#### Gravitational-wave probes of planetaryand asteroid-mass primordial black holes

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Miller et al. Phys.Dark Univ. 32 (2021) 100836; Miller et al. Phys.Rev.D 105 (2022) 6, 062008; LVK: Phys.Rev.D 106 (2022) 10, 102008

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## Motivation

- > Many GW efforts to detect PBHs focus on "sub-solar mass" regime, O(0.1 M₀)
- ➤ However, GWs from PBHs with masses [10-7 - 10-3]M<sub>☉</sub> have not been searched for
- > Matched filtering in this mass range is extremely computationally challenging
- Signal for binaries in this mass range resemble continuous waves, since these systems will inspiral for a very long time and spin up very slowly

100.0% 01-02 01-03a  $\longrightarrow$  01-O2 HMR:  $m_1 \simeq 37 M_{\odot}$ 10.0%  $\widetilde{f}_{\mathrm{PBH}}^{\mathrm{upper}}$ 1.0% 0.1% 1.00.1Primordial Black Hole Mass  $[M_{\odot}]$ 

> Nitz & Wang: Phys.Rev.Lett. 127 (2021) 15, 151101. LVK: Phys.Rev.Lett. 129 (2022) 6, 061104 LVK: arXiv: 2212.01477



## Continuous waves: what and why

- Quasi-monochromatic, persistent signals, subject 5 to Doppler modulations
- > The longer we can observe for, the more sensitive we are to these kinds of signals
- A completely new GW signal type: first detection will be another major milestone
- > We can keep observing continuous-wave (CW) sources!
- So many sources, and so many algorithms not based on matched filtering!



# LVK all-sky search in O3 data

- Four pipelines used
- These limits represent the minimum detectable amplitude h<sub>0</sub> as a function of gravitational-wave frequency, at 95% confidence
- They are generic, in the sense that they can be interpreted to be for *any* system that follows a linear frequency evolution over time



LVK: Phys.Rev.D 106 (2022) 10, 102008



# Interpretation of upper limits

- All-sky continuous-wave upper limits give  $h_0(f)$  averaged over all sky positions
- Searches specify a maximum spin-up, consider systems with linear frequency evolutions
- There could be binary systems with very small masses (PBHs) that are inspiraling and "chirping" linearly

Reinterpret upper limits for chosen chirp masses M

$$\dot{f}_{\rm gw} = \frac{96}{5} \pi^{8/3} \left(\frac{G\mathcal{M}}{c^3}\right)^{5/3} f_{\rm gw}^{11/3}$$

$$f_{\rm gw,L}(t) = f_0 + \dot{f}_{\rm gw}$$

 $d = \frac{4}{h_0} \left(\frac{G\mathcal{M}}{c^2}\right)^{5/3} \left(\frac{\pi f_{gw}}{c}\right)$ 

 $(t - t_0)$ 

#### Under the conditions:

$$\dot{f}_{\rm gw} \leq 2 \times 10^{-9} \, {\rm Hz/s}$$

$$|f_{\rm gw,L}(T_{\rm obs}) - f_{\rm gw}(T_{\rm obs})| \le \Delta f = \frac{1}{T_{\rm FFT}}$$

searched

Linear evolution

#### Miller et al. Phys.Rev.D 105 (2022) 6, 062008



# Enforcing linearity and spin-up

 Chirp mass determines spin-up
 Search ranges for spin-up and frequency restrict which upper

limits we can include here

> Distance & no detection -> rates

$$N_{\rm bin}(f) \simeq rac{4}{3} \pi d(f)^3 RT,$$



T: how long binary spends in the band R: rate N<sub>bin</sub> : number of detected binaries (< 1)

 $N_{\rm bin}^{\rm tot} = \sum N_{\rm bin}(f_i)$ 



#### O3 Constraints, Asteroid-mass PBHs

- > CW limits  $h_0^{95\%}(f) \to d_{\text{bin}}^{95\%}$
- Only allow linear frequency evolution:  $f(t) = f_0 + \dot{f}(t - t_0)$
- > Maximum  $\dot{f}$  searched over ( ~ 1 $\mu$ Hz/s) -> fixed ranges for  $\mathcal{M}, f$
- > Distance & no detection -> rates
- Constarints shown for asymmetric mass ratio binaries,  $m_1 = 2.5 M_{\odot}$



LVK: Phys.Rev.D 106 (2022) 10, 102008; arXiv:2201.00697













# Constraints can be improved

- > All-sky searches provide decent upper limits, but they are not the optimal way to search for inspiraling PBHs
- > In fact, inspiraling compact objects with chirp masses  $< O(10^{-3}) M_{\odot}$ would exhibit power-law frequency evolutions over time
- > Methods to detect "transient" continuous waves, such as those used to search for a remnant of GW170817, can be tuned to search for inspiraling PBHs
- > Note that searches/methods do not assume PBHs exist

# GWs from inspiraling PBHs

- $\sim [10^{-7} 10^{-3}] M_{\odot}$  give rise to signals that are long lasting, compared to those detected from  $O(M_{\odot})$  black holes
- The evolution of these binaries can be described as quasi-Newtonian circular orbits
- > Techniques used in GW data analysis for quasi-monochromatic or power-law signals can also be applied to detect PBHs





# Generalized Frequency-Hough

- Attempt to detect power-law signals that slowly "chirp" in time
- Input: points in time/frequency plane to lines [7,11]

$$\dot{f}_{\rm gw} = \frac{96}{5} \pi^{8/3} \left(\frac{G\mathcal{M}}{c^3}\right)^{5/3} f_{\rm gw}^{11/3}$$



# Generalized Frequency-Hough

Generalization of method to search for isolated neutron stars

Output: two-dimensional histogram in the frequency/ chirp mass plane of the source



# Projected constraints, PBH binaries

- Current CW search results limited -> designed new method to search for light PBHs that tracks the "chirp" inspiral
- Projected constraints on equalmass PBHs, for monochromatic and thermal mass functions (solid and dashed lines, respectively)









#### Conclusions

- masses

CW limits can be mapped to constraints on rates as a function of chirp mass for inspiraling systems, without any model assumptions

> When assuming particular models for PBH formation rates, these limits can be translated to the fraction of dark matter that PBHs compose, for both equal-mass and asymmetric-mass ratio binaries

Future observing runs, especially Einstein Telescope, can probe physical values for the fraction of dark matter at planet and asteroid

## Multi-Messenger CW Workshop

11-13 July 2023, Amsterdam
Registration is open!
<u>https://indico.nikhef.nl/e/mmcw</u>







Back-up slides

# Primordial Black Holes (PBHs)

- Low spins of LIGO/Virgo black hole mergers, and merging rate inferences have revived the interest in PBHs
- > Black holes that formed in the early universe can take on a wide range of masses depending on when they formed
- Could be linked to Dark Matter



#### Current constraints on PBHs

- Asteroid-mass region not well constrained
- Limits based on uniform distribution of PBHs, while GWs could probe clusters of PBHs
- Constraints can be evaded in particular PBH formation models, e.g. if PBHs form in clusters
- > Wide range of masses



### PBH Formation

- Could have formed in early universe which contained inhomogenities that stopped too dense regions from expanding, causing collapses
- > Quantum fluctuations in various inflation theories
- A binary could form if two PBHs form independently but are kept from merging by the gravitational pull of other nearby PBH
  - Binaries would take age of the universe to merge
- > Also, formation could occur through capture in a PBH halo
- > Many more....





#### Distance reach

- > For chosen chirp masses, compute the distance reach at each frequency in the upper limit
- > Frequencies >  $\sim 250$  Hz give rise to signals with spin-ups greater than the maximum allowed by the search
- > Typically, at a fixed chirp mass, higher frequencies  $\rightarrow$  larger distance reach

$$h_0 = \frac{4}{d} \left(\frac{G\mathcal{M}}{c^2}\right)^{5/3} \left(\frac{\pi f_{\rm gw}}{c}\right)$$





# PBH merging rates constraints

- > T is the time that a binary system would spend in a frequency range f to  $f + \delta f$ , where  $\delta f = \dot{f}_{\text{max}} T_{\text{obs}}$ ;  $T = max(\Delta T, T_{obs})$
- > To obtain constraints, sum over binaries present at each frequency, and demand that the total number of binaries observed is less than one

$$N_{\rm bin}(f) \simeq rac{4}{3}\pi d(f)^3 RT_{\rm c}$$

$$N_{
m bin}^{
m tot} = \sum_i N_{
m bin}(f_i)$$

$$\Delta T = \frac{5}{256} \pi^{-8/3} \left(\frac{c^3}{G\mathcal{M}}\right)^{5/3} \left[f^{-8/3} - (f+\delta f)^{-8/3}\right]^{-8/3} \left[f^{-8/3} - (f+\delta f)^{-8/3}\right]^{-8/$$



#### Constraints for equal-mass binaries

- Can derive model-independent rates on the basis of upper limits
- Constraints are made on effective parameter *f* that equals *f*<sub>pbh</sub> for a monochromatic mass function
- Assume rate model from G. Hutsi+ (2021) & Raidal (2019)





 $R = 1.04 \times 10^{-6} \text{kpc}^{-3} \text{yr}^{-1} f(m)^2 \left(\frac{m}{M_{\odot}}\right)^{-32/37} \left(f_{\text{PBH}}\right)^{53/37}$ 

