High precision at low energy: the case of MUonE

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HP2 2022 High Precision for Hard Processes

Discovery Museum, Newcastle upon Tyne, 20-22 September 2022



largest Th. unc. source: HVP



see talk by T. Teubner

B. Abi et al., Phys. Rev. Lett. 126 (2021) 14, 141801 [arXiv:2104.03281[hep-ex]]

An additional puzzle in theoretical predictions



T. Aoyama et al. Phys.Rept. 887 (2020) 1-166

B. Abi et al. [Muon g-2], Phys. Rev. Lett. 126 (2021) no.14, 141801.

Borsanyi, S. et al. Nature 593, 51-55 (2021).



G. Colangelo et al. arXiv:2203.1581 (Snowmass 2021)

recent new developments e.g.

- Lattice 2022 (8-13 August 2022)
- Fifth Plenary Workshop of the Muon g-2 Theory Initiative (5-9 September 2022)

A third independent determination more than welcome



- G. Abbiendi, C.M. Carloni Calame, U. Marconi, C. Matteuzzi, G. Montagna, O. Nicrosini, M. Passera, F. Piccinini, R. Tenchini, L. Trentadue, G. Venanzoni, *Measuring the leading hadronic contribution to the muon g-2 via μe scattering*
 - Eur. Phys. J. C 77 (2017) no.3, 139 arXiv:1609.08987 [hep-ph]
- * C. M. Carloni Calame, M. Passera, L. Trentadue and G. Venanzoni, A new approach to evaluate the leading hadronic corrections to the muon g-2

Phys. Lett. B 746 (2015) 325 - arXiv:1504.02228 [hep-ph]

Master formula

• Alternatively (exchanging s and x integrations in $a_{\mu}^{\rm HLO}$)

$$a_{\mu}^{\text{HLO}} = \frac{\alpha}{\pi} \int_{0}^{1} dx (1-x) \Delta \alpha_{\text{had}}[t(x)]$$
$$t(x) = \frac{x^{2} m_{\mu}^{2}}{x-1} < 0$$



e.g. Lautrup, Peterman, De Rafael, Phys. Rept. 3 (1972) 193

- \rightsquigarrow The hadronic VP correction to the running of α enters
- \leadsto Essentially the same formula used in lattice QCD calculation of $a_{\mu}^{
 m HLO}$
- * $\Delta \alpha_{had}(t)$ (and a_{μ}^{HLO}) can be directly measured in a (single) experiment involving a space-like scattering process

Carloni Calame, Passera, Trentadue, Venanzoni PLB 746 (2015) 325

- * Still a data-driven evaluation of a_{μ}^{HLO} , but with space-like data
- By modifying the kernel function $\frac{\alpha}{\pi}(1-x)$, also $a_{\mu}^{\rm HNLO}$ and $a_{\mu}^{\rm HNNLO}$ can be provided

Balzani, Laporta, Passera, arXiv:2112.05704 [hep-ph]

From time-like to space-like evaluation of $a_{\mu}^{\rm HLO}$



Smooth function

- \mapsto Time-like: combination of many experimental data sets, control of RCs better than O(1%) on hadronic channels required
- → Space-like: in principle, one single experiment, it's a one-loop effect, very high accuracy needed

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Abbiendi et al., EPJC 77 (2017) 3, 139

Abbiendi et al., Letter of Intent: the MUonE project, CERN-SPSC-2019-026, SPSC-I-252 (2019)

- Scattering μ's on e's in a low Z target looks like an ideal process (fixed target experiment)
- \rightsquigarrow The M2 muon beam ($E_{\mu} \simeq 160$ GeV) is available at CERN
- $\rightsquigarrow \sqrt{s} \simeq 0.4 \text{ GeV}$ and $-0.143 < t < 0 \text{ GeV}^2$
- \rightsquigarrow We can cover 87% of the a_{μ}^{HLO} space-like integral (and extrapolate to $x \rightarrow 1$)
- \rightsquigarrow With ~ 3 years of data taking, a statistical accuracy of 0.3% on $a_{\mu}^{\rm HLO}$ can be achieved

$$rac{1}{2}rac{\delta\sigma}{\sigma}\simeqrac{\deltalpha}{lpha}\simeq\delta\Deltalpha_{\mathsf{had}}$$

 $\Delta \alpha_{had}$ is a 0.1% effect in this region \rightarrow to measure it at 1%, σ must be controlled at the 10^{-5} level

• statistics: CERN muon beam M2 (E = 150 GeV), $1.3 \cdot 10^7 \ \mu/s$ with a target (Be/C) with total thickness of 60 cm $\implies L \sim 1.5 \cdot 10^7 \text{nb}^{-1} \implies$ statistical sensitivity $\sim 0.3\%$ on a_{μ}^{HLO} ($\sim 20 \cdot 10^{-11}$) in about 3 yrs of data taking

Sistematics

- (main) experimental sources
 - multiple scattering: E_e in normalization region much lower than in signal region Effect $\sim 1/E \implies$ it affects signal and normalization in different way
 - absolute μ beam energy scale, 5 MeV $\Longrightarrow 10^{-5}$ effect
 - angular intrinsec resolution ($\sim 1\%$)
 - longitudinal alignment ($\sim 10 \mu m$)
- theoretical: higher order radiative corrections modify the shapes
 - order of magnitude estimate, barring infrared logs and setting $c_{i,j} \sim 10$
 - $c_{1,1}\left(\frac{\alpha}{\pi}\right)L \sim 0.2$ $c_{1,0}\left(\frac{\alpha}{\pi}\right) \sim 2.5 \cdot 10^{-2}$ • $c_{2,2}\left(\frac{\alpha}{\pi}\right)^2 L^2 \sim 5 \cdot 10^{-3}$ $c_{2,1}\left(\frac{\alpha}{\pi}\right)^2 L \sim 5 \cdot 10^{-4}$ $c_{2,0}\left(\frac{\alpha}{\pi}\right)^2 \sim 5 \cdot 10^{-5}$
 - $c_{3,3} \left(\frac{\alpha}{\pi}\right)^3 L^3 \sim 1.5 \cdot 10^{-4}$ $c_{3,1} \left(\frac{\alpha}{\pi}\right)^3 L^2 \sim 1.5 \cdot 10^{-5}$ $c_{3,0} \left(\frac{\alpha}{\pi}\right)^3 L \sim 1.5 \cdot 10^{-6}$
 - the most advanced technologies for NNLO calculations and higher order resummation and matching are needed

• a modular apparatus has been proposed (40 independent tracking stations)





- whole acceptance covered with a 10×10 cm² silicon sensor
- thin targets equivalent to 60 cm
- ECal and Muon filter after last station, for PID and background rejection
- two Beam Tests already done at CERN (2017 and 2018)
 - Multiple Scattering measurements
 - 2 selection of a clean sample of elastic events
- Further Beam Test in October 2022
- 3 weeks Test Run in 2023 (proof of concept of the experimental proposal)
- 10 stations before LHC LS3 (2026) with first measurements of $a_{\mu}^{\rm HVP}$ with \sim 1% accuracy

G. Abbiendi et al., arXiv:1905.11677

G. Abbiendi et al., arXiv:2021.11111

First step towards precision: QED NLO



NLO virtual diagrams

analytical expression for tree level

$$\frac{d\sigma}{dt} = \frac{4\pi\alpha^2}{\lambda(s, m_{\mu}^2, m_e^2)} \left[\frac{(s - m_{\mu}^2 - m_e^2)^2}{t^2} + \frac{s}{t} + \frac{1}{2} \right]$$

- VP gauge invariant subset of NLO rad. corr.
- factorized over tree-level: $\alpha \rightarrow \alpha(t)$

(Van Nieuwenhuizen 1971, D'Ambrosio 1983, Kukhto et al. 1987, Bardin, Kalinovskaya 1997)

and corresponding real emission diagrams

• NLO matrix elements calculated with finite m_{μ} and m_{e} mass effects and a Monte Carlo program, MESMER, has been taylored to the fixed target kinematics

Alacevich, Carloni Calame, Chiesa, Montagna, Nicrosini, Piccinini, arXiv:1811.06743; JHEP 02 (2019) 155

Weak interaction effects (LO and NLO)



Alacevich, Carloni Calame, Chiesa, Montagna, Nicrosini, Piccinini, arXiv:1811.06743

- tree-level Z-exchange important at the 10^{-5} level
- purely weak RCs (in QED NLO units) at a few 10^{-6} level

Second step, towards *photonic* radiative corrections at NNLO

- | NLO virtual diagrams |²
- interference of LO $\mu e \rightarrow \mu e \gamma$ amplitude with



- + many others
- interference of LO $\mu e \rightarrow \mu e$ amplitude with



2-loop QED vertex form factors borrowed from Mastrolia and Remiddi, NPB 664 (2003) 341



NNLO double-virtual amplitudes where at least 2 photons connect the *e* and μ lines are approximated according to the Yennie-Frautschi-Suura ('61) formalism to catch the infra-red divergent structure

$$\widetilde{\mathcal{A}}^{\alpha^{2}} = \underbrace{\mathcal{A}_{e}^{\alpha^{2}} + \mathcal{A}_{\mu}^{\alpha^{2}} + \mathcal{A}_{e\mu, 1L \times 1L}^{\alpha^{2}}}_{\text{exact}} + \underbrace{\frac{1}{2}Y_{e\mu}^{2}\mathcal{T} + Y_{e\mu}\left(Y_{e} + Y_{\mu}\right)\mathcal{T} + \left(Y_{e} + Y_{\mu}\right)\mathcal{A}_{e\mu}^{\alpha^{1},\mathsf{R}} + Y_{e\mu}\mathcal{A}^{\alpha^{1},\mathsf{R}}}_{\text{YFS approximated}}$$



--- we estimate the subset of amplitudes in YFS approximation to miss terms of order

 $\left(rac{lpha}{\pi}
ight)^2 \ln^2\left(m_\mu^2/m_e^2
ight) \simeq 5 imes 10^{-4}$

going beyond this requires the full two-loop virtual amplitudes

R. Bonciani et al., PRL 128 (2022) 2; see talks by J. Ronca and Y. Ulrich

--- detailed comparisons ongoing with the independent Monte Carlo code McMule (PSI)

Showing

$$\Delta_{\rm NNLO}^{i} \equiv 100 \times \frac{d\sigma_{\rm NNLO}^{i} - d\sigma_{\rm NLO}^{i}}{d\sigma_{\rm LO}}$$

--- exact NNLO radiation from electron or muon leg, with or without acoplanarity cut







¹ of course with "double boxes" in YFS approximation

NNLO Virtual leptonic pairs (vacuum polarization insertions)

- any lepton (and hadron) in the VP blobs
- interfered with $\mu e \rightarrow \mu e$ or $\mu e \rightarrow \mu e \gamma$ amplitudes



Here the 2-loop integral is evaluated with dispersion relation techniques used in the past for Bhabha: Actis et al., Phys. Rev. Lett. 100 (2008) 131602; Carloni Calame et al., JHEP 07 (2011) 126

and for hadr. corr. in MUonE: Fael & Passera, PRL 122 (2019) 19

$$\frac{g_{\mu\nu}}{q^2 + i\epsilon} \to g_{\mu\nu}\frac{\alpha}{3\pi}\int_{4m_\ell^2}^{\infty} \frac{dz}{z}\frac{R_\ell(z)}{q^2 - z + i\epsilon} = g_{\mu\nu}\frac{\alpha}{3\pi}\int_{4m_\ell^2}^{\infty} \frac{dz}{z}\frac{1}{q^2 - z + i\epsilon}\left(1 + \frac{4m_\ell^2}{2z}\right)\sqrt{1 - \frac{4m_\ell^2}{z}}$$

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Real pair emissions

- they also contribute at NNLO
- squared absolute vaule of



- the emission of an extra electron pair μe → μe e⁺e⁻ is potentially a dramatically large (reducible) background, because of the presence of "peripheral" diagrams
- → A set of experimental cuts is needed to get rid of it. In addition to basic cuts (exactly one muon-like and one electron-like, with E ≥ 1 GeV, particle in the detector), we consider
 - 1. $\theta_{\mu\text{-like}}, \theta_{e\text{-like}} \geq \theta_c = 0.2 \text{ mrad}$
 - 2. acoplanarity $\leq 3.5 \text{ mrad}$
 - 3. geometric distance from the elastic curve in the $[\theta_{\mu}, \theta_{e}]$ plane < 0.2 mrad

Real e^+e^- pairs



 \rightsquigarrow only 0.007% of $\mu e \rightarrow \mu e \ e^+ e^-$ events survives the combination of the three cuts

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• π^0 production $p_2 \longrightarrow p_4 \longrightarrow p_4$

- The process $\mu e \rightarrow \mu e \pi^0$ with $\pi^0 \rightarrow \gamma \gamma$ as possible background, using a phenomenological model for the $\gamma^* \gamma^* \pi^0$ effective vertex
- → not an issue in the signal region

E. Budassi et al., PLB 829 (2022) 137138

- --- perhaps to be considered for NP searches in phase space region outside the signal one
- robustness of the measurement against possible New Physics "contamination" has been studied

A. Masiero, P. Paradisi and M. Passera, arXiv:2002.05418

P.S.B. Dev, W. Rodejohann, X.-J. Xu and Y. Zhang, arXiv:2002.04822

interesting proposals for New Physics searches at MUonE (new light mediators)

• invisibly decaying light Z' in $\mu e \rightarrow \mu e Z'$

Asai et al., arXiv:2109.10093

long-lived mediators with displaced vertex signatures

Galon et al., arXiv:2202.08843

• through scattering off the target nuclei $\mu N \rightarrow \mu N X$

Grilli di Cortona and E. Nardi, arXiv:2204.04227

Summary

- ---- Carloni Calame et al., PLB 746 (2015), 325
- --- Abbiendi et al., Eur. Phys. J. C77 (2017), 139
- → Mastrolia et al., JHEP 11 (2017) 198
- → Di Vita et al., JHEP 09 (2018) 016
- → Alacevich et al., JHEP 02 (2019) 155
- --- Fael and Passera, PRL 122 (2019) 19, 192001
- → Carloni Calame et al., JHEP 11 (2020) 028
- → Banerjee et al., SciPost Phys. 9 (2020), 027
- → Banerjee et al., EPJC 80 (2020) 6, 591
- → Budassi et al., JHEP 11 (2021) 098
- Balzani et al., arXiv:2112.05704 [hep-ph]
- ---- Bonciani et al., PRL 128 (2022) 2, 022002
- → Budassi et al., PLB 829 (2022) 137138

- → A lively theory community is active to provide state-of-the-art calculations to match the required accuracy for meaningful data analysis
- → Independent numerical codes (Monte Carlo generators and/or integrators) are developed and cross-checked to validate high-precision calculations. Chiefly
 - ✓ Mesmer in Pavia

github.com/cm-cc/mesmer

✓ McMule at PSI/IPPP

gitlab.com/mule-tools/mcmule

 $\mapsto\,$ An international MUonE collaboration is growing

THANK YOU

SPARES



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Virtual leptonic (and hadronic) NNLO VP corrections



