Jet calculations with the Sector-improved residue subtraction scheme

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Established by the European Commission

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Sector-improved residue subtraction scheme

$$\sigma_{h_{1}h_{2}\to X} = \sum_{ij} \int_{0}^{1} \int_{0}^{1} dx_{1} dx_{2} \phi_{i,h_{1}}(x_{1},\mu_{F}^{2}) \phi_{j/h_{2}}(x_{2},\mu_{F}^{2}) \hat{\sigma}_{ij\to X}(\alpha_{s}(\mu_{R}^{2}),\mu_{R}^{2},\mu_{F}^{2})$$

$$PDFs$$

$$\hat{\sigma}_{ab\to X} = \hat{\sigma}_{ab\to X}^{(0)} + \hat{\sigma}_{ab\to X}^{(1)} + \hat{\sigma}_{ab\to X}^{(2)} + \mathcal{O}(\alpha_{s}^{3}) \quad \text{NNLO:} \quad \hat{\sigma}_{ab}^{(2)} = \hat{\sigma}_{ab}^{\text{RR}} + \hat{\sigma}_{ab}^{\text{RV}} + \hat{\sigma}_{ab}^{\text{C2}} + \hat{\sigma}_{ab}^{\text{C1}}$$

Divide and conquer IR singularities by sector decomposition:

$$\hat{\sigma}_{ab}^{RR} = \frac{1}{2\hat{s}} \int d\Phi_{n+2} \sum_{i,j} \left[\sum_{k} S_{ij,k} + \sum_{k,l} S_{i,k;j,l} \right] \left\langle \mathcal{M}_{n+2}^{(0)} \middle| \mathcal{M}_{n+2}^{(0)} \right\rangle F_{n+2}$$
Sector parameterization: $\hat{\eta}_i = \frac{1}{2} (1 - \cos \theta_{ir}) \in [0, 1]$ $\hat{\xi}_i = \frac{u_i^0}{u_{max}^0} \in [0, 1]$
Overlapping triple collinear singularities: sub-sectors
Apply Master-formula: $x^{-1-b\epsilon} = \underbrace{-1}_{b\epsilon} + \underbrace{[x^{-1-b\epsilon}]_{+}}_{reg. + sub.}$
 \Rightarrow All contributions numerical finite integrals
 \Rightarrow Numerical ϵ – poles cancellation between real&virtual
$$\left[\begin{array}{c} \text{Czakon'10} \\ \text{I} \\ \hat{\sigma}_{2} \rightarrow \xi_{2} \\ \hat{\sigma}_{2} \rightarrow \xi_$$

Applications of sector-improved residue subtraction

Top-quark pairs + NWA decay

High-precision differential predictions for top-quark pairs at the LHC Czakon, Heymes, Mitov 1511.00549

Jets

<mark>Single-jet inclusive rates with exact color at O(α</mark>s⁴) Czakon, van Hameren, Mitov, **Poncelet** 1907.12911

Photons

NNLO QCD corrections to three-photon production at the LHC Chawdhry, Czakon, Mitov, Poncelet 1911.00479

Weak-bosons/Higgs

NNLO QCD study of polarised W+W- production at the LHC Poncelet, Popescu 2102.13583

Weak-boson+jets

NNLO QCD predictions for W+c-jet production at the LHC Czakon, Mitov, Pellen, Poncelet 2011.01011 Polarised W+j production at the LHC: a study at NNLO QCD accuracy Pellen, Poncelet, Popescu 2109.14336 Angular coefficients in W+j production at the LHC with high precision Pellen, Poncelet, Popescu, Vitos 2204.12394

B-hadrons

<mark>B-hadron production in NNLO QCD: application to LHC ttbar events with leptonic decays,</mark> Czakon, Generet, Mitov and **Poncelet**, 2102.08267

Higher order corrections to spin correlations in top quark pair production at the LHC Behring, Czakon, Mitov, Papanastasiou, Poncelet 1901.05407 NNLO QCD corrections to leptonic observables in top-quark pair production and decay Czakon, Mitov, Poncelet 2008.11133

Next-to-Next-to-Leading Order Study of Three-Jet Production at the LHC Czakon, Mitov, Poncelet 2106.05331

NNLO QCD corrections to diphoton production with an additional jet at the LHC Chawdhry, Czakon, Mitov, **Poncelet** 2105.06940

Exact Top-Quark Mass Dependence in Hadronic Higgs Production Czakon, Harlander, Klappert, Niggetiedt 2105.04436

Infrared-safe flavoured anti-kT jets, Czakon, Mitov, Poncelet 2205.11879

NNLO QCD corrections to Wbb production at the LHC Hartanto, Poncelet, Popescu, Zoia 2205.01687 Flavour anti-kT algorithm applied to Wbb production at the LHC Hartanto, Poncelet, Popescu, Zoia 2209.03280

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NNLO QCD for three jets

Next-to-Next-to-Leading Order Study of Three-Jet Production at the LHC Czakon, Mitov, Poncelet 2106.05331

NNLO QCD prediction for three jet rates and three-to-two jet ratios:

$$R_{3/2}(X,\mu_R,\mu_F) = \frac{\mathrm{d}\sigma_3(\mu_R,\mu_F)/\mathrm{d}X}{\mathrm{d}\sigma_2(\mu_R,\mu_F)/\mathrm{d}X} \sim \alpha_s$$

Full colour real radiation + LC double virtual

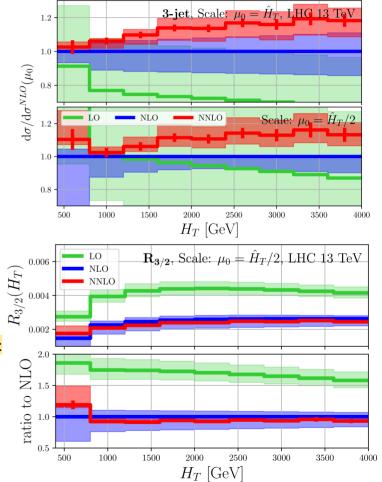
$$\mathcal{R}^{(2)}(\mu_R^2) = 2\operatorname{Re}\left[\mathcal{M}^{\dagger(0)}\mathcal{F}^{(2)}\right](\mu_R^2) + \left|\mathcal{F}^{(1)}\right|^2(\mu_R^2) \equiv \mathcal{R}^{(2)}(s_{12}) + \sum_{i=1}^4 c_i \ln^i\left(\frac{\mu_R^2}{s_{12}}\right)$$

 $\mathcal{R}^{(2)}(s_{12}) \approx \mathcal{R}^{(2)l.c.}(s_{12})$

Leading-color two-loop QCD corrections for three-jet production at hadron colliders Abreu, Cordero, Ita, Page, Sotnikov 2102.13609

Updated plots due to missing colour factor in VV: → VV sizable contribution ~10% → naive sub-leading ~1%

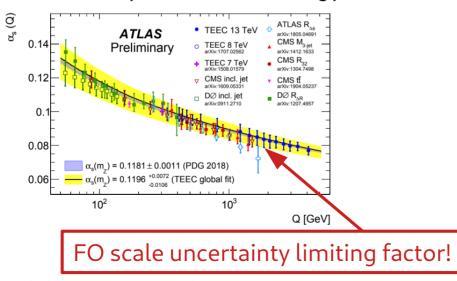
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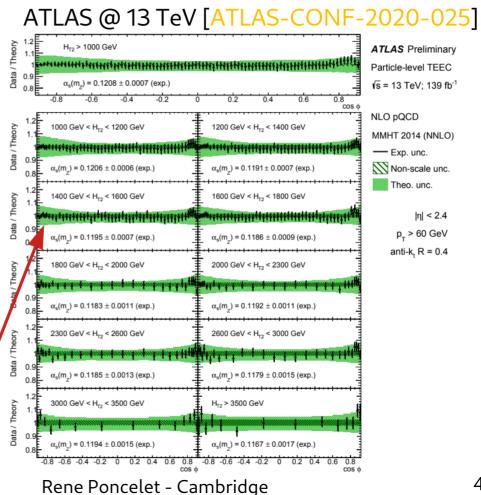
Transverse Energy-Energy Correlator @ LHC

TEEC: Transverse Energy-Energy Correlation

- $\frac{1}{\sigma} \frac{\mathrm{d}\Sigma}{\mathrm{d}\cos\phi} = \frac{1}{N} \sum_{A=1}^{N} \sum_{ij} \frac{E_{\perp,i}^{A} E_{\perp,j}^{A}}{\left(\sum_{k} E_{T,k}^{A}\right)^{2}} \delta(\cos\phi \cos\phi_{ij})$
- HT2 bins from 1000 GeV to 3500 GeV $\rightarrow \alpha_{s}$ sensitivity at different energy scales

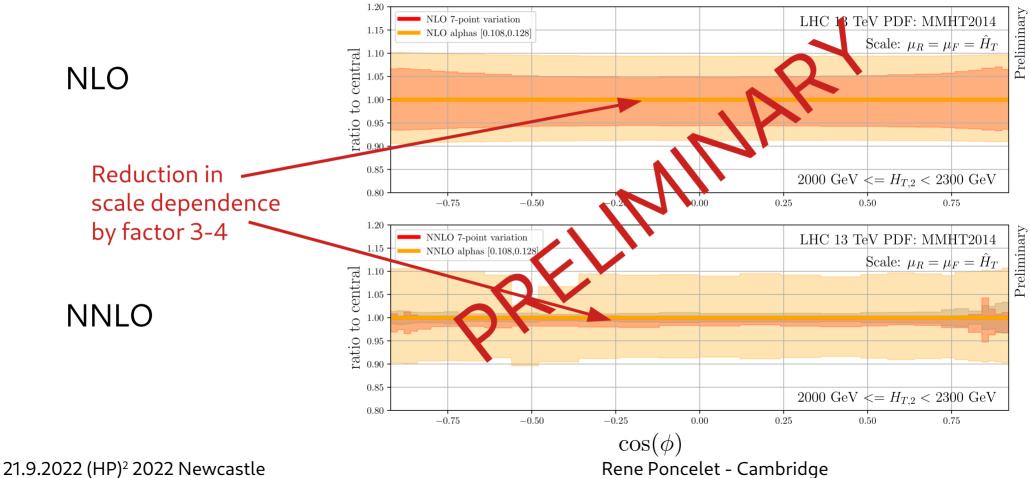


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NNLO QCD corrections to TEEC @ LHC

Massive thanks to Manuel Alvarez and Javier Llorente for computing support!



ATLAS measurement of event shapes @ 13 TeV using multi-jet events (139fb⁻¹) in HT2 bins and high pT jets (> 100 GeV): [ATLAS:2007.12600]

Jet-based
Transverse Thrust:
$$\tau_{T} = 1 - \frac{\sum_{i}^{jets} |\vec{p}_{T,i} \cdot \hat{n}|}{\sum_{i}^{jets} |\vec{p}_{T,i}|}$$
Back-to-Back
Jet-based
Thrust Minor:
$$T_{m} = \frac{\sum_{i}^{jets} |\vec{p}_{T,i} \times \hat{n}|}{\sum_{i}^{jets} |\vec{p}_{T,i}|}$$
More quantities based on eigenvalues of
(transverse) linearised sphericity tensor:

$$\mathcal{M}_{xyz} = \frac{1}{\sum_{i}^{jets} |\vec{p}_{i}|} \sum_{i}^{jets} \frac{1}{|\vec{p}_{i}|} \begin{pmatrix} p_{x,i}^{2} & p_{x,i}p_{y,i} & p_{x,i}p_{z,i} \\ p_{y,i}p_{x,i} & p_{y,i}^{2} & p_{y,i}p_{z,i} \\ p_{z,i}p_{x,i} & p_{z,i}p_{x,i} & p_{z,i}^{2} \end{pmatrix}$$

NNLO QCD corrections to event shapes

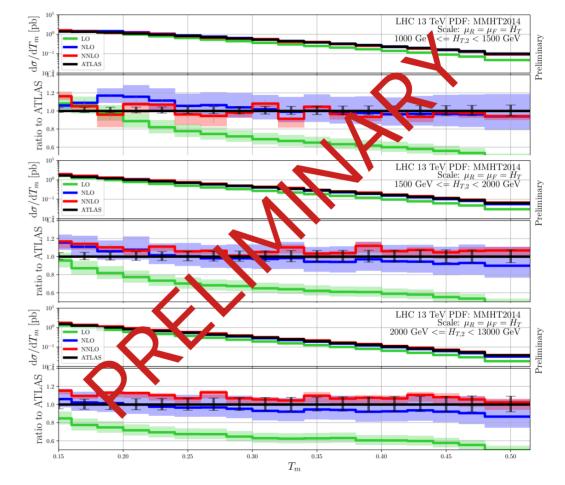
Comparison of public data from HEPdata

Example Thrust-Minor:

- Beautiful perturbative convergence
- Significant reduction of perturbative corrections

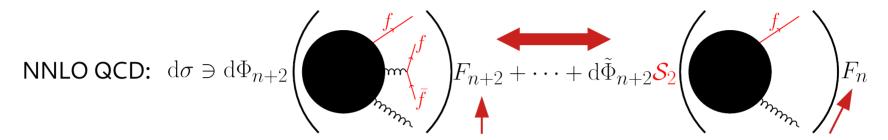
Question to be addressed:

- Resummation effects
- Non-perturbative corrections



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IR-safety of flavour observables



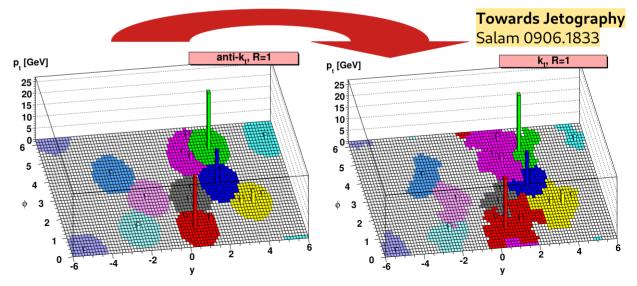
IR safety implies correlated treatment of kinematics and flavour information:

Infrared safe definition of jet flavor, Banfi, Salam, Zanderighi hep-ph/0601139

Open issues:

- → Requires flavour-information
- → (flavour-) kT and anti-kT cluster partonic jets differently

A proper comparison requires to unfold experimental data



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Old problem, new approaches

Renewed interest:

• Anti-kT + flv.-kT flavour matching:

QCD-aware partonic jet clustering for truth-jet flavourPractical Jet Flavour Through NNLOA dress of flavour to suit any jetlabelling Buckley, Pollard 1507.00508Caletti, Larkoski, Marzani, Reichelt 2205.01109Gauld, Huss, Stagnitto 2208.11138

• Fixed-order fragmentation:

B-hadron production in NNLO QCD: application to LHC ttbar events with leptonic decays,Czakon, Generet, Mitov and Poncelet, 2102.08267

A Fragmentation Approach to Jet Flavor Caletti, Larkoski, Marzani, Reichelt 2205.01117

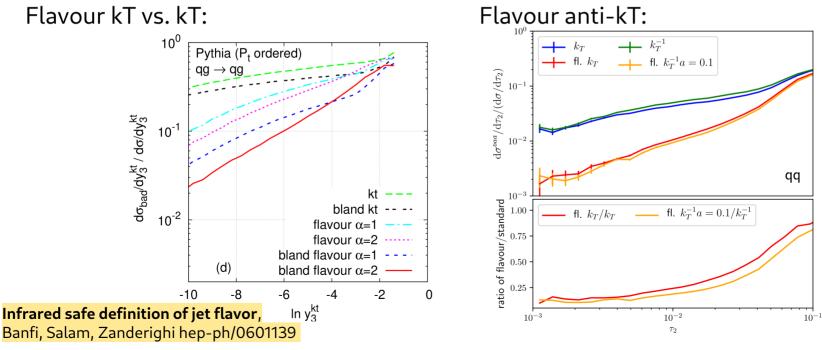
• Modified anti-kT algorithm:

Infrared-safe flavoured anti-kT jets, Czakon, Mitov, Poncelet 2205.11879

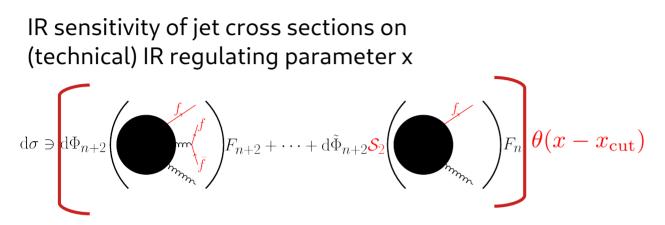
Proposed modification: A soft term designed to modify the distance of flavoured pairs. $d_{ij}^{(F)} = d_{ij} \begin{cases} S_{ij} & \text{i,j is flavoured pair} \\ 1 & \text{else} \end{cases}$ $S_{ij} = 1 - \theta(1-x) \cos\left(\frac{\pi}{2}x\right) \quad \text{with} \quad x = \frac{k_{T,i}^2 + k_{T,j}^2}{2ak_{T,\max}^2}$

Tests of IR safety with parton showers

Dress tree-level di-jet events (definite flavour structure: "qq", "qg" or "gg") with radiation and study jet flavour (q or g) as function of kinematics. In the di-jet limit the flavour needs to correspond to tree level flavours → misidentification rate needs to vanish in di-jet back-to-back limit

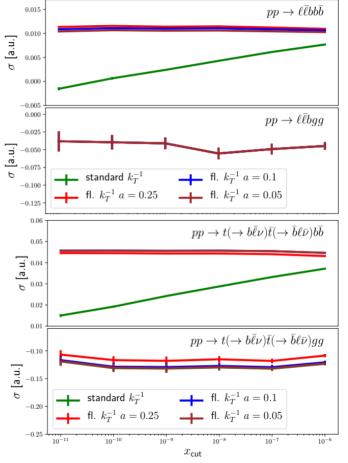


Tests of IR safety with NNLO FO computations



In the limit $x_{cut} \rightarrow 0$:

IR safe jet flavour \rightarrow no dependence on x_cut IR non-safe jet flavour \rightarrow logarithmic divergent



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Z+b-jet Phenomenology: Tunable parameter

Flavour anti-kT: a = 0.1

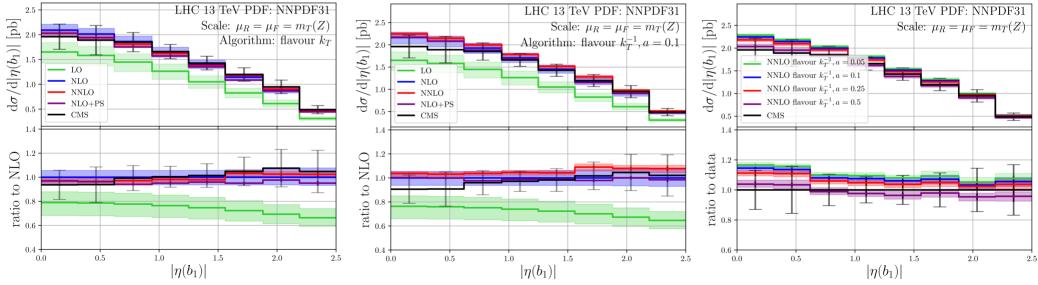
Benchmark process: $pp \rightarrow Z(ll) + b$ -jet

Tunable parameter a:

Flavour kT:

- Limit a → 0 <=> original anti-kT (IR unsafe)
- Large a <=> large modification of cluster sequence

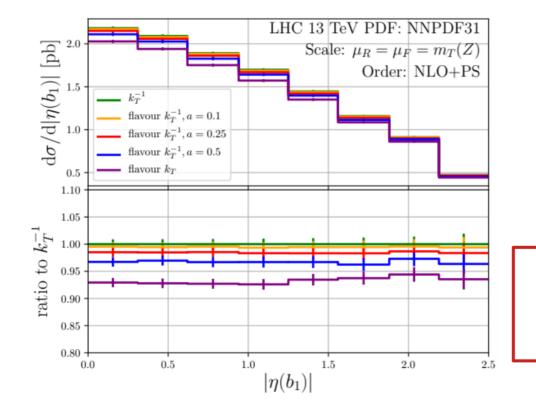
Comparison of different parameter a to data:



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Z+b-jet Phenomenology: Tunable parameter II

What happens in the presence of many flavoured partons? \rightarrow NLO PS



Tunable parameter a:

- Small a: Flavour anti-kT results are more similar to standard anti-kT
 → small unfolding factors
- Larger a: Larger modification of clustering

Good FO perturbative convergence + Small difference to standard anti-kT → a~0.1 is a good candidate

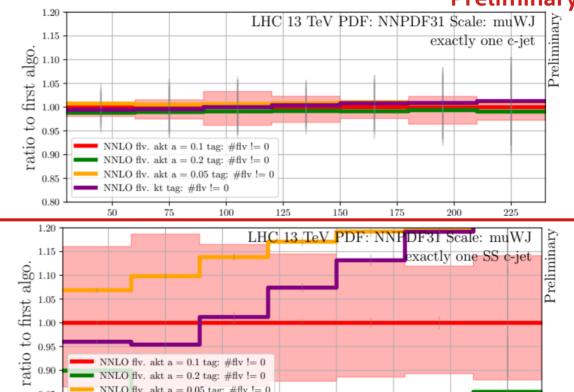
W+c-jet with flavour anti-kT

In collaboration with: Czakon, Mitov, Pellen

Preliminary

Exactly one c-jet requirement:

- Comparison of parameters a: → small dependence < 2%
- Comparison to flv kT: → small dependence @ NNLO < 2%



ONLY large effect in SS contribution

- Exactly one c-jet of SS type: Larger dependence ~15% (roughly size of NNLO scale band)
- BUT: SS contribution ~2-5%
- => OS ~0.2-0.5% dependence

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125

 $p_T(\ell)$ [GeV]

150

175

200

225

100

NNLO fly, akt a = 0.2 tag; #fly != 0NNLO fly. akt a = 0.05 tag: #fly != 0

NNLO fly. kt tag: #fly != 0

75

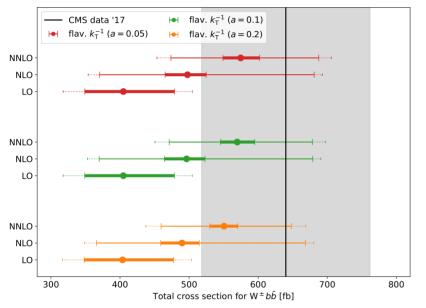
50

0.85

0.80

W+2 bjets: flavour anti-kT

Flavour anti-kT algorithm applied to Wbb production at the LHC Hartanto, Poncelet, Popescu, Zoia 2209.03280

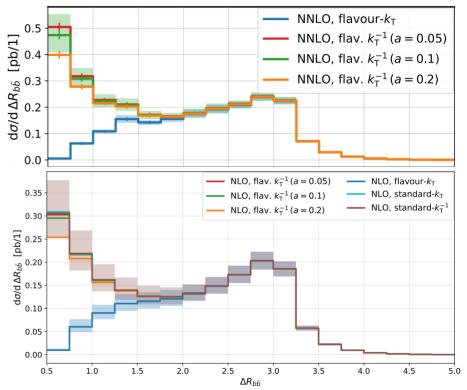


Comparison to data

Measurement of the production cross section of a W boson in association with two b jets in pp collisions at \sqrt{s} = 8 TeV, CMS 1608.07561

(assumes small unfolding corrections → wip)

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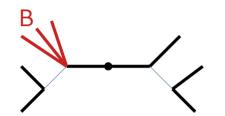


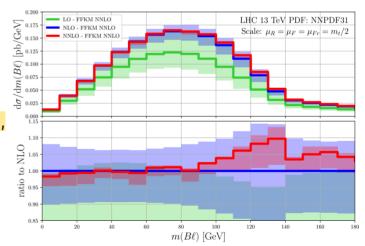
Significant differences between kT and anti-kT In small DeltaR(bb) region? Beam-function?!

Flavour tagging and fixed order fragmentation

- Fixed order QCD predictions with a final state hadron
- Partonic computation + transition of parton to hadron (collinear fragmentation of massless partons)
- Non-perturbative fragmentation function (similar to PDFs): Probability to find a hadron with a fraction x of a parton
- Advantage is that the hadrons momentum is measurable
 → usage as b-tag?
- Implementation in the STRIPPER framework through NNLO QCD:
 B-hadron production in NNLO QCD: application to LHC ttbar events with leptonic decays, Czakon, Generet, Mitov and Poncelet, 2102.08267

 $pp \to t\bar{t} \to B\ell\bar{\ell}\nu\bar{\nu}b + X$

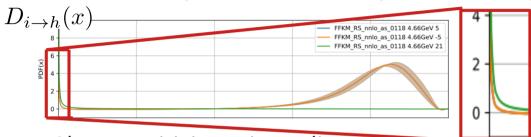




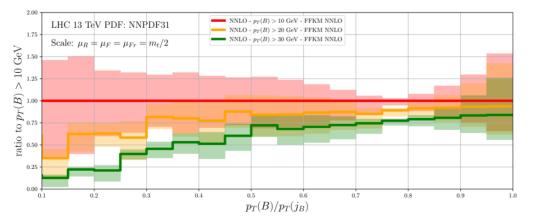
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Subtleties

 pT(B) requirement necessary since NNLO fragmentation function divergent for x → 0 due to g → bbar splitting:

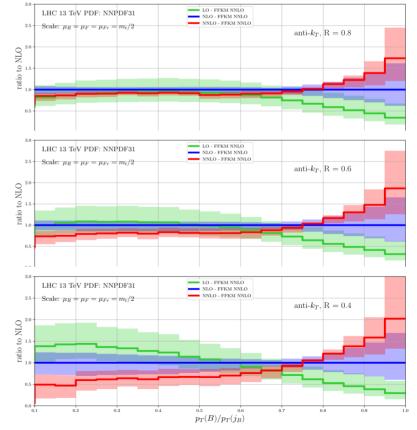


- Also: sensitivity to jet radius
 - → Usage as b-tag needs tuning



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Jet radius variation R = 0.8,0.6,0.4



Summary & Outlook

Summary:

- Many phenomenological applications of the sector-residue subtraction scheme
- NNLO QCD predictions for three-jet observables
- Flavoured jets → proposal for modified anti-kT algorithm
 Outlook:
- More three jet pheno results: TEEC, eventshapes
- More applications to $2 \rightarrow 3$ processes with 1 external mass
- More studies of flavoured jet algorithms