

Neutrinos: the Ghosts of Christmas Past and Present



$$m_\nu = \frac{Y_\nu^2 v^2}{M_N}$$



$$\sigma = \frac{G_F^2 S}{\pi}$$

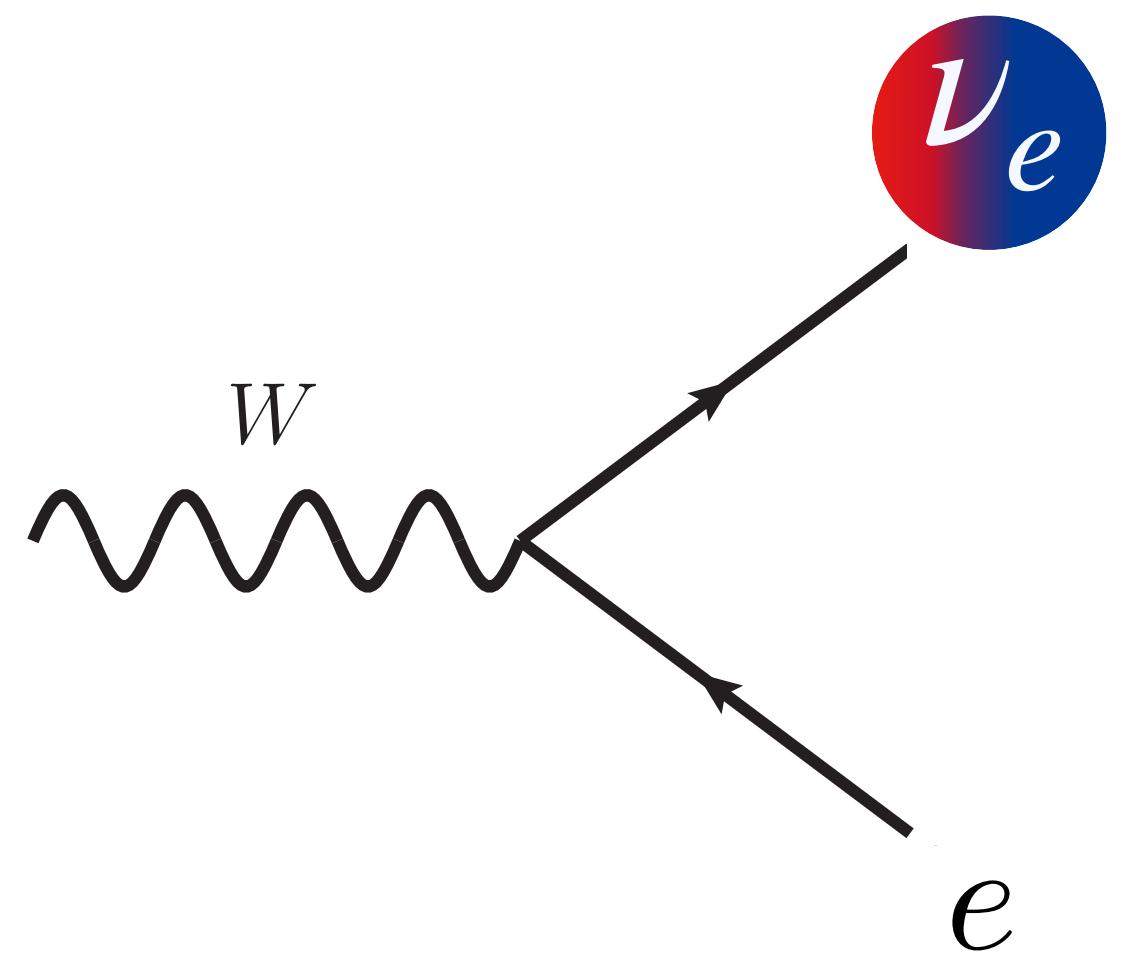


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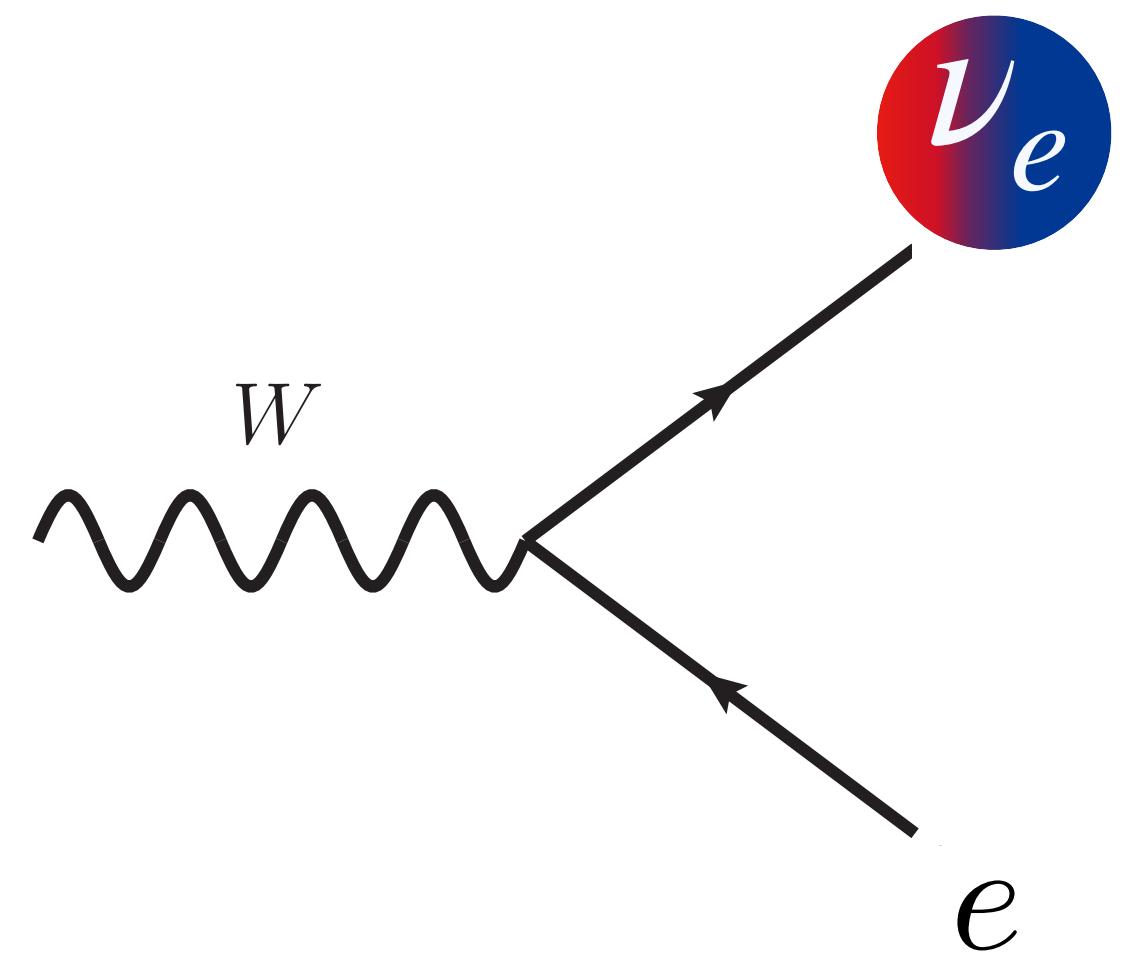
Neutrino Oscillations

Production

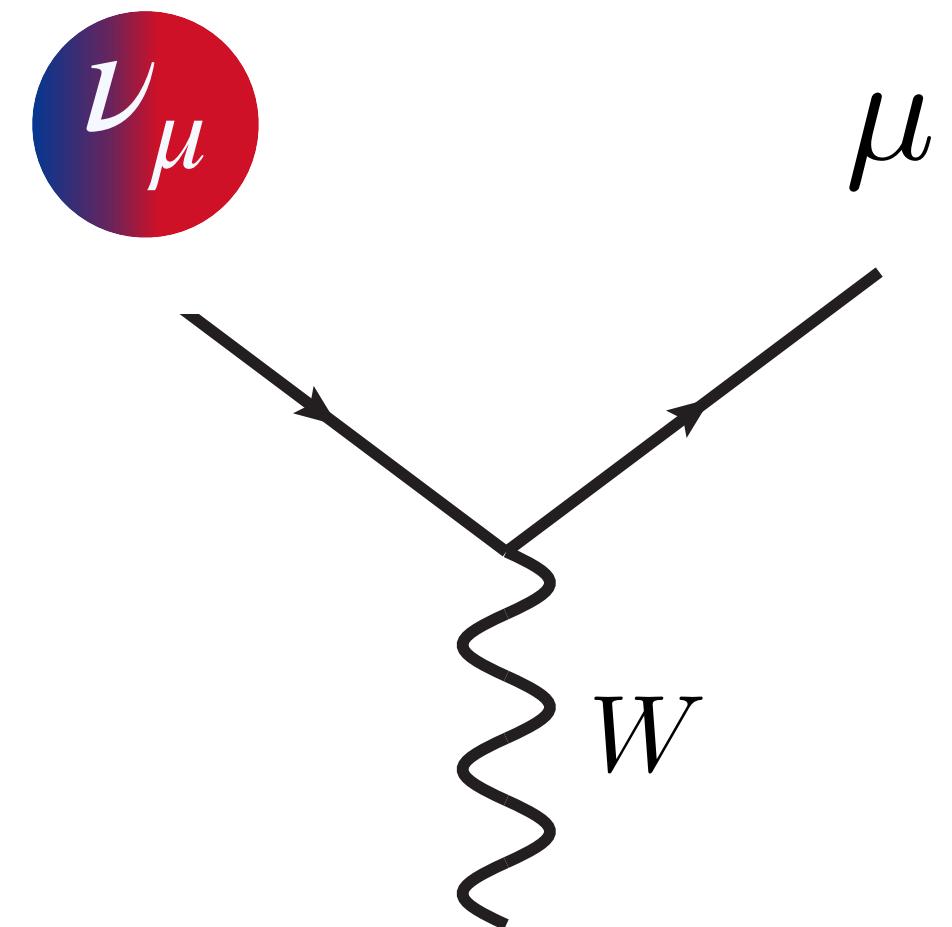


Neutrino Oscillations

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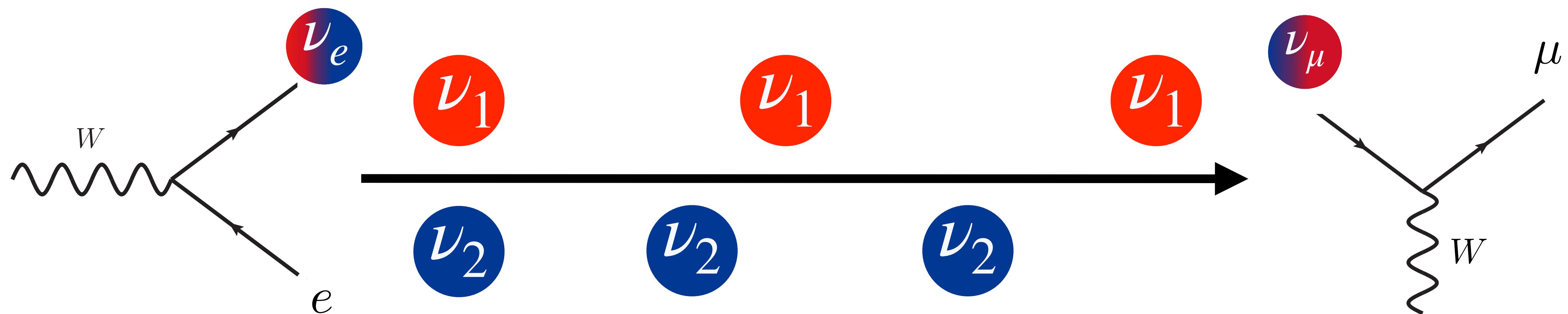


Detection



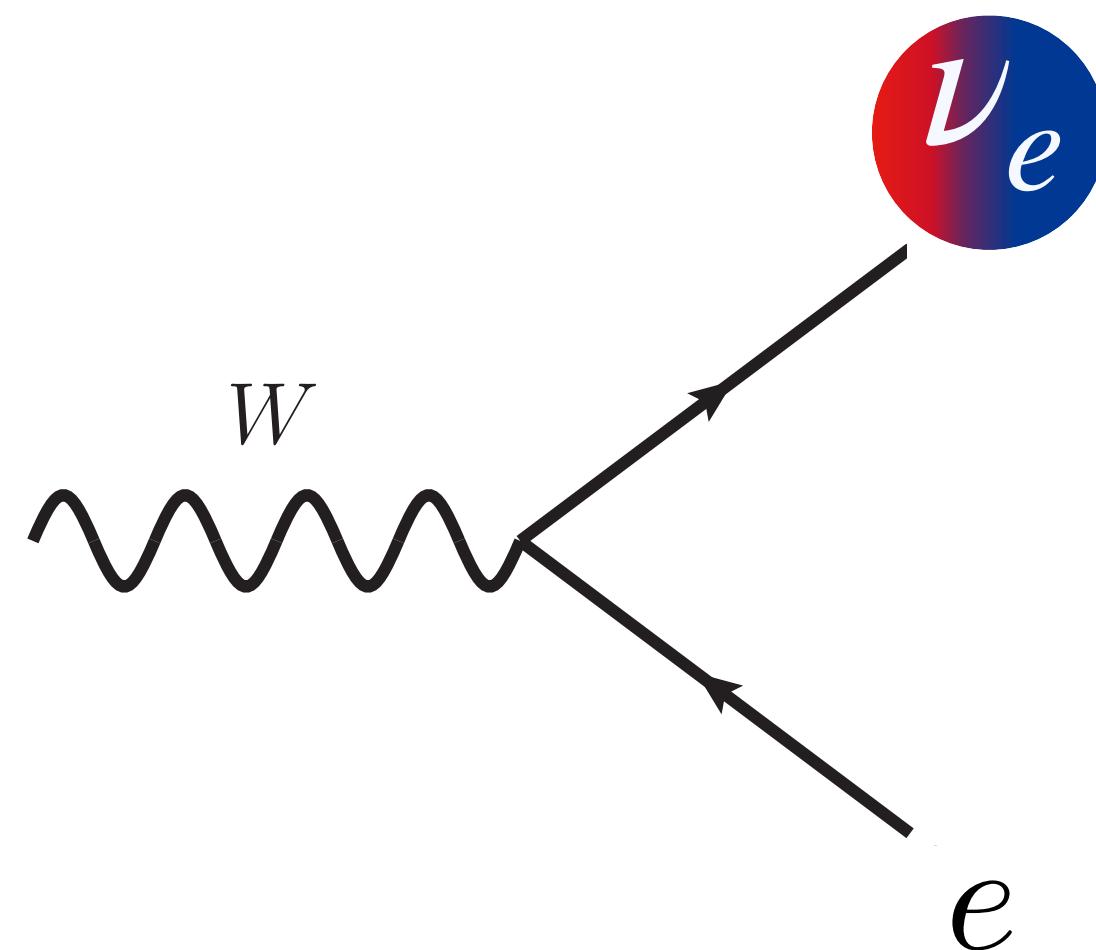
Neutrino Oscillations

Production Propagation Detection

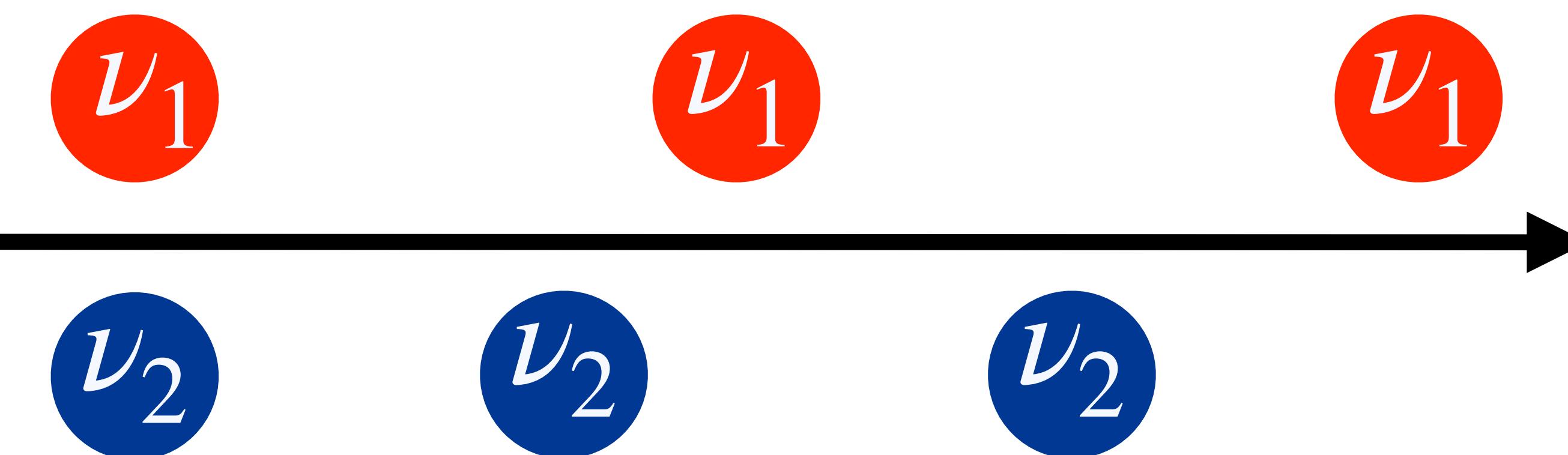


Neutrino Oscillations

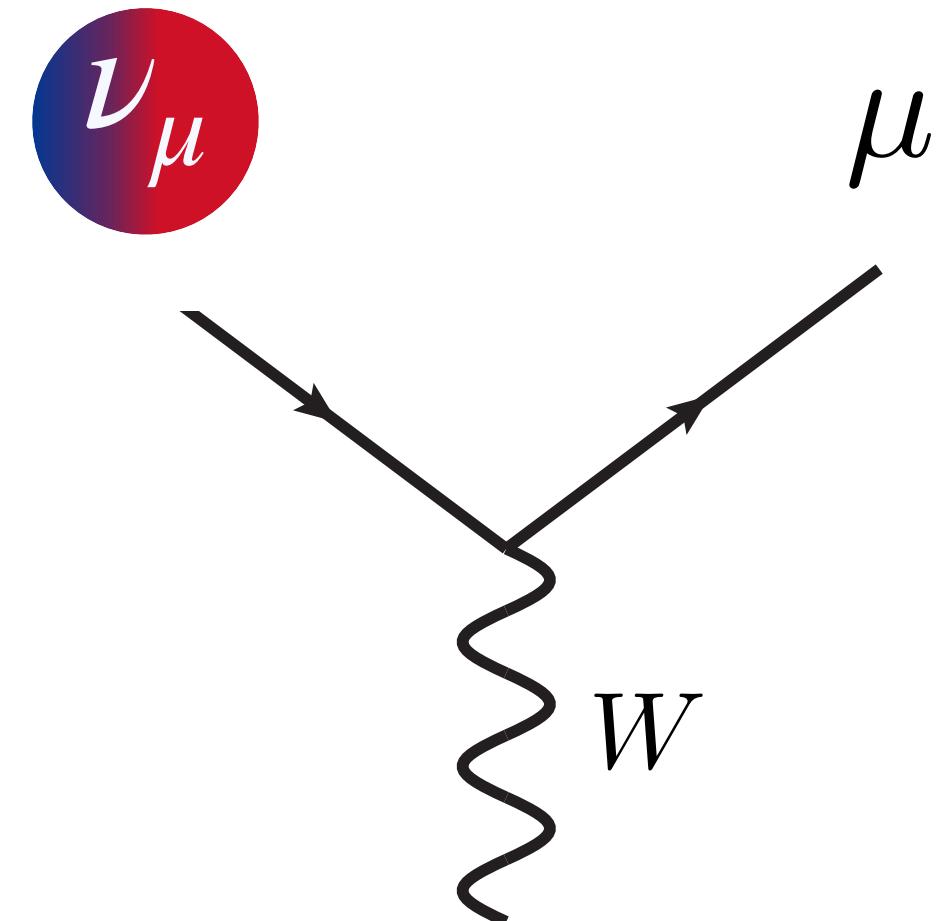
Production



Propagation



Detection



$$|\nu_e\rangle = \sum_i U_{ei} |\nu_i\rangle$$

Flavour states
superposition
massive states

$$\begin{aligned}\nu_1 &: e^{-iE_1 t} \\ \nu_2 &: e^{-iE_2 t} \\ \nu_3 &: e^{-iE_3 t}\end{aligned}$$

Massive states:
Hamiltonian
eigenstates

$$\langle \nu_\mu |$$

Projection over
flavour states

Neutrino Oscillations

Simplified two neutrino case

$$|\nu(t)\rangle = -\sin(\theta)e^{-iE_1 t}|\nu_1\rangle + \cos(\theta)e^{-iE_2 t}|\nu_2\rangle$$

Probability for ν_e to transform to ν_μ

$$P(\nu_e \rightarrow \nu_\mu) = \sin^2(2\theta) \sin^2 \frac{(m_2^2 - m_1^2) L}{4E}$$

Neutrino Oscillations

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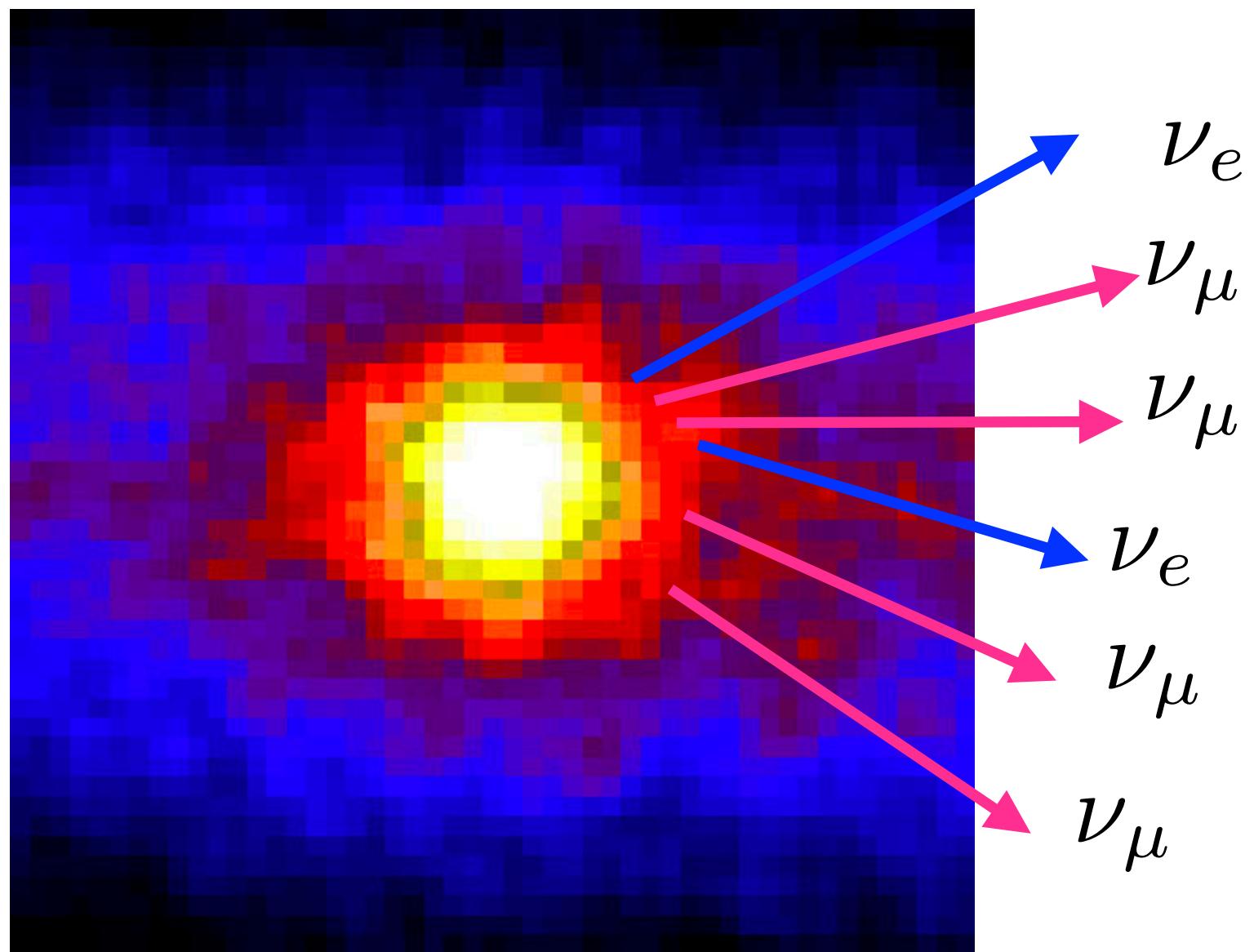
$$P(\nu_e \rightarrow \nu_\mu) = \sin^2(2\theta) \sin^2 \frac{(m_2^2 - m_1^2) L}{4E}$$

Mixing angles: misalignment
between massive and flavour
states

Neutrino masses

Neutrinos: Ghosts of Christmas Past

Ghosts of Christmas Past: Solar Neutrinos



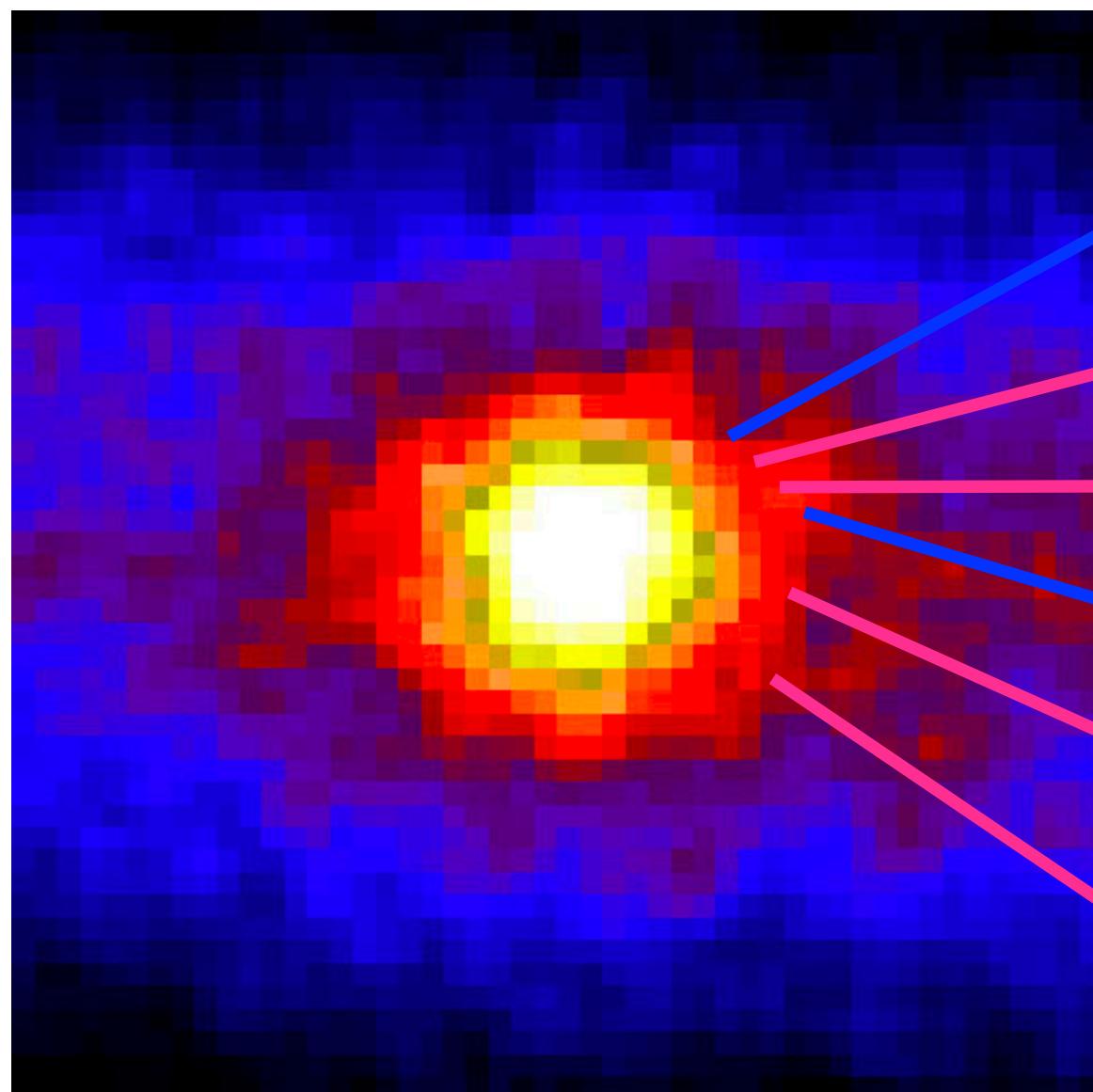
- Sun emits $\mathcal{O}(\text{MeV}) \nu_e$ via nuclear fusion
- Homestake searched for solar neutrinos (1970)

$$\nu_e + {}^{37}\text{Cl} \longrightarrow {}^{37}\text{Ar}^+ + e^-$$
- Bahcall calculated solar neutrino flux. Counting radiative argon \implies deficit

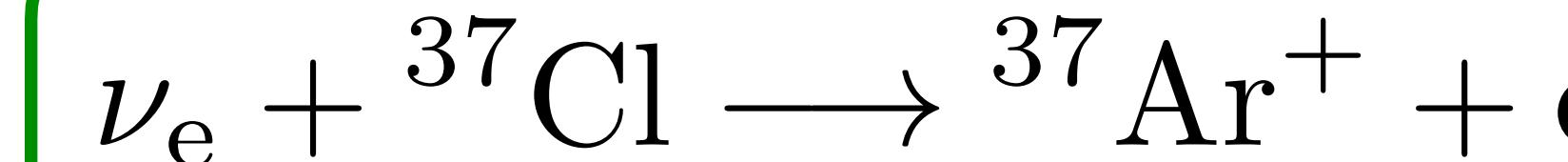


Nobel 2002

Ghosts of Christmas Past: Solar Neutrinos

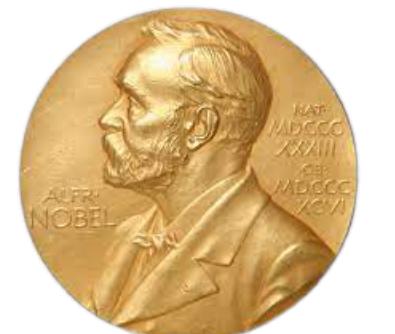
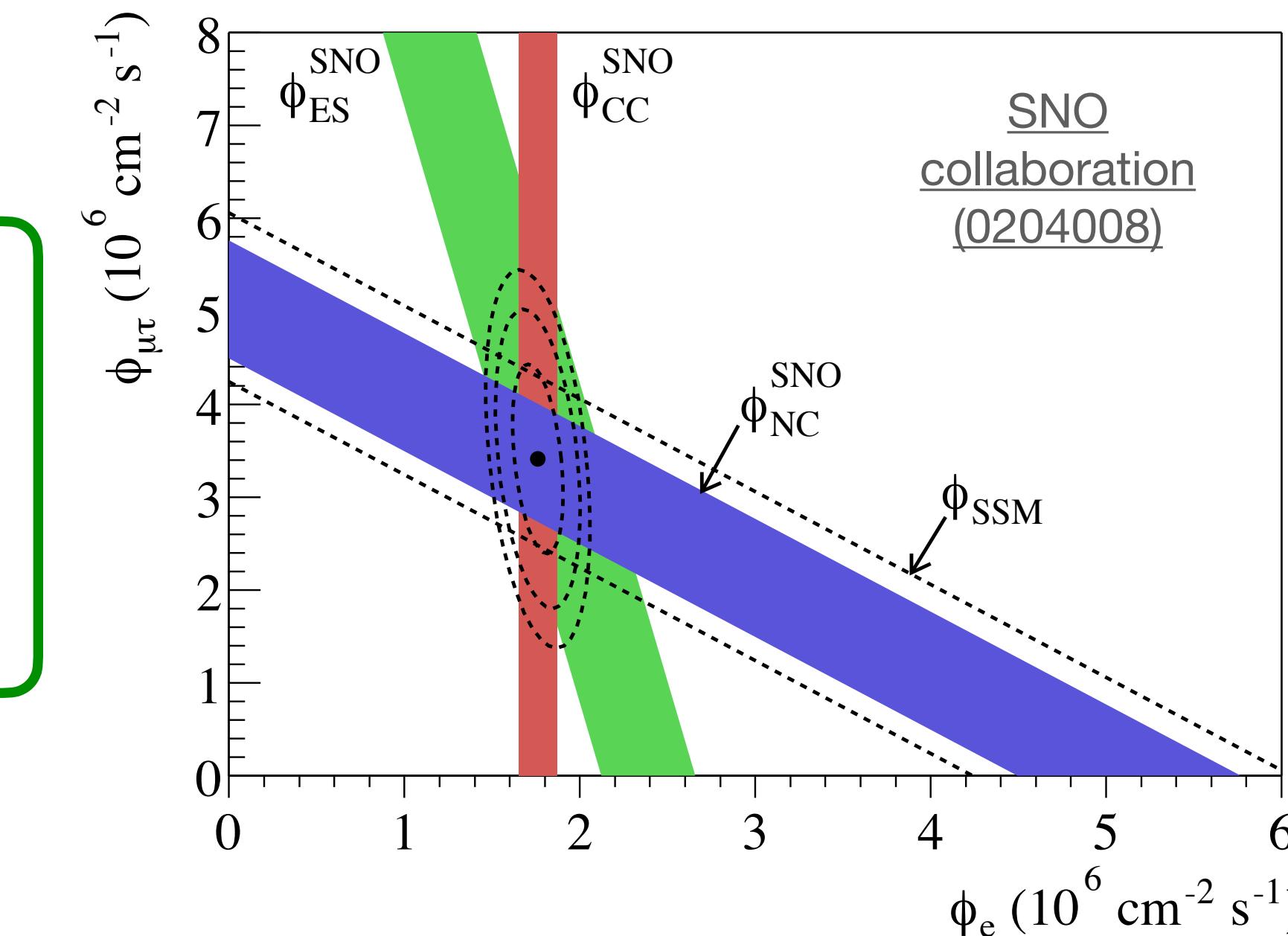
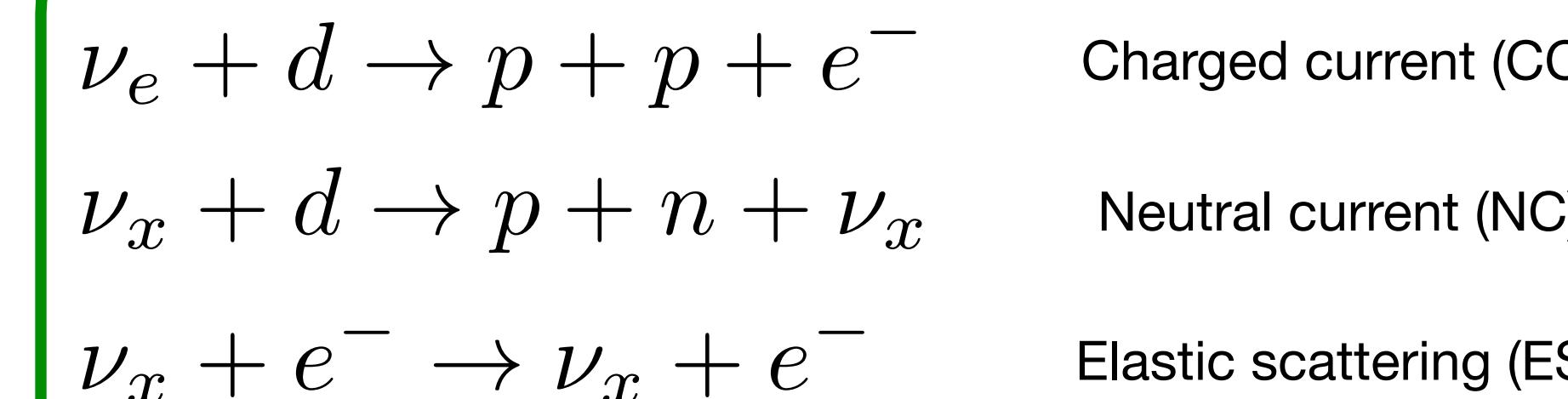


- Sun emits $\mathcal{O}(\text{MeV}) \nu_e$ via nuclear fusion
- Homestake searched for solar neutrinos (1970)



Nobel 2002

- SNO (2002) resolved solar neutrino problem by detecting solar ν_e appearance and solar ν_e disappearance



Nobel 2015

Ghosts of Christmas Past: Atmospheric Neutrinos

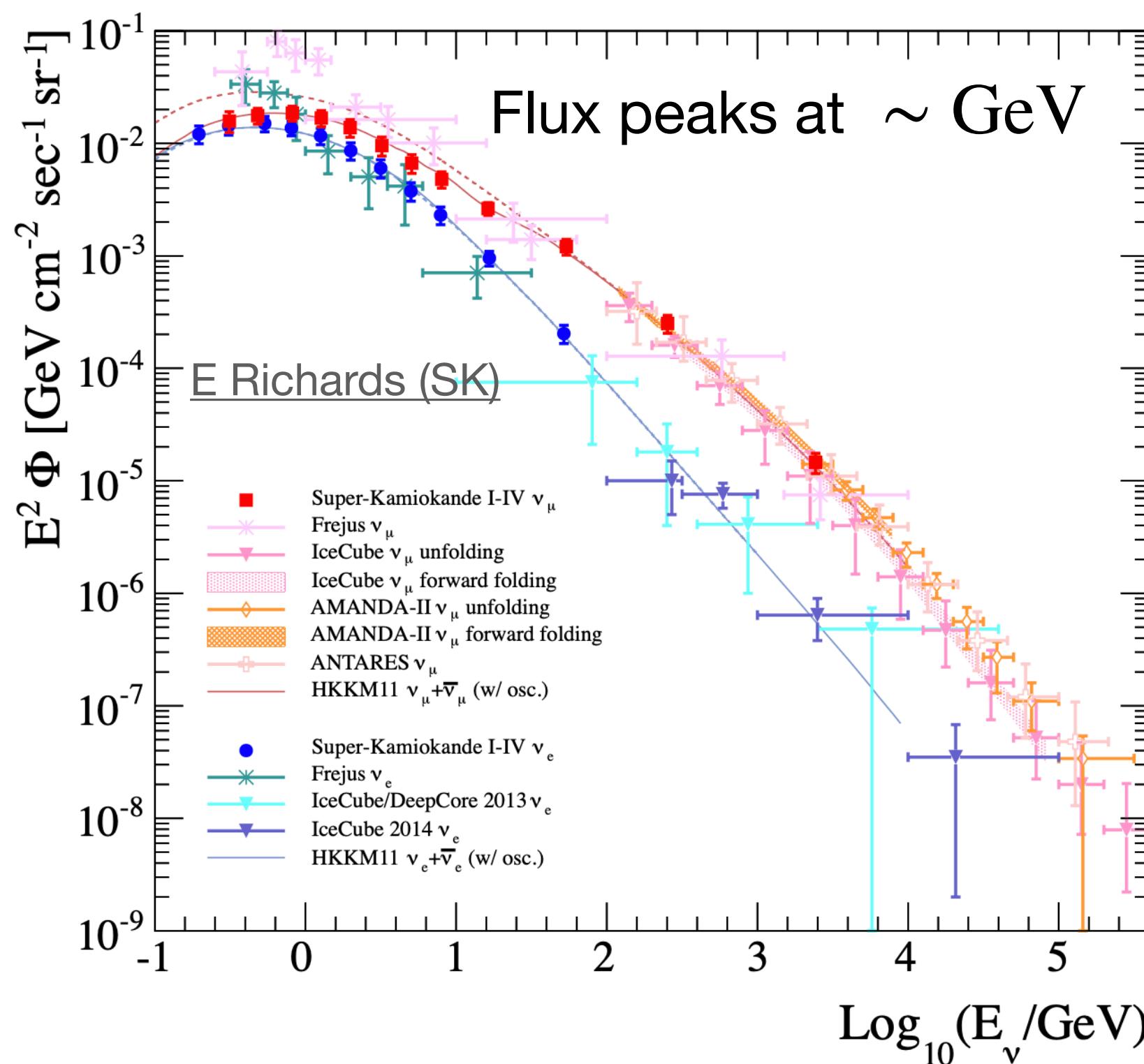
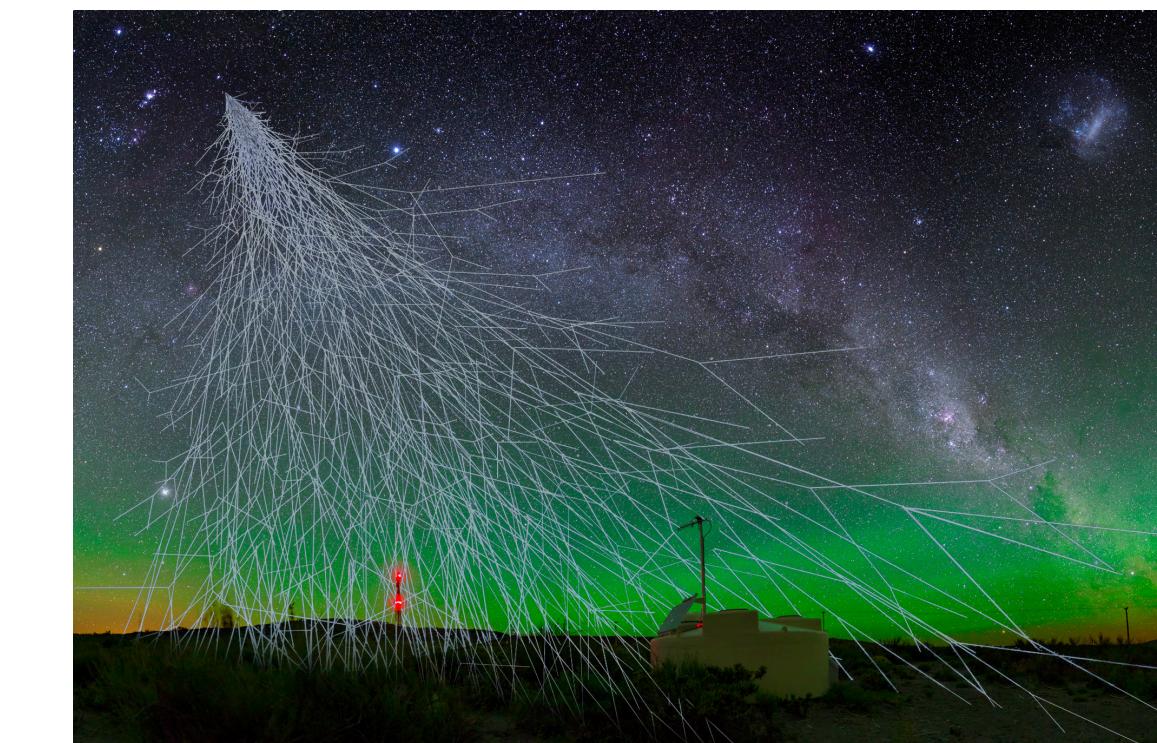
Decay of mesons produce flux of atmospheric neutrinos which spans MeV - TeV scale

$$\pi^\pm \rightarrow \mu^\pm + \nu_\mu (\bar{\nu}_\mu)$$

$$K^\pm \rightarrow \mu^\pm + \nu_\mu (\bar{\nu}_\mu)$$

$$\mu^\pm \rightarrow e^\pm + \nu_e (\bar{\nu}_e) + \bar{\nu}_\mu (\nu_\mu)$$

$$\frac{\nu_e + \bar{\nu}_e}{\nu_\mu + \bar{\nu}_\mu} \sim \frac{1}{2} \quad E_\nu \lesssim 1 \text{ GeV}$$



DETECTION OF MUONS PRODUCED BY COSMIC RAY NEUTRINOS DEEP UNDERGROUND

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Tata Institute of Fundamental Research, Colaba, Bombay

K. HINOTANI and S. MIYAKE,
Osaka City University, Osaka, Japan

D. R. CREED, J. L. OSBORNE, J. B. M. PATTISON and A. W. WOLFENDALE
University of Durham, Durham, U.K.

Received 12 July 1965

First detection in 1965

Ghosts of Christmas Past: Atmospheric Neutrinos

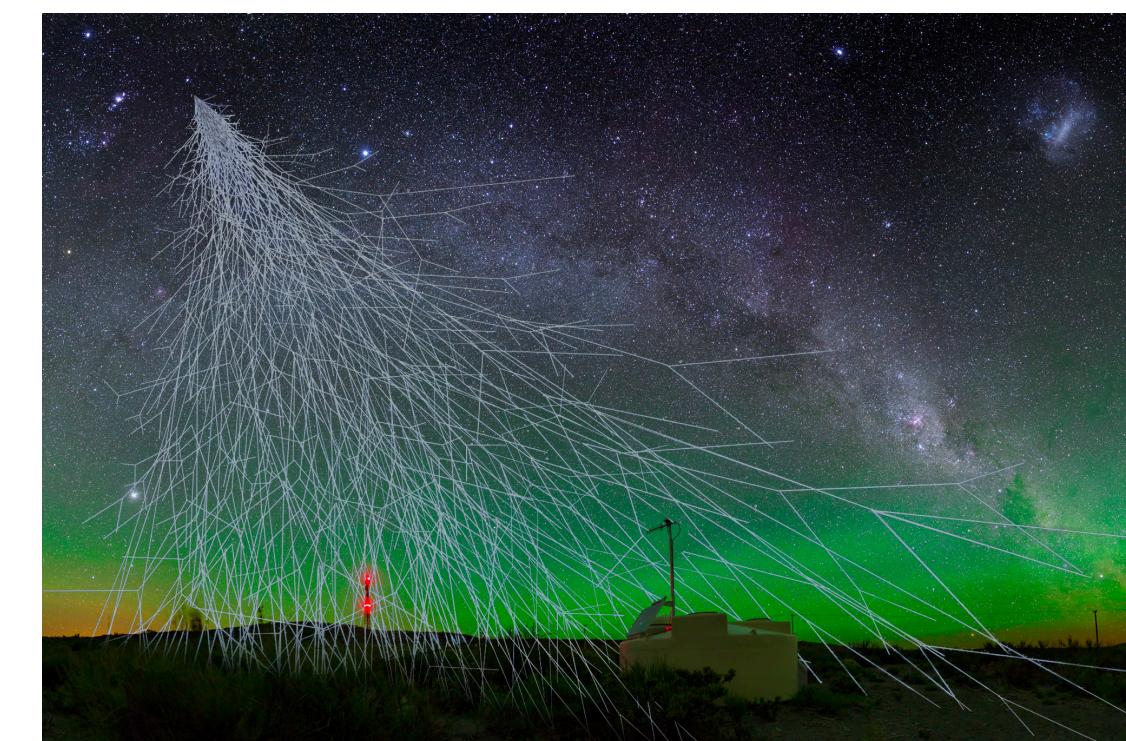
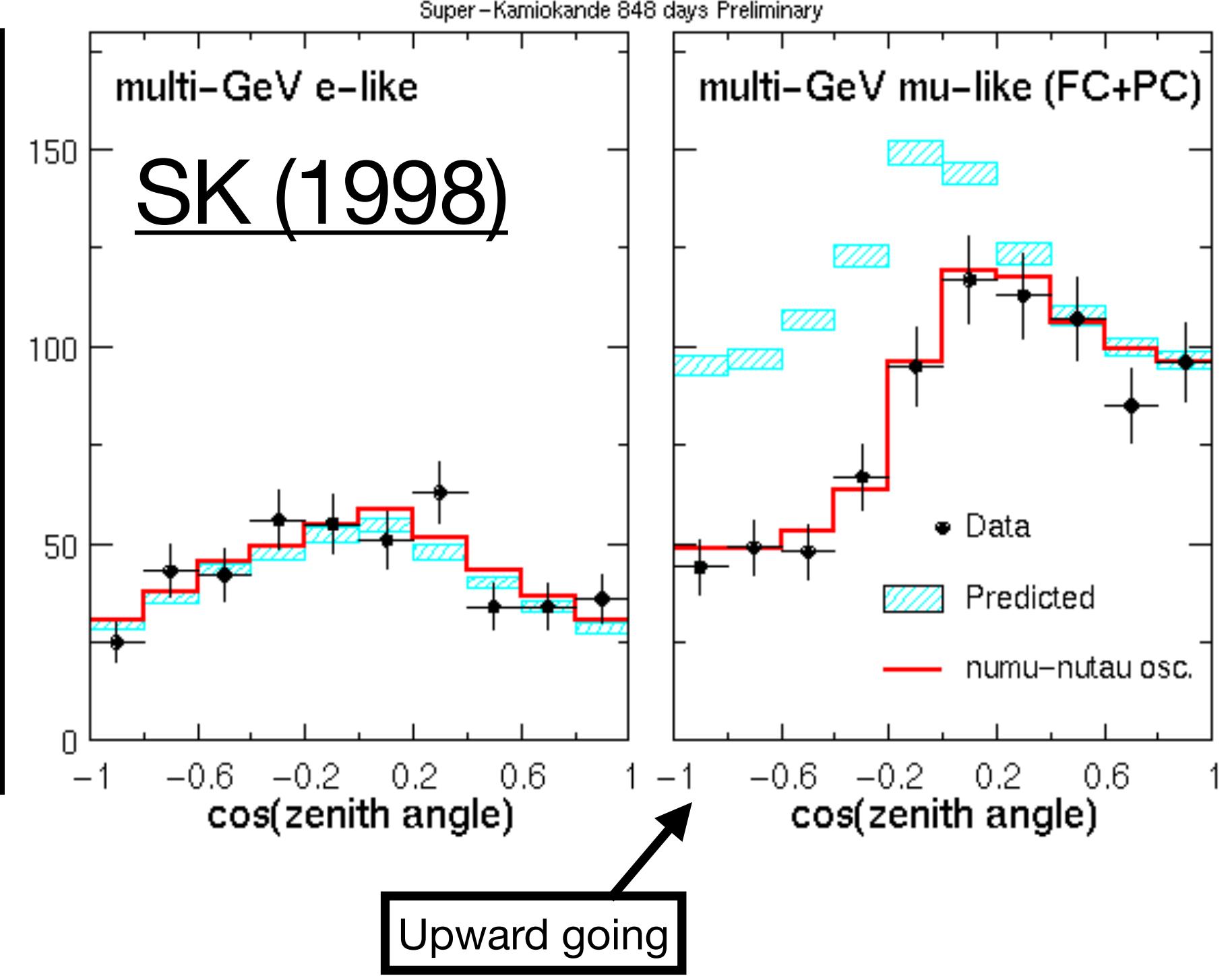
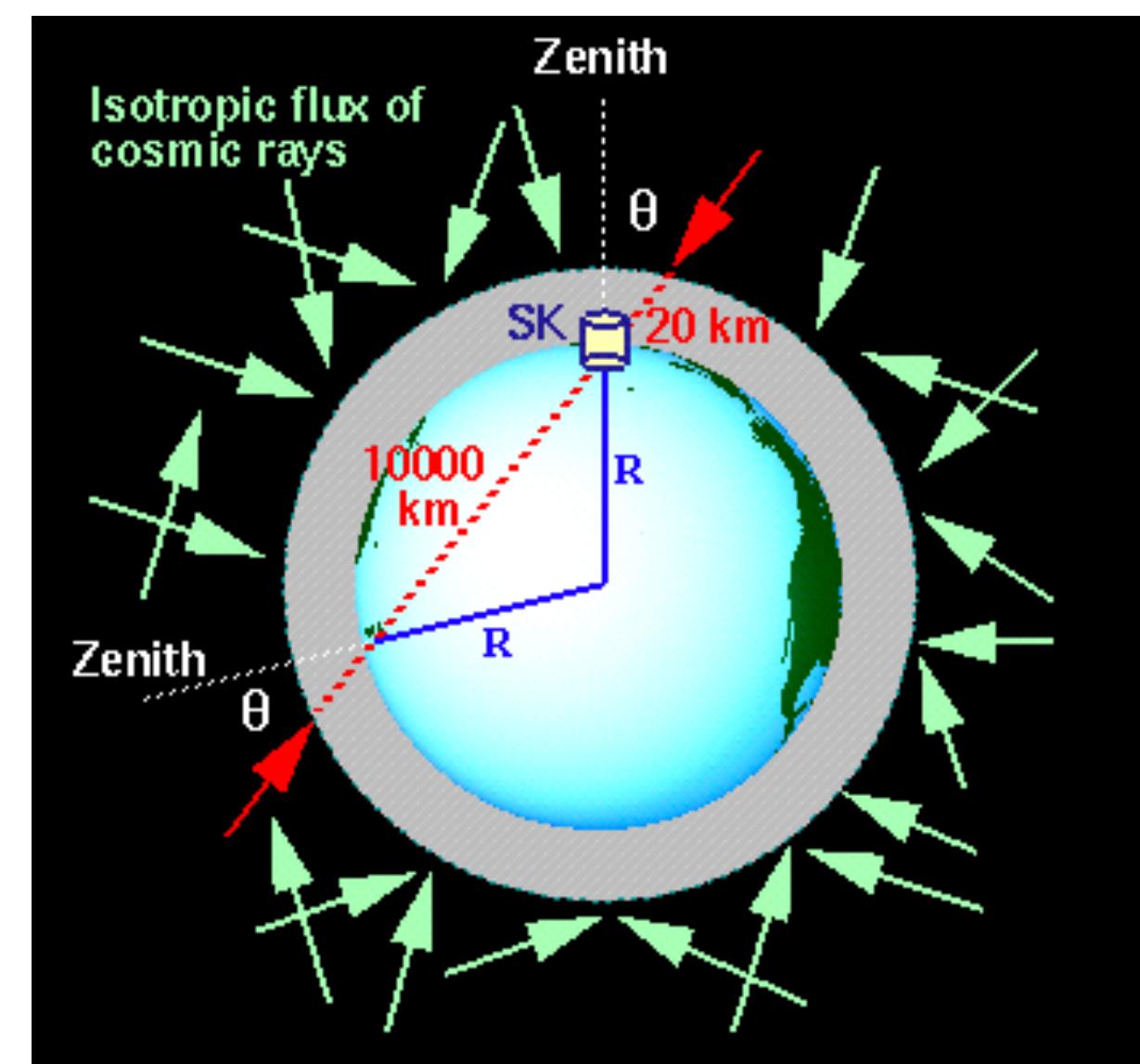
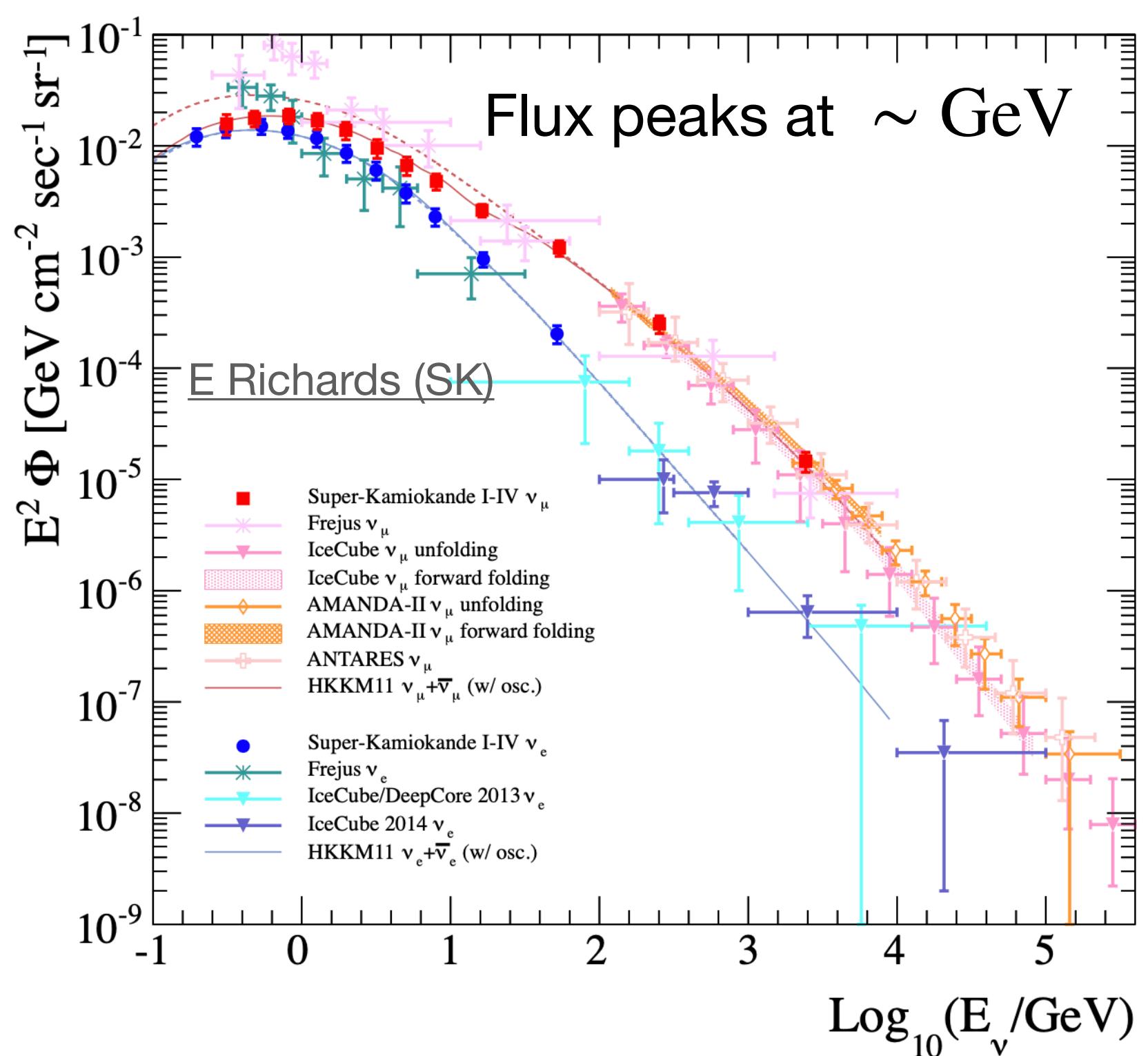
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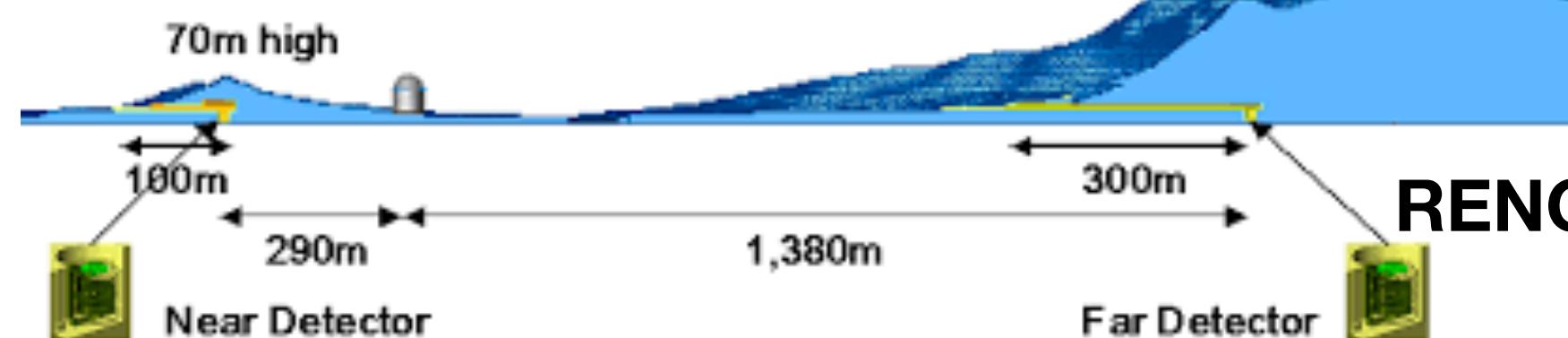
Nobel 2015

Ghosts of Christmas Past: Reactor Neutrinos

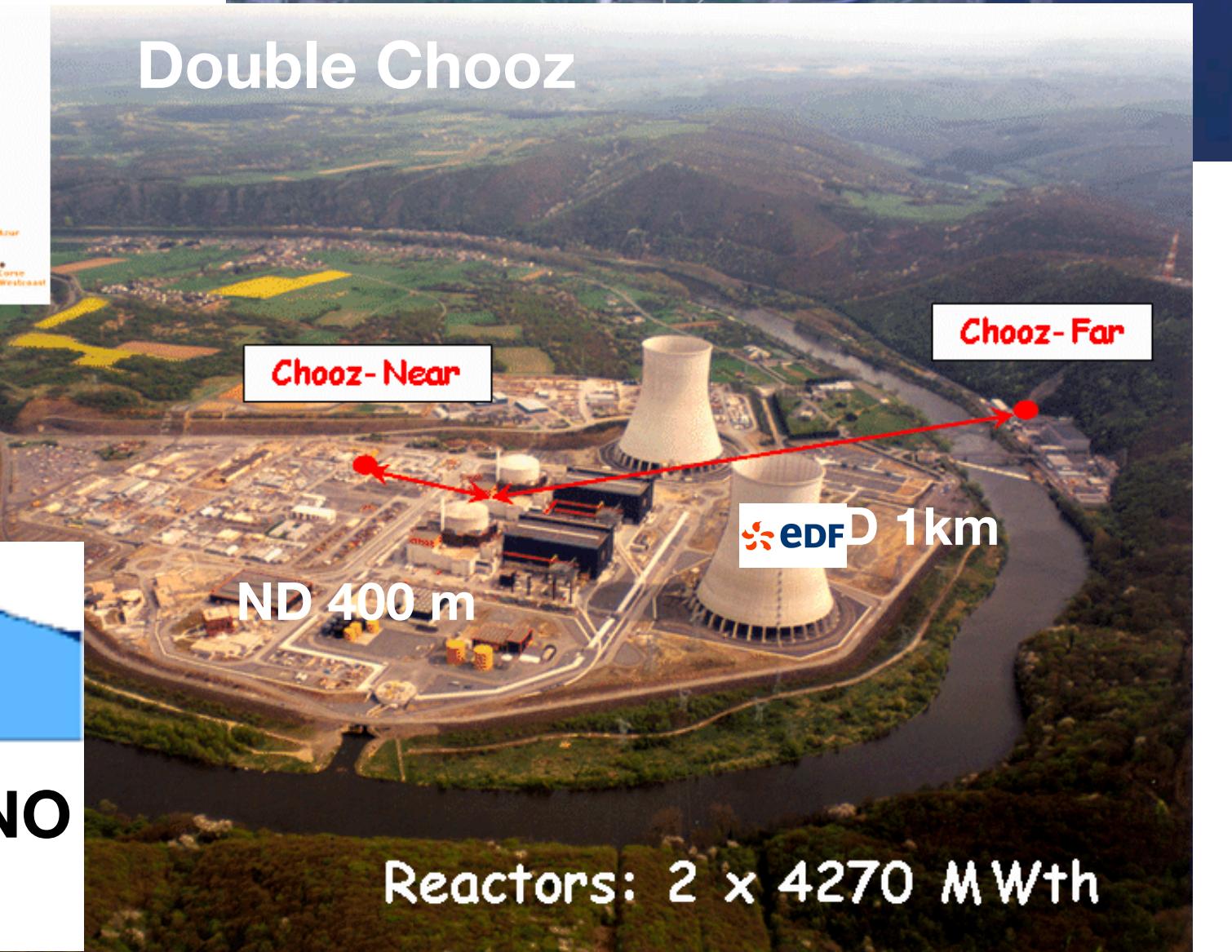
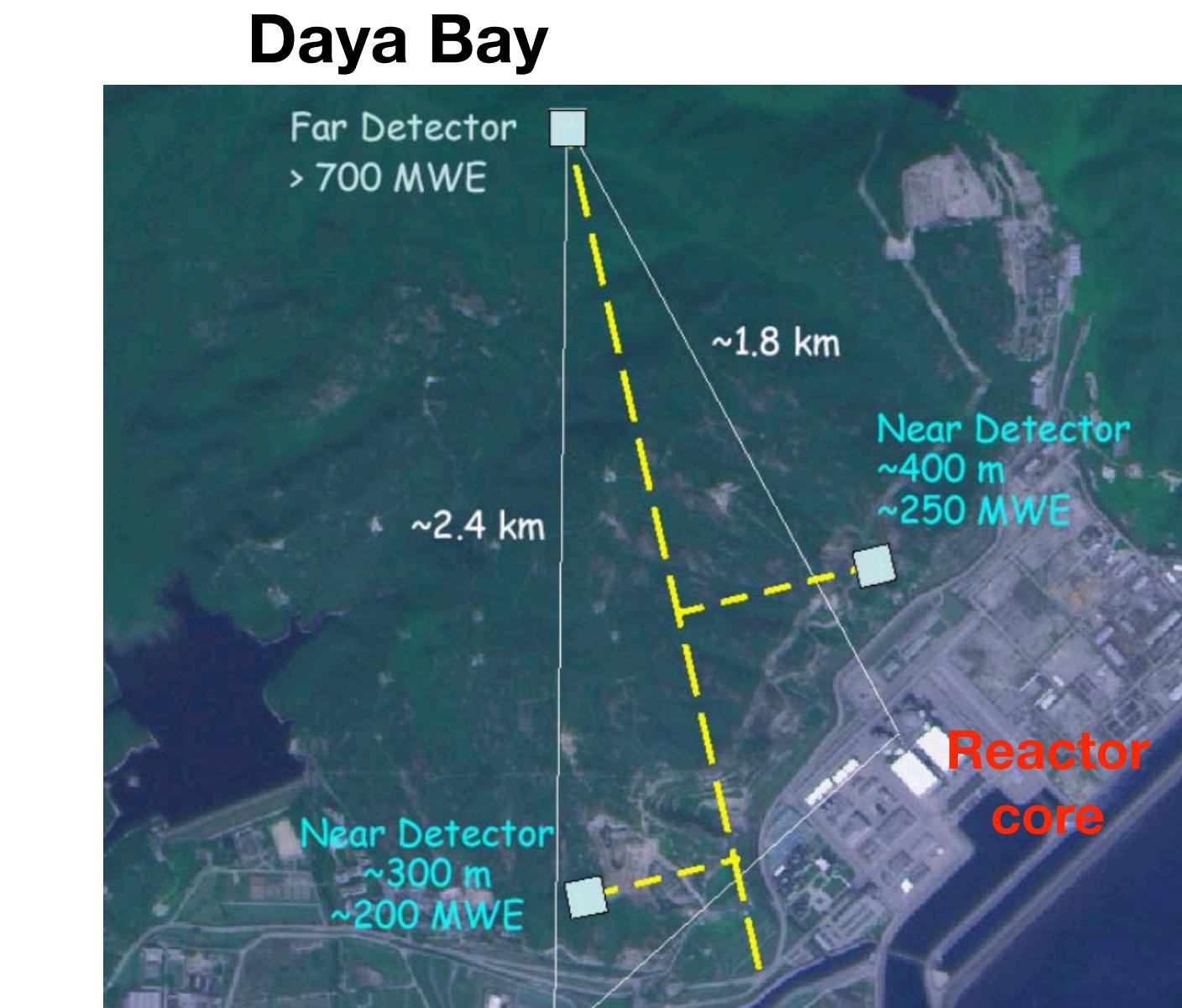
- Reactors produce MeV scale $\bar{\nu}_e$
- Detectors km baseline using IBD to detect neutrinos: $\bar{\nu}_e + p \longrightarrow e^+ + n$

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2(2\theta_{13}) \sin^2 \frac{\Delta m_{31}^2 L}{4E}$$

- 2012 big year for particle physics precision measurement on $\theta_{13} \approx 8^\circ$ by reactor experiments

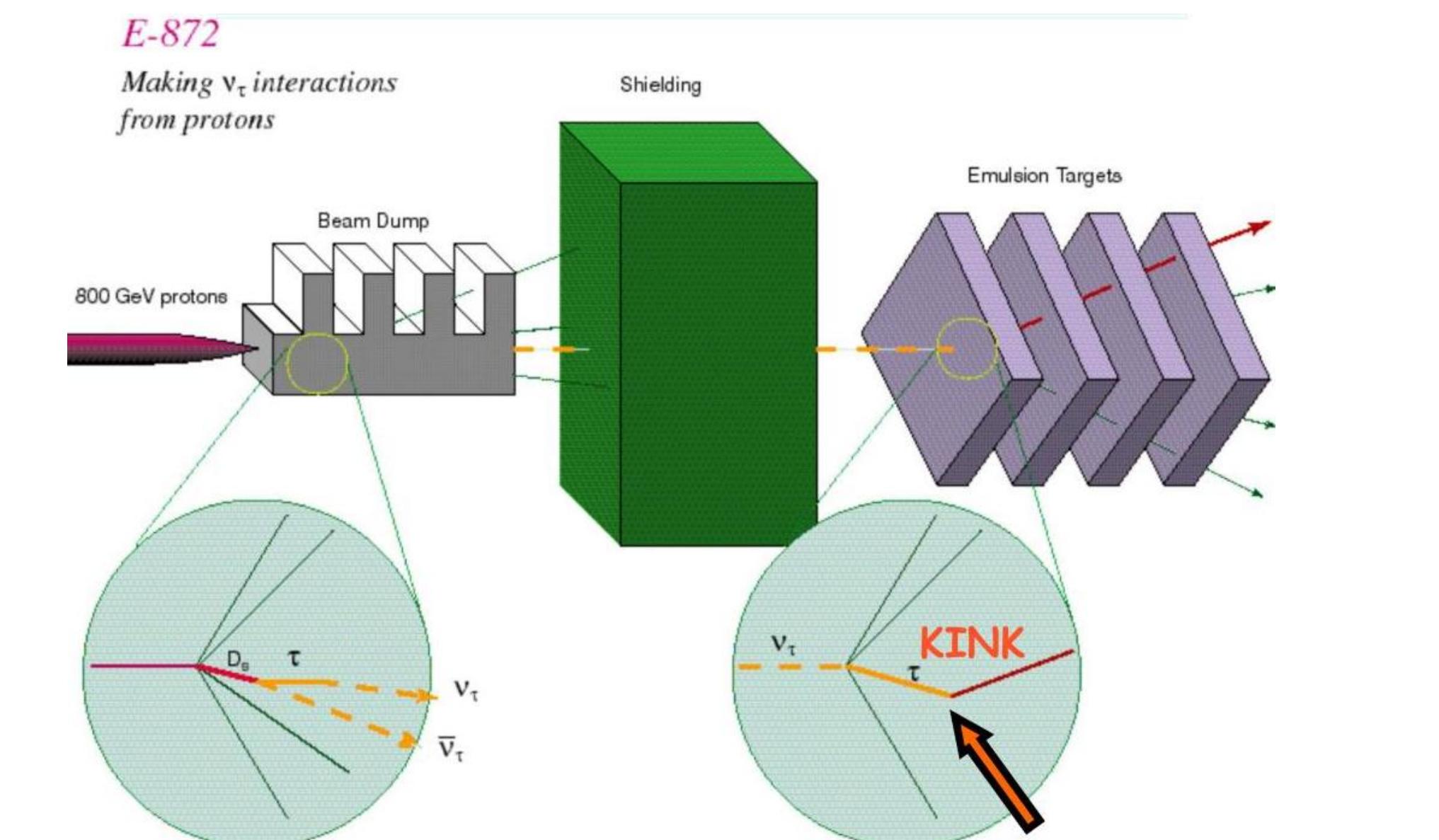
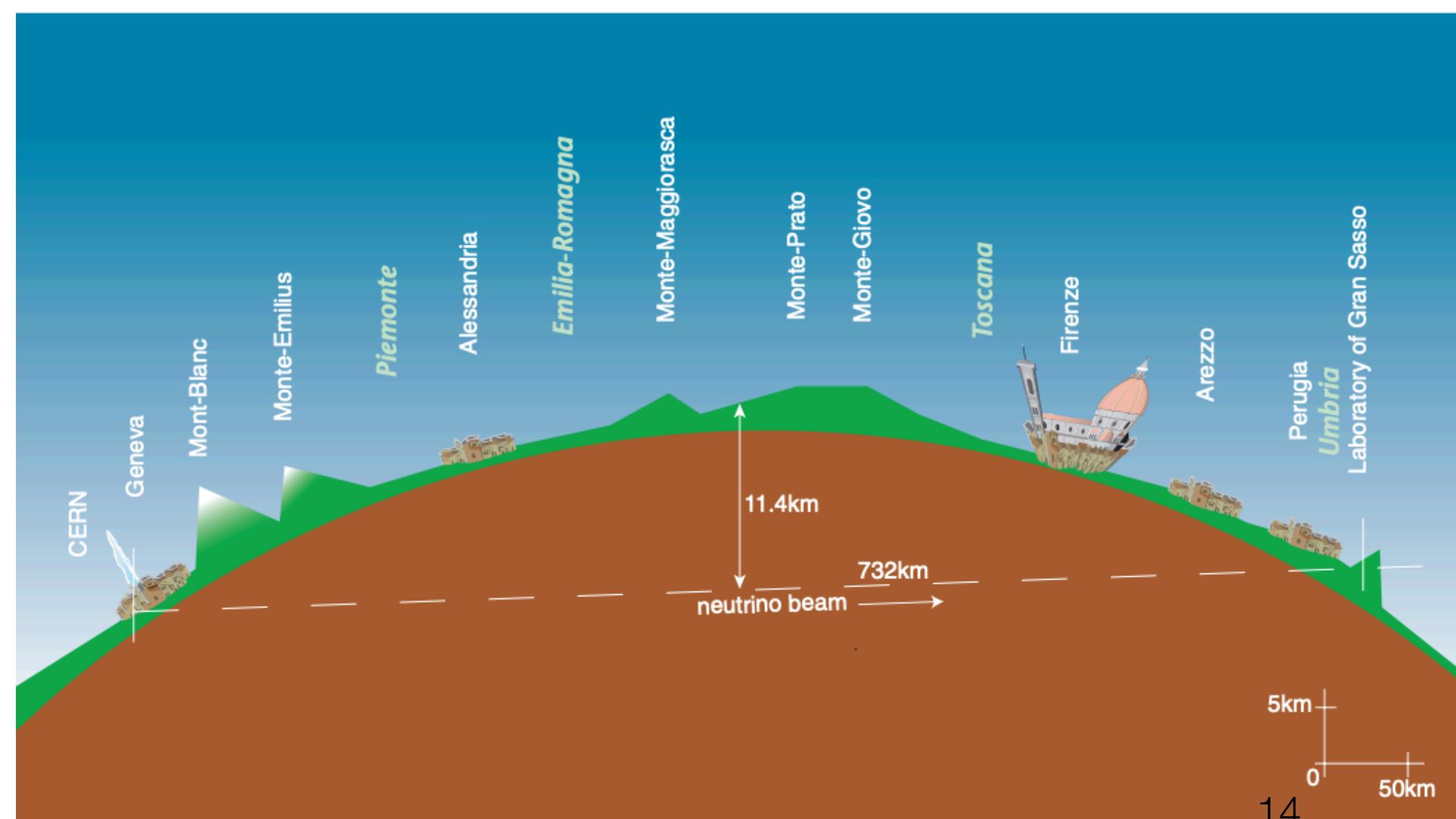


13

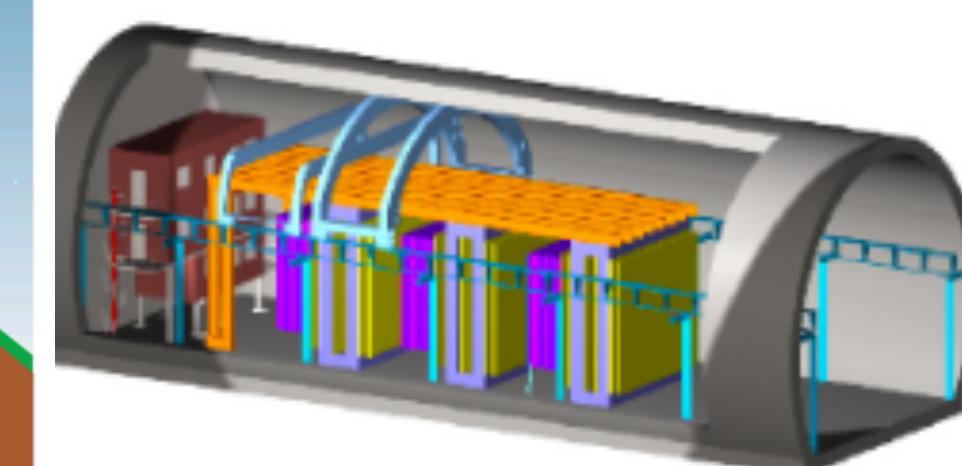


Ghosts of Christmas Past: Tau Neutrinos

- Only 14 ν_τ have ever been observed: 4 by DONUT, direct observation (1997-2000)
- 10 ν_τ observed by OPERA (2010-2012): ν_μ beam oscillates to ν_τ



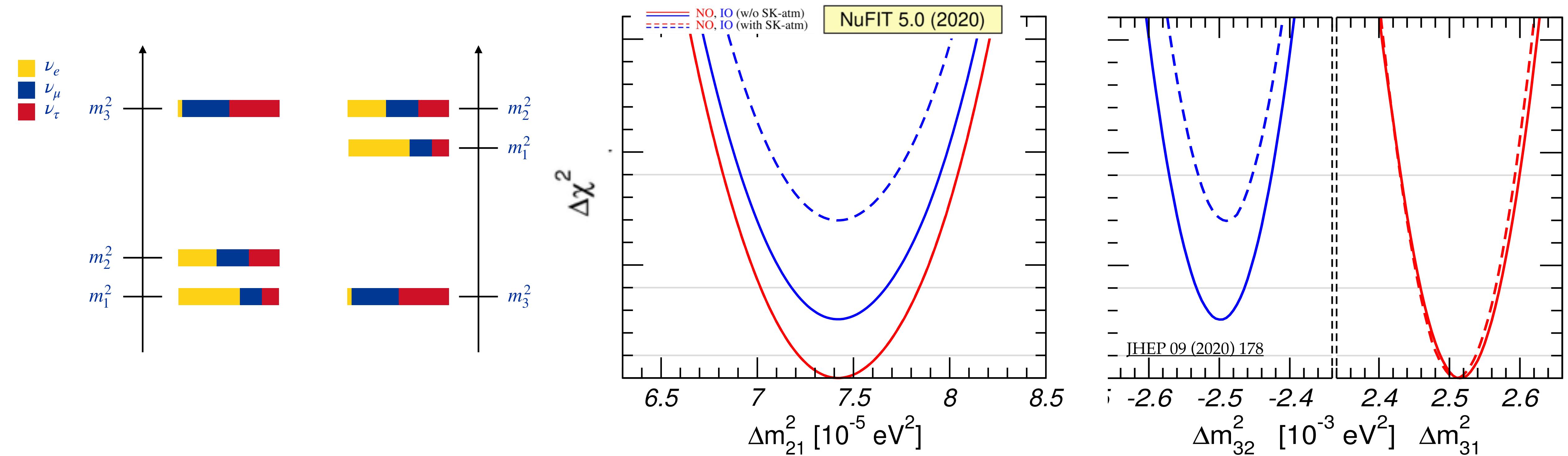
OPERA detector
emulsion cloud chamber



Both DONUT
And OPERA could
Observe mm-scale
Tau kink

*Neutrinos: Ghosts of
Christmas Present*

Ghost of Christmas Present: Current Global Picture



Many neutrino experiments, here $\Delta\chi^2$ from NuFIT collaboration but there are also [Valencia \(2006.11237\)](#) and [Bari \(2003.08511\)](#) groups

Ghost of Christmas Present: Current Global Picture

Atmospheric

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

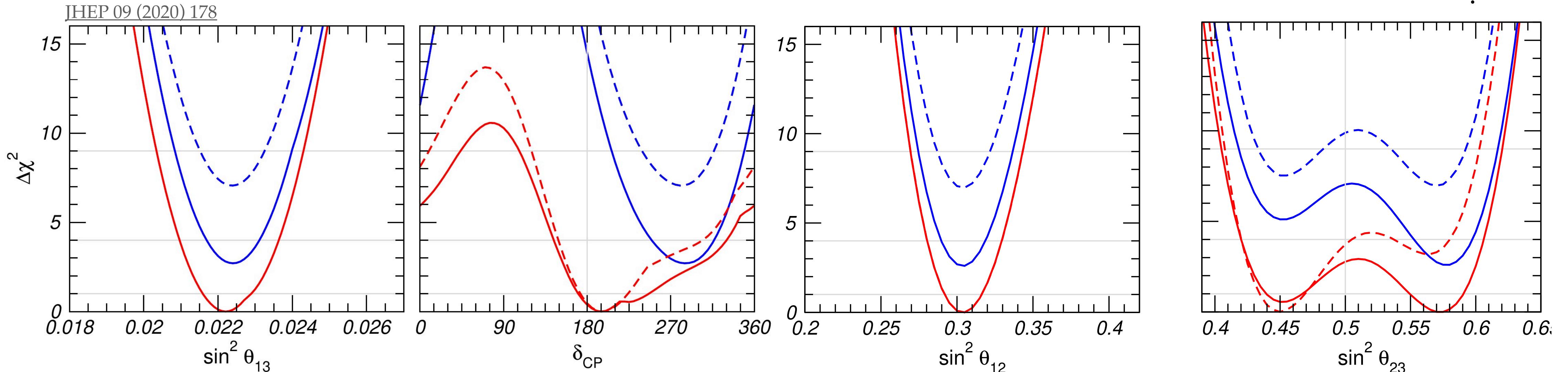
Reactor

$$\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix}$$

Solar

$$\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\underbrace{\begin{pmatrix} e^{in_1} & 0 & 0 \\ 0 & e^{in_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_?$$



$$\theta_{13,\text{bf}} = 8.57^\circ$$

$$\delta_{CP,\text{bf}} = 194^\circ$$

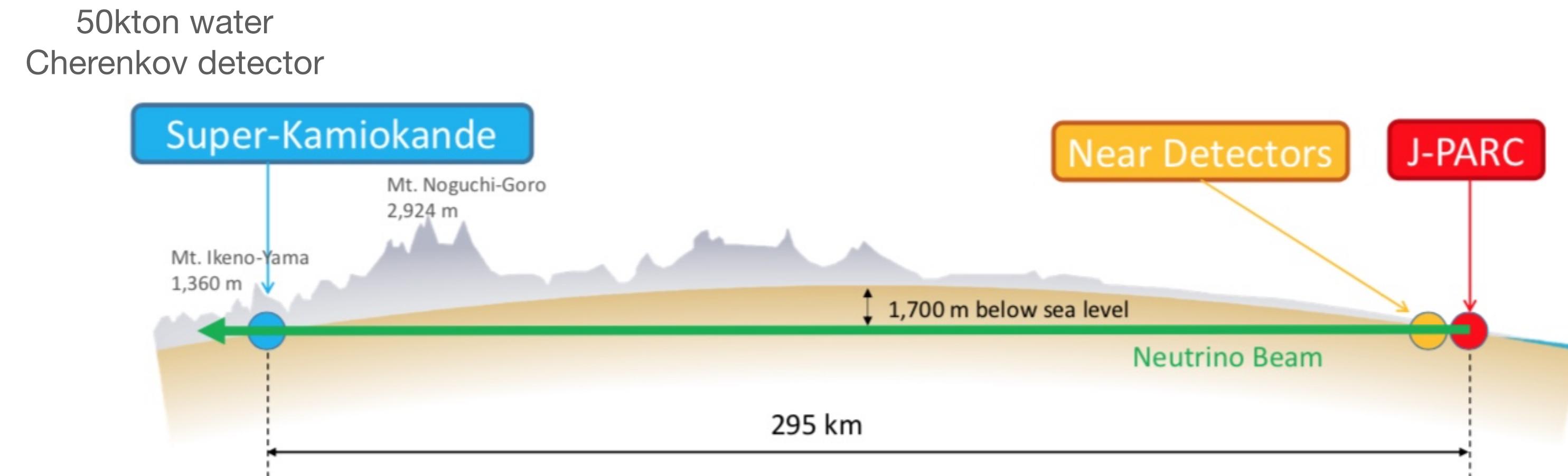
$$\theta_{12,\text{bf}} = 33.44^\circ$$

$$\theta_{23,\text{bf}} = 49.2^\circ$$

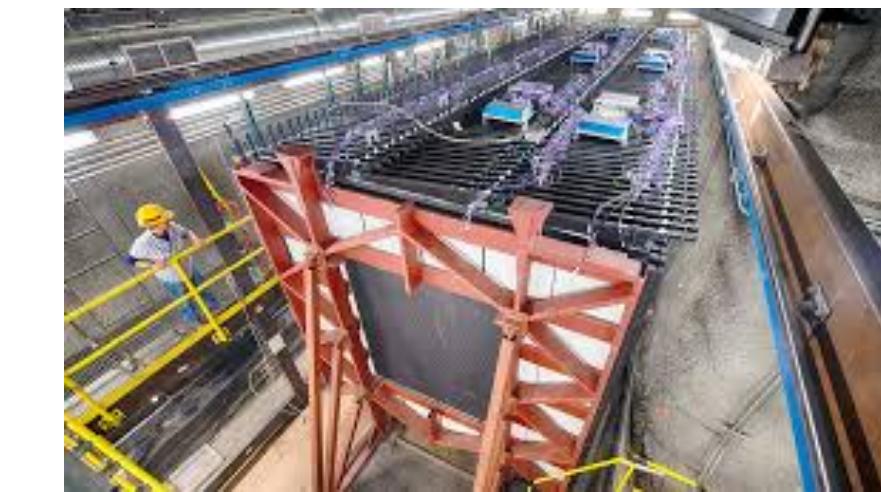
Little direct observation of tau neutrino contributes significantly to poor knowledge of mixing angles and if mixing is unitary.

\mathcal{CP} -violating phase & mass ordering

NOvA and T2K are **long baseline, accelerator oscillation** experiment measure δ_{CP} and the mass ordering.



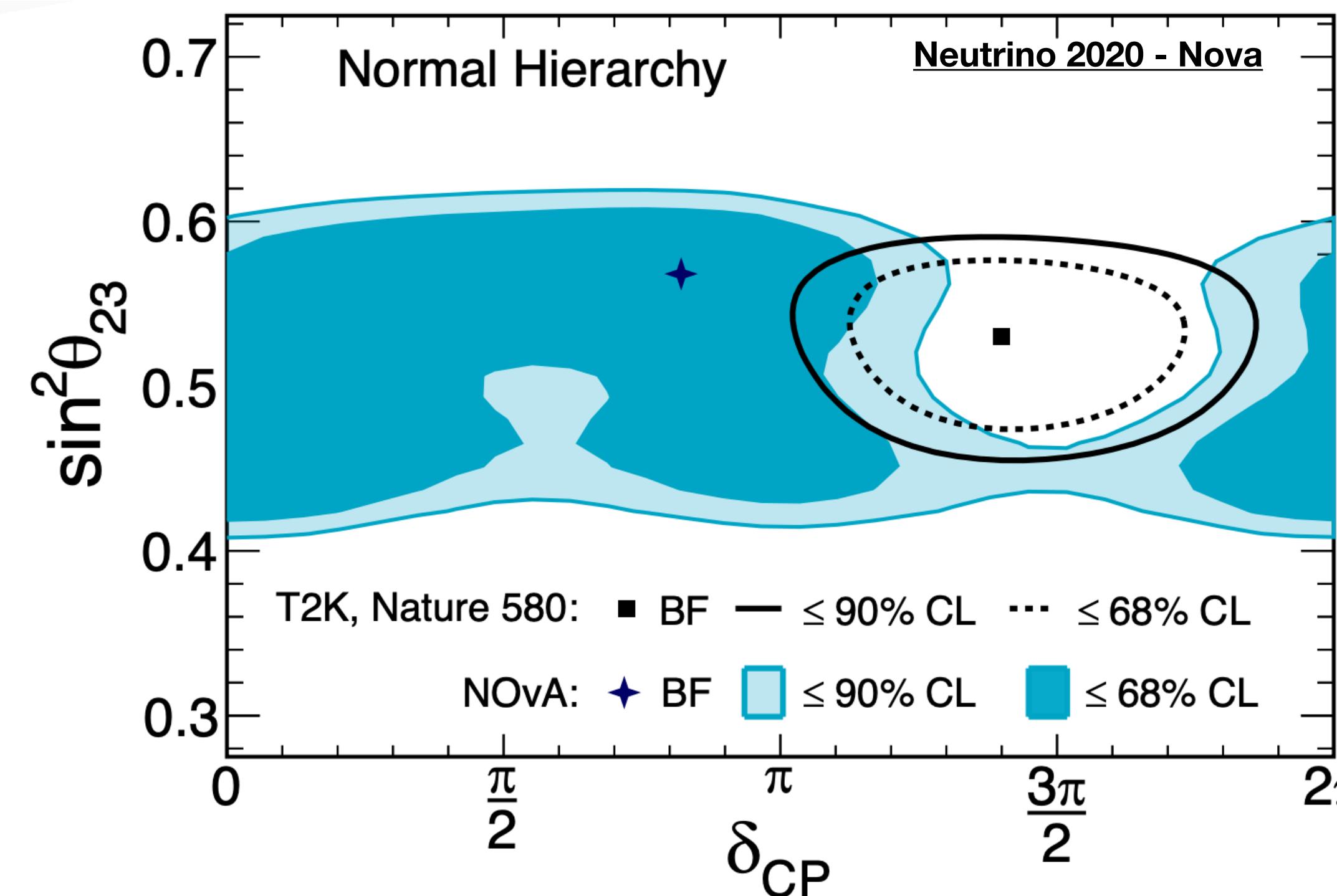
14kton liquid scintillator
detector



$$P(\nu_\mu \rightarrow \nu_e) - P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \propto \delta_{CP}$$

\mathcal{CP} -violating phase & mass ordering

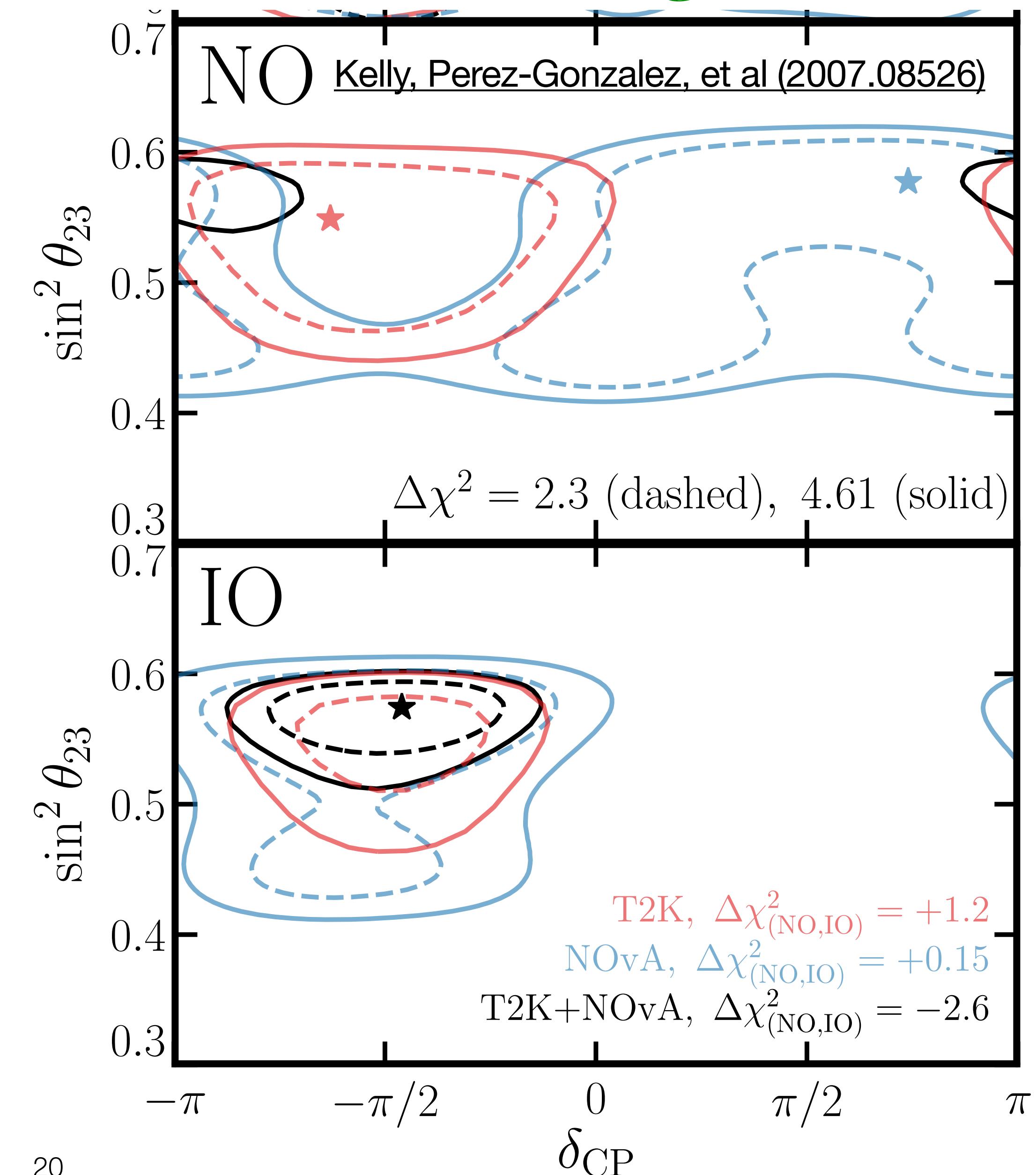
NOvA and T2K are **long baseline, accelerator oscillation** experiment measure δ_{CP} and the mass ordering.



Both experiments **favour normal ordering** with T2K favouring maximal CP violation and NOvA showing preference for near CP conservation

\mathcal{CP} -violating phase & mass ordering

- Recent combination of T2K and NOvA data \Rightarrow inverted ordering preferred
- Preference NO comes from degeneracies between MO, octant, and δ_{CP}
- Not first disagreement between T2K and NOvA (octants differed at more than 3σ circa 2016).



Ghost of Christmas Present: Current Global Picture

Atmospheric

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

Reactor

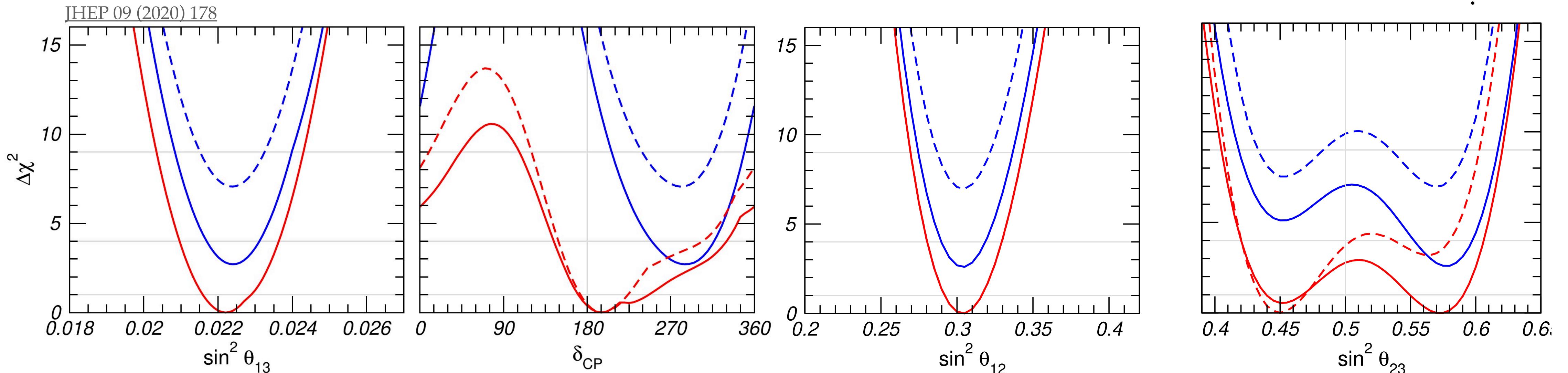
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Solar

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$$\begin{pmatrix} e^{in_1} & 0 & 0 \\ 0 & e^{in_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

?



$$\theta_{13,\text{bf}} = 8.57^\circ$$

$$\delta_{CP,\text{bf}} = 194^\circ$$

$$\theta_{12,\text{bf}} = 33.44^\circ$$

$$\theta_{23,\text{bf}} = 49.2^\circ$$

η_1, η_2 Majorana phases which are observable if neutrinos are their own antiparticle. Possible given neutrinos are electrically neutral.

Nature of Neutrinos

- Charged particles can be distinguished from their antiparticle.
- Neutrinos electrically neutral \implies Majorana ($\nu = C\bar{\nu}^T$) or Dirac.

Dirac Neutrinos

- Add right-handed neutrinos
- Assume that lepton number conserved
- Higgs mechanism provides mass to neutrinos

$$\mathcal{L}_\nu \supset Y_{\alpha i} L^\alpha H N^i \quad Y \sim 10^{-12}$$

Justify small Yukawas....

Nature of Neutrinos

Yukawa coupling forbidden at tree-level as RHN is charged under a new (gauge/global) symmetry spontaneously broken

$$\lambda_{\alpha i} L^\alpha H N^i \rightarrow \frac{\kappa_{\alpha i}}{\Lambda} (L^\alpha H) (N^i \Phi)$$

$$\Phi \text{ spontaneously breaks the new symmetry at } v_\Phi \implies \lambda = \frac{\kappa v_\Phi}{\Lambda}$$

[De Gouvea & Hernandez \(1507.00916\)](#)

Gauged chiral symmetry for RHNS \implies no Majorana masses allowed
Heavy messenger sector including dark matter and Z'

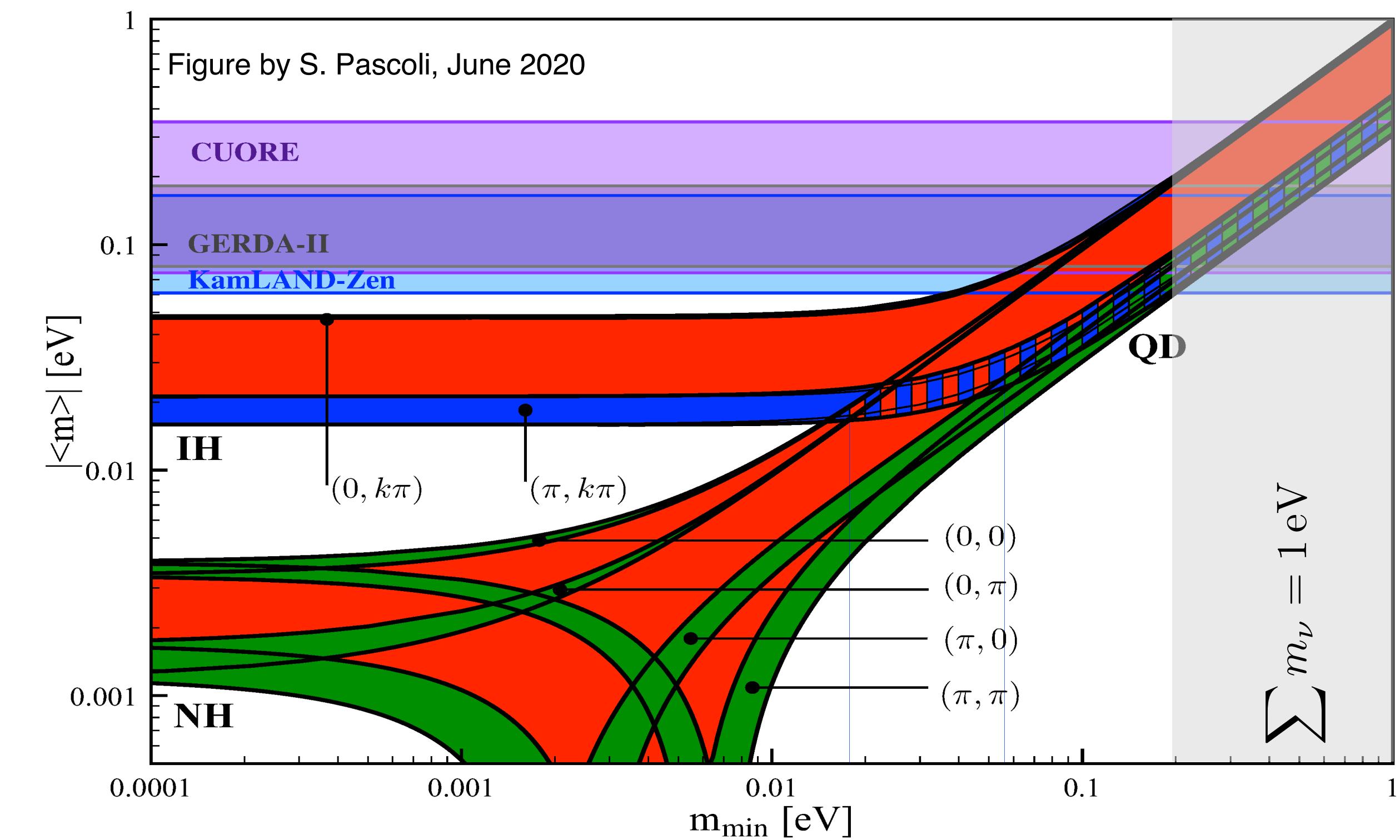
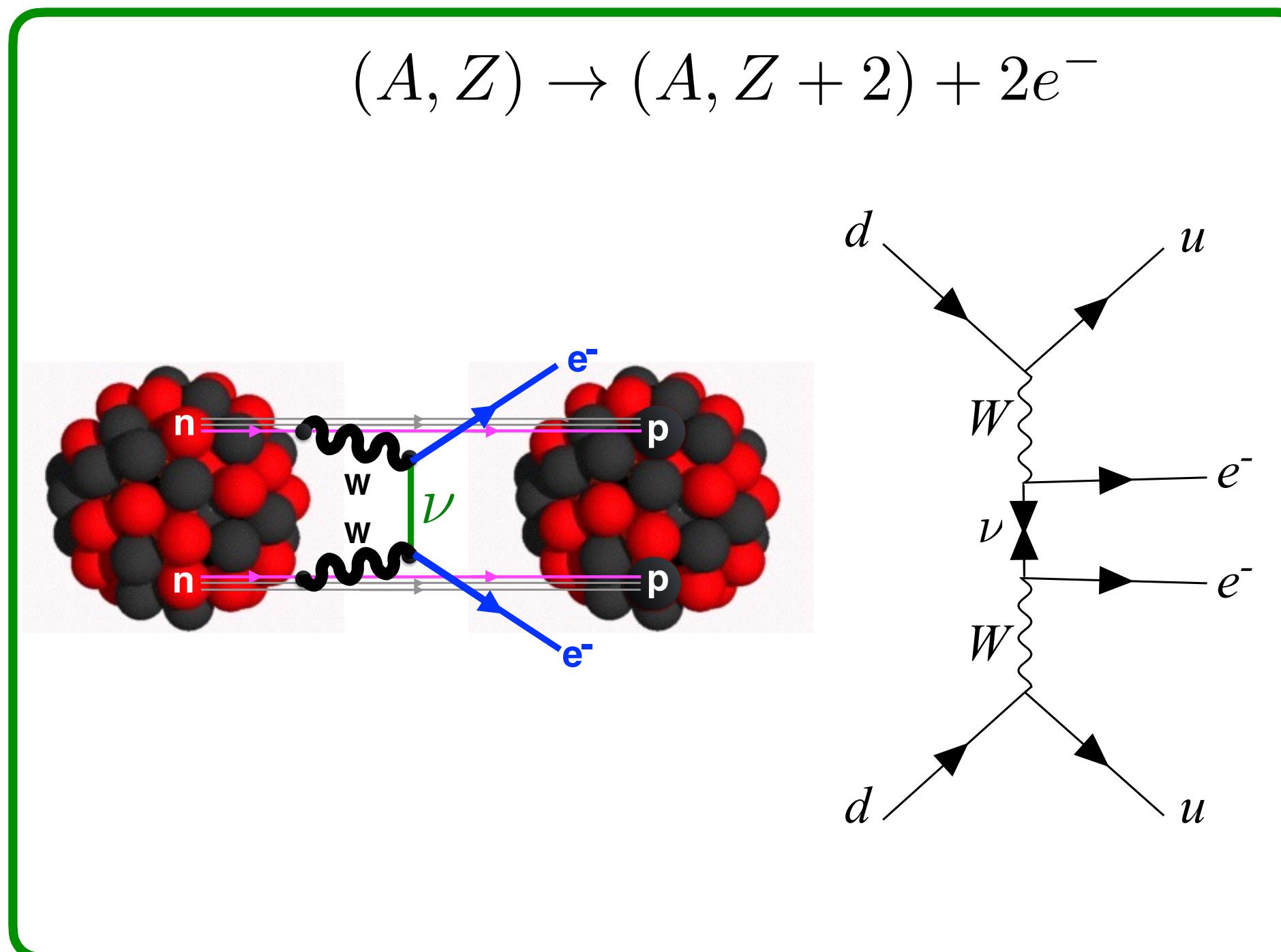
Can also add new Higgs doublet (very small vev) that only couples to neutrinos

[Davidson & Logan \(0906.3335\)](#)

[Machado, Perez-Gonzalez et al\(1507.07550\)](#)

Lepton Number Violation

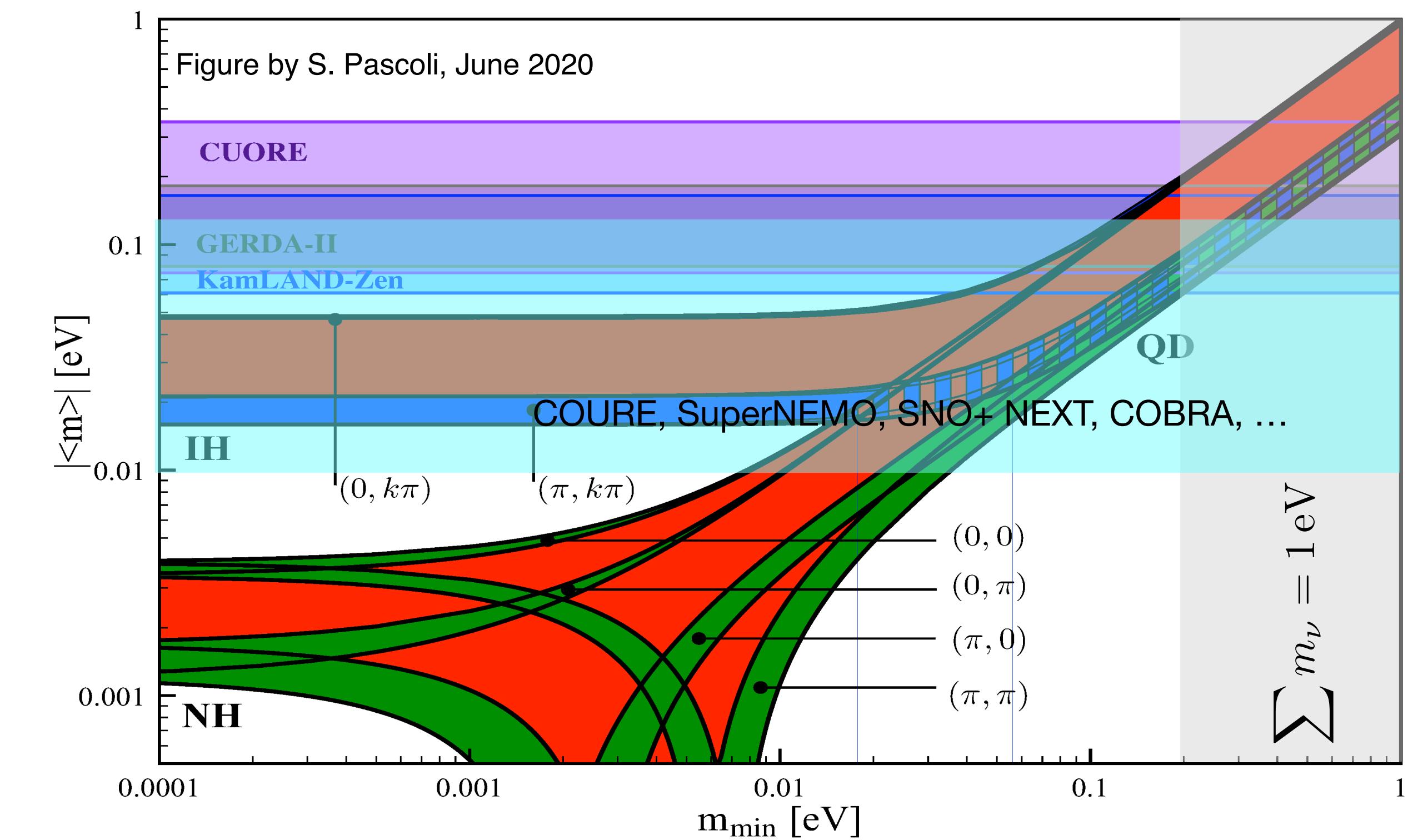
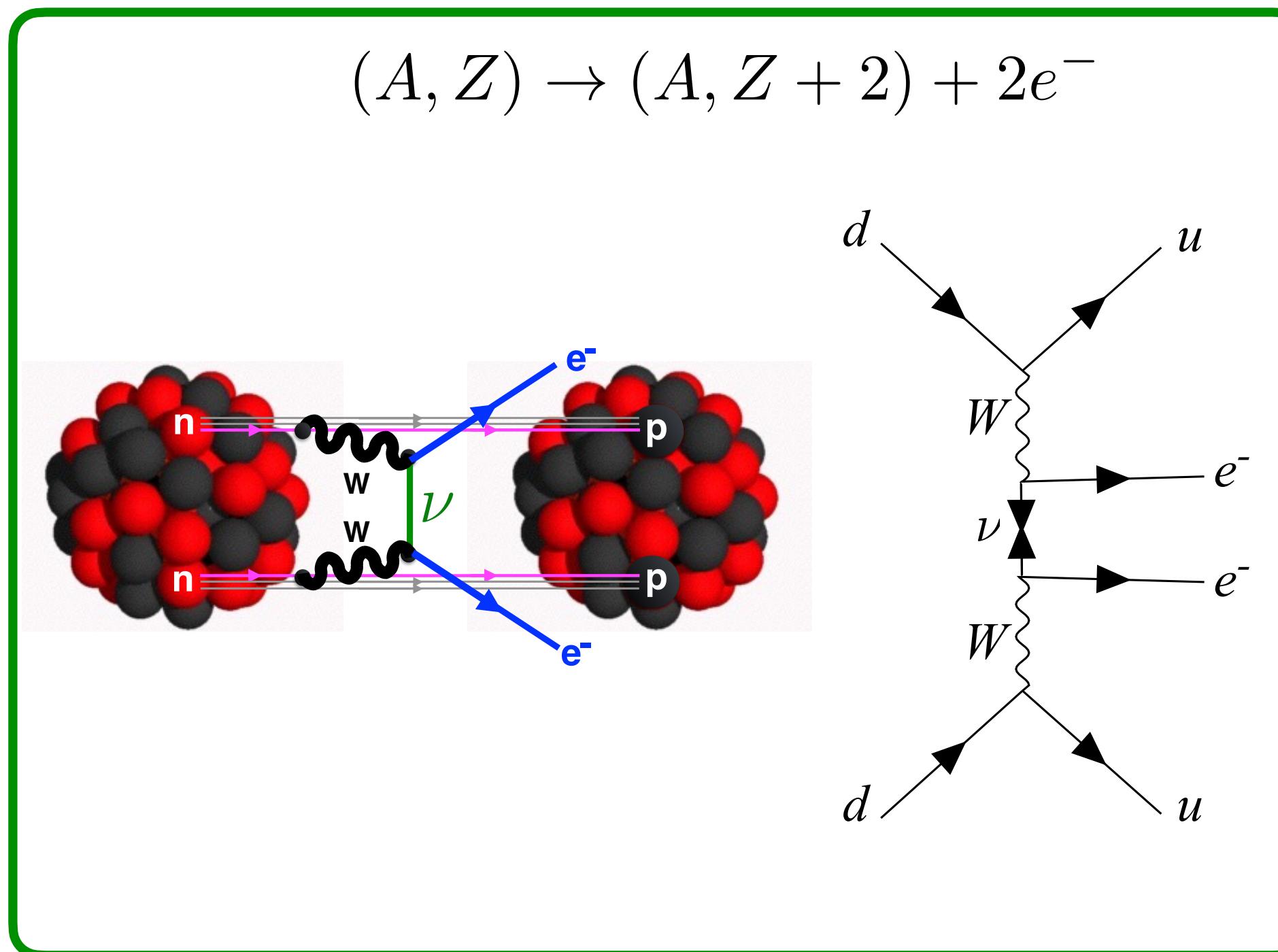
- If neutrinos are Majorana \implies lepton number not conserved.



$$| \langle m \rangle | = | m_1 \sin^2 \theta_{12} + m_2 \cos^2 \theta_{12} e^{i\alpha_{21}} + m_3 \sin^2 \theta_{13} e^{i\alpha_{31}} |$$

Lepton Number Violation

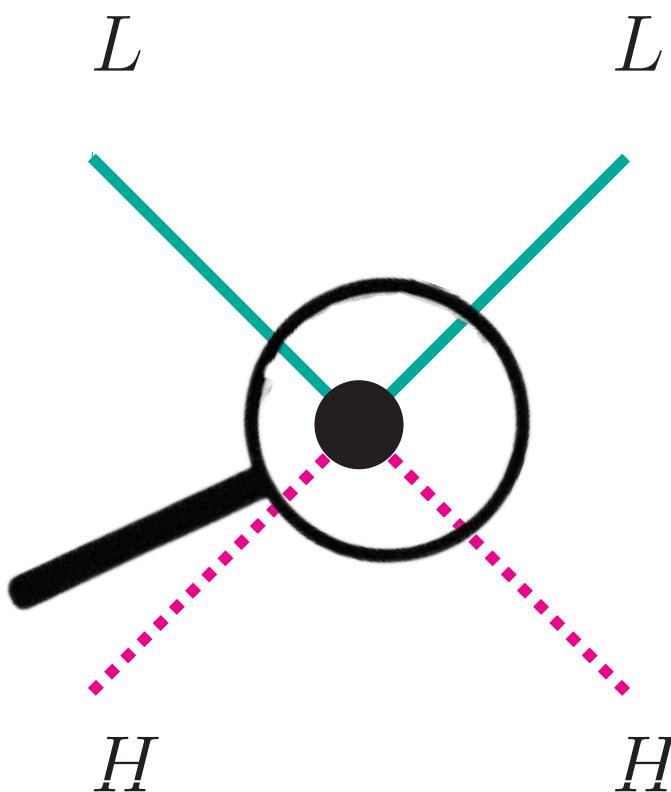
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Majorana Neutrinos

Treat SM as an effective field theory

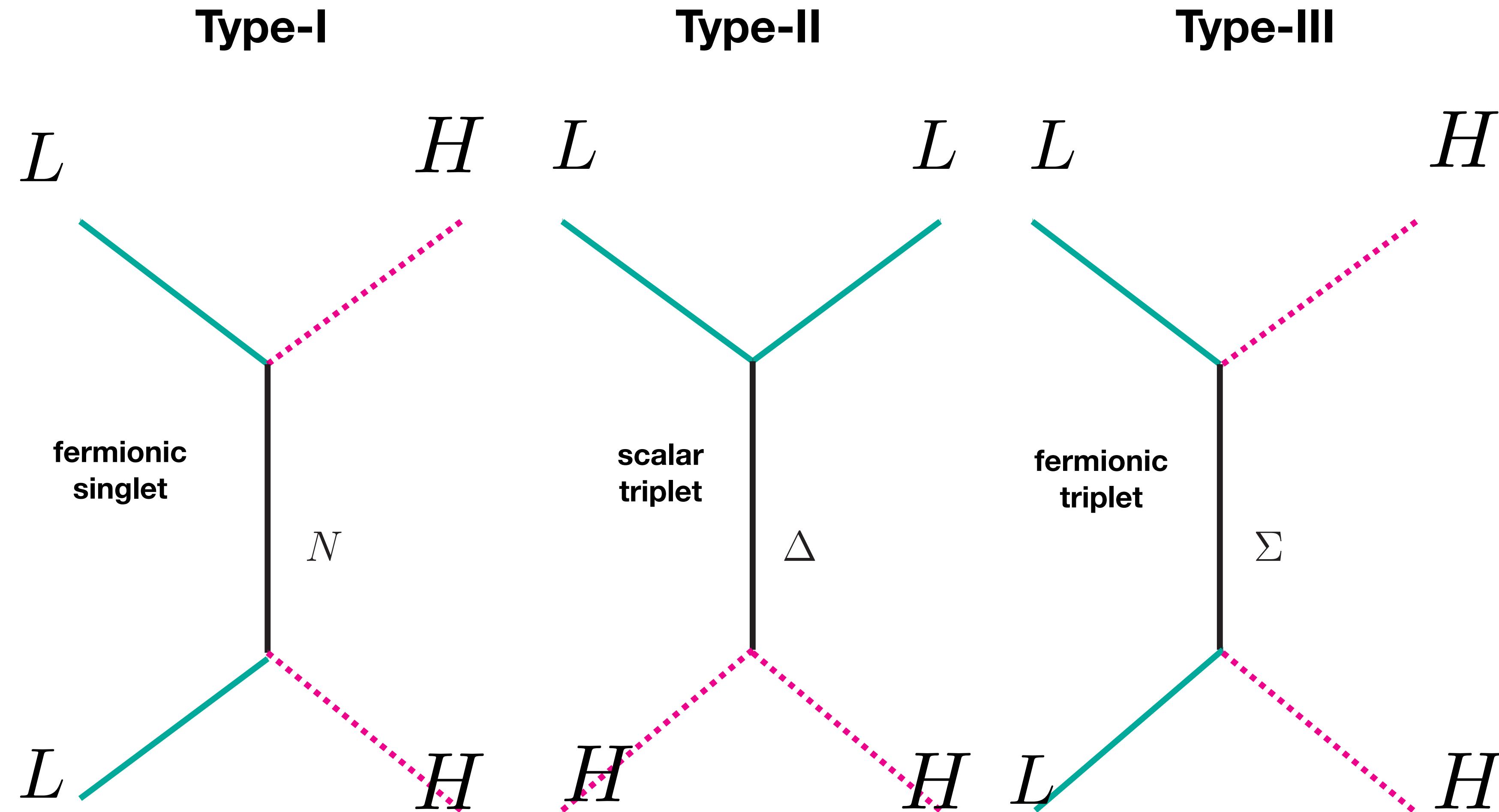


$$\mathcal{L}_\nu \supset Y_{\alpha\beta} \frac{L^\alpha H L^\beta H}{\Lambda} + \mathcal{O}\left(\frac{1}{\Lambda^2}\right) + \dots$$

Weinberg

- Dimension-5 operator violated lepton number by two units
- After EWSB: $m_{\alpha\beta} = Y_{\alpha\beta} v^2 / \Lambda$

Majorana Neutrinos



- Dimension-5 operator violated lepton number by two units
- After EWSB: $m_\nu = \frac{Y^2 v^2}{M}$

Majorana Neutrinos

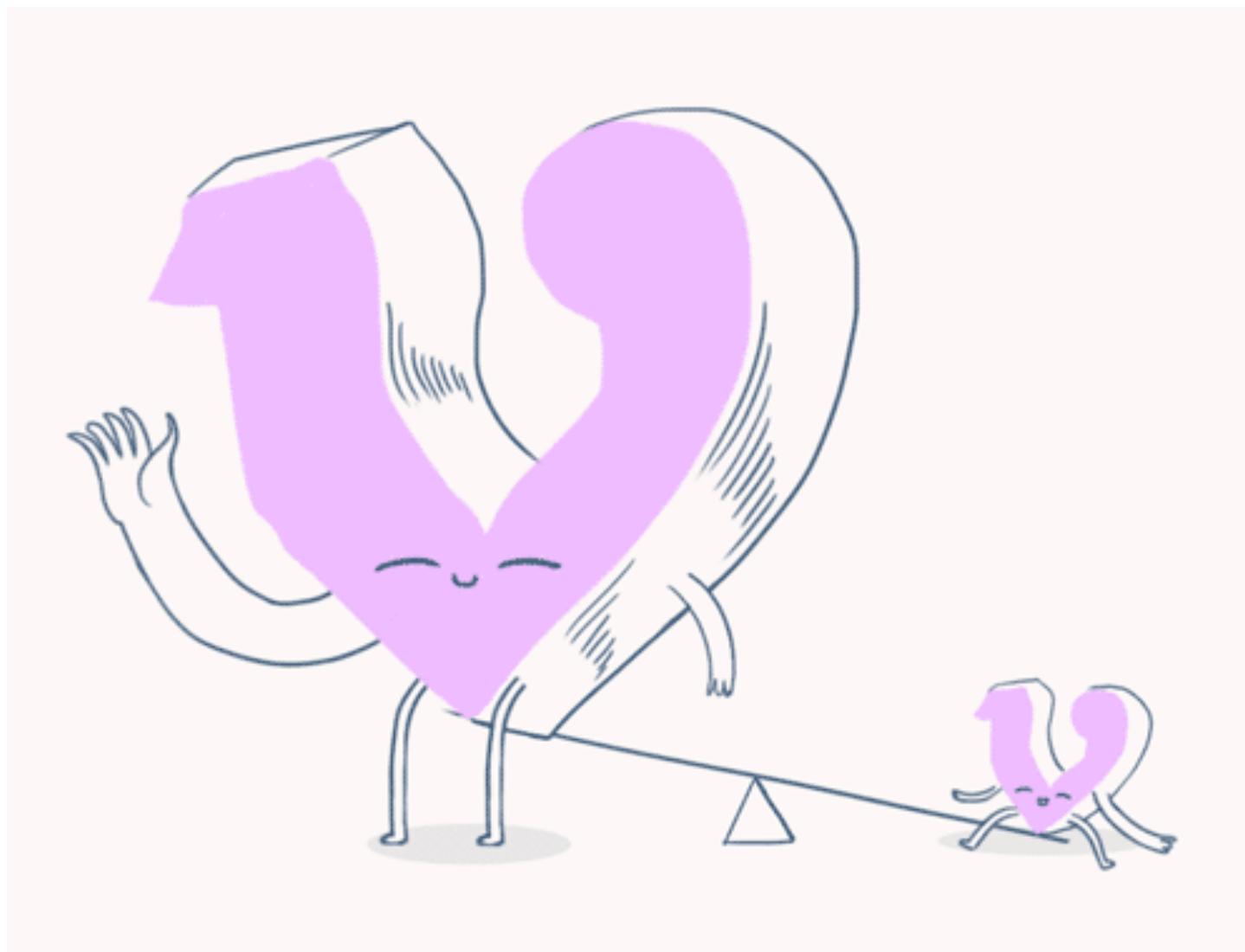
$$m_\nu = \frac{Y^2 v^2}{M}$$

$$Y \sim 1 \implies M \sim 10^{14} \text{ GeV}$$

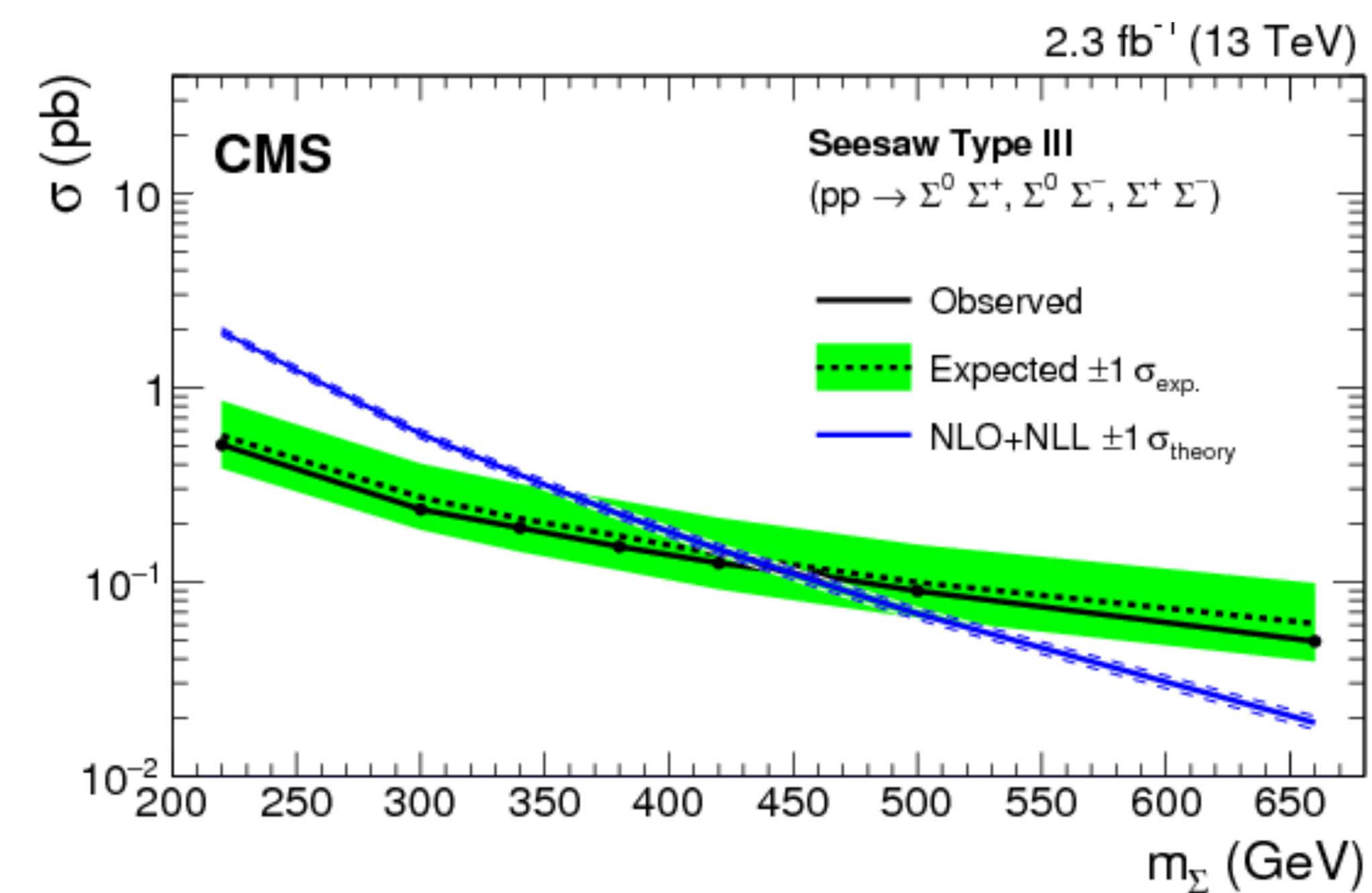
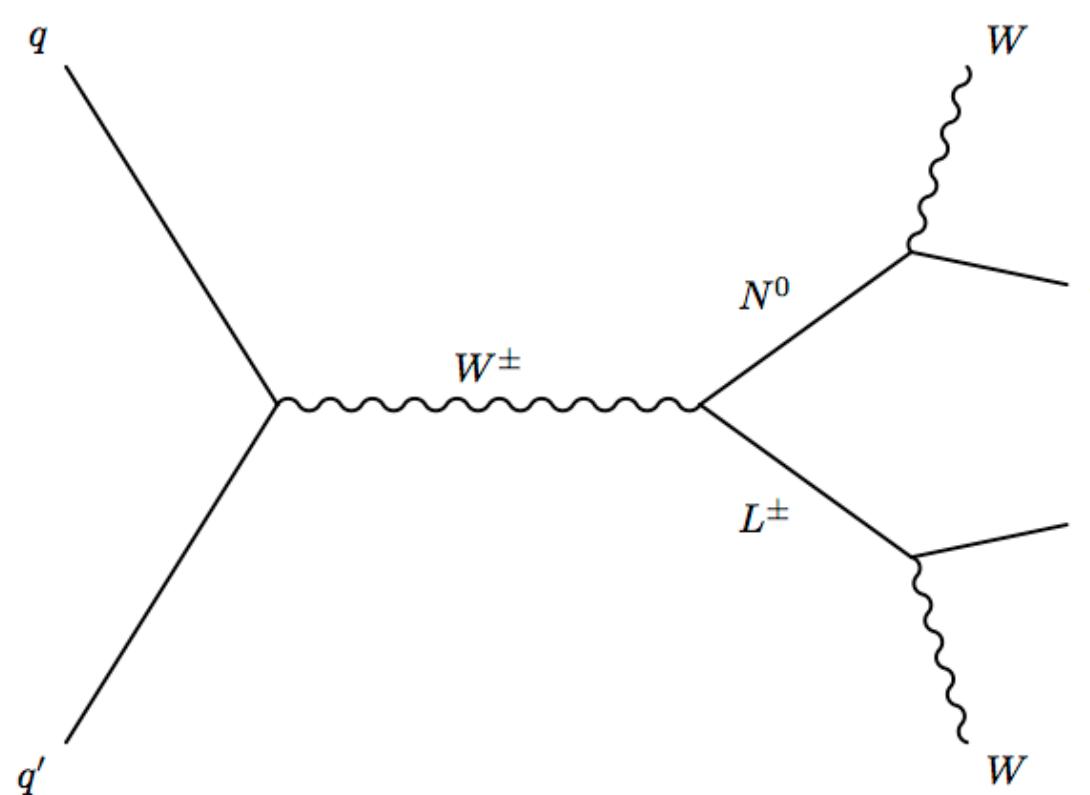
LHC range

$$Y \sim 10^{-6} \implies M \sim 100 \text{ GeV}$$

Type -II & III easier to test as because new states EW charged



Fermionic triplet - Type III



Majorana Neutrino Masses: Tree-level

Type-I seesaw

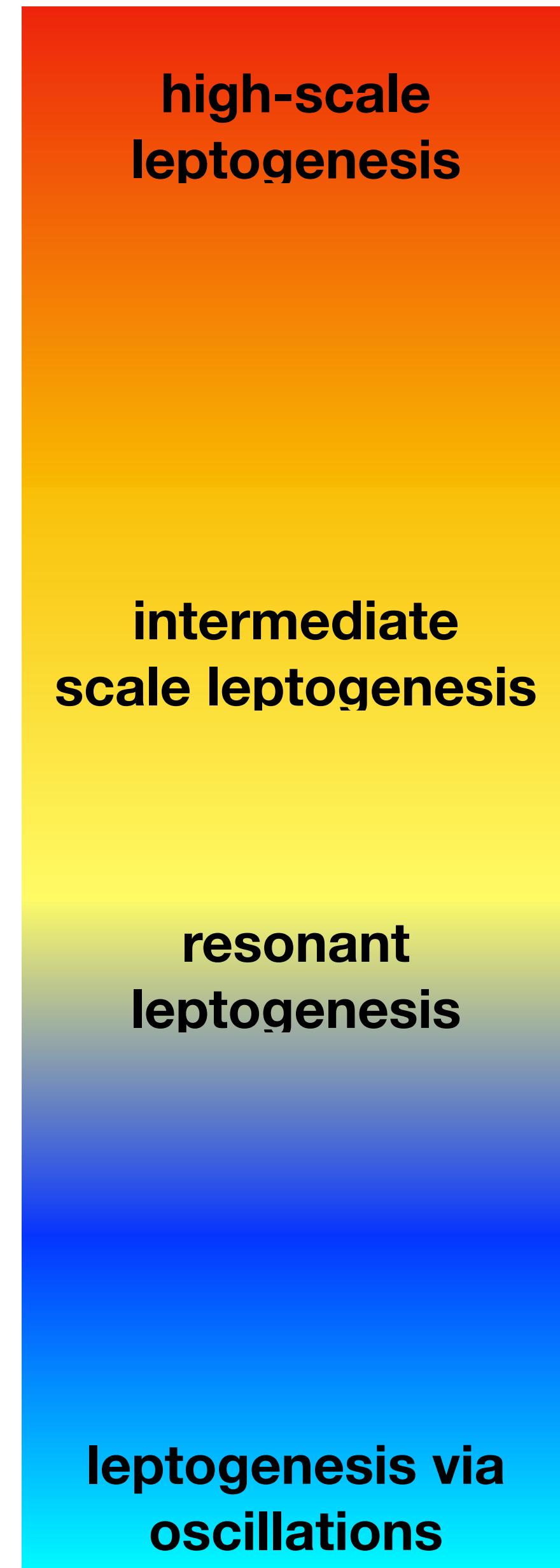
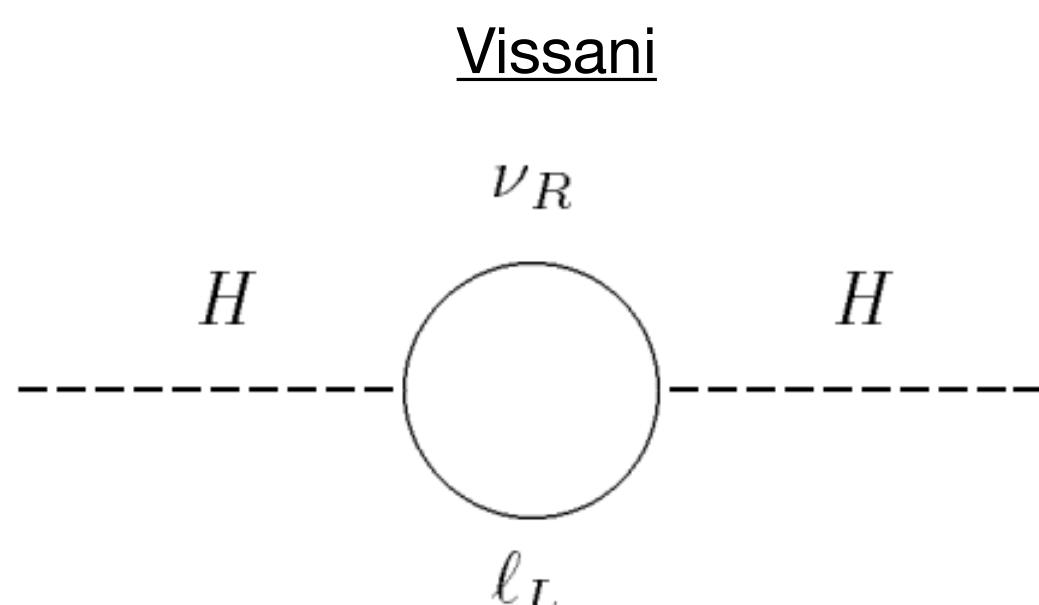
pros

- Simple
- Can explain BAU
- Predicted by GUTs
- Some circumstantial evidence
- Can radiatively generate electroweak scale

Trott & Brivio

cons

- In most regimes tough to test
- Exacerbates hierarchy problem



Mass RHN
 $\mathcal{O}(10^{12})$ GeV

Fukugida & Yanagida *Phys.Lett. B17* 45-47 (1986)

$\mathcal{O}(10^6)$ GeV

Racker, Rius & Pena *JCAP 1207* 030 (2013)

$\mathcal{O}(10^3)$ GeV

Pilaftsis & Underwood *Nucl.Phys. B692* 303 (2004)

$\mathcal{O}(1)$ GeV

Akhmedov, Rubakov & Smirnov
Phys.Rev.Lett. 81 1359-1362 (1998)

Majorana Neutrino Masses: Tree-level

Type-I seesaw

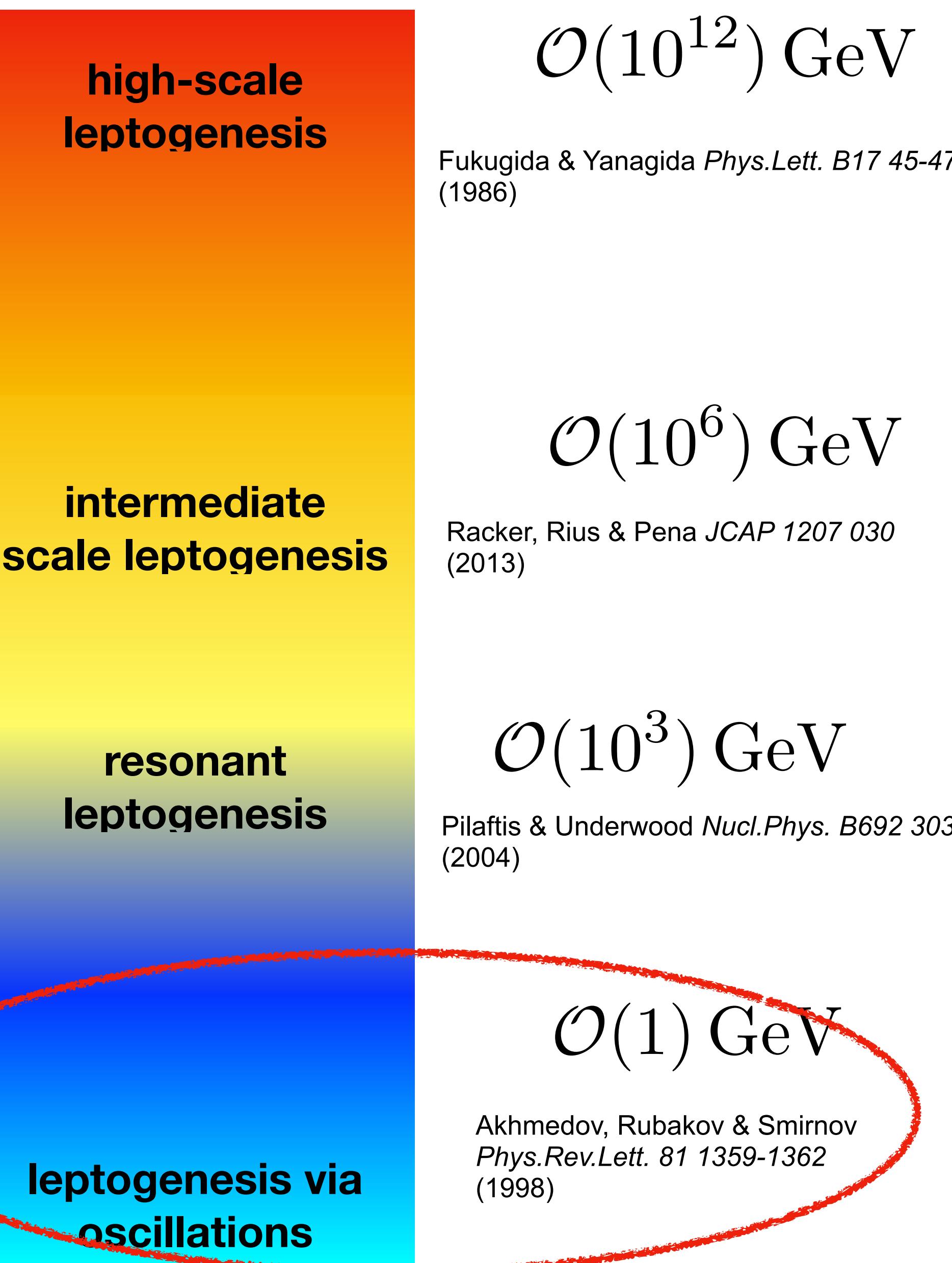
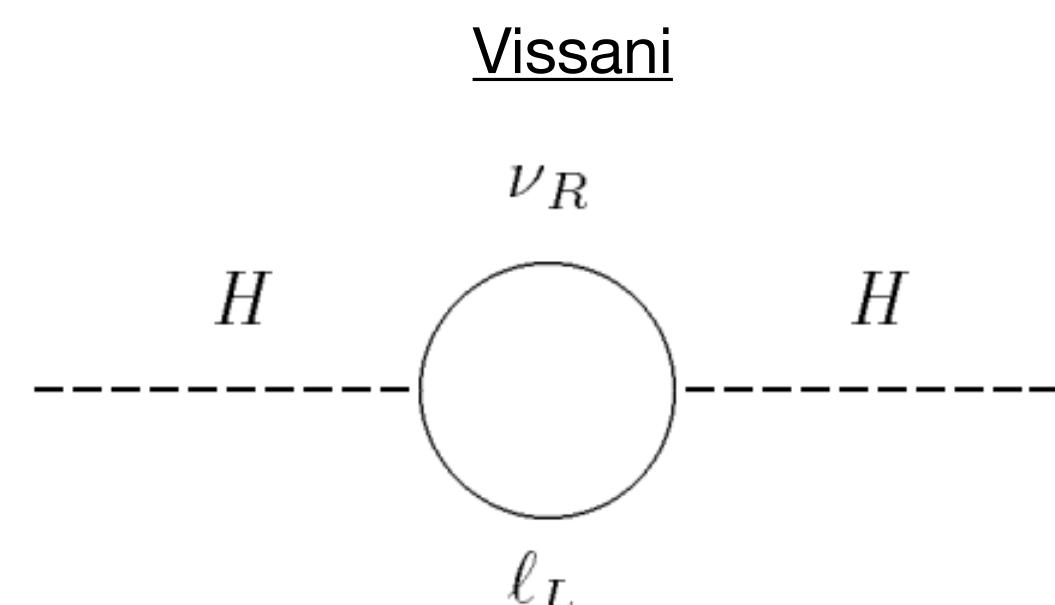
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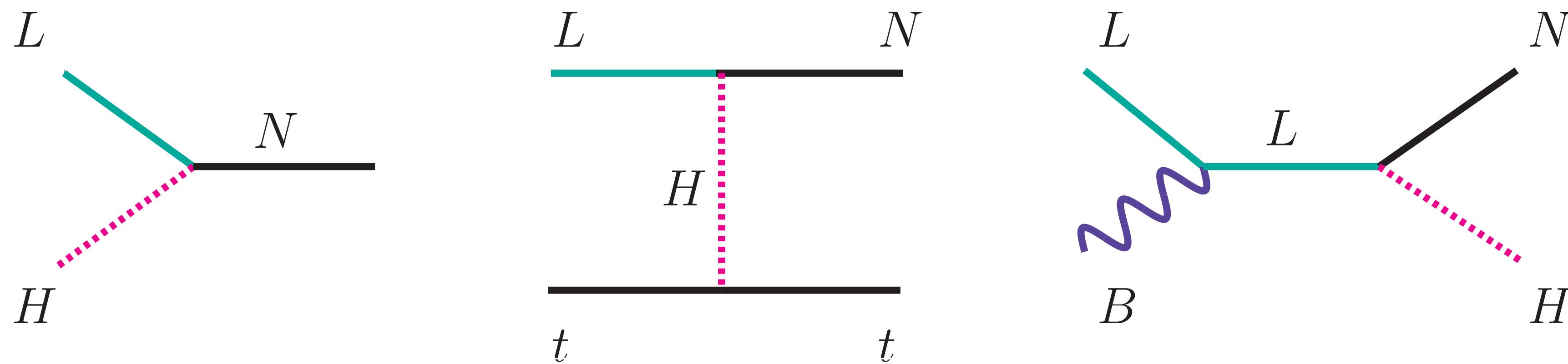
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Leptogenesis via oscillations

- highly degenerate RHNs produced via scattering at $T > T_{EW}$



Akhmedov, Rubakov & Smirnov (1998)

- small Yukawa couplings \rightarrow RHNs may not have equilibrated by the EWPT
- RHNs CP-violating oscillations \rightarrow source of lepton number and flavour asymmetry.

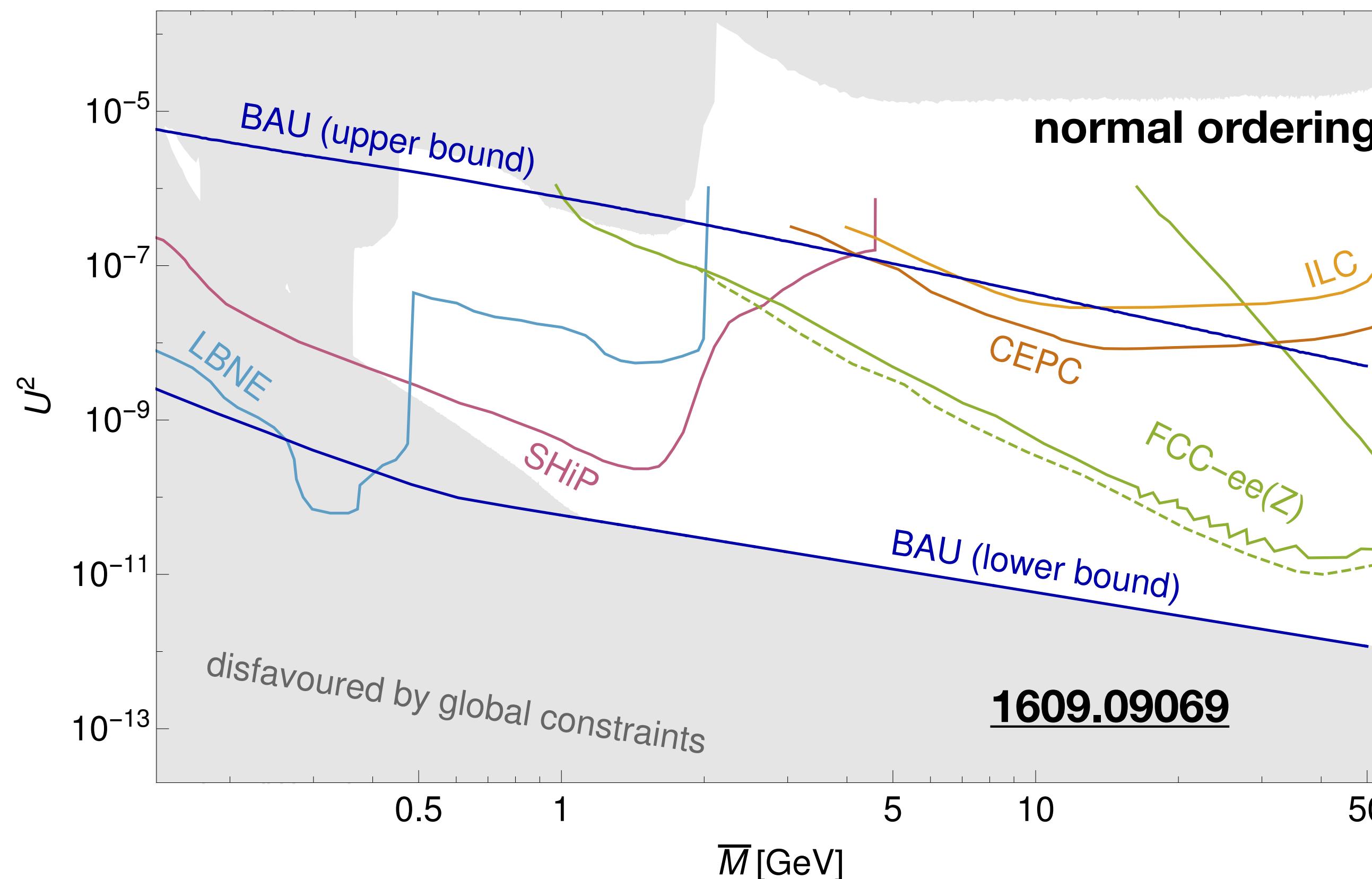
Leptogenesis via Oscillations with 2 RHNs

- GeV-scale RHNs → rich phenomenology

$$Y = \frac{1}{v} U \sqrt{m} R^T \sqrt{M} \quad \text{4 masses, 4 angles, 3 phases (2 masses + 3 angles measured)}$$

Casas & Ibarra, *Nucl.Phys. B618* (2001) 171-204

$$\nu_\alpha = U_{\alpha i} \nu_i + \Theta_{\alpha I} N_I^c \quad |U|^2 = \sum_{\alpha I} |\Theta_{\alpha I}|^2 \quad \overline{M} = \frac{M_1 + M_2}{2}$$



Majorana Neutrino Masses: Tree-level

Type-I seesaw

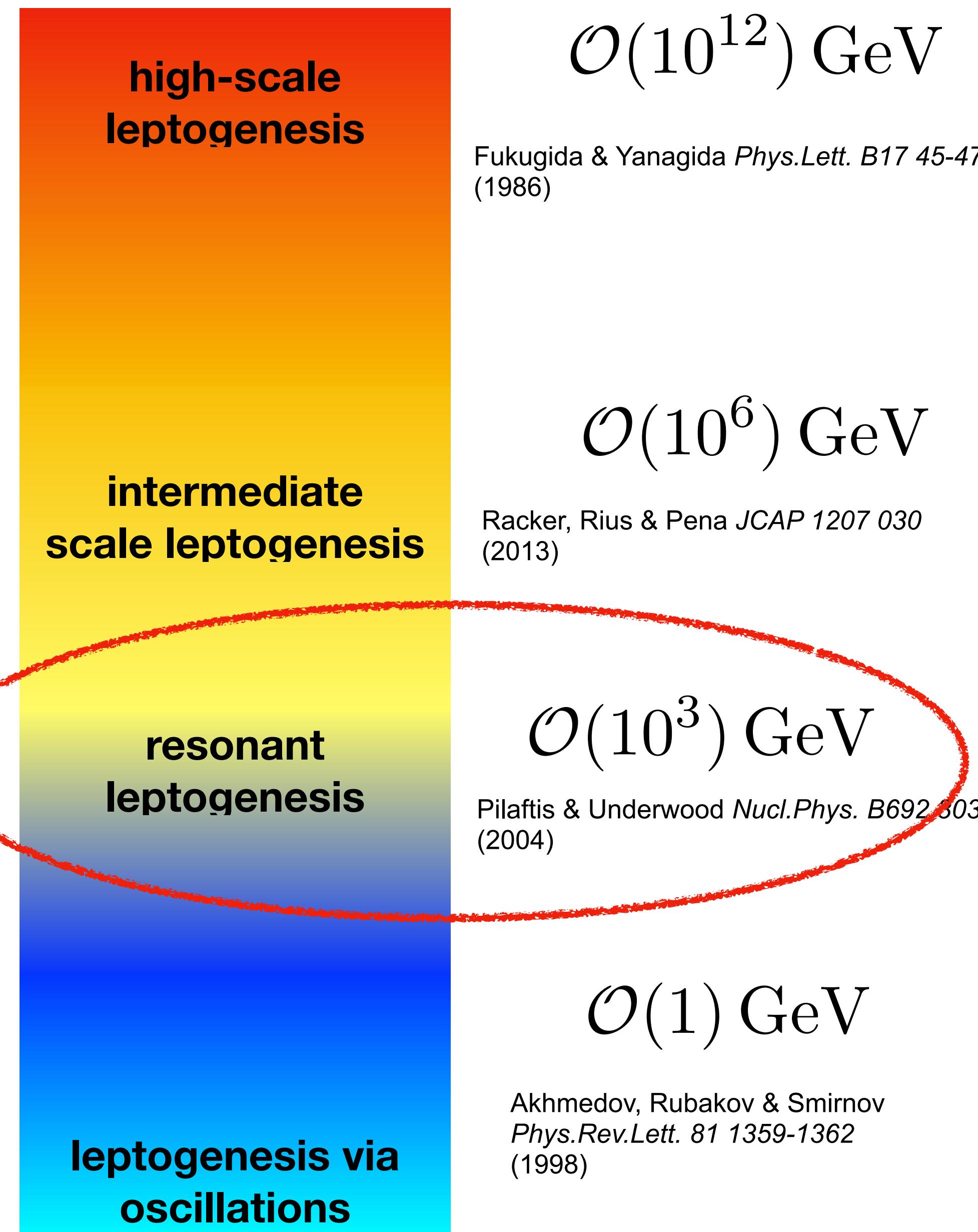
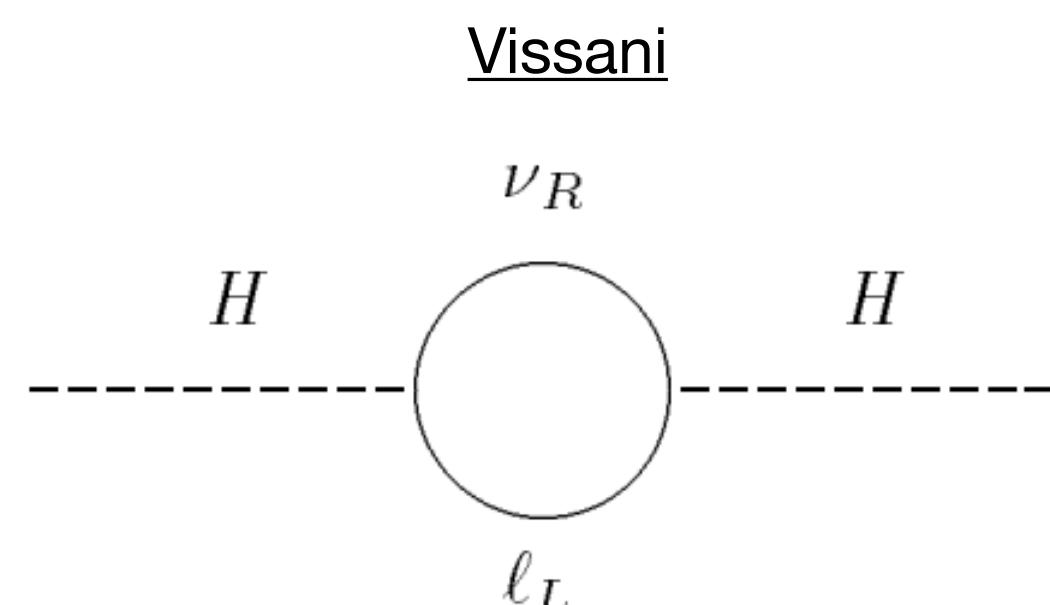
pros

- Simple
- Can explain BAU
- Predicted by GUTs
- Some circumstantial evidence
- Can radiatively generate electroweak scale

Trott & Brivio

cons

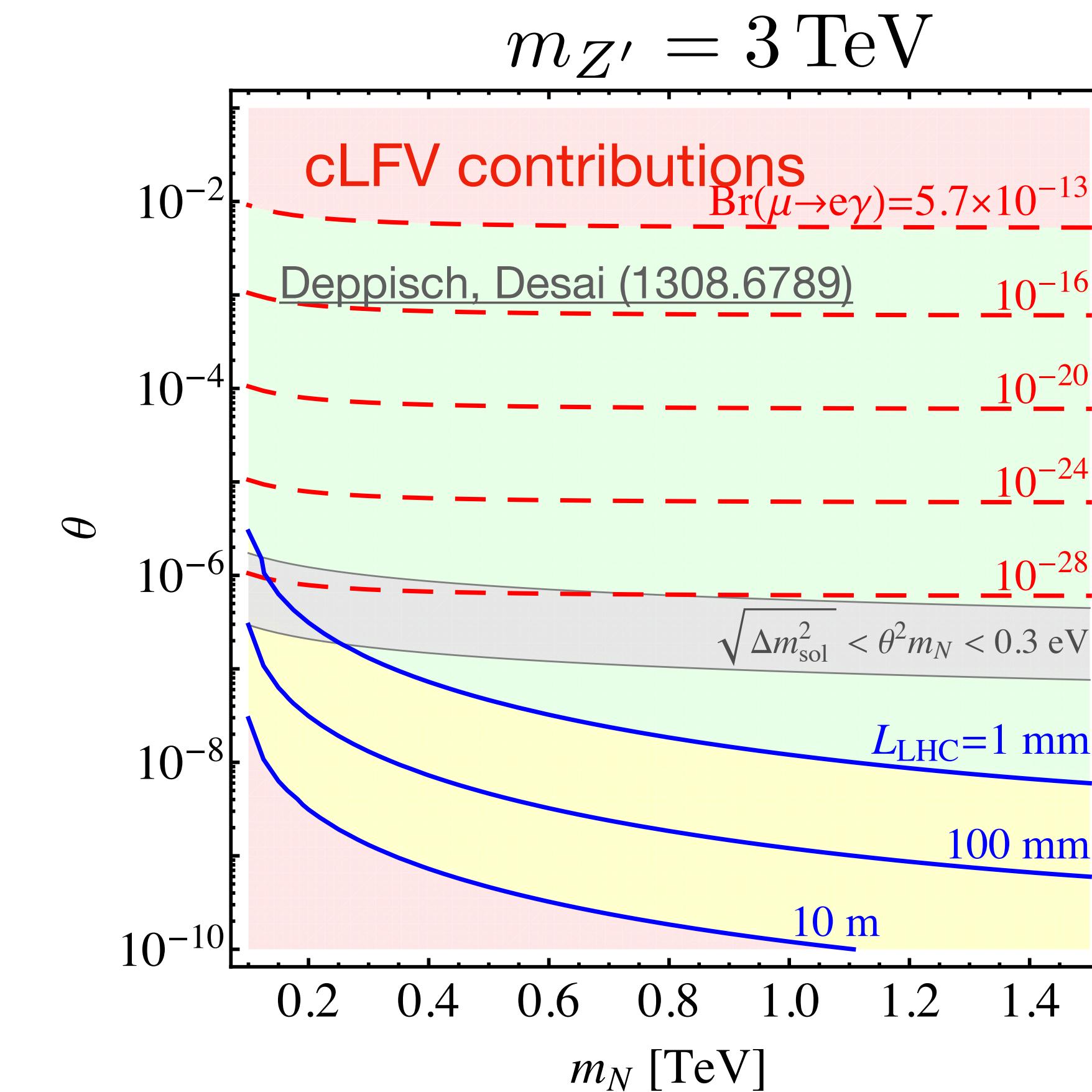
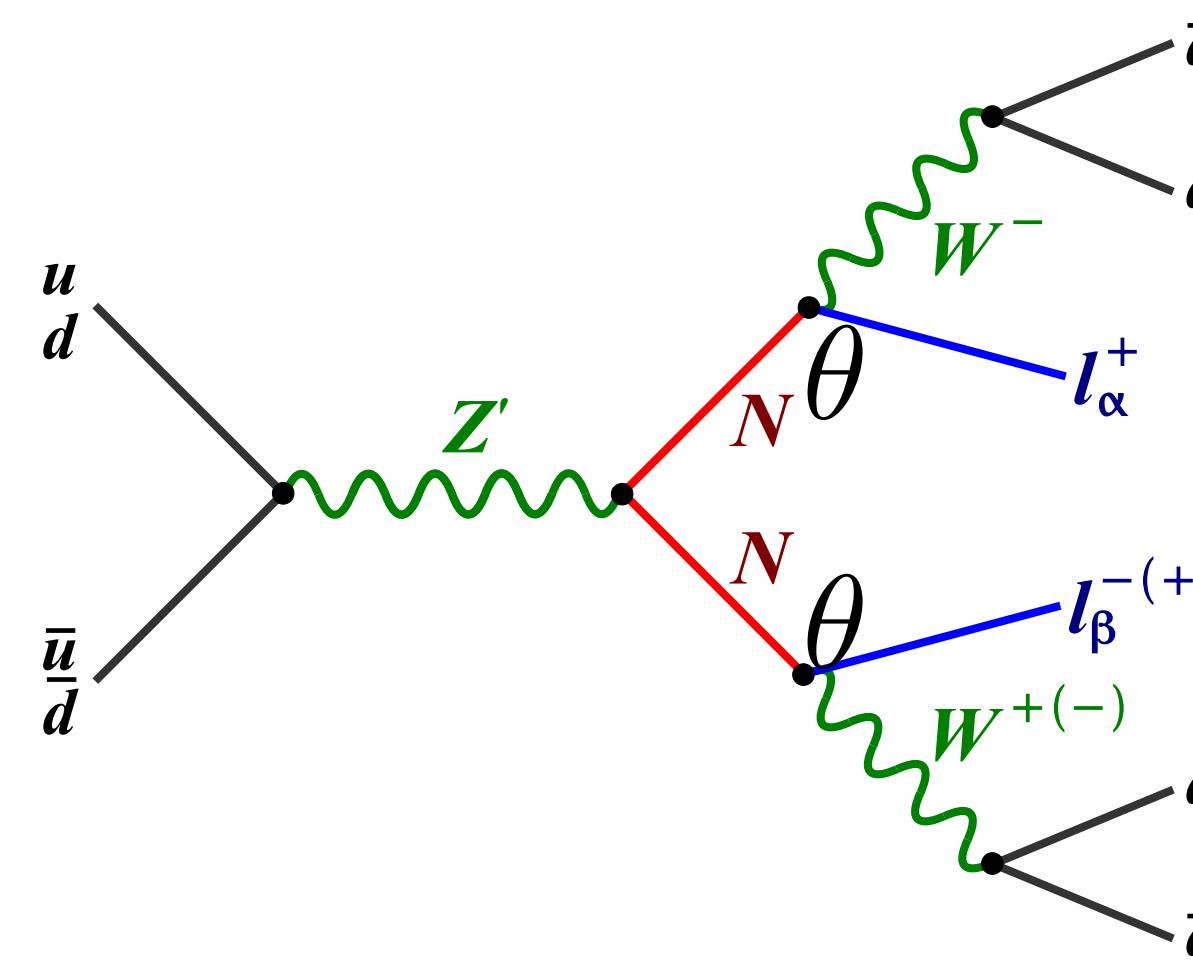
- In most regimes tough to test
- Exacerbates hierarchy problem



Majorana Neutrino Masses: Tree-level

Regime where RHNs decay width similar to their mass differences. Mass range $\sim \text{TeV}$

Harder to test. RHN masses explained by additional $U(1)_{B-L}$ symmetry and can be sufficiently long-lived \rightarrow displaced-vertex signature searched for at LHC, MATHUSLA or SHiP.



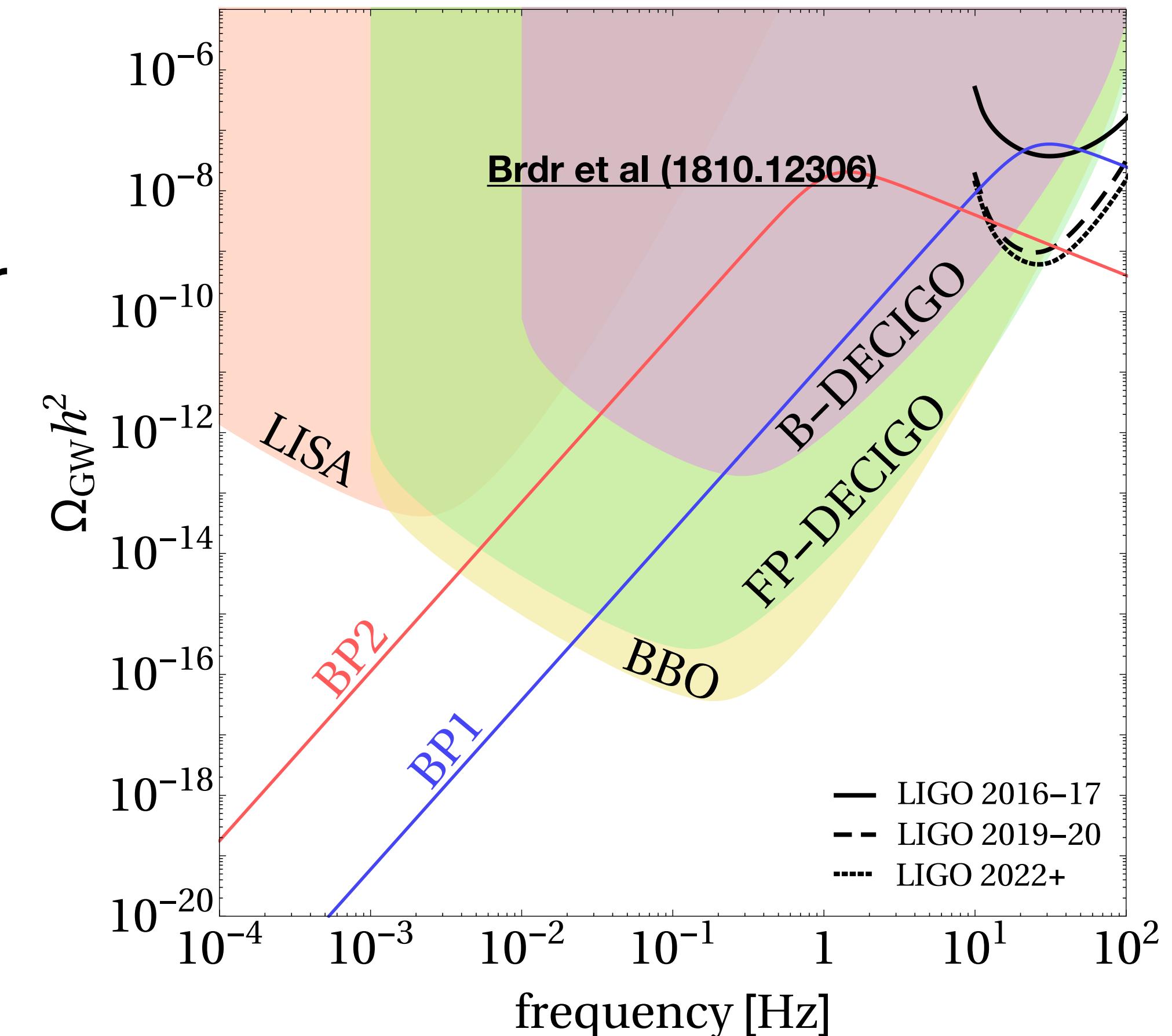
Majorana Neutrino Masses: Tree-level

Increasing interest with using gravitational waves to high scale neutrino mass models

UV-completion based on scale invariance above mass of RHNs.

New scalar breaks scale-invariance \Rightarrow
generates mass for RHNs and strong first order
phase transition

Resonant Leptogenesis +
dynamical generation of EW scale



Majorana Neutrino Masses: Radiative low scale

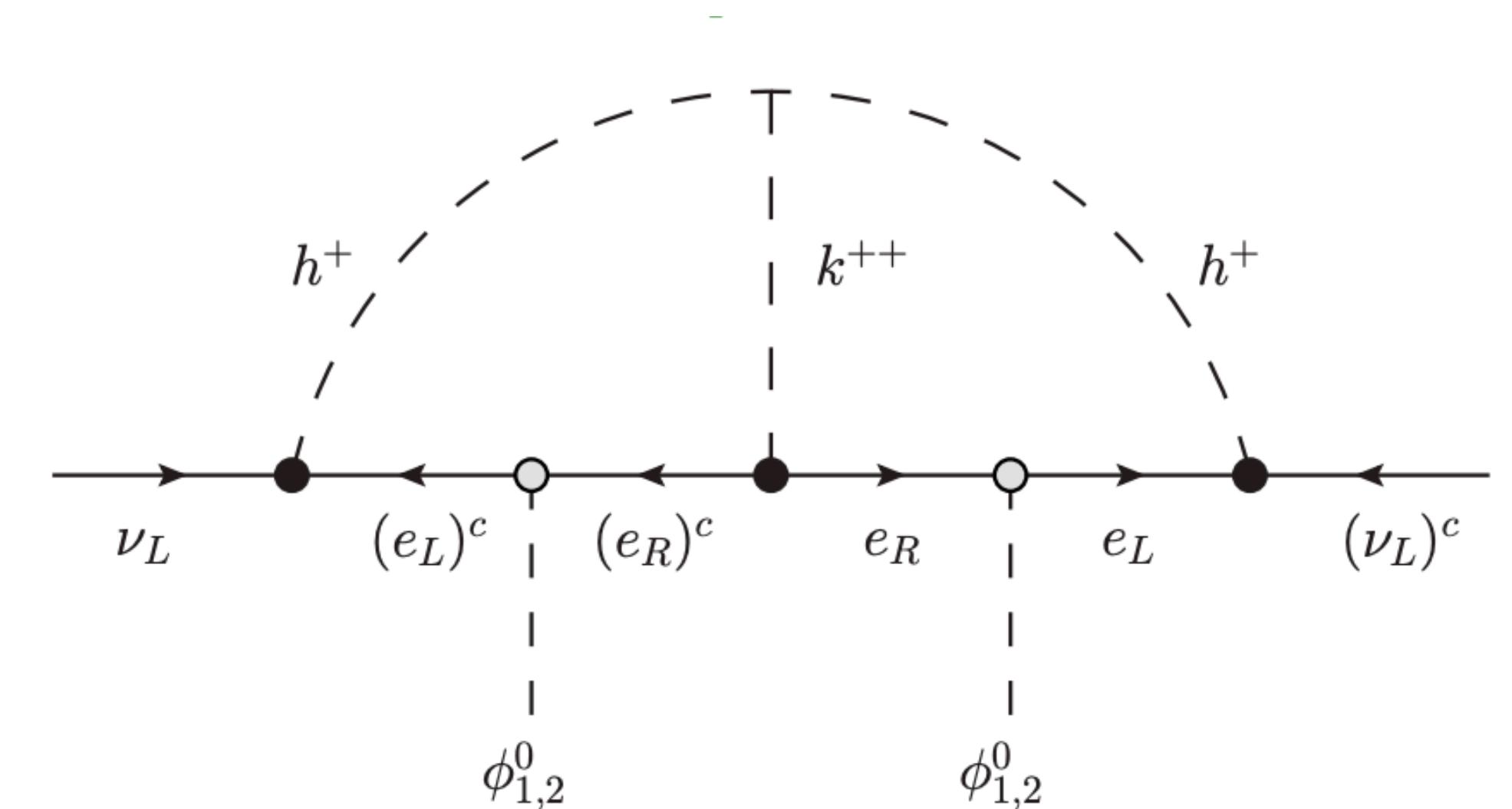
- Neutrino masses zero at tree-level, non-zero at n-loop level
- Loop suppression \implies small neutrino mass new physics is around **TeV-scale**

$$m_\nu \sim \frac{v^2}{M} \times \left(\frac{1}{16\pi^2} \right)^n$$

Zee-Babu model

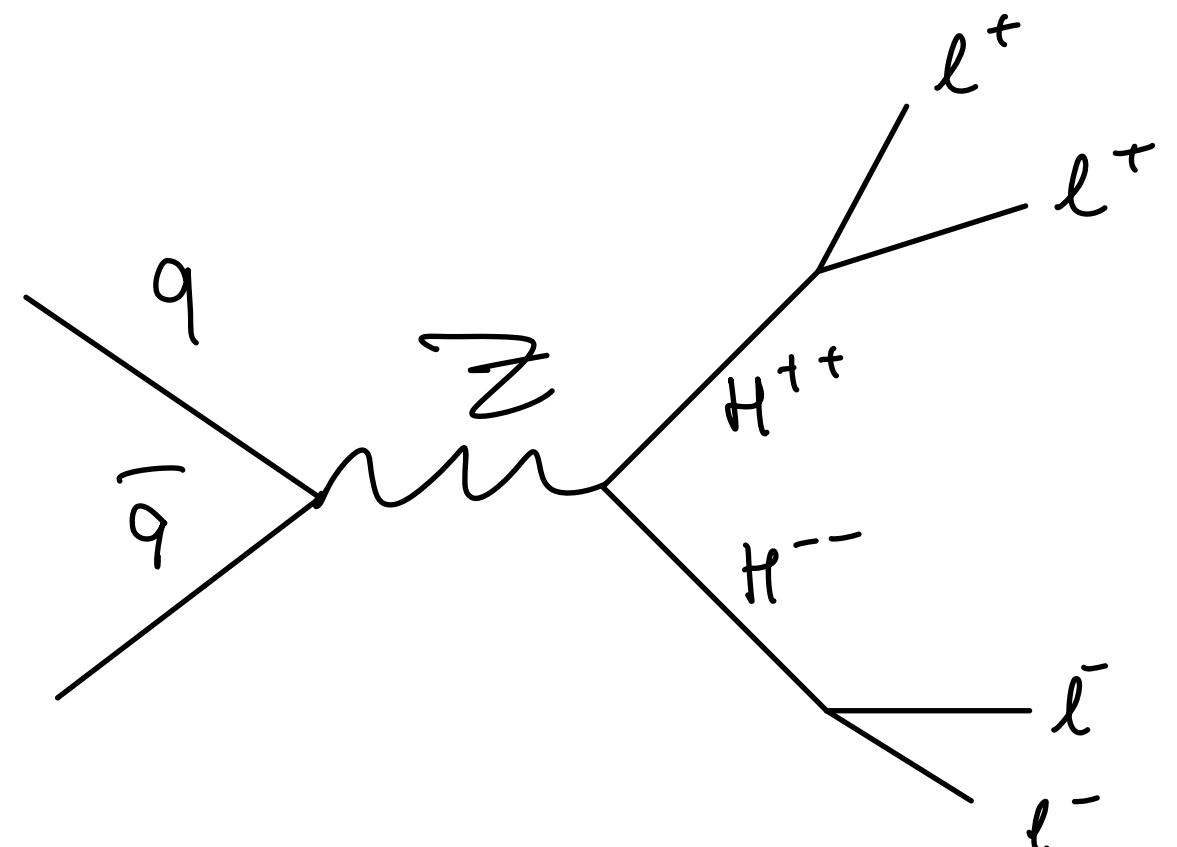
- Rich scalar sector: h^+ , k^{++}

$$\mathcal{L}_{ZB} = f \bar{\tilde{\ell}}_L \ell_L h^+ + g \bar{e}_R^c e_R k^{++} + \mu_{ZB} h^+ h^+ k^{--}$$



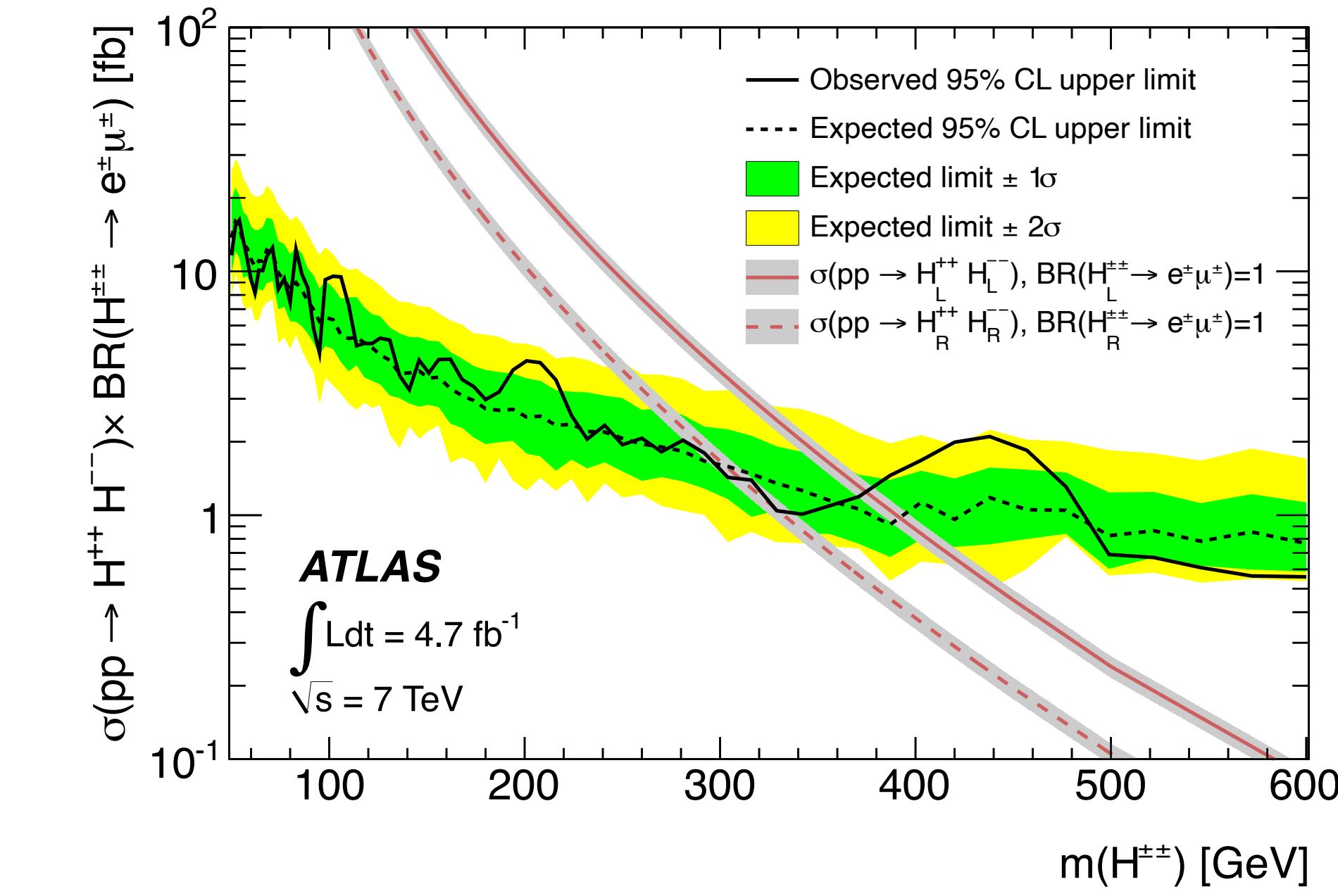
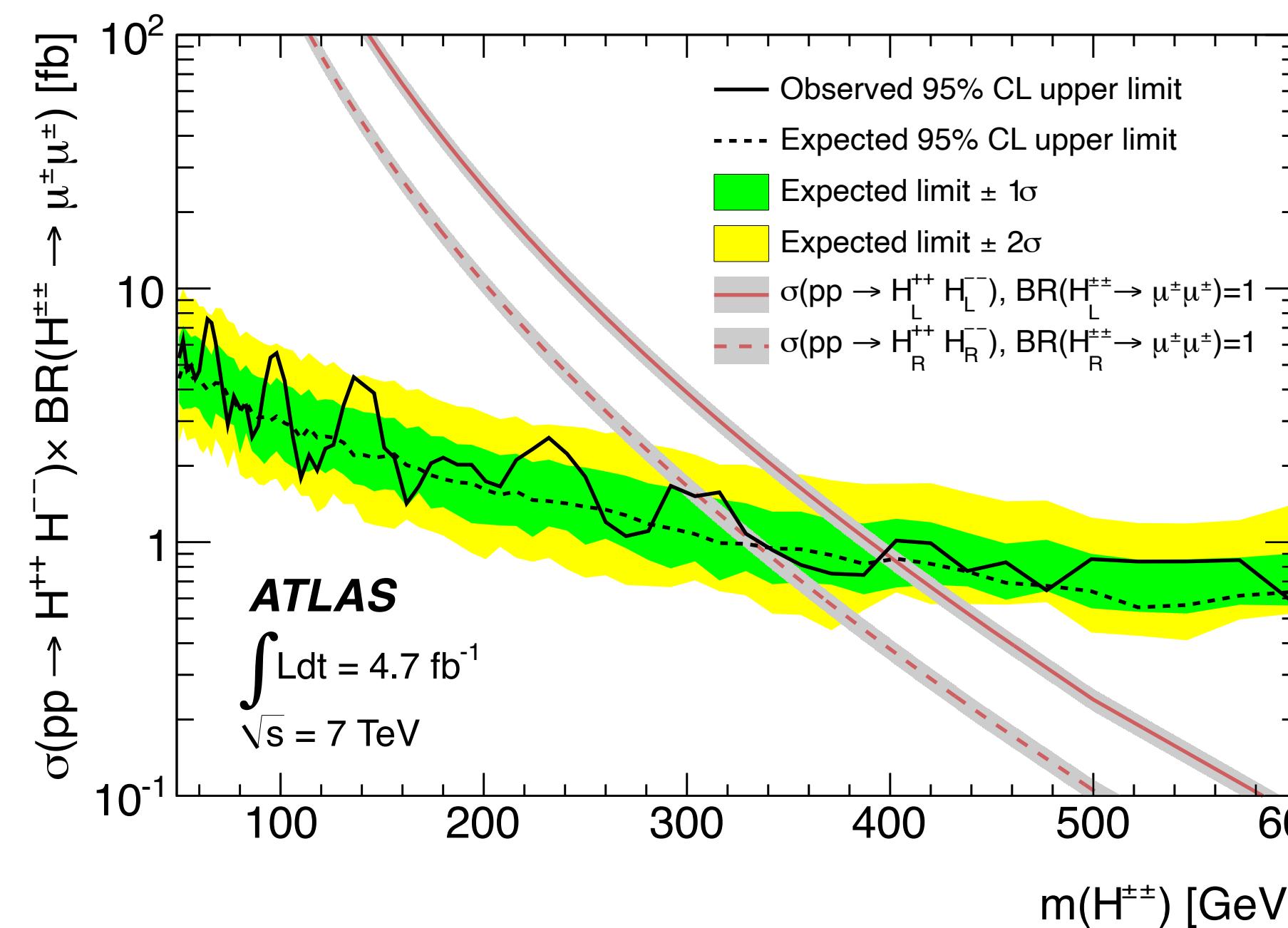
Majorana Neutrino Masses: Radiative low scale

- Neutrino masses zero at tree-level, non-zero at n-loop level
- Loop suppression \implies small neutrino mass new physics is around **TeV-scale**



Doubly charged scalars can be searched for at the LHC (like-sign dilepton searches)

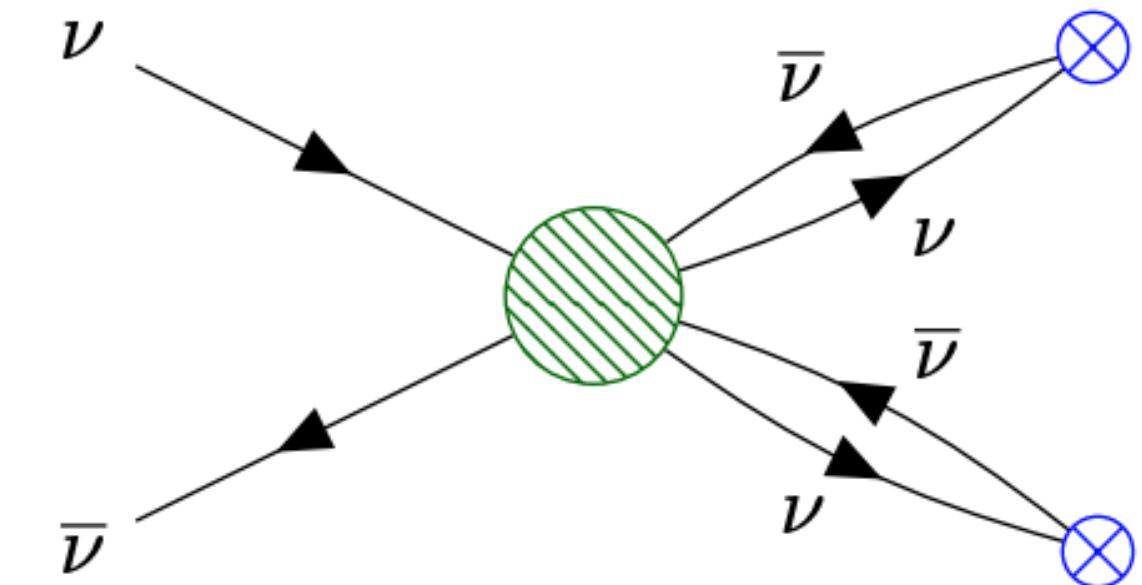
Zee-Babu model



What I have reviewed so far are popular models

Dvali & Funcke [1602.03191](#)

Assume gravity has a theta term: $\mathcal{L}_G \supset \theta_G \tilde{R} R$



Postulate neutrinos have zero bare mass and condense via NP gravitational effects in analogue with QCD:

$$\mathcal{L} \supset g_\nu v \bar{\nu} \nu$$

$$\Lambda_G \sim v \sim m_\nu \sim m_{\eta_\nu}$$

Neutrinos Dirac/ Majorana, rich phenomenology due to late time mass generation

Use Schwinger-Dyson methods to calculate condensation strength.

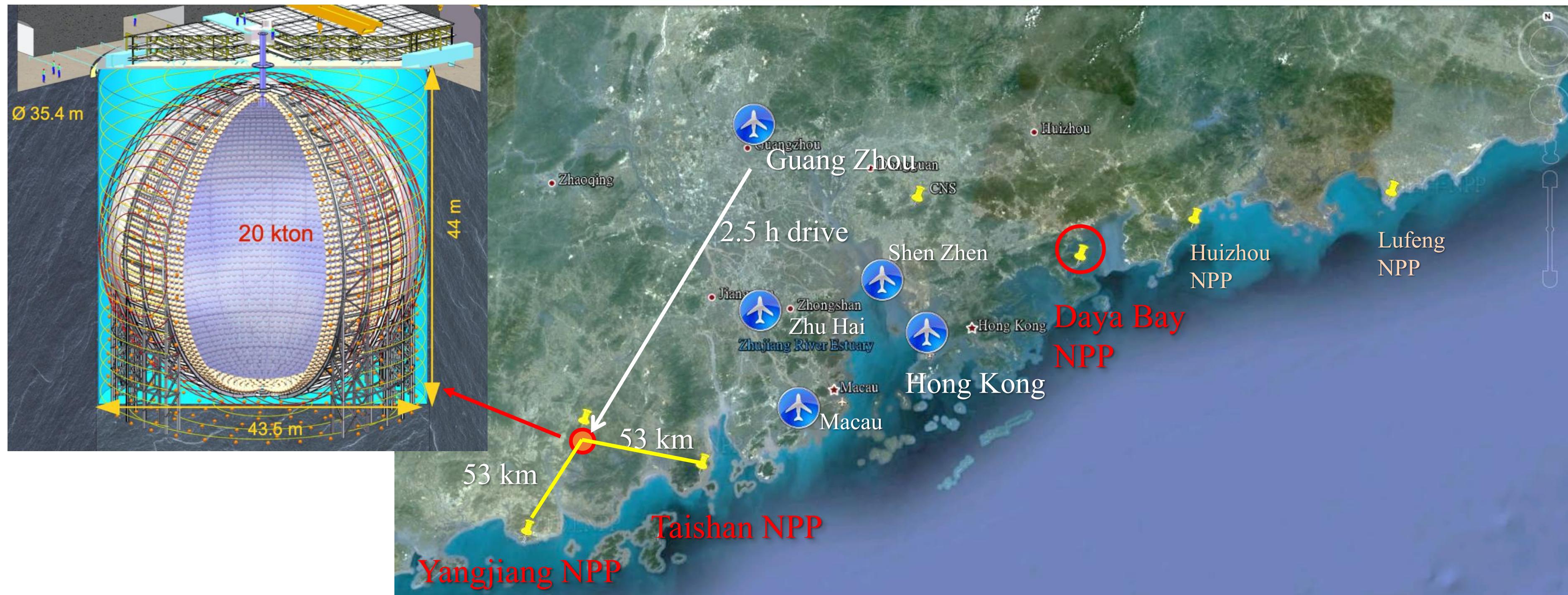
Need a lot of new particle degrees of freedom or face tuning of $\Lambda_{\text{Planck}}/m_\nu$

Barenboim, JT & Zhou [1909.04675](#)

Neutrinos: Yet to come

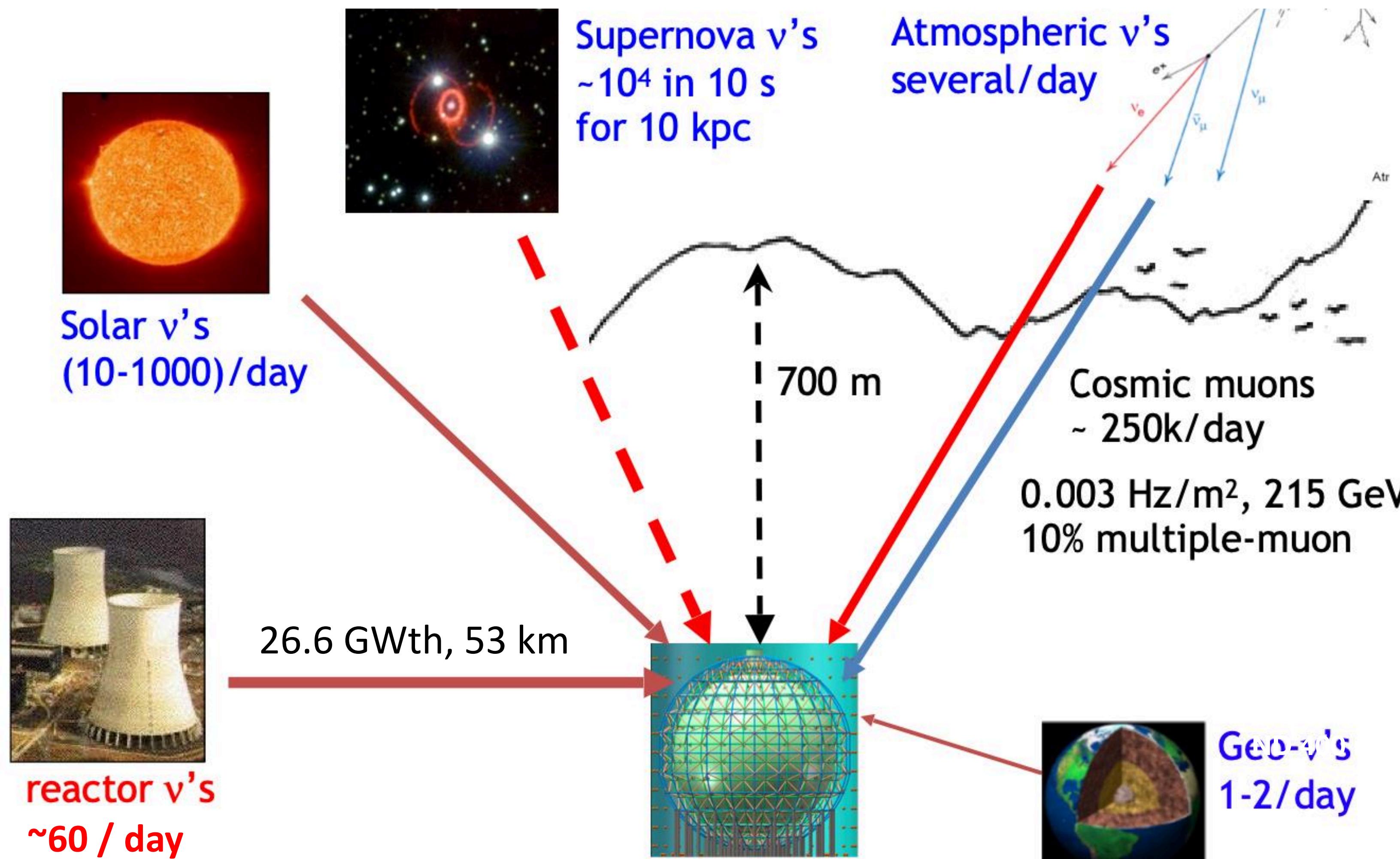
The Future of Neutrinos: Mass Ordering

- JUNO medium baseline reactor experiment will start data taking next year.



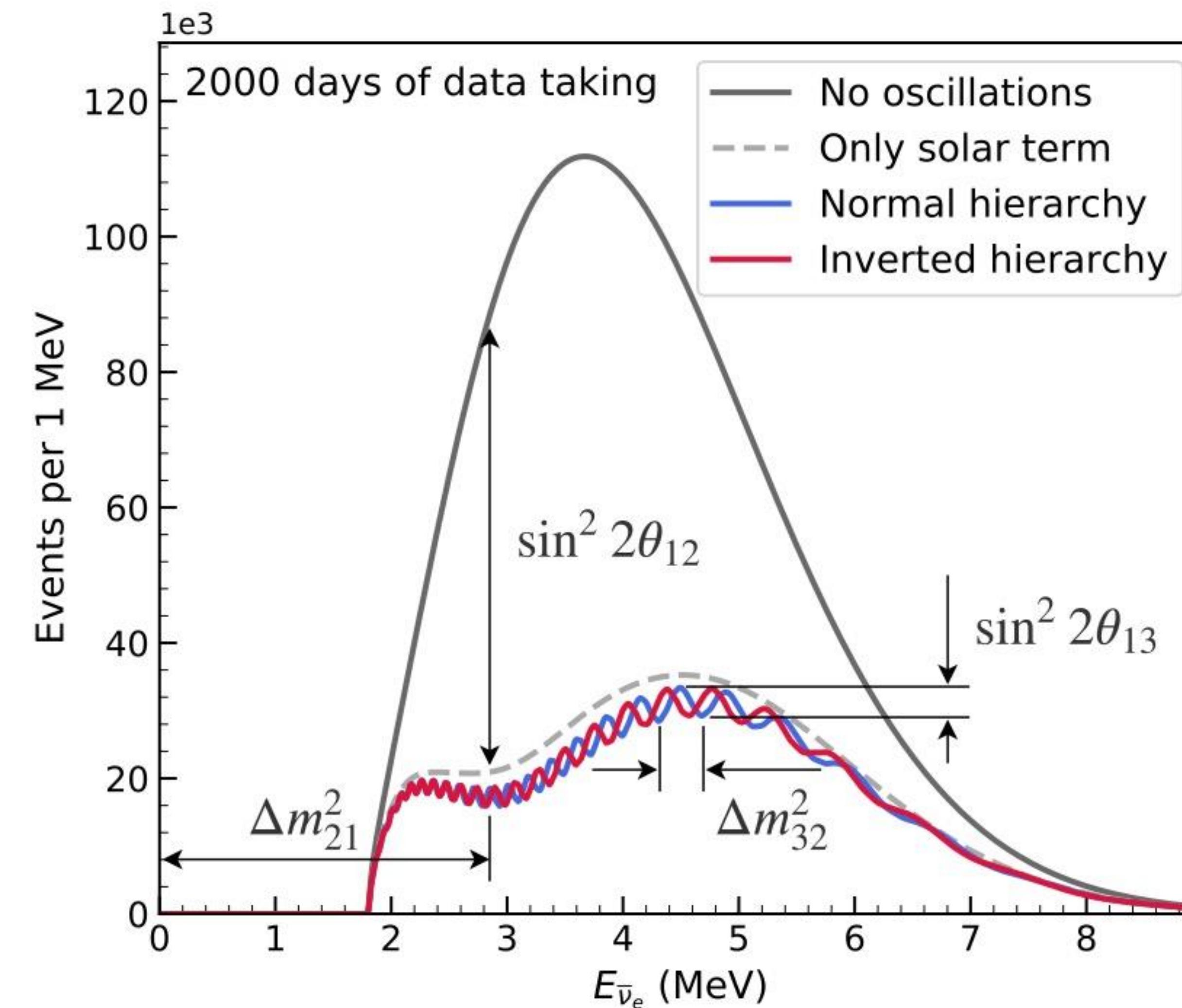
The Future of Neutrinos: Mass Ordering

- JUNO medium baseline reactor experiment will start data taking next year.



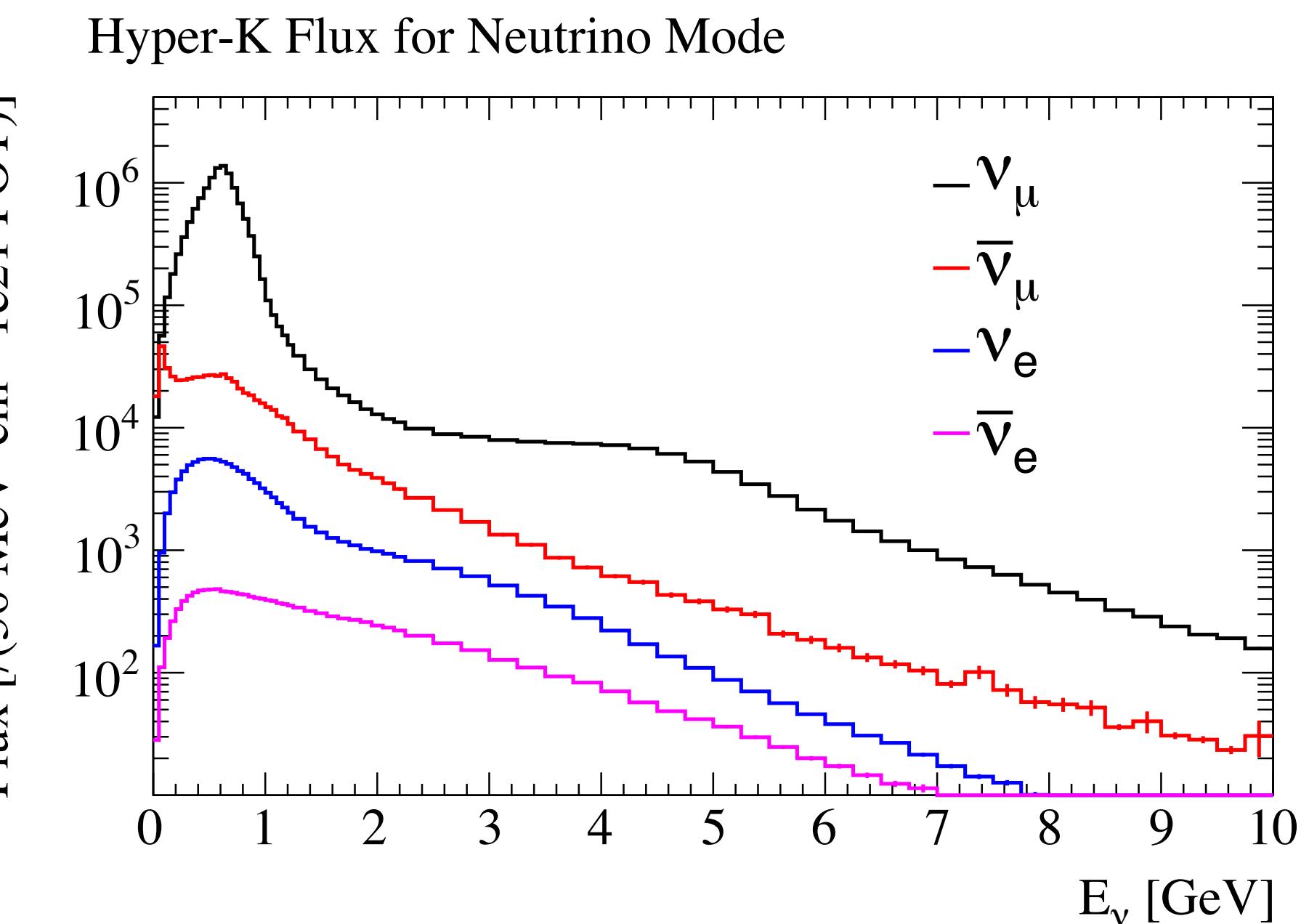
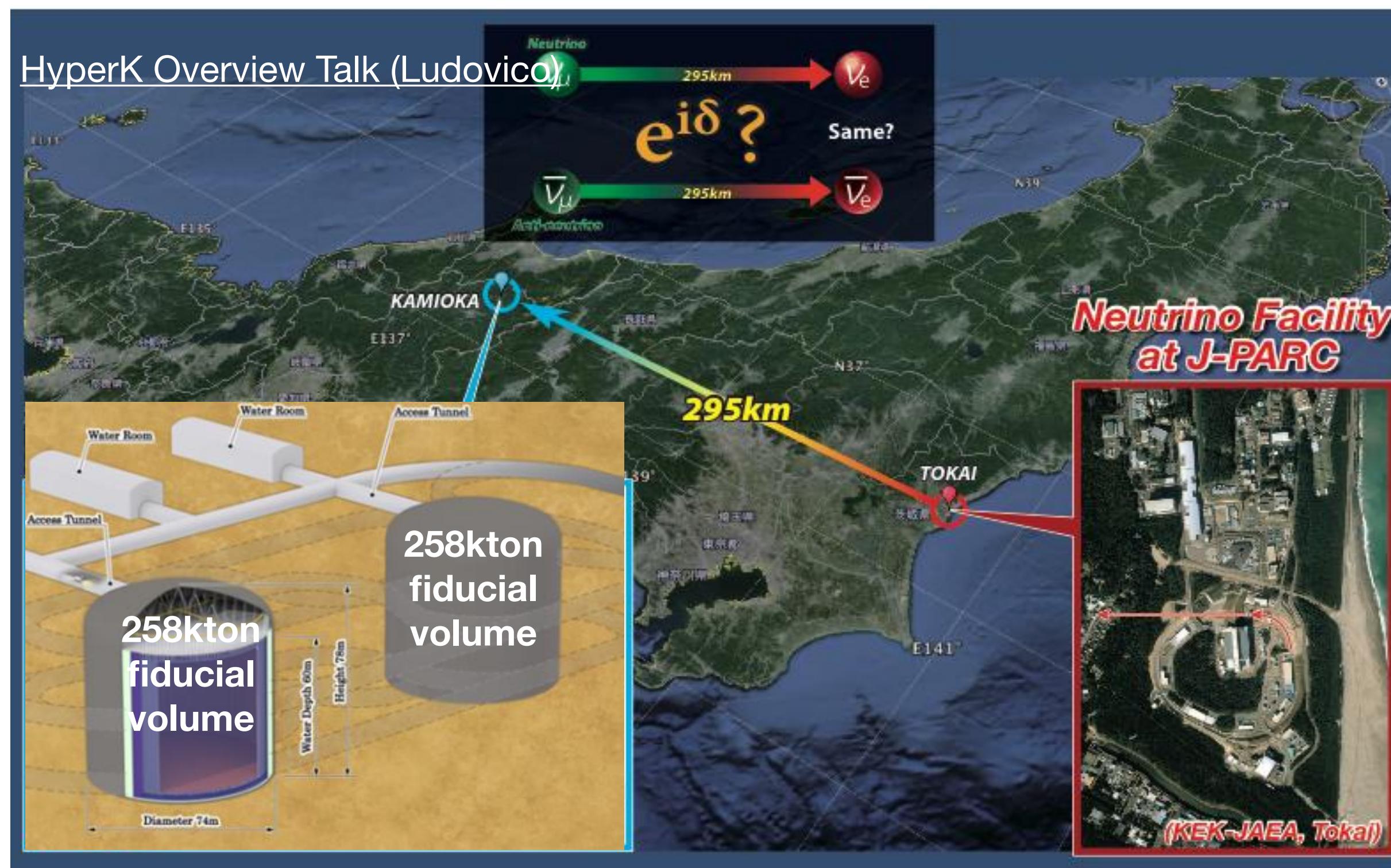
The Future of Neutrinos: Mass Ordering

Precision measurement of $P(\bar{\nu}_e \rightarrow \bar{\nu}_e)$ will resolve mass ordering with 6 years of data taking.



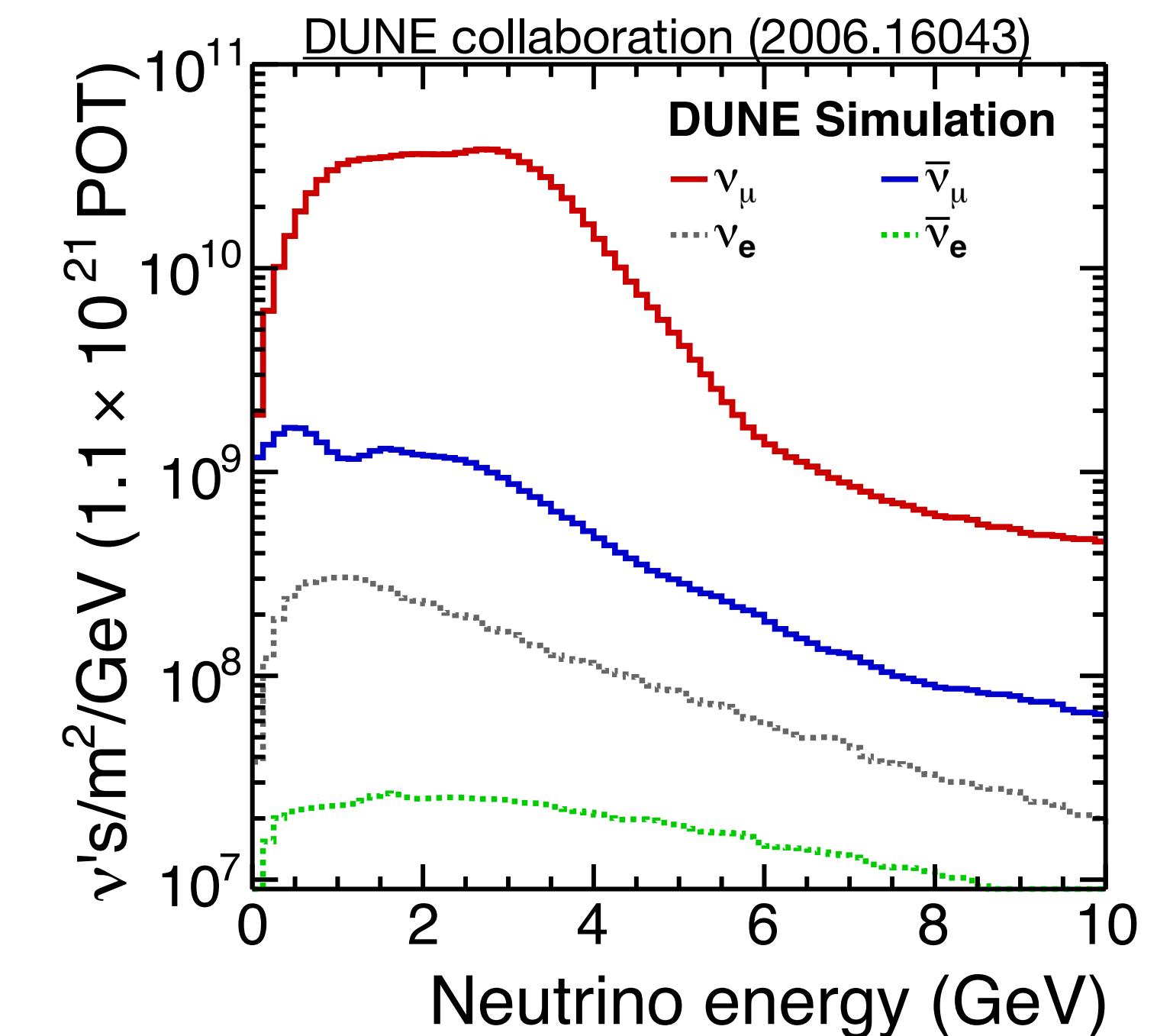
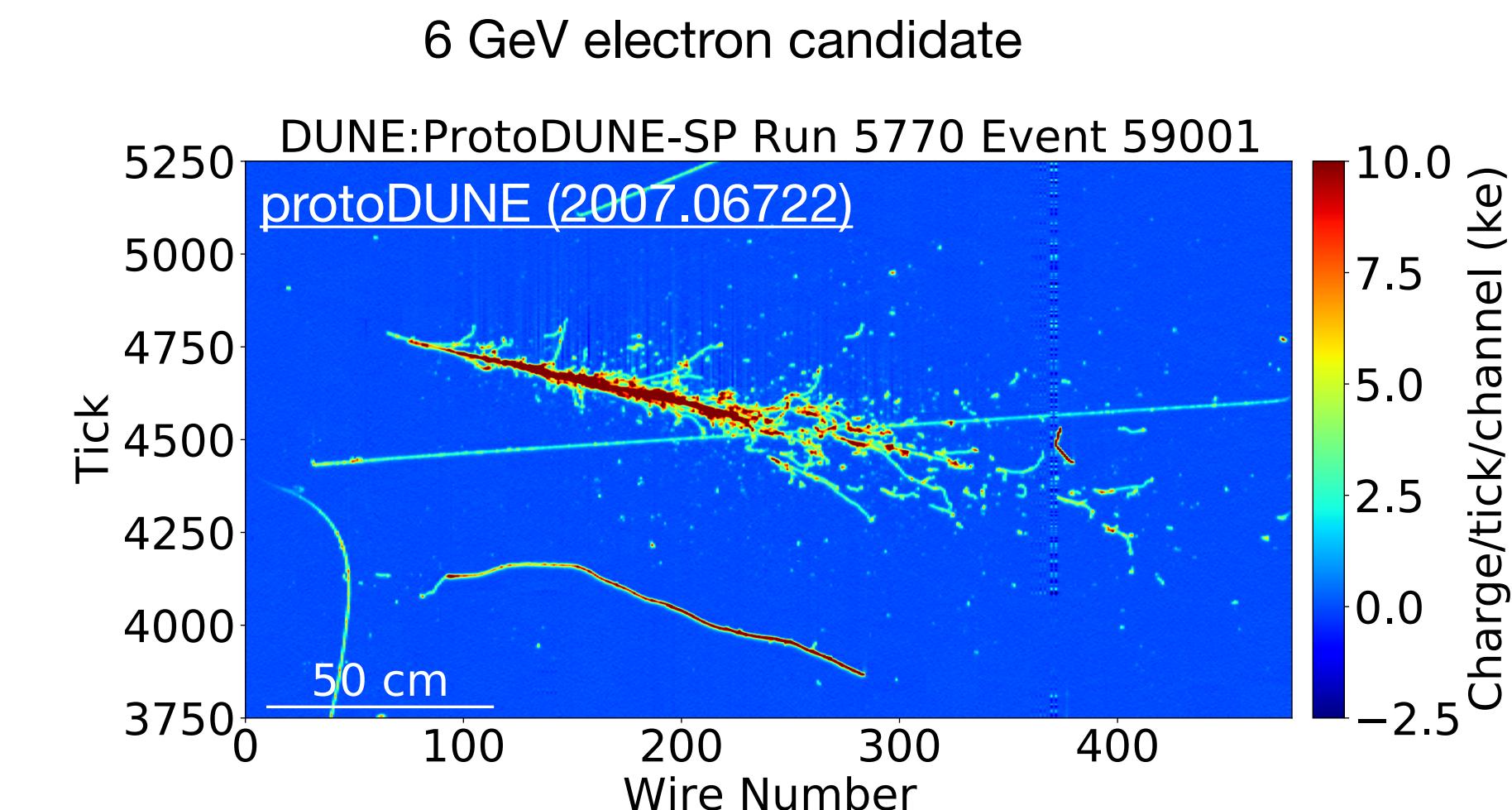
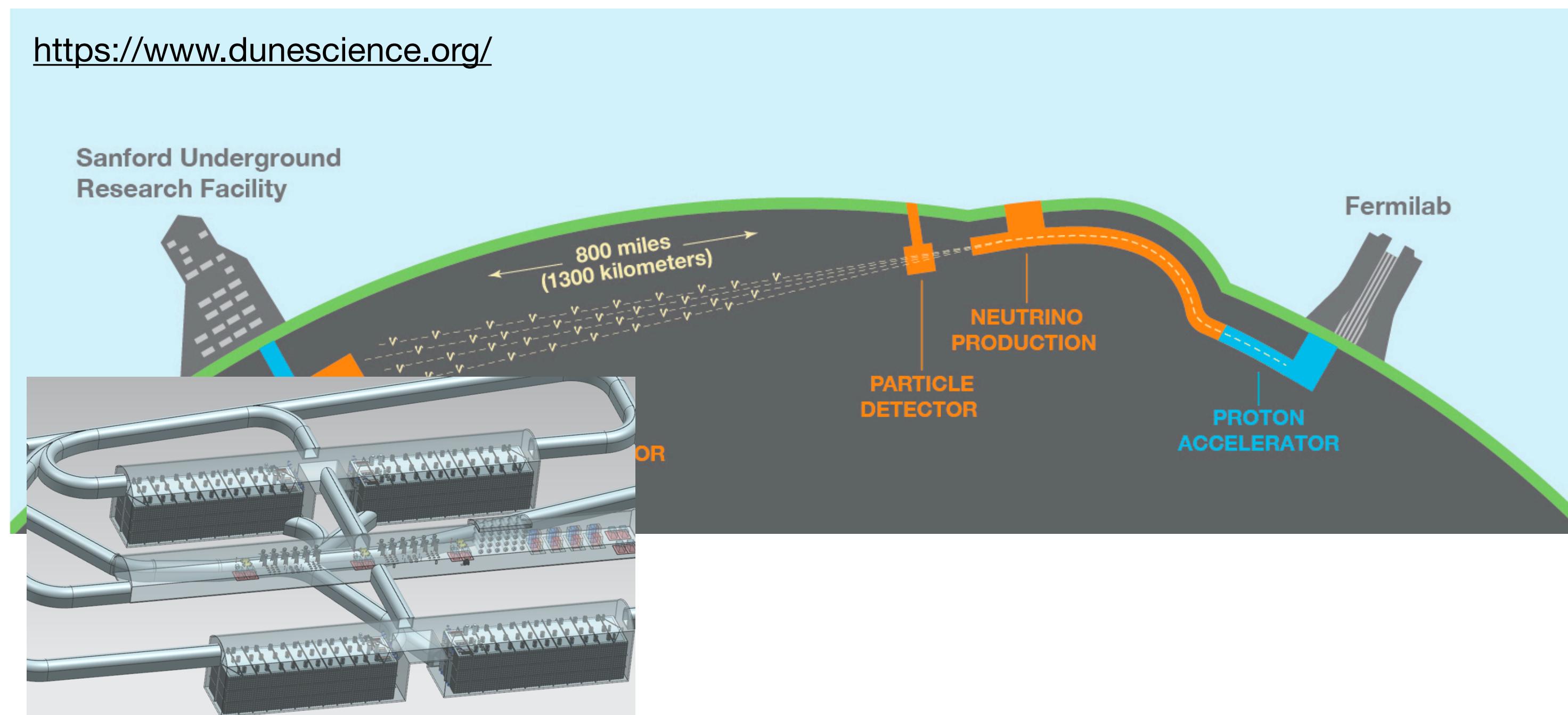
The Future of Neutrinos

- HyperK next generation multipurpose water Cherenkov detector will start data-taking 2027. Similar technology to T2K.
- **physics goals:** CP-violation, mass ordering, proton decay, detect solar, atmospheric, and geoneutrinos...

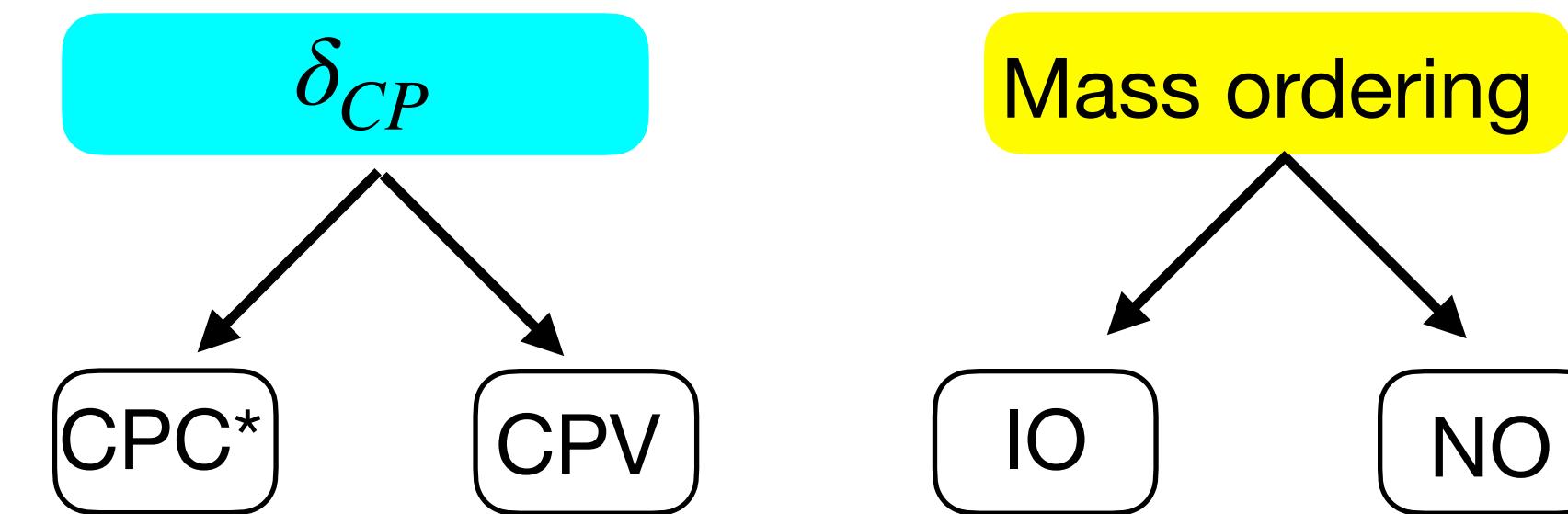


The Future of Neutrinos

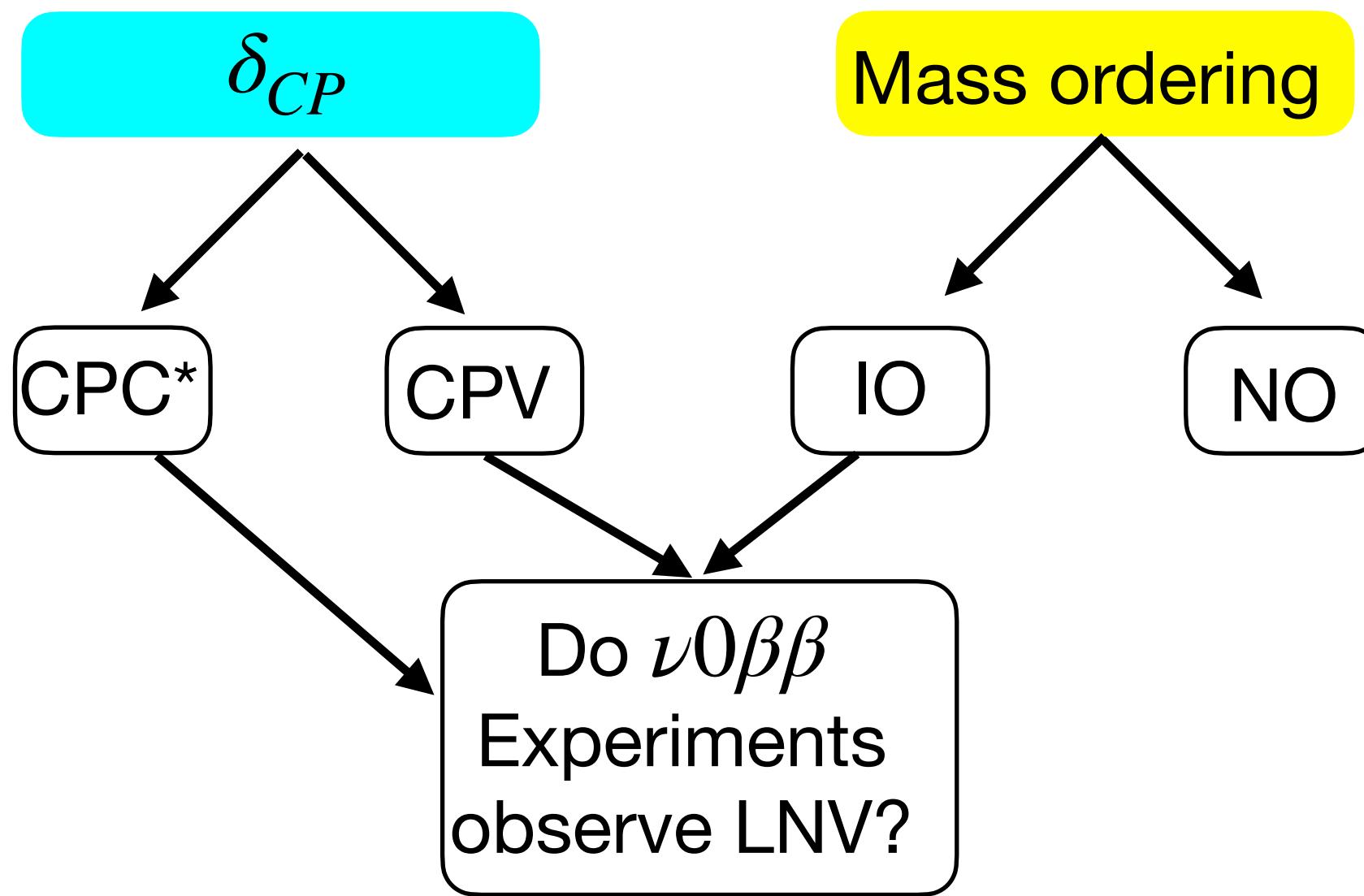
- DUNE LArTPC detector will start running 2026 (3 years no beam). Technology demonstrated by the Fermilab Short Baseline Neutrino Program (Kirsty Duffy's talk).
- **Physics goals:** very similar to HyperK



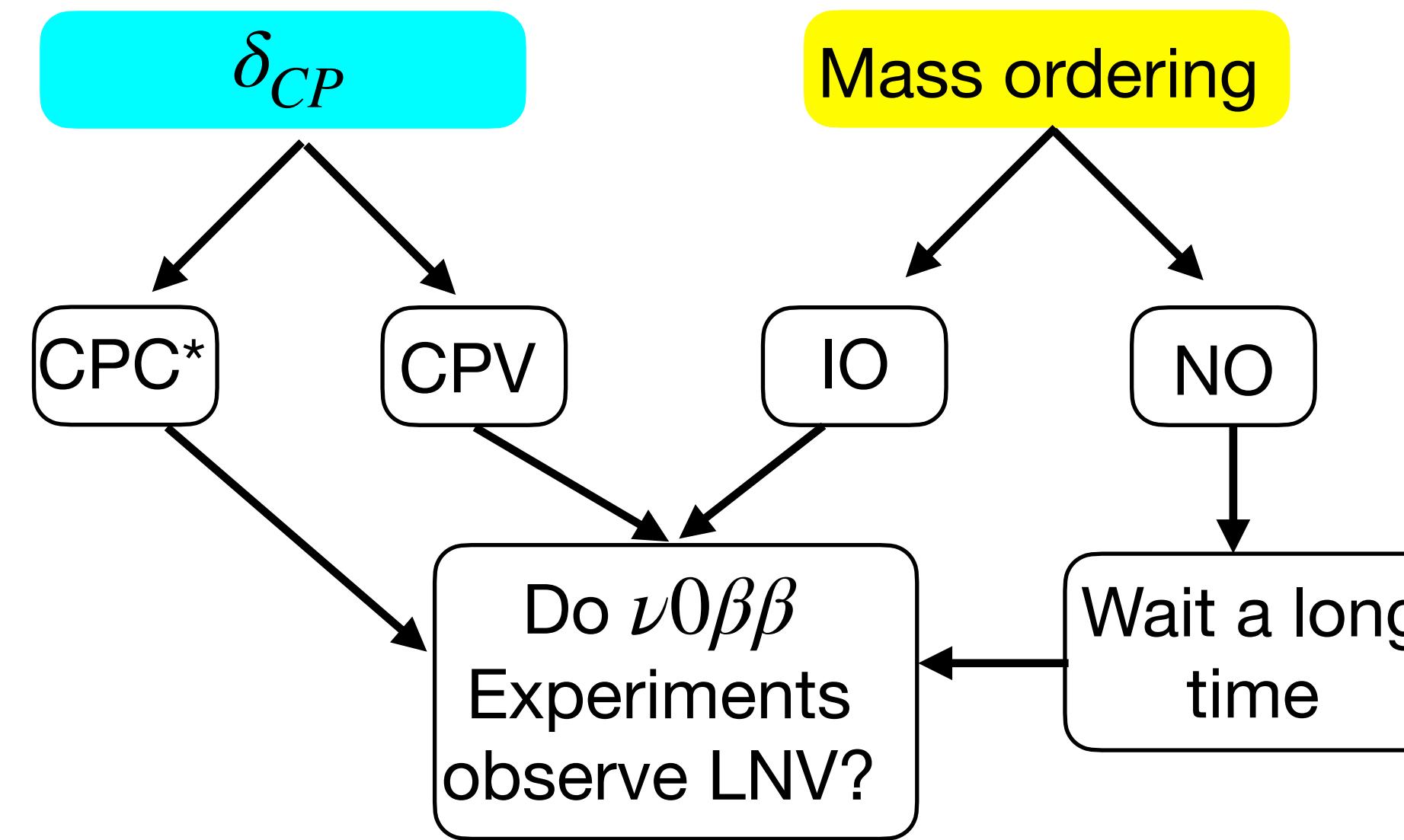
The Future of Neutrinos



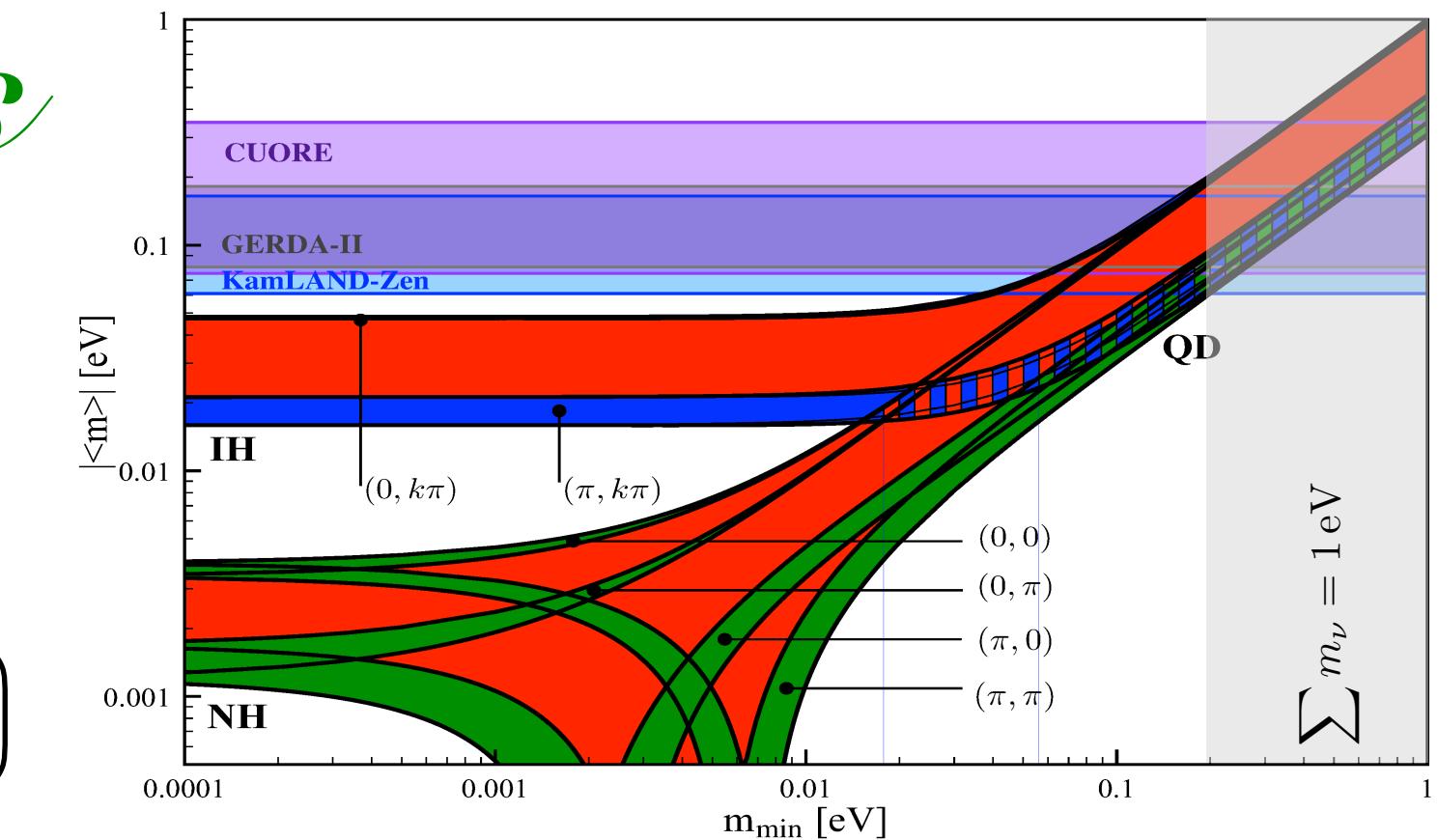
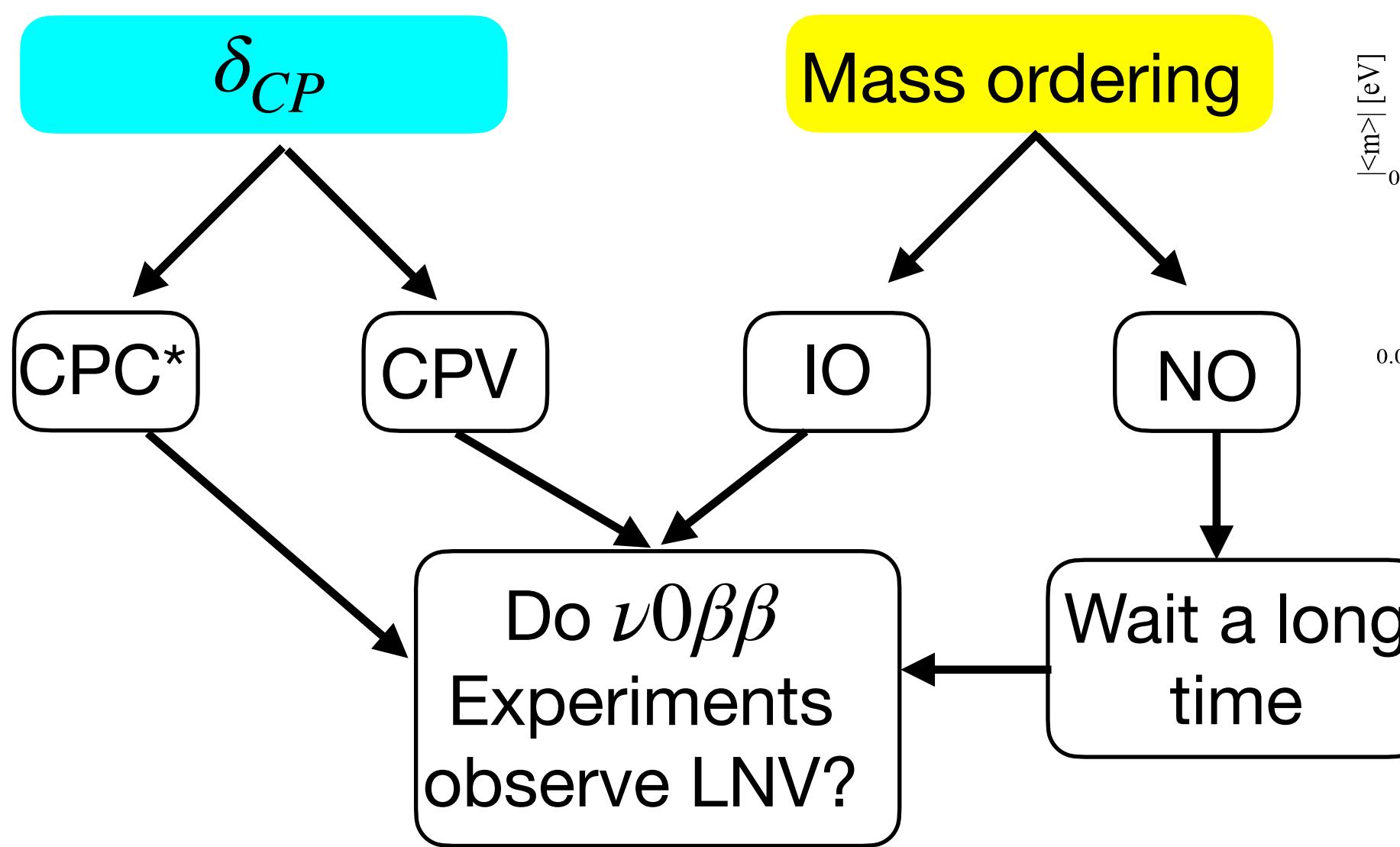
The Future of Neutrinos



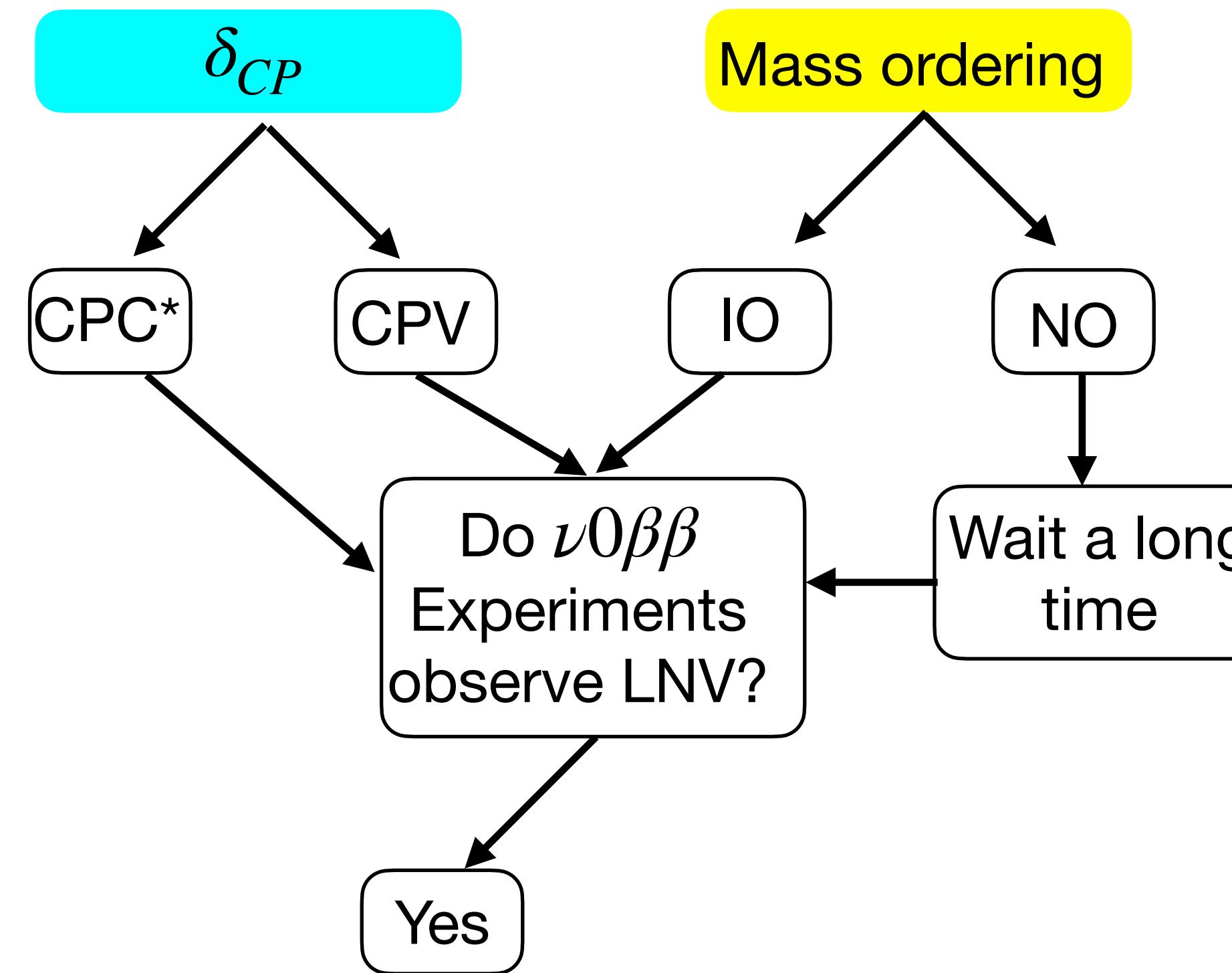
The Future of Neutrinos



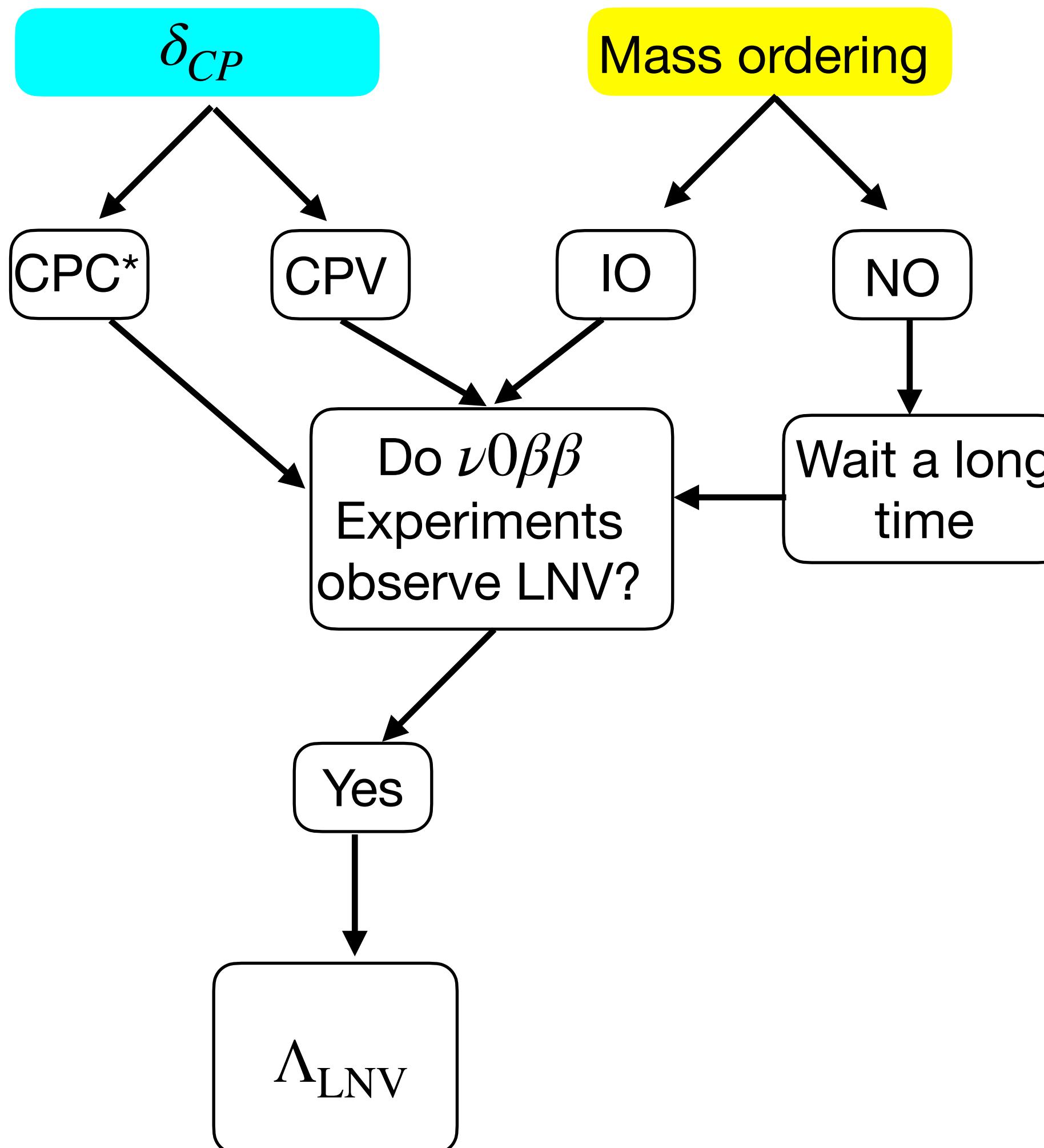
The Future of Neutrinos



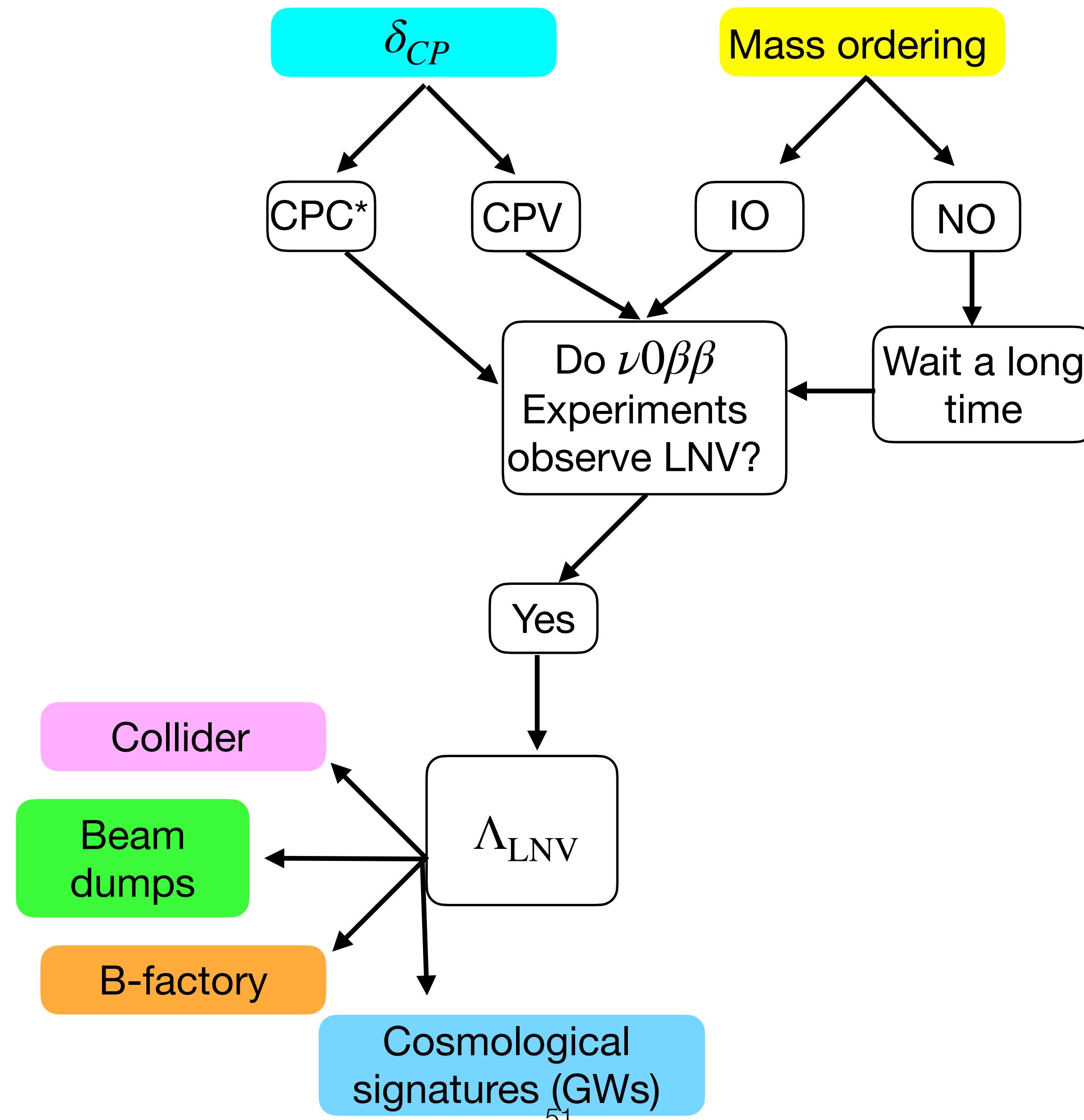
The Future of Neutrinos



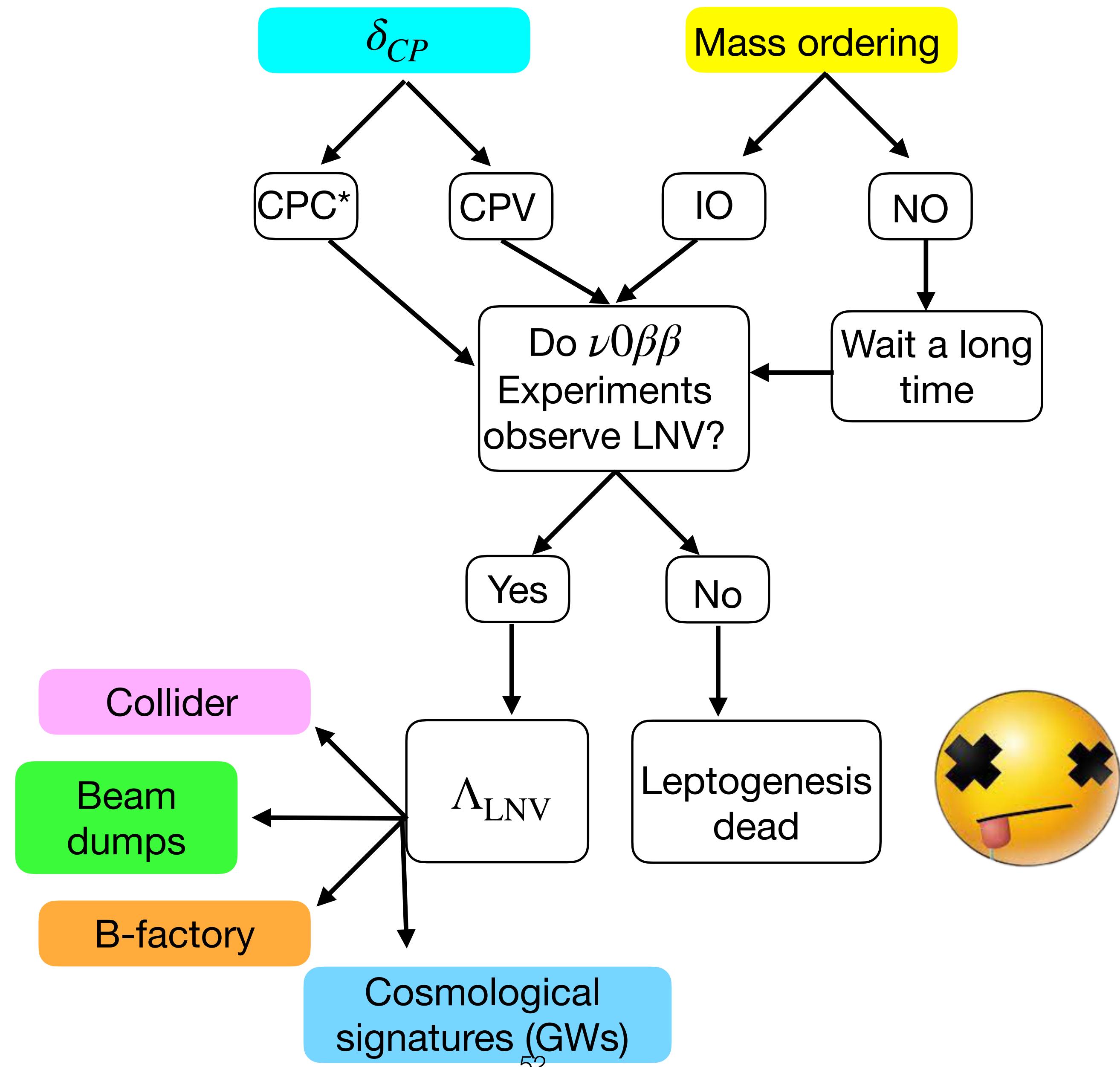
The Future of Neutrinos



The Future of Neutrinos



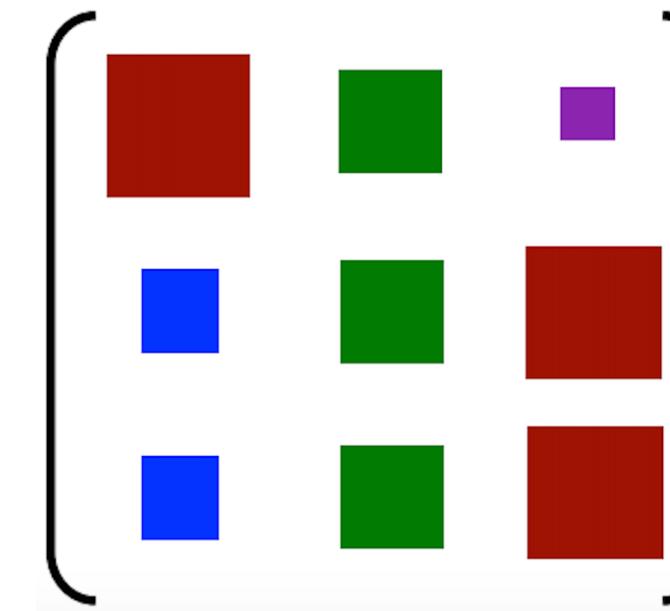
The Future of Neutrinos



The Future of Neutrinos

[Frogatt & Nielsen](#)

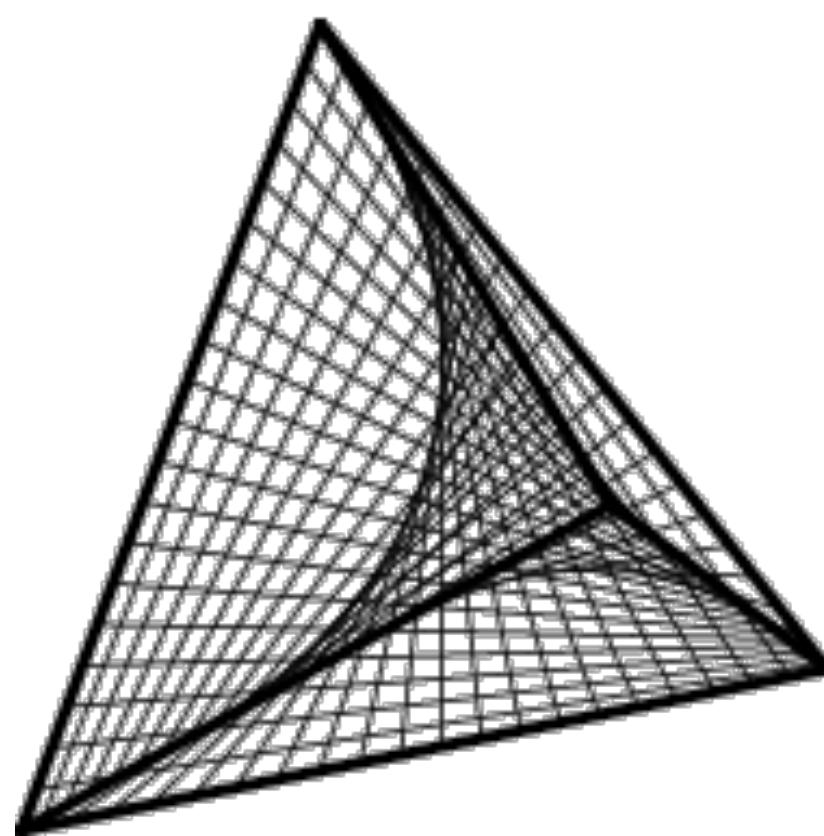
Precision determination of mixing angles



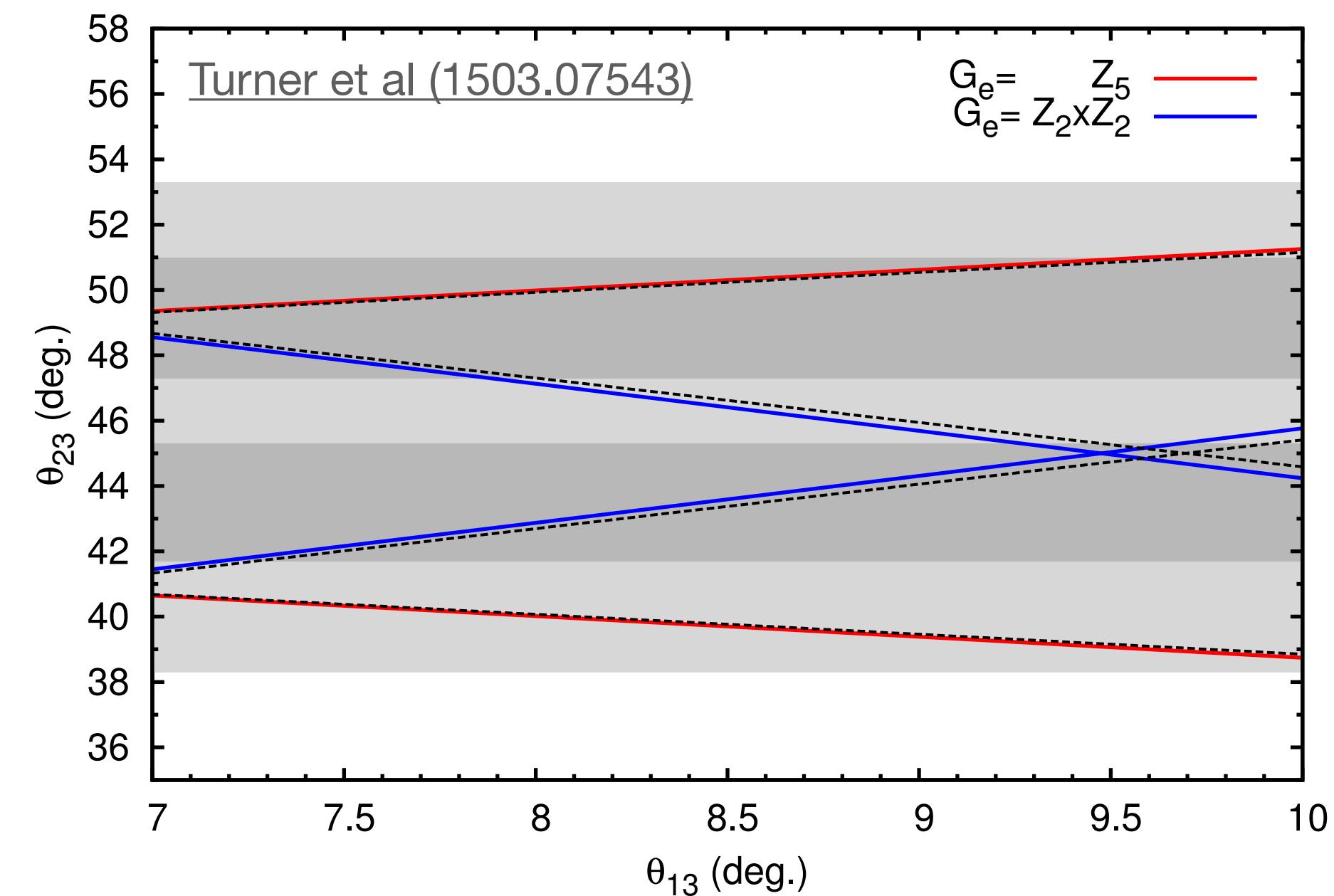
Flavour
Anarchy

[De Gouvea & Murayama](#)

Discrete symmetries popular way to address “flavour puzzle”



$$U_{\text{TBM}} = \begin{pmatrix} \frac{2}{\sqrt{6}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

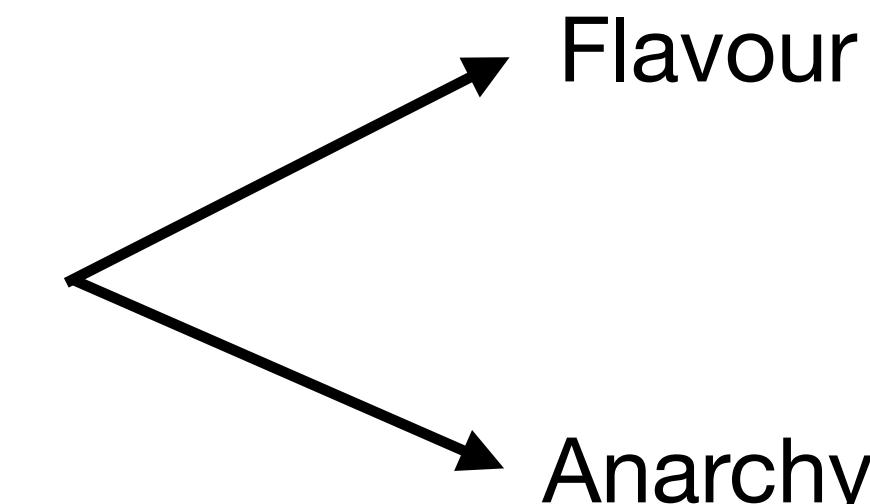
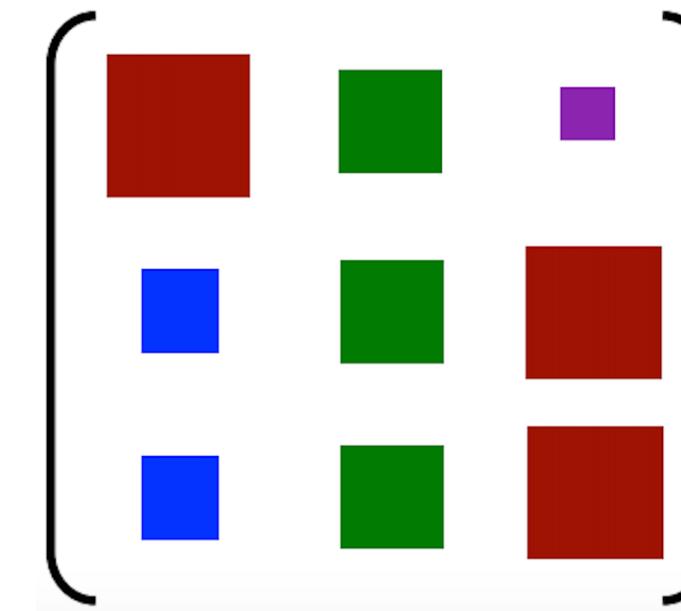


If flavour breaking scale is not high, we can search for cLFV, collider signatures

The Future of Neutrinos

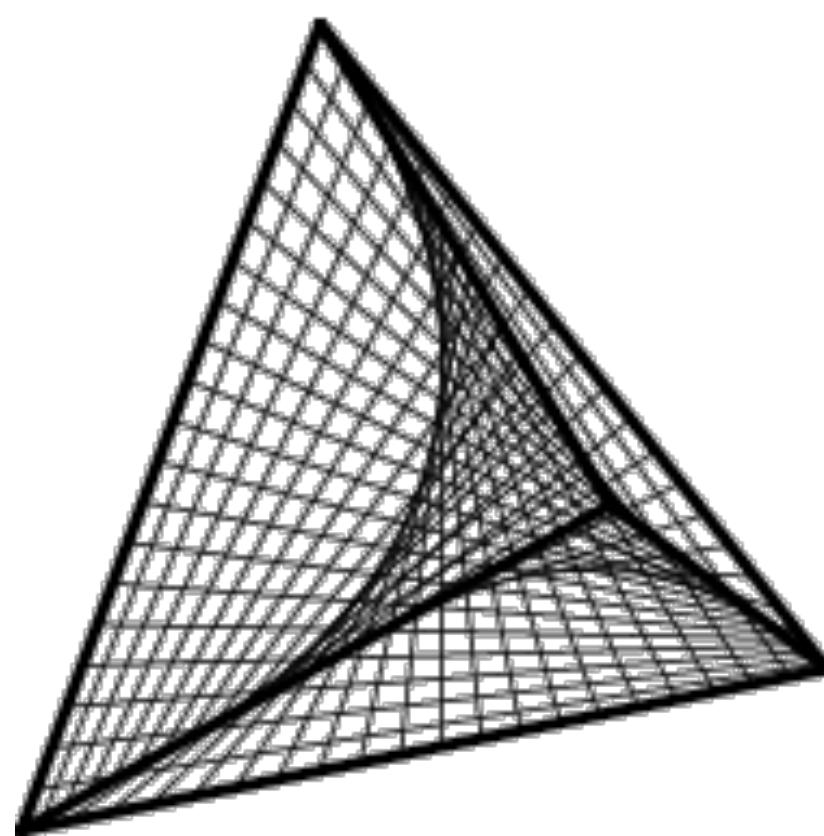
[Frogatt & Nielsen](#)

Precision determination of mixing angles



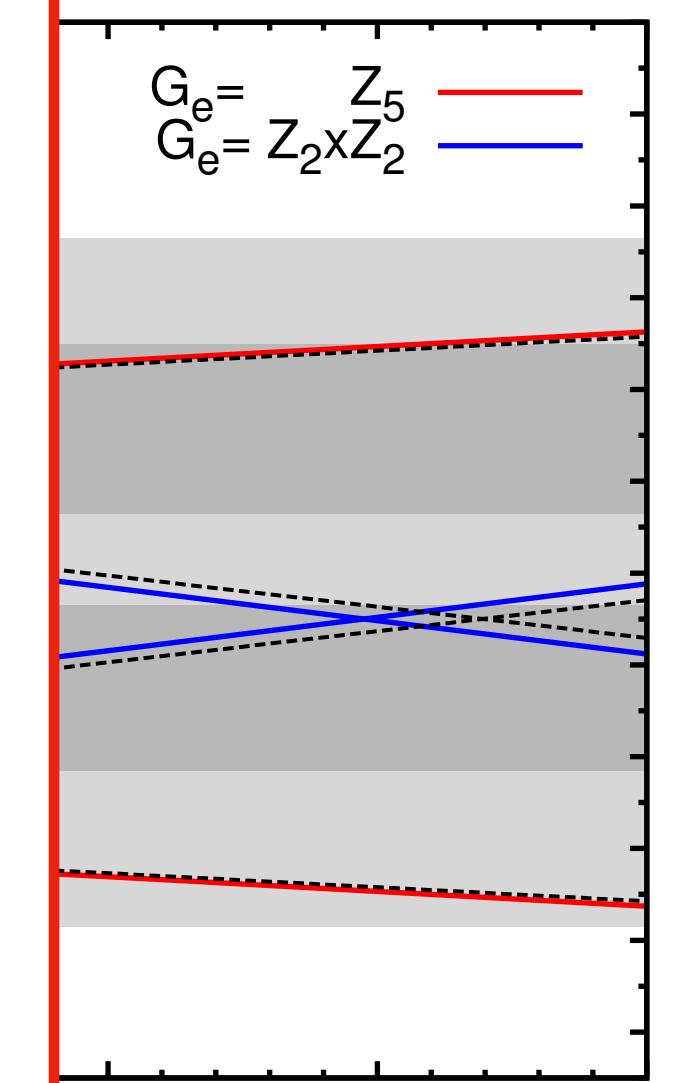
[De Gouvea & Murayama](#)

Discrete symmetries popular way to address “flavour puzzle”



Interesting connection with GWs: if there was a discrete flavour symmetry you would get domain walls in the early Universe. The explicit breaking of the residual flavour symmetries can cause the domain walls to decay giving a GW signal

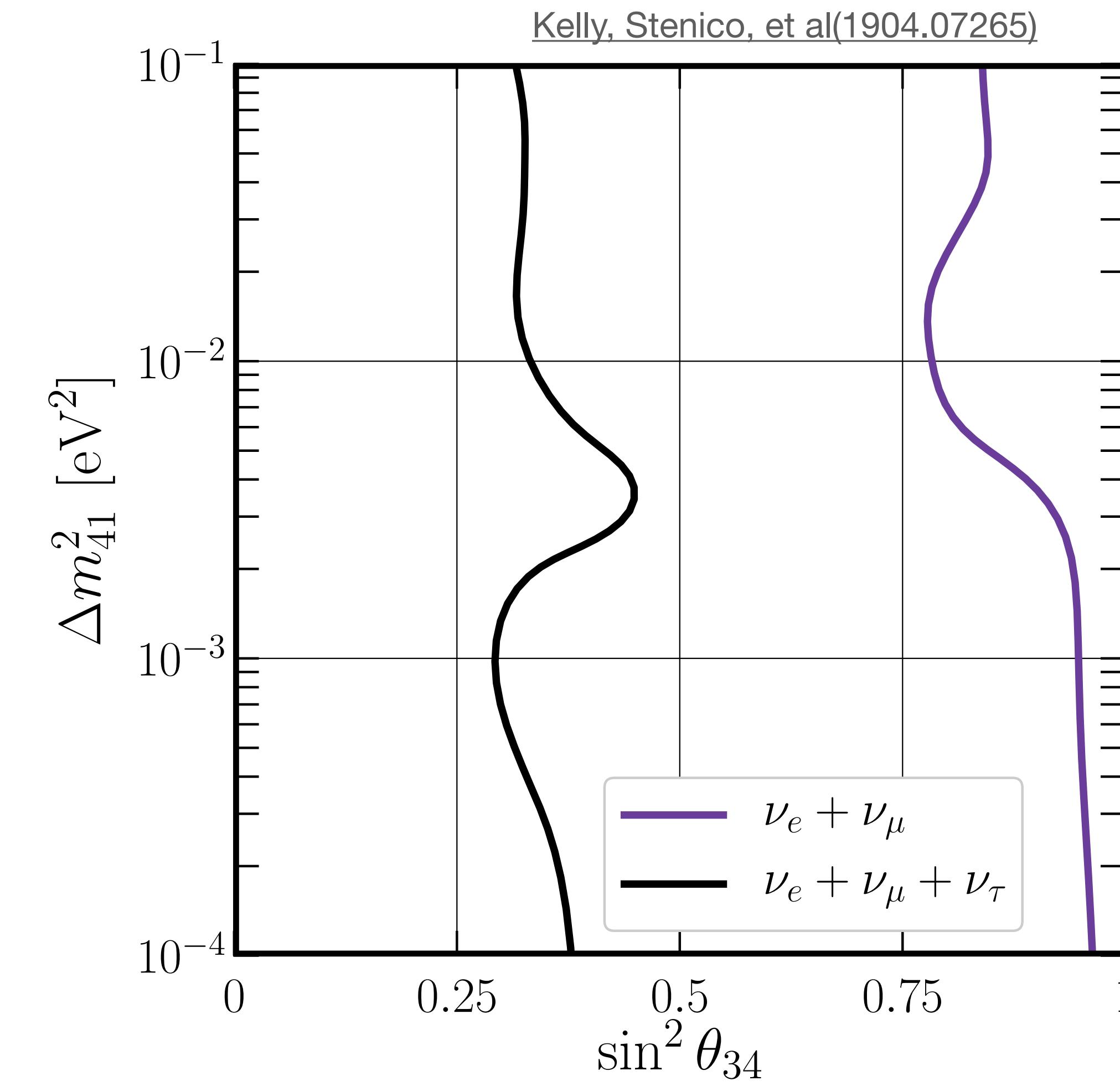
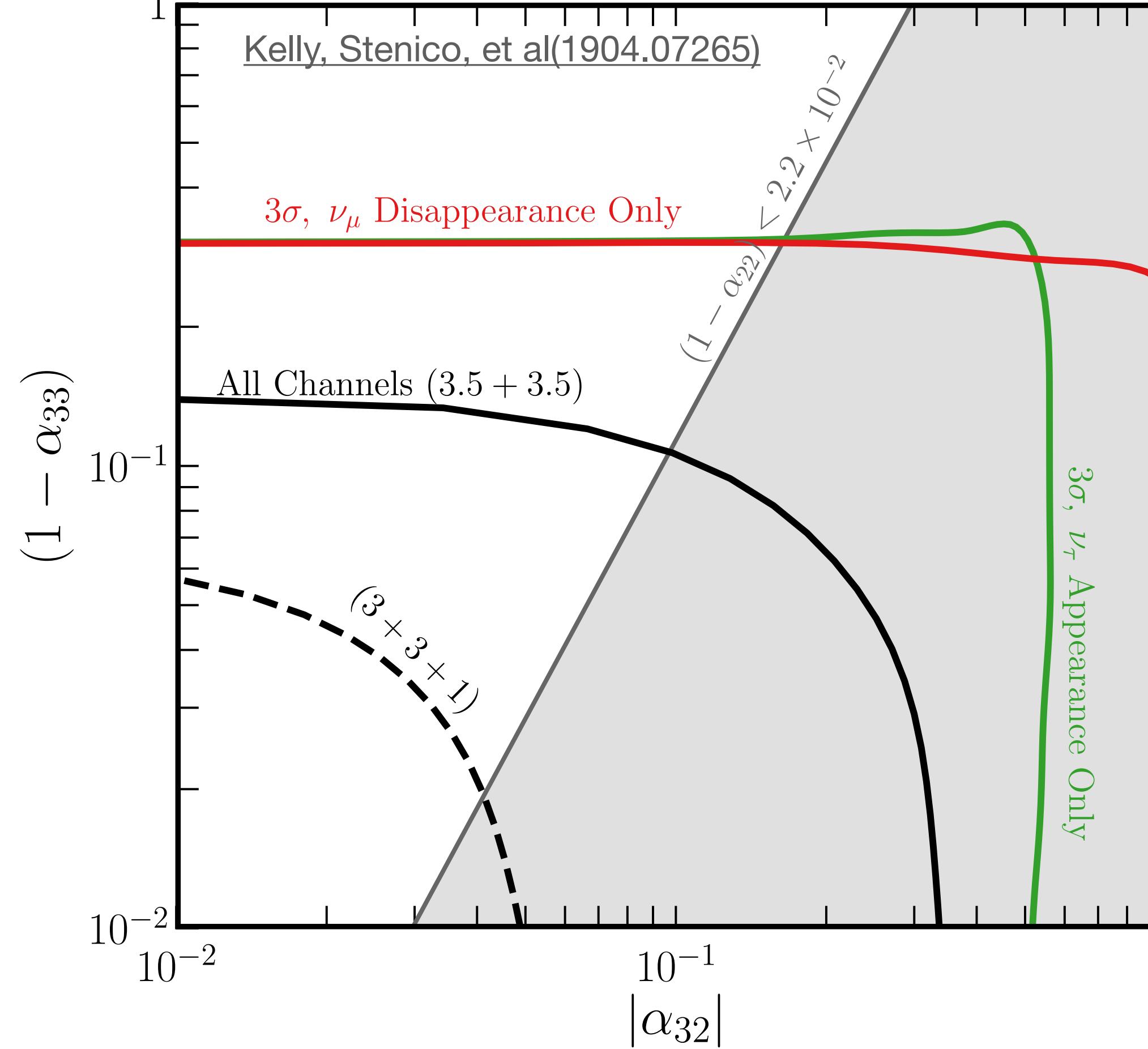
[Pascoli, Zhou et al \(2009.01903\)](#)



If flavour breaking scale is not high, we can search for cLFV, collider signatures

The Future of Neutrinos

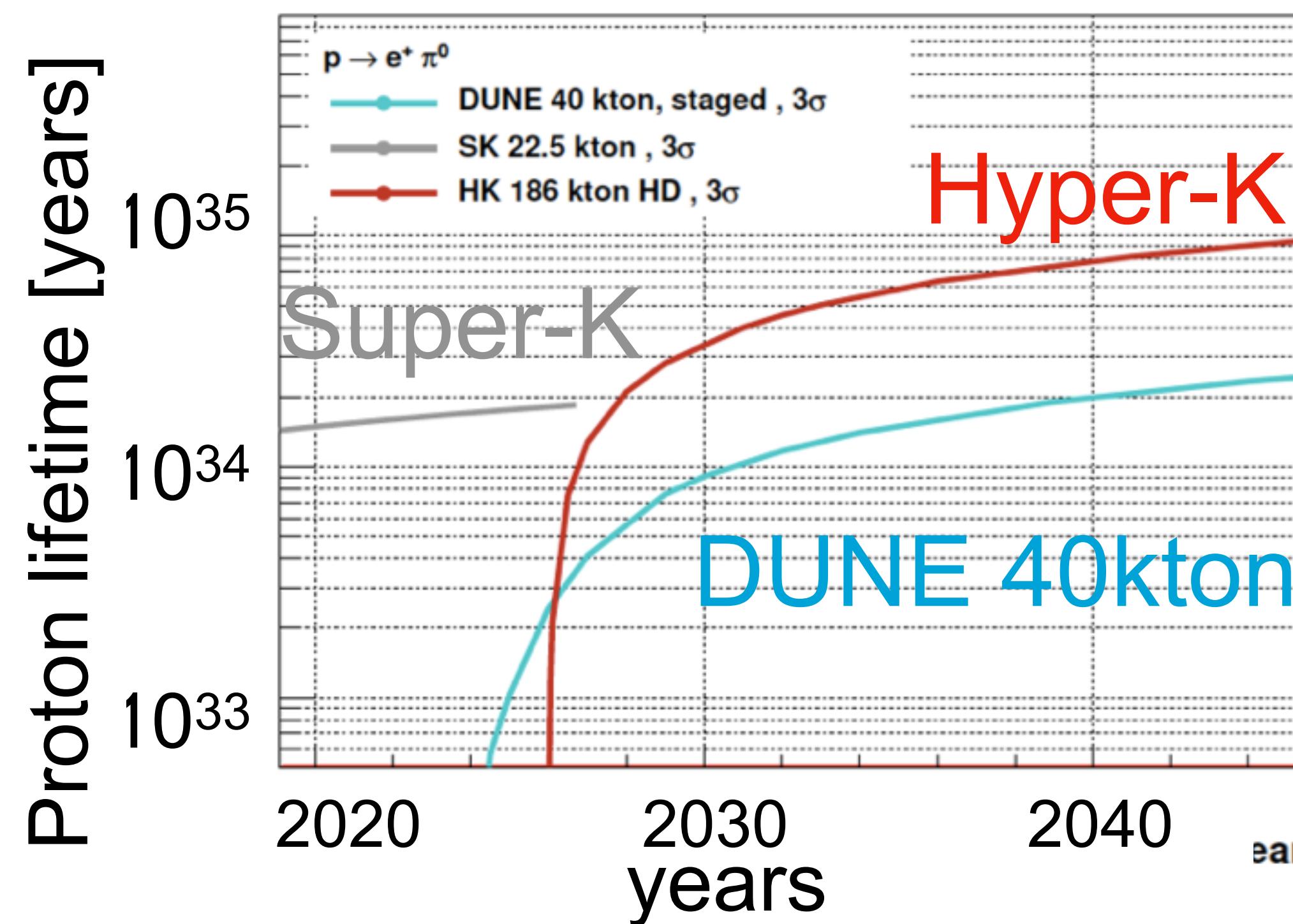
Precision determination of mixing angles will also be improved by measuring tau neutrino appearance.



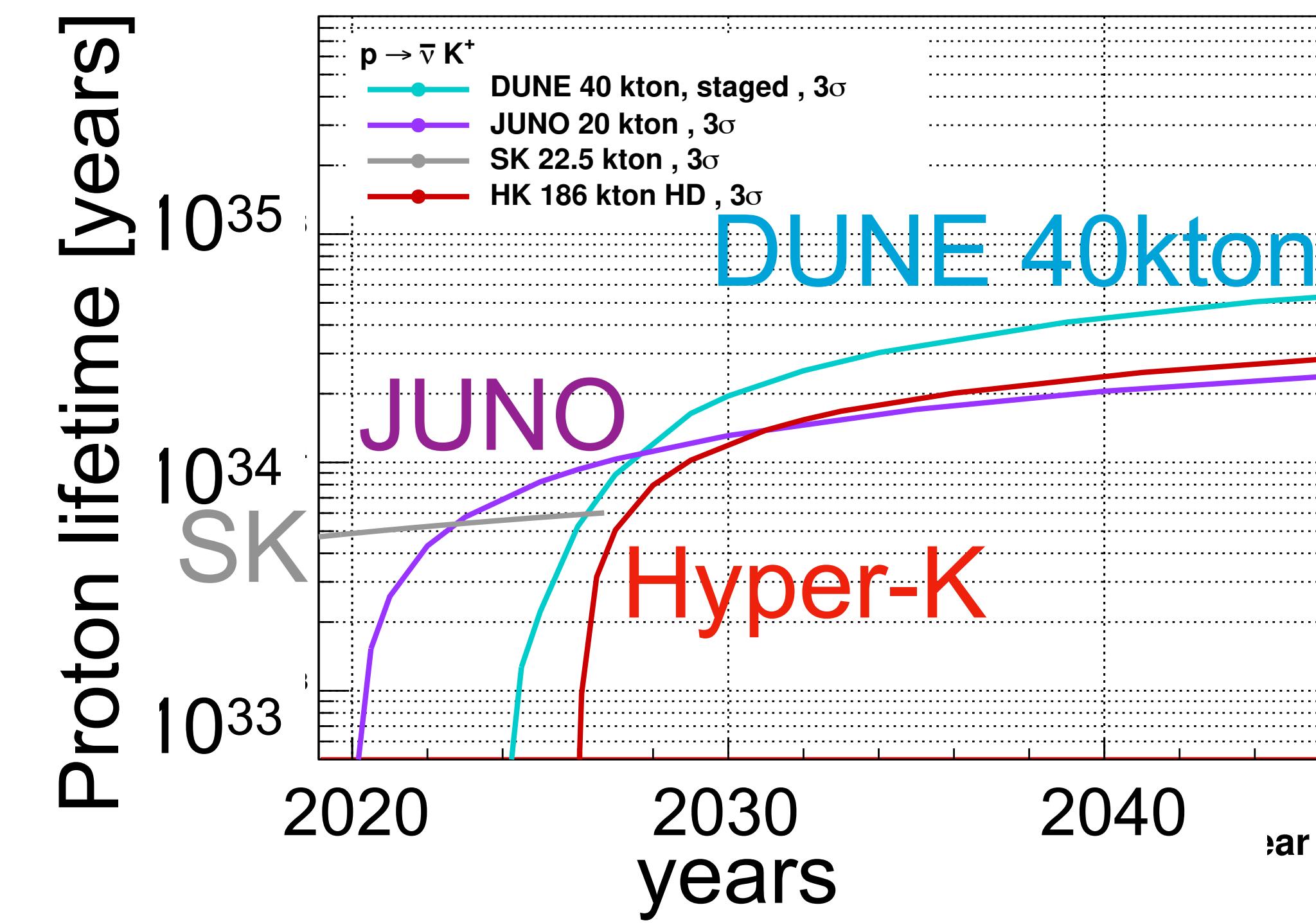
The Future of Neutrinos

Shiozawa (neutrino townhall 2018)

$p \rightarrow e^+ \pi^0$ 3σ discovery



$p \rightarrow \nu K^+$ 3σ discovery

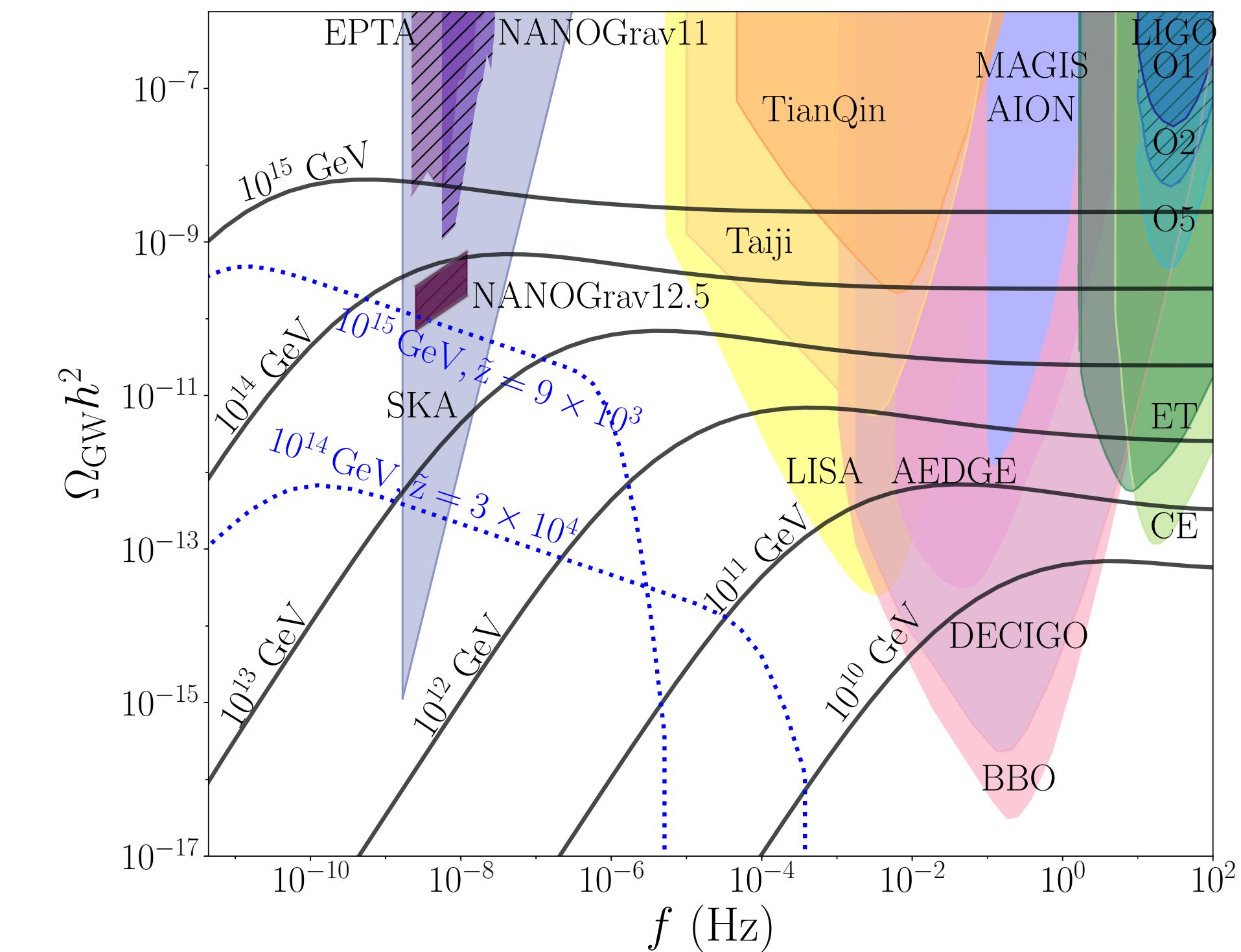
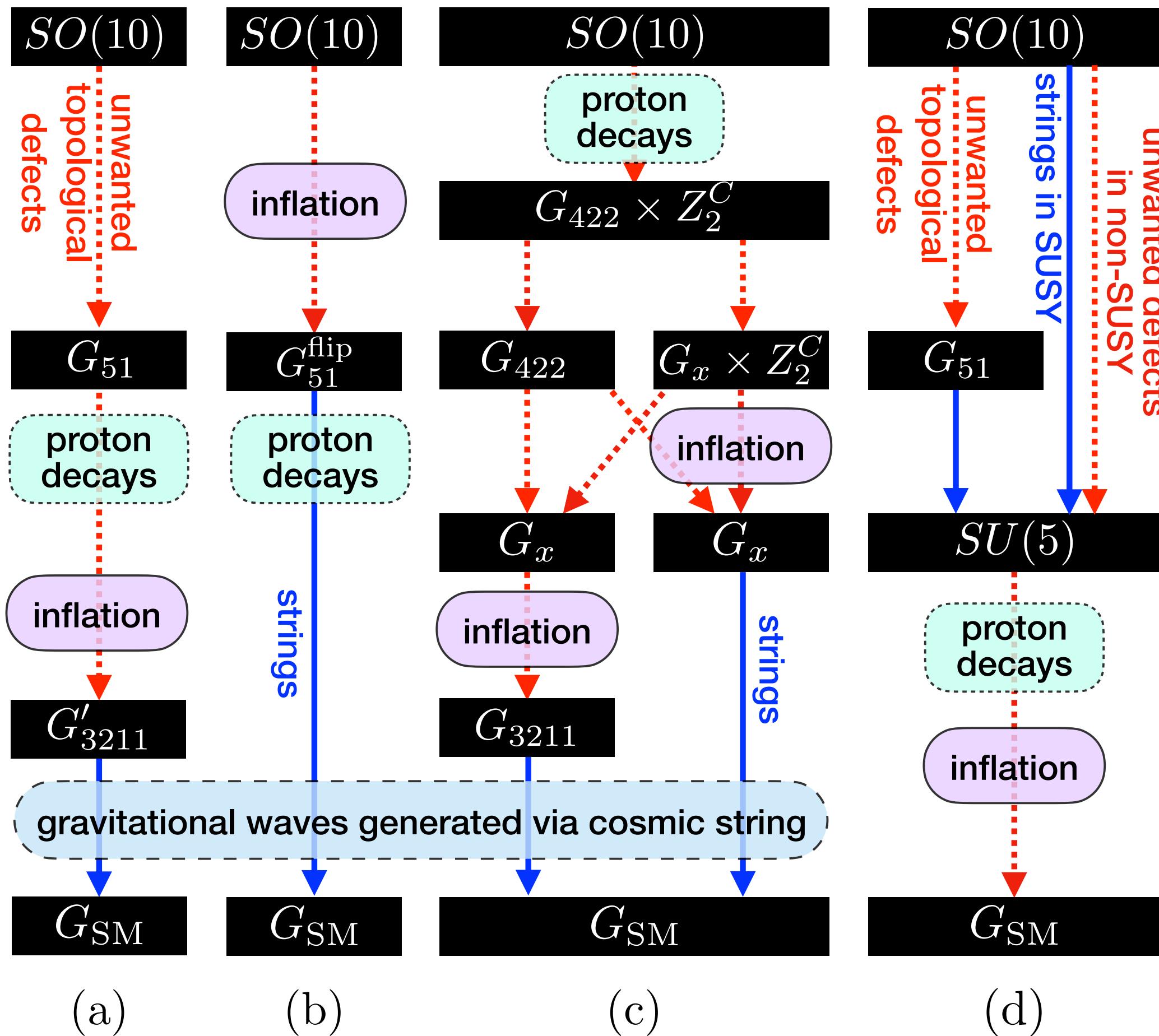


GUTs produce cosmic string networks \implies gravitational waves

Correlate limits (or observations!) from PD with observations/non-observations of GWs to assess symmetry breaking chains

The Future of Neutrinos

[King, Pascoli, Turner, Zhou \(2005.13549\)](#)



Observables		Proton decays	
GWs	Observed	$p \rightarrow \pi^0 e^+$ observed \Rightarrow non-SUSY contribution indicated	
	Marginal	<ul style="list-style-type: none"> types (a) and (c) favoured types (b) and (d) excluded 	
		<ul style="list-style-type: none"> type (a) allowed if $p \rightarrow K^+ \bar{\nu}$ not observed and $\Lambda_{\text{pd}} \sim \Lambda_{\text{cs}}$ 	

The Future of Neutrinos

- Axions at the DUNE near detector [Kelly, Liu et al\(2011.05995\)](#)
- Primordial black holes at DUNE [De Romeri et al \(2106.05013\)](#)
- Earth's tomography using DUNE [Kelly, Perez-Gonzalez et al \(2110.00003\)](#)
- Measuring the diffuse neutrino background at HyperK [De Gouvea, Perez-Gonzalez et al \(2007.13748\)](#)
- Boosted Dark Matter at HyperK & DUNE [Necib et al \(1610.03486\)](#)
- Lorentz violation HyperK [Litchfield\(NuFact 2019\)](#)
- Non-standard interactions (ν EFT) at HyperK and DUNE [NSI report \(1907.00991\)](#)
- Supernovae neutrino detection [DUNE report](#) [HyperK \(2101.05269\)](#)
- Time varying neutrino masses [Krnjaic, Machado, Necib \(1705.06740\)](#)

Experiments

Filter this page
(Note: The process can take some time.)

Search Reset

Types of Neutrino Experiments

- Accelerator Beam Dump
- Astrophysical Neutrinos
- Atmospheric Neutrinos
- Electron Neutrino
- Long-Baseline Neutrino Oscillations
- Muon Neutrino
- Neutrinoless Double Beta Decay
- Short-Baseline Neutrino Oscillations
- Solar Neutrinos
- Supernova Neutrinos
- Tau Neutrino

EXPAND ALL **COMPRESS ALL**

Neutrino Experiments

- AMANDA High-Energy Astrophysical Neutrinos, Supernova Neutrinos (Home, INSPIRE) References
- ANTARES High-Energy Astrophysical Neutrinos (Home, INSPIRE) References
- ARIANNA High-Energy Astrophysical Neutrinos (Home, INSPIRE) References
- ArgoNeuT Neutrino Interactions (INSPIRE) References
- Baikal High-Energy Astrophysical Neutrinos, Supernova Neutrinos (Home, INSPIRE) References
- Baksan Atmospheric and Supernova SN1987A Neutrinos (Home, INSPIRE) References
- BOREXino Solar Neutrinos (Home, INSPIRE) References
- BECB Accelerator SBL Oscillations, Neutrino Interactions (INSPIRE) References
- BNL-E-734 Neutrino Interactions (INSPIRE) References
- BNL-E-776 Accelerator SBL Oscillations (INSPIRE) References
- BNL-E-816 Accelerator SBL Oscillations (INSPIRE) References
- Bugey Reactor SBL Oscillations (INSPIRE) References
- CCFR Accelerator SBL Oscillations, Muon Neutrino - Nucleon Scattering (Home, INSPIRE) References
- CDHSW Accelerator SBL Oscillations, Muon Neutrino - Nucleon Scattering (INSPIRE) References
- CERN-PS-191 Accelerator SBL Oscillations (INSPIRE) References
- CHARM Accelerator SBL Oscillations (INSPIRE) References
- Chooz Reactor LBL Oscillations (Home, INSPIRE) References
- CHORUS Accelerator SBL Oscillations (Home, INSPIRE) References
- COBRA Double Beta Decay (^{70}Zn , ^{116}Cd , ^{128}Te , ^{130}Te) (Home) References
- CUORE Double Beta Decay (^{130}Te) (Home) References
- CUORICINO Double Beta Decay (^{130}Te) (Home) References
- DANSS Reactor SBL Oscillations (INSPIRE) References
- Daya Bay Reactor LBL Oscillations (Home) References
- DONUT Tau Neutrino Interactions (Home, INSPIRE) References
- Double Chooz Reactor LBL Oscillations (Home) References
- ELEGANT Double Beta Decay (^{48}Ca , ^{100}Mo) References
- EXO Double Beta Decay (^{136}Xe) (Home) References
- FNAL-E-0053 Accelerator SBL Oscillations (INSPIRE) References
- FNAL-E-0531 Accelerator SBL Oscillations (INSPIRE) References
- FNAL-E-0613 Accelerator SBL Oscillations (INSPIRE) References
- Frejus Atmospheric Neutrinos, Proton Decay (INSPIRE) References
- GALLEX Solar Neutrinos, SBL Oscillations with ^{51}Cr Neutrino Source (Home, INSPIRE) References
- Gargamelle Accelerator SBL Oscillations, Neutrino Interactions (INSPIRE) References
- Genova ^{187}Re Electron Neutrino Mass (Home, INSPIRE) References
- GERDA Double Beta Decay (^{76}Ge) (Home) References
- GEMMA Reactor Electron Antineutrino - Electron Scattering References
- GLUE High-Energy Astrophysical Neutrinos (Home, INSPIRE) References
- GNO Solar Neutrinos (Home, INSPIRE) References
- Gosgen Reactor SBL Oscillations (INSPIRE) References
- Gotthard Double Beta Decay (^{136}Xe) (Home) References
- Heidelberg-Moscow Double Beta Decay (^{76}Ge) (Home) References
- Homestake Solar Neutrinos (INSPIRE) References
- ICARUS Accelerator Long Baseline Oscillations, Supernova Neutrinos, Proton Decay (Home, INSPIRE) References
- IceCube High-Energy Astrophysical Neutrinos (Home, INSPIRE) References
- IGEX Double Beta Decay (^{76}Ge) (INSPIRE) References
- IHEP-JINR Accelerator SBL Oscillations, Neutrino-Nucleon Interactions (INSPIRE) References
- ILL Reactor SBL Oscillations (INSPIRE) References
- IMB Atmospheric Neutrinos, Supernova SN1987A Neutrinos, Proton Decay (Home, INSPIRE) References
- K2K Accelerator Long Baseline Oscillations (Home, INSPIRE) References
- Kamiokande Solar, Atmospheric and Supernova SN1987A Neutrinos, Proton Decay (Home, INSPIRE) References
- KamLAND Reactor Long Baseline Oscillations, Supernova Neutrinos (Home, INSPIRE) References
- KamLAND-Zen Double Beta Decay (^{136}Xe) (Home) References
- KARMEN Accelerator SBL Oscillations (Home, INSPIRE) References
- Krasnoyarsk Reactor SBL Oscillations (INSPIRE) References
- LAMPF-0645 Accelerator SBL Oscillations (INSPIRE) References
- LAMPF-0764 Accelerator SBL Oscillations (INSPIRE) References
- LSD Supernova SN1987A Neutrinos, Astrophysical Neutrinos (INSPIRE) References
- LSND Accelerator SBL Oscillations (Home, INSPIRE) References
- LVD Supernova Neutrinos, Astrophysical Neutrinos (Home, INSPIRE) References
- MACRO Atmospheric Neutrinos (Home, INSPIRE) References
- Mainz Electron Neutrino Mass (Home, INSPIRE) References
- MiBeta Electron Neutrino Mass (Home) References
- MicroBooNE Accelerator SBL Oscillations, Neutrino Interactions (Home) References
- MINERvA Neutrino Interactions (Home, INSPIRE) References
- MiniBooNE Accelerator SBL Oscillations, Supernova Neutrinos (Home, INSPIRE) References
- MINOS Accelerator Long Baseline Oscillations, Atmospheric Neutrinos (Home, INSPIRE) References
- MUNU Reactor Electron Antineutrino - Electron Scattering (Home, INSPIRE) References
- NEMO Double Beta Decay (^{82}Se , ^{100}Mo , ^{150}Nd) / γ -rays, TCDDE References

There are many more interesting theories/physics cases and Experiments I didn't have time to discuss.

Carlo Giunti's website: <http://www.nu.to.infn.it/exp/> where he keeps a list of ongoing/future neutrino experiments

Progress is contingent upon understanding neutrino nucleon interactions better and improved modelling of underlying nuclear physics!

Thank you for your attention



$$m_\nu = \frac{Y_\nu^2 v^2}{M_N}$$



$$\sigma = \frac{G^2 F^S}{\pi}$$



And enjoy the holidays!

Backup slides leptogenesis

Mass RHN

$\mathcal{O}(10^{12})$ GeV

Fukugida & Yanagida *Phys.Lett. B17* 45-47
(1986) Buchmuller, Di Bari & Plumacher
New J.Phys. 6 105 (2004)
Barbieri, Creminelli, Strumia &
Tetradis *Nucl.Phys. B575* 61-77 (2000)

$\mathcal{O}(10^6)$ GeV

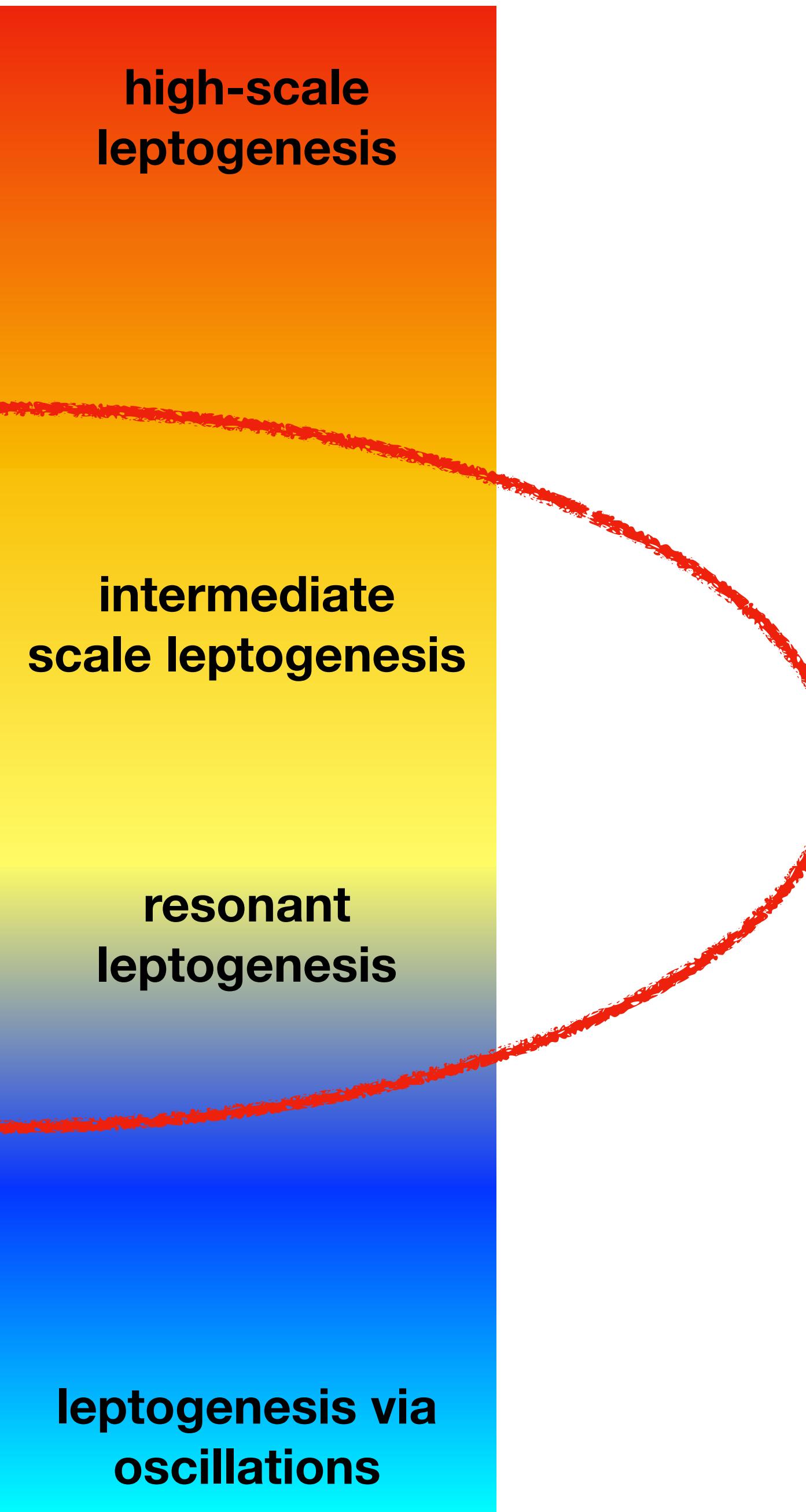
Racker, Rius & Pena *JCAP 1207* 030
(2013) Moffat, Petcov, Pascoli, Schulz & Turner
Phys.Rev. D98 no. 1, 015036 (2018)

$\mathcal{O}(10^3)$ GeV

Pilaftsis & Underwood *Nucl.Phys. B692* 303-345
(2004) Abada, Aissaoui, Losada
Nucl.Phys. B728 55-66 (2005)

$\mathcal{O}(1)$ GeV

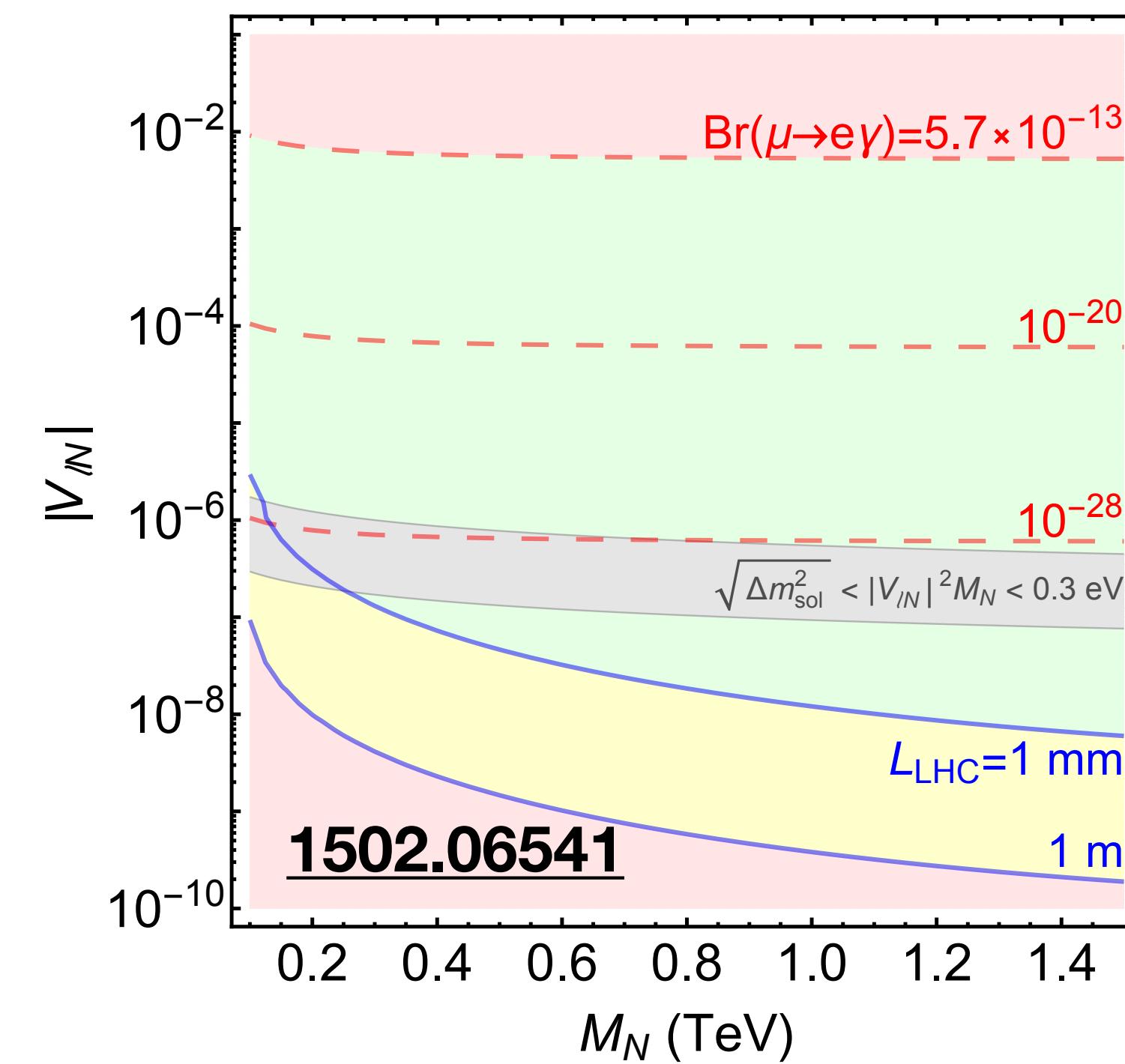
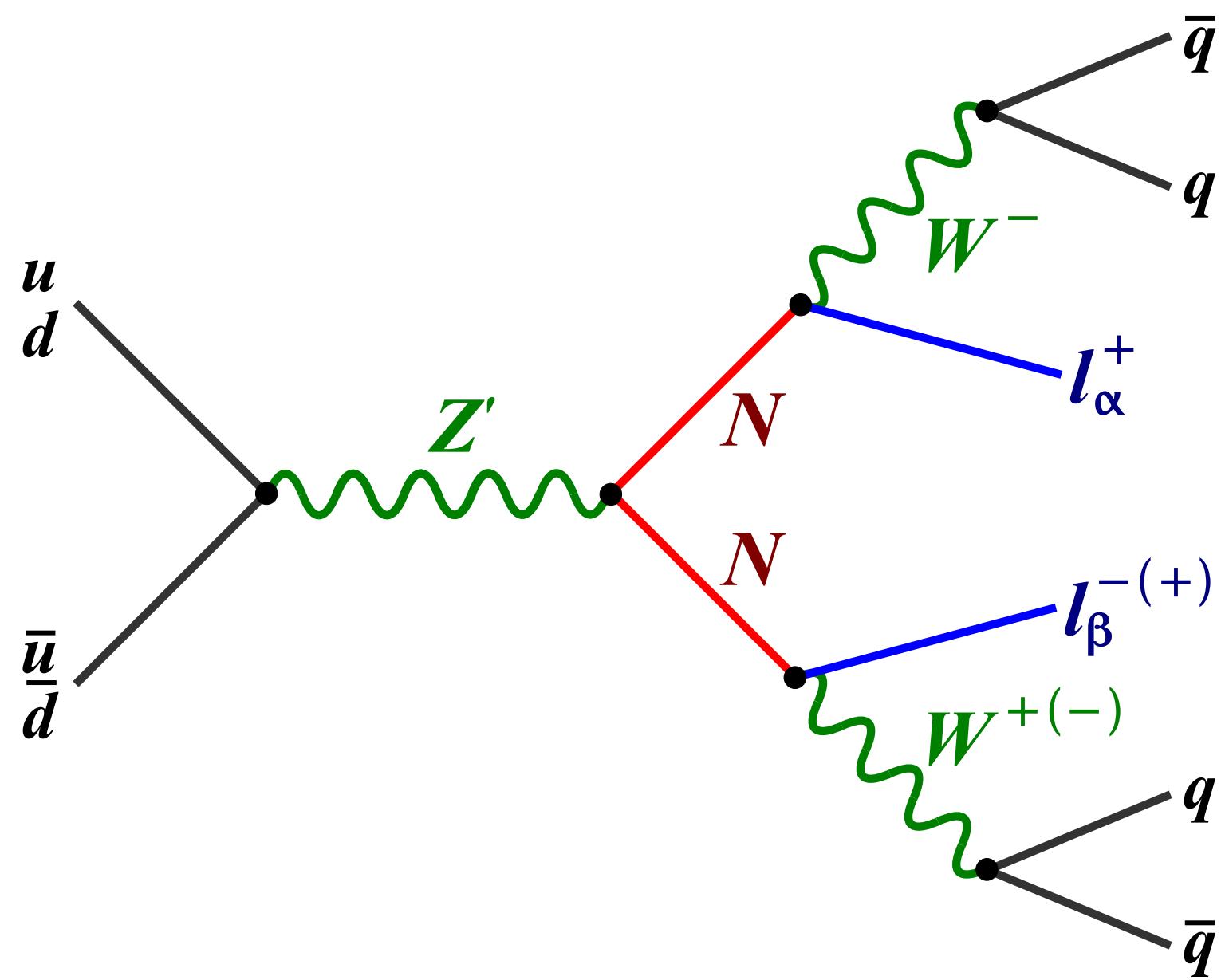
Akhmedov, Rubakov & Smirnov
Phys.Rev.Lett. 81 1359-1362
(1998) Asaka & Shaposhnikov
Phys.Lett. B620 17-26 (2005)
Asaka, Eijima & Ishida
JHEP 1104 011(2011)



Resonant Leptogenesis

Pilaftsis & Underwood *Nucl.Phys. B* 692 303-345(2004) Abada, Aissaoui, Losada *Nucl.Phys. B* 728 55-66 (2005)

- RHNs decay width similar to their mass differences. Mass range $\sim \text{TeV}$
- RHN masses explained by additional $U(1)_{B-L}$ symmetry and can be sufficiently long-lived \rightarrow displaced-vertex signature searched for at LHC, MATHUSLA or SHiP.



Deppisch, Dev & Pilaftsis *New J.Phys.* 17 no.7, 075019 (2015)
 Helo, Kovalenko & Hirsch *Phys.Rev. D* 89 073005 (2014)
 Gago, Hernández, Jones-Pérez, Losada & Briceño
Nucl.Part.Phys.Proc. 273-275 2693-2695 (2016)
 Antusch, Cazzato & Fischer *JHEP* 1612 007 (2016)

Deppisch, Dev, Pilaftsis, New J.Phys. 17 (2015) no.7, 075019

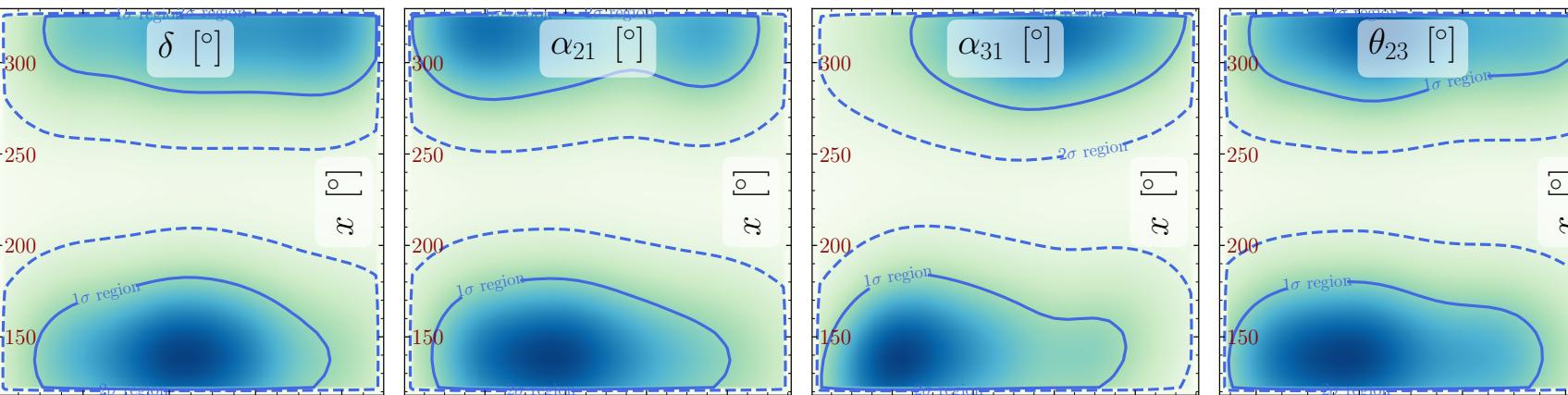
Resonant Leptogenesis in the Neutrino Option

- Assume Higgs potential vanishes at M
- Integrate out TeV RHN and RG evolve: Higgs potential produced for $M \sim 10^3$ TeV

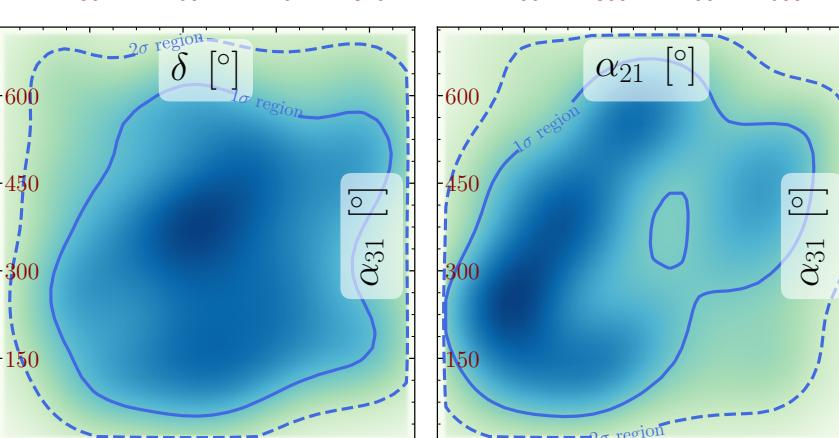
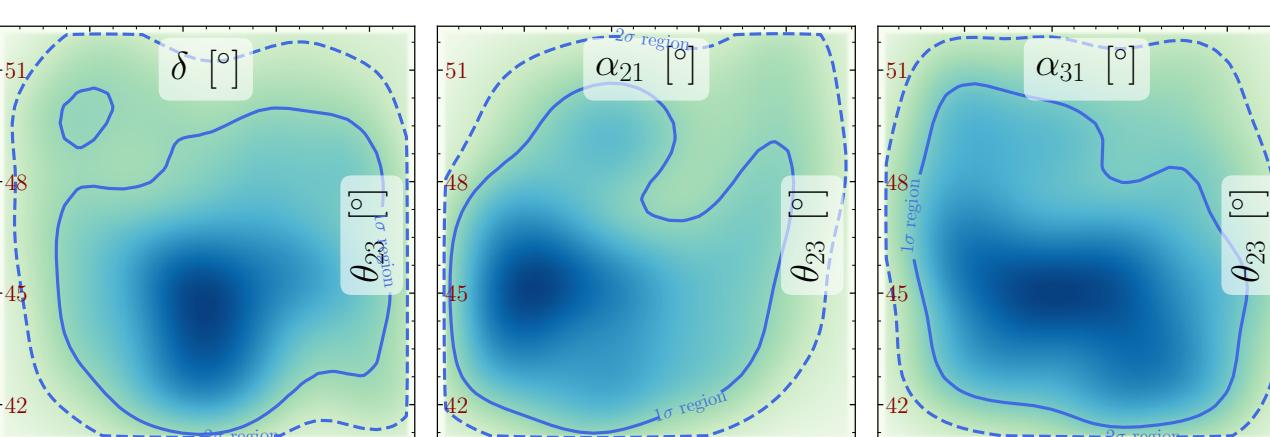
Brdar, Hemboldt, Iwamoto, Schmitz *Phys.Rev. D100 075029 (2019)*
 Brivio, Moffat, Pascoli, Petcov, Turner *JHEP 1910 059 (2019)*

**Brivio et al
1905.12642**

Normal Ordering

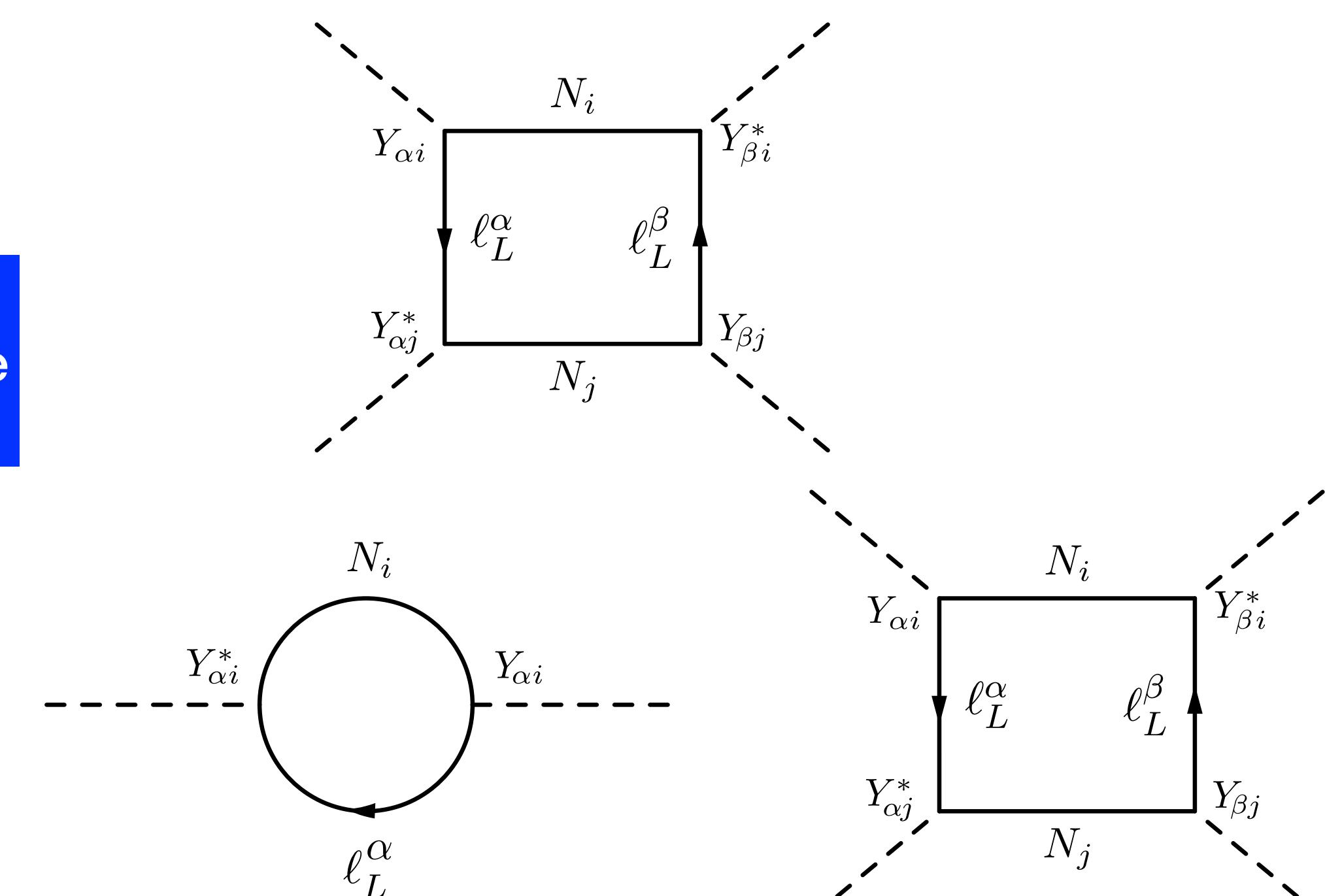


slight preference for atmospheric angle to be in upper octant.



$$\frac{\Delta M}{M} \sim 10^{-8}$$

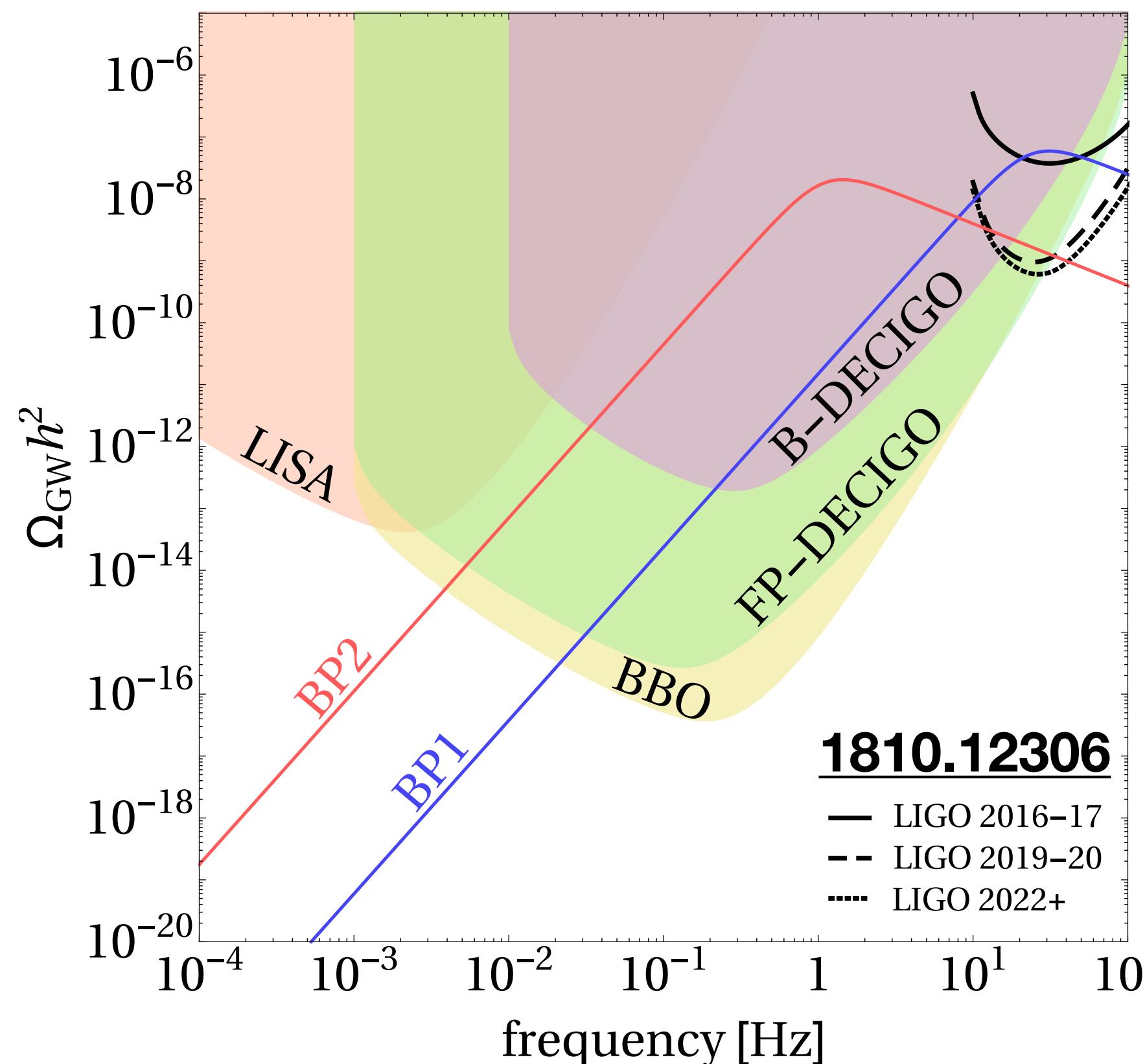
$$\overline{M} = 1.2 \times 10^6 \text{ GeV}$$



Scale invariance broken at quantum level

Resonant Leptogenesis in the Neutrino Option

- UV-completion of Neutrino Option (Brdar, Emonds, Helmboldt, Lindner) minimal renormalisable model based on classical scale invariance
- New scalar breaks scale-invariance → generates mass for RHNs and strong first order phase transition



Brdar, Emonds, Helmboldt, Lindner
Phys.Rev. D99 (2019) no.5, 055014

Mass RHN

$\mathcal{O}(10^{12})$ GeV

Fukugida & Yanagida *Phys.Lett. B17* 45-47 (1986) Buchmuller, Di Bari & Plumacher *New J.Phys.* 6 105 (2004) Barbieri, Creminelli, Strumia & Ternadic *Nucl.Phys. B575* 61-77 (2000)

$\mathcal{O}(10^6)$ GeV

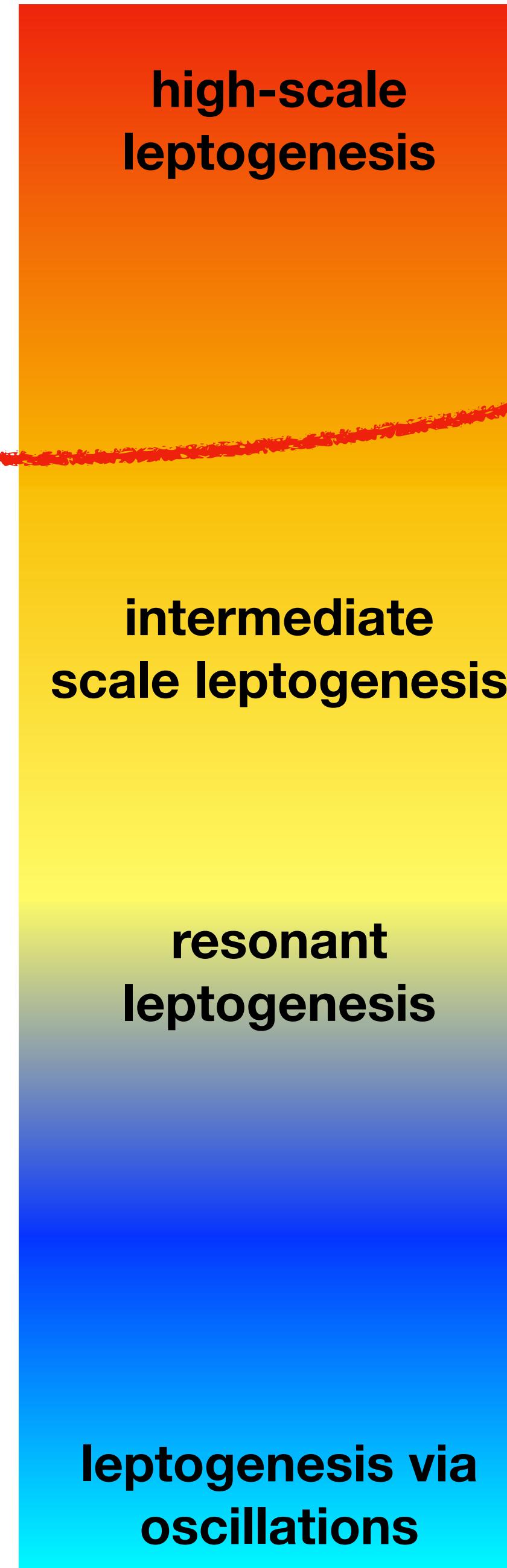
Racker, Rius & Pena *JCAP 1207* 030 (2013) Moffat, Petcov, Pascoli, Schulz & Turner *Phys.Rev. D98 no. 1*, 015036 (2018)

$\mathcal{O}(10^3)$ GeV

Pilaftis & Underwood *Nucl.Phys. B692* 303-345 (2004) Abada, Aissaoui, Losada *Nucl.Phys. B728* 55-66 (2005)

$\mathcal{O}(1)$ GeV

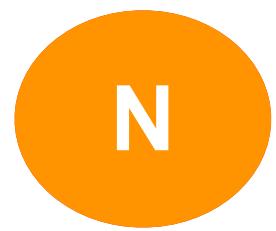
Akhmedov, Rubakov & Smirnov *Phys.Rev.Lett.* 81 1359-1362 (1998) Asaka & Shaposhnikov *Phys.Lett. B620* 17-26 (2005) Asaka, Eijima & Ishida *JHEP 1104* 011(2011)



Difficult to test as RHNs very heavy however gravitational waves offer an additional telescope on high-scale leptogenesis

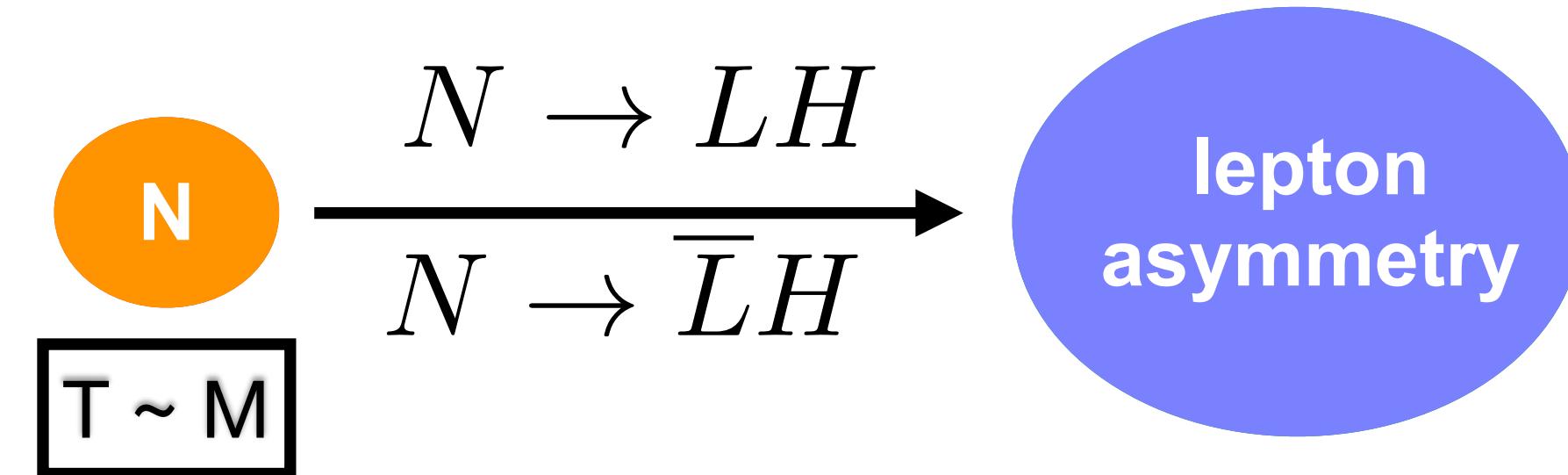
Thermal leptogenesis

Fukugida, Yanagida



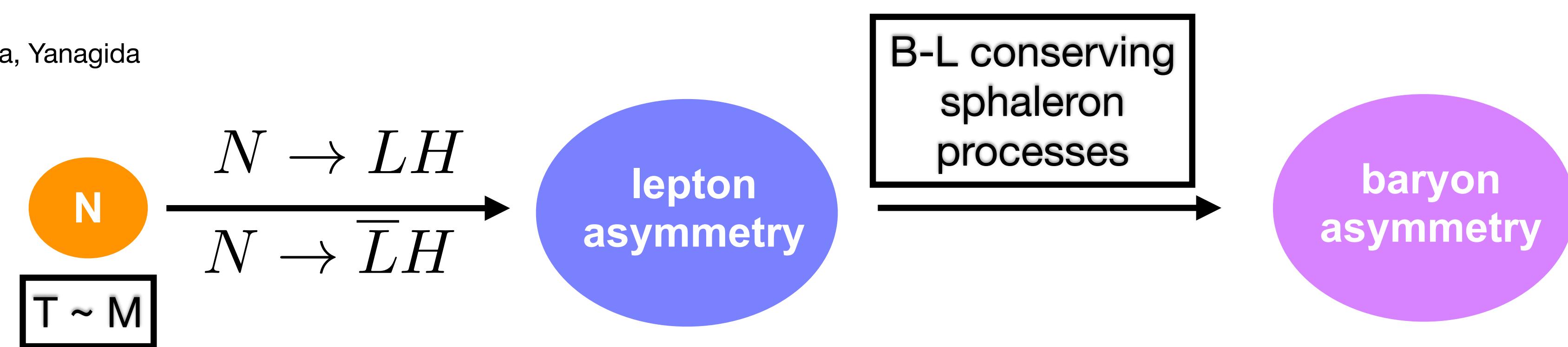
Thermal leptogenesis

Fukugida, Yanagida



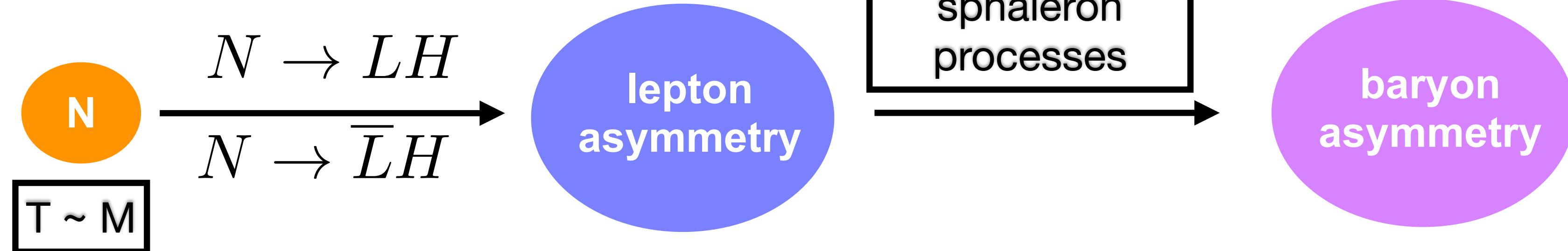
Thermal leptogenesis

Fukugida, Yanagida



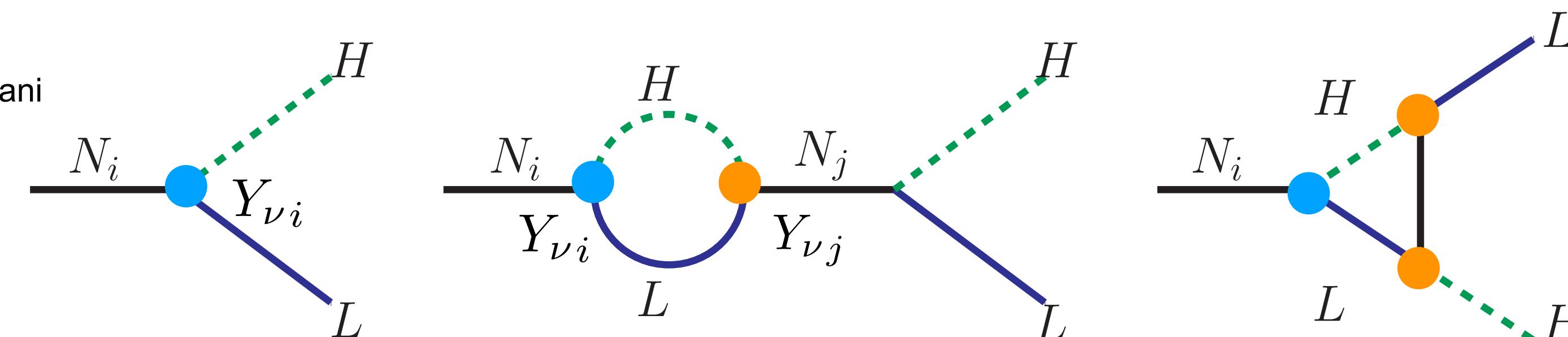
Thermal leptogenesis

Fukugida, Yanagida



**Decay asymmetry from interference between tree
and loop level diagrams**

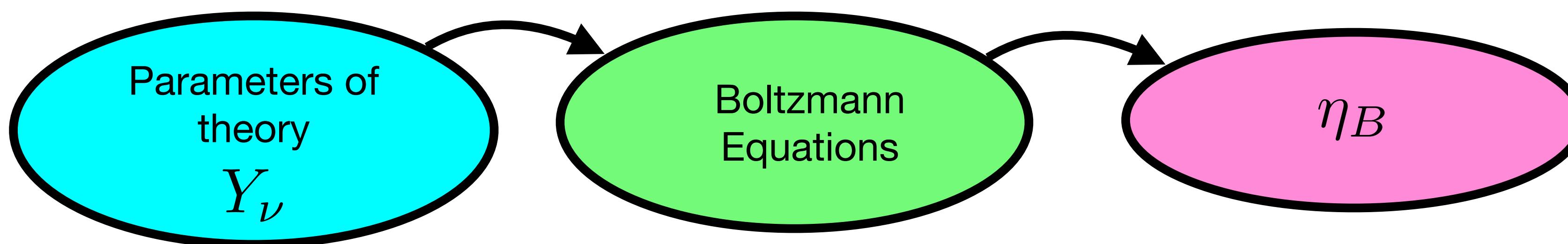
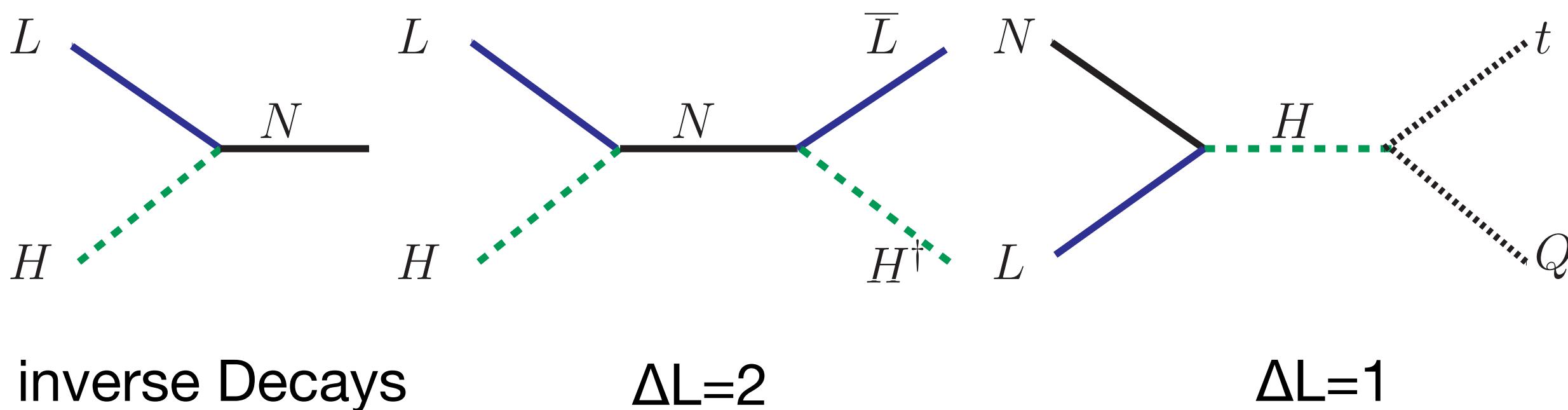
Covi, Roulet, Vissani



$$\epsilon_i = \frac{\Gamma_i - \overline{\Gamma}_i}{\Gamma_i + \overline{\Gamma}_i}$$

Thermal leptogenesis

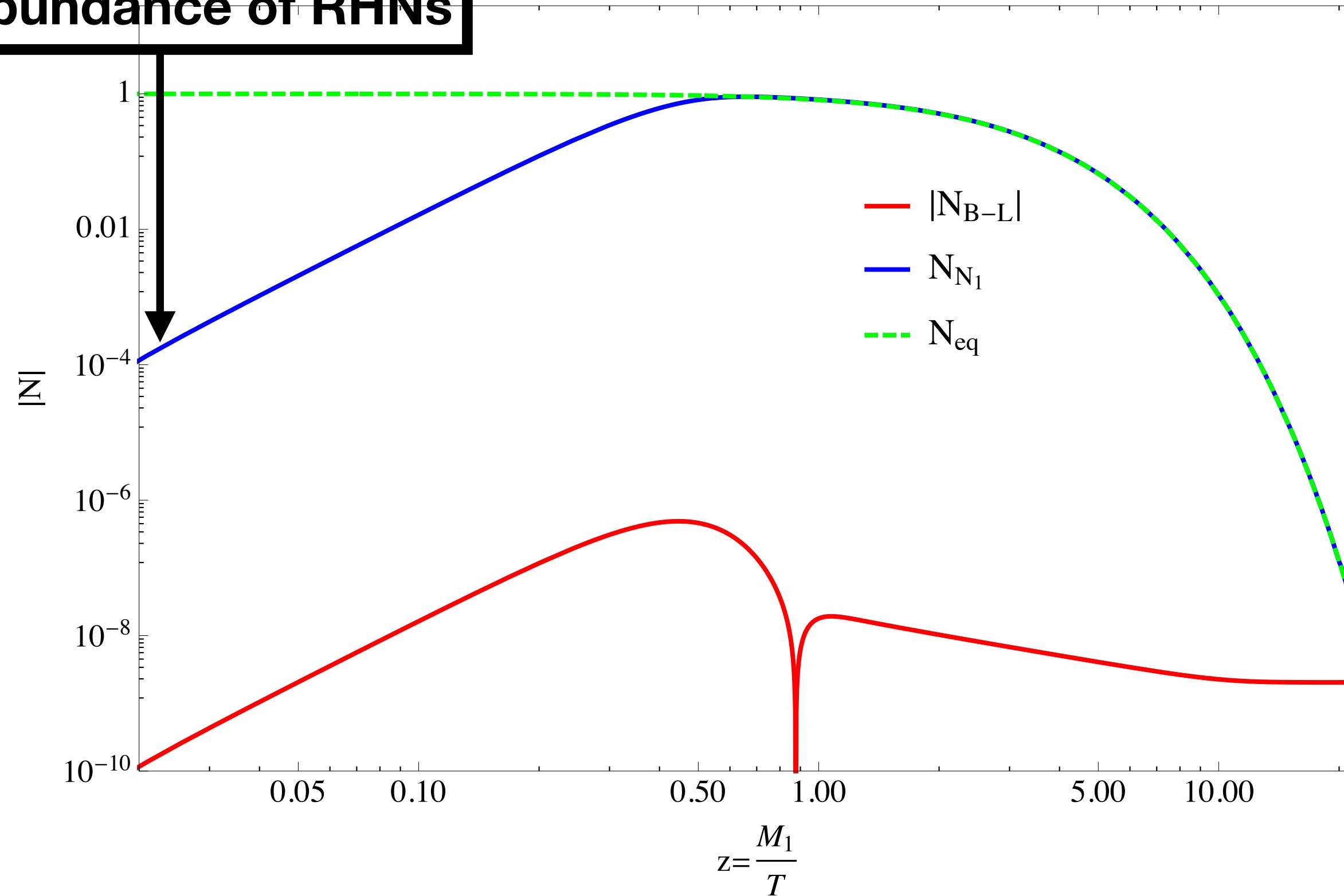
Washout and scattering processes

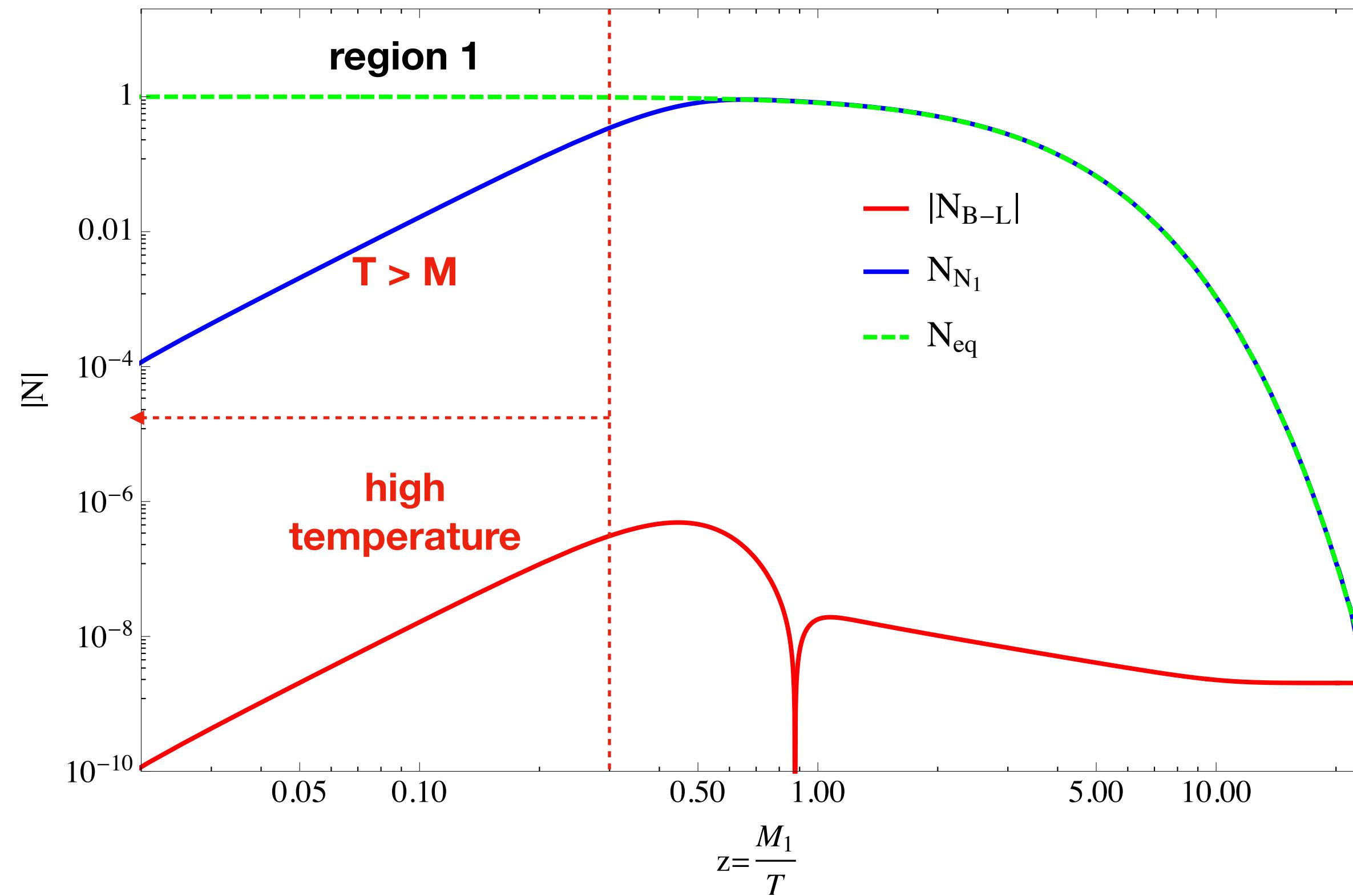


$$\frac{dn_{N_i}}{dz} = - D_i (n_{N_i} - n_{N_i}^{\text{eq}}),$$

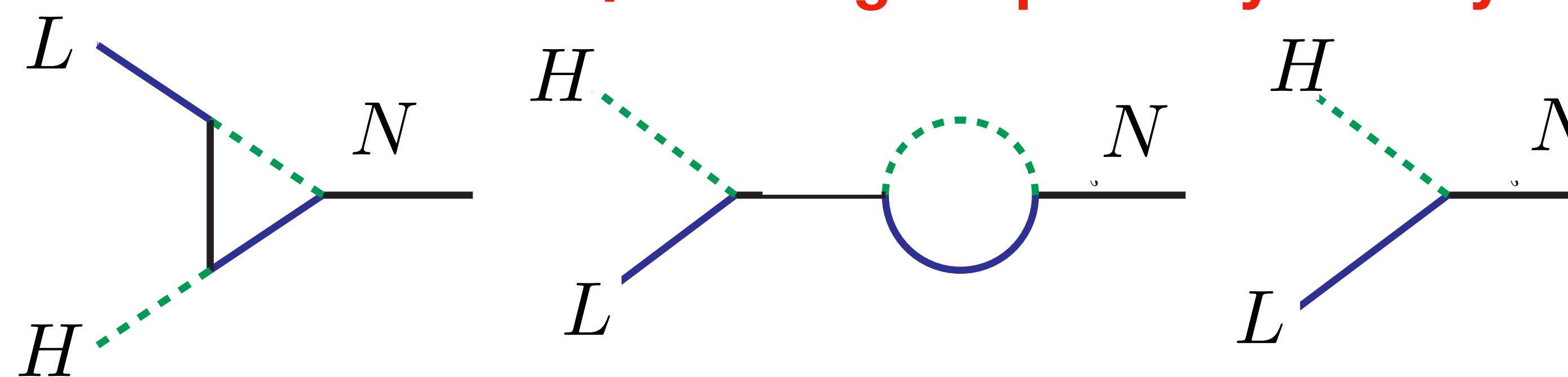
$$\frac{dn_{B-L}}{dz} = \sum_{i=1}^3 \left(\epsilon^{(i)} D_i (n_{N_i} - n_{N_i}^{\text{eq}}) - W_i n_{B-L} \right).$$

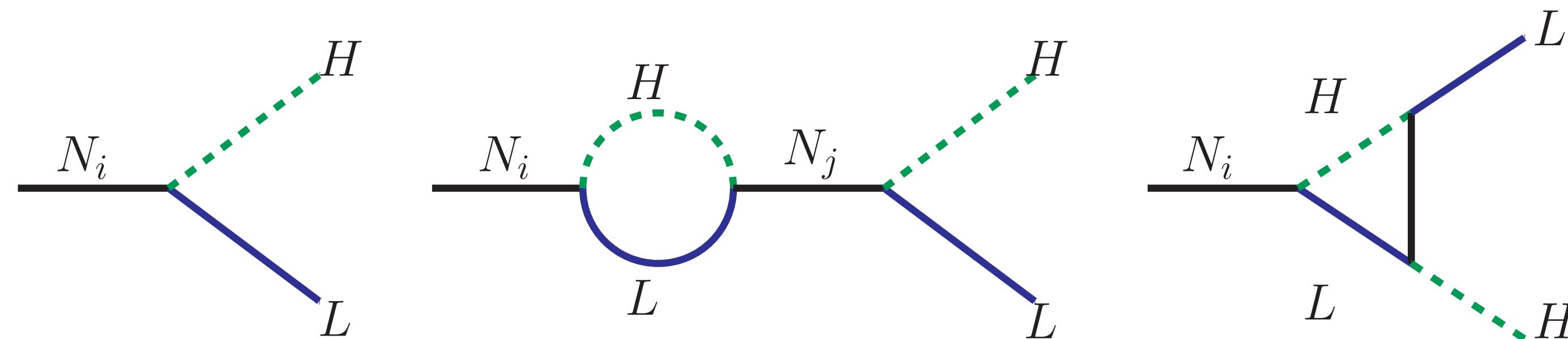
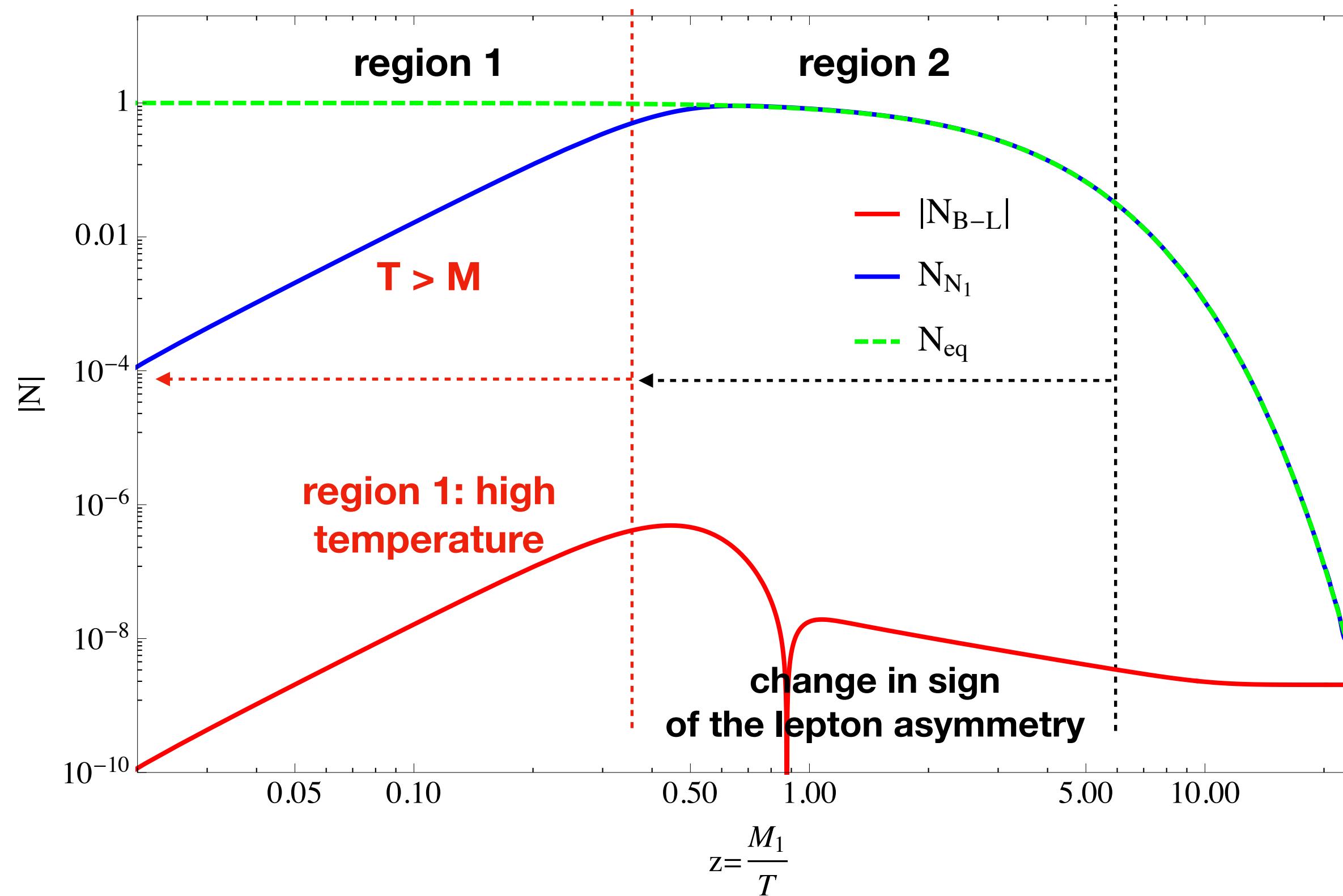
**assume zero initial
abundance of RHNs**

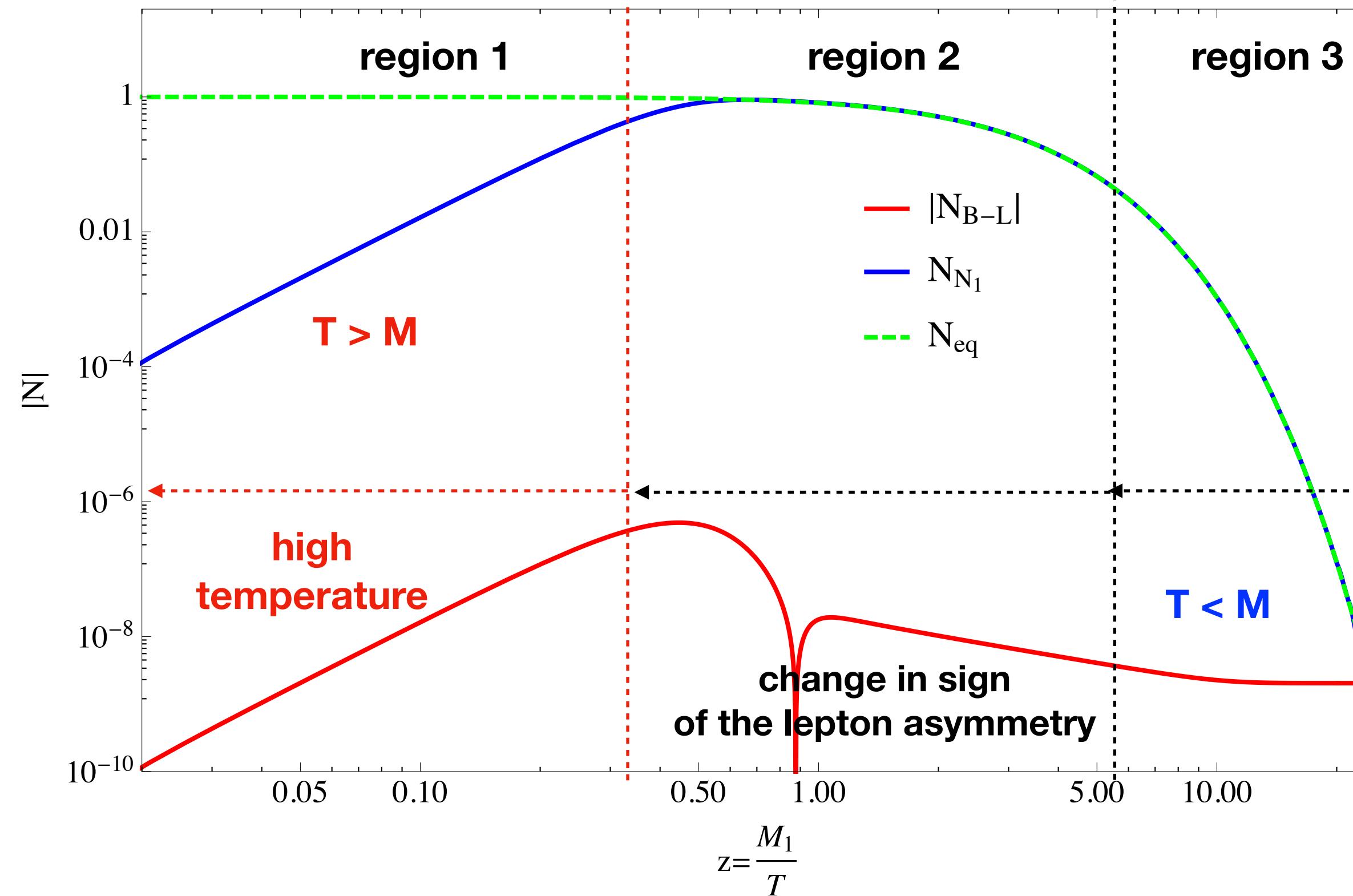




Region 1: leptons and Higgs have enough energy to inverse decay creating a lepton asymmetry







Region 3: At $T < M$, RHN abundance is depleted. Lepton asymmetry freezes out.

Primordial Black holes induced leptogenesis

Work in collaboration with **Yuber Perez Gonzalez**: [2010.03565](#)

Astrophysical BHs require $M > 3M_{\odot}$

For smaller BH mass (between Planck and solar mass scale) require large perturbations in the early Universe : **bubble collision, collapse of density perturbations...**

Carr et al, 0912.5297

$r_S \sim \lambda_C \longrightarrow$ PBHs evaporate by emitting particles

Hawking, 1975

$$\frac{dM}{dt} = - \sum_a \frac{g_a}{2\pi^2} \int_0^\infty \frac{\sigma_{\text{abs}}^{s_a}(GMp) p^3 dp}{\exp[E_a(p)/T_{\text{BH}}] - (-1)^{2s_a}}$$

$$T_{\text{BH}} = \frac{1}{8\pi GM} \approx 1.06 \left(\frac{10^{13} \text{ g}}{M} \right) \text{ GeV}.$$

PBHs are totally indiscriminate in their particle production: just need T_{BH} to be close to particle mass