Comparison of public codes for Drell-Yan processes at NNLO accuracy

Adam Kardos in collaboration with Sergey Alekhin, Sven-Olaf Moch and Zoltán Trócsányi Based on arXiv:2104.02400 (EPJC)





#### LHC is performing really well:

- +  $\mathcal{O}(1-2\%)$  uncertainty for colorless final states (W $^\pm$  and Z)
- ATLAS Drell-Yan (fiducial) measurements at  $7\,{\rm TeV}\colon$ 
  - W $^+ 
    ightarrow \ell^+ \, 
    u 
    ho: 0.6\%$  (stat. unc.)
  - W $^- 
    ightarrow \ell^- ar{
    u} 
    ceil{:} 0.5\%$  (stat. unc.)
  - Z/ $\gamma^* 
    ightarrow \ell^+ \, \ell^- \colon 0.32\%$  (stat. unc.)
  - Normalization uncertainty: 1.8% (luminosity)
- Colorless measurements are systematics





To make best of the measurements high accuracy predictions are needed

- In p p collisions (typically) massless final states are available at NNLO accuracy in QCD
- NLO accuracy in EW corrections
- These predictions can be obtained by various tools using local subtractions or global slicing methods





Subtractions or slicing, does it make any difference? (See also Kirill's talk from Tuesday) The original integral:

$$= \lim_{\epsilon \to 0} \left[ \int_0^1 \frac{\mathrm{d} \mathbf{x}}{\mathbf{x}} \mathbf{x}^{\epsilon} \mathsf{F}(\mathbf{x}) - \frac{1}{\epsilon} \mathsf{F}(0) \right]$$

• Subtraction: zero is added in a clever(ish) way:

$$I = \lim_{\epsilon \to 0} \left[ \int_0^1 \frac{\mathrm{d}x}{x} x^\epsilon \left( \mathsf{F}(x) - \mathsf{F}(0) \right) + \int_0^1 \frac{\mathrm{d}x}{x} x^\epsilon \mathsf{F}(0) - \frac{1}{\epsilon} \mathsf{F}(0) \right] = \int_0^1 \frac{\mathrm{d}x}{x} \left( \mathsf{F}(x) - \mathsf{F}(0) \right)$$

$$= \int_0^1 \frac{\mathrm{d}x}{x} \left( \mathsf{F}(x) - \mathsf{F}(0) \right)$$
UNIVERSITY of DEBRECEN

• Slicing: we alter the singular piece of the calculation

$$\begin{split} & \mathsf{I} \sim \lim_{\epsilon \to 0} \left[ \mathsf{F}(0) \int_0^\delta \frac{\mathrm{d} \mathsf{x}}{\mathsf{x}} \mathsf{x}^\epsilon + \int_\delta^1 \frac{\mathrm{d} \mathsf{x}}{\mathsf{x}} \mathsf{x}^\epsilon \mathsf{F}(\mathsf{x}) - \frac{1}{\epsilon} \mathsf{F}(0) \right] = \\ & = \mathsf{F}(0) \log \delta + \int_\delta^1 \frac{\mathrm{d} \mathsf{x}}{\mathsf{x}} \mathsf{F}(\mathsf{x}) \end{split}$$

To not to change the value too much  $\delta$  should be chosen small!

Drell-Yan processes are calculated at NNLO in QCD both by subtraction and by slicing methods



#### Public codes implementing the DY process:

- DYNNLO http://theory.fi.infn.it/grazzini/dy.html
  - q<sub>T</sub> subtraction: global slicing method
- FEWZ https://www.hep.anl.gov/fpetriello/FEWZ.html
  - Sector decomposition: local subtraction method
- MATRIX https://matrix.hepforge.org/
  - q<sub>T</sub> subtraction: global slicing method
- MCFM https://mcfm.fnal.gov/

FRRECEN

- N-jettiness subtraction: global slicing method
- DYTurbo https://dyturbo.hepforge.org/
  - Reimplementation of DYNNLO with resummation, at fixed order same results as DYNNLO
     UNIVERSITY of



How to compete with experimental accuracy?

- Theory predictions are numbers stemming from many uncertainties:
  - PDF uncertainty: better PDF fits
  - Dependence on non-physical scales ( $\mu_{\rm R}$ ,  $\mu_{\rm F}$ ): going to higher orders
  - Statistical uncertainties:
    - More PS points in integration
    - Better integrator
    - Optimize over special aspects of calculation (dynamics, subtractions, &c.)
  - Method-dependent parameters:
    - Non-physical
    - Result cannot depend on them
    - No dependence or dependence smaller than statistical uncertainty
    - $\Rightarrow$  High precision runs might need parameter refinement



- Experimental data reached very high accuracy
- Computational power also increased
- Became possible to deliver (numerically) very precise NNLO computations
- ATLAS reported (arXiv:1612.03016):

is observed. For the fiducial and differential cross-section measurements with additional kinematic requirements on the lepton transverse momenta and rapidities, however, poorer agreement is found: for the integrated fiducial  $W^+$ ,  $W^-$ ,  $Z/\gamma^*$  cross sections, the differences between FEWZ and DYNNLO predictions calculated with the ATLAS-epWZ12 PDF set amount to (+1.2, +0.7, +0.2)%, which may be compared to the experimental uncertainties of ±(0.6, 0.5, 0.32)%, respectively.<sup>3</sup>

- Computational tools are black boxes for experiments
- ⇒ Better to check consistency



Idea:

- Take publicly available codes for DY
- Fix parameters
- Validate analysis and parameters through LO and NLO
  - At LO checking parameters through dynamics
  - At NLO most of the tools use Catani-Seymour subtraction
  - $\Rightarrow$  At NLO checking numerical integration
- Target cross section precision is aimed at  $\mathcal{O}(0.1\%)$  for each bin
  - Aim is not to compare accuracy for the NNLO contribution
  - Physical cross section should have high accuracy (this is measured)



# Calculational setup

Would like to test programs in realistic environment:

- ATLAS data for W $^{\pm}$  and Z/ $\gamma^{*}$  at 7 TeV [arXiv:1612.03016].
  - Pseudo-rapidities for decay leptons (e^ $\pm$  and  $\mu^{\pm}$ ) (W^{\pm}) and for decay lepton-pairs [Z/ $\gamma^*$ ]
  - Cuts on lepton  $p_{\perp}$  and pseudo-rapidities
  - -~ For Z/ $\gamma^*$  central and forward region are considered
- D0 data for W $^\pm$  at  $1.96\,{\rm TeV}$  [arXiv:1412.2862]
  - Electron charge asymmetry (A<sup>e</sup>) measured in electron pseudo-rapidity
  - Also in forward region
  - Symmetric  $p_{\perp}$  cuts:  $p_{\perp}^{\nu} > 25 \, \mathrm{GeV}$ ,  $p_{\perp}^{\ell} > 25 \, \mathrm{GeV}$
  - Staggered p<sub> $\perp$ </sub> cuts:  $p_{\perp}^{\nu} > 25 \,\mathrm{GeV}$ ,  $p_{\perp}^{\ell} > 35 \,\mathrm{GeV}$
- EW parameters were chosen to minimize NLO EW corrections (irrelevant for



#### Theory tools

Used publicly available tools for the comparisons:

- DYNNLO (version 1.5)
  - Legacy code
  - Used by experiments
  - Superseded by MATRIX
- FEWZ (version 3.1)
  - Used as baseline due to being local subtraction
- MATRIX (version 1.0.4)
- MCFM (version 9.0)



#### NLO comparisons

NI O check at 7 TeV

- W<sup>±</sup> in central
- $Z/\gamma^*$  in central

•  $Z/\gamma^*$  in forward region





2

3

n.

 $\eta_{\parallel}$ 

# DYNNLO NNLO comparisons



Default  $r_{\rm cut}^{\rm min}=q_{\perp}^{\rm min}/M_V=0.8\%$  slicing parameter was used UNIVERSITY of DEBRECEN

### MATRIX NNLO comparisons

- MATRIX employes the  $\mathsf{q}_\perp$  subtraction: a global slicing method
- $\Rightarrow\,$  Slicing parameter should be selected carefully:
  - -~ Default slicing parameter for W^{\pm}:  $r_{\rm cut}^{\rm min}=0.15\%$
  - Default slicing parameter for Z/ $\gamma^*$ :  $\vec{r}_{\rm cut}^{\rm min}=0.05\%$
  - + MATRIX offers to extrapolate  $r_{\rm cut}^{\rm min}$  to 0



# MATRIX NNLO comparisons





Note: no extrapolation applied. Extrapolation is not enough to eradicate UNIVERSITY of differences DEBRECEN

### MCFM NNLO comparisons

- MCFM employs the N-jettiness subtraction: a global slicing method
- $\Rightarrow$  Slicing parameter should be selected carefully:
  - Default slicing parameter:  $\tau_{\rm cut} = 6\cdot 10^{-3}$
  - Decreased as much as possible to be still in reasonable run times

— Tried 
$$au_{
m cut}=10^{-3}$$
 and  $au_{
m cut}=4\cdot10^{-4}$ 



# MCFM NNLO comparisons



17/23

## NNLO comparisons



# FEWS vs. NNLOJET



# NNLO comparisons

- MC integration was pushed to have negligible (stat.) uncertainty on plots
- Decreasing slicing parameters and extrapolation helped to bring predictions closer
- Specialties of the experimental cuts put pressure on theory calculations
  - Symmetric cuts
  - Forward region
- Staggered cuts in central region gives the best agreement between codes
- Missing power corrections can be a reason for deviations



## Conclusions

- NNLO calculations are really tough
- People put extraordinary efforts to deliver NNLO predictions to experiments
- NNLO predictions are driving forces behind activity at colliders
- All NNLO methods and calculations mirror the exceptional genius of people behind them
- Experiments reached really high precision
- Experiments can probe regions of phase space challenging to some methods
- Luckily we have several methods
- $\Rightarrow\,$  Can be decided which is best for each scenario
- $\Rightarrow$  Can fine-tune methods to cope with fiducial cut challenges



# The Galileo Galilei Institute For Theoretical Physics

Centro Nazionale di Studi Avanzati dell'Istituto Nazionale di Fisica Nucleare

Arcetri, Firenze



#### Workshop

Theory Challenges in the Precision Era of the Large Hadron Collider

Aug 28, 2023 - Oct 13, 2023

Apply (deadline: May 01, 2023 )

Related events Theory Challenges in the Precision Era of the Large Hadron Collider (Conference) - Aug 28, 2023

Theory Challenges in the Precision Era of the Large Hadron Collider (Training Week) - Sep 25, 2023

# Thank you for your attention!

# Back-up slides



#### NLO comparisons at D0





# NNLO comparisons at D0 for DYNNLO





### NNLO comparisons at D0 for MATRIX



4/12

### NNLO comparisons at D0 for MCFM



5/12

# NNLO comparisons at D0 for staggered cuts





CPU times

#### MATRIX CPU times for ${\sf r}_{cut}^{min}=0.05\%$ :

Region	CPU time [h]
Central	200.000
Forward	350.000

#### MCFM CPU times for $\tau_{\rm cut} = 4 \cdot 10^{-4}$

Process	CPU time [h]
W±	180.000
${\sf Z}/\gamma^*$ Central	160.000
${\sf Z}/\gamma^*$ Forward	50.000



# MATRIX improvements

#### In arXiv:2111.13661:

- "Transverse-momentum cuts on undistinguished particles in two-body final states induce an enhanced sensitivity to low momentum scales"
- Linear power corrections are implemented to "circumventing the numerical instabilities related to the use of a tiny value of the slicing parameter"
- Inclusion of linPCs resulted in an agreement
   with FEWZ





• Power corrections can be defined through slicing parameters:

$$\tau = \frac{\mathsf{q}_{\perp}^2}{\mathsf{Q}^2} = \frac{\left(\sum_i \mathbf{k}_{\perp,i}\right)^2}{\mathsf{Q}^2} \left(\mathsf{q}_{\perp} \operatorname{slicing}\right), \quad \tau = \frac{\sum_i 2\min\left(\mathsf{p}_{\mathsf{a}} \cdot \mathsf{k}_i, \mathsf{p}_{\mathsf{b}} \cdot \mathsf{k}_i\right)}{\mathsf{Q}^2} \left(\operatorname{jettiness slicing}\right)$$

• The phase space is partitioned into two disjoint regions:

$$\sigma = \int \mathrm{d}\tau \frac{\mathrm{d}\sigma}{\mathrm{d}\tau} = \int_{\tau_{\rm cut}} \mathrm{d}\tau \frac{\mathrm{d}\sigma}{\mathrm{d}\tau} + \int^{\tau_{\rm cut}} \mathrm{d}\tau \frac{\mathrm{d}\sigma}{\mathrm{d}\tau} = \sigma(\tau_{\rm cut}) + \int^{\tau_{\rm cut}} \mathrm{d}\tau \frac{\mathrm{d}\sigma}{\mathrm{d}\tau}$$

• Analytical integration in first term using universal QCD factorization:

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\tau} \sim \delta(\tau) + \sum_{i} \left[ \frac{\log^{i} \tau}{\tau} \right]_{+} + \sum_{j} \tau^{\mathsf{p}-1} \log^{j} \tau + \mathcal{O}(\tau^{\mathsf{p}})$$
(ERSITY of BRECEN

- Fiducial cuts affect decay phase space
- $\Rightarrow\,$  Decay phase space has effect on power corrections due to  $q_{\perp}$  dependent terms
  - Can get idea about linear power corrections through decay phase spaces [arXiv:1911.08486, arXiv:2006.11382]
    - $\ \Phi_L(\textbf{q}_\perp)$  leptonic phase space with  $\textbf{q}_\perp$  transverse momentum
    - $\Phi_{\mathsf{L}}(0)$  leptonic phase space with zero transverse momentum
  - Linear power corrections are estimated through difference from Born decay phase space:

$$\left|1-\frac{\Phi_{\mathsf{L}}(\mathsf{q}_{\perp})}{\Phi_{\mathsf{L}}(0)}\right|$$









