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# Signatures of primordial black hole dark matter at DUNE and THEIA

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Credits: SXS Lensing, nasa.gov

## Primordial black hole dark matter

Primordial black holes are formed from collapse of primordial inhomogeneities in the early Universe before the BBN epoch. They constitute one of the earliest proposed and appealing DM candidates.



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Hawking (1971), Zeldovich, Novikov (1966), Y. B. . N.

### Primordial Black Holes evaporation

Black holes are thought to evaporate via Hawking radiation, emitting particles near their event horizons.

► Very light non-rotating (maximally rotating) PBHs with  $M_{PBH} < 5$  (7) × 10<sup>14</sup> g are expected to have completely evaporated by now. D. N. Page (1976), MacGibbon+ (2008)

► A large number of constraints have been derived on the abundance of light PBH in the  $10^{15}$ – $10^{17}$  g mass range, from the non-observation of their evaporation products:  $\gamma$  rays, electrons/positrons and neutrinos.

Further constraints: from CMB and BBN, the EDGES measurements of 21cm absorption, or the heating of interstellar medium.

Carr (1976), Lehoucq+ (2009), Wright (1996), Ballesteros+ (2019), Laha+ (2020), Boudaud&Cirelli (2019), Dasgupta+ (2019), DeRocco& Graham (2019), Laha (2019), Clark+ (2017), Stöcker+ (2018), Acharya&Khatri (2020), Keith+ (2020), Clark+ (2018), Hektor+ (2018), Carr+ (2010), Coogan+ (2020), Bernal+ (2022), Auffinger (2022) ...



B. J. Kavanagh, "PBH bounds." https://github.com/bradkav/PBHbounds.

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### Neutrinos from PBH evaporation

A BH can emit elementary particles following a blackbody-like distribution.

We rely on the public tool *BlackHawk* to compute the neutrino spectra from an evaporating PBH of a given mass and spin. Arbey and Auffinger (2019)

Two contributions:

- Primary: neutrinos directly emitted in the evaporation of the PBH
- Secondary: stems from the hadronization and subsequent decay of the primary particles

Both PBHs in the galactic halo and outside contribute to the differential neutrino flux from PBH evaporation:

Galactic contribution

Extragalactic contribution

$$\frac{d\phi_{\text{gal}}^{\nu}}{dE} = \frac{f_{\text{PBH}}}{M_{\text{PBH}}} \frac{d^2 N}{dEdt} \int \frac{1}{4\pi} d\Omega \int \rho_{\text{MW}} \left[ r(\ell, \psi) \right] d\ell$$

$$\frac{d\phi_{\rm exg}^{\nu}}{dE} = \frac{f_{\rm PBH}\,\bar{\rho}_{\rm DM}}{M_{\rm PBH}} \int_{t_{\rm min}}^{t_{\rm max}} dt [1+z(t)] \, \frac{d^2N}{dE_0 dt} \left| \begin{array}{c} {}^{\rm (E_0\,=\,neutrino\,\,energy}{}_{\rm at\,\,the\,\,source)} \\ E_0 = [1+z(t)]E \end{array} \right|^{(E_0\,=\,neutrino\,\,energy}{}_{\rm at\,\,the\,\,source)} \right|^{(E_0\,=\,neutrino\,\,energy}{}_{\rm at\,\,the\,\,source)}$$

Assumption: the DM is distributed isotropically at sufficiently large scales. The exg differential neutrino flux over the full sky is estimated as the redshifted sum over all epochs emission.

### Neutrinos from PBH evaporation

Total differential  $v_e$  flux from PBHs. We assume a monochromatic PBH mass distribution and different PBH spins.

- Heavier PBHs lead to reduced fluxes, shifted towards smaller neutrino energies.
- The dominant (smooth) contribution to the flux stems from particles emitted with no additional angular momentum apart from their intrinsic spin.
  Peaks at higher energies: radiation of particles with higher angular momentum, leading to a substantial enhancement of the flux if the latter is aligned with the PBH one.



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### Neutrinos from PBH evaporation

The PBH evaporation rate depends on the total number of degrees of freedom of the emitted particles, as well as on their masses.

- Neutrinos are considered to be massless (non-zero masses would make a very small difference) and with a Majorana nature.
- We do not include the effect of neutrino mixing in our analysis.
- Flavour conversions would only affect the secondary components.
- Considering neutrino mixing would reduce the total differential flux by less than ~ 2%



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# Monochromatic vs extended mass distribution

A large variety of formation mechanisms predicts extended mass distributions for PBHs. A common approximation to the true mass distribution generated from a peak in the power spectrum of primordial fluctuations is the log-normal distribution for the comoving number density:



the flux is still sizeable at larger energies



E [MeV]

### Event spectra at future facilities

proce



Predicted  $v_e$  event rate from PBH evaporation at the DUNE far detector, mainly sensitive to the  $v_e$ component of the PBH neutrino flux via the

ss 
$$\nu_e + {}^{40}\mathrm{Ar} \to e^- + {}^{40}K$$

Bckgs: solar neutrinos (<16 MeV) and CC atm neutrinos (>30 MeV).
SNOwGLoBES

DUNE collab. Abi+ (2020) <u>http://www.phy.duke.edu/~schol/snowglobes</u> DSNB: de Gouvêa+ Phys. Rev. D 102 (2020) 123012

THEIA is a proposal for a water-based liquid scintillator (obtained by adding ultra-pure water to an organic scintillator). The dominant reaction channel will be IBD:  $\bar{\nu}_e + p \rightarrow e^+ + n$ 

Bckgs: solar neutrinos, CC atm neutrinos, reactor neutrinos and NC high-energy atmospheric neutrinos.

Theia collab. M. Askins+, EPJC 80 5 (2020)



### Statistical analysis

We assume the following least-squared functions:

$$\chi^{2}_{\text{DUNE}} = \sum_{i} \left( \frac{N_{i}^{\text{PBH}} + (1+\alpha)N_{i}^{\text{atmCC}} + (1+\beta)N_{i}^{\text{DSNB}} - N_{i}^{\text{atmCC}} - N_{i}^{\text{DSNB}}}{\sqrt{N_{i}^{\text{PBH}} + N_{i}^{\text{atmCC}} + N_{i}^{\text{DSNB}}}} \right)^{2} + \left(\frac{\alpha}{\sigma_{\alpha}}\right)^{2} + \left(\frac{\beta}{\sigma_{\beta}}\right)^{2}$$

The nuisance parameters  $\alpha$  and  $\beta$  account for uncertainties on the backgrounds (DSNB and atm CC). We fix their standard deviations to  $\sigma_{\alpha} = 0.35$  and  $\sigma_{\beta} = 0.35$ .

Battistoni+, Astropart. Phys. 23 (2005) 526–534 Sawatzki+, Phys. Rev. D 103 no. 2, (2021) Guo, Phys. Rev. D 99, 073007 De Gouvêa+, Phys. Rev. D 102 (2020) 123012

$$\chi^{2}_{\text{THEIA}} = \sum_{i} \left( \frac{N_{i}^{\text{PBH}} + (1+\alpha)N_{i}^{\text{atmCC}} + (1+\beta)N_{i}^{\text{DSNB}} + (1+\gamma)N_{i}^{\text{atmNC}} - N_{i}^{\text{atmCC}} - N_{i}^{\text{DSNB}} - N_{i}^{\text{atmNC}}}{\sqrt{N_{i}^{\text{PBH}} + N_{i}^{\text{atmCC}} + N_{i}^{\text{DSNB}} + N_{i}^{\text{atmNC}}}} + \left(\frac{\alpha}{\sigma_{\alpha}}\right)^{2} + \left(\frac{\beta}{\sigma_{\beta}}\right)^{2} + \left(\frac{\gamma}{\sigma_{\gamma}}\right)^{2} . \quad (12)$$

 $\sigma_{\gamma}$  refers to the atm NC background, and is fixed to 30%.

KamLAND Collaboration, Astrophys. J. 745 (2012) 193 Cheng+, Phys. Rev. D 103 (Mar, 2021) 053002

### PBH sensitivities at DUNE



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Super-Kamiokande: current bound from neutrinos  $\begin{array}{c} Dasgupta+, Phys. Rev. Lett. 125, 101101\\ See also: Bernal+ 2203.14979\end{array}$ Evaporation bounds: mainly extragalactic  $\gamma$ -ray background, CMB and MeV extragalactic  $\gamma$ -ray background and soft  $\gamma$  rays from COMPTEL.

JUNO: Wang+, Phys. Rev. D 103 no. 4, (2021) 043010

### PBH sensitivities at Theia



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### Conclusions

Primordial black holes are a viable dark matter candidate, whose masses can span over several orders of magnitude.

Light PBHs with masses in the range 10<sup>15</sup> – 10<sup>17</sup> g would evaporate and produce sizeable fluxes of MeV neutrinos via Hawking radiation.

We have explored the possibility of detecting neutrinos from PBH evaporation with the future neutrino experiments DUNE and THEIA.

Our analysis shows that both DUNE and THEIA could considerably improve the existing bounds from Super-Kamiokande on the abundance of PBHs with masses between 10<sup>15</sup> – 10<sup>16</sup> g, and even allow to probe heavier PBHs up to ~ 5 × 10<sup>16</sup> g.

Next-generation neutrino facilities could set competitive bounds on the PBH parameter space, complementary to existing limits from other cosmic messengers.

The quest for PBH dark matter at future neutrino experiments like DUNE and THEIA could be easily included in their physics programmes, leveraging existing supernova neutrino analyses.



### PBH sensitivities at DUNE



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