

Three loop corrections for the QCD and QED form factors

Results and code interface

4th Workstop/Thinkstart: Towards $N^3\text{LO}$ for $\gamma^* \rightarrow \ell\ell$ | IPPP Durham | August 3, 2022
Matteo Fael | with Fabian Lange, Kay Schönwald and Matthias Steinhauser

Outline

- Results for QED
- C/Fortran code of the Form Factors
- Hadronic contributions at $O(\alpha^3)$?

Renormalization and IR subtraction

- UV renormalization in the on-shell scheme. Mass and wave function renormalization.
Melnikov, van Ritbergen, Phys.Lett.B 482 (2000) 99;
 Chetyrkin, Steinhauser, Nucl.Phys.B 573 (2000) 617-651
- Coupling constant renormalized in \overline{MS} . We use $\alpha_s^{(n_f)}$ as expansion parameter.
- Structure of IR poles is universal

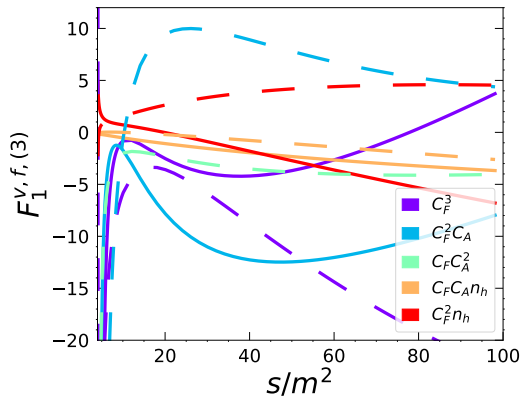
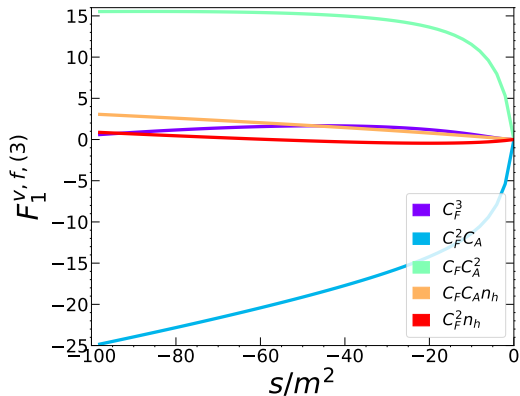
$$F_i^{\text{UV ren}}(s) = Z_{\text{IR}} F_i^{\text{f}}(s)$$

with Z_{IR} given in term of the $\Gamma \equiv \Gamma_{\text{cusp}}$ anomalous dimension

$$\begin{aligned} \log Z_{\text{IR}} = & -\frac{\alpha_s(\mu)}{\pi} \frac{\Gamma_{\text{cusp}}^{(1)}}{2\epsilon} + \left(\frac{\alpha_s(\mu)}{\pi}\right)^2 \left(\frac{\beta_0 \Gamma_{\text{cusp}}^{(1)}}{16\epsilon^2} - \frac{\Gamma_{\text{cusp}}^{(2)}}{4\epsilon}\right) \\ & + \left(\frac{\alpha_s(\mu)}{\pi}\right)^3 \left(-\frac{\beta_0^2 \Gamma_{\text{cusp}}^{(1)}}{96\epsilon^3} + \frac{\beta_1 \Gamma_{\text{cusp}}^{(1)} + 4\beta_0 \Gamma_{\text{cusp}}^{(2)}}{96\epsilon^2} - \frac{\Gamma_{\text{cusp}}^{(3)}}{6\epsilon}\right) \end{aligned}$$

Grozin, Henn, Korchemsky, Marquard, Phys.Rev.Lett. 114 (2015) 6, 062006; JHEP 01 (2016) 140.

Results in QCD



Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

QEDfication

- $C_F \rightarrow 1$

$$(T_{\text{fun}}^A T_{\text{fun}}^A)_{ij} = C_F \delta_{ij}$$

- $C_A \rightarrow 0$

$$f^{ACD} f^{BCD} = \text{Tr}(T_{\text{adj}}^A T_{\text{adj}}^B) = C_A \delta_{AB}$$

- $T_F \rightarrow 1$

$$\text{Tr}(T_{\text{fun}}^A T_{\text{fun}}^B) = T_F \delta_{AB}$$

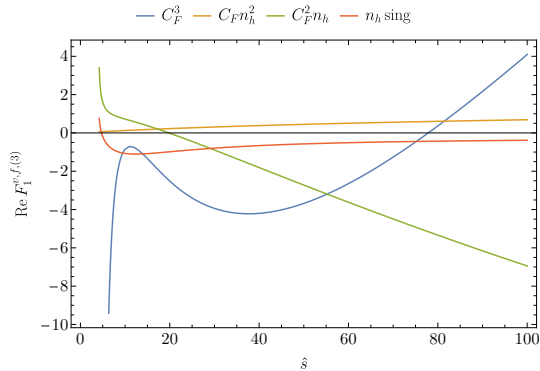
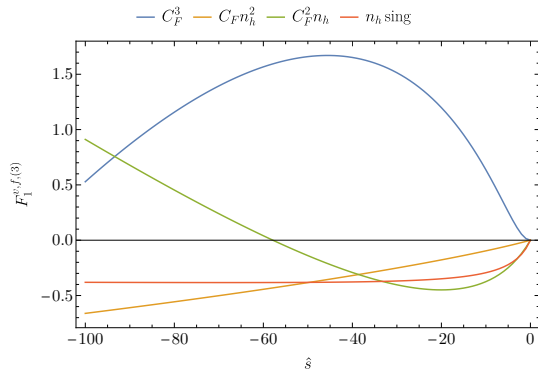
- Number of massless quarks: $n_l \rightarrow 0$ (or massification).

- Massive quark (equal internal and external): $n_h \rightarrow 1$.

- $\alpha_s^{(n_l)}(\mu_s) \rightarrow \bar{\alpha}(\mu_s)$.

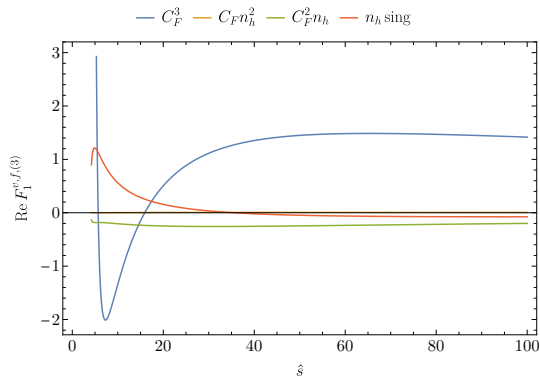
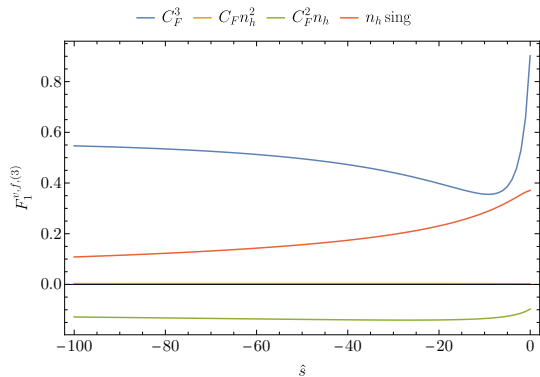
on-shell $\overline{\text{MS}}$ relation for QED: [Baikov, Chetyrkin, Kuhn, Sturm, Nucl.Phys.B 867 \(2013\) 182](#)

Results for QED



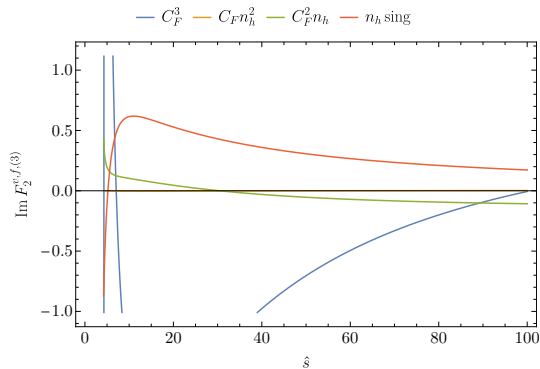
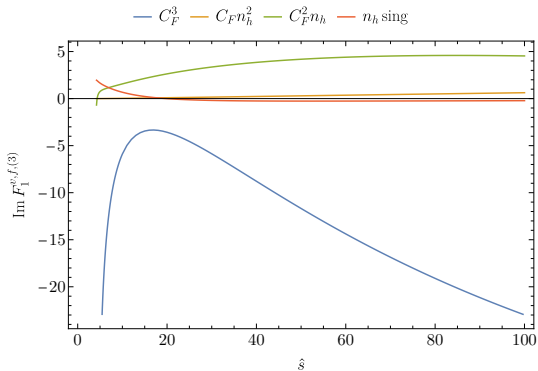
Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

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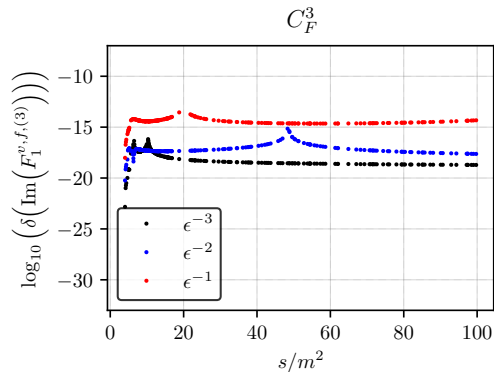
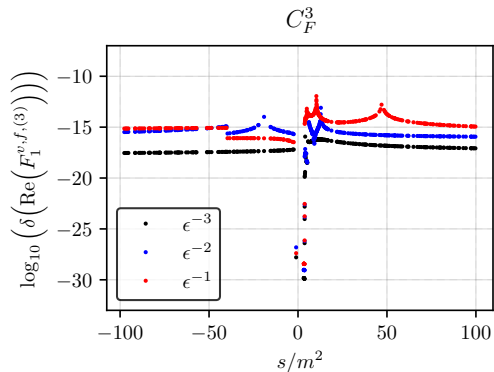
Estimation of the accuracy

- Bare amplitude up to α_s^2 written in term of HPLs (ginac). Exact UV and IR counter terms.
- We estimate the precision from the numerical pole cancellations of the renormalized and infrared-subtracted form factor:

$$\delta(F^{f,(3)}|_{\epsilon^i}) = \frac{F^{(3)}|_{\epsilon^i} + F^{(CT+Z)}|_{\epsilon^i}}{F^{(CT+Z)}|_{\epsilon^i}}$$

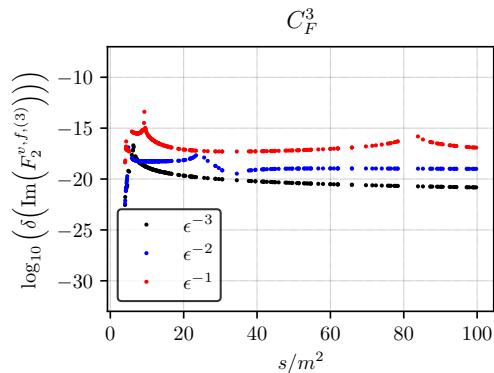
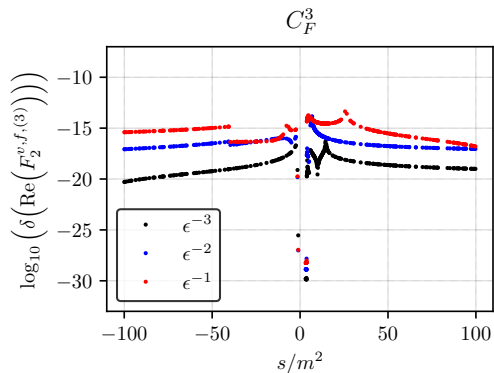
- Close to the zeros of the form factors, the relative precision deteriorates.
- We declare a minimal precision by removing all points whose coefficient size is smaller than 2.5 % of the average in each range
 - $s/m^2 < 0$
 - $0 < s/m^2 < 3.95$
 - $4.05 < s/m^2 < 16$
 - $s/m^2 > 16$
- n_h singlet is finite.

Pole cancellation in QED



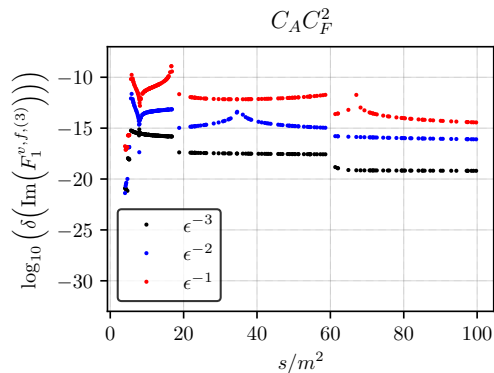
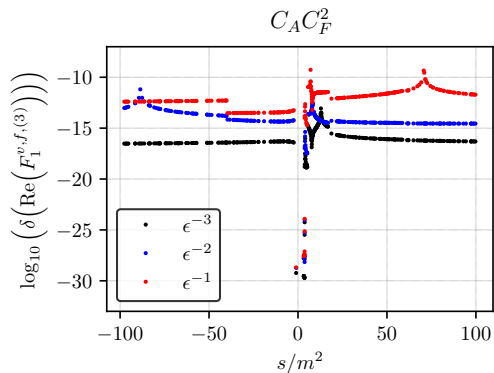
Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

Pole cancellation in QED



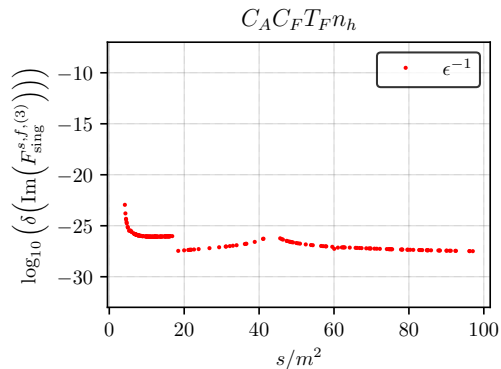
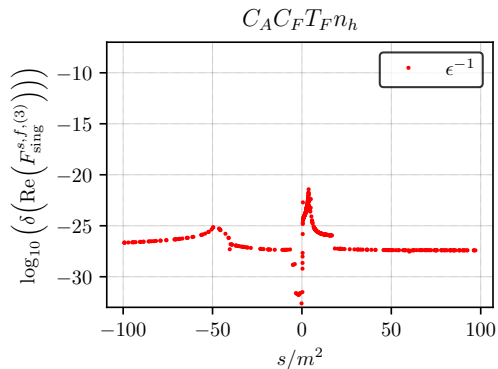
Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

Pole cancellation in QCD



Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

Pole cancellation for singlet contribution



Lange, MF, Schönwald, Steinhauser, Phys.Rev.Lett. 128 (2022) 172003, hep-ph/2207.00027.

Towards N³LO Monte Carlo

Experiment	process	F^ℓ	\sqrt{s} [GeV]	Range	
MUonE	$\mu e \rightarrow \mu e$	e	0.405	$t \in [0.14, -10^{-3}] \text{ GeV}^2$	$t/m_e^2 \in [-5 \times 10^5, -4 \times 10^3]$
		μ	0.405	$t \in [0.14, -10^{-3}] \text{ GeV}^2$	$t/m_\mu^2 \in [-12.5, -0.1]$
BESIII	$ee \rightarrow ll$	e	$\sim 4 \text{ GeV}$	$s = 16 \text{ GeV}^2$	$s/m_e^2 \simeq 6 \times 10^7$
		μ	$\sim 4 \text{ GeV}$	$s = 16 \text{ GeV}^2$	$s/m_\mu^2 \simeq 1400$
		τ	$\sim 4 \text{ GeV}$	$s = 16 \text{ GeV}^2$	$s/m_\tau^2 \simeq 5$
P2	$ep \rightarrow ep$	e	1.08 GeV	$t \in [-13, -5] \times 10^{-3} \text{ GeV}^2$	$t/m_e^2 \in [-5, -2] \times 10^4$

Mathematica Package

- <https://gitlab.com/formfactors3l/formfactors3l>
- What is implemented:
 - Bare and renormalized form factors.
 - Non-singlet and n_h singlet
 - Interpolation grid in the ranges $-40 < s/m^2 < 3.75$ and $4.25 < s/m^2 < 60$.
 - In the remaining regions, it uses special series expansion around $s = \pm\infty$ and $s = 4$.
- Complete QCD color structures. For QED
 - $C_F \rightarrow 1$, denoted by cR
 - $T_F \rightarrow 1$, denoted by I2R
 - $C_A \rightarrow 0$, denoted by cA
 - $n_\ell \rightarrow 0$, no massless lepton.
 - $n_h = 1$, internal and external leptons have the same mass.
- 10^4 random points, single core, in about 120s (full QCD) for each form factor.

- Compute bare form factor F_1 and F_2 at order ϵ^0 at $s/m^2 = -1$:

```
In[] := Get["PATH/FormFactors3l.m"]
```

```
In[] := FormFactorBareNonSing[veF1, 0, -1]
```

```
Out[]:= 77.0506 cA^2 cR+95.0634 cA cR^2+0.467466 cR^3-21.9243 cA cR I2R nh  
-11.5582 cR^2 I2R nh+0.751403 cR I2R^2 nh^2-62.6063 cA cR I2R nI  
-45.5408 cR^2 I2R nI+9.35837 cR I2R^2 nh nI+11.8102 cR I2R^2 nI^2
```

```
In[] := FormFactorBareNonSing[veF2, 0, -1]
```

```
Out[]:= 36.8004 cA^2 cR+21.3251 cA cR^2+31.335 cR^3-10.8947 cA cR I2R nh  
-4.17381 cR^2 I2R nh+0.434916 cR I2R^2 nh^2-26.8267 cA cR I2R nI  
-12.9507 cR^2 I2R nI+3.90787 cR I2R^2 nh nI+4.58439 cR I2R^2 nI^2
```


■ n_h -singlet contribution

```
In[] := FormFactorBareNhSing[veF1, 0, -1]
```

```
Out[]:= -0.075237 d33 nh/nc
```

■ Renormalized form factor

```
In[] := FormFactorRenNonSing[veF1, -1]
```

```
Out[]:= 3.10714 cA^2 cR-3.23413 cA cR^2+0.0144347 cR^3+0.0435081 cA cR I2R nh  
-0.0640418 cR^2 I2R nh-0.0107609 cR I2R^2 nh^2-2.59041 cA cR I2R nl  
+1.02032 cR^2 I2R nl+0.000282528 cR I2R^2 nh nl+0.494057 cR I2R^2 nl^2
```

Towards a compiled C/Fortran version

- Use C: interpolation routines from GNU Scientific Libraries.
- Use Fortran: native support for complex numbers.

Features

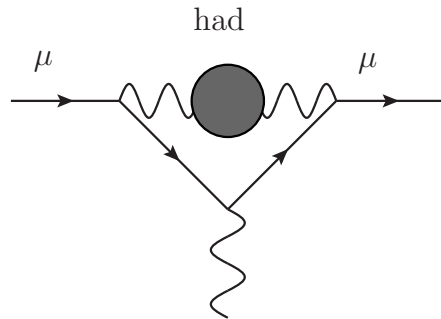
- Should we have a `FormFactors3l.F` or `FormFactors3l_qed.F`?
- For QED we can aim at 10^{-10} relative precision at $O(\epsilon^0)$.
- Implementation of bare and renormalized form factors?
- Standalone implementation of the $s = \pm\infty$ expansion?
- Do you need an exact (but slow) evaluation of the CTs?
- Parallellizable?

NLO Hadronic contribution to the form factors

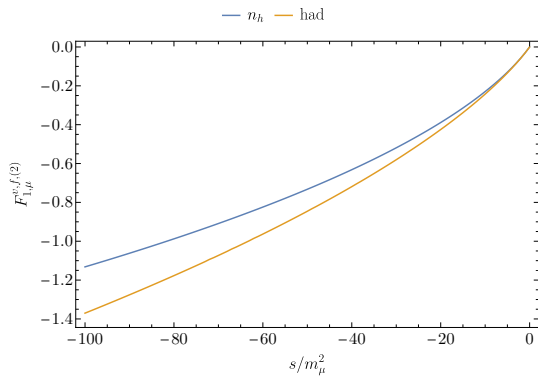
$$F_i^{\ell, \text{vp}}(s) = -\frac{\alpha}{\pi} \int_0^1 dx \Pi \left(\frac{m_\ell^2 x^2}{x-1} \right) f_i(x, s/m_\ell^2)$$

MF, JHEP 02 (2019) 027

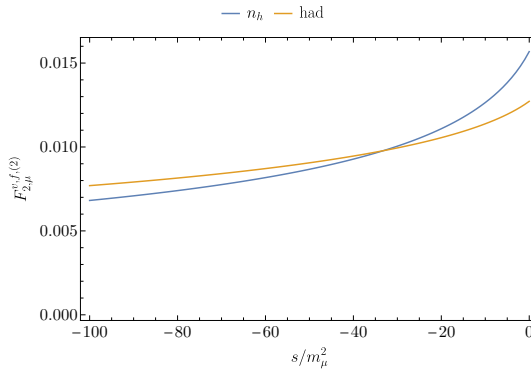
- Are hadronic corrections really needed?
- For electron form factors negligible at low energies.
- For muon, n_h and hadronic corrections are comparable.
- NLO corrections are doable.
- Light-by-light appears at NLO also.



Muon form factors



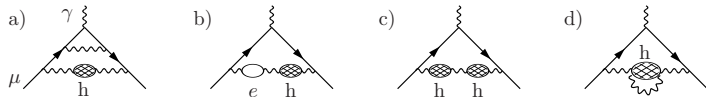
for hadronic vp: Jegerlehner, alpha0EDc17



NLO Hadronic corrections

- Dispersion relation

$$F_i^{\ell, \nu p}(s) = \left(\frac{\alpha}{\pi}\right)^3 \int_{4m_\pi^2}^{\infty} \frac{dz}{z} R(z) K(z, s)$$



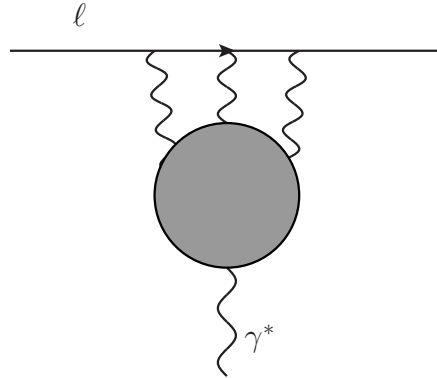
Jegerlehner, Nyffeler, Phys.Rept. 477 (2009) 1

- Diagrams with heavy photon with mass \sqrt{z} .
- Expansion in $m_\ell^2 \ll z$.

Example for a_μ^{HNLO}

$$K^{[(a)]}(s/m^2) = \frac{m^2}{s} \left[\frac{223}{54} - 2\zeta(2) - \frac{23}{36} \ln \frac{s}{m^2} \right] + \frac{m^4}{s^2} \left[\frac{8785}{1152} - \frac{37}{8} \zeta(2) - \frac{367}{216} \ln \frac{s}{m^2} + \frac{19}{144} \ln^2 \frac{s}{m^2} \right] \\ + \frac{m^6}{s^3} \left[\frac{13072841}{432000} - \frac{883}{40} \zeta(2) - \frac{10079}{3600} \ln \frac{s}{m^2} + \frac{141}{80} \ln^2 \frac{s}{m^2} \right] + \dots$$

Light-by-light



Conclusions

- Non-singlet and n_h -singlet contributions to the QED/QCD massive form factors at N³LO.
- Currents: vector, axial-vector, scalar, pseudoscalar.
- Bare form factors are determined as expansions around certain regular and singular kinematical points.
- We provide a Mathematica package with interpolation grids and expansions around $s = \pm\infty$ and $s = 4$.
- Precision (QCD): at least 7 digits.
- Precision (QED), vector form factors: at least 10 digits.
- We provide one of the building blocks for $\gamma^* \rightarrow \ell^+\ell^-$ differential Monte Carlo with dominant N³LO corrections.