A path to break degeneracies among New Physics patterns in $b \rightarrow s \ell \ell$

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Motivation and goals

My starting point:

The results on global fits to 254 observables presented in Bernat's talk

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Problem to face in this talk:

Only by adding more and more statistics blindly in some observables of the global analyses it will be very difficult to disentangle **the scenario realized in Nature among the preferred ones.** Our goal:

To identify the **key observables** that can help to truly change the present paradigm guiding global fits and breaking degeneracies.

... and learn also about hadronic physics (if not marginal).

$b \rightarrow s \ell \ell$ in the weak effective theory

► Effective Hamiltonian at scale m_b : $\mathcal{H}_{eff}^{bs\ell\ell} = \mathcal{H}_{eff, SM}^{bs\ell\ell} + \mathcal{H}_{eff, NP}^{bs\ell\ell}$

$$\mathcal{H}_{\text{eff, NP}}^{bs\ell\ell} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \mathcal{C}_i \mathcal{O}_i + \text{h.c.}$$

We also include a small λ_u contribution.

► From the set of operators (ℓ = e, µ)

$$\begin{split} O_{7}^{bs} &= \frac{e}{16\pi^{2}} m_{b} (\bar{s}\sigma_{\mu\nu}P_{R}b)F^{\mu\nu} , \qquad O_{7}^{\prime bs} = \frac{e}{16\pi^{2}} m_{b} (\bar{s}\sigma_{\mu\nu}P_{L}b)F^{\mu\nu} , \\ O_{9}^{bs\ell\ell} &= \frac{e^{2}}{16\pi^{2}} (\bar{s}\gamma_{\mu}P_{L}b)(\bar{\ell}\gamma^{\mu}\ell) , \qquad O_{9}^{\prime bs\ell\ell} = \frac{e^{2}}{16\pi^{2}} (\bar{s}\gamma_{\mu}P_{R}b)(\bar{\ell}\gamma^{\mu}\ell) , \\ O_{10}^{bs\ell\ell} &= \frac{e^{2}}{16\pi^{2}} (\bar{s}\gamma_{\mu}P_{L}b)(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell) , \quad O_{10}^{\prime bs\ell\ell} = \frac{e^{2}}{16\pi^{2}} (\bar{s}\gamma_{\mu}P_{R}b)(\bar{\ell}\gamma^{\mu}\gamma_{5}\ell) , \end{split}$$

O7,7' (even if constrained by radiative decays important to let them float)

1 $\mathcal{B}_{B_{s} \to \mu^+ \mu^-}$ exhibits a small (but persistent) deviation from the SM. It requires $C_{10\mu}^{\rm NP}$ positive (small) or $C_{10\mu}^{\rm NP}$ negative or both or a scalar contribution.

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For a long time there was a discussion among two scenarios:

a $C_{9\mu}$ large (abs.val.) and negative that can explain P'_5 and all the anomalies $(\mathcal{B}_{B_5 \rightarrow \mu^+ \mu^-}$ requires a scalar contribution).**Caveat**: difficult for model-building. b $C_{9\mu} = -C_{10\mu}$ small that can explain some anomalies but **totally fails** with P'_5 $(\mathcal{B}_{B_5 \rightarrow \mu^+ \mu^-}$ is explained with $C_{10\mu}$). Model-building welcome.

.... now this discussion is superseded (we will see why)

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The global fits to 254 observables presented in Bernat's talk \Rightarrow

Two main conclusions:

- ► Iteration after iteration the Pull_{SM} of most hypotheseshave been **confirmed or increased**.
- The preferred scenarios in terms of Pull_{SM}:
 - All contain the coefficient of the semileptonic operator $O_{o}^{bs\ell\ell}$
 - They remain packed within a narrow range

1D hypothesis:

For the complete fit, the $\mathsf{Pull}_{\mathrm{SM}}$:

- $C_{9\mu}^{\text{NP}}$ has been 4.5 [2016], 5.8 [2018], 5.6 [2019], 6.3 [2020], 7.0 [2021].
- $C_{9\mu}^{\rm NP} = -C_{10\mu}^{\rm NP}$: 4.2 [2016], 5.3 [2018], 5.2 [2019], 5.8 [2020], 6.5 [2021].

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	$\mathcal{C}_7^{\mathrm{NP}}$	$\mathcal{C}_{9\mu}^{\mathrm{NP}}$	$\mathcal{C}^{\mathrm{NP}}_{10\mu}$	$\mathcal{C}_{7'}$	$C_{9'\mu}$	$C_{10'\mu}$
Bfp	+0.01	-1.21	+0.15	+0.01	+0.37	-0.21
1σ	[-0.02, +0.04]	[-1.38, -1.01]	[+0.00, +0.34]	[-0.02, +0.03]	[-0.12, +0.80]	[-0.42, +0.02]
2σ	[-0.04, +0.06]	[-1.52, -0.83]	[-0.11, +0.49]	[-0.03, +0.05]	[-0.51, +1.12]	[-0.60, +0.23]

Or for the 6D fits:

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Some of the most relevant 2D hypotheses:

Hyp V*: ($C_{9\mu} = -1.15$, $C_{9'\mu} = -C_{10'\mu} = +0.17$) are **5.9** [2019], **6.6** [2020], **7.1** [2021]

Scn 8: $(C_{9\mu}^{V} = -C_{10\mu}^{V} = -0.34, C_{9}^{U} = -0.82)$ are **5.7** [2019], **6.5** [2020], **7.2** [2021]

Notice that RHCs are only $\mathcal{O}(15\%)$ w.r.t. the dominant $\mathcal{C}_{9\mu}$

This brings three main intertwined questions:

- Q1. What is **the reason** of this degeneracy?
- Q2. Can we **break** this degeneracy?

Q3. Is the degeneracy **linked** to the NP structure of $C_{9\mu}$?

Q1: Reason of the degeneracy

The LFUV central observables:

..., R_{κ} and $R_{\kappa*}$ (to a lesser extent) have the structure in the bin [1.1,6] GeV²

$$\begin{aligned} R_{\mathcal{K}} &= \left[1 - \mathbf{0.26}(\mathcal{C}_{10\mu} + \mathcal{C}_{10'\mu}) + \mathbf{0.23}(\mathcal{C}_{9\mu} + \mathcal{C}_{9'\mu}) + \mathbf{S}_{\mu}^{\mathcal{K}}\right] / \\ & \left[1 - \mathbf{0.26}(\mathcal{C}_{10e} + \mathcal{C}_{10'e}) + \mathbf{0.23}(\mathcal{C}_{9e} + \mathcal{C}_{9'e}) + \mathbf{S}_{e}^{\mathcal{K}}\right], \end{aligned}$$

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To find an observable with a different structure is required....

Q3: The structure of $C_{9\mu}$

Let's focus on the different contributions of the semileptonic coefficient $C_{9\mu}$

$$\mathcal{C}_{9\mu}^{\mathrm{eff}} \to \mathcal{C}_{9\mu}^{\mathrm{eff}} = \mathcal{C}_{9\mu\,\mathrm{pert}}^{\mathrm{SM}} + \frac{\mathcal{C}_{9\mu}^{\mathrm{NP}}}{\mathcal{C}_{9\mu j}} + \mathcal{C}_{9\mu j}^{\mathrm{cc}\bar{c}\,B \to K^*}$$

- $\blacktriangleright \ {\cal C}_{9\mu\,{\rm pert}}^{\rm SM}$ stands for the perturbative SM contribution

where in our conservative parametrization:

 $C_{9j}^{c\bar{c}B \to K^*} = s_j C_{9j \text{ KMPW}}^{c\bar{c}B \to K^*}$ with s_i nuisance parameters from -1 to 1

Other estimates of long-distance charm would require a separated analysis.

Q2 and Q3: The structure of $C_{9\mu}$

In [Algueró, Capdevila, Descotes-Genon, Masjuan, JM, PRD'19, 1809.08447] it was proposed: to remove hypothesis that NP is purely LFUV

$$egin{array}{rcl} C_{ie}^{\mathrm{NP}} &=& egin{array}{cc} C_{i}^{\mathrm{NP}} &=& egin{array}{cc} C_{i\mu}^{\mathrm{NP}\,\mathrm{V}} &=& egin{array}{cc} C_{i\mu}^{\mathrm{NP}\,\mathrm{V}} + egin{array}{cc} C_{i\mu}^{\mathrm{U}} \end{array} \end{array}$$

• Common contribution C_i^{U} to all charged leptons.

In particular (i = 9):

• $C_{9\mu}^{\rm NP}$ is then splited in two pieces: $C_{9\mu}^{\rm NP\,V}$ and $C_{9\mu}^{\rm U}$

Can we disentangle/measure in an efficient way these two different pieces?

$\mathcal{C}_{9\mu}^{\rm V}$ Lepton Flavour Universal Violating piece

 Q^2 and Q^3 : $\mathcal{C}_{9\mu}^{V}$

Solution: We need an LFUV $C_{9\mu}^V$ -dominated observable.

This is Q_5 : \rightarrow excellent disentangling properties like R_K (contrary to R_{K^*})

$$\begin{array}{lll} Q_5 & = & Q_5^{\rm SM} - \textbf{0.25} \mathcal{C}_{9\mu}^{V} + \textbf{0.20} \mathcal{C}_{10\mu}^{V} - 0.02 \mathcal{C}_{10\mu}^{V} \\ & & -0.04 \mathcal{C}_{9\mu}^{V} + 0.03 \mathcal{C}_{9\mu}^{V}^{-2} - 0.03 \mathcal{C}_{9\mu}^{V} \mathcal{C}_{9'\mu}^{Q} + \mathcal{R}_{\mathcal{C}} \end{array}$$

▶ $C_{i\mu}^{V} = C_{i\mu} - C_{ie}$, R_{C} is marginal, redefine $\tilde{C}_{9\mu}^{V} = C_{9\mu}^{V} - 0.8C_{10\mu}^{V}$ (absorb RHC).

where $\tilde{\mathcal{C}}_{\mathbf{9}\mu}^{\mathrm{V}}$ is obtained using

$$Q_5\simeq Q_5^{
m SM}-0.25 {\tilde {\cal C}}_{9\mu}^{
m V}~~and~~Q_5^{
m SM}=-0.0074\pm 0.0007$$

But in global fits strong correlation of sign of $C_{10'}\sim -0.2$ and $C_{9\mu}\sim -1$ Lower bound on the absolute value of $|\mathcal{C}_{9\mu}^V|$

$$|\mathcal{C}_{9\mu}^{\rm NP\,V}|\simeq |\tilde{\mathcal{C}}_{9\mu}^{\rm V}+0.8\mathcal{C}_{10'\mu}^{\rm V}|\geq |\tilde{\mathcal{C}}_{9\mu}^{\rm V}|$$

Q₅ classifies scenarios in two groups:

- ▶ Q_5 large \rightarrow all scenarios with large $C_{9\mu}^V$ (negligible C_9^U)
- Q₅ small → scenarios with small C^V_{9µ} and large C^U₉.



All top scenarios $C_{9\mu}^V \sim -1$

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Most preferred scenarios predicts R_{k} far away from one.

One of the most interesting scenarios (8) predicts Q_5 quite close to $0 \Rightarrow P_5'^{\mu} \sim P_5'^{e} \neq SM$.

 $R_K \sim 0.8$ and $Q_5 \sim 0.1$ with tiny experimental uncertainties would be great news!!!



 $\mathcal{C}_{9\mu}^{\text{V}}$ is an undoubtedly signal of New Physics

$\mathcal{C}_9^{\textit{U}}$ Lepton Flavour Universal piece

Q3: Measuring the whole $C_{9\mu} = C_{9\mu}^{V} + C_{9}^{U}$ to get C_{9}^{U}

The observable that best measures the whole NP piece of the semileptonic coefficient $C_{9\mu}$ is $P'_{5\mu}$

$$\begin{array}{ll} P_{5\mu}^{\prime\,[1.1,6]} &=& P_{5\mu}^{\prime\,\mathrm{SM}} - \textbf{0.25}\mathcal{C}_{9\mu} + \textbf{0.20}\mathcal{C}_{10^{\prime}\mu} - \textbf{0.02}\mathcal{C}_{10\mu} \\ && -0.04\mathcal{C}_{9^{\prime}\mu} + \textbf{0.03}\mathcal{C}_{9\mu}^2 - \textbf{0.03}\mathcal{C}_{9\mu}\mathcal{C}_{9^{\prime}\mu} \end{array}$$

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Approximating $P_5^{\prime\mu} \simeq P_5^{\prime\mu\,\rm SM} - 0.25\tilde{\mathcal{C}}_{9\mu}$ (with $\tilde{\mathcal{C}}_{9\mu} = \mathcal{C}_{9\mu} - 0.8\mathcal{C}_{10'\mu}$)

J. Matias (UAB)

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A C_{9µ} of O(−1) is required to explain P'^[1,1,6]_{5µ} and C_{10µ} is marginal in P'₅
 C_{9µ} = −C_{10µ} fails ⇐ it forces a too small C_{9µ} due to B_s → µµ constrain on C_{10µ}.

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• A $C_{9\mu}$ of $\mathcal{O}(-1)$ is required to explain $P_{5\mu}^{\prime [1.1,6]}$ and $C_{10\mu}$ is marginal in P_5^{\prime}

► $C_{9\mu} = -C_{10\mu}$ fails \Leftarrow it forces a too small $C_{9\mu}$ due to $B_s \to \mu\mu$ constrain on $C_{10\mu}$.

a) If data only from Q_5 and $P_5'^{\mu}$ we obtain a lower bound on the universal piece:

$$|\mathcal{C}_9^{\textit{U}}| \simeq |\tilde{\mathcal{C}}_{9\mu} - \tilde{\mathcal{C}}_{9\mu}^{\textit{V}} + 0.8(\mathcal{C}_{10'\mu} - \mathcal{C}_{10'\mu}^{\textit{V}})| \geq |\tilde{\mathcal{C}}_{9\mu} - \tilde{\mathcal{C}}_{9\mu}^{\textit{V}}|$$

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b) if also data on $P_5^{\prime e}$ is provided by LHCb then

$$|\mathcal{C}_9^{\textit{U}}| \simeq |\tilde{\mathcal{C}}_{9e} + 0.8 \mathcal{C}_{10'e}| \geq |\tilde{\mathcal{C}}_{9e}|$$

 $P_1(Q_1)$ can also help in the disentangling strategy:

c) informing on presence of RHCs (Hyp.V) d) turning $\tilde{\mathcal{C}} \to \mathcal{C}$ & bound \to measurem.

This observable measures $\tilde{\mathcal{C}}_{9\mu}^{\rm NPV} + \tilde{\mathcal{C}}_{9}^{U}$ that in absence of RHC is $\mathcal{C}_{9\mu}^{\rm NPV} + \mathcal{C}_{9}^{U}$

An explanation of P'_5 purely in terms of hadronic SM is simply wrong ($C_{9\mu}^{\rm NP\,V} \neq 0$ for R_{κ}).



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Non-exclusive (and non-exhaustive) possibilities

- Direct NP contribution.... a Z'?
- Loop-effects with SM particles
 - τ -loops with or without connection with $R_{D(*)}$.
- Contrived hadronic contribution beyond those already included
 - Unlikely considering recent theoretical progress.
 - ▶ Nicely it can be bounded using info on Q_5 , $b \rightarrow s\tau\tau$ (or other LFU NP sources).

LFU via tau-loops: An EFT interpretation SMEFT

Connect $b \to s\ell\ell$ and $b \to c\ell\nu$ anomalies within SMEFT ($\Lambda_{NP} \gg m_{t,W,Z}$) $\mathcal{L}_{SMEET} = \mathcal{L}_{SM} + \mathcal{L}_{d > A}$ with higher-dim ops involving only SM fields

[Grzadkowski, Iskrzynski, Misiak, Rosiek ; Alonso, Grinstein, Camalich]

Two ops. with left-handed doublets

$$\mathcal{O}_{ijkl}^{(1)} = [\bar{Q}_i \gamma_\mu Q_j] [\bar{L}_k \gamma^\mu L_l] \qquad \mathcal{O}_{ijkl}^{(3)} = [\bar{Q}_i \gamma_\mu \vec{\sigma} Q_j] [\bar{L}_k \gamma^\mu \vec{\sigma} L_l]$$

▶ FCCC part of $\mathcal{O}_{2333}^{(3)}$ can describe $R_{D(*)}$ (rescaling of G_F for $b \to c\tau\nu$)

- FCNC part of $\mathcal{O}_{2333}^{(1,3)}$ with $C_{2333}^{(1)} = C_{2333}^{(3)}$ (assumed $C_{33}^{(3)} = 0$) [Capdevila, Crivellin, SDG, Hofer, Matias]
 - Avoids bounds from $B \to K^{(*)}\nu\nu$, Z decays, direct production in $\tau\tau$
 - Large NP contribution b → sττ through C^V_{9τ} = −C^V_{10τ}
 - Through radiative effects, NP contribution to C^U₀



Hyp: link between charged-neutral anomalies in SMEFT in scenario 8: $C_9^V = -C_{10}^V, C_9^U$ Huge $B_s \to \tau^+ \tau^-$ and $B \to K^{(*)} \tau^+ \tau^ \downarrow \downarrow$ off-shell photon penguin with $O_9^{\tau\tau}$ generates LFU-NP in C_9 Combining all leads to:



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Hypothesis: link between charged and neutral anomalies in SMEFT in scenario 8 Huge $B_s \rightarrow \tau^+ \tau^-$ and $B \rightarrow K^{(*)} \tau^+ \tau^ \downarrow \downarrow$ off-shell photon penguin with $\mathcal{O}_9^{\tau\tau}$ generates LFU-NP in \mathcal{C}_9

Furthermore if the link with $R_{D(*)}$ exists it should fulfill:



... also without R_D connection

Also a more general link between C_9^U and $b \to s \tau \tau$ is possible(without R_D connection)

Mixing of $\mathcal{O}_9^{\tau\tau}$ into \mathcal{O}_9^U

$$\mathcal{C}_9^{\tau\tau} \to \mathcal{C}_9^U$$
 (1)

IF assumed that LQs is the most plausible solution also $C_{10}^{\tau\tau}$ is present

•
$$C_9^{\tau\tau} = C_{10}^{\tau\tau}$$
 in case of S_2 LQ

• $C_9^{\tau\tau} = -C_{10}^{\tau\tau}$ in case of U_1 or $S_1 + S_3$ LQs

In both cases same plots of C_9^U versus $\mathcal{B}_{B_s \to \tau \tau}$ and $\mathcal{B}_{B \to K^* \tau \tau}$ holds (but not $R_{D(*)}$)

Anything else?

Once identified:

- $C_{9\mu}^V$ from Q_5
- C_9^U from $(P_5'^{\mu} Q_5 = P_5'^e)$

 Δ between C_9^U from $P_5^{\prime e}$ and $C_9^{\rm NP U}$ generated from τ -loops or other LFU NP sources:

quantitative upper bound on the existence (or not) of some marginal non-perturbative contribution.

A raw estimate of it is the difference between:

- C_9^U from global fit without the connection with R_{D,D^*}
- C_9^{NPU} from τ -loop with R_{D,D^*} link (at a scale of $\Lambda = 2 \text{ TeV}$)

With present data this amount in scenario 8 to an upper bound of $|\Delta| < 0.3$

- Compatible with zero at 1σ
- The larger the A scale the smaller the contribution
- Other LFU-NP sources may contribute and further reduce Δ

Disentangling scenarios



Conclusions

$C_{9\mu}$ decision tree



or in words...

- We pointed out that:
 - the degeneracy among dominant scenarios (Pull_{SM} > 7σ) is due to the similar weight of WC entering R_X observables.
 - $\blacktriangleright\,$ the breaking of this degeneracy requires a $\mathcal{C}_{9\mu}\text{-}\text{dominated observable}$
- Understanding the structure of NP of $C_{9\mu}^{NP} = C_{9\mu}^{V} + C_{9}^{U}$ is crucial to disentangle the two main scenarios:
 - Hypothesis V \rightarrow ($C_{9\mu}, C_{9'\mu} = -C_{10'\mu}$)
 - ▶ Scenario 8 → $(C_{9\mu}^V = -C_{10\mu}^V, C_9^U)$
- ▶ Size of $C_{9\mu}^V$ can be determined (up to a small RHC) by Q_5
- Size of C_9^U can be determined (up to a small RHC) by $P_5^{\prime e}$ (or Q_5 and $P_5^{\prime \mu}$)
 - If the link between $\frac{R_D^{(*)}}{R_D^{(*)}SM}$ and $b \to s\ell\ell$ holds a large $b \to s\tau\tau$ is generated.
 - ▶ C_9^{UNP} can be induced by τ -loops and connected with $\mathcal{B}_{B_s \to \tau^+ \tau^-}$ and $\mathcal{B}_{B \to K^* \tau^+ \tau^-}$
 - ▶ Other NP sources of LFU are possible (Z').
- ► The difference between C₉^U and C₉^{UNP} (if any) is all the possible space for a marginal extra hadronic contribution.

Questions:

- ▶ What is the expected precision achievable for Q₅ in the short and long-term?
- ▶ How large have to be $B_s \rightarrow \tau \tau$ and $B \rightarrow K^* \tau \tau$ to be measurable at LHCb? What are the short term perspectives?