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Astrophysical Constraints on Axions

David J. E. Marsh Recent Progress in Axion Physics, Durham, Sept. 2022



Science and Technology Facilities Council



COSMIC STRUCTURES

Axion Cosmology

Eq. of state, $w \rightarrow$ homo. pressure. Appears in Friedmann \rightarrow affects expansion. Sound speed, $c_s \rightarrow$ pressure perts. Affects growth of structure.

	CDM	Baryons	Photons	Λ
W	0	0	1/3	-1
C _s ²	0	0 (late) 1/3 (early)	1/3	Does not cluster (-1)

Cosmology: measure n-point functions on the sky, normally in Fourier space. Compare to theory predictions by Bayesian parameter fitting (typically MCMC).

CMB measures the expansion rate via Sachs Wolfe (late) and Silk Damping (early). Late time observables sensitive to late time expansion + growth of structure.

CMB Constraints

Estimate

Consistent w/ Λ CDM expansion rate to $z^{10^5} \rightarrow$ DM formed before this time.

 $m > 2.6 \times 10^{-25} \text{ eV}$ =H(10⁵)

Precision

Solve Boltzmann equations & MCMC 8+ parameters with Planck likelihood using cosmosis+axionCAMB.



Also use polarization & lensing anisotropies + correlations. Vary UBDM and CDM density simultaneously.

CMB Constraints

Estimate

Consistent w/ LCDM expansion rate to $z^{10^5} \rightarrow$ DM formed before this time.

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Solve Boltzmann equations & MCMC 8+ parameters with Planck likelihood using cosmosis+axionCAMB.



Talk to Johannes Esklit

Birefringence

 $\beta=0.30^{\rm o}\pm 0.11^{\rm o}$

Minami & Komatsu (2020) Planck collab. (2022)

Isotropic birefringence can be caused by an ultralight axion via:

$$\mathcal{L} = g\phi F_{\mu\nu}\tilde{F}^{\mu\nu} \Rightarrow \beta = \int_{\eta_{\rm CMB}}^{\eta_0} g\frac{d\phi}{d\eta}d\eta$$

This fixes the axion mass to a range also probed by primary anisotropies and lensing:

$$10^{-33} \text{ eV} \lesssim m \lesssim 10^{-28} \text{ eV}$$

$$_{\text{H}_0} \qquad \qquad \text{H}_{\text{CMB}}$$

Isocurvature in the ultralight axion also induces anisotropic birefringence and large angle BB with amplitude fixed by scale of inflation.



Birefringence constraints can be highly complementary/synergistic to direct searches:



SO opportunities: foregrounds, anisotropic component, ultralight DM bounds...

Fig: Planck (2018)



Power spectrum = F.T. of two-point correlation of density perts.



Estimating Bounds

DJEM & Hoof (2021)



Power Spectrum Constraints

Example 1: DES lensing (Dentler, DJEM et al, 2021)

- Measure "galaxy shear correlation function" consistent with CDM.
- Compute from P(k) using axionCAMB (linear theory) + "halo model" non-linearities.
- DES-Y1 Bayesian \rightarrow m>10⁻²³ eV.
- DES-Y3 forecast 10⁻²² eV, Euclid 10⁻²⁰ eV.

Lensing advantage: directly measure DM



Example 2: BOSS Lyman-alpha forest (Rogers & Peiris, 2020). Measure flux power spectrum consistent with CDM. Non-linear and gas physics using N-body +hydro+"emulator" for parameter dependence. Limit m>2x10⁻²⁰ eV.

INSIDE GALAXIES

Schrödinger-Poisson

Non. relativistic limit of the Klein-Gordon-Einstein equations:

$$\phi = \frac{1}{\sqrt{2m^2}} \left(\psi e^{imt} + \psi^* e^{-imt} \right)$$

Real field \rightarrow Complex

$$\psi = \sqrt{\rho} e^{iS} , \vec{v} = \nabla S$$

Madelung/fluid interpretation



"Fuzzy DM physics" different from CDM/WDM etc on scales ~ de Broglie. Waves → interference. Gradient pressure → stable solutions & coherence. Challenge: simulate this equation over a wide range of scales. Cosmology covers Gpc \rightarrow kpc lengths, and overdensities ~ 10⁵.



Soliton+ halo structure also in non-linear optics. Called "incoherent soliton"

Note: this is all classical physics

Fig. Philip Mocz

Dwarf galaxies 10⁸ M_{sol}, particle mass ~10⁻²² eV



Veltmaat et al (2018)

"Inflaton halos" ~ 100g, particle mass ~10¹⁴ GeV





Schive et al (2014)





Mixed CDM+gas

Mocz et al (2019)

Lague et al (forthcoming)

Mixed CDM+ FDM

Opportunity/challenge: physics consequences of interference fringes.

Note: this is all classical physics

Inside a Halo: Dynamics

Velocity field obeys Maxwell-Boltzmann. Coherence length and time:

$$L = \frac{2\pi}{m\langle v \rangle} \qquad \tau = \frac{2\pi}{m\langle v \rangle^2}$$

Coherent patches \sim quasi-particles \rightarrow gravitationally scatter and heat/cool.

 $t_{\rm relax} \sim 10^{10} m_{22}^3 v_{100}^2 R_5^4 \text{ yr}$

Survival of old star cluster in Eridanus-II → exclude too much heating

 \rightarrow Lower bound m>10⁻¹⁹ eV

Avoid with mixed DM models, and v low masses.

DJEM & Niemeyer (2019)



Movies made with data from Veltmaat et al (2018)

INTENSITY MAPPING

The future of cosmology...

Bauer, DJEM et al (2020); Hotinli, DJEM & Kamionkowski (2021)

Neutral Hydrogen: z=1100 to 0

Line intensity of neutral hydrogen gas \rightarrow possibility to map Universe in 3d. CMB modes ~ L² IM tomography modes ~ k³ \rightarrow huge increase in available information



Fig: Mao et al (2008)

Neutral Hydrogen: z=1100 to 0

1100<z<10 "dark ages": neutral hydrogen everywhere, nearly linear.

10<z<0 "post-reionization": neutral hydrogen inside dark matter halos



Post-Reion.: Peaks and Halos

"Halo model" \rightarrow dark matter distribution and statistics from "peak theory".

"HI halo model": assert a relation between halo mass and hydrogen + density profile.



Two-point stats: correlations between halos versus within. Standard calculation.



Dark Ages: "VAOs"

Theory: Tseliakhovitch & Hirata (2014) Simulation: 21cmvFAST, Munoz (2019)

DM-baryon relative velocity (vBC): coherent on MPc scales + baryon acoustic oscillations. First stars collapsing HI at Jeans scale ~ kpc. Star formation different in vBC coherence patches \rightarrow couple small-large scales \rightarrow "VAO".



Bias of Mixed Dark Matter

Suppressed clustering \rightarrow fewer halos. But, fixed total amount of hydrogen \rightarrow increase the bias.

Consequence: increase the HI power spectrum on large scales. Consistent with N-body simulations.



Forecasts: SKA+CMB



Fisher matrix = inverse covariance $F_{ij} = \sum_{\ell} \frac{1}{(\Delta C_{\ell})^2} \frac{\partial C_{\ell}}{\partial p_i} \frac{\partial C_{\ell}}{\partial p_j},$ $(\Delta C_{\ell})^2 = \frac{2}{(2\ell+1)f_{\rm sky}}(C_{\ell} + N_{\ell})^2,$ $m_a = 10^{-28} \,\mathrm{eV}$ $m_a = 10^{-24} \,\mathrm{eV}$ 0.10single dish, full survey interf., linear scales only 0.08 interf., full survey Ω_a/Ω_d 0.040.02 0.00 0.220.240.260.280.22 0.240.26 0.28

 $\Omega_{\rm CDM}$

 $\Omega_{\rm CDM}$

VAO Signature & Bias



P(k) cut-off with k<kJ \rightarrow remove first stars \rightarrow drastically suppress VAO amplitude for m<10⁻¹⁸ eV.

HERA Forecasts

Detect VAO signature at ~20 σ with CDM, thus very sensitive to P(k) cut-off with axions.



Foreground modeling.

Baryonic feedback modeling.

Summary



Post-reionization \rightarrow orders of magnitude improvement over CMB. Test GUT scale predictions. Dark ages \rightarrow increase lower limit on (fuzzy)DM particle mass \rightarrow close gaps astro – black holes.



Test string theory predictions of Cicoli et al (2021)? Let's go further...

AXION STARS/SOLITONS

Solitons from Schrödinger

Ground state solutions of the SP equations \rightarrow one parameter family.

Ansatz

$$\chi'' + \frac{2\chi'}{r} = 2(V - \gamma)\chi$$
 Dimensionless variables with $\chi(0)=1.$
$$V'' + \frac{2V'}{r} = \chi^2$$

 $\psi(r,t) = \chi(r)e^{-i\gamma t}$

→ Boundary value problem with eigenvalue γ . Solved for:

$$\gamma = -0.69, V_0 = -1.34$$

Zero oscillation, lowest energy solution.



Soliton Formation

Solitons cores of DM halos form during "violent relaxation" from initial coherence. Second formation mechanism: gravitational BEC in "kinetic regime".



Gravitationally stable simulation box.

Condensation and growth time predictable from scattering theory.



Gravitationally unstable: form a DM halo around the nucleated axion star. Virial equilibrium \rightarrow core size+growth quench.

Soliton Distribution

Soliton dists+merger rates: Du, DJEM et al (in prep)

 10^{15} Equilibrium between core+halo \rightarrow relation 10^{11} between host mass and soliton: $M_{\rm sol} \propto M_h^{1/3}$ $F_S [Mpc^ 10^{7}$ 10^{3} Compute "halo mass function" \rightarrow predict soliton mass function. 10^{-1} $F_S := \frac{dn_{\rm sol}}{d\ln M}$ $10^{-5}_{10^{-8}}$ number density 10^{-6} 10^{-2} 10^{-4} Axion Star Mass, $M_S \left[M_{\odot} (10^{-11} \text{ eV}/m_a) \right]$

Theory and simulation work around the initial formation in relatively low mass halos → we can do cosmology, merger rates, etc. Challenge: what is the distribution in e.g. the Milky Way?

Boson Stars: Full GR

Spin-0 bosons \rightarrow oscillating metric on time scale m⁻¹ \rightarrow "oscillatons". Simulate in full 3D, with axion interaction potential using GRChombo.

No interactions: BH formation at "Kaup mass":

$$M_{\rm Kaup} \approx \frac{M_{pl}^2}{m} \Rightarrow \phi_0 \sim M_{pl}$$

Finite interactions given by f → explosion in "axion Nova":

$$M_{\rm Nova} \approx \frac{M_{pl} f_a}{m} \Rightarrow \phi_0 \sim f_a$$

Interactions → more instabilities e.g. radio emission by F*F (in prep.).



MINICLUSTERS



Aesthetic advantage over pre-inflation SSB: no free parameter of initial field value. Disadvantage: very hard computational problem. Little consensus on the relic density. QCD axion axion mass in meV range + miniclusters. Stats: see Riess, DJEM & Hoof (2021)



Fig: Buschmann et al (2022)

Largely unknown: role of domain walls.

Very similar situation occurs for dark photon DM from isocurvature.

Miniclusters

String network \rightarrow enhanced fluctuations:

— Minicluster

- · · Adiabatic

 $P(k) \propto k^{-6.89}$

109

1010





0.00 0.05 0.10 0.15 0.20 0.25 0.30 0.35 $x \, [\mathrm{pc}/h]$

Minicluster Pheno

Possible pheno: microlensing, appearance in haloscopes, collision w/ neutron stars. Vital questions: distribution, density profiles, survival. End-to-end sims impossible.



MCs must be dense and heavy.Earth encounters rare, tidal $m^1 meV$ window?streams possible \rightarrow astronomy.

M31: NS pops \rightarrow only low m possible.

SUPERRADIANCE!

Black Hole Superradiance

Review: Brito et al (2015)

Solve for instabilities of KG equation on

Kerr:
$$\Box \phi - \partial_{\phi} V(\phi) = 0$$

Non-relativistic limit in "tortoise coords", find instability ($\omega < 0$):

$$\frac{d^2\psi_{lm}}{dr^{*2}} = \left[\omega^2 - V(r,\omega)\right]\psi_{lm} \,.$$

Ergo-region Barrier region Potential Well Exponential growth region "Mirror" at r~1/μ Fig: Arvanitaki & Dubovsky (2010) → Black Hole Horizon r* Physical picture: "Penrose process/ black hole bomb"



Resonant bosons extract spin from astrophysical BHs, if $\Gamma_{\rm SR}$ > $\Gamma_{\rm others}$

Spin-0 Fields, No ϕ^4 Term

GIF by Matthew J. Stott



"Exclusion probability" is marginal likelihood. Statistically robust constraints.



Stott & DJEM (2018)

φ⁴ Instability: "Bosenova"

Yoshino & Kodama (2012); Arvanitaki+(2014); Stott (2018)

Bose enhanced 2-2 scattering in superradiant cloud can have a rate $\Gamma_4 > \Gamma_{SR}$. Shuts off SR by cloud collapse above critical value of $\lambda \phi^4$ coupling, $\lambda = m^2/f_{\text{pert.}}^2$



advanced rate calcs. Quantitatively similar.

STRING AXIVERSE

Arvanitaki et al (2009); Mehta, DJEM et al (2021)

The KS Axiverse

Triangulate (FRST) KS polytopes \rightarrow CY₃. 1000's of CYs at large h¹¹ in laptop-time! Pick point in Kähler cone (no stabilization). Kähler metric and axion potential computed:

$$\mathcal{L} = -\frac{M_{\rm pl}^2}{8\pi^2} K_{ij} g^{\mu\nu} \partial_\mu \theta^i \partial_\nu \theta^j -\sum_{a=1}^\infty \Lambda_a^4 \Big\{ 1 - \cos(\mathcal{Q}^a{}_i \theta^i + \delta^a) \Big\}$$

For an astrophysicist: databases of K, Λ , Q sampling KS, triangulations, and Kähler cone. Aim: rigorously exclude CY's based only on vacuum properties of the axions.

Demirtas et al (1808.01282)





Axion Spectra

Find vacua of V(f) in fundamental domain. Expand to quartic order \rightarrow masses +quartics ("fpert").

Trends: Kähler cones become very narrow at large h I I \rightarrow cycles in the CY have large volumes \rightarrow (ultra)light axions and smaller decay constants.



Constraints on IIB CY Vacua

Ensemble of O(10⁵) CYs. All up to h11=5. 100 per h11 up to 176. Up to 100 per h11 to 491.



Beyond Superradiance

- Why superrradiance? Vacuum process, no cosmological assumptions. Only need the axion potential.
- Why go further? Large h11, and moduli space away from Kahler cone tip \rightarrow larger volume \rightarrow lower decay constants \rightarrow superradiance shuts off.
- What observables will be best? Ideally vacuum processes, cosmology independent, exploit massless fields.
- What is a bad observable? Unfortunately, axion DM from realignment: too many cosmological assumptions.

Demirtas et al (2021)

Visible Sector Couplings

- Choose divisor for QCD. Dilate V until divisor volume $\rightarrow \alpha_{\rm QCD}$. Demand geometric.
- Axion masses and f's by using hierarchy of instanton scales + Kähler metric.
- Strong-CP: $\Sigma_i \theta_i$ in front of CS must be small.
- Assuming a GUT you can find $g_{a\gamma}$:

$$\mathcal{L}_{\rm EM,CS} = \sum_{i} c_i \frac{\alpha_{\rm EM}}{2\pi f_i} \theta_i F \tilde{F} \equiv \sum_{i} g_i \theta_i F \tilde{F}$$
(ci's known)

Effective axion-photon coupling for massless $g_{\rm eff}$ = linear combination

$$= \sqrt{\sum_{i \in \{m=0\}} g_i^2}$$

3e4 random CYs Computationally limited at large h11



Axion-Like PROpagation

X-ray Spectrum Oscillations

e.g. Matthews et al (2022); Reynolds et al (2020) "other David Marsh"

Photon-axion conversion in cluster B-fields \rightarrow spectrum oscillations. Vanishing if $m_a > \omega_p$. Need to marginalize random magnetic field models. Fit Chandra satellite data.





Freeze-in DM & Decays

Axion production via Primakoff process off SM charged particles from vacuum initial state. Irreducible contribution to DM, all fields with m<reheat temperature. Must not "overclose".



Perturbative decay \rightarrow photodissociation of elements, cosmic reionization etc. even for $\xi \sim 10^{-10}$.

$$\Gamma_a = \frac{g_{a\gamma}^2 m_a^3}{64\pi} = \frac{1}{1.32 \times 10^8 \text{ s}} \left(\frac{g_{a\gamma}}{10^{-12} \text{ GeV}^{-1}}\right)^2 \left(\frac{m_a}{10 \text{ MeV}}\right)^3 \,.$$

(very) Preliminary Results

String Datasets:

- PQAxiverse: conditioned on correct QCD coupling and strong-CP. Masses and visible sector couplings estimated. O(10) CYs at h11=491.
- h11=491: exploratory dataset, O(10⁵) CYs. Zeroth order assumption: massless axions, estimate couplings from Kähler eigs only. Tip of Kähler cone → smaller couplings than the bulk → limits conservative.

Astrophysical Datasets:

- Chandra: X-rays effective likelihood for single massless axion. Ignore CYs with any resonant axions. Future: modify ALPro for one axion in resonant region + massless.
- Freeze-in DM: crude cuts on overclosure and decaying DM (PQAxiverse only). Future: build cosmological likelihoods using micrOMEGAS and GAMBIT.





20% favoured if h^{11} <200. Spectrum overfitting. Zero allowed manifolds h^{11} >250.



All randomly sampled CYs at 491 are strongly excluded (saturate likelihood). Recall: tip of SKC \rightarrow limits are likely conservative over moduli space (!)

Speculations and Hopes

- Triangulations (CYs) of 491 polytope obeying constraints are VERY rare. Probably true of most of moduli space. Astrophysics prefers "boundaries" of KS axiverse.
- Can we use ML techniques to find allowed models at 491?
- Is there a symmetry underlying this? Can we use to ML to find it?
- If there is a symmetry \rightarrow restrict combinatorics that count the CYs at 491 \rightarrow make a brute force exploration of remaining 491 landscape possible?

LEVERHULME TRUST_____

Topology from cosmology: axions, astrophysics, and machine learning



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