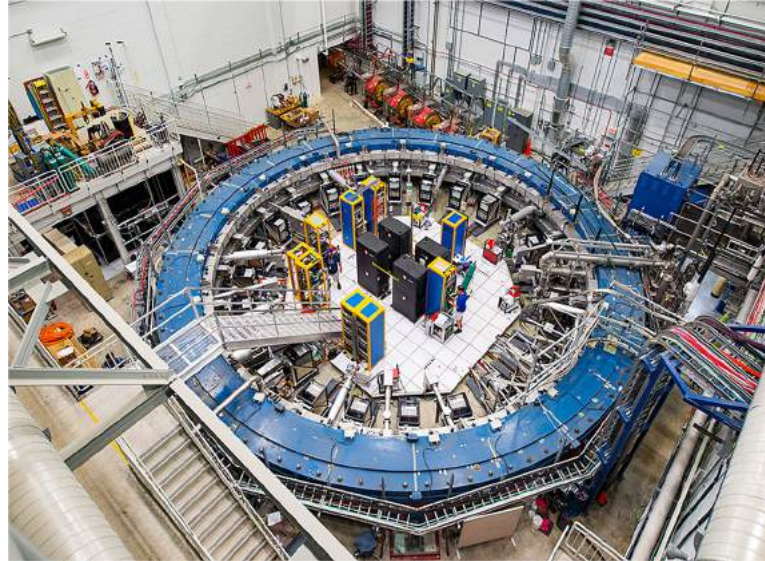


Muon g-2 Experiments



Alex Keshavarzi

 @AlexKeshavarzi

UK HEP Forum 2021

23rd November 2021



Magnetic moments

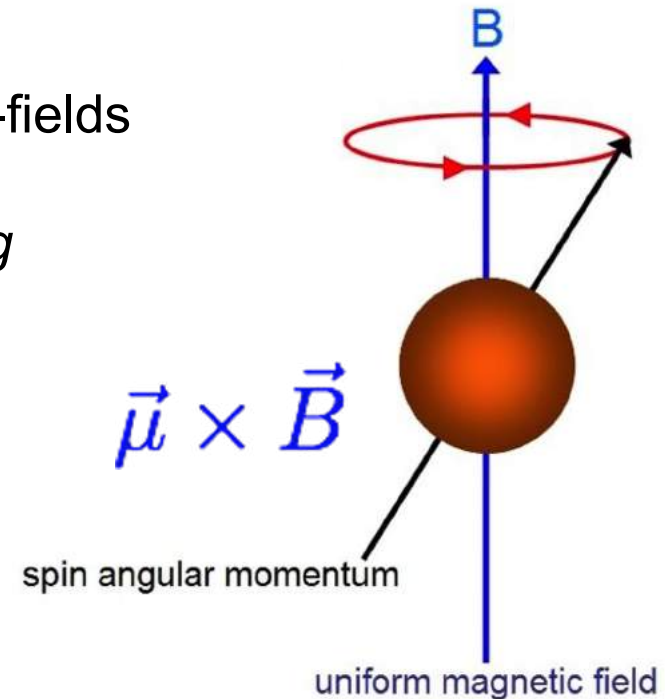
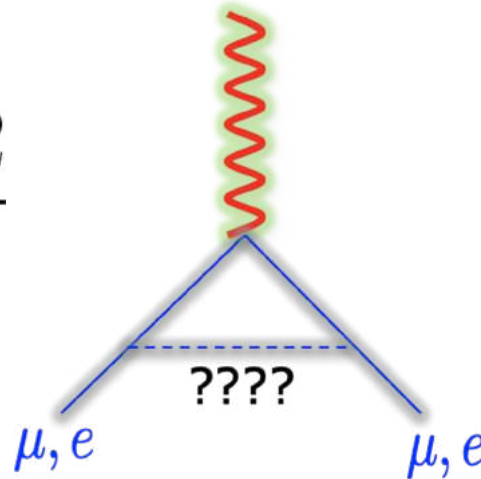
The muon has an intrinsic magnetic moment that is coupled to its spin via the gyromagnetic ratio g :

$$\vec{\mu} = g \frac{e}{2m_\mu} \vec{S}$$

Magnetic moment (spin) interacts with external B-fields

Makes spin precess at frequency determined by g

$$a_\mu = \frac{g - 2}{2}$$



Muon g-2 in the SM

$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$

Contribution	Value $\times 10^{11}$
Experiment (E821)	116 592 089(63)
HVP LO (e^+e^-)	6931(40)
HVP NLO (e^+e^-)	-98.3(7)
HVP NNLO (e^+e^-)	12.4(1)
HVP LO (lattice, $udsc$)	7116(184)
HLbL (phenomenology)	92(19)
HLbL NLO (phenomenology)	2(1)
HLbL (lattice, uds)	79(35)
HLbL (phenomenology + lattice)	90(17)
QED	116 584 718.931(104)
Electroweak	153.6(1.0)
HVP (e^+e^- , LO + NLO + NNLO)	6845(40)
HLbL (phenomenology + lattice + NLO)	92(18)
Total SM Value	116 591 810(43)
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	279(76)

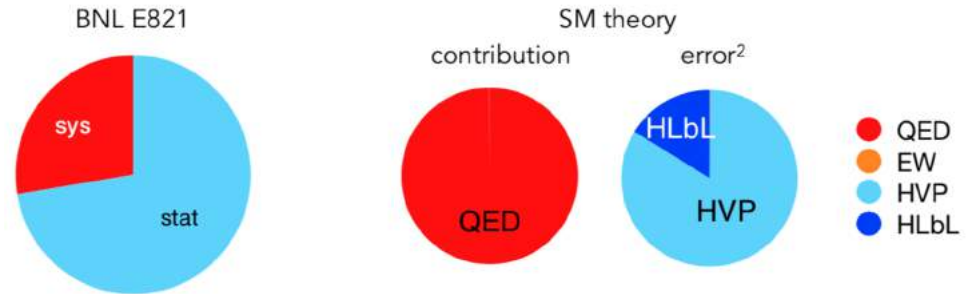
arXiv.org > hep-ph > arXiv:2006.04822

High Energy Physics - Phenomenology

(Submitted on 8 Jun 2020)

The anomalous magnetic moment of the muon in the Standard Model

T. Aoyama, N. Asmussen, M. Benayoun, J. Bijnens, T. Blum, M. Bruno, I. Caprini, C. M. Carloni Calame, M. Cè, G. Colangelo, F. Cuciarello, H. Czyż, I. Danilkin, M. Davies, C. T. H. Davies, M. Della Morte, S. I. Eidelman, A. X. El-Khadra, A. Gérardin, D. Giusti, M. Golterman, Steven Gottlieb, V. Gülpers, F. Hagelstein, M. Hayakawa, G. Herdoíza, D. W. Hertzog, A. Hoecker, M. Hoferichter, B.-L. Hoid, R. J. Hudspeth, F. Ignatov, T. Izubuchi, F. Jegerlehner, L. Jin, A. Keshavarzi, T. Kinoshita, S. Kubis, A. Kupich, A. Kupść, L. Laub, C. Lehner, L. Lellouch, I. Logashenko, B. Malaescu, K. Maltman, M. K. Marinković, P. Masjuan, A. S. Meyer, H. B. Meyer, T. Mibe, K. Miura, S. E. Müller, M. Nio, D. Nomura, A. Nyffeler, V. Pascalutsa, M. Passera, E. Perez del Rio, S. Peris, A. Portelli, M. Procura, C. F. Redmer, B. L. Roberts, P. Sánchez-Puertas, S. Serednyakov, B. Shwartz, S. Simula, D. Stöckinger, H. Stöckinger-Kim, P. Stoffer, T. Teubner, R. Van de Water, M. Vanderhaeghen, G. Venanzoni, G. von Hippel, H. Wittig, Z. Zhang, M. N. Achasov, A. Bashir, N. Cardoso, B. Chakraborty, E.-H. Chao, J. Charles, A. Crivellin, O. Deineka, A. Denig, C. DeTar, C. A. Dominguez, A. E. Dorokhov, V. P. Druzhinin, G. Eichmann, M. Fael, C. S. Fischer, E. Gámiz, Z. Gelzer, J. R. Green, S. Guellati-Khelifa, D. Hatton, N. Hermansson-Truedsson et al. (32 additional authors not shown)



$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65) \text{ ppm}$$

Muon g-2 theory initiative recommended result:

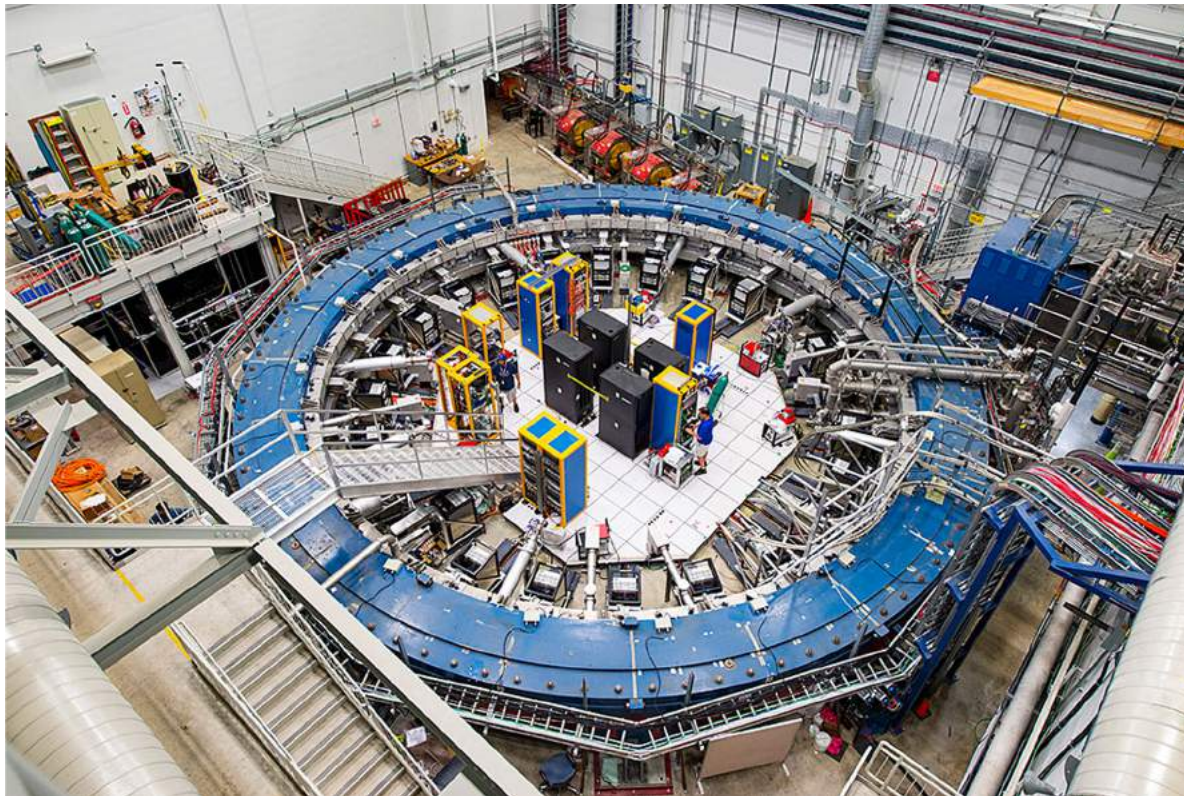
$$a_\mu^{\text{SM}} = 116\,591\,810(43) \times 10^{-11} \text{ (0.37 ppm)}$$

Results in 3.7σ discrepancy when compared to BNL measurement.

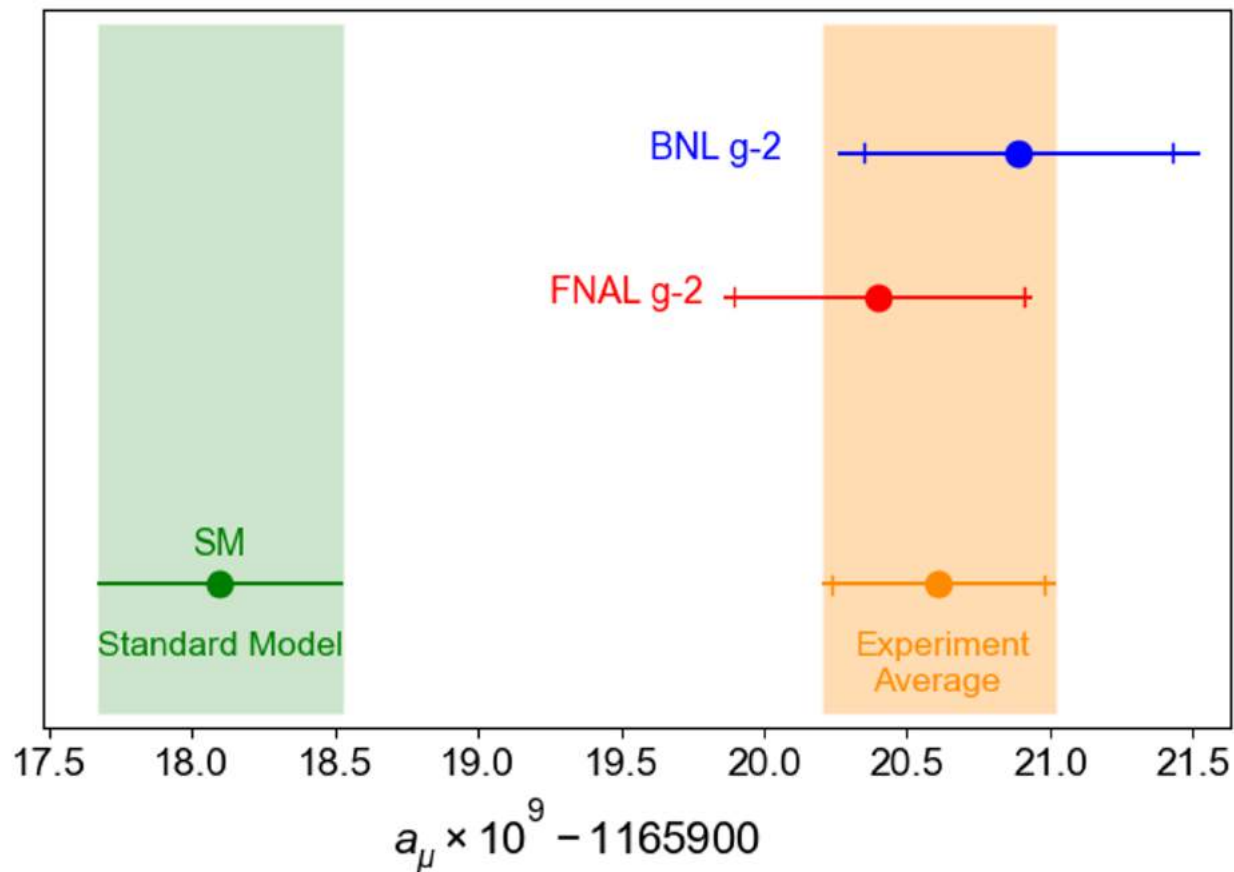
See Andreas Juettner's Muon g-2 theory talk direct after this!



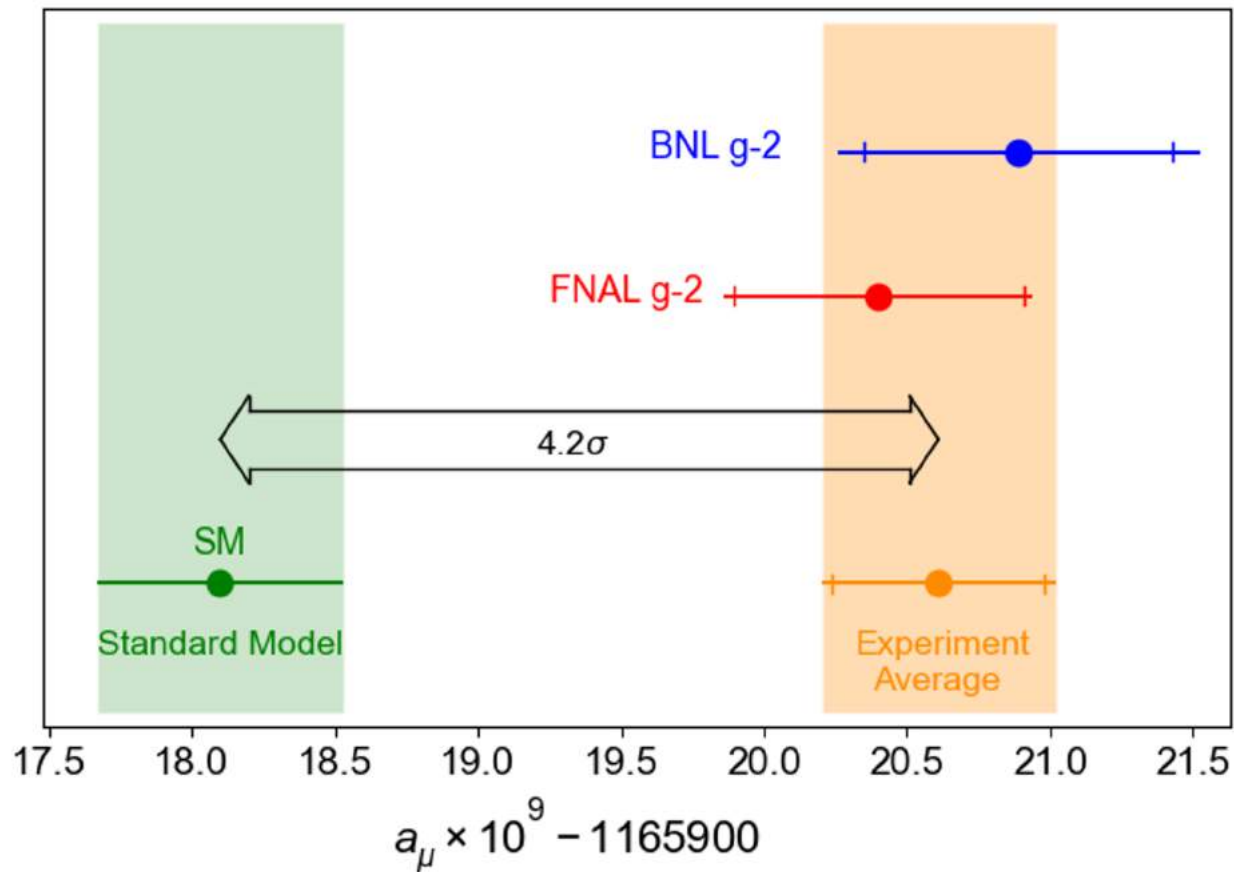
The Fermilab Muon g-2 Experiment



Unblinded result



Unblinded result



Systematic Uncertainties

~ 80 effects considered significant in determining the systematic uncertainty.
Dedicated runs taken for some of them e.g. at different beam momentum.
Documented in 98 pages of PRDs.

Total systematic uncertainty 157 ppb. Those above 30 ppb are below

Source	Systematic Uncertainty (ppb)	Improvements undertaken
Calorimeter pileup	35	
Beam Mean Momentum & Spread	53	Increased kicker voltage: 130-161 kV
Drift of beam over measurement	75	Replaced damaged quadrupole resistors
Transient B-field (from kicker)	37	Improved magnetometer
Transient B-field (from quadrupoles)	92	More extensive measurements / damping
Total	140	

Other effects at 10-20 ppb also significantly improved by better temperature control in the experimental hall.

Measurement principle

- Inject polarised muon beam into magnetic storage ring
- Measure difference between spin precession and cyclotron frequencies

$$g = 2, \omega_a = 0$$

- $g \neq 2, \omega_a \propto a_\mu$

$$\omega_a = \omega_s - \omega_c = a_\mu \frac{eB}{mc}$$

Spin precession freq.

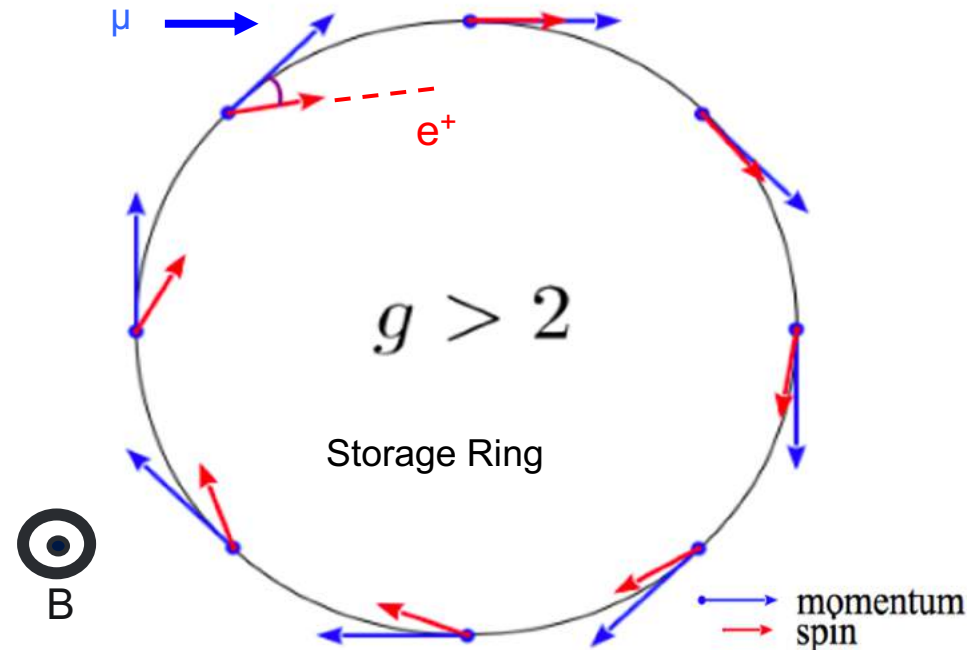
$$\omega_s = \frac{geB}{2mc} + (1 - \gamma) \frac{eB}{\gamma mc}$$

Larmor precession

Cyclotron freq.

$$\omega_c = \frac{eB}{\gamma mc}$$

Thomas precession



Measurement details

The experiment actually measures two frequencies

$$a_\mu = \frac{\omega_a}{\tilde{\omega}_p} \frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$$

What we measure

3ppb

0.0003ppb

22ppb

$$\mathcal{R}'_\mu = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

Unblinding conversion factor

Measured $g - 2$ frequency

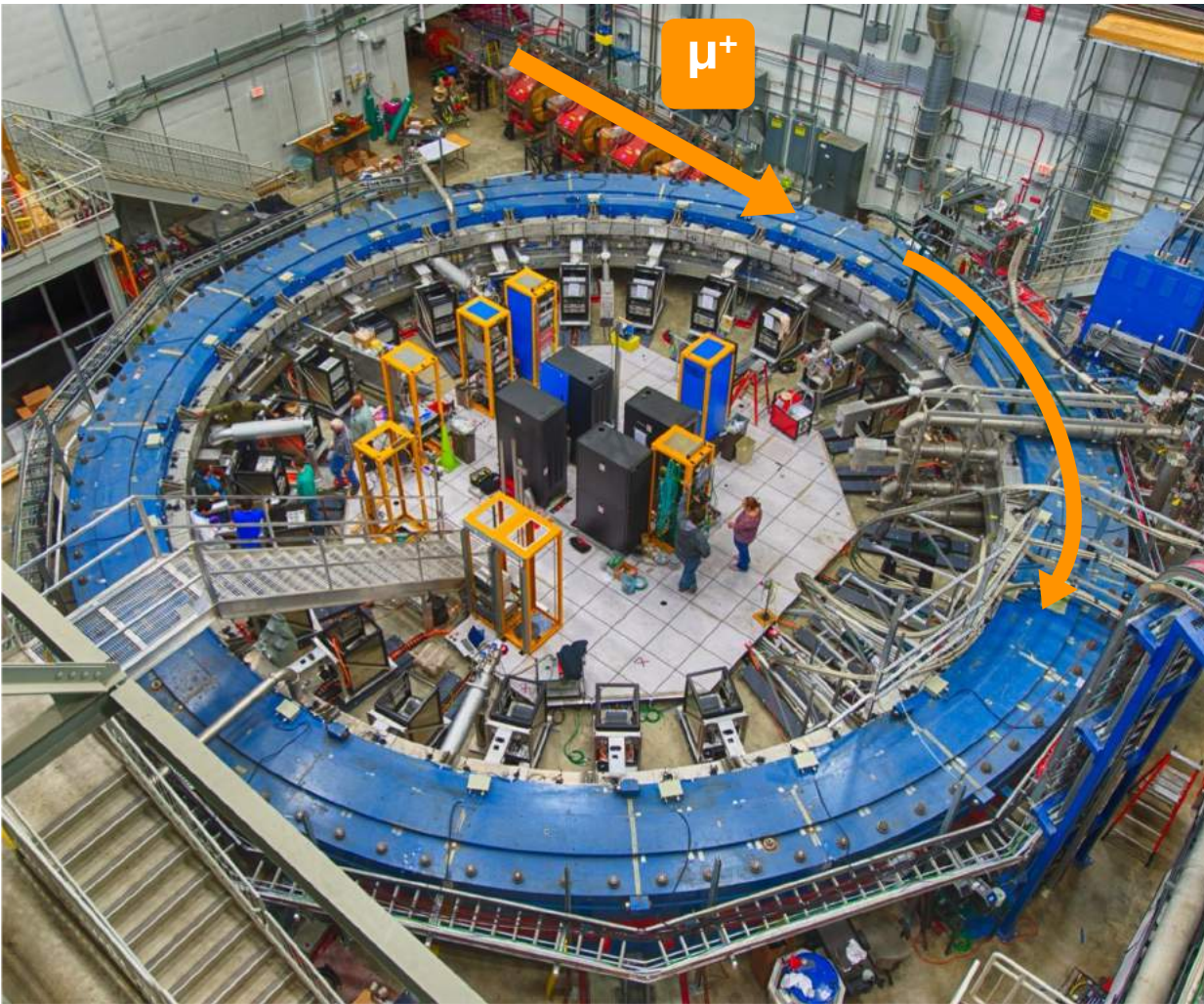
Corrections from the beam dynamics systematic effects

NMR probe calibration factor

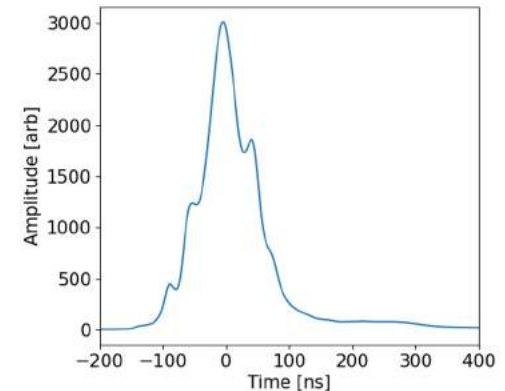
Magnetic field weighted over the muon distribution and azimuthally averaged

Corrections from the transient magnetic field

Beam injection

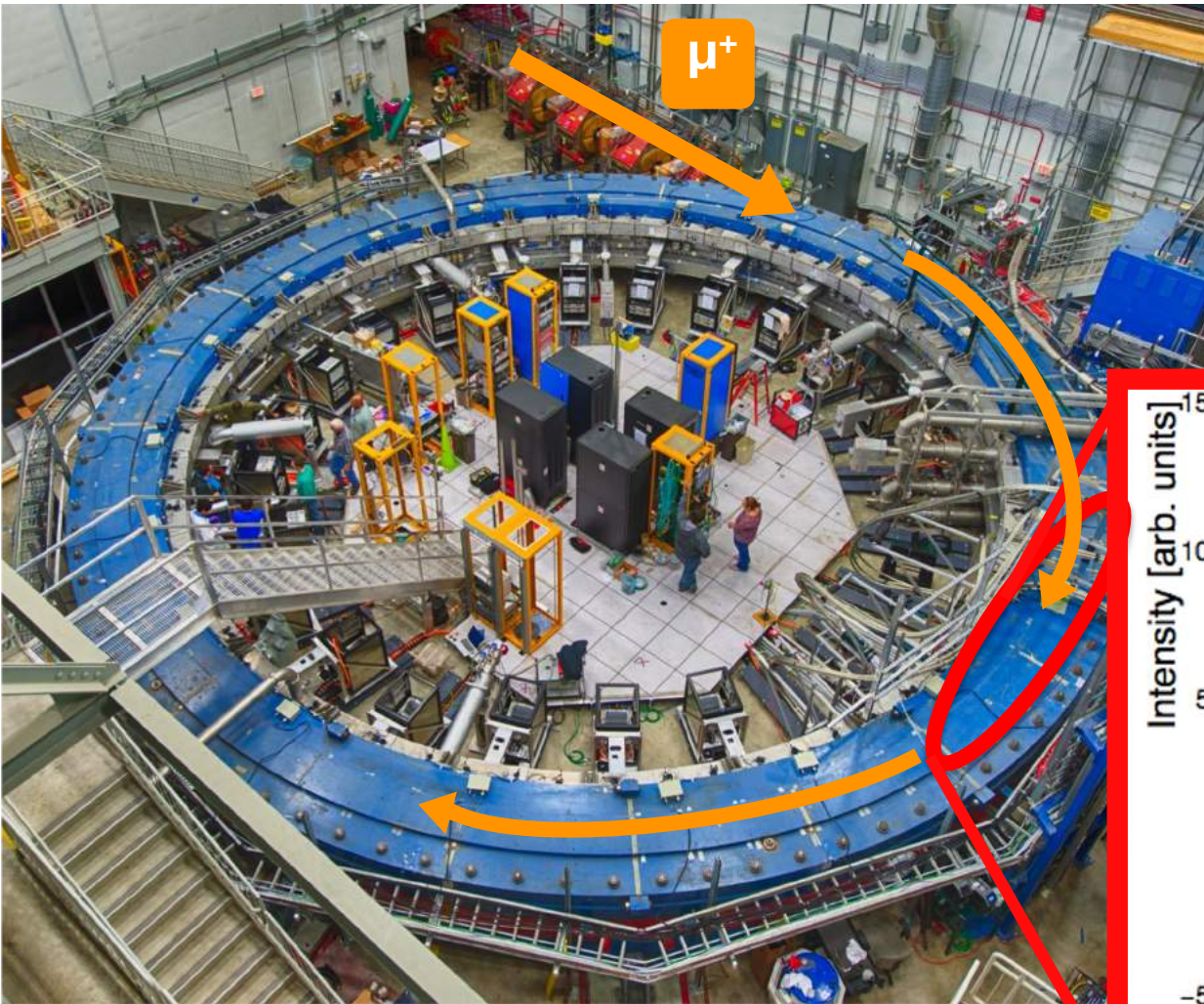


- Monitor beam profile before entrance with scintillating X and Y fibres
- Get time profile of beam using scintillating pad

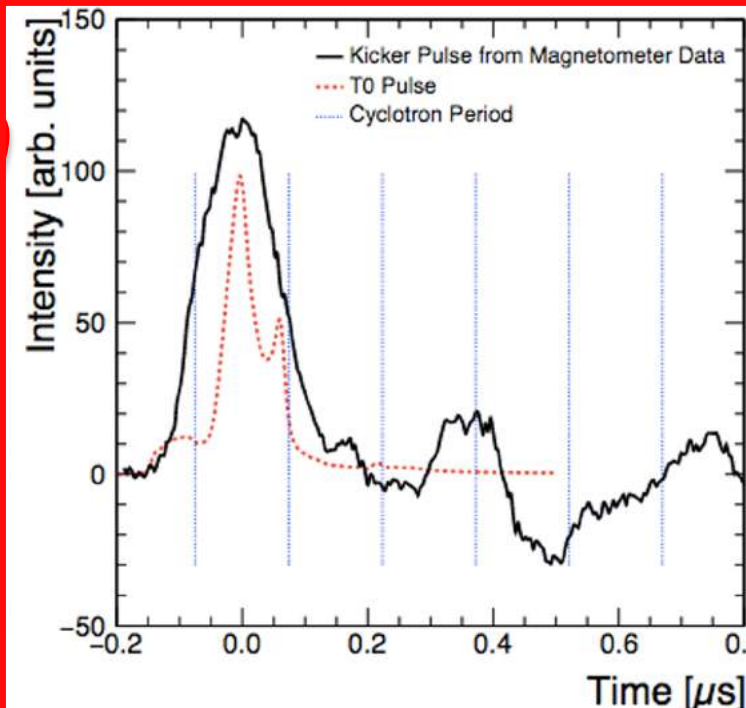


- ~125ns wide
- Cancel B-field during injection using Inflector, so muons can get into the ring

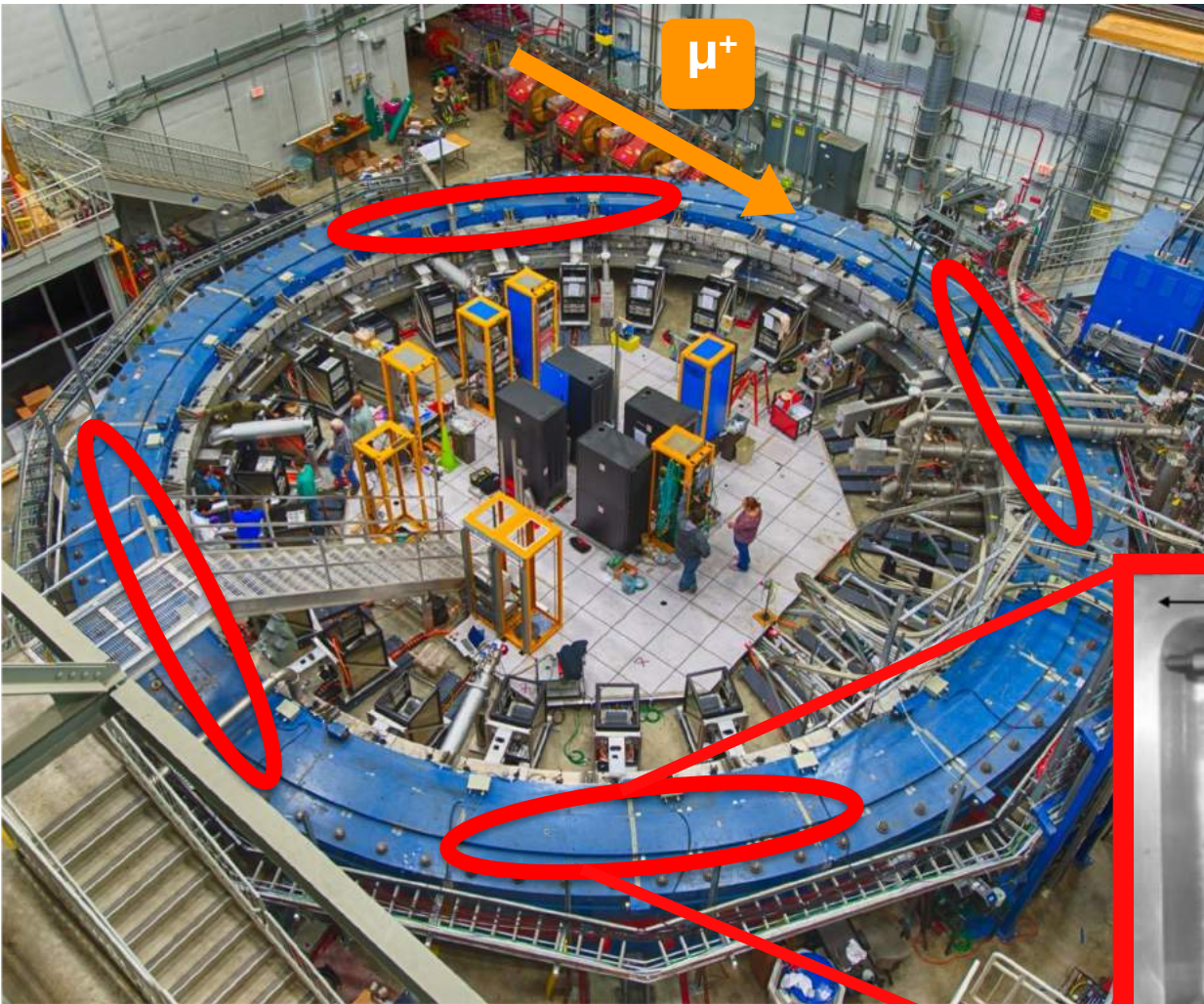
'Kick' onto correct orbit



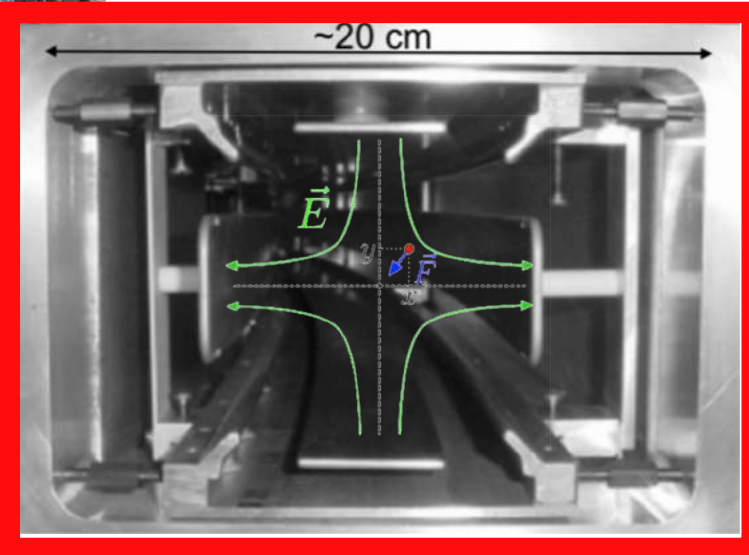
- After Inflector muons are 77mm away from ideal radius
- Apply short magnetic pulse to 'kick' muons onto the correct orbit



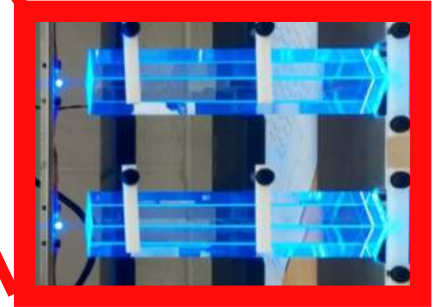
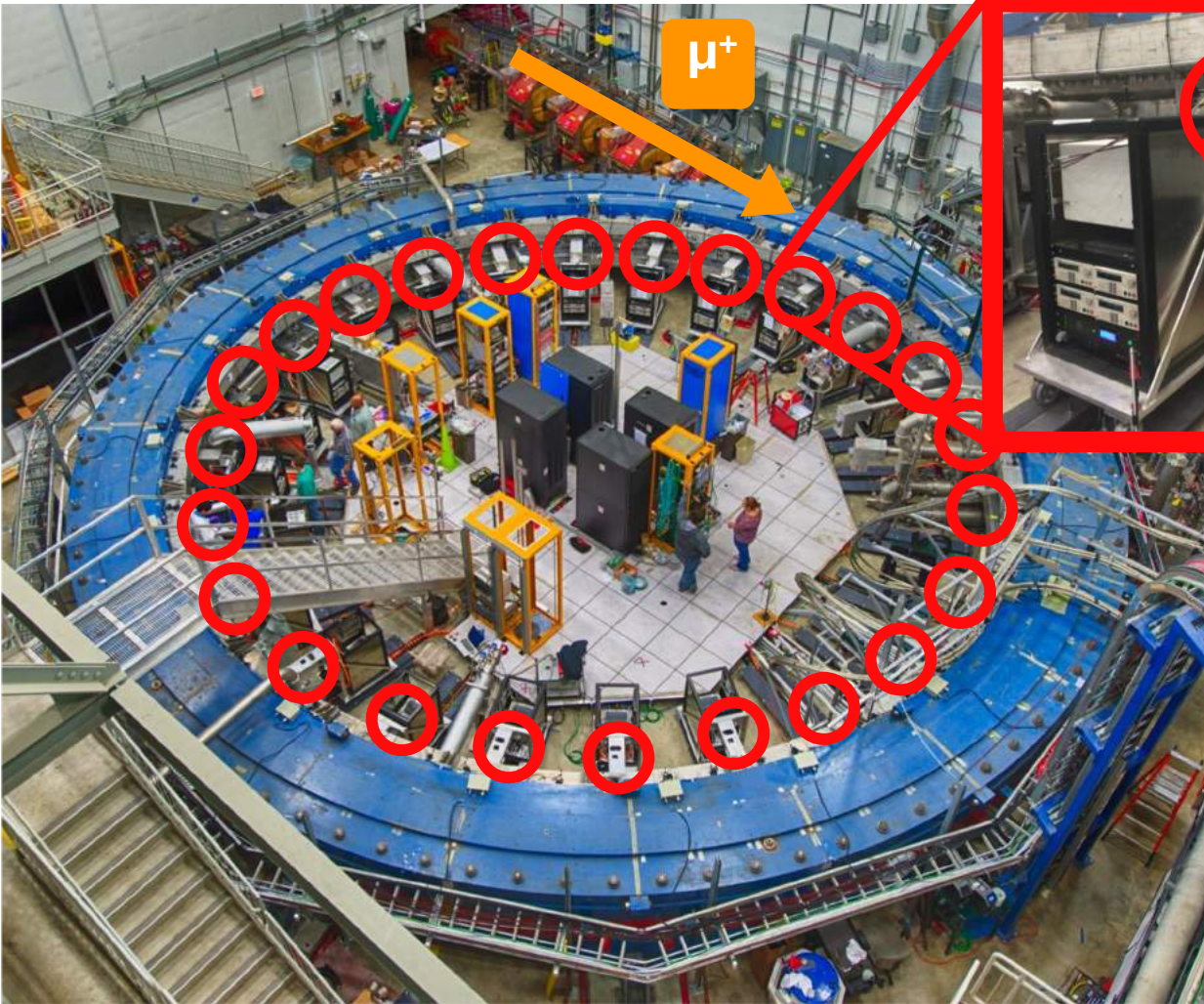
Beam focusing



- Focus the muons vertically
- Aluminium electrodes cover $\sim 43\%$ of total circumference



Calorimeters

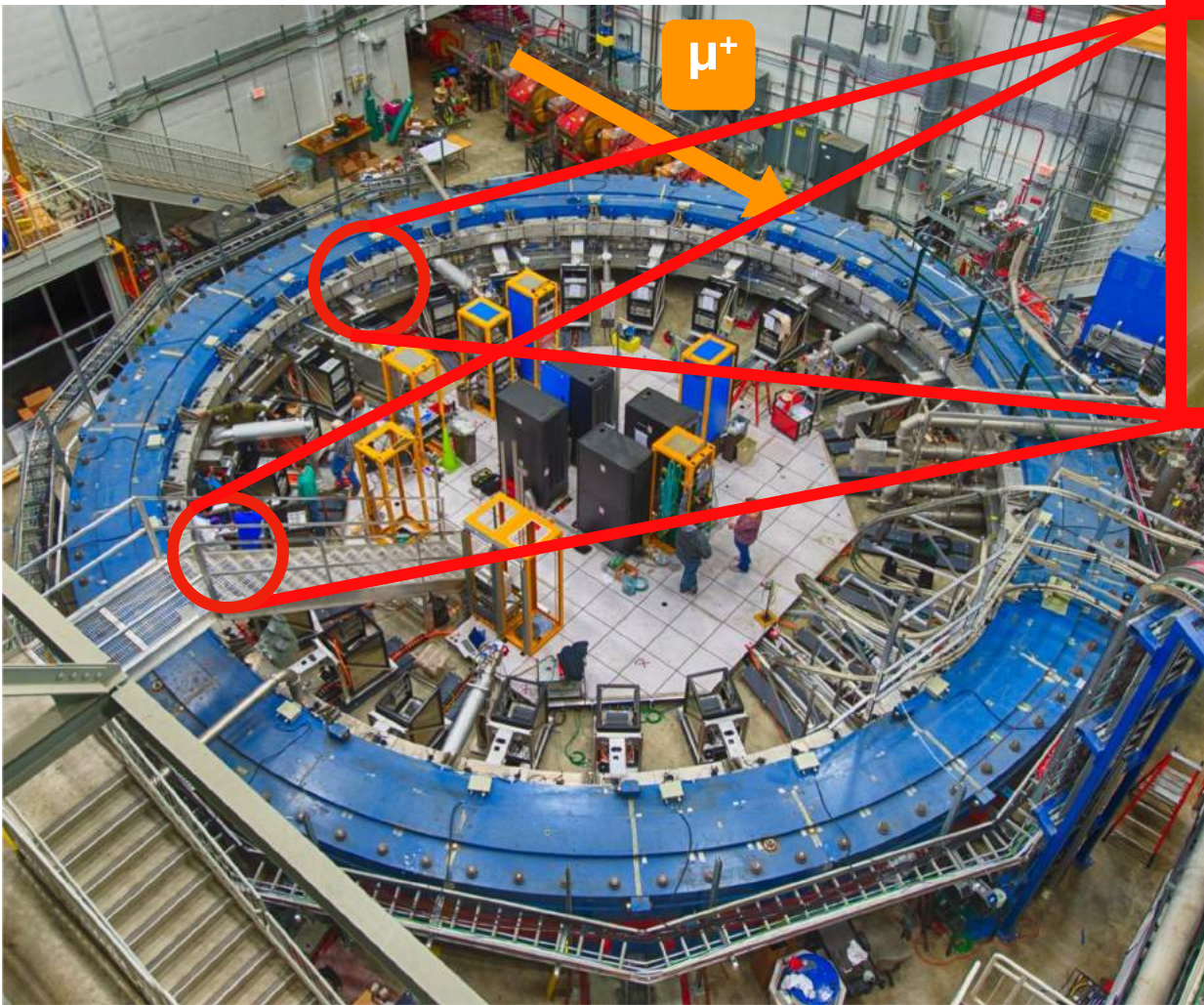


24 Calorimeters

Arrays of 6 x 9 PbF₂ crystals
2.5 x 2.5 cm² x 14 cm (15X₀)

Readout by SiPMs to 800
MHz WFDs

Tracking Detectors



2 Tracking stations

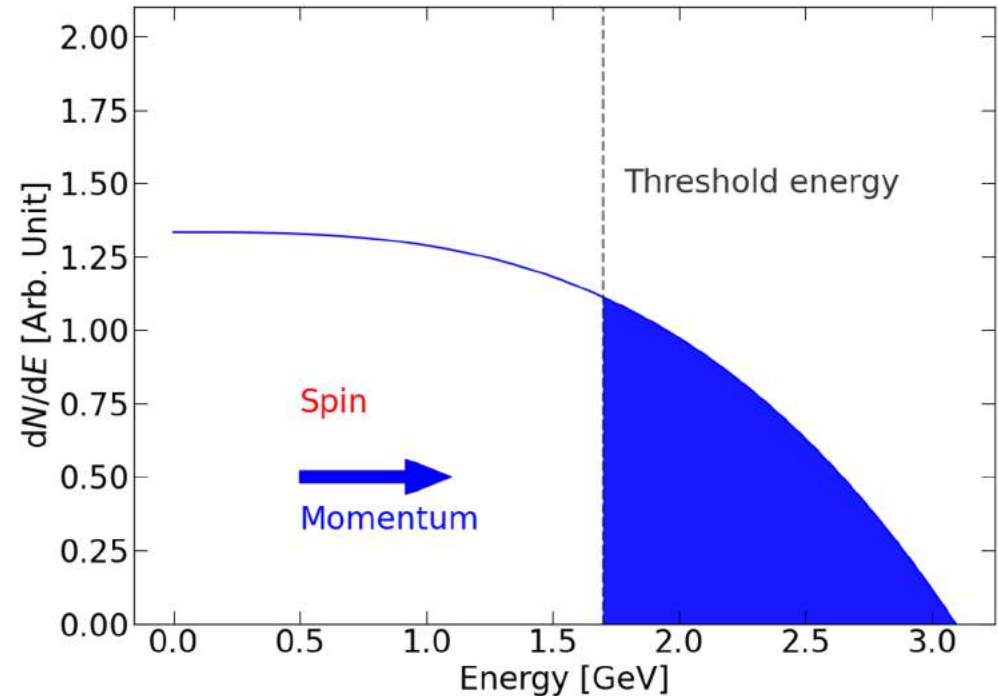
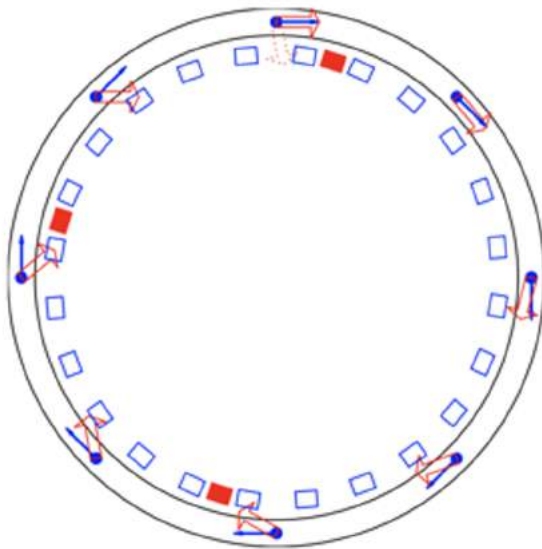
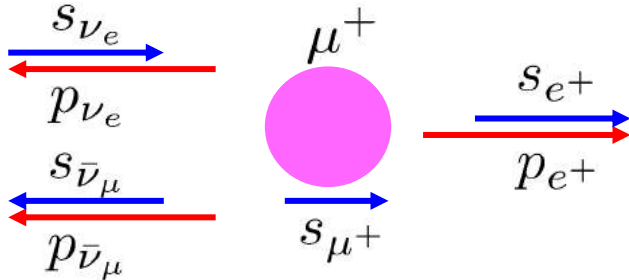
Each contain 8 modules

128 gas filled straws in each module

Traceback positrons to their decay point

Measuring ω_a

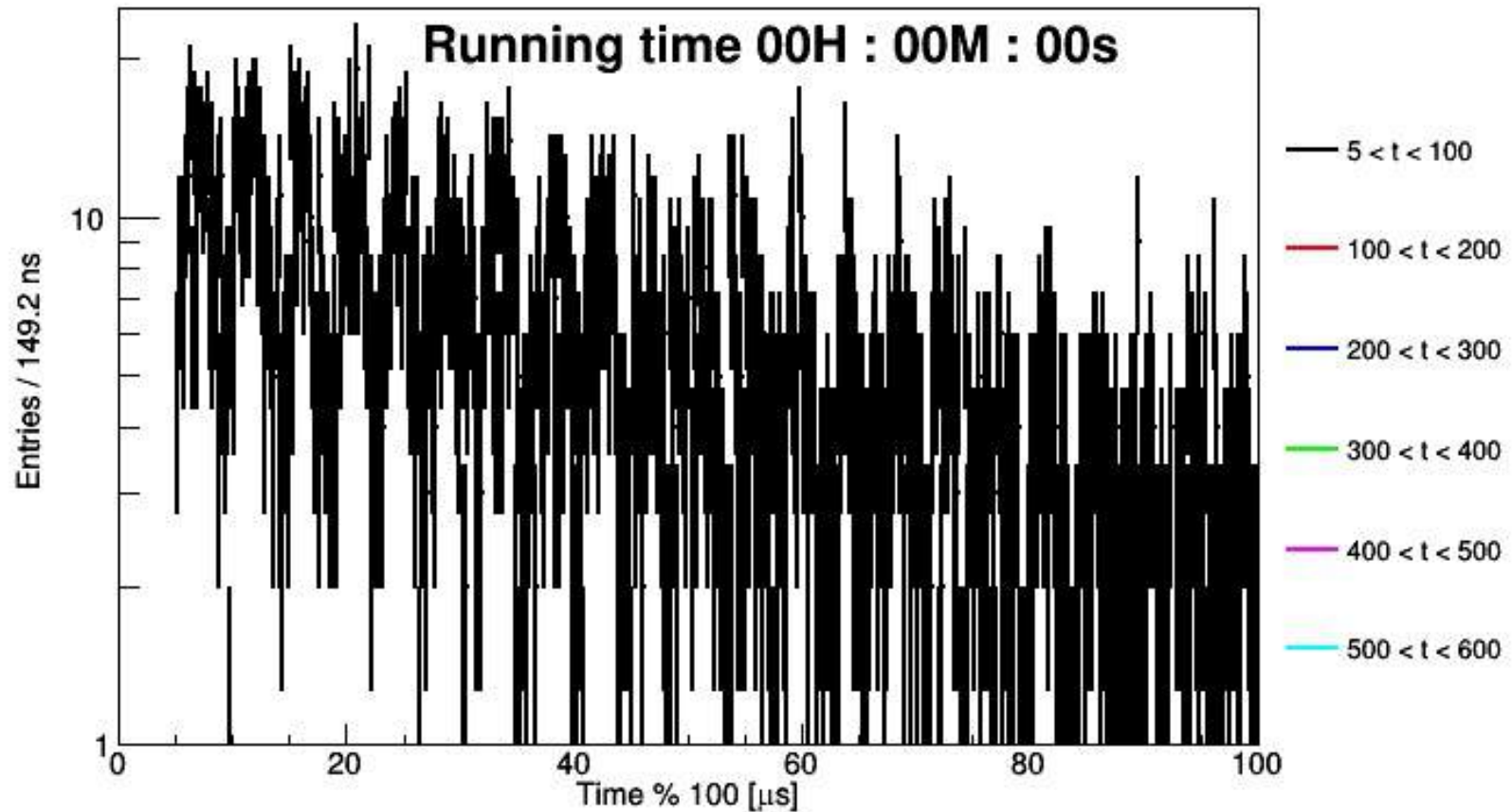
- e^+ preferentially emitted in direction of muon spin



The number of high momentum positrons above a fixed energy threshold oscillates at precession frequency

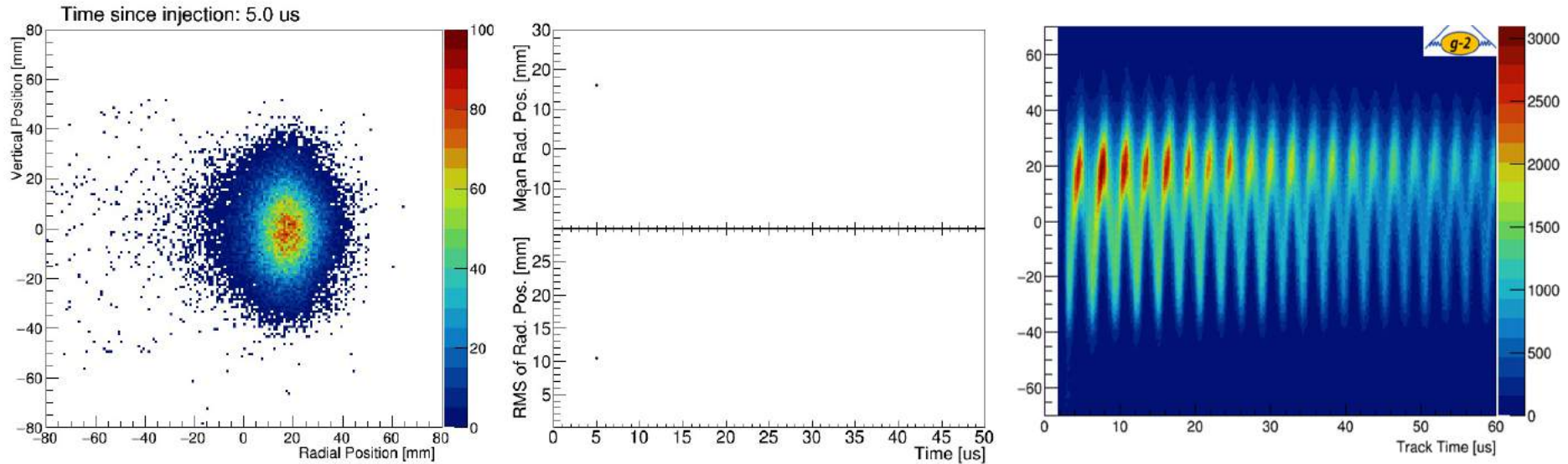
Simply count the number above an energy threshold vs time

Precession in 1 hour of data



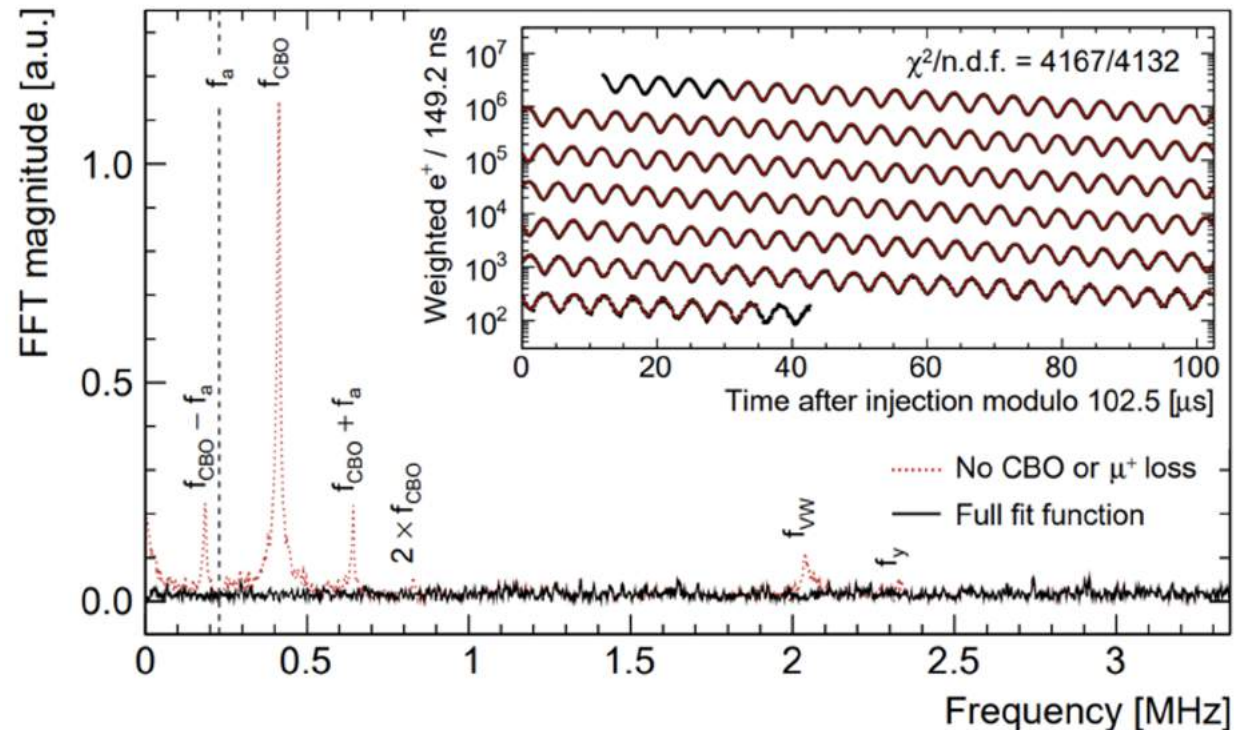
$$N_e(t) \simeq N_0 e^{-\frac{t}{\gamma\tau}} [1 - A \cos(\omega_a t + \phi_a)]$$

Beam Measurements



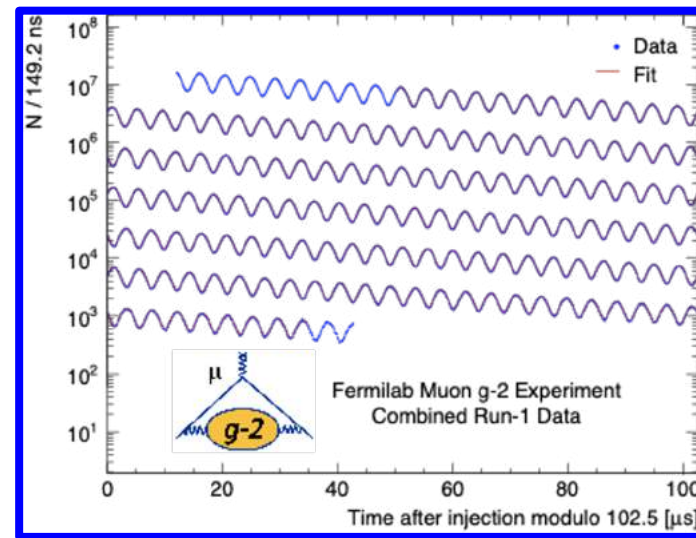
- Use the tracking detectors to measure the decay positrons to infer the decay position
- Muons oscillate radially and vertically at different frequencies, according to the quadrupole strength

- A fourier transform of the residuals to the fit shows contributions from the movements of the beam, pileup and muon losses



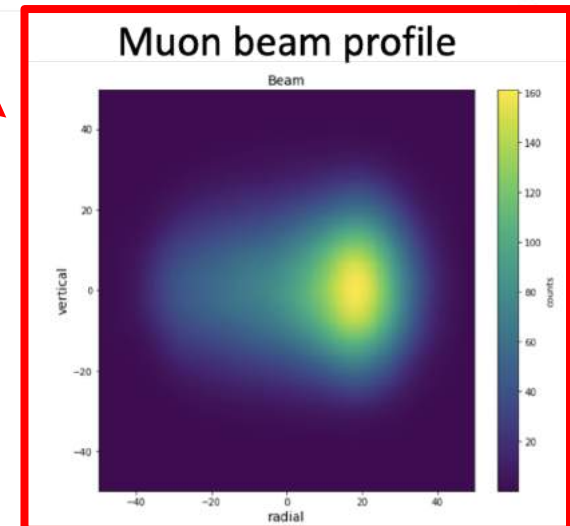
- To account for these effects additional terms are included in the final 24 parameter fit function

Field measurement



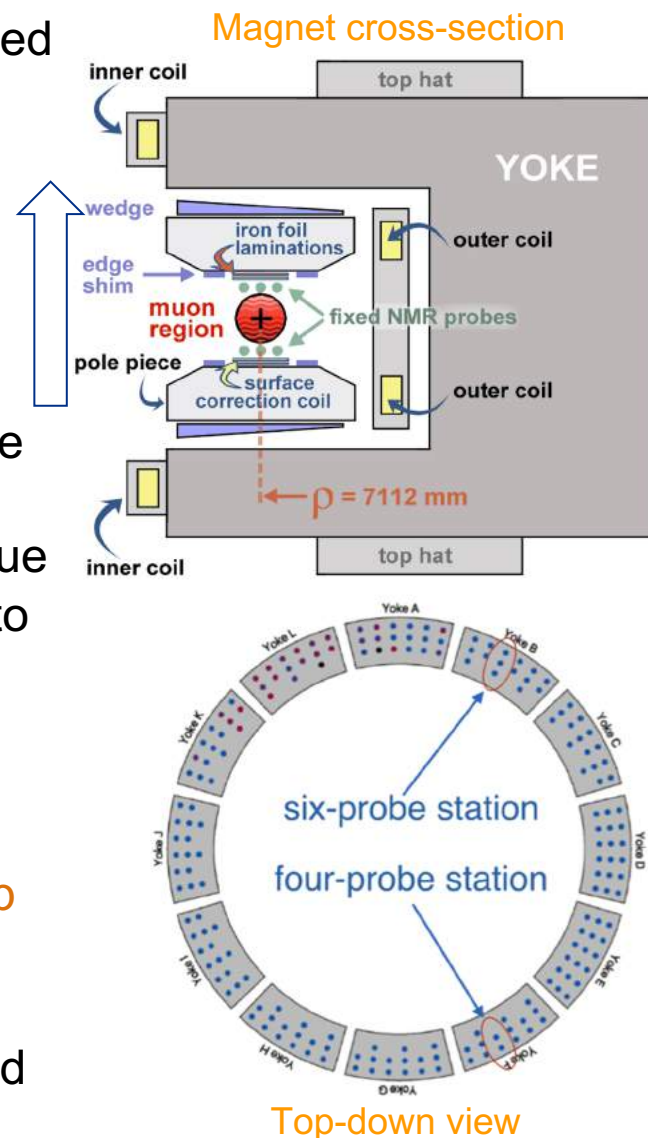
$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

Measuring the magnetic field is the last piece



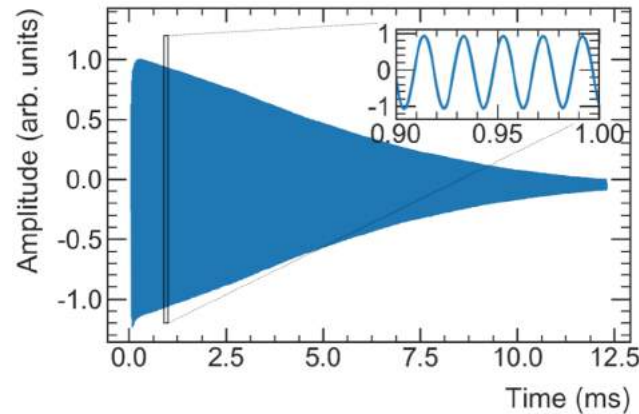
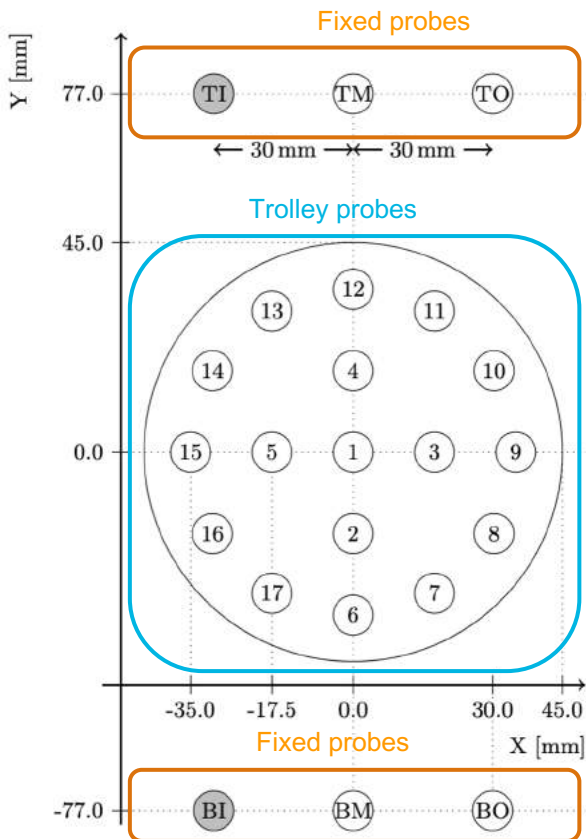
The g-2 storage ring magnet

- 7.112 m radius 'C'-shape magnet with vertically-aligned field $B = 1.45$ T
- Dipole field has ppm-level uniformity (14 ppm RMS across the full azimuth)
- Tiny (ppm) changes in magnet geometry, driven by temperature changes, cause the field to drift over time
- Measured using pulsed NMR – a well-known technique that is routinely used in a wide range of applications to measure magnetic fields at the ppb level
- 378 'fixed' NMR probes, built for this experiment, around the ring measure the drift continuously, and provide feedback to the magnet power supply to keep the dipole (vertical) term constant
- Shimming devices minimise gradients (transverse and azimuthal field components).

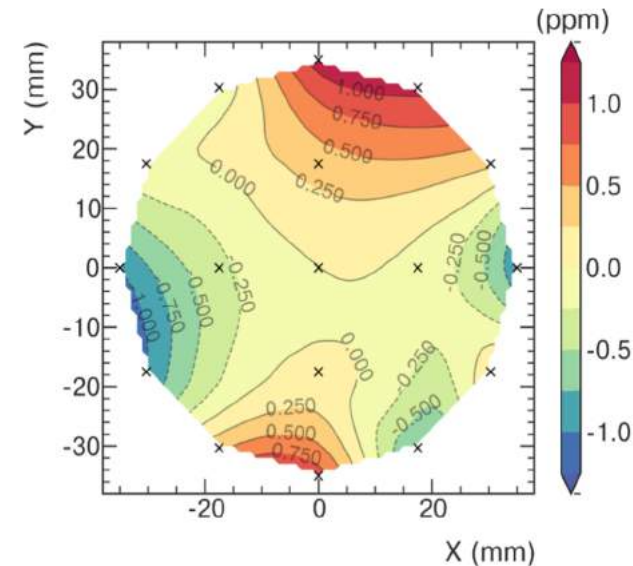


Measuring the field: the NMR Trolley

- An in-vacuum trolley with 17 NMR probes drives around the ring every ~ 3 days, mapping out the field components



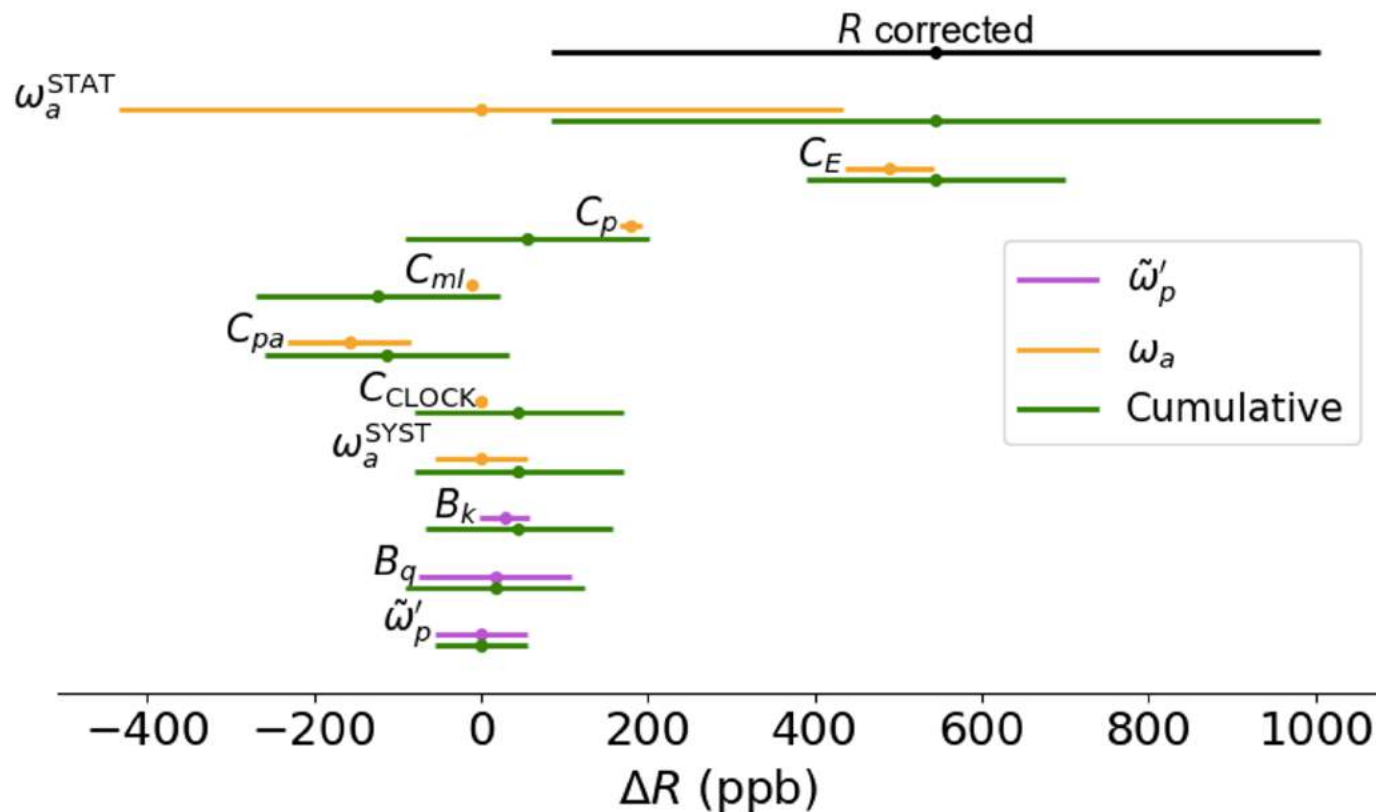
Field measured by extracting frequency from a Free Induction Decay (FID) spectrum



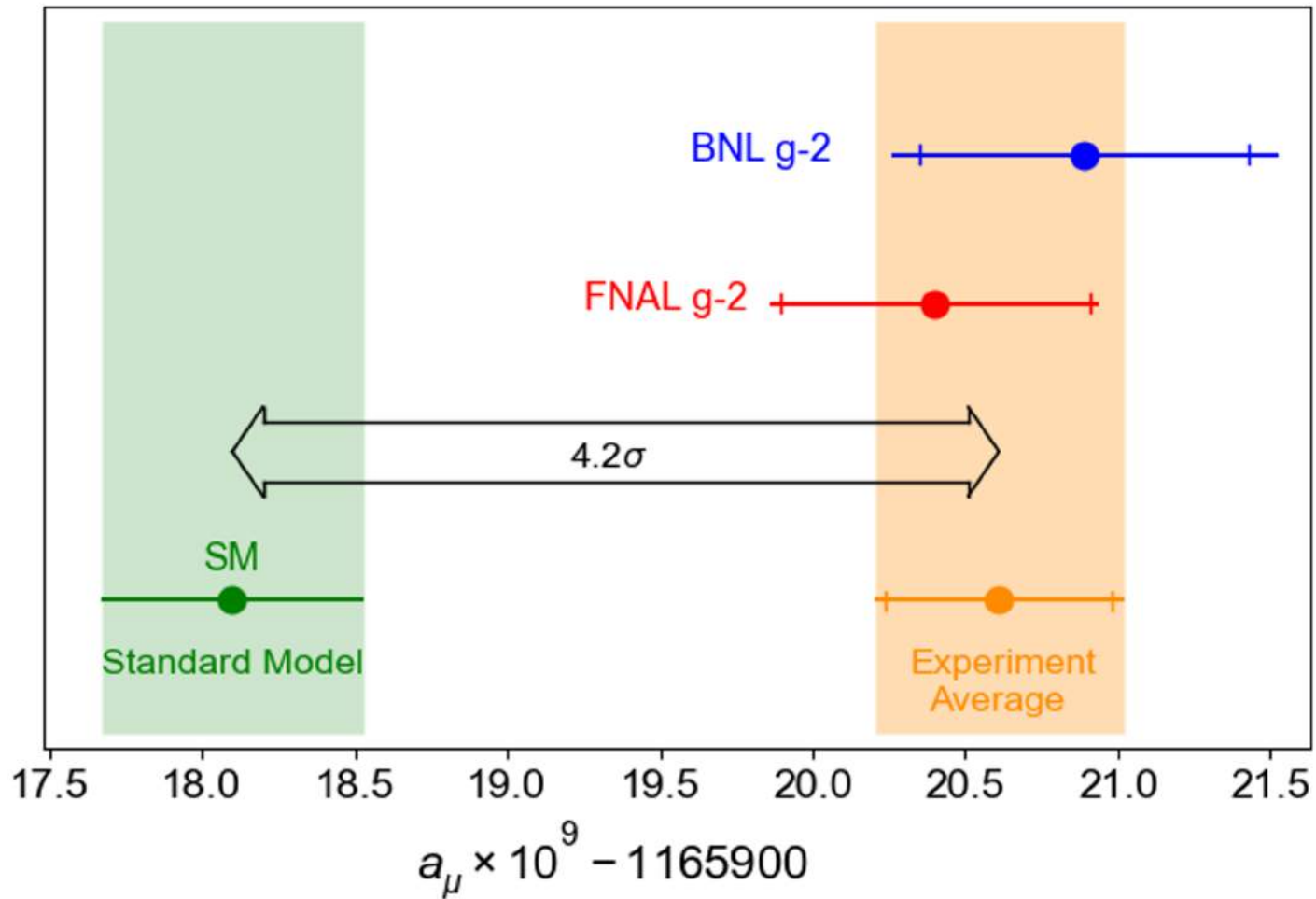
At ~ 8000 azimuthal locations, obtain a field contour plot from the 17 probes

Correcting Measured R

$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

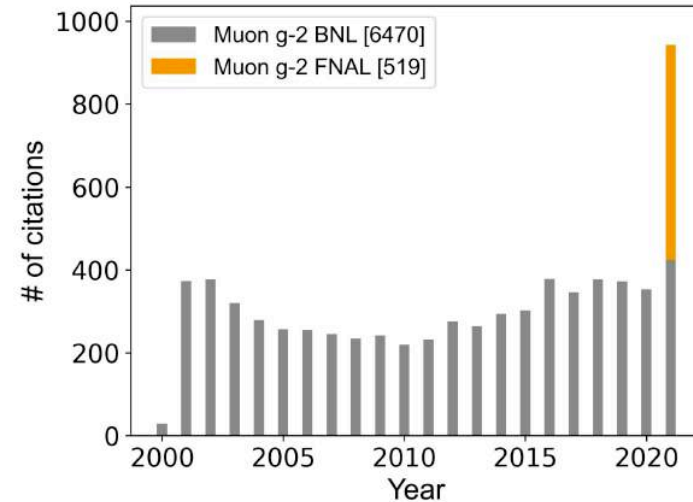
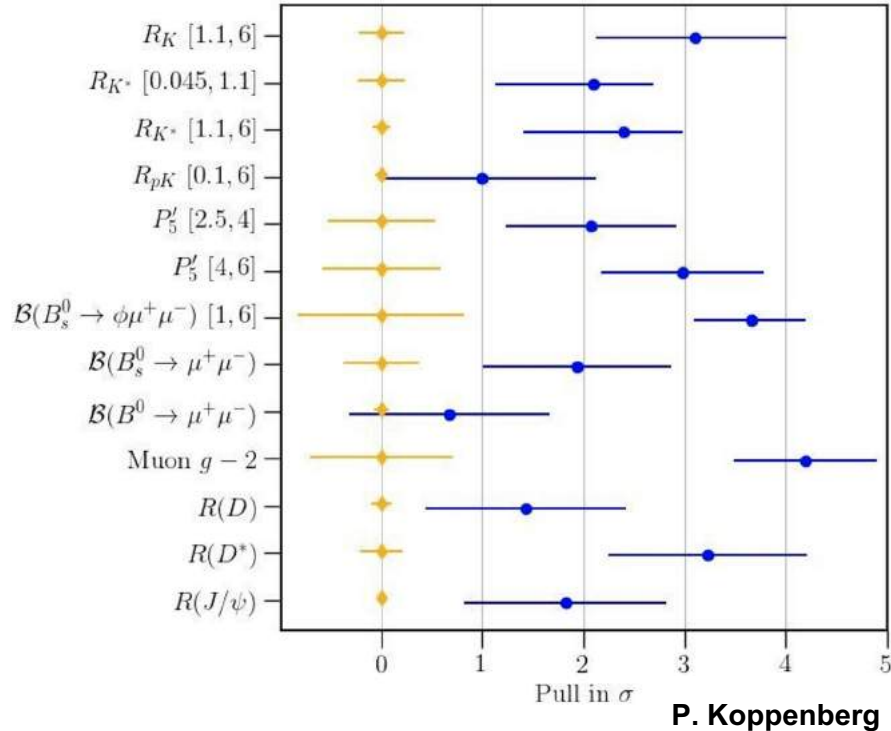


The result



Interpretation

Needs more precision

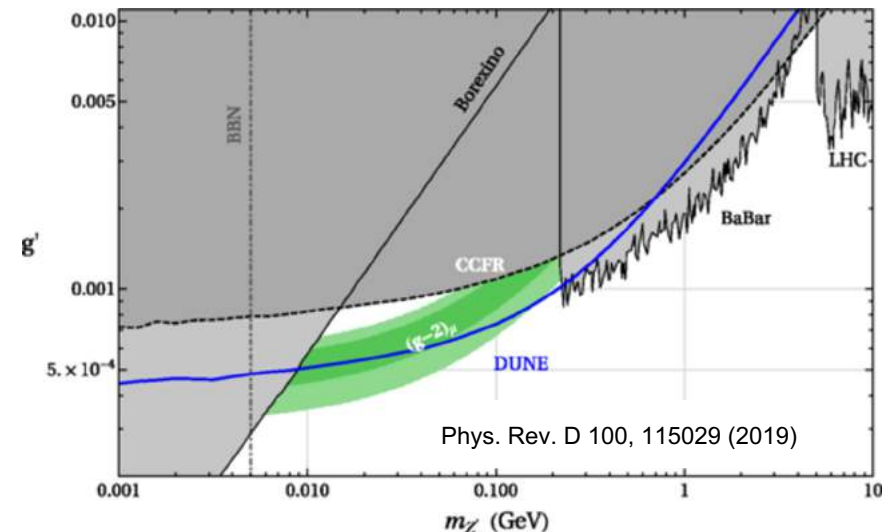


TeV Leptoquarks

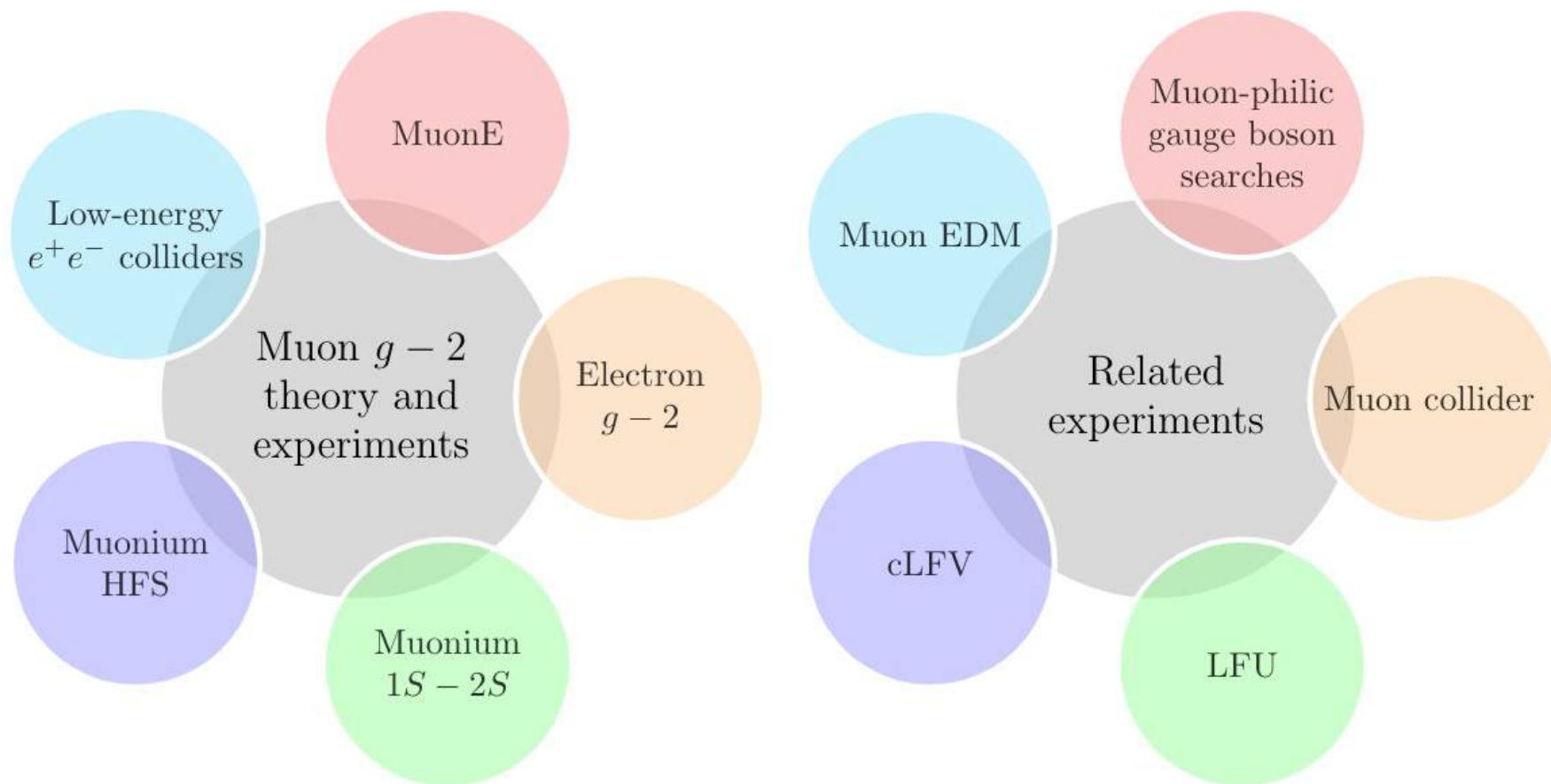
Z', ALPs

LHC evading SUSY

Tweaked Higgs extensions ...



Connected experiments



J-PARC g-2

- Aims to provide an independent g-2 measurement with a completely new approach in terms of muon beam line, storage ring conditions and positron detection.
- Will use a low-emittance 300 MeV/c (ultra-cold) muon beam, eliminating a need for a strong focusing electric field.

general form of spin precession vector:

$$\vec{\omega} = -\frac{e}{m} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta}{2} \left(\vec{\beta} \times \vec{B} \cdot \frac{\vec{E}}{c} \right) \right]$$

BNL E821 approach
 $\gamma=30$ ($P=3$ GeV/c)

$$\vec{\omega} = -\frac{e}{m} \left[a_\mu \vec{B} + \frac{\eta}{2} \left(\vec{\beta} \times \vec{B} + \frac{\vec{E}}{c} \right) \right]$$

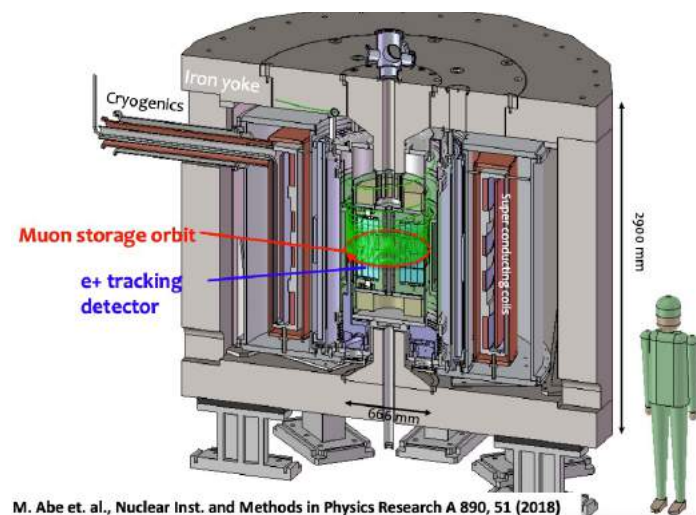
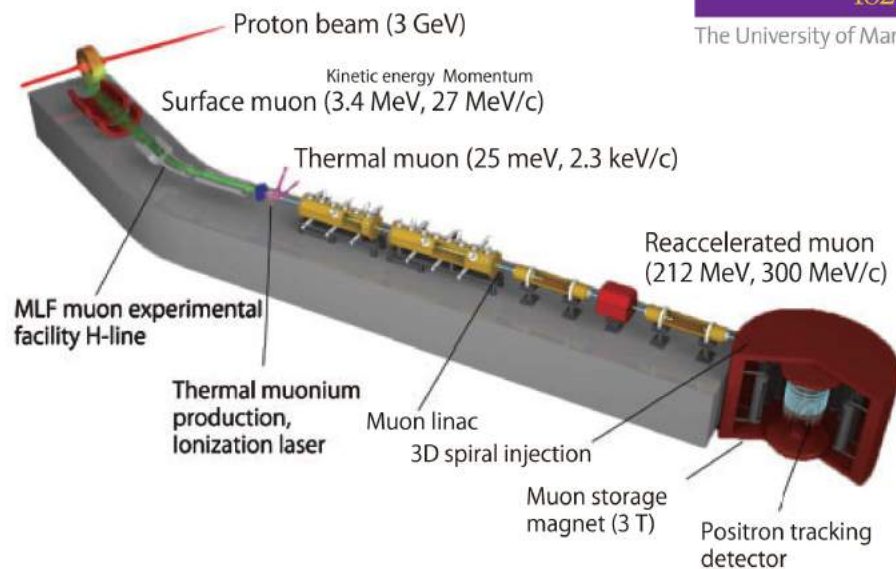
FNAL E989

J-PARC approach
 $E = 0$ at any γ

$$\vec{\omega} = -\frac{e}{m} \left[a_\mu \vec{B} + \frac{\eta}{2} (\vec{\beta} \times \vec{B}) \right]$$

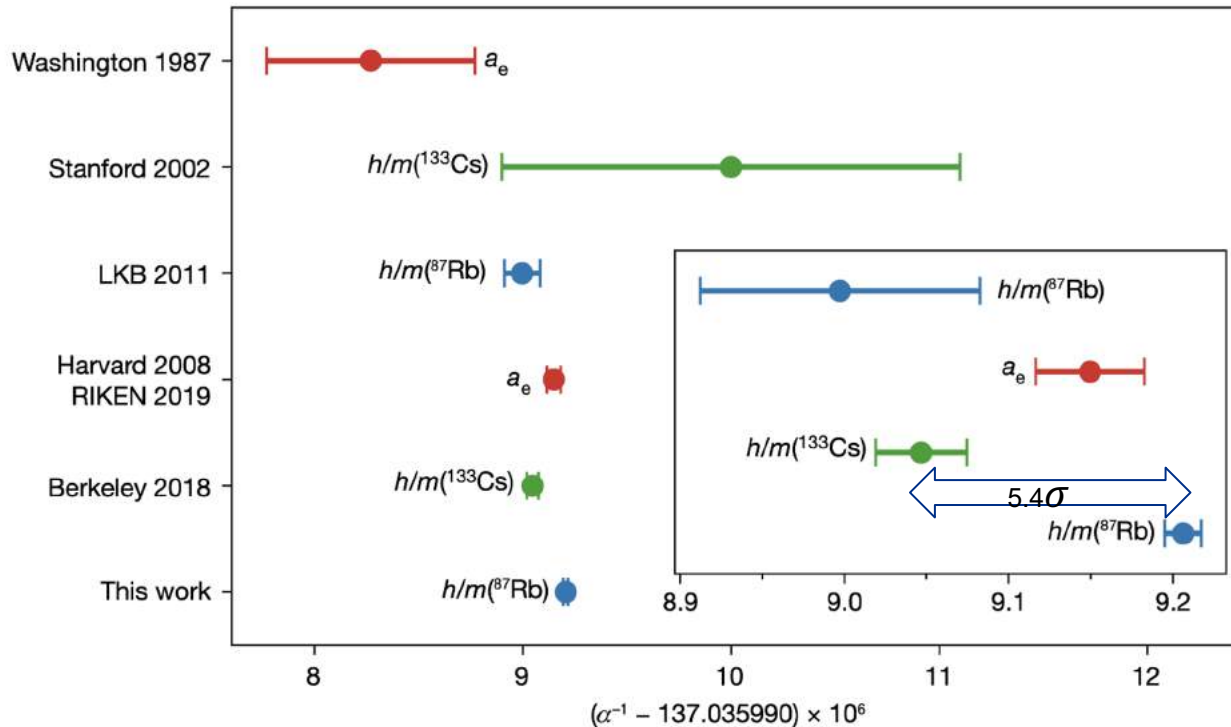
J-PARC E34

Phase-I expected to begin in 2025 with a target precision of about 455 ppb (stat = 450 ppb, sys = 70 ppb)
[Fermilab Run-1 = 460 ppb]



M. Abe et. al., Nuclear Inst. and Methods in Physics Research A 890, 51 (2018)

Electron g-2



Historically it has been used to determine α assuming the SM is correct.

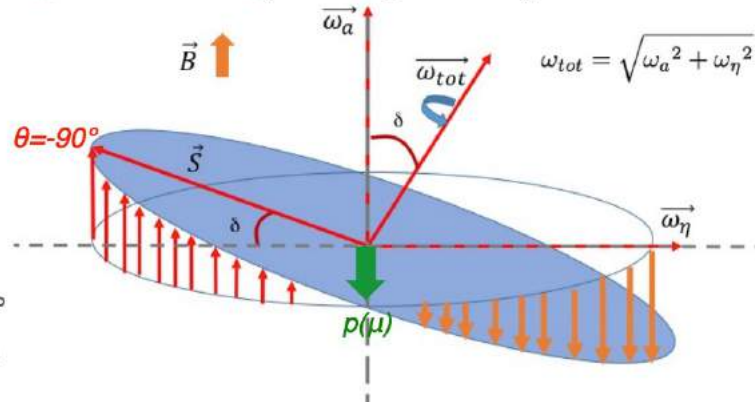
SM $(g-2)_e$ determined from these two α differ from measured $(g-2)_e$ by -2.5σ and $+1.6\sigma$ meaning it cannot presently be used to constrain BSM.

But has potential if can improve experimental $(g-2)$ measurement and α measurement

Any observed muon EDM would be an unambiguous signal of CP violation and BSM physics.
→ And an explanation for the universe's matter-antimatter asymmetry.

$$\vec{\omega}_{a\eta} = \vec{\omega}_a + \vec{\omega}_\eta = -\frac{Qe}{m} \left[a\vec{B} - \left(a - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right] - \eta \frac{Qe}{2m} \left[\frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right] \quad \vec{d} = \eta \left(\frac{Qe}{2mc} \right) \vec{s}$$

- Causes an increase in muon precession frequency
- Precession plane tilts towards center of ring
- Vertical oscillation is 90° out of phase with the a_μ oscillation



Both the Fermilab and J-PARC g-2 experiments will also search for a muon EDM in their storage ring experiments.

BNL: $|d_\mu| \approx 2.5 \times 10^{-19} \text{ e.cm}$

FNAL: $|d_\mu| \approx 3.0 \times 10^{-21} \text{ e.cm}$

J-PARC: $|d_\mu| \approx 1.0 \times 10^{-21} \text{ e.cm}$

PSI – a dedicated muon EDM experiment

- Cancel anomalous precession with matched E-field:

$$E \cong aBc\beta\gamma^2$$

- Spin remains parallel to orbit
- No “contamination” from anomalous spin precession

$$s_z \propto \eta E^* \cdot t$$

$$\vec{\omega} = \frac{q}{m} \left[a\vec{B} + \left(\frac{1}{1 - \gamma^2} - a \right) \frac{\vec{\beta} \times \vec{E}}{c} + \frac{\eta_d}{2} \left(\frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} \right) \right]$$

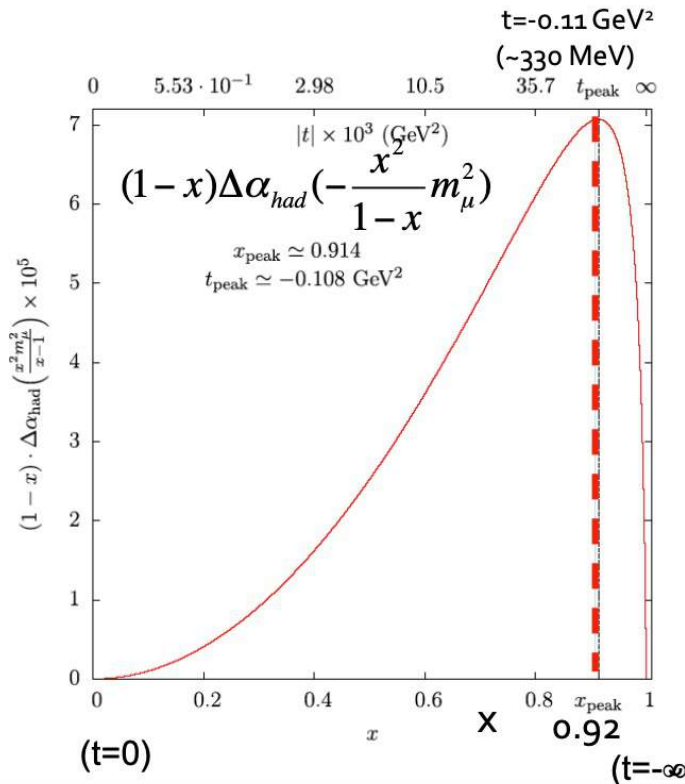
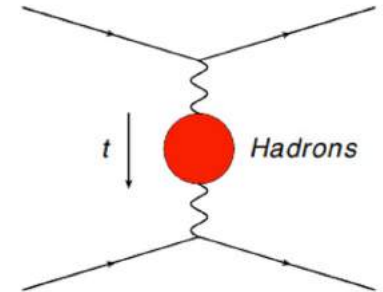
PSI Sensitivity (1 year): $\sigma(d_\mu) < 5 \times 10^{-23} \text{ e.cm}$

arXiv:2102.08838

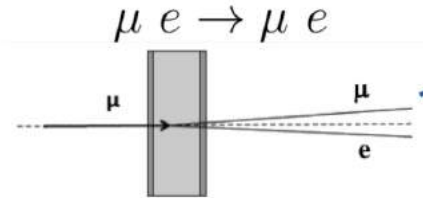
Measurement of $\Delta\alpha_{had}(t)$: hadronic contribution to the running of the electromagnetic coupling constant.

$$a_{\mu}^{HLO} = \frac{\alpha_0}{\pi} \int_0^1 dx (1-x) \Delta\alpha_{had}[t(x)]$$

$$t(x) = \frac{x^2 m_{\mu}^2}{x-1} < 0$$



Extraction of $\Delta\alpha_{had}(t)$ from the «shape» of the $\mu+e \rightarrow \mu+e$ elastic differential cross section



$$\frac{d\sigma_{data}/dt}{\frac{d\sigma_{MC}^{no VP}/dt} = \frac{1}{|1 - \frac{\Delta\alpha_{lep}(t)}{1} - \frac{\Delta\alpha_{had}(t)}{1}|^2}}$$

From theoretical calculation To be measured

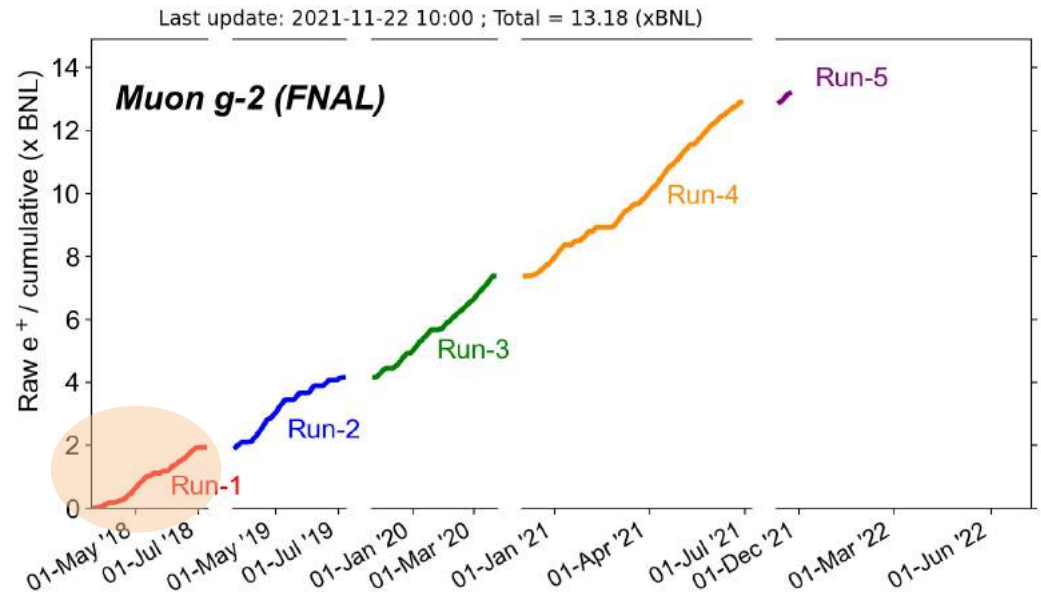
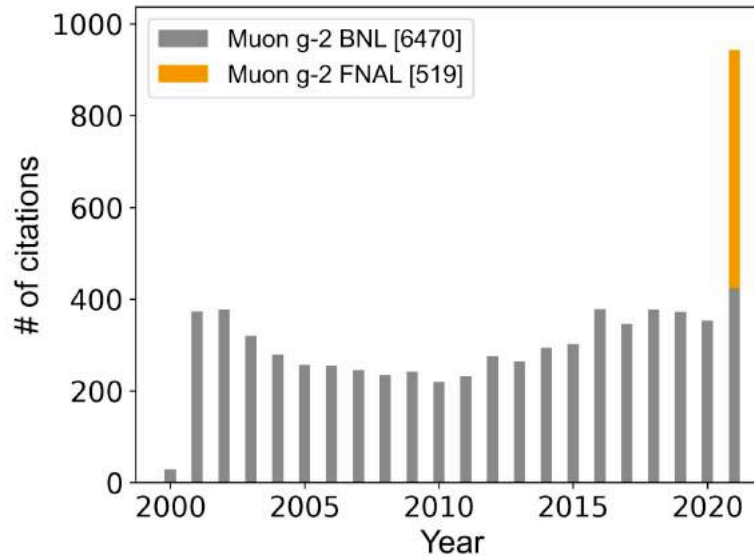
CERN M2 beam of 160 GeV muons allows to cover the whole a_{μ}^{HLO} .

Experiment expected to begin in 2024 with a target precision of about 0.3% stat and 10 ppm systematic accuracy.

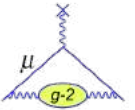
Also requires significant theory improvements that have already been achieved, e.g. a full NNLO MC generator for μ -e scattering (10^{-5} accuracy).

Conclusions

- The analysis of the Run-1 data produced a result with 460 ppb precision.
- Strengthened evidence for deviation from SM in muon $g-2$: 4.2σ tension with the theoretical prediction.
- There is a lot more data to analyse - expect a factor 2 improvement for Run-2/3 analysis, still statistics limited.
- Run-5 will give us a total dataset $\sim x20$ of the first publication and will become systematics limited.



Thank you

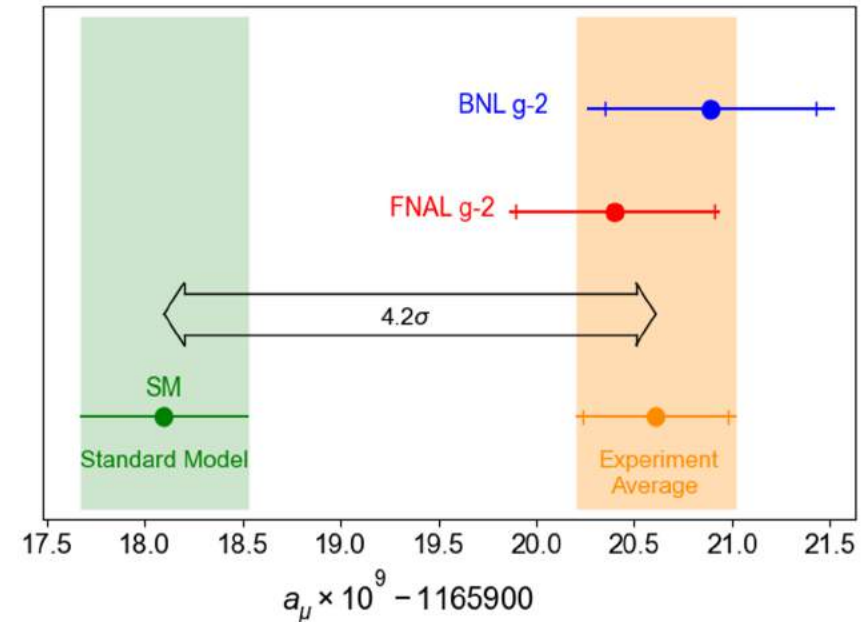


Science and
Technology
Facilities Council



Horizon 2020

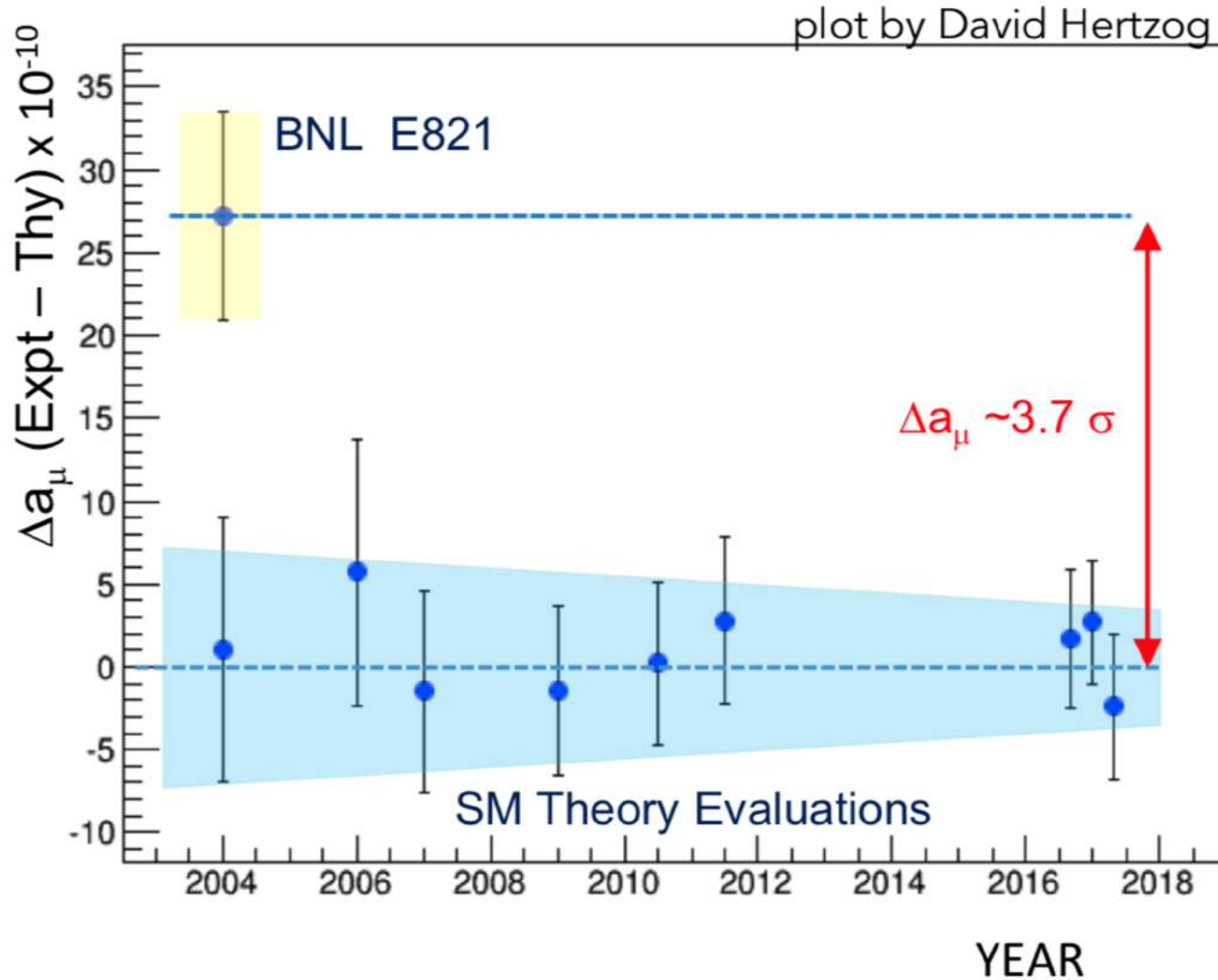
- FNAL Main: [Phys.Rev.Lett. 126 \(2021\) 141801](#)
- FNAL ω_a : [Phys.Rev.D 103 \(2021\) 072002](#)
- FNAL Field: [Phys.Rev.A 103 \(2021\) 042208](#)
- FNAL Beam Dynamics: [arXiv:2104.03240 \(2021\)](#)



- Muon g-2 Theory Initiative (all contributions within): [Phys.Rept. 887 \(2020\) 1-166](#), <https://muon-gm2-theory.illinois.edu/white-paper/>
- BNL Final: [Phys.Rev.D 73 \(2006\) 072003](#)
- Dune/g-2 Z' sensitivity: [Phys. Rev. D 100 \(2019\) 115029](#)
- BSM g-2: [arXiv:2104.03691 \(2021\)](#)

Backups

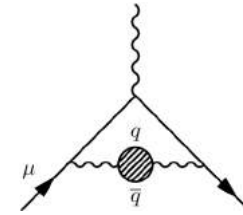
Muon g-2: History of Experiment vs Theory



Muon g-2 in the SM: HVP

$$\Delta a_\mu = 279(76) \times 10^{-11} \rightarrow 2.39(0.65)$$

- Hadronic Vacuum Polarisation - hadronic blob coupled to 2 photons.
- Two point function - in principal, much easier than HLbL.
- Most precisely calculated from $e^+e^- \rightarrow$ hadrons cross section data.



Lattice (error ~ 1.6 ppm of a_μ^{SM})

- Uncertainties dominated by finite volume, discretisation and isospin breaking systematics.

Data-driven (error ~ 0.3 ppm of a_μ^{SM})

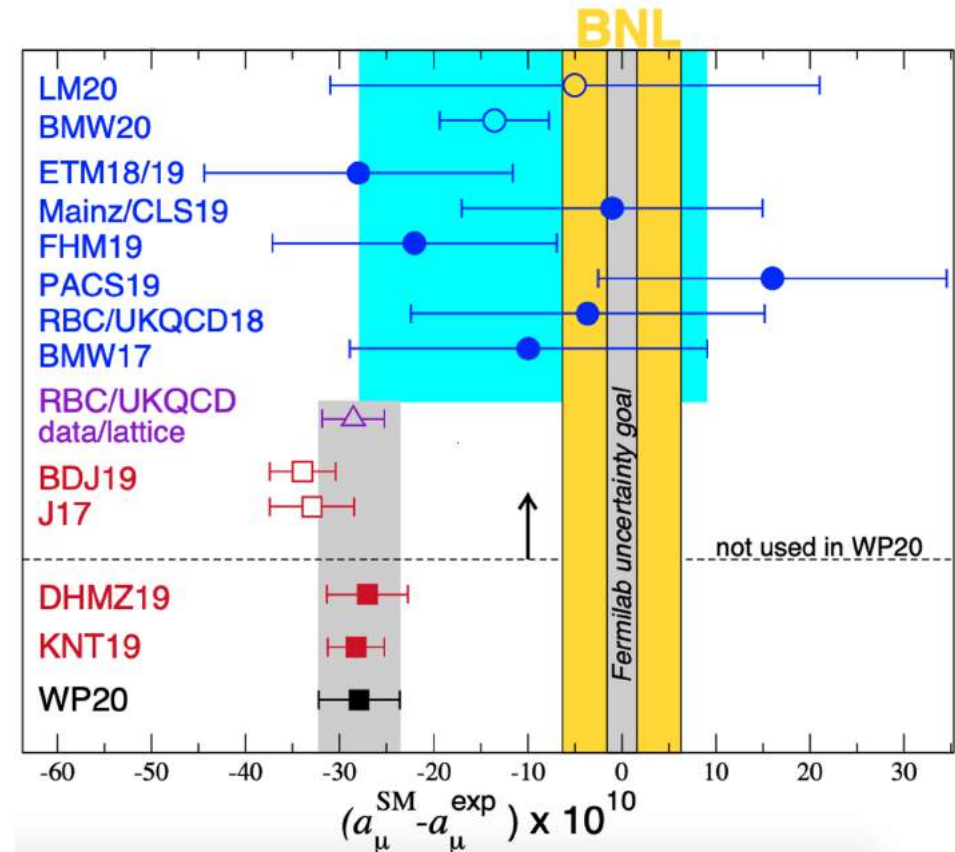
- Cross section data consistently combined and input into dispersion integral:

$$a_\mu^{\text{LOHVP}} = \frac{1}{4\pi^3} \int_{s_{th}}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

- Several groups have achieved this (most precisely in the UK).

Recommended Muon g-2 TI value from data-driven result:

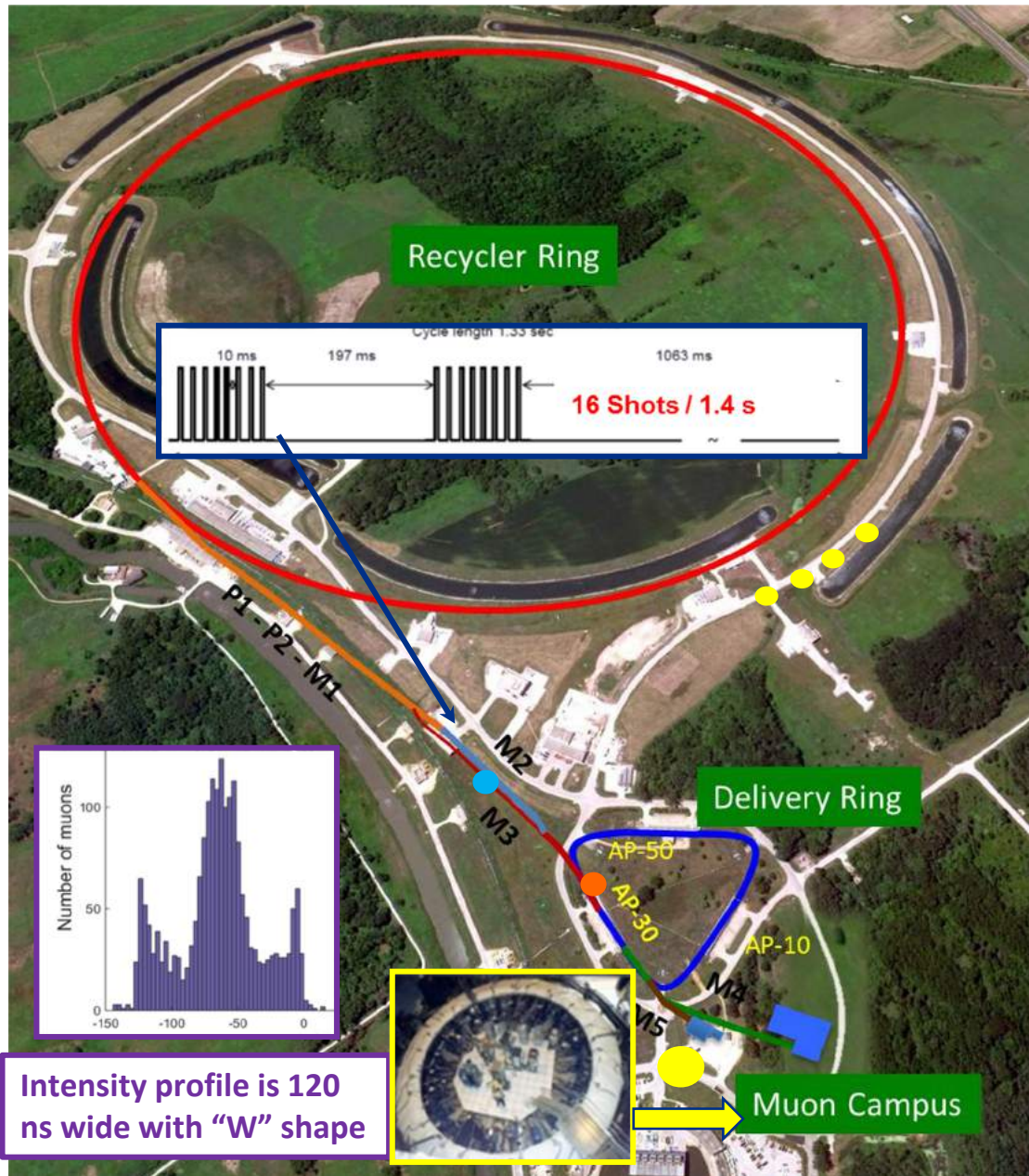
$$a_\mu^{\text{HVP}} = 6845(40) \times 10^{-11}$$



See Thomas Lenz's talk, tomorrow, 12.30pm:
"Experimental input to the Standard Model prediction of g-2"

Muons at Fermilab

- Deliver two 4×10^{12} 8 GeV proton batches to the Main Injector Recycler (graphic shows one)
- Batches are split into four bunches
- One bunch extracted every 10 msec to AP0 target hall
- 3.1 GeV pions are selected and focused by Li lens
- Transported through dense FODO lattice to Delivery Ring
- Several passes around Delivery Ring to remove protons by time-of-flight.
- Muons are focused and injected into the Muon g-2 storage ring.
- Whole cycle repeats twice every 1.4s



Muons at Fermilab



Lower instantaneous rate but
larger integrated rate than BNL



$\sim 10,000\mu^+$ (from 10^{12} p) at 3.1 GeV every 10 ms

(g-2): $\frac{1}{3}$ of proton cycles, neutrino expts: $\frac{2}{3}$

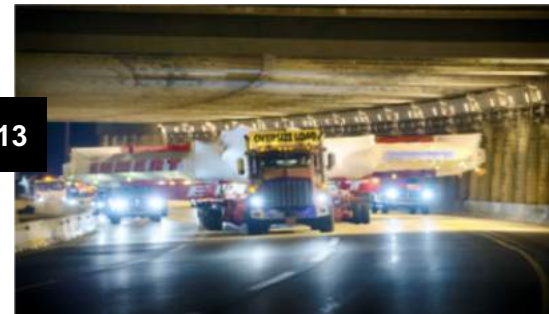
Extra 900m of instrumented beamlines

4 years to build (2 years magnet 'shimming' ...)

May 2013



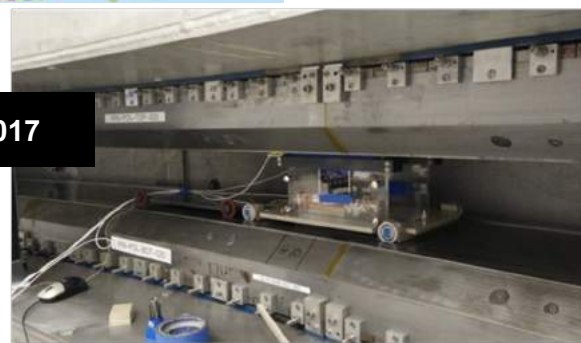
2013



2015



2016-2017



May 31 2017 (g-2)

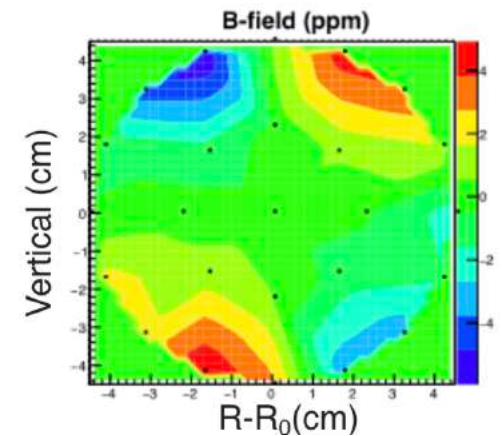
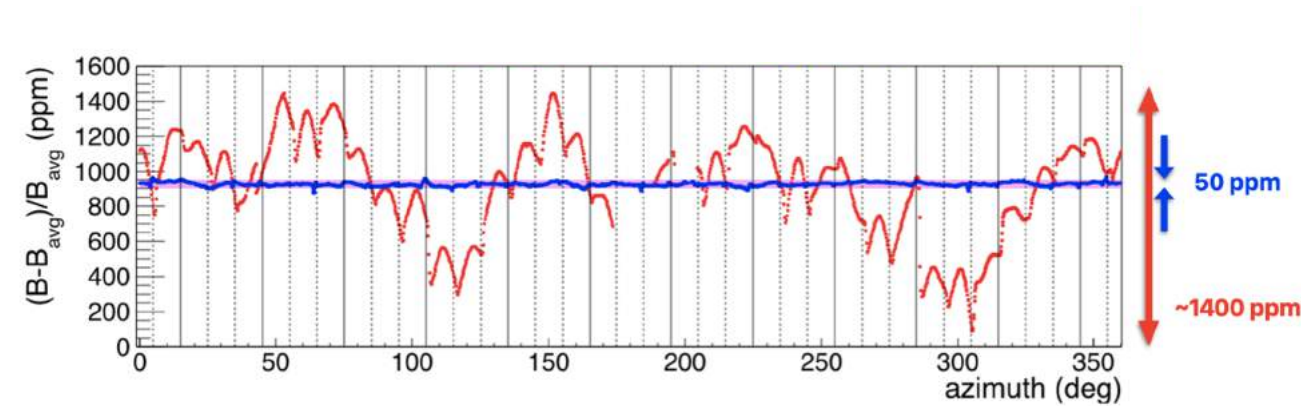


Run-1 data taking started Feb. 2018

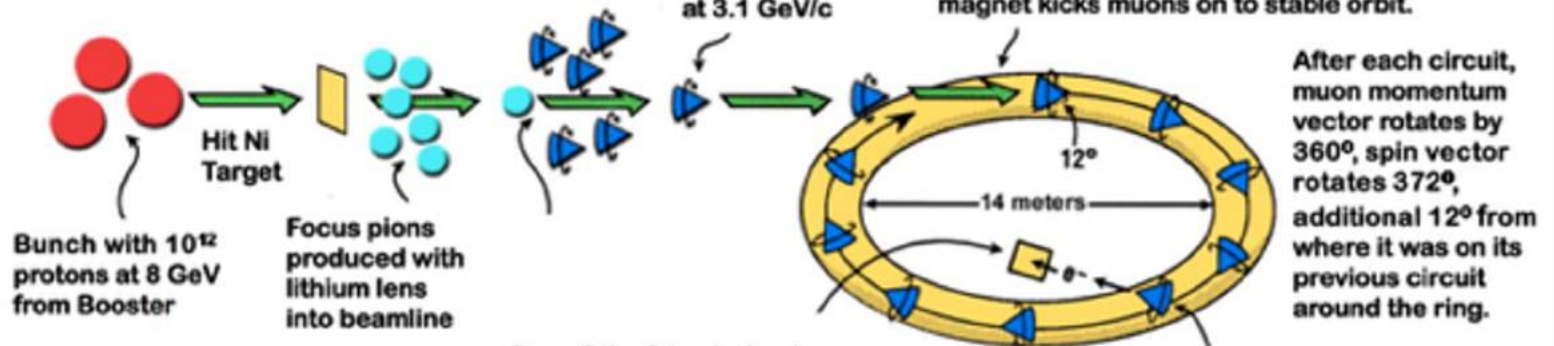
The g-2 ring



Magnetic field uniformity 3 times better than the goal (BNL)



Overview of the g-2 experiment

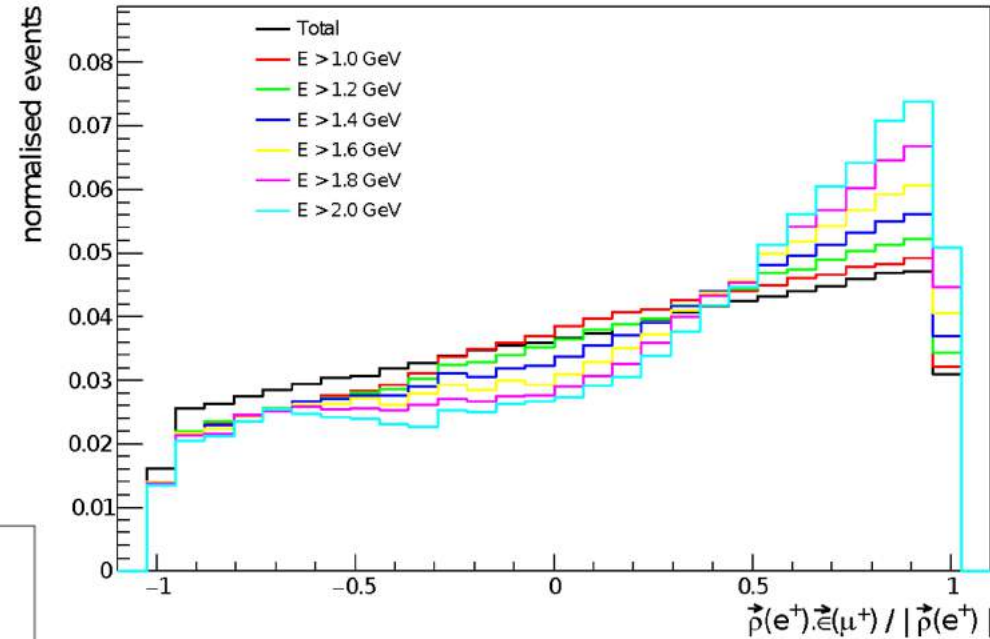
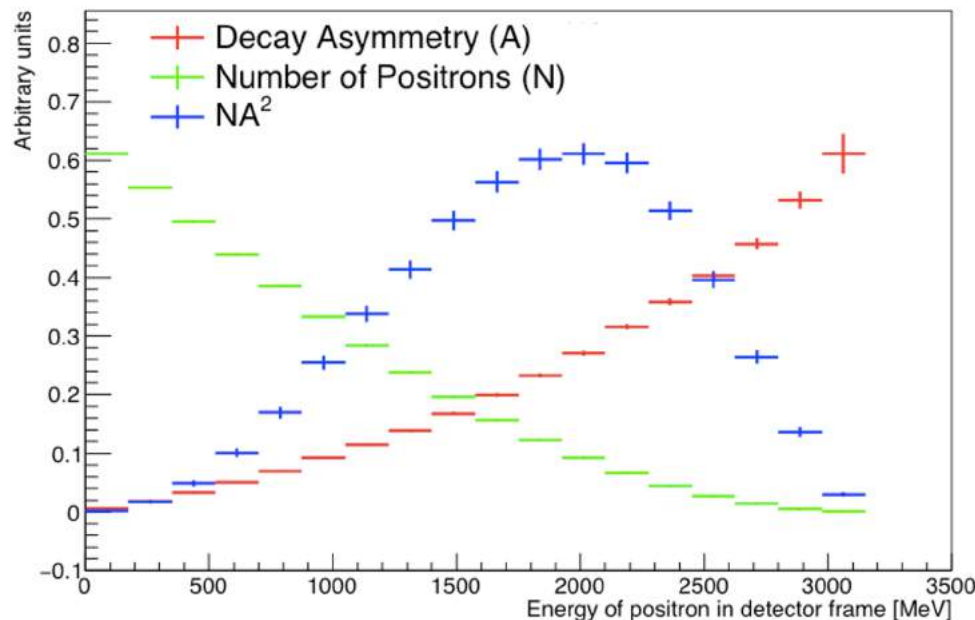
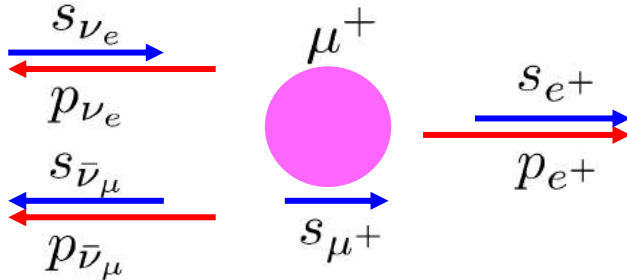


Fermilab statistics advantages

- Long decay channel for $\pi \rightarrow \mu$
 - Reduced p and π in ring
 - Factor 20 reduction in hadronic flash
 - 4x higher fill frequency than BNL
- 21 times more positrons detected than at BNL

Muon decay

- e^+ preferentially emitted in direction of muon spin



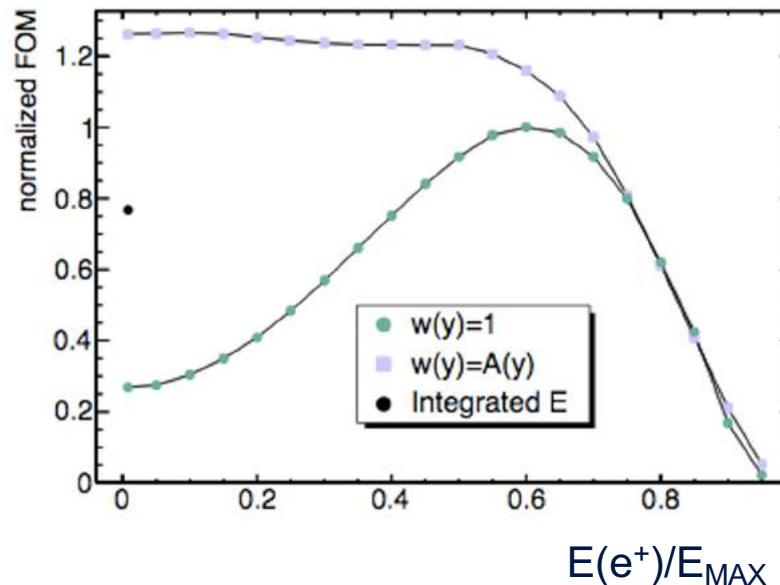
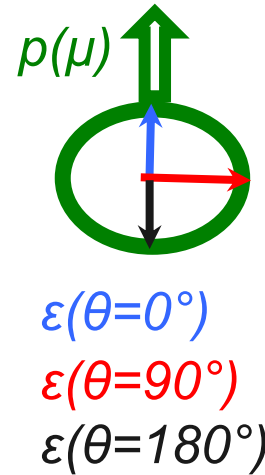
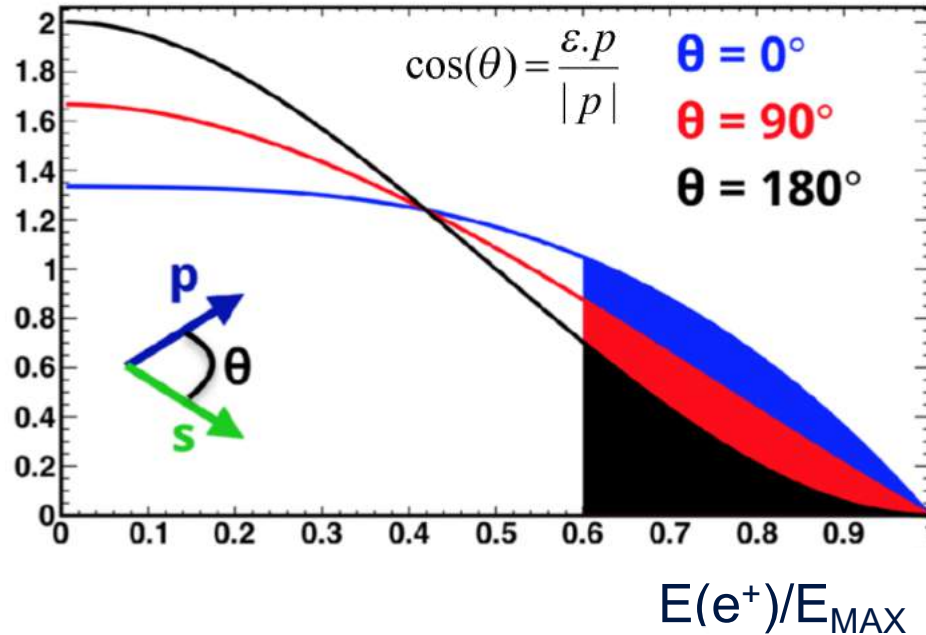
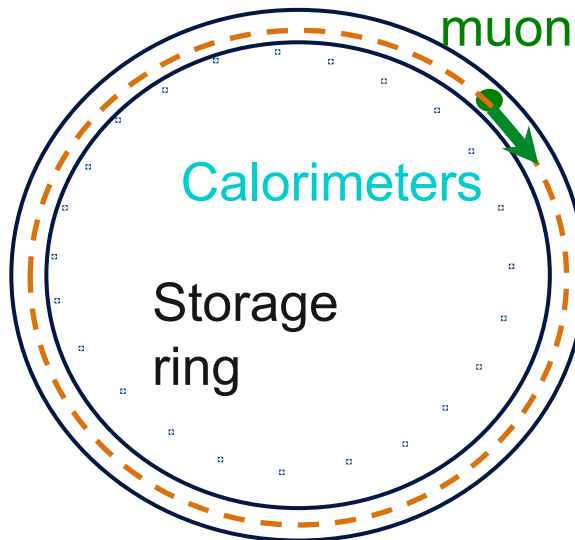
$$\frac{\Delta\omega_a}{\omega_a} \propto \frac{1}{\gamma_\mu \sqrt{NA^2}}$$

- Asymmetry is larger for higher momentum e^+
- Optimal cut at ~ 1.7 GeV

Measuring ω_a

The number of high momentum positrons above a fixed energy threshold oscillates at precession frequency

TOO MUCH IN THIS SLIDE



Simply measure the time and energy of decay positrons

Count the number above at each energy and weight based on asymmetry

- Injected beam has a small vertical component
- Need to use electrostatic quadrupoles to focus the beam vertically

$$\vec{\omega}_a = \frac{e}{mc} \left[a_\mu \vec{B} - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) \vec{\beta} \right]$$

- This introduces 2 additional terms reducing the precession frequency
- **We can minimise the first by choosing $\gamma = 29.3$ to give $p_\mu = 3.1 \text{ GeV}$**
- For a 1.45T field, this sets the radius of the ring to 7.11m
- However we now have 2 corrections to make to a_μ because:

Not all muons are at the 'magic' momentum of 3.1GeV

Vertical momentum component aligned with B field

E-field
correction

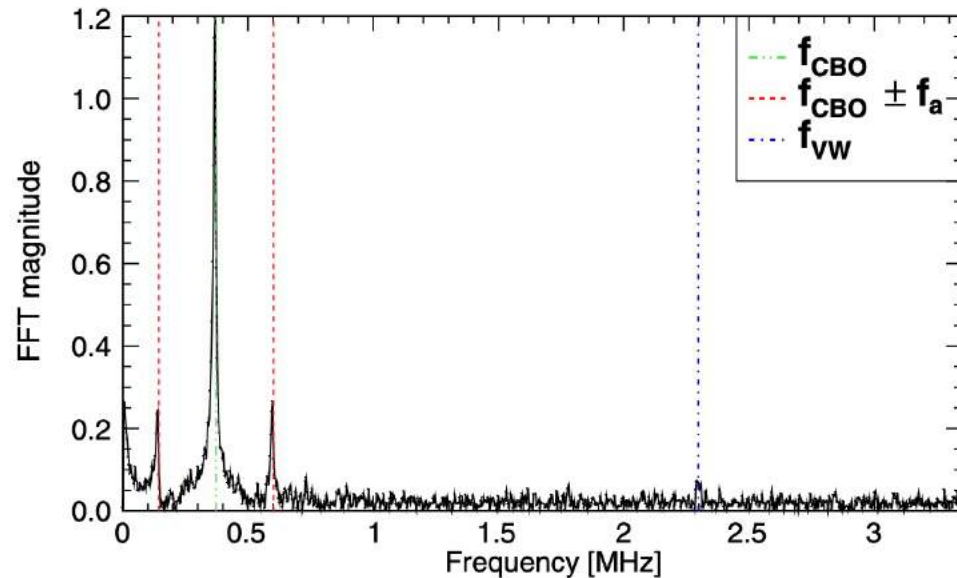
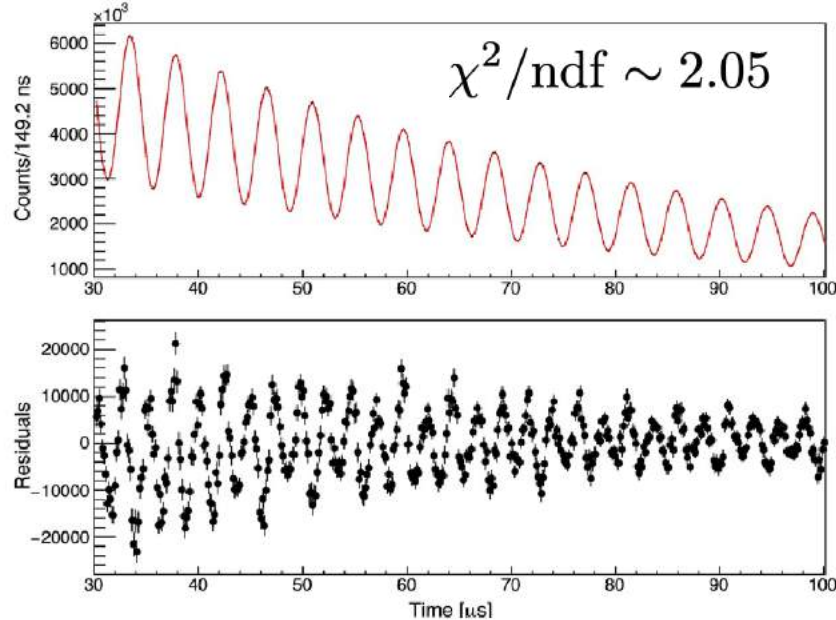
$$C_E = \frac{\Delta\omega_a}{\omega_a}$$

Pitch
correction

$$C_P = \frac{\Delta\omega_a}{\omega_a}$$

- Both corrections depend on the quadrupole field strength, and are $< 0.5 \text{ ppm}$

Results of 5 parameter fit ...

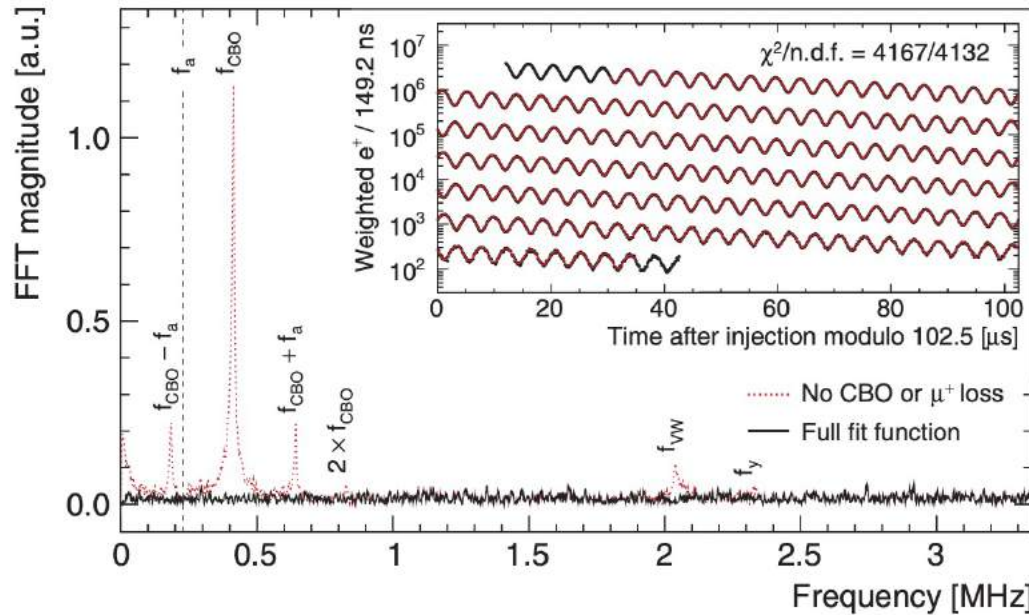


Add additional 17 terms in fit to describe:

- Muons lost from storage ring not by decaying
- Pileup (concurrent, multiple e^+ in same calorimeter crystal)
- Vertical and radial beam motion

And get $\chi^2/\text{ndf} \sim 1.008$

Resulting 22 Parameter Fit



Phys. Rev. D 103, 072002 (2021)

Statistical uncertainty from this fit : 434 ppb

Largest correction to data is : 489 ppb (total correction is 456 ppb)

Total systematic uncertainty is : 157 ppb (aim was 100 ppb)

Deviation from SM (with BNL) : 2400 ppb

Two Sets of corrections to get to (g-2)

1. Our NMR frequencies are multiplied by: $\frac{\mu_p}{\mu_e} \frac{m_\mu}{m_e} \frac{g_e}{2}$ to determine the B-field

The uncertainty on this correction is very small (24 ppb) from CODATA and external to experiment

2. Corrections because the beam and apparatus are not perfect

- beam has a few MeV momentum spread
- beam has a small vertical momentum component
- mean position (and rms) of beam drifts slowly over measurement period
- there are transient magnetic fields in addition to the 1.45 T storage ring field
- muon population is depleted other than by decay & is momentum dependent

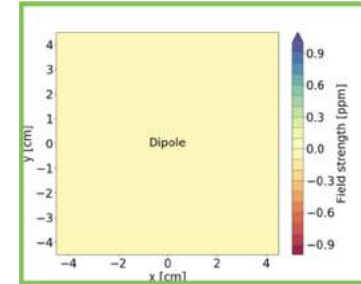
The uncertainty in these corrections largely determines the systematic uncertainty

Spatial dependence of B

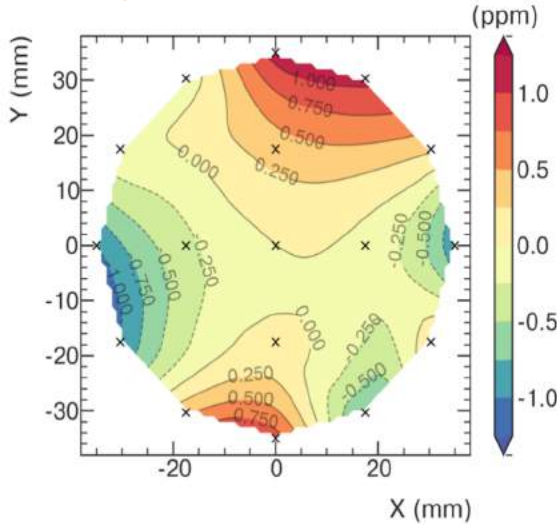
- Extract terms from a multipole (m) expansion of B in r and θ :

$$B \approx B_y = A_0 + \sum_{n=1} \left(\frac{r}{r_0} \right)^n (A_n \cos(n\theta) + B_n \sin(n\theta))$$

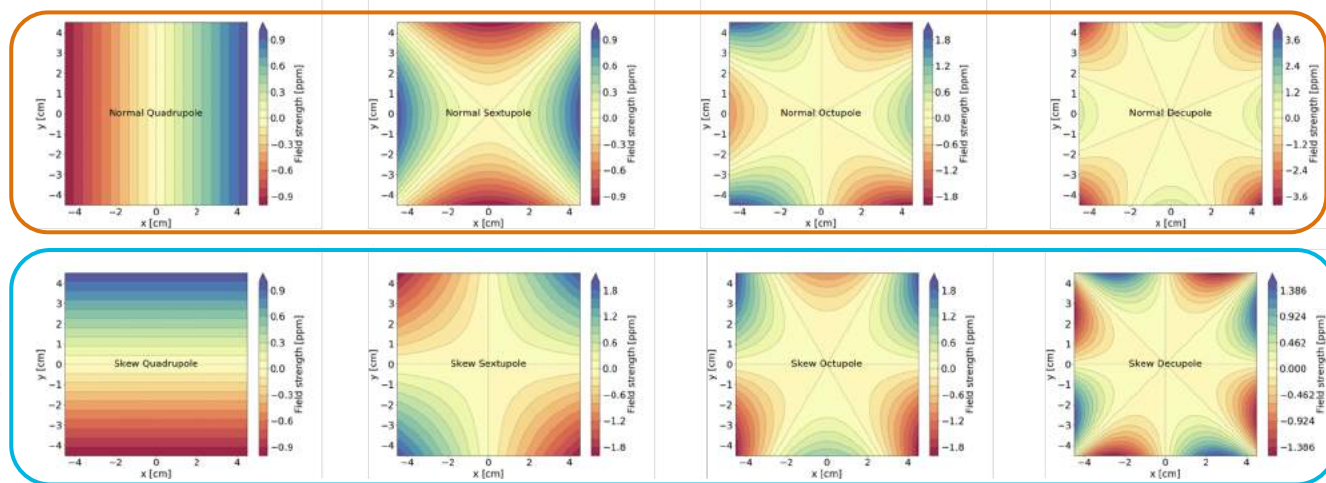
Dipole (m1)



Field gradients in an azimuthal “slice”



“Normal” terms: m2, m4, m6, m8, ...

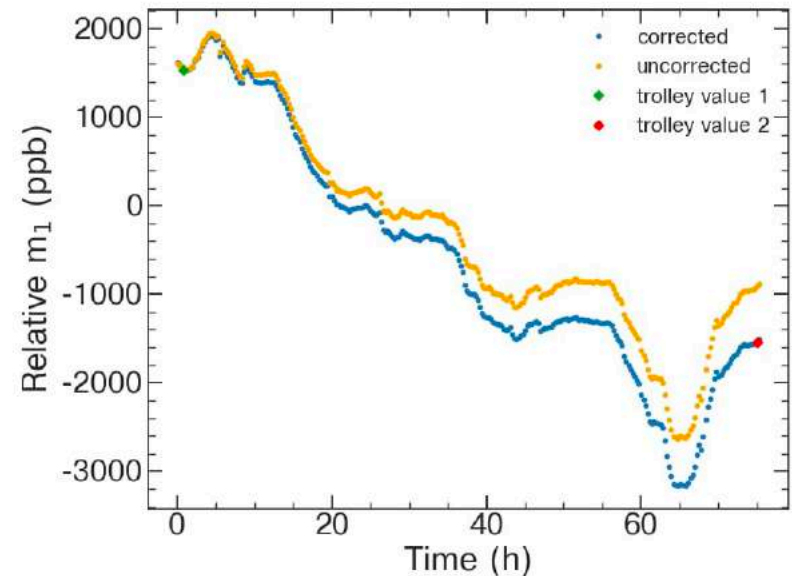


“Skew” terms: m3, m5, m7, m9, ...

- Trolley: Fit the 2D contour plot to extract the multipole terms (m1, m2, m3, ...)
- Fixed probes: extract terms from geometric combination of probe frequencies
- Fixed probes can track m1, m2, m3, m4 only

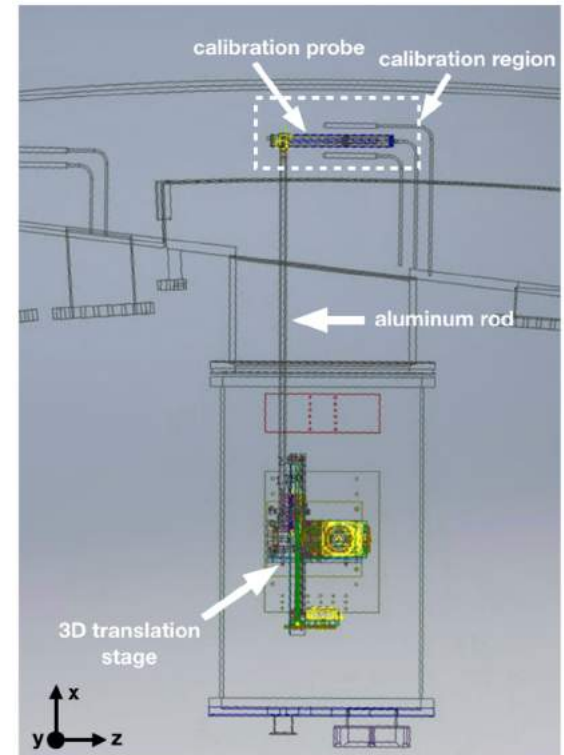
Interpolating between trolley runs

- Need to know the field experienced by the muons, but the trolley cannot take data when the muons are present. **One trolley run takes 3 hours, every ~3 days.**
- Fixed probes take data continuously during muon fills. Use this data to **interpolate** between trolley runs.
- There are 72 fixed probe 'stations' around the ring, every ~5 degrees
- The fixed probe measurements are calibrated using the trolley measurements both times the trolley passes
- Calibration drifts over time, due to changes in higher-order terms that cannot be tracked by the fixed probes
- Leads to the **tracking error uncertainty** (22 - 43 ppb in the run 1 datasets)



Absolute calibration

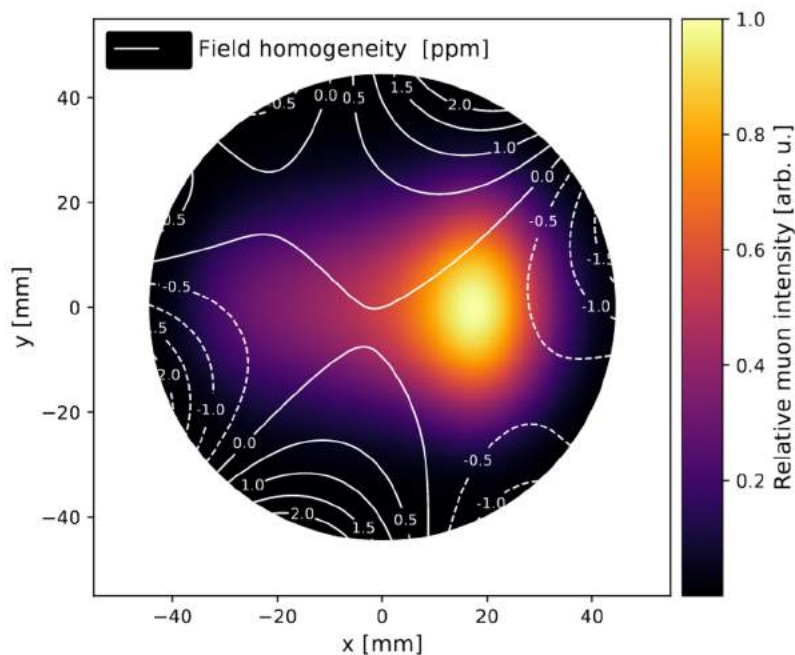
- Trolley and fixed NMR probes use **petroleum jelly** as the proton sample. Chosen for low volatility.
- Must calibrate with protons in a water sample (measurement standard) in order to measure a_μ
- A dedicated **calibration probe with a cylindrical H_2O sample** is installed inside the vacuum chamber.
- In a dedicated calibration campaign, trolley and calibration probes switch places to repeatedly measure the same field in the same place
- Calibration probe is calibrated against a different probe with a **spherical water sample**.
- Both calibration probes were cross-checked with a **spherical 3He sample** (different systematics)



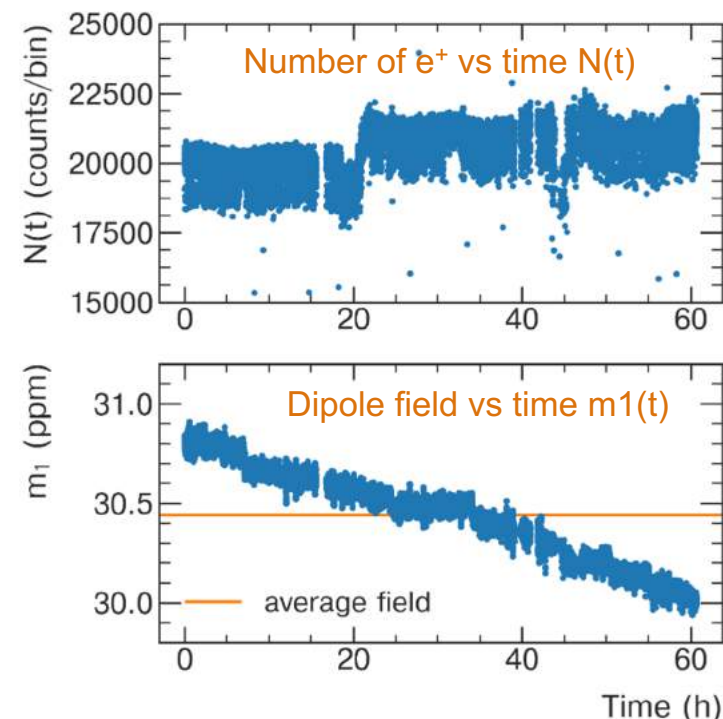
Agreement between all three calibration probes at 10 ppb level

The muon-weighted field

- To obtain the field experience by the muons, the magnetic field distribution as a function of time must be weighted by:
 - The number of muons as a function of time, $N(t)$
 - The beam distribution as a function of time



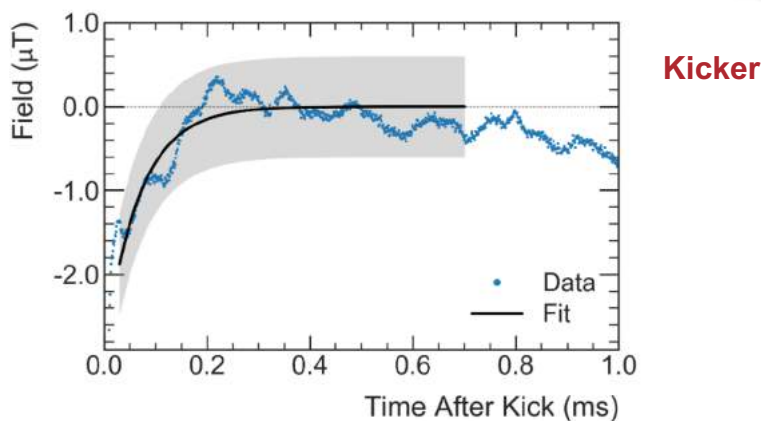
The field is weighted by the 2D beam distribution. An average beam distribution for every 3 hours is used.



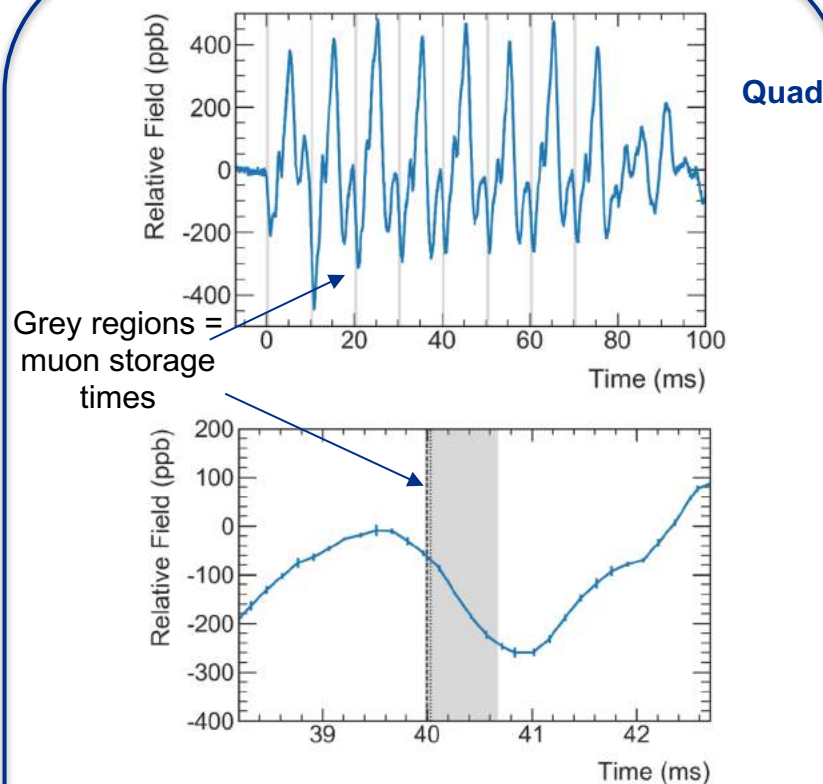
Measured field (every 1.7 s) is weighted by the number of detected e^+

Transient fields

- Largest uncertainties come from “fast transient” fields generated by the pulsed systems (kickers and quads)
- Muons experience a field change which the fixed probes do not see (due to shielding)
- Effects were measured separately during dedicated measurement campaigns.



- Kicker pulse of **22 mT for 150 ns** just after muon injection.
- Field change caused by residual field after kicker pulse. Muons present from **30 μs to 700 μs** after the kick (fit region)
- Kicker correction: **-27 (37) ppb**



- Measured with a dedicated in-vacuum NMR probe located between quad plates during pulsing
- Quad correction: **-17 (92) ppb**

Systematic Uncertainties

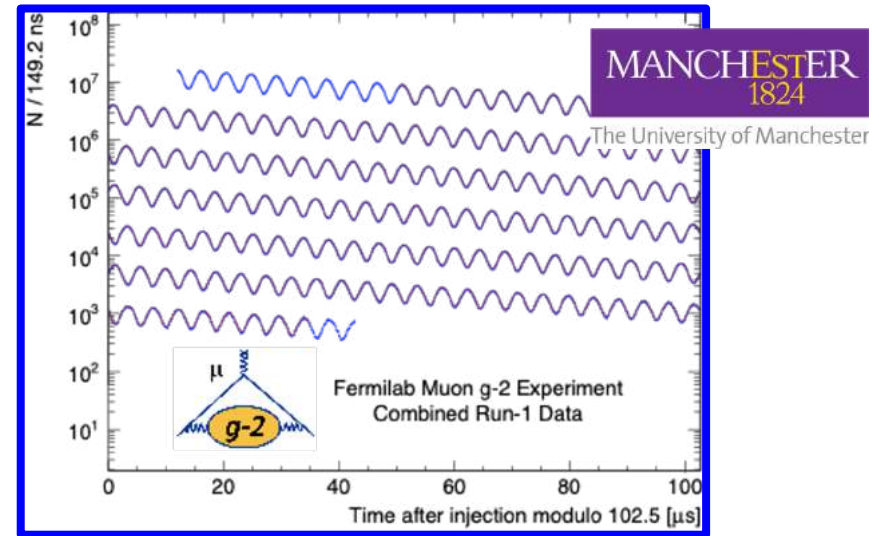
~ 80 effects considered significant in determining the systematic uncertainty.
Dedicated runs taken for some of them e.g. at different beam momentum.
Documented in 98 pages of PRDs.

Total systematic uncertainty 157 ppb. Those above 30 ppb are below

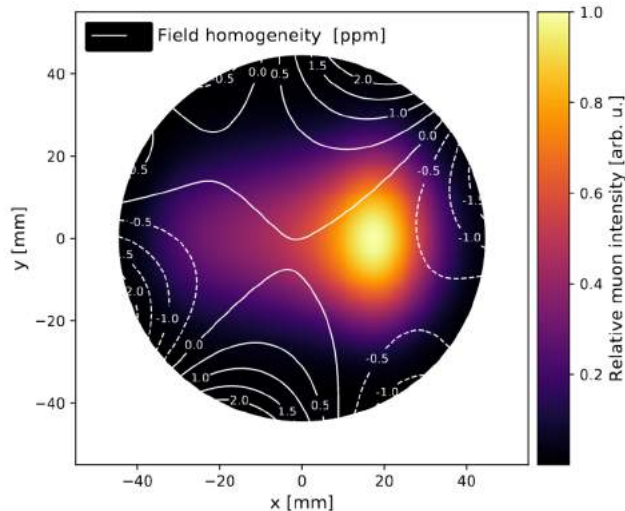
Source	Systematic Uncertainty (ppb)	Improvements undertaken
Calorimeter pileup	35	
Beam Mean Momentum & Spread	53	Increased kicker voltage: 130-161 kV
Drift of beam over measurement	75	Replaced damaged quadrupole resistors
Transient B-field (from kicker)	37	Improved magnetometer
Transient B-field (from quadrupoles)	92	More extensive measurements / damping
Total	140	

Other effects at 10-20 ppb also significantly improved by better temperature control in the experimental hall.

Bringing it all together



$$\mathcal{R}'_{\mu} = \frac{\omega_a}{\tilde{\omega}'_p(T_r)} = \frac{f_{\text{clock}} \omega_a^m (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$



Corrections

Combination with BNL



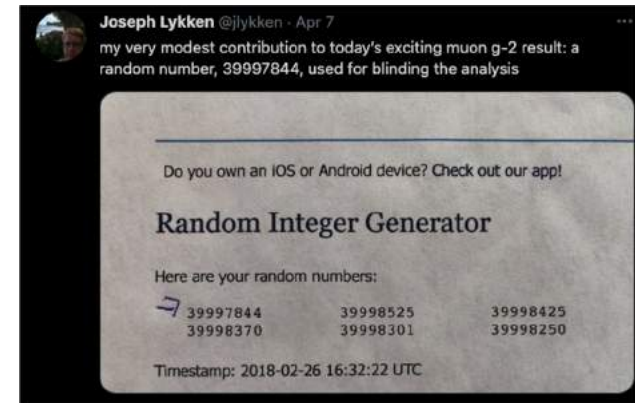
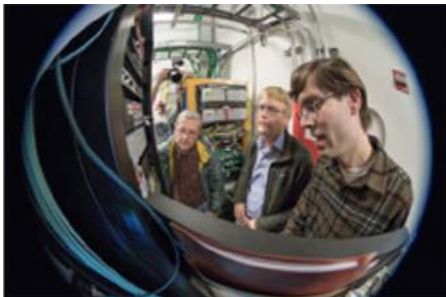
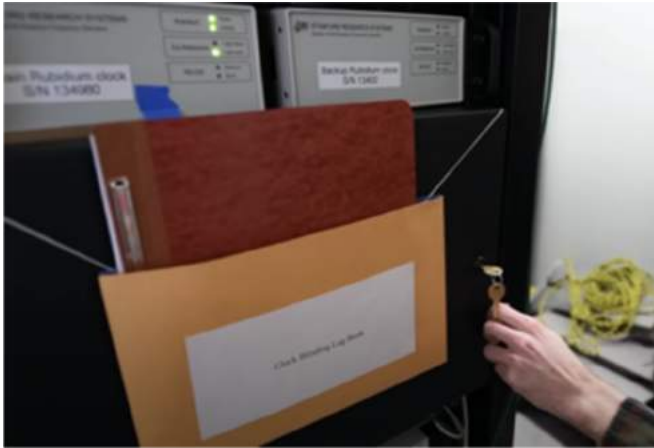
The superconducting coil, pole pieces, yoke and inflector magnet remain from BNL experiment

The underlying experimental methodology is very similar to CERN-III and BNL

- New NMR systems for magnetic field measurement
- New higher resolution calorimeter & straw trackers
- New quadrupole and kicker system
- A laser calibration system for the calorimeters (plus tracker E/p)
- Zero pion contamination and less pileup
- Full GEANT simulation and additional beam transport simulations
- More independent analyses

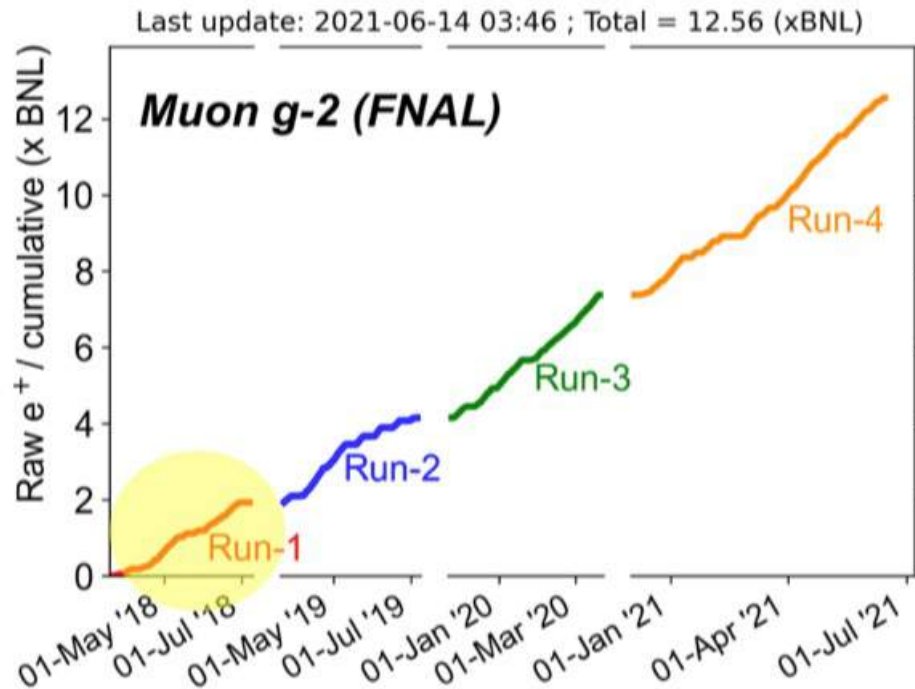
Clock Blinding

- The clock is hardware blinded to have a frequency of $(40 \pm \epsilon)$ MHz
- Only 2 people outside of the collaboration set and know the number
- Blinding offset was ± 25 ppm (approx $\times 10$ BNL-SM difference)



- Additionally each analysis is blinded in software

What next to improve/cross-check this result...



Run-4 ends in two weeks.

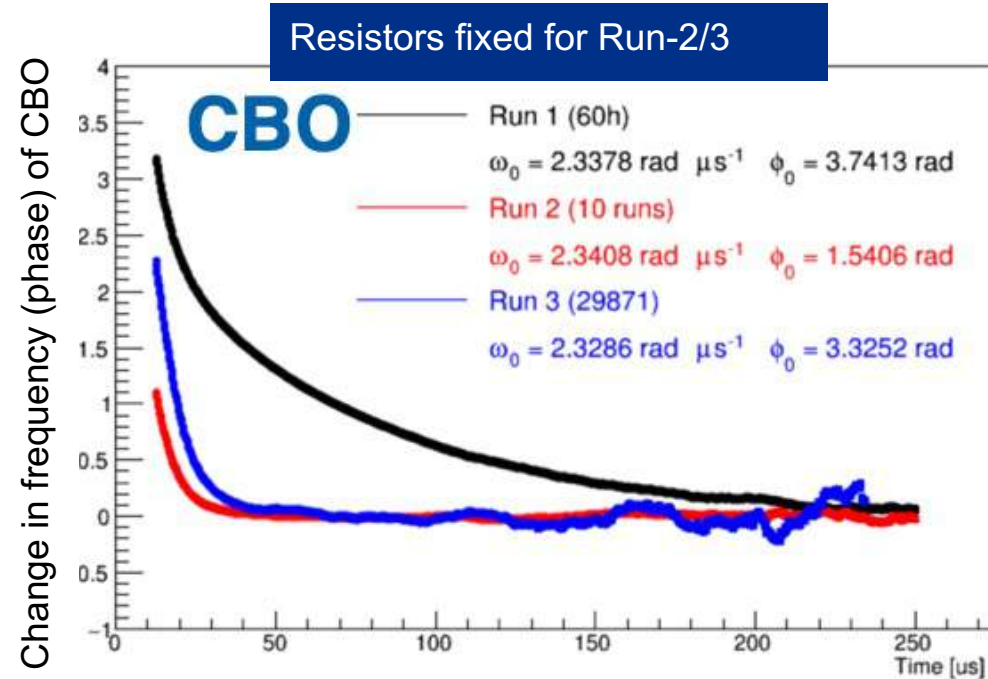
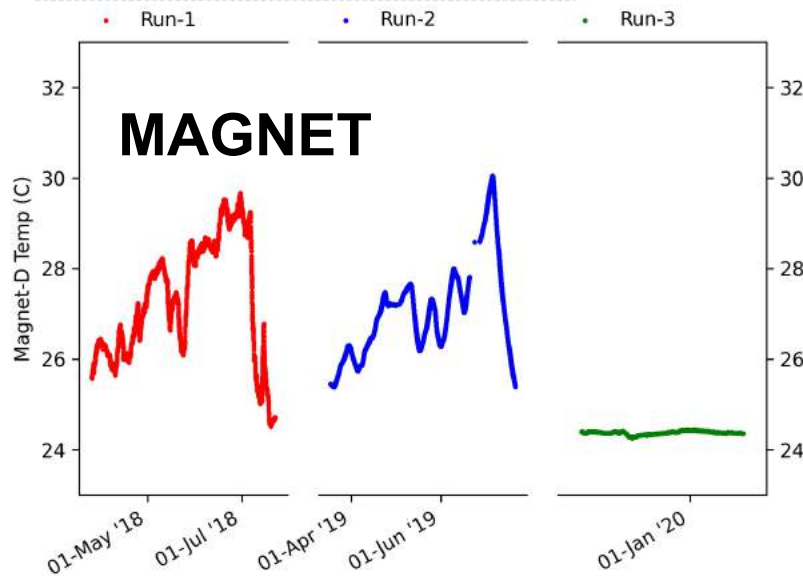
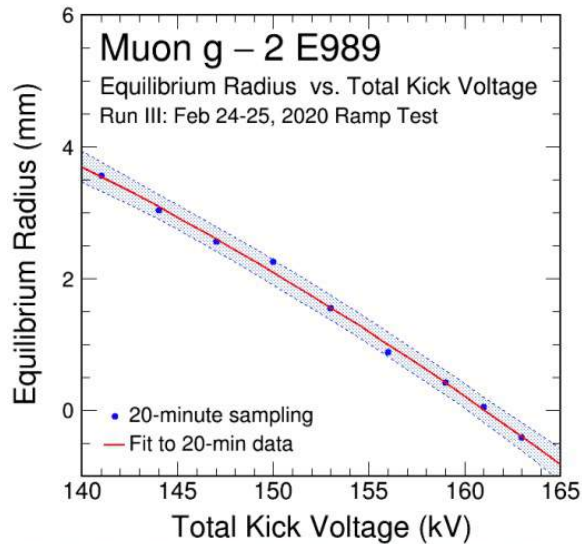
A final Run-5 will give us a total dataset ~ x20 that of the the first publication

Current publication (Run-1) based on dataset ~ 1 BNL (~ 10B muon decays)

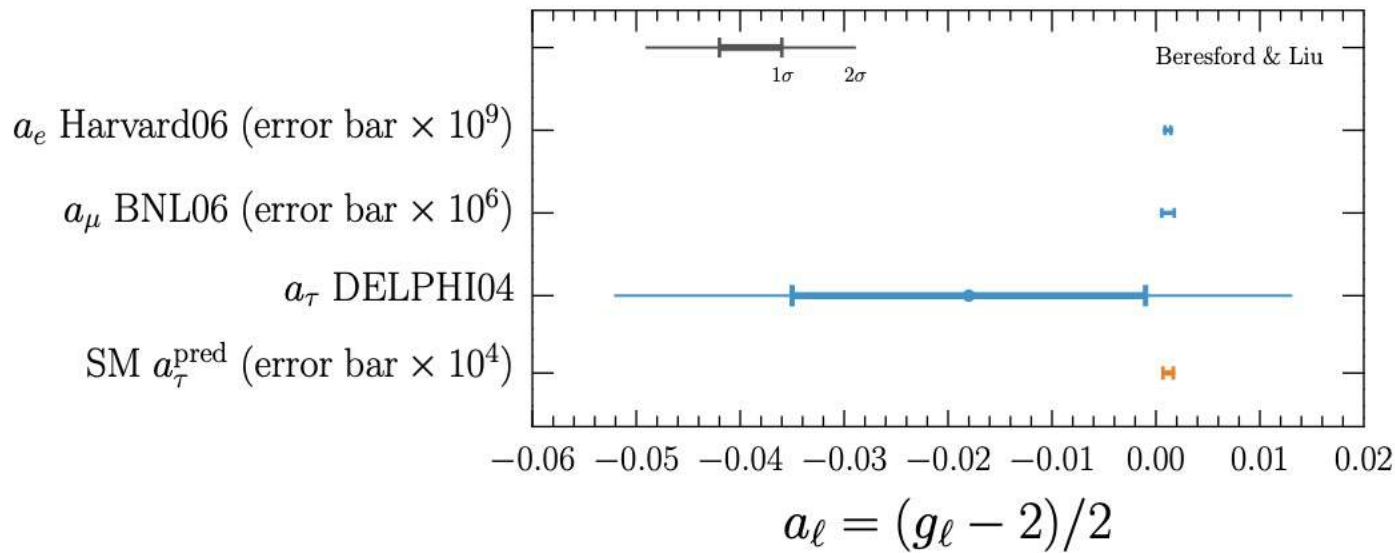
Now analysing Run-2/Run-3 : should reduce statistical uncertainty by 2 (~ 220 ppb) and expect to reach systematic uncertainty goal of 100 ppb : still statistics limited

With full dataset (upto Run-5) likely we become systematics limited

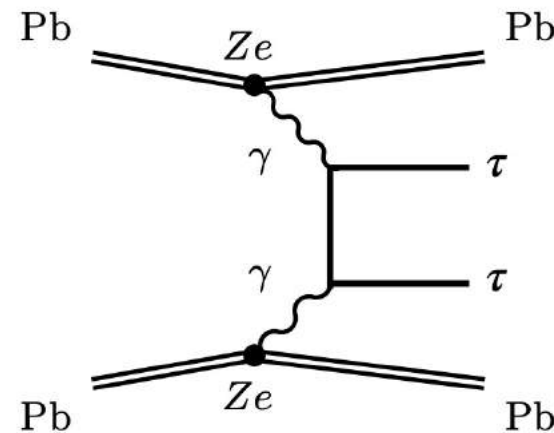
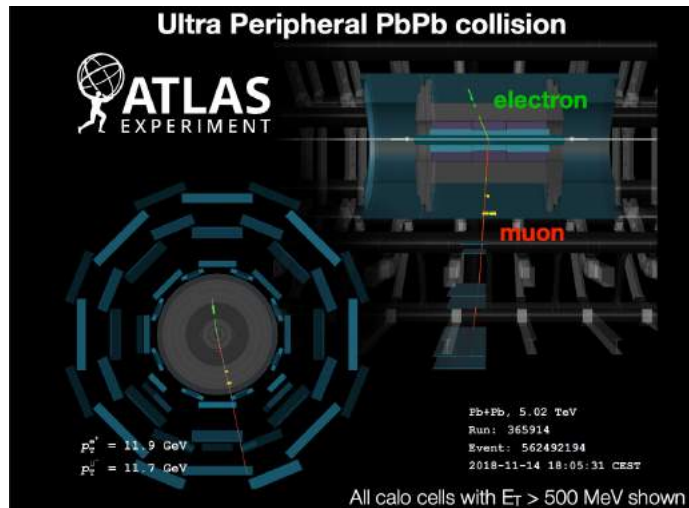
Run-2/3 Improvements

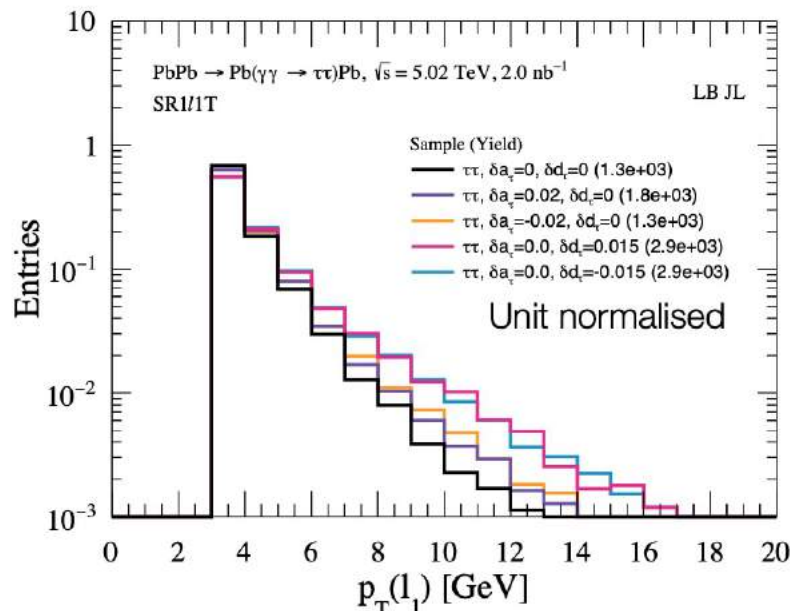


Tau g-2



$$\Delta(g - 2)_{\text{BSM}} \propto \frac{m_l^2}{M_{\text{BSM}}^2}$$





Improve on DELPHI by factor of 3.

(g-2) modifies shape of tau p_T spectrum

a_e Harvard06 (error bar $\times 10^9$)

a_μ BNL06 (error bar $\times 10^6$)

PDG a_τ DELPHI04

a_τ 2 nb^{-1} , 10% syst

a_τ 2 nb^{-1} , 5% syst

a_τ 20 nb^{-1} , 5% syst

SM a_τ^{pred} (error bar $\times 10^4$)

SMEFT a_τ^{pred} , $C_{\tau B} = -1$

