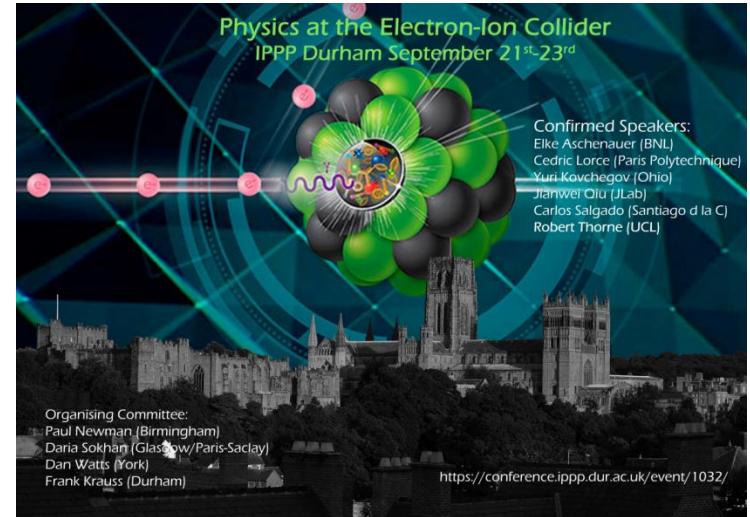
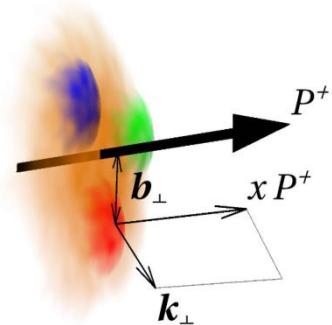




IP PARIS



TMDs and their friends

Cédric Lorcé



September 21

Outline

- Motivations
- Transverse-momentum dependent PDFs
- Generalized PDFs
- Wigner distributions

Nucleon structure

Non-relativistic picture

dominated by **constituents**

Spectroscopy



$$\mu \approx \Lambda_{\text{QCD}}$$

Mass

$$M_N c^2 \sim \sum_Q M_Q c^2 + E_{\text{binding}}$$

~ 102 %

~ - 2 %

Spin

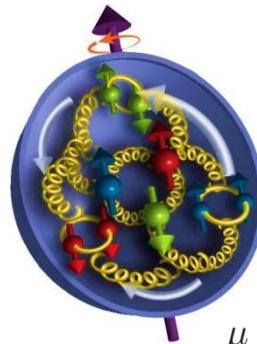
$$J_z^N \sim \sum_Q S_z^Q$$

~ 100 %

Relativistic picture

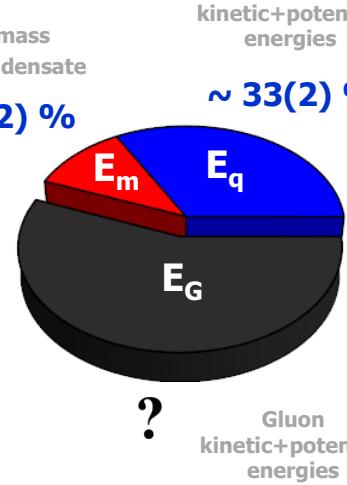
dominated by **dynamics**

High-energy scattering



$$\mu \approx 2 \text{ GeV}$$

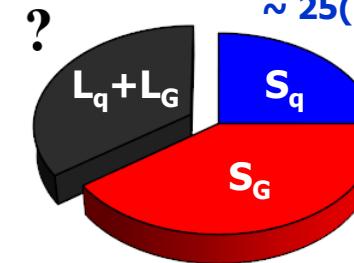
Quark mass
& QCD condensate
~ 11(2) %



Quark
kinetic+potential
energies

Orbital angular
momentum

Quark spin



~ 25(10) %

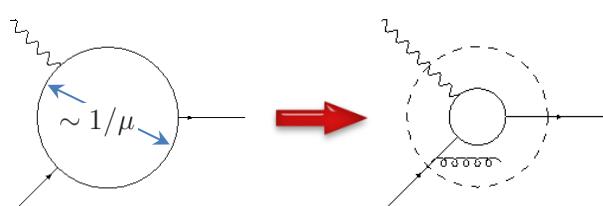
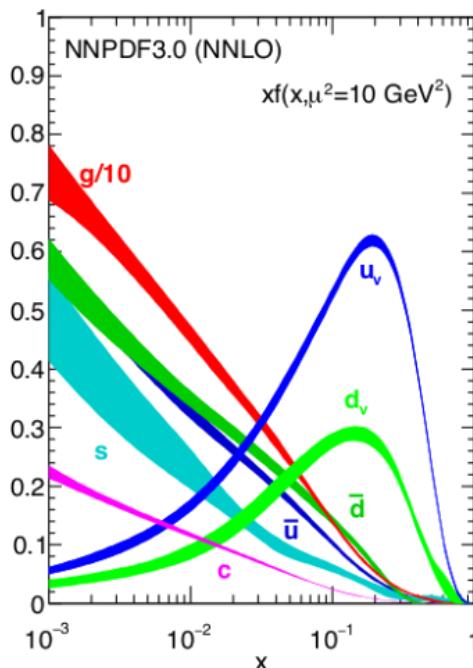
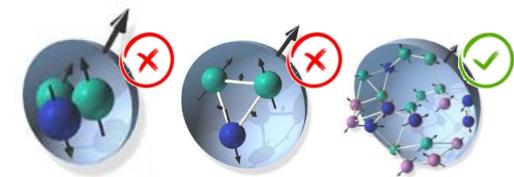
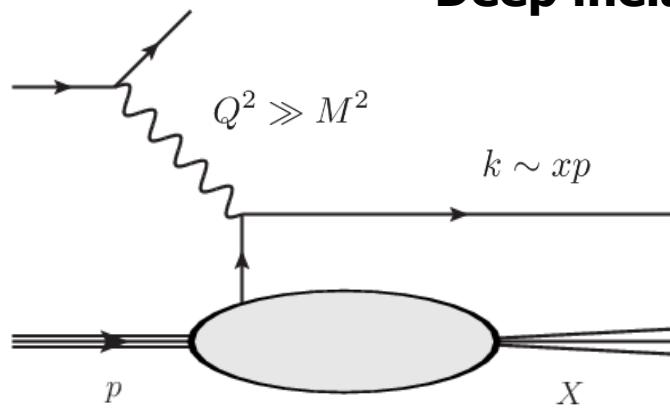
~ 40(?) %

Gluon spin

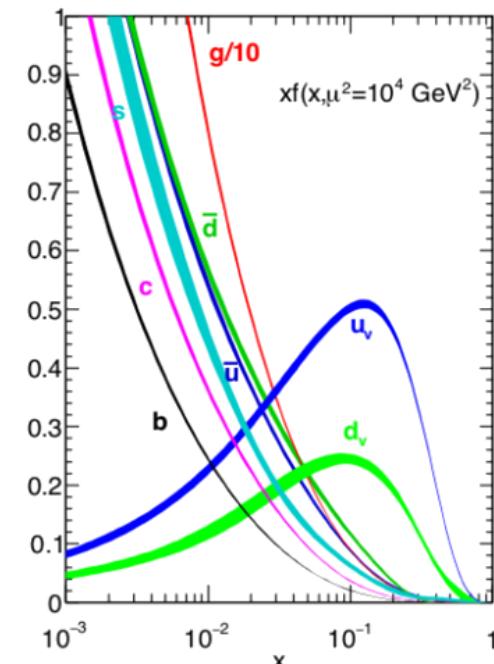
→ **EIC !**

Parton distribution functions (PDFs)

Deep inelastic scattering (DIS)

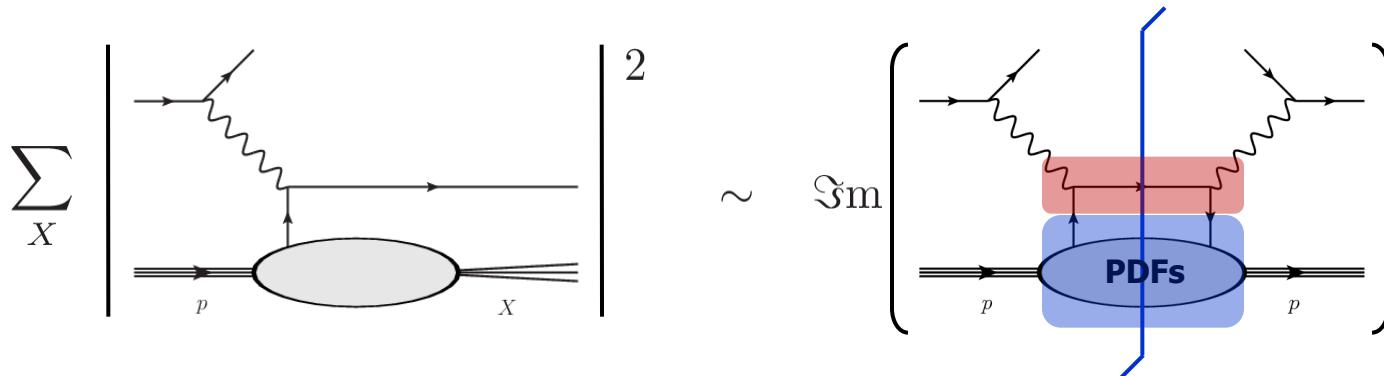


Perturbative QCD evolution
(DGLAP)



Parton distribution functions (PDFs)

Optical theorem



$$d\sigma^{lp \rightarrow lX} \sim \sum_i \int_x^1 \frac{dy}{y} C_i\left(\frac{x}{y}, \frac{Q^2}{\mu_F^2}; \alpha_S(\mu_R)\right) f_i(y, \mu_F^2)$$

Perturbative
Process-dependent
Non-perturbative
Process-independent

$$i = q_f, \bar{q}_f, g$$

$$\mu_F \approx \mu_R$$

Parton density (in $A^+ = 0$ gauge)

$$a^\pm = \frac{1}{\sqrt{2}}(a^0 \pm a^3)$$

$$f(x) = \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle P | \bar{\psi}(-\frac{z}{2}) \gamma^+ \psi(\frac{z}{2}) | P \rangle \Big|_{z^+ = z_\perp = 0}$$

$$k^+ = xP^+$$

$$\sim \frac{1}{2x(2\pi)} \sum_\lambda \int \frac{d^2 k_\perp}{(2\pi)^2} \frac{\langle P | b_{k,\lambda}^\dagger b_{k,\lambda} | P \rangle}{\langle P | P \rangle}$$

$$\langle p' | p \rangle = 2p^+(2\pi)^3 \delta(p'^+ - p^+) \delta^{(2)}(\vec{p}'_\perp - \vec{p}_\perp)$$

Transverse-momentum dependent PDFs

Going beyond the collinear partonic picture

Quark Wigner operator (contour gauge)

$$j^\mu(r, k) = \int \frac{d^4 z}{(2\pi)^4} e^{ik \cdot z} \bar{\psi}(r - \frac{z}{2}) \gamma^\mu \psi(r + \frac{z}{2})$$

Non-local in both longitudinal and transverse directions !

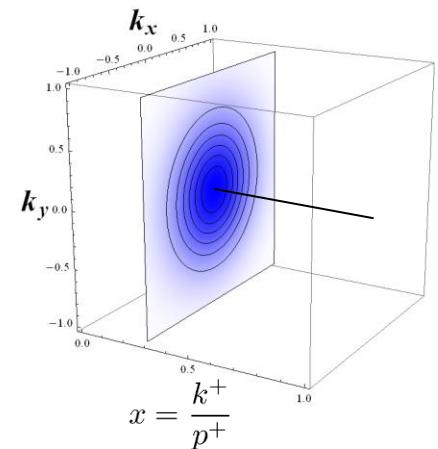
Charge current operator

$$j^\mu(r) = \int d^4 k j^\mu(r, k)$$

Unintegrated distributions

$$\text{PDF}(x) \sim \frac{1}{2} \int dk^- d^2 k_\perp \langle p | j^+(0, k) | p \rangle$$

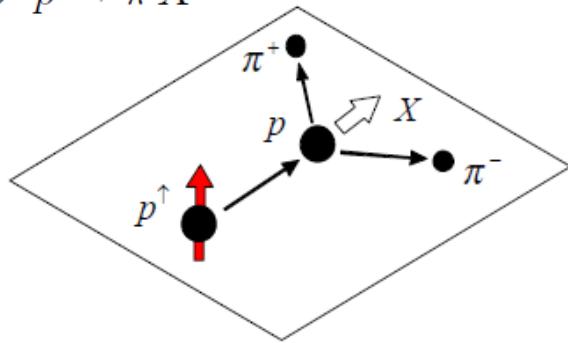
$$\text{TMD}(x, \vec{k}_\perp) \sim \frac{1}{2} \int dk^- \langle p | j^+(0, k) | p \rangle$$



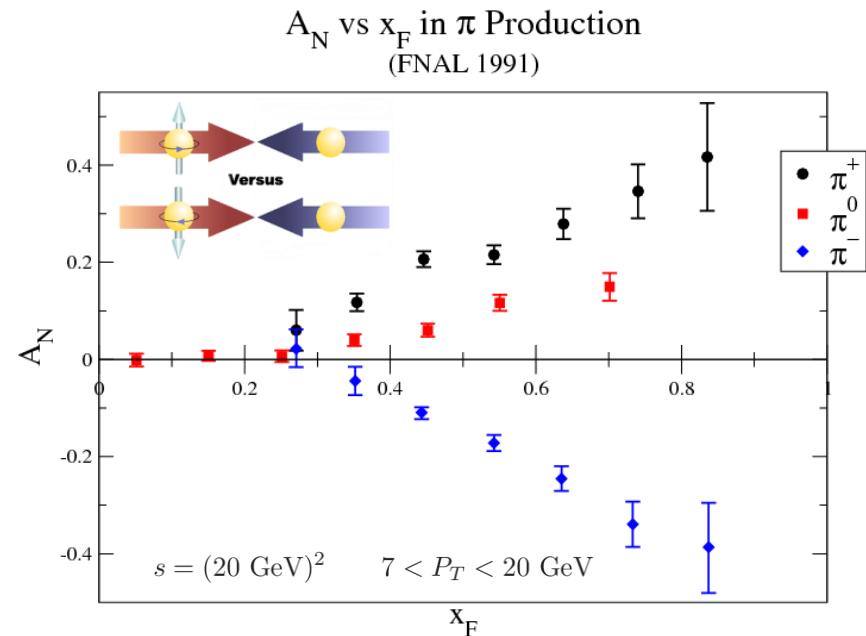
Transverse-momentum dependent PDFs

Single-spin asymmetries (SSA)

E.g. $p^\uparrow p \rightarrow \pi X$



$$A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} \propto \vec{S}_N \cdot (\vec{p}_N \times \vec{P}_T)$$



Collinear twist-2



$$A_N \propto \frac{m_q}{E_q} \alpha$$

Price for helicity flip

Too small !

[Kane, Pumplin, Repko (1978)]

Collinear twist-3

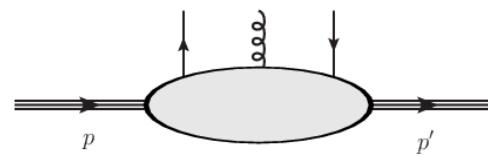


$$Q \sim P_T \gg \Lambda_{\text{QCD}}$$

Intrinsic k_T

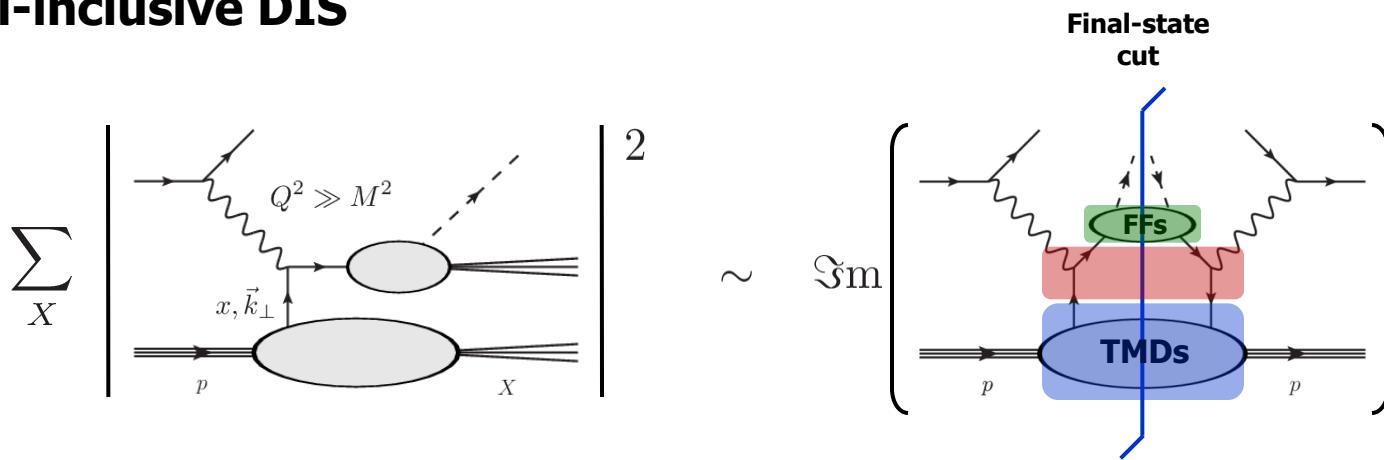


$$Q \gg P_T \sim \Lambda_{\text{QCD}}$$

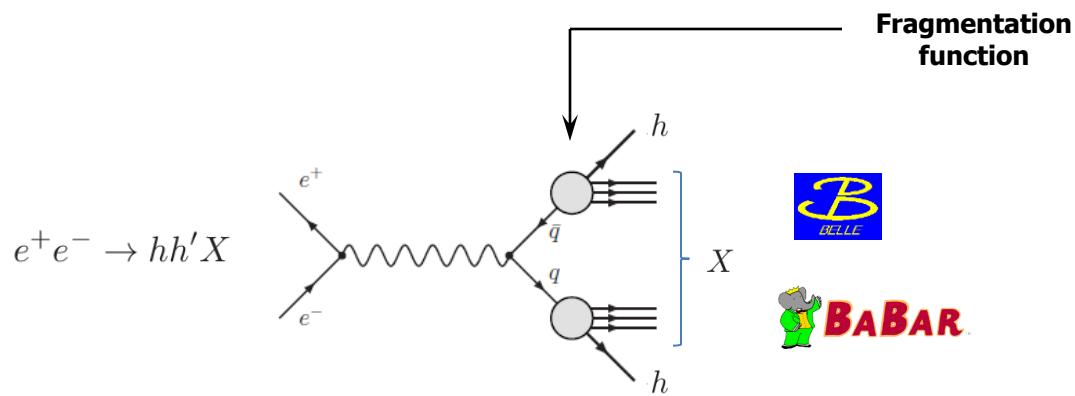


Transverse-momentum dependent PDFs

Semi-inclusive DIS

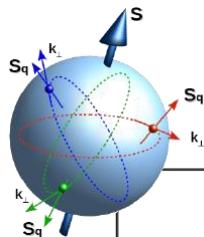


$$d\sigma \sim \sum \left[\text{TMD}(x, \vec{k}_\perp) \otimes d\hat{\sigma}_{hard} \otimes \text{FF}(z, \vec{p}_\perp) + \mathcal{O}\left(\frac{P_T}{Q}\right) \right]$$



Transverse-momentum dependent PDFs

Rich spin structure (leading twist)



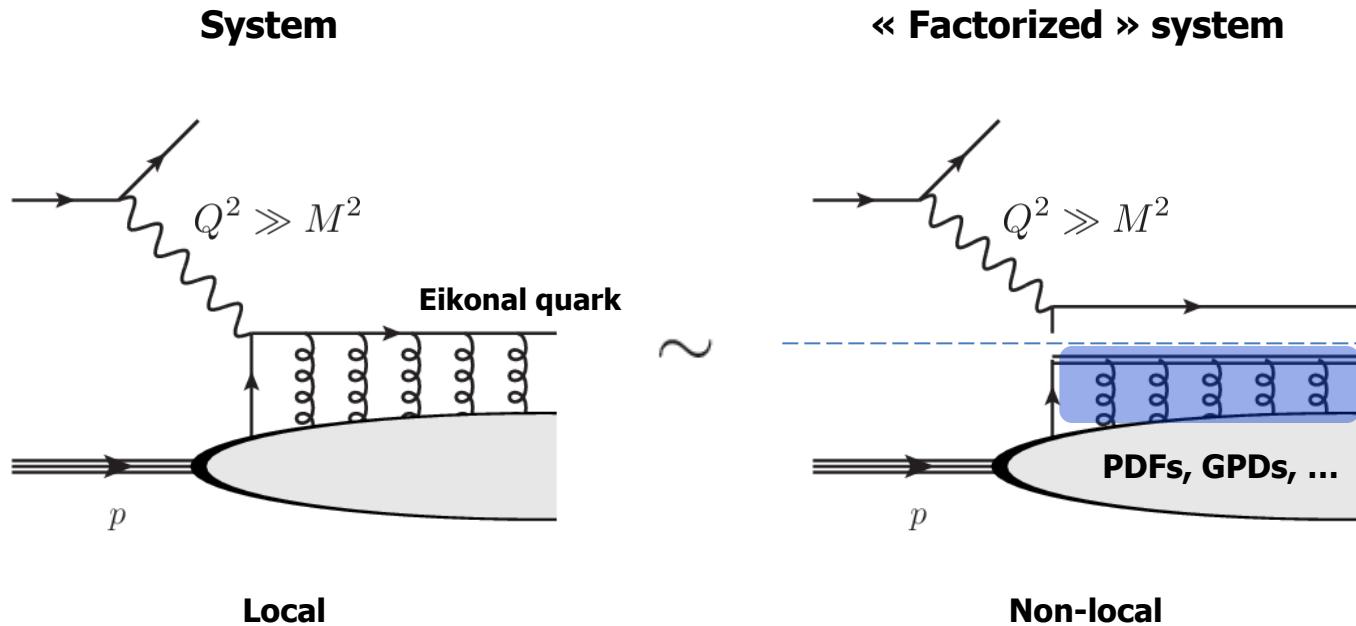
Quark polarization

	U	T_x	T_y	L
U	f_1	$\frac{k_y}{M} h_1^\perp$	$-\frac{k_x}{M} h_1^\perp$	
T_x	$\frac{k_y}{M} f_{1T}^\perp$	$h_1 + \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$\frac{k_x}{M} g_{1T}$
T_y	$-\frac{k_x}{M} f_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$h_1 - \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_y}{M} g_{1T}$
L		$\frac{k_x}{M} h_{1L}^\perp$	$\frac{k_y}{M} h_{1L}^\perp$	g_{1L}

Survives integration over transverse momentum

Transverse-momentum dependent PDFs

Gauge symmetry



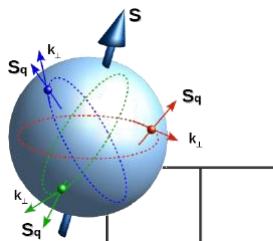
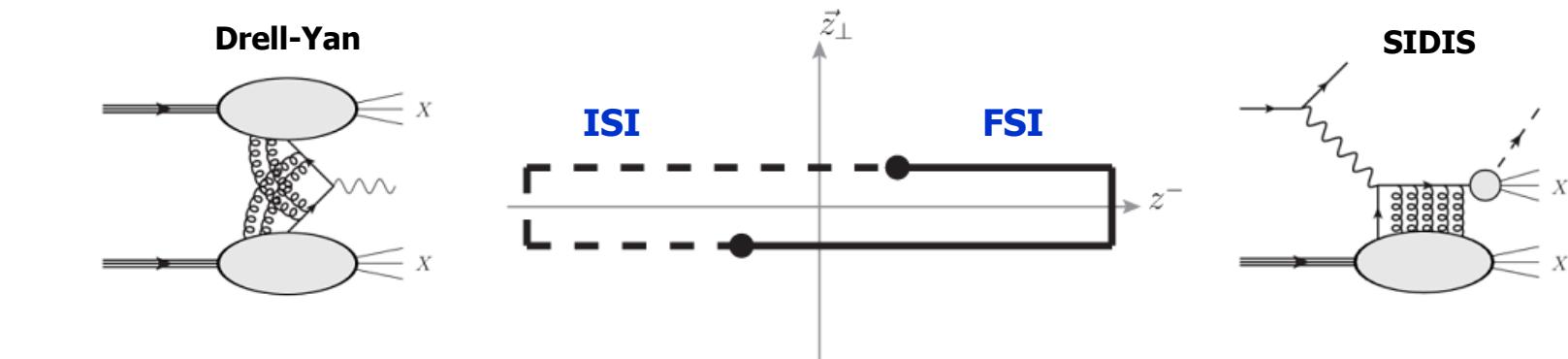
$$j^\mu(r, k) = \int \frac{d^4 z}{(2\pi)^4} e^{ik \cdot z} \bar{\psi}(r - \frac{z}{2}) \gamma^\mu \mathcal{W}(r - \frac{z}{2}, r + \frac{z}{2}) \psi(r + \frac{z}{2})$$

Wilson line

$$\mathcal{W}(b, a) = \mathcal{P} \left[e^{ig \int_a^b A^\mu dz_\mu} \right]$$

Transverse-momentum dependent PDFs

Process dependence



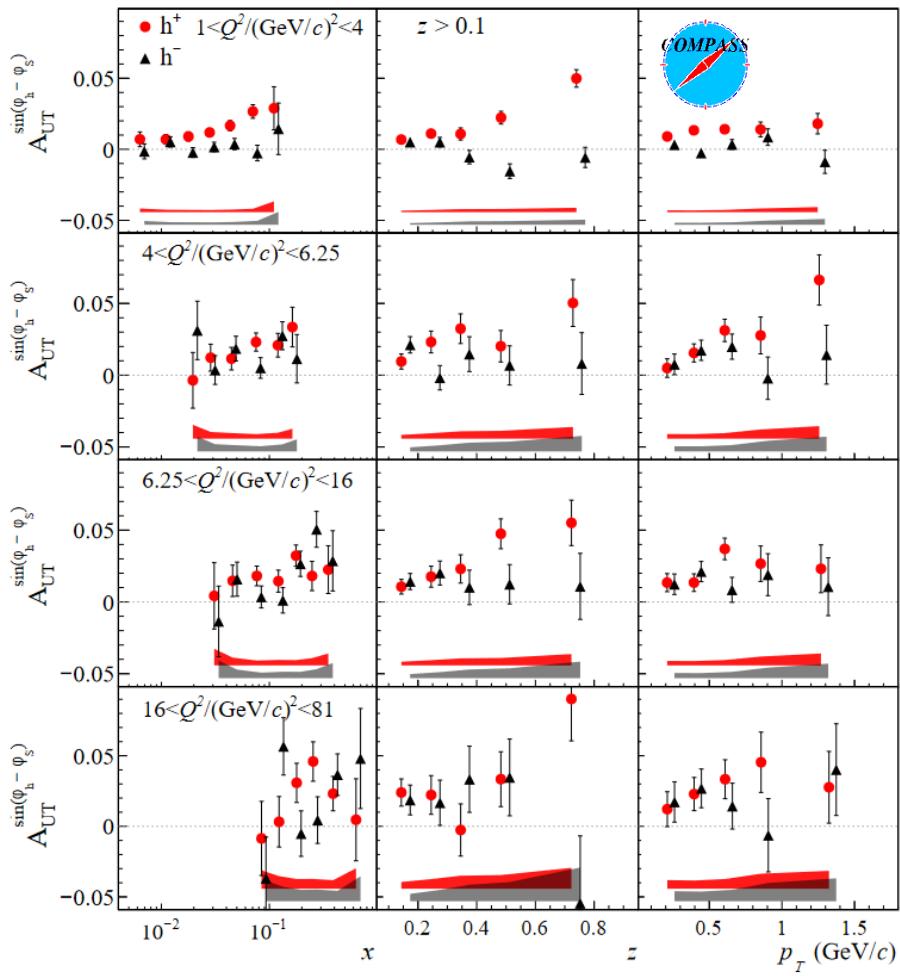
Quark polarization

	U	T_x	T_y	L
U	f_1	$\frac{k_y}{M} h_{1L}^\perp$	$-\frac{k_x}{M} h_{1L}^\perp$	
T_x	$\frac{k_y}{M} f_{1T}^\perp$	$h_1 + \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$\frac{k_x}{M} g_{1T}$
T_y	$-\frac{k_x}{M} f_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$h_1 - \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_y}{M} g_{1T}$
L		$\frac{k_x}{M} h_{1L}^\perp$	$\frac{k_y}{M} h_{1L}^\perp$	g_{1L}

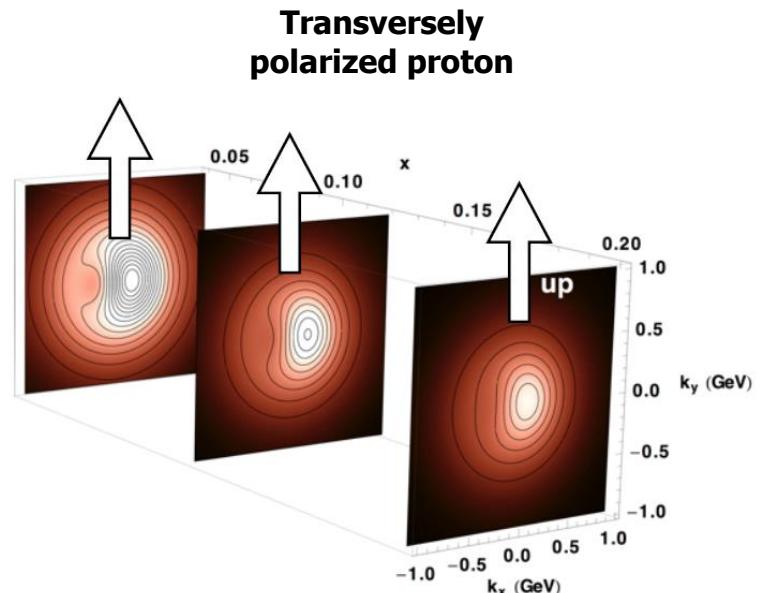
Changes sign
under ISI \leftrightarrow FSI
(naive T-odd)

Transverse-momentum dependent PDFs

Clear SIDIS experimental signal



[Adolph *et al.*, PLB770 (2017) 138]



Courtesy of A. Bacchetta

Parton distribution zoo

Nonlocal quark operator

$$\tilde{O}(z) = \bar{\psi}(-\frac{z}{2}) \Gamma \mathcal{W}(-\frac{z}{2}, \frac{z}{2}) \psi(\frac{z}{2})$$



$$O \sim \tilde{O}(0)$$

$$O(x) \sim \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \tilde{O}(z)|_{z^+=z_\perp=0}$$

$$O(x, \vec{k}_\perp) \sim \int \frac{dz^- d^2 z_\perp}{(2\pi)^3} e^{ixP^+z^- - i\vec{k}_\perp \cdot \vec{z}_\perp} \tilde{O}(z)|_{z^+=0}$$

$$\langle p | O(x, \vec{k}_\perp) | p \rangle$$

TMDs

$$x, \vec{k}_\perp$$

$$\langle p | O(x) | p \rangle$$

PDFs

$$x$$

$$\langle p | O | p \rangle$$

Charges

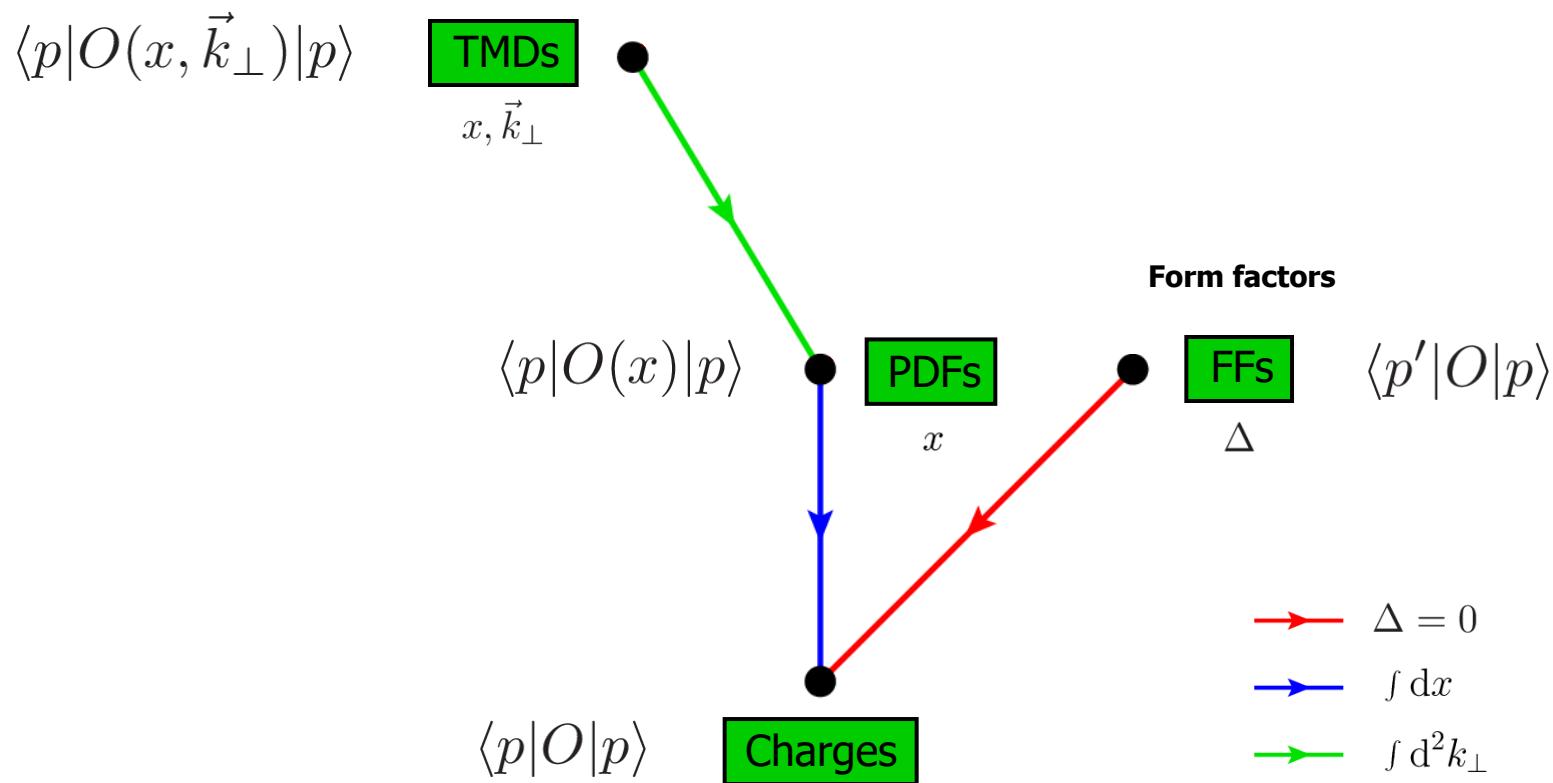
- $\int dx$
- $\int d^2 k_\perp$

Parton distribution zoo

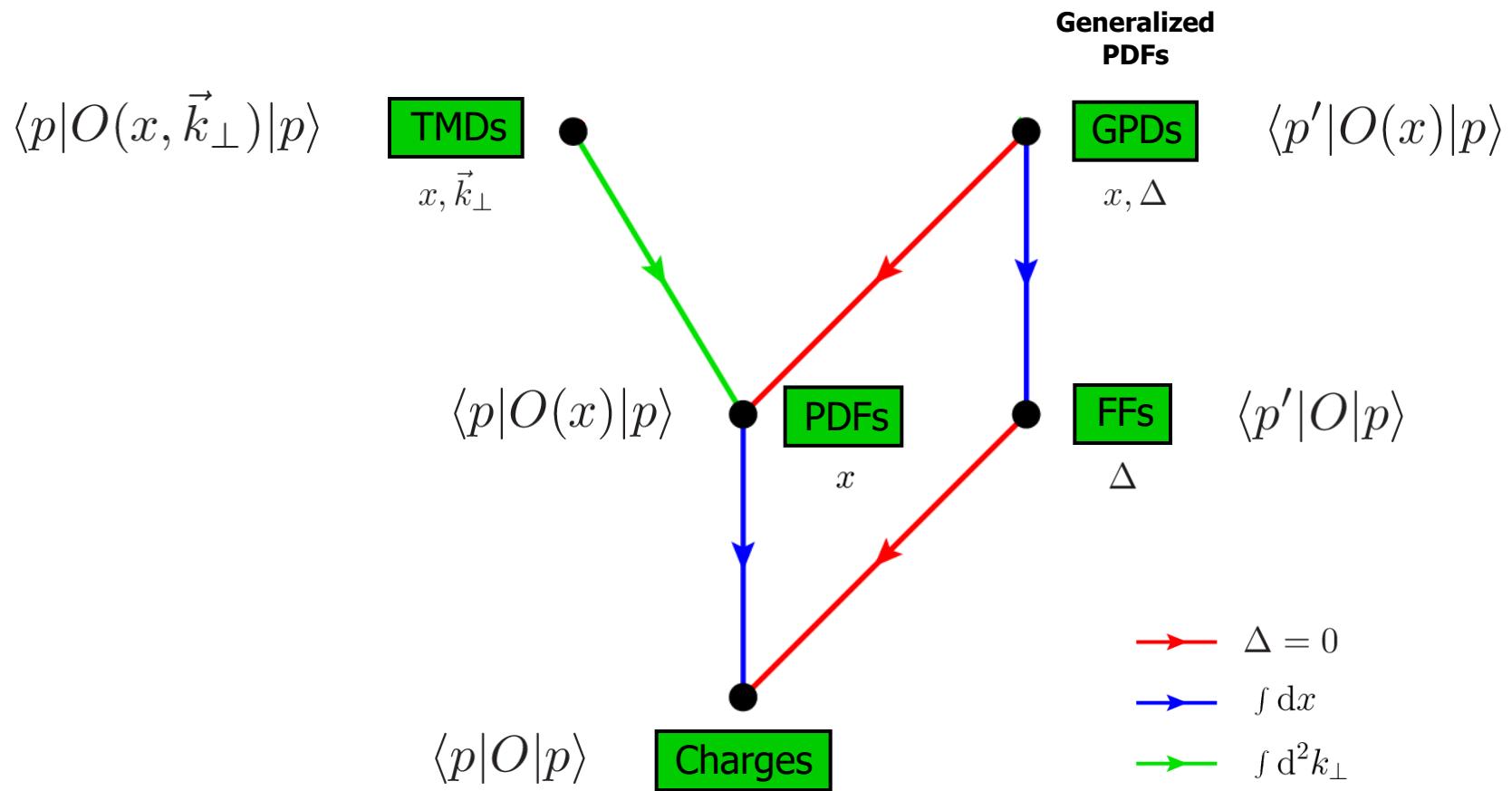
Off-forward amplitudes

$$\langle p | \cdots | p \rangle \longrightarrow \langle p' | \cdots | p \rangle$$

$$P = \frac{1}{2}(p' + p), \quad \Delta = p' - p$$



Parton distribution zoo



Parton distribution zoo



Generalized
TMDs

GTMDs

$\langle p' | O(x, \vec{k}_\perp) | p \rangle$

x, \vec{k}_\perp, Δ

$\langle p | O(x, \vec{k}_\perp) | p \rangle$

TMDs

x, \vec{k}_\perp

GPDs

$\langle p' | O(x) | p \rangle$

x, Δ

$\langle p | O(x) | p \rangle$

PDFs

x

$\langle p' | O | p \rangle$

FFs

Δ

$\langle p | O | p \rangle$

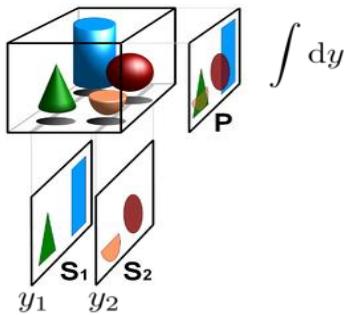
Charges

→ $\Delta = 0$

→ $\int dx$

→ $\int d^2 k_\perp$

Parton distribution zoo



« Mother »
distributions

GTMDs

x, \vec{k}_\perp, Δ

TMDs

x, \vec{k}_\perp

GPDs

x, Δ

PDFs

x

FFs

Δ

Charges

→ $\Delta = 0$

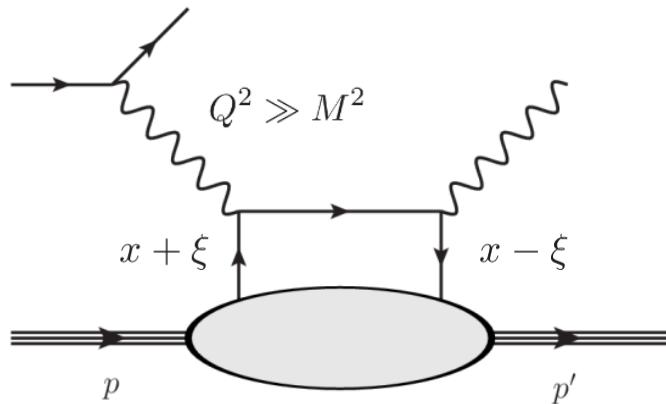
→ $\int dx$

→ $\int d^2 k_\perp$

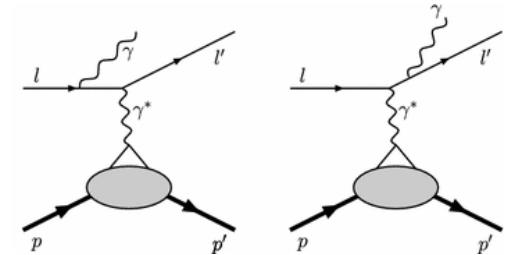
[Meissner, Metz, Schlegel (2009)]
[C.L., Pasquini, Vanderhaeghen (2011)]

Generalized PDFs

Deeply virtual Compton scattering (DVCS)



interferes with



Correlator (in $A^+ = 0$ gauge)

$$\begin{aligned} \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle p' | \bar{\psi}(-\frac{z}{2}) \gamma^+ \psi(\frac{z}{2}) | p \rangle \Big|_{z^+=z_\perp=0} \\ = \frac{1}{2P^+} \bar{u}(p') \left[\gamma^+ H(x, \xi, t) + \frac{i\sigma^{+\mu} \Delta_\mu}{2M} E(x, \xi, t) \right] u(p) \end{aligned}$$

Momentum transfer variables

$$\xi = -\frac{\Delta^+}{2P^+} = \frac{p^+ - p'^+}{p^+ + p'^+}, \quad t = \Delta^2 = (p' - p)^2$$

Generalized PDFs

Link with other non-perturbative functions

$$H(x, 0, 0) = f(x)$$

PDF

$$\int dx H(x, \xi, t) = F_1(t)$$

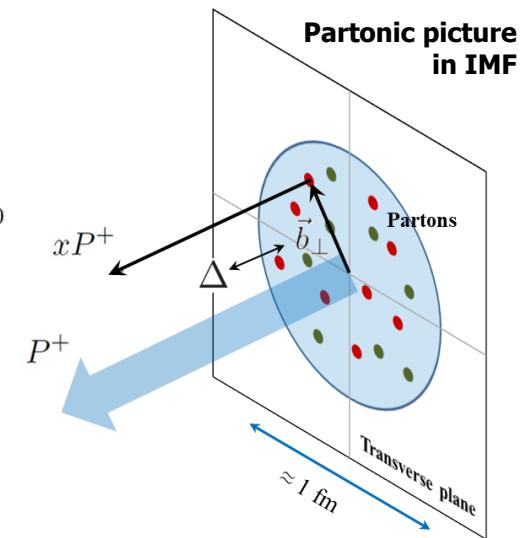
$$\int dx E(x, \xi, t) = F_2(t)$$

Electromagnetic
form factors

2+1D imaging (in $A^+ = 0$ gauge)

$$\rho(x, \vec{b}_\perp) = P^+ \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i \vec{\Delta}_\perp \cdot \vec{b}_\perp} \frac{\langle p' | j^+(xP^+) | p \rangle}{2P^+} \Big|_{\Delta^+ = 0}$$

$$j^\mu(k^+) = \int \frac{dz^-}{2\pi} e^{ik^+ z^-} \bar{\psi}(-\frac{z}{2}) \gamma^\mu \psi(\frac{z}{2}) \Big|_{z^+ = z_\perp = 0}$$



Generalized PDFs

Link with gravitational form factors

$$a^{\{\mu} b^{\nu\}} = \frac{1}{2}(a^\mu b^\nu + a^\nu b^\mu)$$
$$a^{[\mu} b^{\nu]} = \frac{1}{2}(a^\mu b^\nu - a^\nu b^\mu)$$

$$\begin{aligned} \langle p' | T^{\mu\nu}(0) | p \rangle &= \bar{u}(p') \left[P^{\{\mu} \gamma^{\nu\}} A(t) + \frac{P^{\{\mu} i\sigma^{\nu\}} \lambda \Delta_{\lambda}}{2M} B(t) \right. \\ &\quad \left. + \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{M} C(t) + M g^{\mu\nu} \bar{C}(t) + \frac{P^{[\mu} i\sigma^{\nu]} \lambda \Delta_{\lambda}}{2M} D(t) \right] u(p) \end{aligned}$$

$$T_q^{++}(0) = (\bar{\psi} \gamma^+ \frac{i}{2} \overset{\leftrightarrow}{D}^+ \psi)(0)$$

$$= 2(P^+)^2 \int dx x \frac{1}{2} \int \frac{dz^-}{2\pi} e^{ixP^+ z^-} \bar{\psi}(-\frac{z}{2}) \gamma^+ \psi(\frac{z}{2}) \Big|_{z^+ = z_\perp = 0}$$



$$\begin{aligned} \int dx x H(x, \xi, t) &= A(t) + 4\xi^2 C(t) \\ \int dx x E(x, \xi, t) &= B(t) - 4\xi^2 C(t) \end{aligned}$$

Generalized PDFs

Poincaré constraints

$$\sum_{a=q,g} A_a(0) = 1$$

$$\sum_{a=q,g} B_a(0) = 0$$

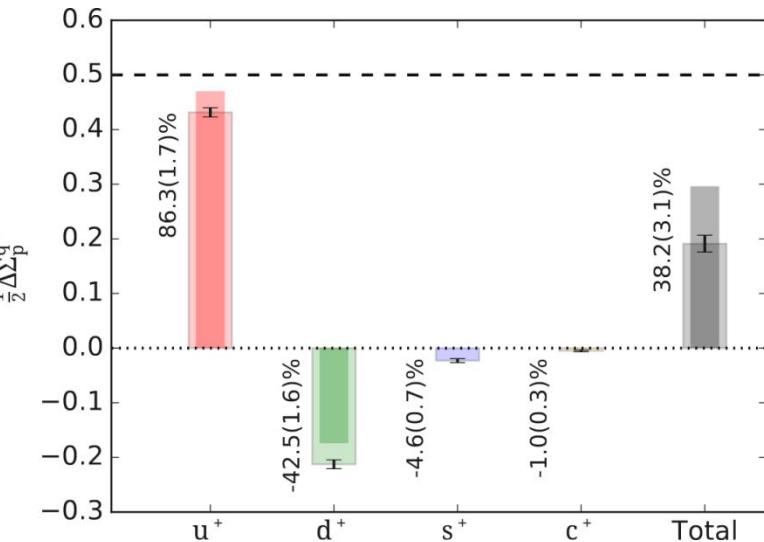
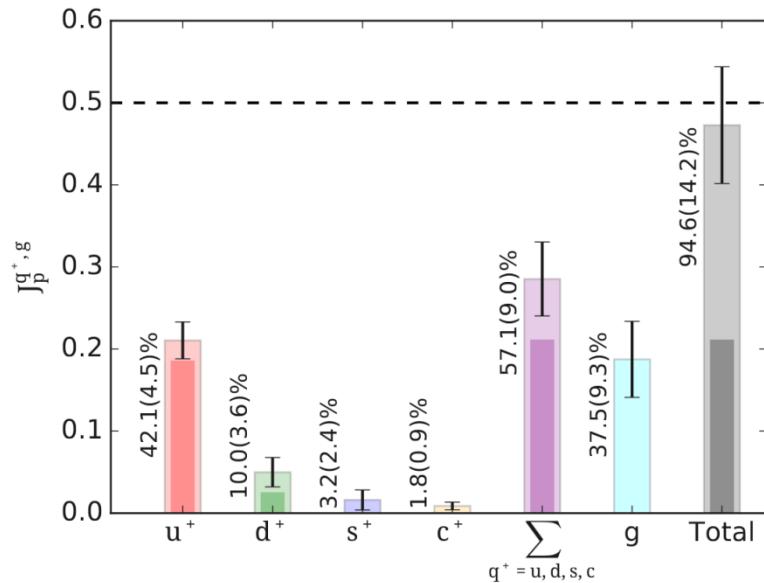
$$\sum_{a=q,g} \bar{C}_a(t) = 0$$

$$D_q(t) = -G_A(t)$$

Angular momentum

$$\langle J_q^z \rangle = \frac{A_q(0) + B_q(0)}{2}$$

$$\langle S_q^z \rangle = \frac{G_A(0)}{2} = \frac{\Delta \Sigma}{2}$$

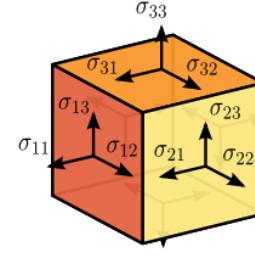


Energy-momentum tensor

Instead of charge current, we can look at four-momentum current

$$T^{\mu\nu} = \begin{bmatrix} T^{00} & & & \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \end{bmatrix}$$

Energy density Momentum density
Energy flux Momentum flux



Shear stress

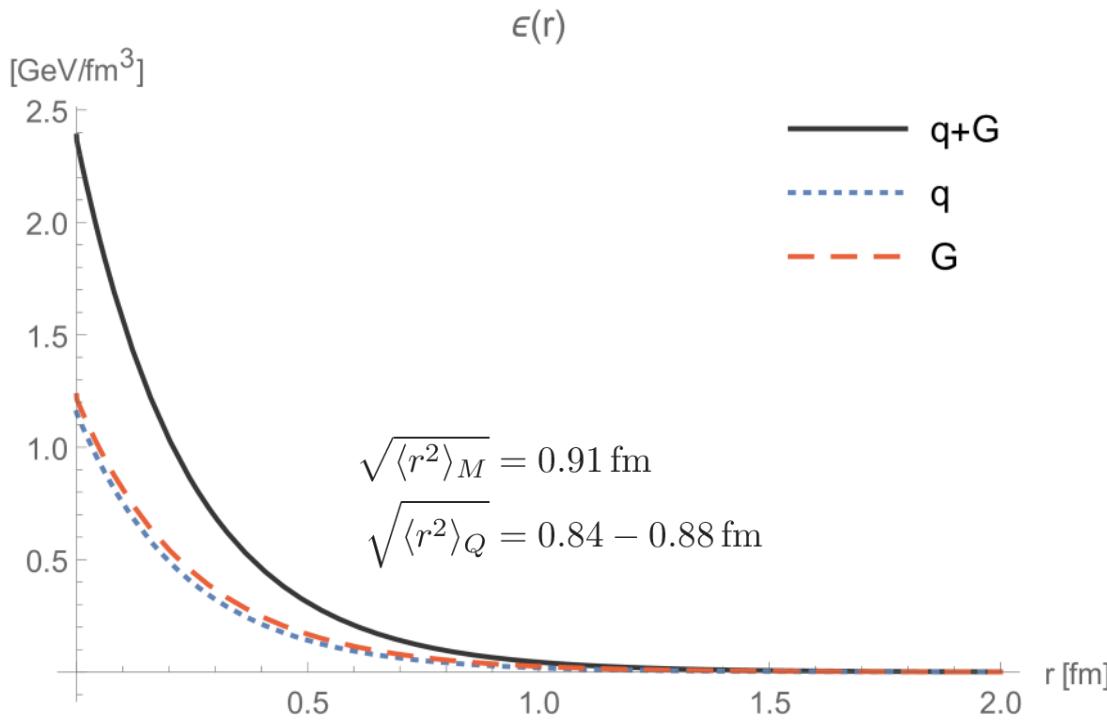
Normal stress (pressure)

Four-momentum operator

$$P_a^\mu = \int d^3r T_a^{0\mu}(r) \quad a = q, g$$

Energy distribution (3D Breit frame)

$$\langle T^{00} \rangle_{\vec{0},\vec{0}}(\vec{r}) = \int \frac{d^3\Delta}{(2\pi)^3} e^{-i\vec{\Delta}\cdot\vec{r}} \left(\frac{\vec{\Delta}}{2} |T^{00}(0)| - \frac{\vec{\Delta}}{2} \right)$$

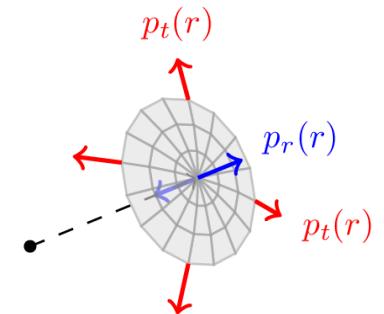


Multipole model for the gravitational form factors

$$F(t) = \frac{F(0)}{(1 + t/\Lambda^2)^n}$$

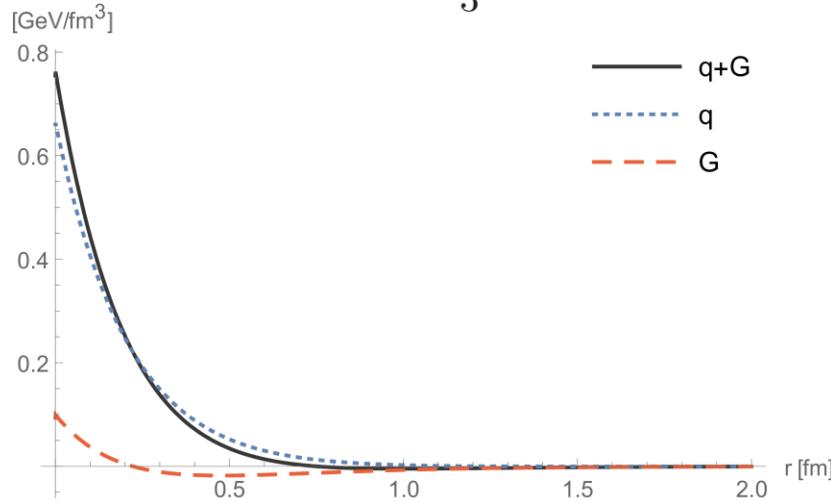
Pressure distributions (3D Breit frame)

$$\langle T^{ij} \rangle_{\vec{0},\vec{0}}(\vec{r}) = \int \frac{d^3\Delta}{(2\pi)^3} e^{-i\vec{\Delta}\cdot\vec{r}} \left\langle \frac{\vec{\Delta}}{2} |T^{ij}(0)| - \frac{\vec{\Delta}}{2} \right\rangle$$



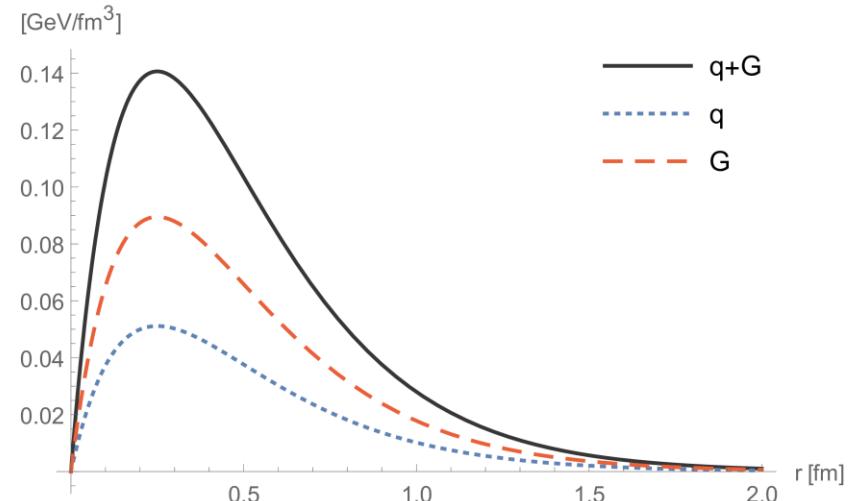
Isotropic pressure

$$p(r) = \frac{p_r(r) + 2p_t(r)}{3}$$



Pressure anisotropy

$$s(r) = p_r(r) - p_t(r)$$

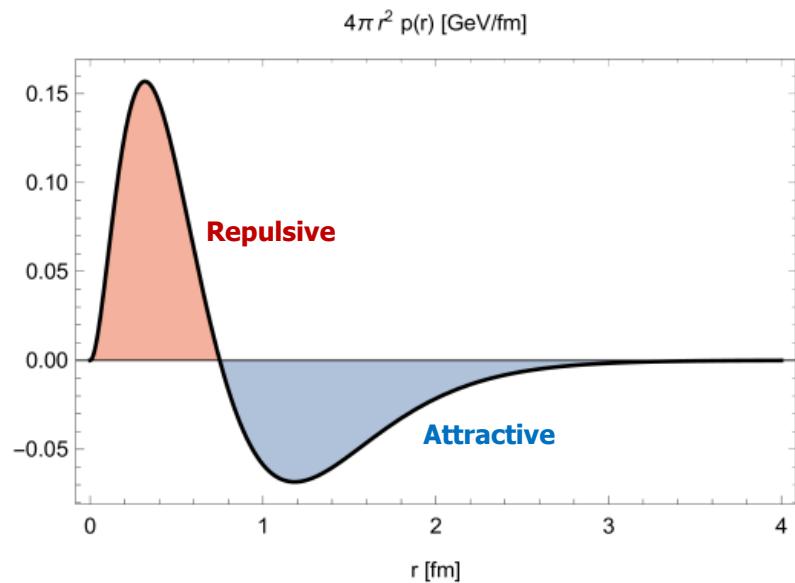


Mechanical equilibrium

$$\nabla^i \langle T^{ij} \rangle_{\vec{0}, \vec{0}}(\vec{r}) = 0 \quad \rightarrow \quad \frac{dp_r(r)}{dr} = -\frac{2s(r)}{r}$$

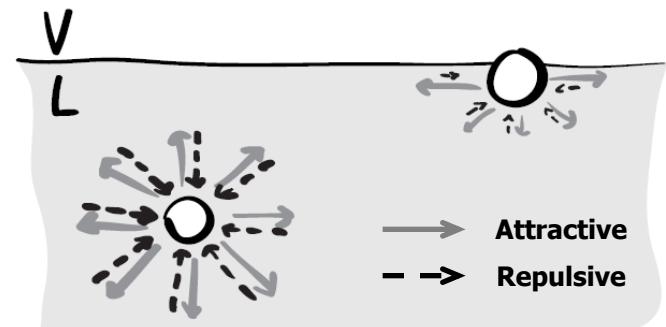
von Laue relation

$$\int_0^\infty dr r^2 p(r) = 0$$



Surface tension

$$\gamma = \int dr s(r)$$

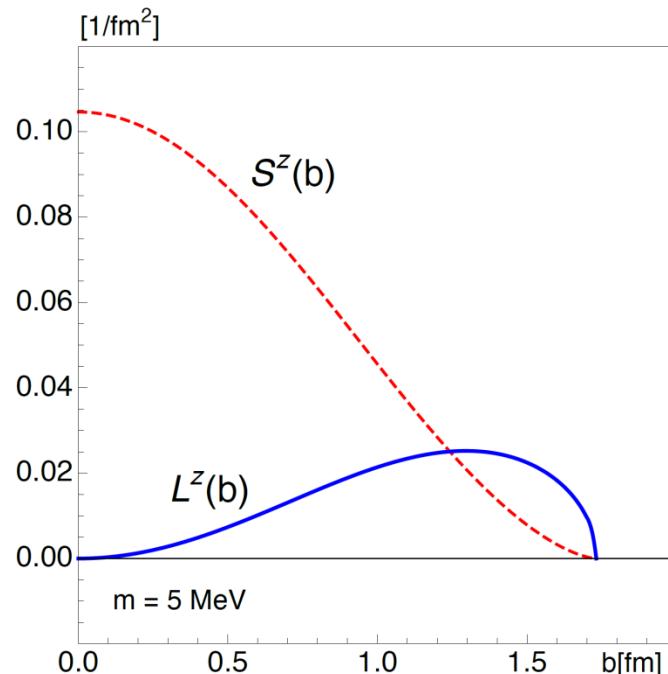


Angular momentum distributions

Orbital *vs* intrinsic

$$L^i(\vec{r}) = \epsilon^{ijk} r^j \langle T^{0k} \rangle_{\vec{0}, \vec{0}}(\vec{r})$$

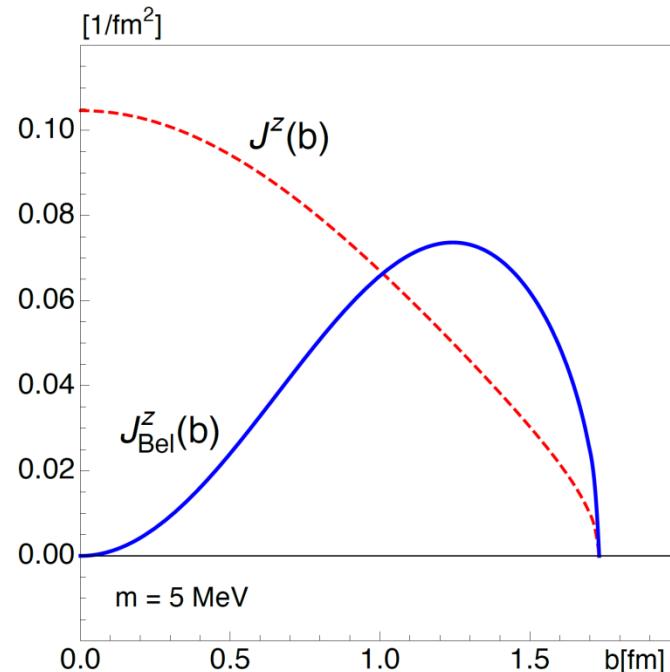
$$S^i(\vec{r}) = \langle \bar{\psi} \gamma^i \gamma_5 \psi \rangle_{\vec{0}, \vec{0}}(\vec{r})$$



Kinetic *vs* Belinfante

$$J^i(\vec{r}) = L^i(\vec{r}) + S^i(\vec{r})$$

$$J_{\text{Bel}}^i(\vec{r}) = \epsilon^{ijk} r^j \langle \frac{1}{2}(T^{0k} + T^{k0}) \rangle_{\vec{0}, \vec{0}}(\vec{r})$$



Large- N_c bag model

Generalized TMDs

$$\text{PDF}(x) \sim \frac{1}{2} \int dk^- d^2 k_\perp \langle p | j^+(0, k) | p \rangle$$

$$\text{TMD}(x, \vec{k}_\perp) \sim \frac{1}{2} \int dk^- \langle p | j^+(0, k) | p \rangle$$

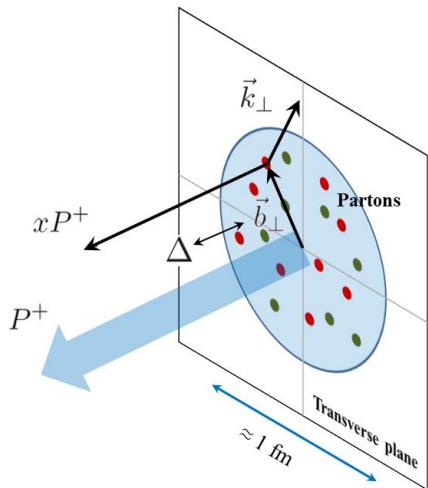
$$\text{GPD}(x, \Delta) \sim \frac{1}{2} \int dk^- d^2 k_\perp \langle p' | j^+(0, k) | p \rangle$$

$$\text{GTMD}(x, \vec{k}_\perp, \Delta) \sim \frac{1}{2} \int dk^- \langle p' | j^+(0, k) | p \rangle$$

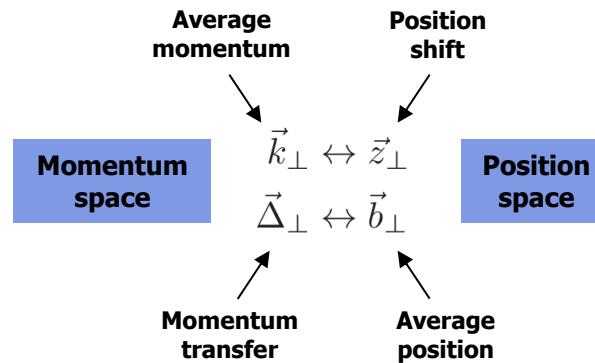
Quark Wigner operator

$$j^\mu(r, k) = \int \frac{d^4 z}{(2\pi)^4} e^{ik \cdot z} \bar{\psi}(r - \frac{z}{2}) \gamma^\mu \mathcal{W}(r - \frac{z}{2}, r + \frac{z}{2}) \psi(r + \frac{z}{2})$$

Wigner (or phase-space) distribution



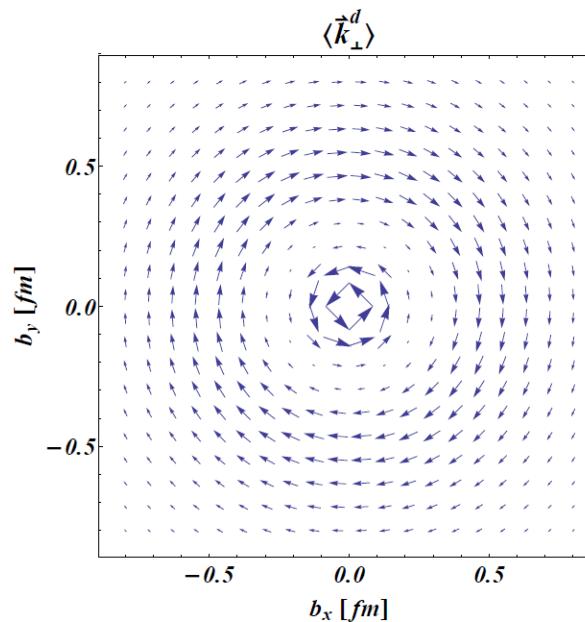
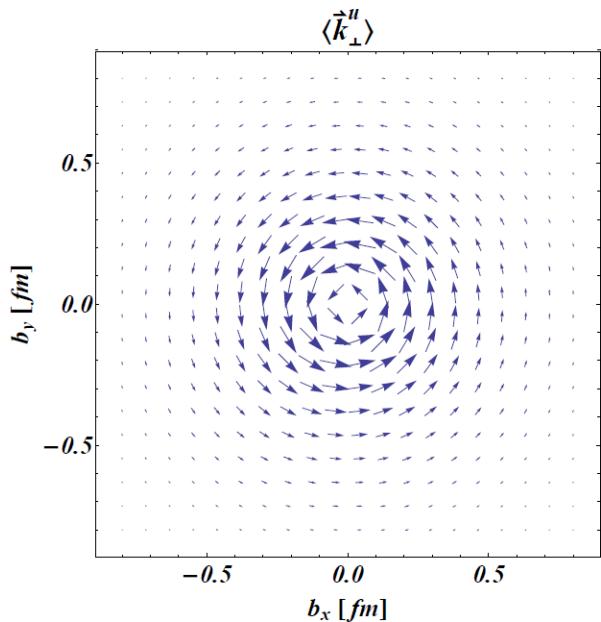
$$\text{WD}(x, \vec{k}_\perp, \vec{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i \vec{\Delta}_\perp \cdot \vec{b}_\perp} \text{GTMD}(x, \vec{k}_\perp, \Delta) \Big|_{\Delta+ = 0}$$



Generalized TMDs

Average transverse quark momentum

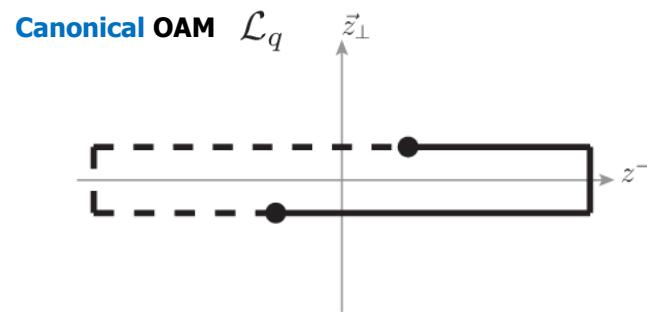
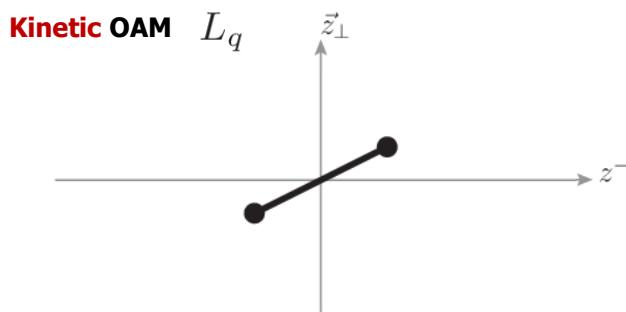
$$\langle \vec{k}_\perp \rangle(\vec{b}_\perp) = \int dx d^2 k_\perp \vec{k}_\perp \rho_{LU}(x, \vec{k}_\perp, \vec{b}_\perp)$$



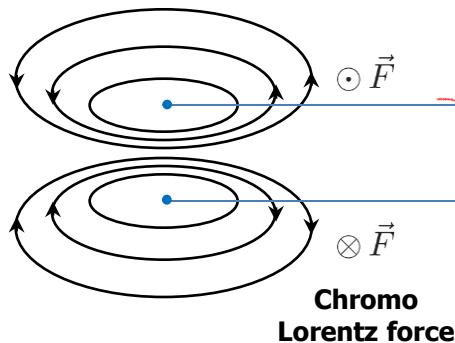
Generalized TMDs

Orbital angular momentum

$$\ell_z = \int d^2 b_\perp \vec{b}_\perp \times \langle \vec{k}_\perp \rangle(\vec{b}_\perp)$$

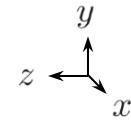
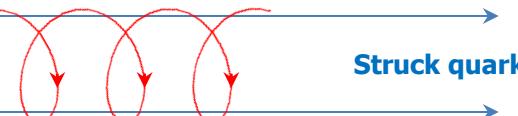


« Inside » the nucleon

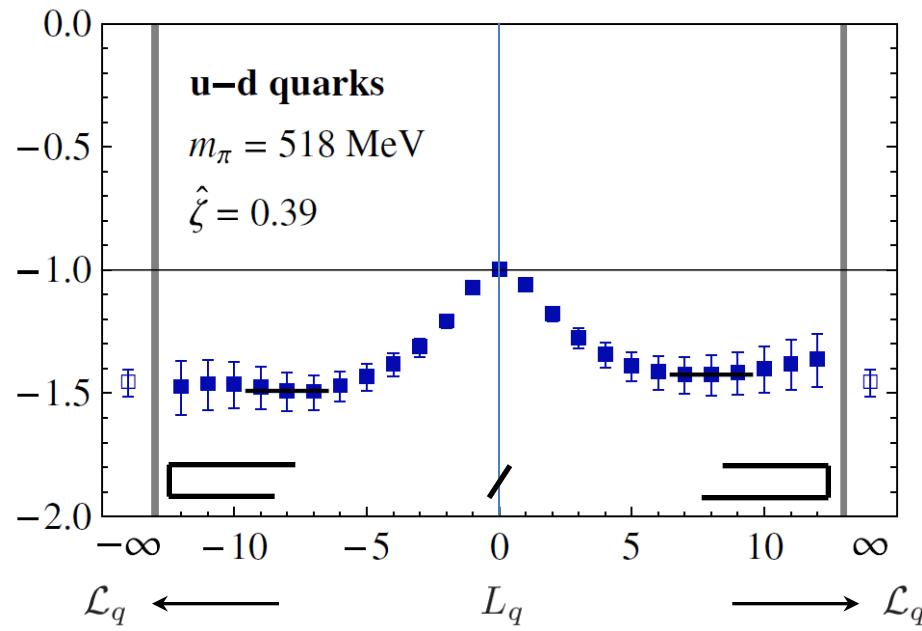


Chromo
Lorentz force

« Outside » the nucleon



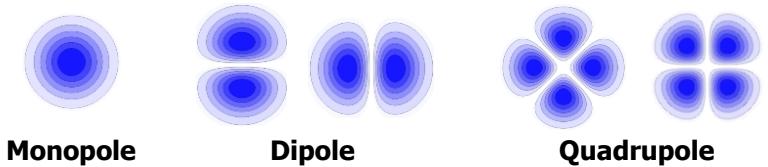
Generalized TMDs



[Engelhardt, PRD95 (2017) 094505]

Correlations

Twist-2



Nucleon polarization

Quark polarization

GTMDs

	U	T_x	T_y	L
U	F_{11}	$\frac{i}{M} (k_y H_{11} + \Delta_y H_{12})$	$-\frac{i}{M} (k_x H_{11} + \Delta_x H_{12})$	$\frac{i(\vec{\Delta}_\perp \times \vec{k}_\perp)_z}{M^2} G_{11}$
T_x	$\frac{i}{M} (k_y F_{12} + \Delta_y \bar{F}_{13})$	$\frac{1}{M} (k_x \bar{G}_{12} + \Delta_x \bar{G}_{13})$
T_y	$-\frac{i}{M} (k_x F_{12} + \Delta_x \bar{F}_{13})$	$\frac{1}{M} (k_y \bar{G}_{12} + \Delta_y \bar{G}_{13})$
L	$-\frac{i(\vec{\Delta}_\perp \times \vec{k}_\perp)_z}{M^2} F_{14}$	$\frac{1}{M} (k_x H_{17} + \Delta_x H_{18})$	$\frac{1}{M} (k_y H_{17} + \Delta_y H_{18})$	G_{14}

$$\Delta = 0$$

$$\int d^2 k_\perp$$

TMDs

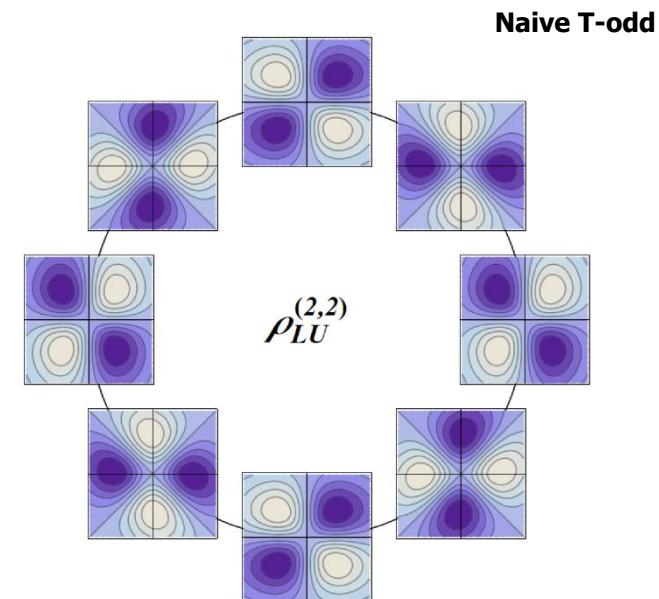
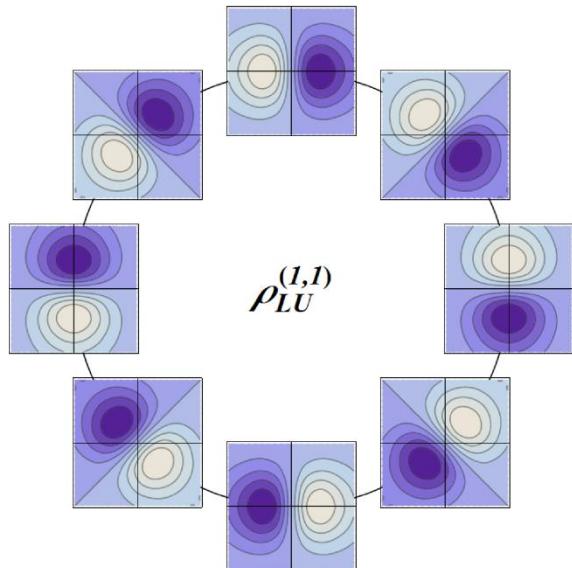
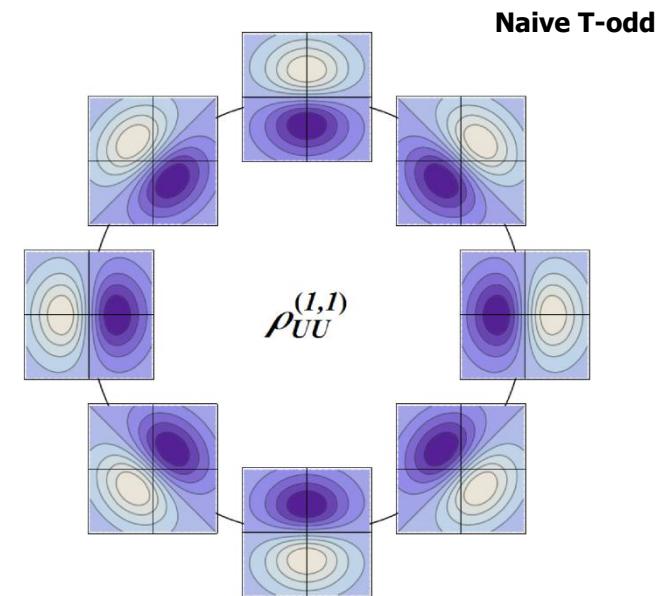
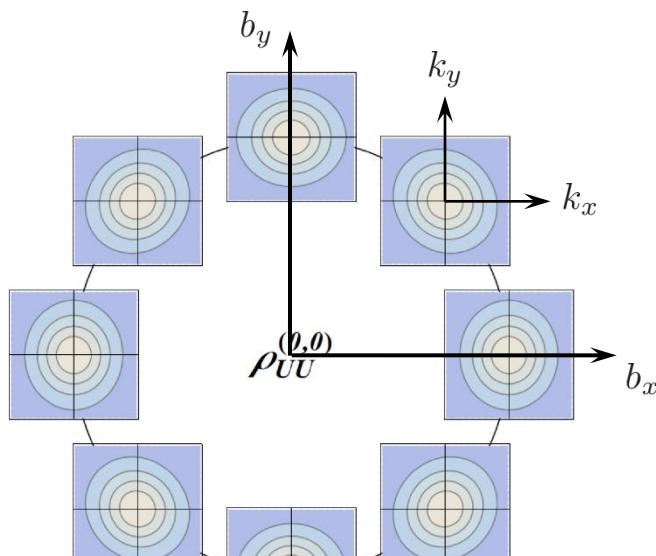
$$\vec{k}_\perp \leftrightarrow i\vec{\Delta}_\perp$$

GPDs

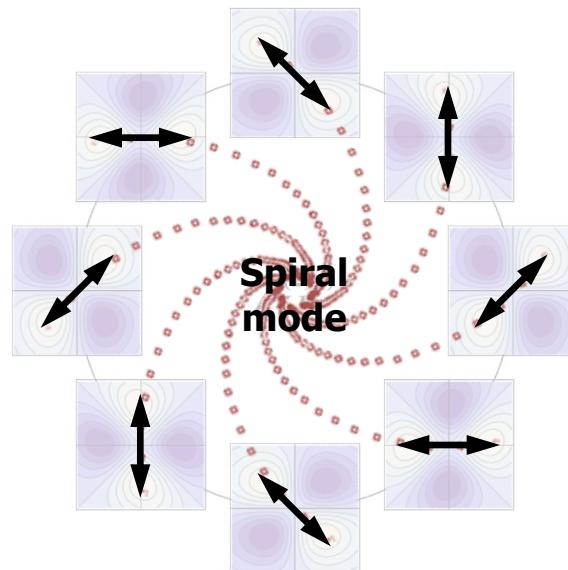
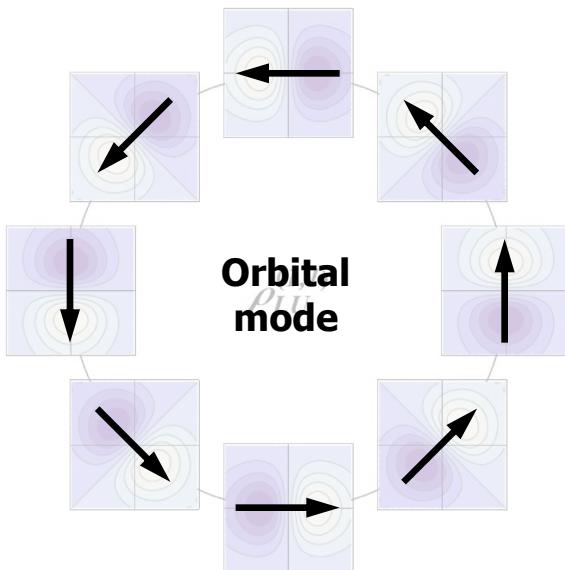
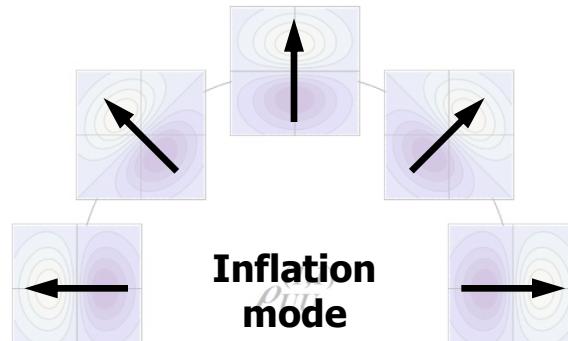
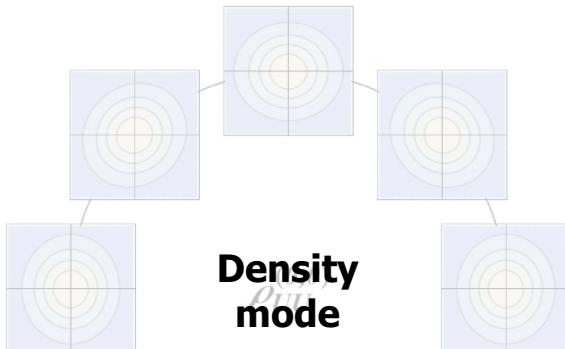
	U	T_x	T_y	L
U	f_1	$\frac{k_y}{M} h_1^\perp$	$-\frac{k_x}{M} h_1^\perp$	
T_x	$\frac{k_y}{M} f_{1T}^\perp$	$h_1 + \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$\frac{k_x}{M} g_{1T}$
T_y	$-\frac{k_x}{M} f_{1T}^\perp$	$\frac{k_x k_y}{M^2} h_{1T}^\perp$	$h_1 - \frac{k_x^2 - k_y^2}{2M^2} h_{1T}^\perp$	$\frac{k_y}{M} g_{1T}$
L		$\frac{k_x}{M} h_{1L}^\perp$	$\frac{k_y}{M} h_{1L}^\perp$	g_{1L}

	U	T_x	T_y	L
U	\mathcal{H}	$i \frac{\Delta_y}{2M} \mathcal{E}_T$	$-i \frac{\Delta_x}{2M} \mathcal{E}_T$	
T_x	$i \frac{\Delta_y}{2M} \mathcal{E}$	$\mathcal{H}_T + \frac{\Delta_x^2 - \Delta_y^2}{2M^2} \tilde{\mathcal{H}}_T$	$\frac{\Delta_x \Delta_y}{M^2} \tilde{\mathcal{H}}_T$	$\frac{\Delta_x}{2M} \tilde{\mathcal{E}}$
T_y	$-i \frac{\Delta_x}{2M} \mathcal{E}$	$\frac{\Delta_x \Delta_y}{M^2} \tilde{\mathcal{H}}_T$	$\mathcal{H}_T - \frac{\Delta_x^2 - \Delta_y^2}{2M^2} \tilde{\mathcal{H}}_T$	$\frac{\Delta_y}{2M} \tilde{\mathcal{E}}$
L		$\frac{\Delta_x}{2M} \tilde{\mathcal{E}}_T$	$\frac{\Delta_y}{2M} \tilde{\mathcal{E}}_T$	$\tilde{\mathcal{H}}$

Phase-space representation



Phase-space representation



How to measure Wigner distributions

Many observables sensitive to GTMDs have recently been proposed

- **eA scattering**

- Dijet production
- Longitudinal SSA

[Hatta, Xiao, Yuan (2016)]

[Hatta, Nakagawa, Yuan, Zhao (2016)]

[Ji, Yuan, Zhao (2016)]

- **pA scattering**

- Double parton scattering (DPS)
- Ultra-peripheral collisions (UPCs)

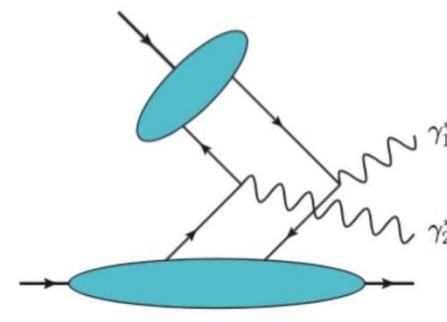
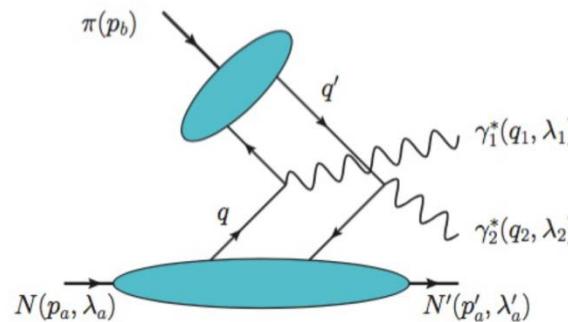
[Hagiwara, Hatta, Xiao, Yuan (2017)]

[Hagiwara *et al.* (2017)]

- **πN scattering**

- Exclusive double Drell-Yan

[Bhattacharya, Metz, Zhou (2017)]

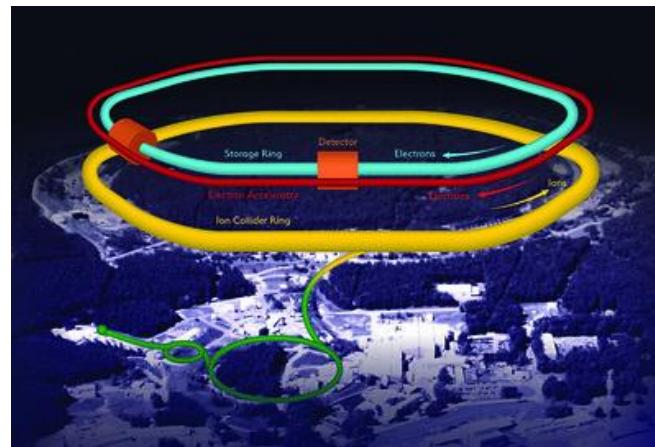


Typically complicated final state →

very small cross sections !

Summary

- **Several generalizations of PDFs** (TMDs, GPDs, GTMDs, ...)
- **Multidimensional imaging of nucleon structure**
- **Indirect access to energy-momentum tensor** (mass, spin, ...)
- **Further extensions** (gluon distributions, higher spin, multiparton scattering, ...)



So much to learn thanks to the EIC !