



Baryons as a Probe for Hadronization

Stefan Kiebacher in collaboration with Stefan Gieseke and Simon Plätzer | September 4, 2023





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Open questions:

- How and with which kinematics are Baryons produced in high energy collisions?
- Which stages of the hadronization impact the kinematics and multiplicity of Baryons?
- What knowledge can we extract from Baryon observables e.g. their Angular Correlations?

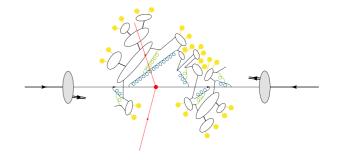








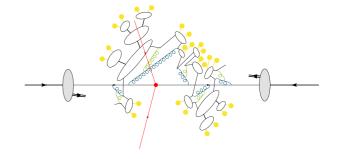
Primordial cluster formation







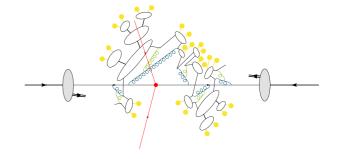
- Primordial cluster formation
- Colour Reconnection







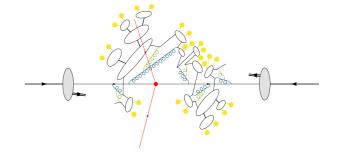
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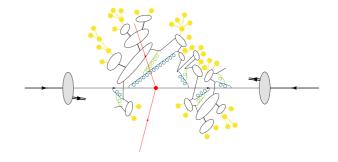
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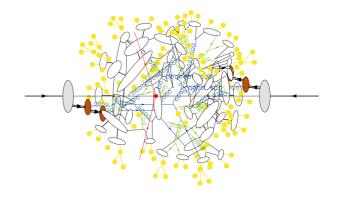
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- Multiple Parton Interactions (MPIs)





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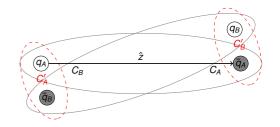


- 1. Describe current Hadronization model in Herwig
- 2. Show the resulting baryon angular correlations
- 3. Present new model for Colour Reconnection and some changes to the Cluster Fission
- 4. Work in Progress: New structure for systematic Kinematics for the Cluster Fission and Decay

Colour Reconnection (CR)



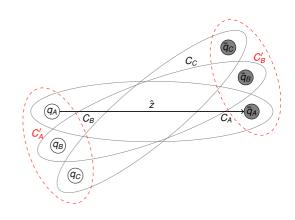
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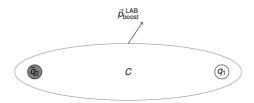


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- Find next to maximal y_{sum} cluster to make baryon-antibaryon pair
 - ⇒ Baryonic Colour Reconnection (BCR) accepted with probability P_B [Gieseke, Kirchgaeßer, and Plätzer 2018]
- Note: Clusters can be light for Baryon Production However a lot of multiplicity is needed!



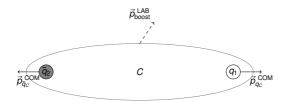


• Fission all clusters $M>M_{\max}(q_1,\bar{q}_2)$ above a threshold $M_{\max}(q_1,\bar{q}_2)$ recursively



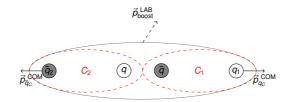


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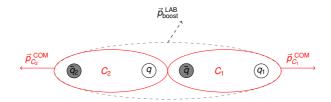


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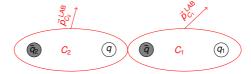


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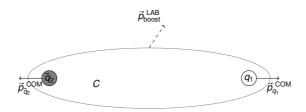


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- 4. Boost C_1 , C_2 back into the lab frame



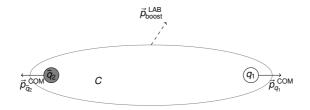


Clusters decay to two hadrons



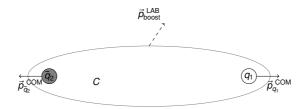


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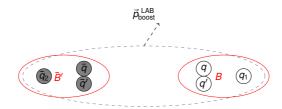


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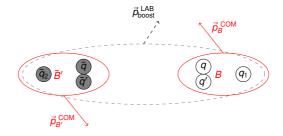


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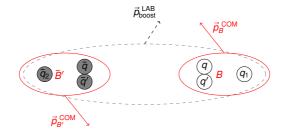


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- Note: Not much multiplicity, but high-mass clusters are needed to produce Baryons!



Baryon Angular Correlations



- Depletion of near-sided baryons only reproduced by Baryonic Colour Reconnection (BCR)
- Cluster Decay (CD)
 baryons are giving opposite features to data
- Cluster Decay baryons are solely responsible for unphysical far-side peak
- BCR alone cannot produce enough baryons especially for low multiplicity events (e.g. at LEP)

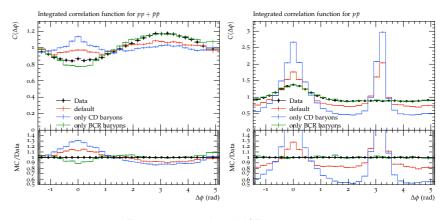


Figure: measured by ALICE [Adam et al. 2017]

Baryon Angular Correlations



Solutions:

 Disable CD baryon production

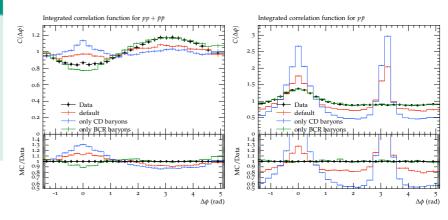


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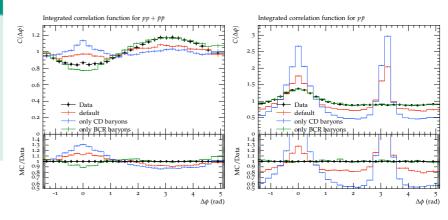


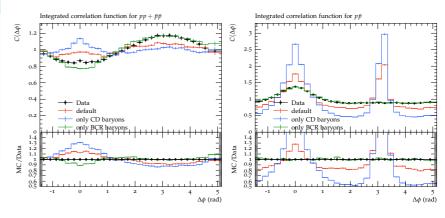
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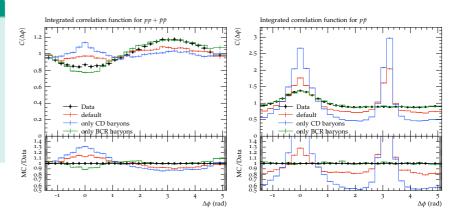


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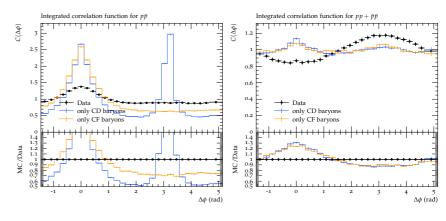
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- 3. New Diquark Colour Reconnection algorithm



Cluster Fission (CF) vs Cluster Decay (CD) Baryons



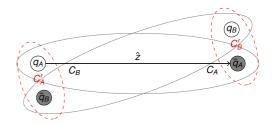
- Near-side depletion not reproduced
 - ⇒ CD and CF are oblivious to other baryons
- Far-side peak is completely gone!
- Near-side still overshoots the data \Rightarrow but this is only one mechanism







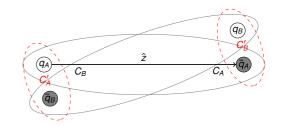
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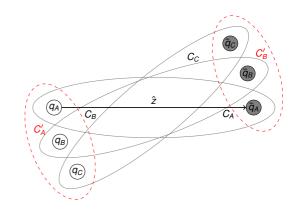


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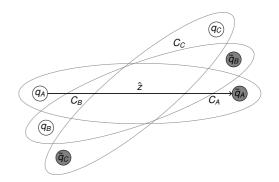
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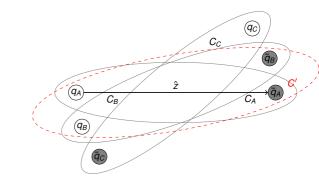


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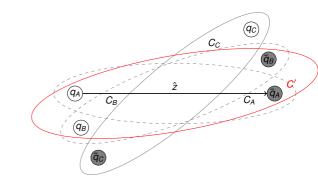


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 ⇒ Diquark Colour Reconnection (DCR) accepted with probability P_D



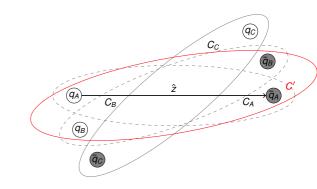


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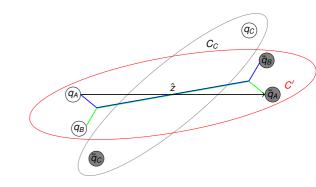


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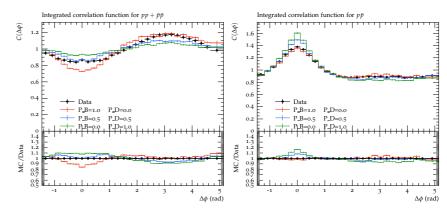
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- Similar to Pythia's String Junction Colour Reconnections [Christiansen and Skands 2015]



Diquark Colour Reconnection



- Purely Diquark-type CR with P_D = 1 has not enough depletion for pp correlations
- Near-sided peak reproduced!
- No far-sided peak for pp
 DCR and good
 phenomenology







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 - \Rightarrow Would like angular correlations of baryons from e^+e^- , where BCR is expected to be small!

Revisiting Cluster Fission Kinematics

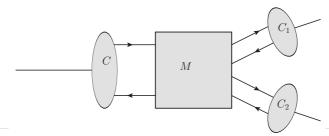


Idea: Cluster Fission is a partonic 2 \rightarrow 4 process [Plätzer 2023].

■ Factorize the process $C(p_i, p_i) \to C_1(q_i, q)$, $C_2(q_i, \bar{q})$ (see Jan Priedigkeit's Bachelor thesis Graz):

$$d\Gamma(C \to C_1, C_2) = \int d^4 \Phi_{q_i} d^4 \Phi_{q} d^4 \Phi_{\bar{q}} d^4 \Phi_{q_j} (2\pi)^4 \delta^4(p_i + p_j - q_i - q - \bar{q} - q_j) |\mathcal{M}(p_i, p_j \to q_i, q, \bar{q}, q_j)|^2$$
(1)

$$d\Gamma(C \to C_1, C_2) = \int dM_1 dM_2 d\Phi_2(P|Q_1, Q_2) d\Phi_2(Q_1|q_i, q) d\Phi_2(Q_2|q_j, \bar{q}) |\mathcal{M}(p_i, p_j \to q_i, q, \bar{q}, q_j)|^2$$
(2)







Revisiting Cluster Fission Kinematics

• (Pre-)Sample Masses M_1 , M_2 from flat Phase Space weight (Jan Priedigkeit's Bachelor thesis Graz) $d\Phi_4 \propto dM_1 dM_2 \sqrt{\lambda(M,M_1,M_2)} \sqrt{\lambda(M_1,m_1,m)} \sqrt{\lambda(M_2,m_2,m)}/(M_1M_2)^2$





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- Rejection sampling of soft $q\bar{q}$ emission diagram, which in the soft limit is given in [Catani and Grazzini 2000] [Plätzer 2022] (up to colour factors) by:





- (Pre-)Sample Masses M_1, M_2 from flat Phase Space weight (Jan Priedigkeit's Bachelor thesis Graz) $d\Phi_4 \propto dM_1 dM_2 \sqrt{\lambda(M, M_1, M_2)} \sqrt{\lambda(M_1, m_1, m)} \sqrt{\lambda(M_2, m_2, m)}/(M_1 M_2)^2$
- Rejection sampling of soft $q\bar{q}$ emission diagram, which in the soft limit is given in [Catani and Grazzini 2000] [Plätzer 2022] (up to colour factors) by:

$$|\mathcal{M}(p_i, p_j \to q_i, q, \bar{q}, q_j)|^2 \propto \frac{2(q_i \cdot q_j)(q\bar{q}) + [q_i \cdot (q - \bar{q})][q_j \cdot (q - \bar{q})]}{2(q \cdot \bar{q})^2[q_i \cdot (q + \bar{q})][q_j \cdot (q + \bar{q})]}$$
(3)



Used the angular correlations of baryons to examine the kinematics of the cluster model





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Outlook for hadronization in Herwig:



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Outlook for hadronization in Herwig:

Dynamic gluon constituent masses (WIP by Daniel Samitz, S. Plätzer)



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Outlook for hadronization in Herwig:

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- Restructure the CF and CD to implement flexible kinematics (WIP with S. Plätzer, S. Gieseke)



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- Make Colour Reconnection dynamic via soft gluon evolution [Gieseke, Kirchgaeßer, Plätzer, and Siodmok 2018; Plätzer 2023] (WIP with S. Plätzer, S. Gieseke) ⇒ reduce the free parameters by 2



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Long term goals for hadronization in Herwig:



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Long term goals for hadronization in Herwig:

Make hadronization model independent of the shower cutoff



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- Restructure the CF and CD to implement flexible kinematics (WIP with S. Plätzer, S. Gieseke)
- Make Colour Reconnection **dynamic** via soft gluon evolution [Gieseke, Kirchgaeßer, Plätzer, and Siodmok 2018; Plätzer 2023] (WIP with S. Plätzer, S. Gieseke) ⇒ reduce the free parameters by 2

Long term goals for hadronization in Herwig:

- Make hadronization model independent of the shower cutoff
- Dynamic hadronization model for generalisation to dark hadrons (by Simon Plätzer, Dominic Stafford et al.)



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Outlook for hadronization in Herwig:

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Long term goals for hadronization in Herwig:

- Make hadronization model independent of the shower cutoff
- Dynamic hadronization model for generalisation to dark hadrons (by Simon Plätzer, Dominic Stafford et al.)
- Convince experimentalists to get more (identified) particle correlation data (also for LEP), since important modelling input





TLDL: Lots of construction sites in the Hadronization model in Herwig ...



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- (2023). "Colour evolution and infrared physics". In: <u>JHEP</u> 07, p. 126. DOI: 10.1007/JHEP07(2023)126. arXiv: 2204.06956 [hep-ph].

Thank You For Your Attention!

Questions? Remarks? Comments?

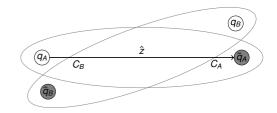


Boost in cluster rest frame



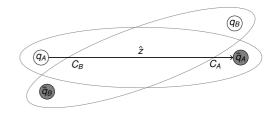


- Boost in cluster rest frame
- Select next cluster at random



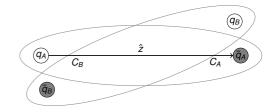


- Boost in cluster rest frame
- Select next cluster at random
- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} ⇒ find maximal y_{sum}



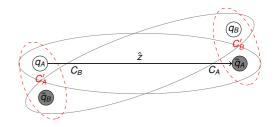


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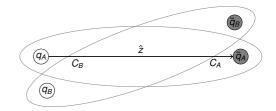


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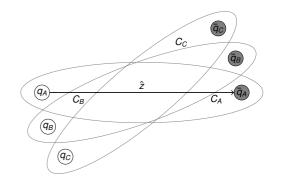


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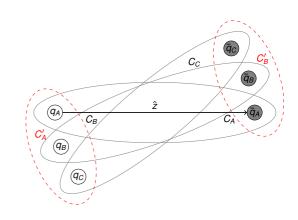


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- Find next to maximal y_{sum} cluster to make baryon-antibaryon pair



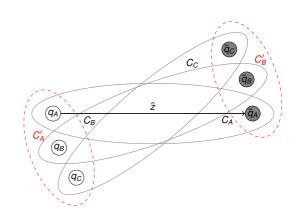


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- Find next to maximal y_{sum} cluster to make baryon-antibaryon pair
 - \Rightarrow Baryonic Colour Reconnection (BCR) accepted with probability P_B



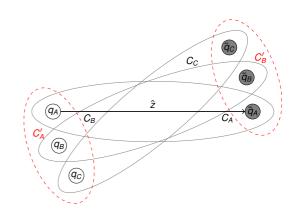


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- Find next to maximal y_{sum} cluster to make baryon-antibaryon pair
 - ⇒ Baryonic Colour Reconnection (BCR) accepted with probability *P_R*
- Note: Clusters can be light for Baryon Production



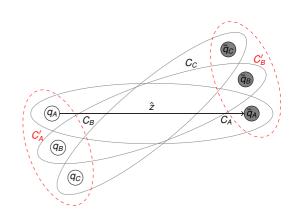


- Boost in cluster rest frame
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- Note: Clusters can be light for Baryon Production However a lot of multiplicity is needed!





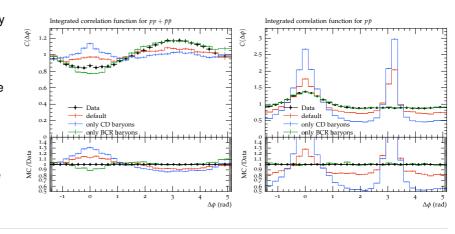
- Boost in cluster rest frame
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 - \Rightarrow Baryonic Colour Reconnection (BCR) accepted with probability P_B
- Note: Clusters can be light for Baryon Production However a lot of multiplicity is needed!
- In fact BCR regulated the over-abundance of high multiplicity events [Gieseke, Kirchgaeßer, and Plätzer 2018]



Backup: Baryon Angular Correlations



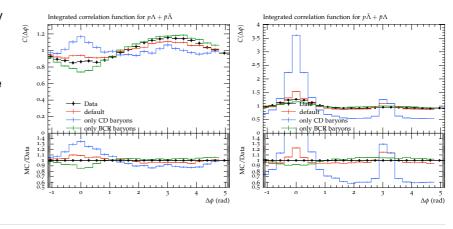
- Depletion of near-sided baryons only reproduced by Baryonic Colour Reconnection (BCR)
- Cluster Decay (CD)
 baryons are giving opposite features to data
- Cluster Decay baryons are solely responsible for unphysical far-side peak
- BCR alone cannot produce enough baryons especially for low multiplicity events (e.g. at LEP)



Backup: Baryon Angular Correlations



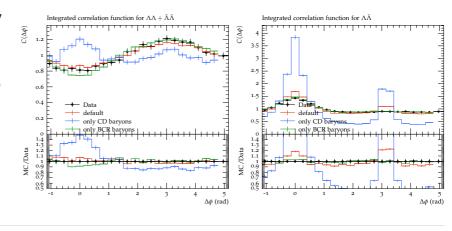
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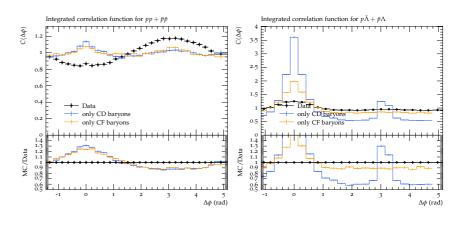
- Depletion of near-sided baryons only reproduced by Baryonic Colour Reconnection (BCR)
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 baryons are giving opposite features to data
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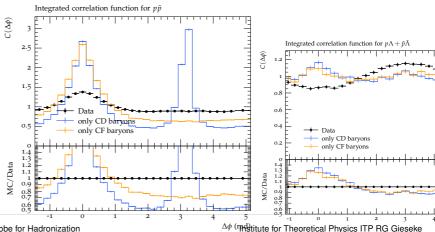
- CD baryon mechanism vs new Cluster Fission (CF) mechanism
- Near-side depletion not reproduced
 - \Rightarrow CD and CF are oblivious to other baryons







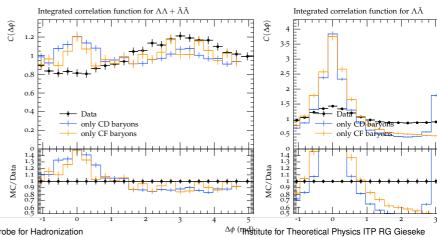
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Boost in cluster rest frame

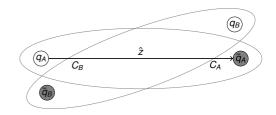




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Backup: Diquark Colour Reconnection Algorithm

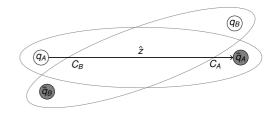
- Boost in cluster rest frame
- Select next cluster at random







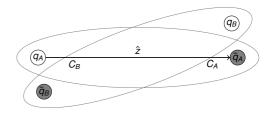
- Boost in cluster rest frame
- Select next cluster at random
- Compute $y_{sum} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} ⇒ find maximal y_{sum}







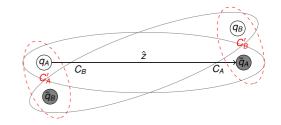
- Boost in cluster rest frame
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- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} ⇒ find maximal y_{sum}
- If $y_{q_B} > 0$ and $y_{\bar{q}_B} < 0$ for $y_{\text{sum}}^{\text{max}}$







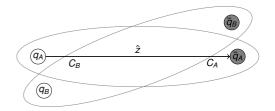
- Boost in cluster rest frame
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- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} ⇒ find maximal y_{sum}
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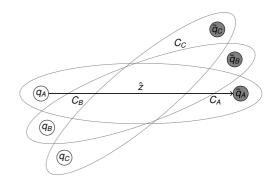
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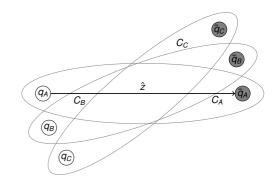
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- If $y_{q_B} < 0$ and $y_{\bar{q}_B} > 0$ for $y_{\text{sum}}^{\text{max}}$
- Find next to maximal $y_{\text{sum}}^{\text{max,2}}$ cluster







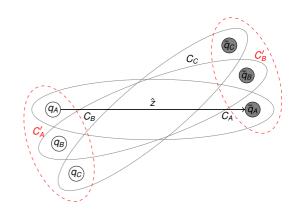
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- Find next to maximal $y_{\text{sum}}^{\text{max,2}}$ cluster
- If $y_{q_R} < 0$ and $y_{\bar{q}_R} > 0$ for $y_{\text{sum}}^{\text{max},2}$







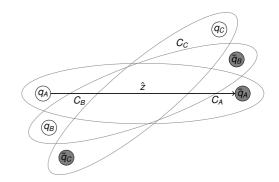
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- If y_{qB} < 0 and y_{qB} > 0 for y_{sum}^{max,2}
 ⇒ Baryonic Colour Reconnection (BCR) accepted with probability P_B







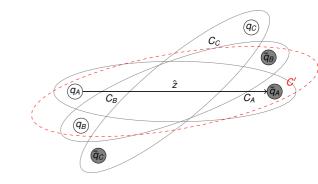
- Boost in cluster rest frame
- Select next cluster at random
- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} ⇒ find maximal y_{sum}
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- Find next to maximal $y_{\text{sum}}^{\text{max,2}}$ cluster
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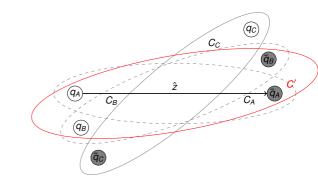
- Boost in cluster rest frame
- Select next cluster at random
- Compute $y_{\text{sum}} = |y_{q_B}| + |y_{\bar{q}_B}|$ with respect to \hat{z} \Rightarrow find maximal y_{sum}
- If $y_{q_B} < 0$ and $y_{\bar{q}_B} > 0$ for $y_{\text{sum}}^{\text{max}}$
- Find next to maximal $y_{\text{sum}}^{\text{max,2}}$ cluster
- If y_{q_B} > 0 and y_{q̄_B} < 0 for y_{sum}^{max,2}
 ⇒ Diquark Colour Reconnection (DCR) accepted with probability P_D







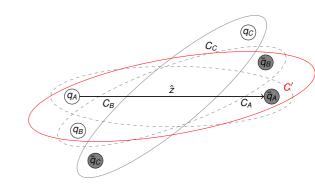
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- Find next to maximal $y_{\text{sum}}^{\text{max,2}}$ cluster
- If $y_{q_B} > 0$ and $y_{\bar{q}_B} < 0$ for $y_{\text{sum}}^{\text{max},2}$ \Rightarrow Diquark Colour Reconnection (DCR) accepted with probability P_D if $M_{C'} > M_{\text{Baryon Pair}}^{\text{Lightest}}$







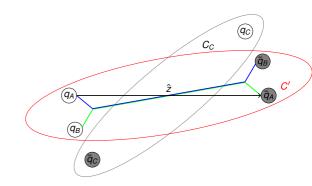
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- Note: Mixed need for multiplicity <u>and</u> existing mass for producing baryons



Backup: Diquark Colour Reconnection Algorithm



- Boost in cluster rest frame
- Select next cluster at random
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- If $y_{q_B} > 0$ and $y_{\bar{q}_B} < 0$ for $y_{\text{sum}}^{\text{max,2}}$ \Rightarrow Diquark Colour Reconnection (DCR) accepted with probability P_D if $M_{C'} > M_{\text{Baryon Pair}}^{\text{Lightest}}$
- Note: Mixed need for multiplicity <u>and</u> existing mass for producing baryons
- Similar to Pythia's String Junction Colour Reconnections [Christiansen and Skands 2015]



Backup: Spectra of Protons



■ Proton p_T-spectra are badly modelled

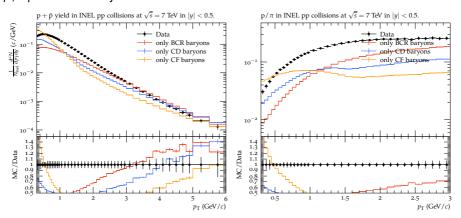


Figure: Compare p_T – spectra of p for only BCR, only CD or only CF baryon mechanisms [Adam et al. 2015]

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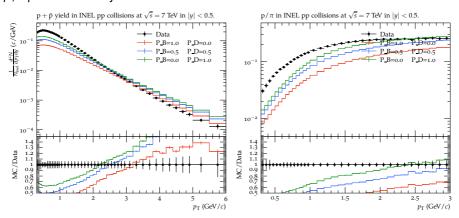


Figure: Compare p_T – spectra of p for only new DCR baryon mechanism with different probabilities [Adam et al. 2015]

Backup: Tuning



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1. Perform a dedicated tune to LEP multiplicities, event shapes and momentum spectra for CF, CR parameters



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Possible solutions: Use of a different "Loss function" than regular χ^2 e.g. $\chi^2 \to \frac{\chi^2}{1+\chi^2}$ or $\tanh(\chi^2)$

Backup: Cluster Fission Details



- A cluster of mass M is fissioned if $M^{\text{Cl}_{pow}} \geq \text{Cl}_{max}^{\text{Cl}_{pow}} + (m_1 + m_2)^{\text{Cl}_{pow}}$, where m_1, m_2 are the masses of the constituents of the cluster
- Currently masses are sampled as follows, where $r_1, r_2 \in [0, 1]$ are uniform random numbers [Bahr et al. 2008]:

$$M_1 = m_1 + (M - m_1 - m_q)r_1^{\frac{1}{P_{\text{split}}}}$$
 (4)

$$M_2 = m_2 + (M - m_2 - m_q) r_2^{\frac{1}{P_{\text{split}}}}$$
 (5)

- Reject samples if $M_1 + M_2 > M$
- Problems: huge dependence on parameters Cl_{max} and especially P_{split}
- Work in progress: Sample masses according to phase space

Backup: Angular Correlations



- The shown plots are showing correlations integrated in $\Delta \eta$ up to $\Delta \eta_{\text{max}} = 1.3$
- The angular correlations are measured via the event mixing [Adam et al. 2017]:

$$C_i(\Delta\phi, \Delta\eta) = \frac{S(\Delta\phi, \Delta\eta)}{B(\Delta\phi, \Delta\eta)} \tag{6}$$

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(7)

$$B_i(\Delta\phi, \Delta\eta) = \frac{1}{N_{\text{pairs}}^{\text{mixed}}} \frac{d^2 N_{\text{pairs}}^{\text{mixed}}}{d\Delta\eta d\Delta\phi}$$
(8)

$$C_i(\Delta\phi) = \int_0^{\Delta\eta_{\text{max}}} C_i(\Delta\phi, \Delta\eta) d\Delta\eta$$
 (9)

Spectra of Protons



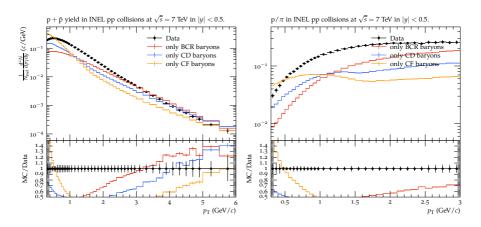


Figure: Compare p_T — spectra of p for only BCR, only CD or only CF baryon mechanisms [Adam et al. 2015]

Spectra of Protons



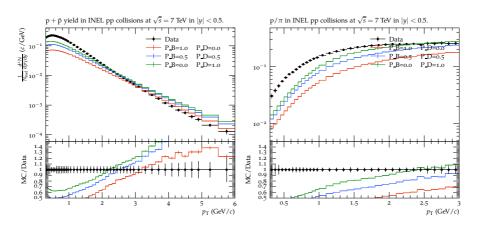


Figure: Compare p_T – spectra of p for only new DCR baryon mechanism with different probabilities [Adam et al. 2015]

Spectra of Strange Baryons



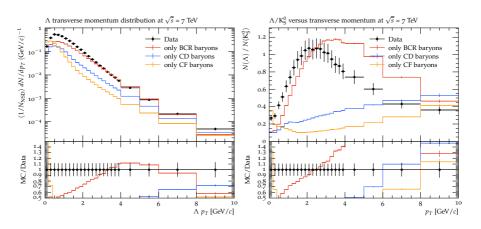


Figure: Compare p_T – spectra of Λ , Ξ for only BCR, only CD or only CF baryon mechanisms [Khachatryan et al. 2011]





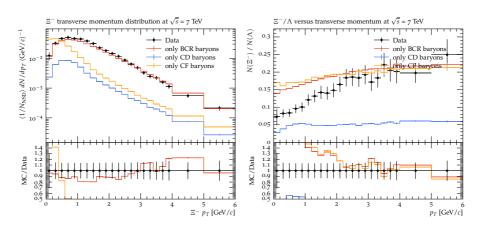


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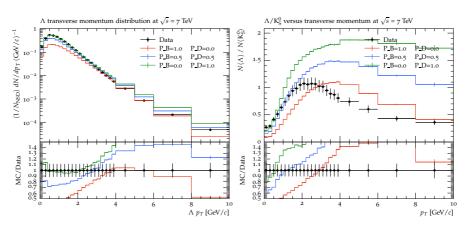


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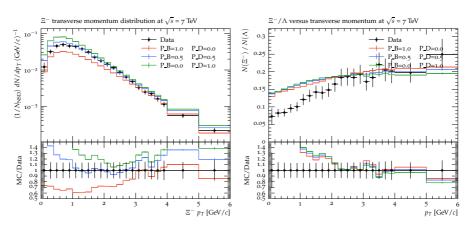


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Consistent Two Particle Boost



- If we boost a two particle system $P = (p_i + p_j)$ into its rest frame $\hat{P} = (\hat{p}_i + \hat{p}_j)$ one needs to be careful to tranform the relative momentum correctly $\hat{P}_{rel} = (\hat{p}_i \hat{p}_i)$
- The naive transformation would be to just use $\Lambda_{(-P)}$, but this would give in general $\hat{P}_{rel} = (\hat{p}_i \hat{p}_j + 2k)$, because $\Lambda \hat{p}_i = p_i + k$ and $\Lambda \hat{p}_j = p_j k$.
- Intuitively the momentum P is completely oblivious to its components and therefore Λ must depend on both the consituents p_i, p_i
- Want a Lorentz Tranformation (matrix or tensor) $\Lambda(p_i, p_j | \hat{p}_i, \hat{p}_j)$ such that $\Lambda \hat{p}_i = p_i$ and $\Lambda \hat{p}_j = p_j$
- Found solution for $\Lambda(p_i, p_j | \hat{p}_i, \hat{p}_j)$, but numerically not very easy
- lacktriangle Work in Progress: Tensor for this trafo $\Lambda^{
 u}_{\ \mu}$