

# Tutorial 1: Neutrino detection

**Lecturer: Thierry Lasserre**

**Tutors: Patrick Foldenauer  
Salvador Rosauero Alcaraz  
Arsenii Titov**

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# $\beta$ -decay: Fermi Theory (1)

We can derive the electron spectrum of the process  $n \rightarrow p + e^- + \bar{\nu}_e$  from first principles in Quantum mechanics:

- Assume free particle states:

$$|i\rangle = |n\rangle \quad \Rightarrow \quad \psi_i(x_1, x_2, x_3, x_4) = e^{i\vec{p}_n \cdot \vec{x}_1}$$

$$|f\rangle = |p, e^-, \bar{\nu}_e\rangle \quad \Rightarrow \quad \psi_f(x_1, x_2, x_3, x_4) = e^{i\vec{p}_p \cdot \vec{x}_2} e^{i\vec{p}_e \cdot \vec{x}_3} e^{i\vec{p}_\nu \cdot \vec{x}_4}$$

- We can assume the interaction to be point-like:

$$V_W(x_1, x_2, x_3, x_4) = G_F \delta^{(3)}(\vec{x}_1 - \vec{x}_2) \delta^{(3)}(\vec{x}_2 - \vec{x}_3) \delta^{(3)}(\vec{x}_3 - \vec{x}_4)$$

New interaction constant



# $\beta$ -decay: Fermi Theory (2)

The matrix element for this process is given by

$$\langle i | V_W | f \rangle = \int d^3x_1 \dots d^3x_4 \psi_i^* V_W(x_1, x_2, x_3, x_4) \psi_f$$

Inserting the expression and integrating over  $x_1, x_2, x_3$  yields

$$\langle i | V_W | f \rangle = G_F \int d^3x^4 e^{i(\vec{p}_p + \vec{p}_e + \vec{p}_\nu - \vec{p}_n) \cdot \vec{x}_4} = (2\pi)^3 G_F \delta^{(3)}(\vec{p}_p + \vec{p}_e + \vec{p}_\nu - \vec{p}_n)$$

With the matrix element we can compute the decay rate of the process using **Fermi's golden rule**:

$$d\Gamma_{if} = 2\pi |V_{if}|^2 \rho(E_i)$$

With the density of states

$$\rho(E_i) = \delta(Q - E_e - E_\nu) \frac{d^3p_p}{(2\pi)^3} \frac{d^3p_e}{(2\pi)^3} \frac{d^3p_\nu}{(2\pi)^3} \quad \text{with} \quad Q = m_n - m_p$$

# $\beta$ -decay: Fermi Theory - Exercise

**Exercise 1.** Using **Fermi's golden rule** find an expression for the electron spectrum in the beta decay, *i.e.*  $d\Gamma_{if}/dE_e$

**Hint 1:** Remember that  $\int dx \delta(x - a) f(x) = f(a)$

**Hint 2:** Integrate first over the proton momentum  $p_p$  and use the relation

$$\int d^3p_p \delta^{(3)}(\vec{p}_p + \vec{p}_e + \vec{p}_\nu - \vec{p}_n) \delta^{(3)}(\vec{p}_p + \vec{p}_e + \vec{p}_\nu - \vec{p}_n) = \delta^{(3)}(0) = \frac{V}{(2\pi)^3},$$

where  $V$  is a normalisation volume that can be set to 1.

**Hint 3:** Integrate over  $p_\nu$  and substitute  $p_e$  using the relations

$$d^3p_\nu = 4\pi E_\nu^2 dE_\nu$$

$$d^3p_e = 4\pi p_e^2 dp_e = 4\pi \sqrt{E_e^2 - m_e^2} E_e dE_e$$

# $\beta$ -decay: Fermi Theory - Exercise

**Exercise 1. Solution:** Starting with **Fermi's golden rule** and using [Hint 2](#) we obtain the expression:

$$d\Gamma_{if} = \frac{G_F^2}{(2\pi)^5} \delta(Q - E_e - E_\nu) d^3p_e d^3p_\nu$$

Next, we integrate over the neutrino momentum  $p_\nu$  using the first part of [Hint 3](#) since it cannot be measured:

$$d\Gamma_{if} = \frac{G_F^2}{(2\pi)^5} \delta(Q - E_e - E_\nu) d^3p_e 4\pi E_\nu^2 dE_\nu = \frac{G_F^2}{8\pi^4} (Q - E_e)^2 d^3p_e$$

Finally, substituting  $dp_e$  using the second part of [Hint 3](#) we arrive at

$$d\Gamma_{if} = \frac{G_F^2}{2\pi^3} (Q - E_e)^2 \sqrt{E_e^2 - m_e^2} E_e dE_e$$

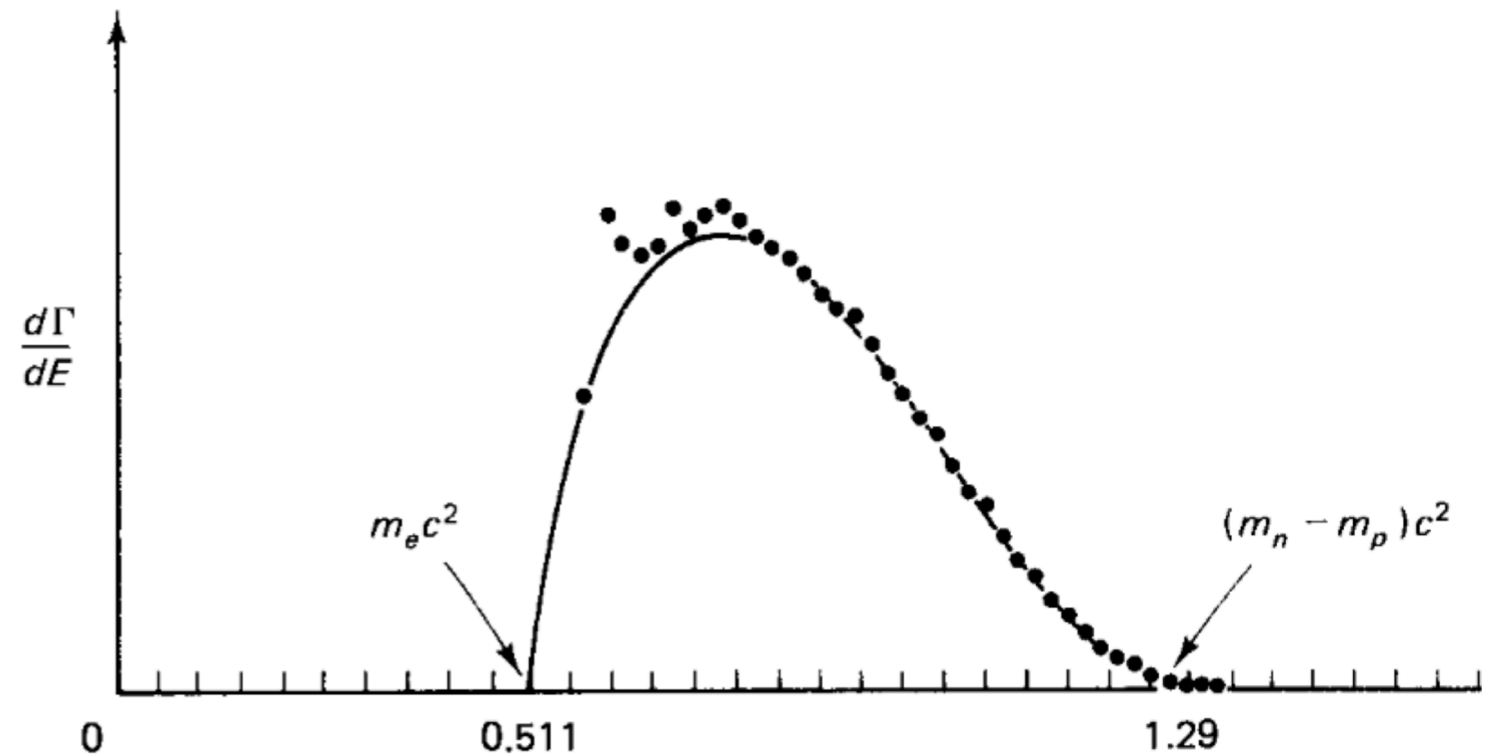
or

$$\frac{d\Gamma_{if}}{dE_e} = \frac{G_F^2}{2\pi^3} (Q - E_e)^2 \sqrt{E_e^2 - m_e^2} E_e$$

# $\beta$ -decay: Fermi Theory (3)

The expression we obtained from our quantum mechanics calculation exactly describes the measured  $\beta$ -decay electron spectrum

$$\frac{d\Gamma_{if}}{dE_e} = \frac{G_F^2}{2\pi^3} (Q - E_e)^2 \sqrt{E_e^2 - m_e^2} E_e$$



This allows to extract the **Fermi constant**  $G_F$  from the spectrum and determine it to

$$G_F \approx 1.15 \times 10^{-5} \text{ GeV}^{-2}$$

# Inverse $\beta$ -decay

With the methods developed before we can also find an expression for the **inverse  $\beta$ -decay**, or neutrino-proton scattering  $\bar{\nu}_e + p \rightarrow n + e^+$ .

We proceed exactly analogous, but with the initial and final states as

$$\begin{aligned} |i\rangle = |\bar{\nu}_e, p\rangle &\Rightarrow \psi_i(x_1, x_2, x_3, x_4) = e^{i\vec{p}_\nu \cdot \vec{x}_1} e^{i\vec{p}_p \cdot \vec{x}_2} \\ |f\rangle = |n, e^+\rangle &\Rightarrow \psi_f(x_1, x_2, x_3, x_4) = e^{i\vec{p}_n \cdot \vec{x}_3} e^{i\vec{p}_e \cdot \vec{x}_4} \end{aligned}$$

Assuming the proton to be at rest  $\vec{p}_p = 0$ , we can derive the cross section from

$$d\sigma = 2\pi |V_{if}|^2 \delta(E_\nu + m_p - E_n - E_e) \frac{d^3p_e}{(2\pi)^3} \frac{d^3p_n}{(2\pi)^3}$$

Inserting the expression for  $V_{ij}$  from before and integrating over  $p_n$  yields the simple expression for the total cross section (neglecting  $m_e$  and taking  $m_n \approx m_p$ )

$$\sigma_{\text{tot}} \approx \frac{G_F^2}{\pi} \frac{E_\nu^2 m_p}{2E_\nu + m_p}$$

# Inverse $\beta$ -decay - Exercise

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**Exercise 2.** Using the expression for the neutrino-proton scattering cross section

$$\sigma_{\text{tot}} \approx \frac{G_F^2}{\pi} \frac{E_\nu^2 m_p}{2E_\nu + m_p}$$

give a rough estimate for the expected cross section of reactor neutrinos with an energy of  $E_\nu \approx 1$  MeV. Express your results in units of barn and  $\text{cm}^2$ .

Compare your result to Thomson scattering of a photon on a non-relativistic electron,  $\sigma_{\text{Thomson}} \approx 10^{-24} \text{ cm}^2$ .

**Hint:**  $1 \text{ barn} \equiv 10^{-24} \text{ cm}^2 = 2.57 \times 10^3 \text{ GeV}^{-2}$



# Inverse $\beta$ -decay - Exercise

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**Exercise 2. Solution** Since  $E_\nu \ll m_p$  we can further approximate (we don't have to)

$$\sigma_{\text{tot}} \approx \frac{G_F^2}{\pi} E_\nu^2$$

Inserting the numbers then yields

$$\sigma_{\text{tot}} \approx 4 \times 10^{-17} \text{ GeV}^{-2} \approx 10^{-5} \text{ fb} \approx 10^{-44} \text{ cm}^2$$

The neutrino-proton scattering cross section is 20 (!) orders of magnitudes smaller than Thomson scattering. We need huge detectors and large fluxes to detect neutrinos!

# Neutrino detection

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**Exercise 3.** Consider a research nuclear reactor with a power of  $P_{\text{reactor}} = 20 \text{ MW}$  which mainly has  ${}_{92}^{235}\text{U}$ . If an average of 6 neutrinos are emitted from each fission, what would be the neutrino flux at a location 150 m away from the core?

**Hint:** The reaction is  ${}_{92}^{235}\text{U} + {}_0^1n \rightarrow {}_{55}^{144}\text{Cs} + {}_{37}^{90}\text{Rb} + 2{}_0^1n$ , and the atomic masses for each element are:

$$m({}^{235}\text{U}) = 235.044u,$$

$$m({}^{144}\text{Cs}) = 143.932u,$$

$$m({}^{90}\text{Rb}) = 89.915u,$$

$$m({}_0^1n) = 1.009u.$$

$$1u = 931.5 \text{ MeV}$$

**Hint 2:**  $1 \text{ MW} \approx 6.25 \cdot 10^{18} \text{ MeV/s}$

# Neutrino detection

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## Exercise 3. Solution

The excess mass is given by the difference between the fission products and the fuel:  $\Delta m = 235.044u + 1.009u - 143.932u - 89.915u - 2.018u = 0.188u$ , and the energy released in each fission is then

$$\Delta E = \Delta mc^2 = 0.188 \cdot 931.5 \text{ MeV} = 175.1 \text{ MeV}$$

The neutrino flux will be given by

$$\phi_\nu = \frac{6 \cdot N_{\text{fission}}}{4\pi L^2 \Delta t} = \frac{6 \cdot P_{\text{reactor}}}{4\pi L^2 \Delta E} = 1.5 \cdot 10^9 \text{ cm}^{-2} \text{ s}^{-1}$$

# Neutrino detection

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**Exercise 4.** If we wanted to detect these reactor neutrinos through inverse beta decay (IBD), how much mass of the liquid scintillator  $C_{18}H_{30}$  would we need to have an event rate of about 10 events/day?

**Hint1:** Remember that the IBD cross section on free protons is  $\sigma_{\text{IBD}} \sim 10^{-44} \text{cm}^2$ , and the event rate would be given by

$$N_{\text{IBD}}/\Delta t = \phi_{\nu} \sigma_{\text{IBD}} N_{\text{targets}}$$

**Hint2:** Assume that the number of free protons per molecule is given by the number of  $H$  atoms

The molar mass for  $C_{18}H_{30}$  is  $M_{C_{18}H_{30}} = 246.4 \text{ g/mol}$

Avogadro's number is  $N_A = 6.022 \times 10^{23} / \text{mol}$

# Neutrino detection

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## Exercise 4. Solution

The number of free protons would be given by the number of hydrogen atoms in the  $C_{18}H_{30}$  molecule times the number of molecules. Thus

$$N_{\text{targets}} = 30 \cdot \frac{m_{\text{detector}}}{M_{C_{18}H_{30}}} \cdot N_A$$

Using the previously calculated flux and cross section, we find

$$m_{\text{detector}} = \frac{N_{\text{IBD}}}{\phi_{\nu} \sigma_{\text{IBD}} \Delta t} \cdot \frac{M_{C_{18}H_{30}}}{30 N_A} \approx 105 \text{ tons}$$

# Neutrino detection

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**Exercise 5.** Apart from IBD, there are other processes through which neutrinos can interact with matter. A particularly interesting one is coherent neutrino nucleus scattering (CEvNS), in which low energy neutrinos interact with the whole atomic nucleus instead of interacting with individual nucleons.

How large would a detector made of CsI(Na) need to be to observe an event rate of about 10 events/day?

Compare this to the previous detector relying on IBD.

**Hint:** The CEvNS cross section on Cs for neutrinos with  $E_\nu \sim 2$  MeV is  $\sigma_{\text{CEvNS}} \sim 10^{-40} \text{ cm}^2$ .

The molar mass of CsI is  $M_{\text{CsI}} = 259.8 \text{ g/mol}$

# Neutrino detection

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## Exercise 5. Solution

Now we only want to know the number of Cs atoms in our detector. Thus

$$N_{\text{targets}} = \frac{m_{\text{detector}}}{M_{\text{CsI}}} \cdot N_A$$

Using the previously calculated flux and cross section, we find

$$m_{\text{detector}} = \frac{N_{\text{CE}\nu\text{NS}}}{\phi_{\nu} \sigma_{\text{CE}\nu\text{NS}} \Delta t} \cdot \frac{M_{\text{CsI}}}{N_A} \approx 333 \text{ kg}$$