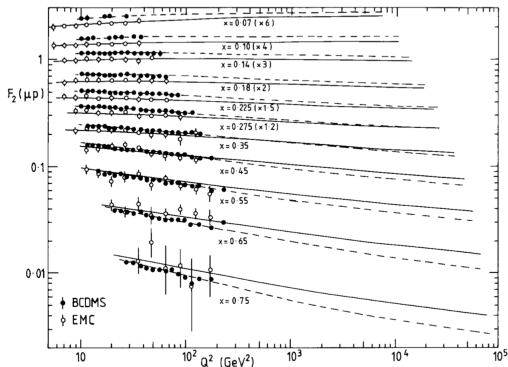


# Lattice Observables for PDFs

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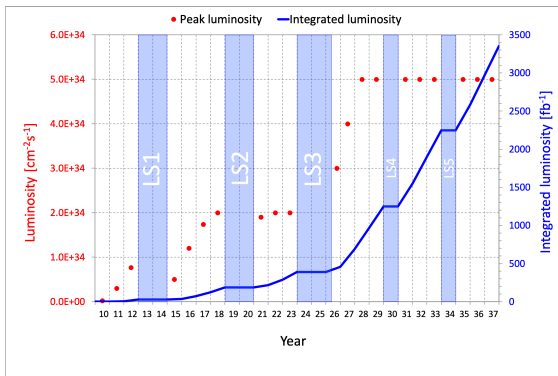


in collaboration with M Wilson, T Giani, C Monahan, K Cichy, J Karpie, K Orginos,  
A Radyushkin, S Zafeiropoulos

# plan

1. PDFs/TMDs/GPDs in QFT
2. lattice observables
3. inverse problems

# run 3/hi-lumi LHC: setting the target for theory predictions



theory predictions need to match the experimental precision

↔ increased *precision and accuracy* of *parton distribution functions*

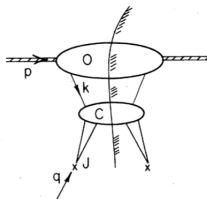
# EIC and strong interactions

- hadron *tomography*
- understanding the nucleon spin
- understanding nucleons/nuclear physics from QCD
- current understanding of nucleon structure:

$$\mathcal{W}(x, k_{\perp}, b_{\perp}) = \int \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{-i\Delta_{\perp} \cdot b_{\perp}} \frac{1}{2} \int \frac{dz^- d^2 z_{\perp}}{(2\pi)^3} e^{i[(xp^+)z^- - k_{\perp} \cdot z_{\perp}]}$$
$$\times \langle p + \frac{\Delta_{\perp}}{2} | \bar{\psi}(-z/2) \Gamma \lambda_A \mathcal{U} \psi(z/2) | p - \frac{\Delta_{\perp}}{2} \rangle$$

# DIS hadronic tensor – factorization

$$H_{\mu\nu} = \int d^D y e^{iq \cdot y} \langle p | J_\mu(y) J_\nu(0) | p \rangle$$



$$H_{\mu\nu} = F_1(x, Q^2) \left( \frac{q_\mu q_\nu}{q^2} - g_{\mu\nu} \right) + F_2(x, Q^2) \left( p + \frac{1}{2x} q \right)_\mu \left( p + \frac{1}{2x} q \right)_\nu$$

expressed in terms of PDFs – physical, finite quantities

$$F_i(x, Q^2) = \int_x^1 \frac{d\xi}{\xi} C_i(\xi, Q^2, \mu^2) f_R(x/\xi, \mu^2) + \mathcal{O}(1/Q^2)$$

$$f(x) = \int \frac{dz^-}{2\pi} e^{i(xp^+)z^-} \langle p | \bar{\psi}(-z^-/2) \Gamma \lambda_A \mathcal{U} \psi(z^-/2) | p \rangle$$

# lattice QCD for PDFs?

- cannot compute light-cone quantities in Euclidean field theory
- Xi. Ji (2013): compute the spatial correlator

[ji 20]

$$\langle p | \bar{\psi}(-z_3/2) \Gamma \lambda_A \mathcal{U} \psi(z_3/2) | p \rangle$$

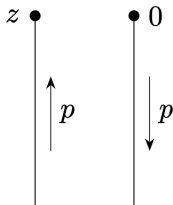
- spun a lot of activity in LQCD community: **quasi-PDF**, **pseudo-PDF**, **loffe Time Distributions** – see recent reviews at Lattice conferences and dedicated workshops

[constantinou et al 20]

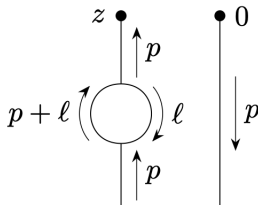
- a factorization formula relates correlators in the spatial direction with the light-cone quantities (after renormalization)

# toy-model computation

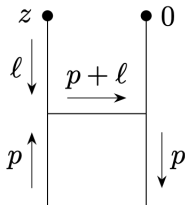
loffe time distribution  $\widehat{\mathcal{M}}(\nu, z^2) = \langle p | \phi(z) \phi(0) | p \rangle$



(a)



(b)



(c)

at tree level

$$\widehat{\mathcal{M}}_a(\nu, z^2) = \exp[-ip \cdot z] = \exp[-i\nu] = \widehat{\mathcal{M}}^{(0)}(\nu, 0)$$



## one-loop in MSbar

- light-cone separation – FT of PDFs

$$\begin{aligned}\widehat{\mathcal{M}}_R(\nu; \mu^2) &= \left[ 1 + \frac{\alpha}{6} \left( \log \frac{m^2}{\mu^2} + b \right) \right] \widehat{\mathcal{M}}^{(0)}(\nu, 0) \\ &+ \alpha \int_0^1 dx (1-x) \log \frac{\mu^2}{m^2(1-x+x^2)} \widehat{\mathcal{M}}^{(0)}(x\nu, 0)\end{aligned}$$

- spatial separation – FT of quasi-PDFs

$$\begin{aligned}\widehat{\mathcal{M}}_R(\nu, z_3^2; \mu^2) &= \left[ 1 + \frac{\alpha}{6} \left( \log \frac{m^2}{\mu^2} + b \right) \right] \widehat{\mathcal{M}}^{(0)}(\nu, 0) \\ &+ \alpha \int_0^1 dx (1-x) 2K_0(mz_3) \widehat{\mathcal{M}}^{(0)}(x\nu, 0)\end{aligned}$$

# factorization theorem

for  $mz_3 \ll 1$  (small distances, large momenta)

$$2K_0(Mz_3) = -\log(m^2 z_3^2) + 2\log(2e^{-\gamma_E}) + \mathcal{O}(m^2 z_3^2)$$

and therefore

$$\widehat{\mathcal{M}}_R(\nu, -z_3^2; \mu^2) = \int_{-1}^1 d\xi \tilde{C}(\xi\nu, \mu^2 z_3^2) \hat{f}_R(\xi, \mu^2) + \mathcal{O}(m^2 z_3^2)$$
$$\tilde{C}(\xi\nu, \mu^2 z_3^2) = e^{i\xi\nu} - \alpha \int_0^1 dx (1-x) \log\left(\mu^2 z_3^2 \frac{e^{2\gamma_E}}{4}\right) e^{ix\xi\nu}$$

Wilson coefficient is IR-safe

$z^2$  dependence only at  $\mathcal{O}(\alpha)$

# QCD matrix elements

$$M_{\Gamma,A}(z) = \bar{\psi}(z) \Gamma \lambda_A P \exp \left( -ig \int_0^z d\eta A(\eta) \right) \psi(0)$$

Ioffe time distributions

$$\mathcal{M}_{\gamma^\mu,A}(z, P) = \langle P | \mathcal{M}_{\gamma^\mu,A}(z) | P \rangle$$

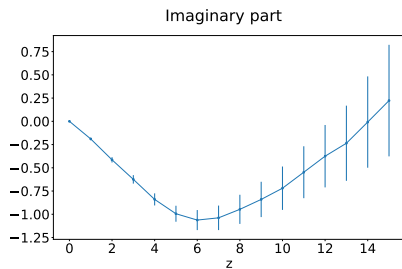
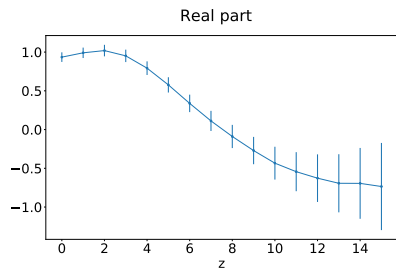
Lorentz covariance

$$\mathcal{M}_{\gamma^\mu,A}(z, P) = P^\mu h_{\gamma^\mu,A}(z \cdot P, z^2) + z^\mu h'_{\gamma^\mu,A}(z \cdot P, z^2)$$

# lattice observables - 1

$$\mathcal{O}_{\gamma^0}^{\text{Re}}(zP_z, z^2) \equiv \text{Re} [h_{\gamma^0,3}(zP_z, z^2)]$$

$$\mathcal{O}_{\gamma^0}^{\text{Im}}(zP_z, z^2) \equiv \text{Im} [h_{\gamma^0,3}(zP_z, z^2)]$$



[C Alexandrou et al 18]

# systematic errors

- cut-off effects
- finite volume effects
- excited states contamination
- truncation effects
- higher-twist terms
- isospin breaking

↔ **renormalized lattice observables at the continuum limit**

lattice QCD yields a discrete set of datapoints to be included in PDF fits

# stochastic solution of inverse problems

unknown model:  $\theta \in \Theta$

Data:  $Y = K\theta + \varepsilon$ ,     where:  $\varepsilon \sim \pi_\varepsilon = \mathcal{N}(0, \Sigma)$

$\theta \in \Theta \implies \text{Prob}(y|\theta) = \pi_\varepsilon(y - K\theta)$

## Bayes theorem

$$\text{Prob}(\theta|y) = \frac{\text{Prob}(y|\theta)\text{Prob}(\theta)}{\text{Prob}(y)}$$

choose a parametrization, find the posterior distribution

# replicas

data:  $y_i, i = 1, \dots, N_{\text{dat}}$  are treated as stochastic variables  $\sim \mathcal{N}(Y, C)$

bootstrap: replicas simulate the fluctuations of  $y$

$$y^{(k)} = Y + \varepsilon^{(k)}, \quad k = 1, \dots, N_{\text{rep}}$$

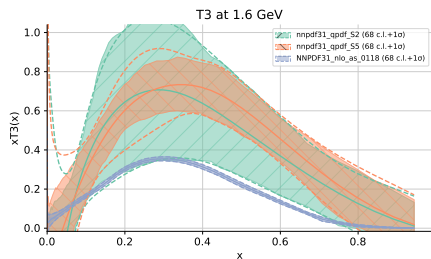
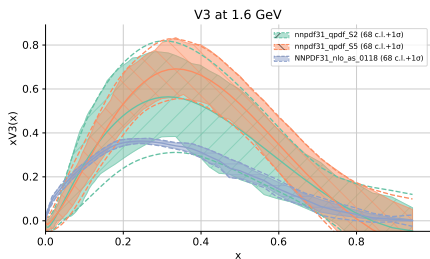
for each replica we minimize:

$$\chi^{2(k)}[\theta] = \frac{1}{N_{\text{dat}}} \sum_{ij} \left( g[\theta]^{(k)} - y^{(k)} \right)_i C_{ij}^{-1} \left( g[\theta]^{(k)} - y^{(k)} \right)_j + \text{priors}$$

$\{\theta^{(k)}, k = 1, \dots, N_{\text{rep}}\}$  yields the probability distribution in model space

# fit results

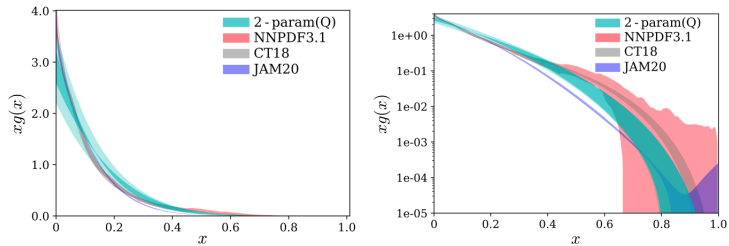
lattice data can be included in NNPDF fits, just like any other data!



[Idd et al 19, using ETMC data]

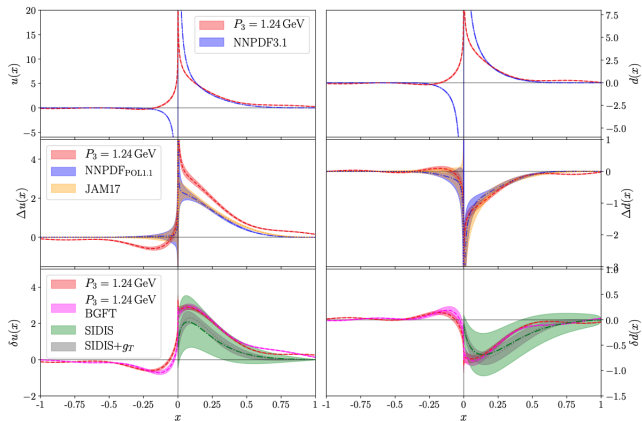


# HadStruc



$$m_{\pi} = 358 \text{ MeV}$$

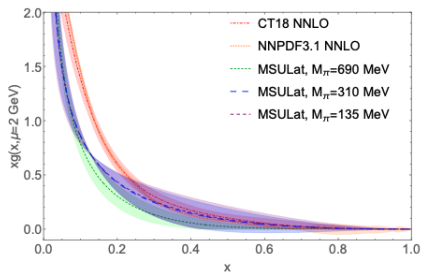
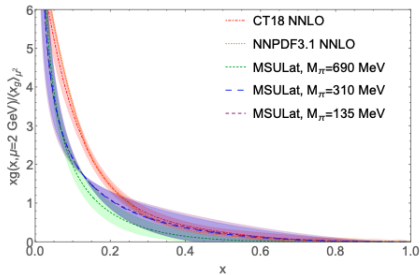
[Khan et al 21]



$$m_\pi = 260 \text{ MeV}$$

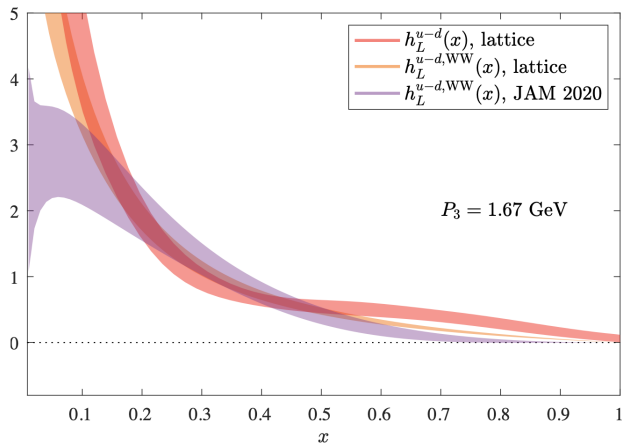
[Alexandrou et al 21]

# MSULat



[Fan et al 21]

# twist-3 and Wandzura-Wilczek



$$m_\pi = 260 \text{ MeV}$$

[Bhattacharya et al 21]

# outlook

- light-cone PDFs + factorization describe the structure of the proton
- current extraction from data is very precise + improving
- lattice quantities needs to be properly renormalised and extrapolated to the continuum limit  
↔ are on the same footing as experimental data
- lattice data provide complementary information, can be included in global fits like any other data (inverse problem)
- identify the areas where a significant phenomenological impact from lattice QCD is possible
- GPDs for EIC