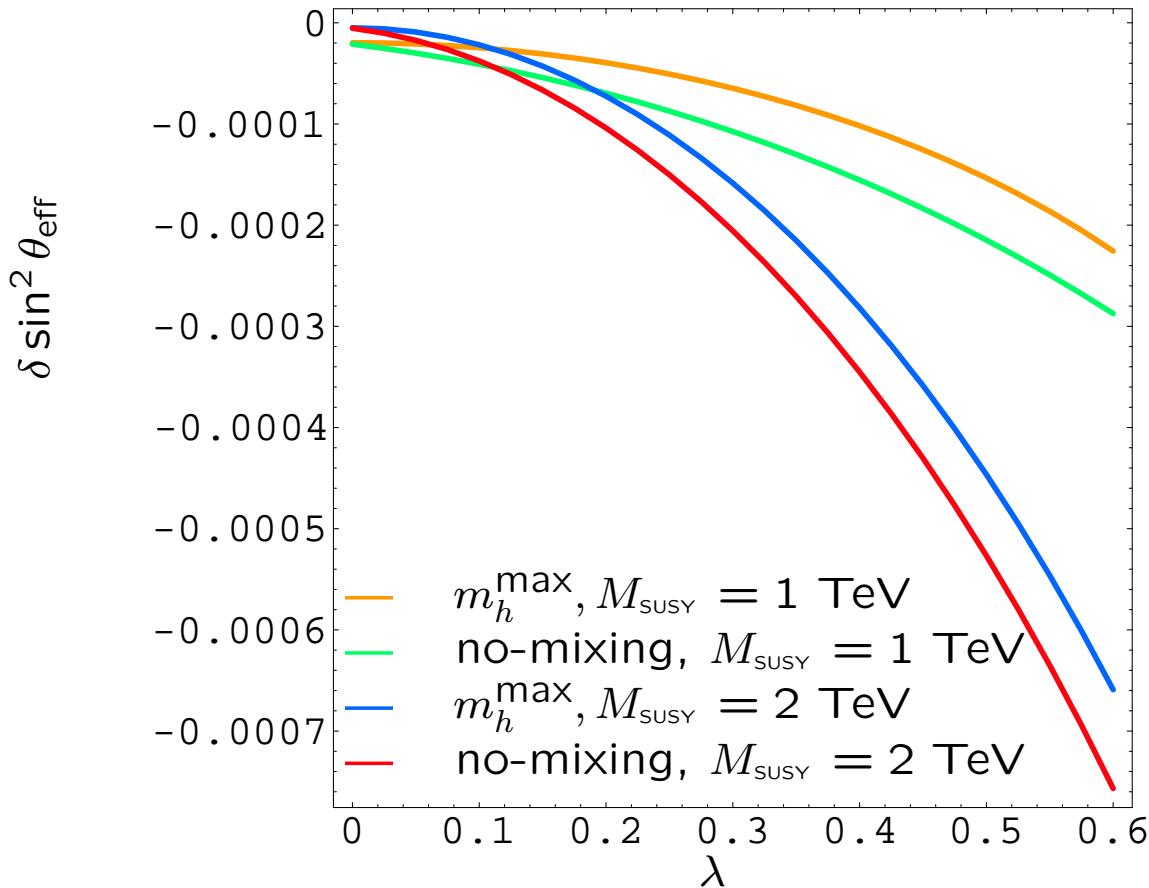
follows the behavior of $\Delta\rho$

$$\delta M_W^{\text{exp,today}} = 34 \text{ MeV}$$

$$\delta M_W^{\text{exp,future}} = 7 \text{ MeV}$$

⇒ extreme parameter
regions already ruled out

$\delta \sin^2 \theta_{\text{eff}}$ as a function of λ :



follows the behavior of $\Delta\rho$

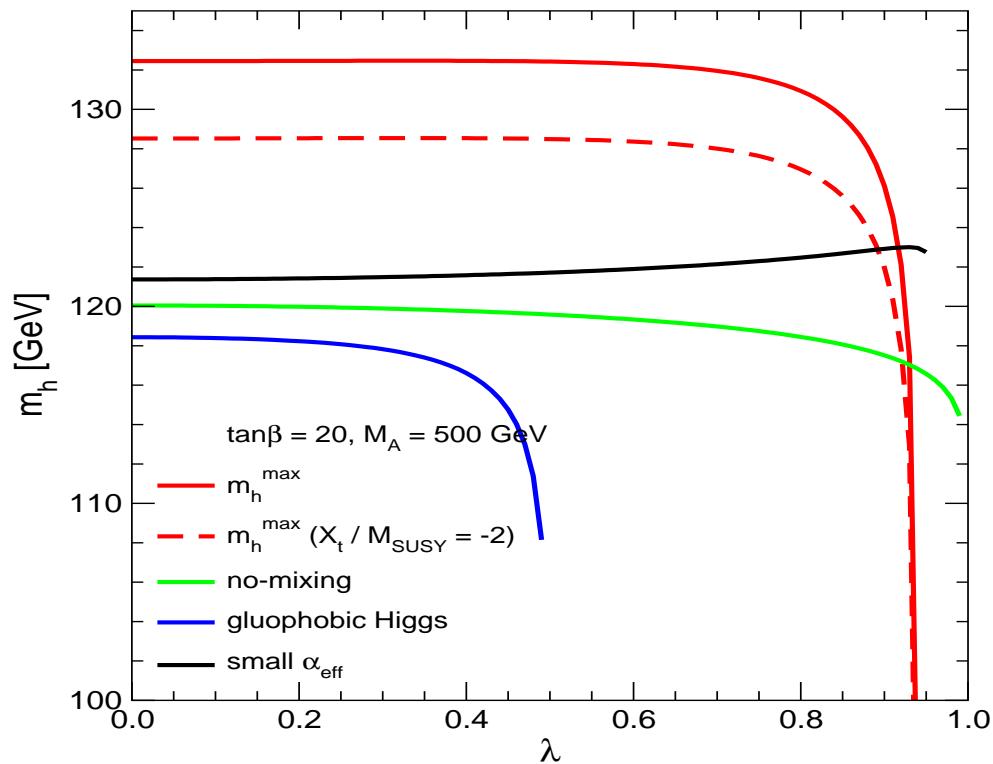
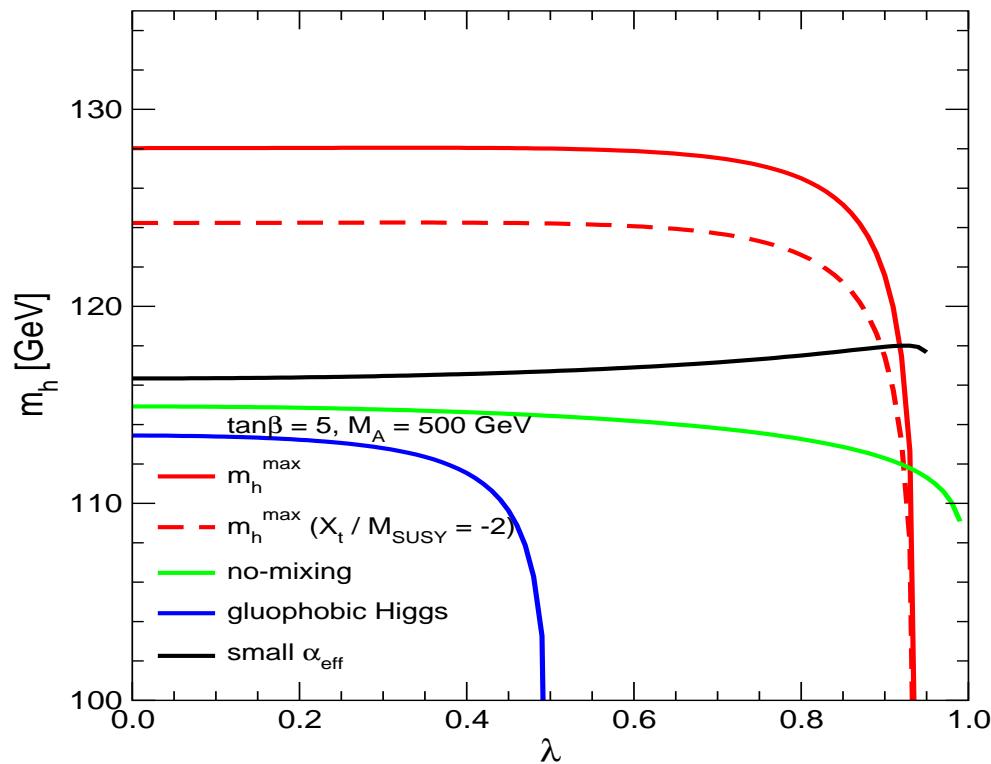
$$\delta \sin^2 \theta_{\text{eff}}^{\text{exp,today}} = 17 \times 10^{-5}$$

$$\delta \sin^2 \theta_{\text{eff}}^{\text{exp,future}} = 1.3 \times 10^{-5}$$

⇒ extreme parameter
regions already ruled out

⇒ highly sensitive test in the
future

Effects in benchmark scenarios:



→ small effects for small/moderate λ

⇒ $\delta m_h = \mathcal{O}(5 \text{ GeV})$ only for very large λ

→ mostly decreasing m_h , but also increase possible
(e.g. in small α_{eff} scenario)

4. Conclusions

- Precision observables can
 - give valuable information about the “true” Lagrangian
 - constrain MSSM parameter space already today
- MSSM with NMHV:
 - mixing in the \tilde{t}/\tilde{c} and in the \tilde{b}/\tilde{s} sector
- \Rightarrow Evaluation of M_W , $\sin^2 \theta_{\text{eff}}$, m_h in NMHV MSSM
- Analytical results: for arbitrary mixing
Numerical results: only for LL mixing, parametrized with λ
corresponds to $(\delta_{LL})_{23}$
- large effects possible for M_W , $\sin^2 \theta_{\text{eff}}$: $\lambda \lesssim 0.2 \Rightarrow \delta M_W \lesssim 20 \text{ MeV}$
 $\lambda \lesssim 0.2 \Rightarrow \delta \sin^2 \theta_{\text{eff}} \lesssim 10^{-4}$
- moderate effects possible for m_h only for large λ