PRIMORDIAL BLACK HOLES AND THE 21 CM LINE

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IFIC, CSIC - U. Valencia







NEHOP

New Horizons in Primordial Black Hole physics

 $Naples, Italy \\ June 19^{th} to June 21^{st} 2023$



Gravitational waves

Dark Matter

Primordial Black Holes

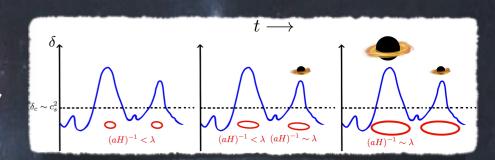
Cosmic radiation backgrounds

Formation: Physics of the Early Universe



Origin of SMBHs

The early universe is very hot and dense: ideal environment for black hole formation



Y. B. Zel'dovich and I. D. Novikov, Sov. Astron. 10:602, 1967

S. Hawking, Mon. Not. R. Astron. Soc. 152:75, 1971

Formed during radiation era from the gravitational collapse of a large fluctuation (at horizon entry) with a mass of the order of the horizon mass... or via collapse of cosmic string loops or a scalar field, bubble collisions...

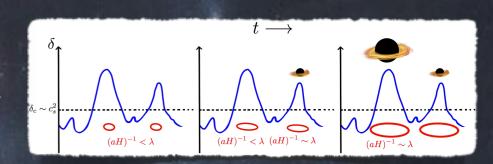
$$M_{\rm PBH} \sim \frac{t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \,\mathrm{s}}\right) \mathrm{g}$$
 $t = 10^{-43} \,\mathrm{s} \rightarrow M_{\rm PBH} \sim 10^{-5} \,\mathrm{g}$ $t = 1 \,\mathrm{s} \rightarrow M_{\rm PBH} \sim 10^{5} \,M_{\odot}$

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see talks by P. Cole, Y. Tada, E. Tomberg, C. Animali, A. Escriva, A. Perez, A. Kusenko, M. Zantedeschi, P. Conzinu, J. I. Juan ...

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Black holes radiate thermally, so they eventually evaporate

S. W. Hawking, Commun. Math. Phys. 43:199, 1975

D. N. Page, Phys. Rev. D13:198, 1976

$$T_{\rm BH} \sim \frac{1}{8 \,\pi \, G M_{\rm BH}} \sim 10 \left(\frac{10^{15} \,\rm g}{M_{\rm BH}}\right) \,\rm MeV$$

$$\tau(M_{\rm BH}) \sim G^2 M_{\rm BH}^3 \sim 100 \left(\frac{M_{\rm BH}}{10^{15} \,\rm g}\right)^3 \,\rm Gyr$$

For masses between 10 $^{-17}$ M $_{\odot}$ and 10 5 M $_{\odot}$, they would be present today

The early universe is very hot and dense: ideal environment for black hole formation

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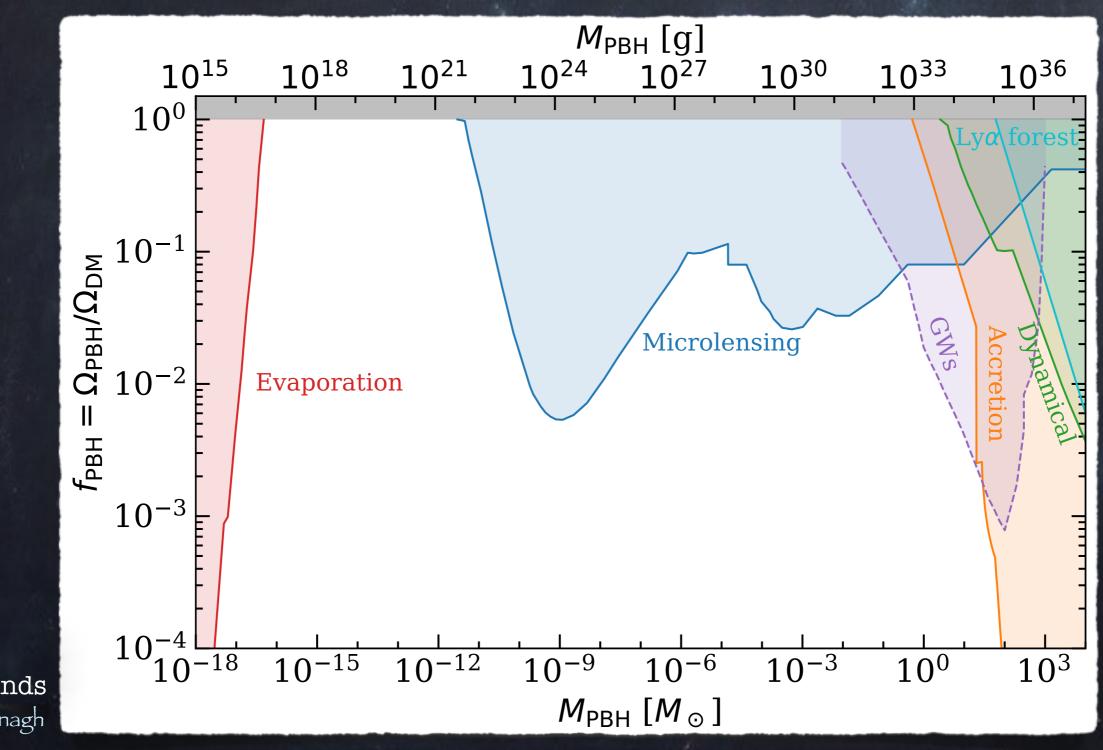
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PBHs would form before BBN, so they would not count as baryonic matter



A DM candidate which is not a new particle (although its formation involves BSM physics)

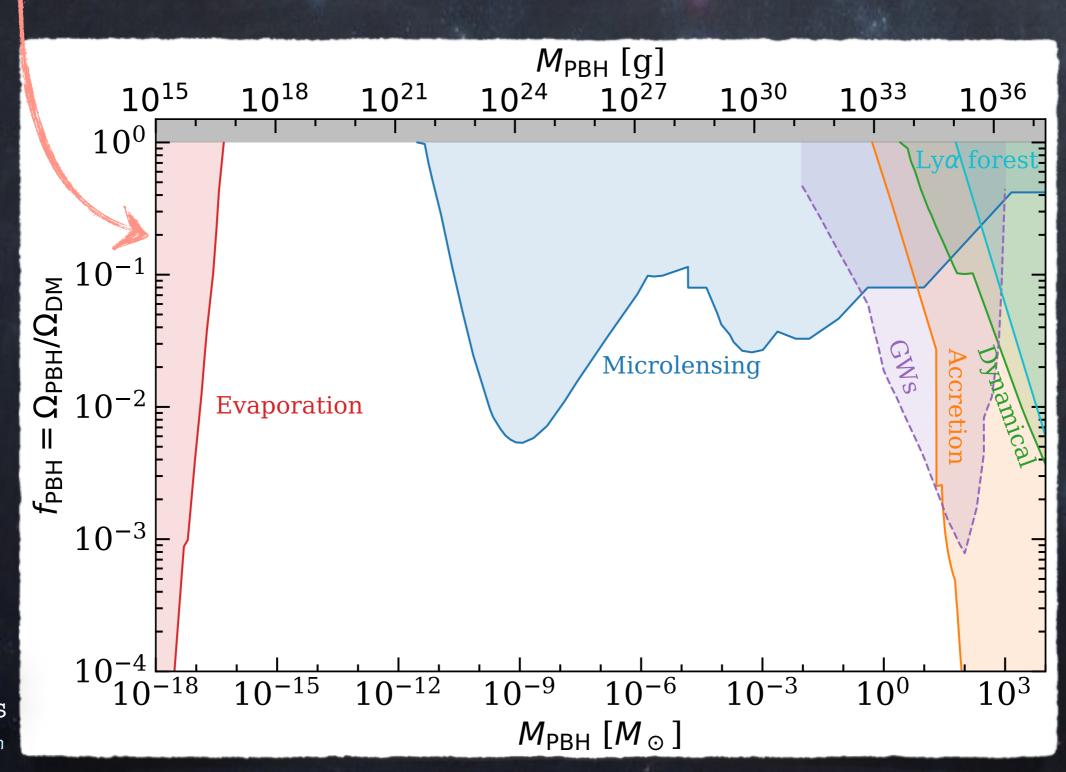


Using PBHbounds B. J. Kavanagh



Partial evaporation

Hawking radiation: cosmic-ray, γ-ray, v bkgs; ionization and thermal history



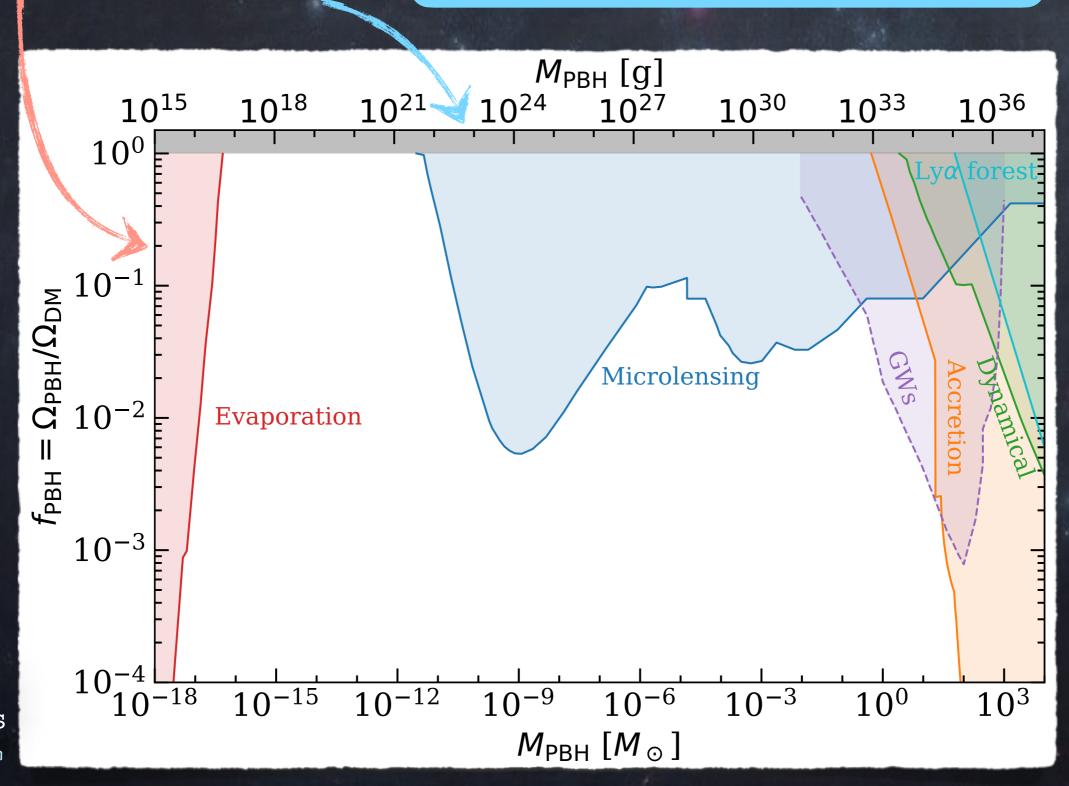
Using PBHbounds B. J. Kavanagh

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(femto, micro, milli) Lensing of GRBs, stars, SN, QSO, pulsars, FRBs

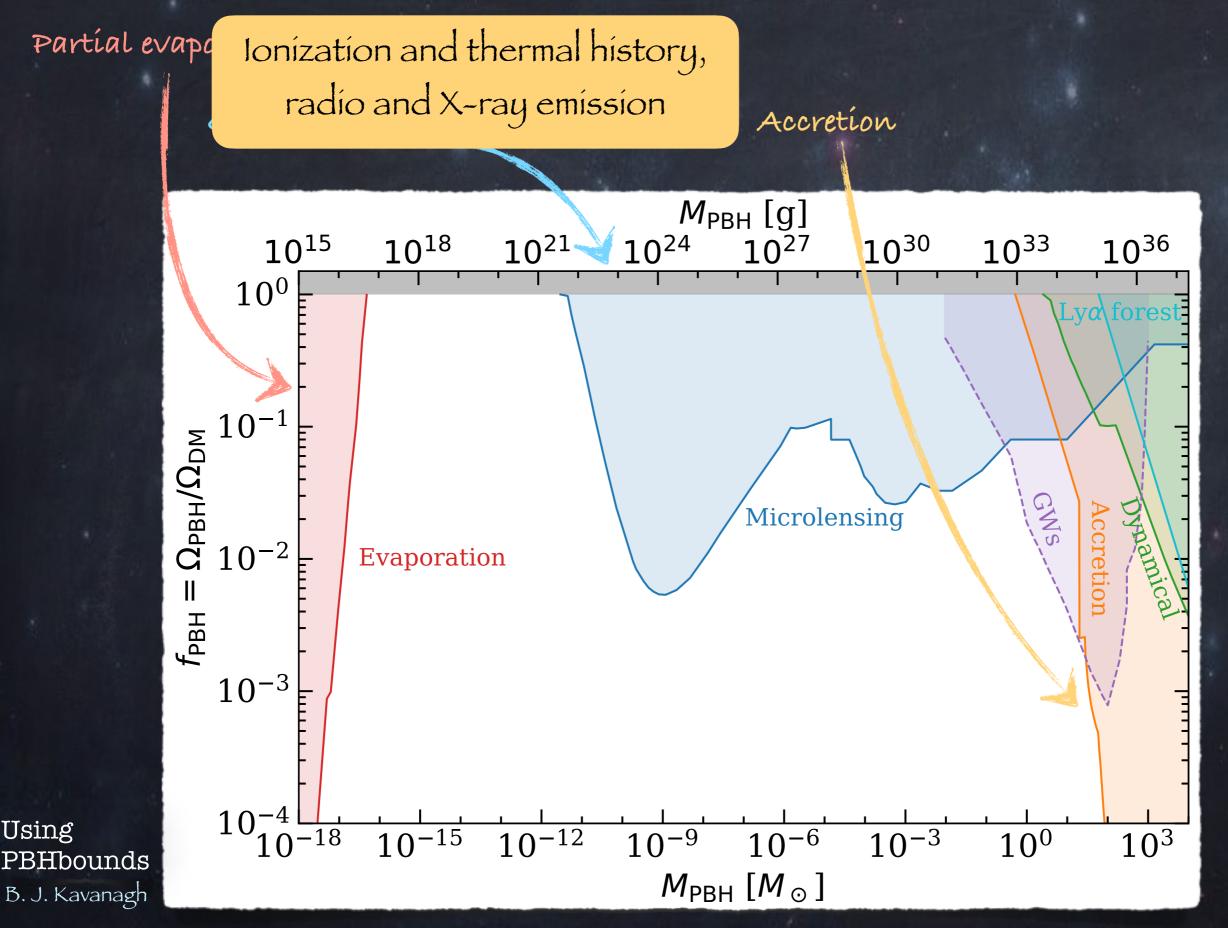
Gravitational lensing



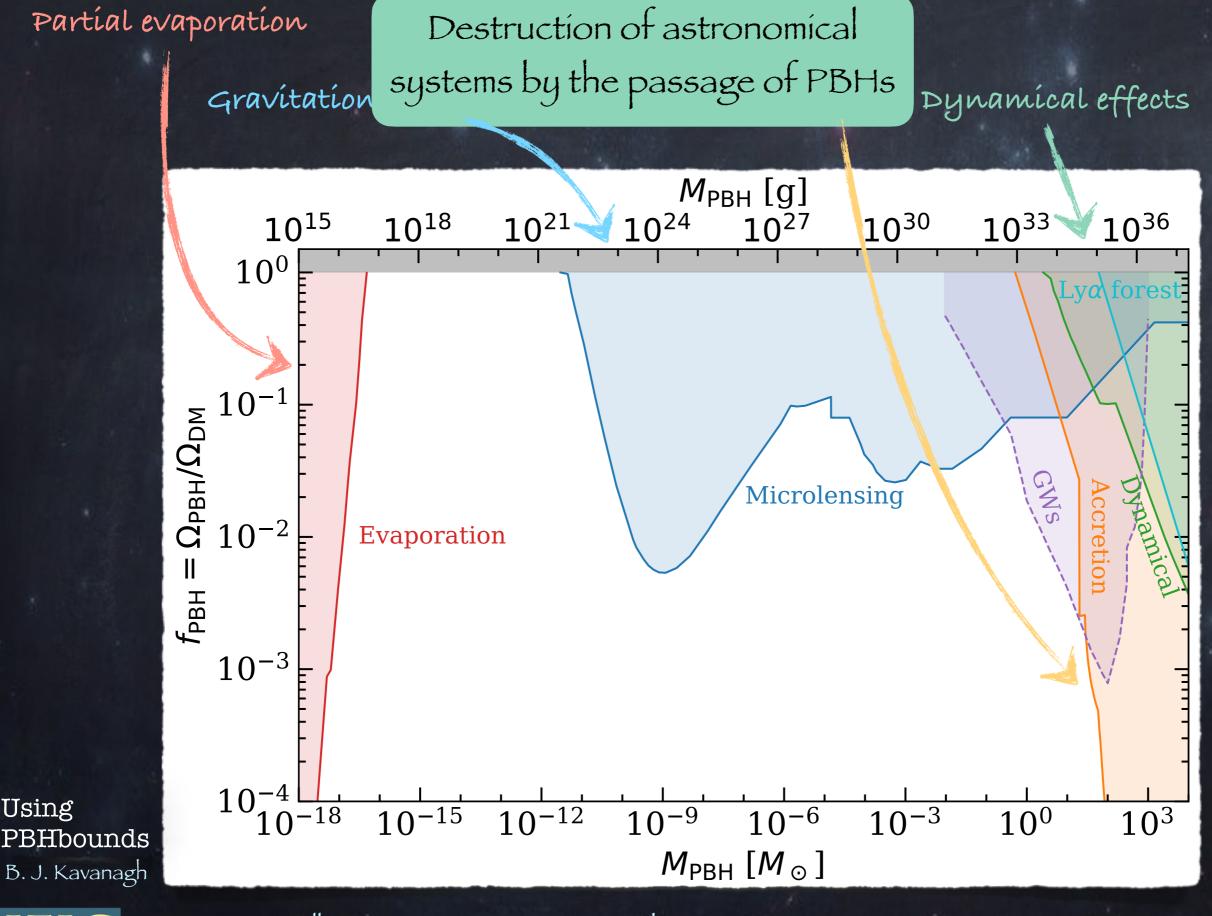
Using PBHbounds B. J. Kavanagh

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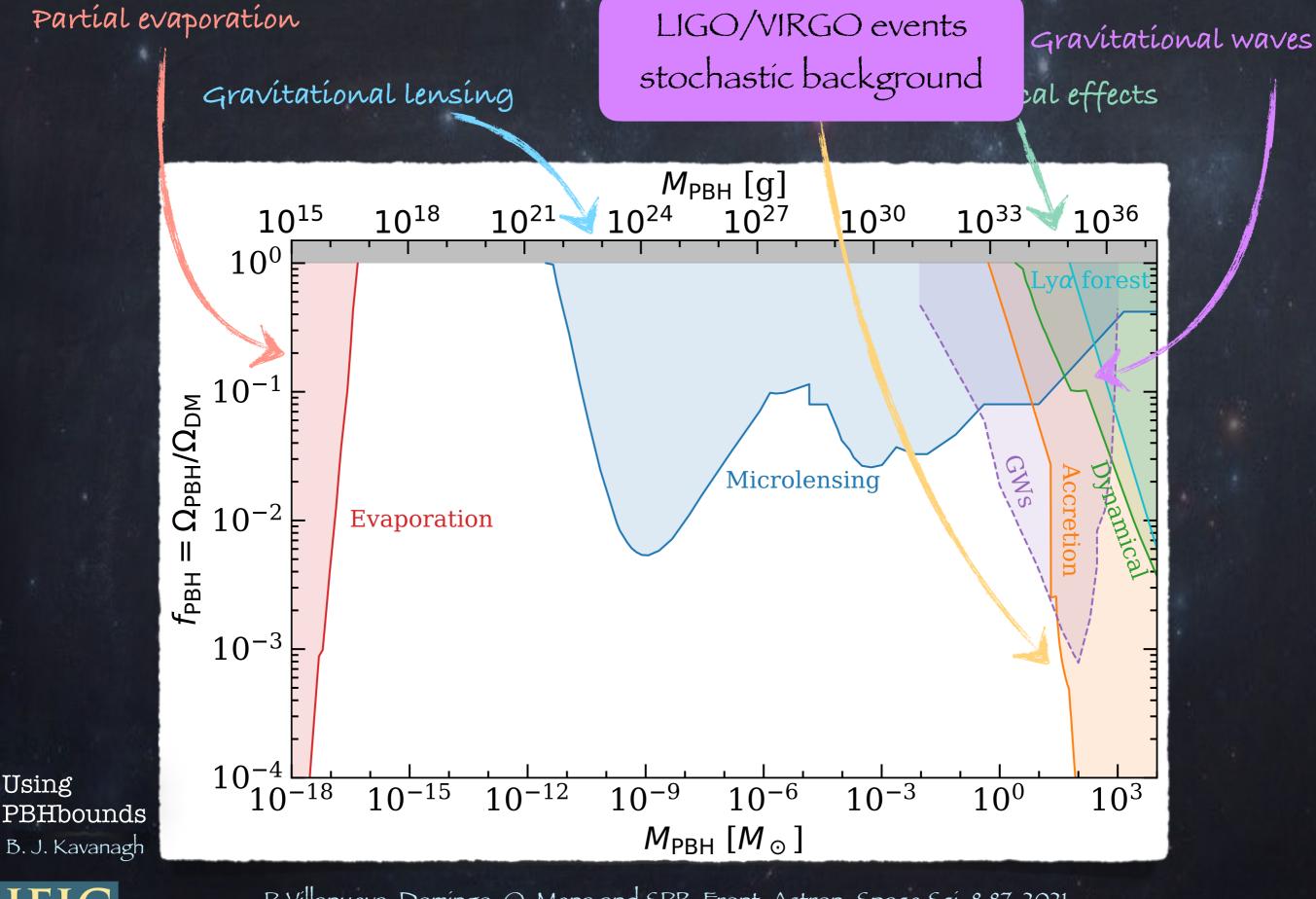
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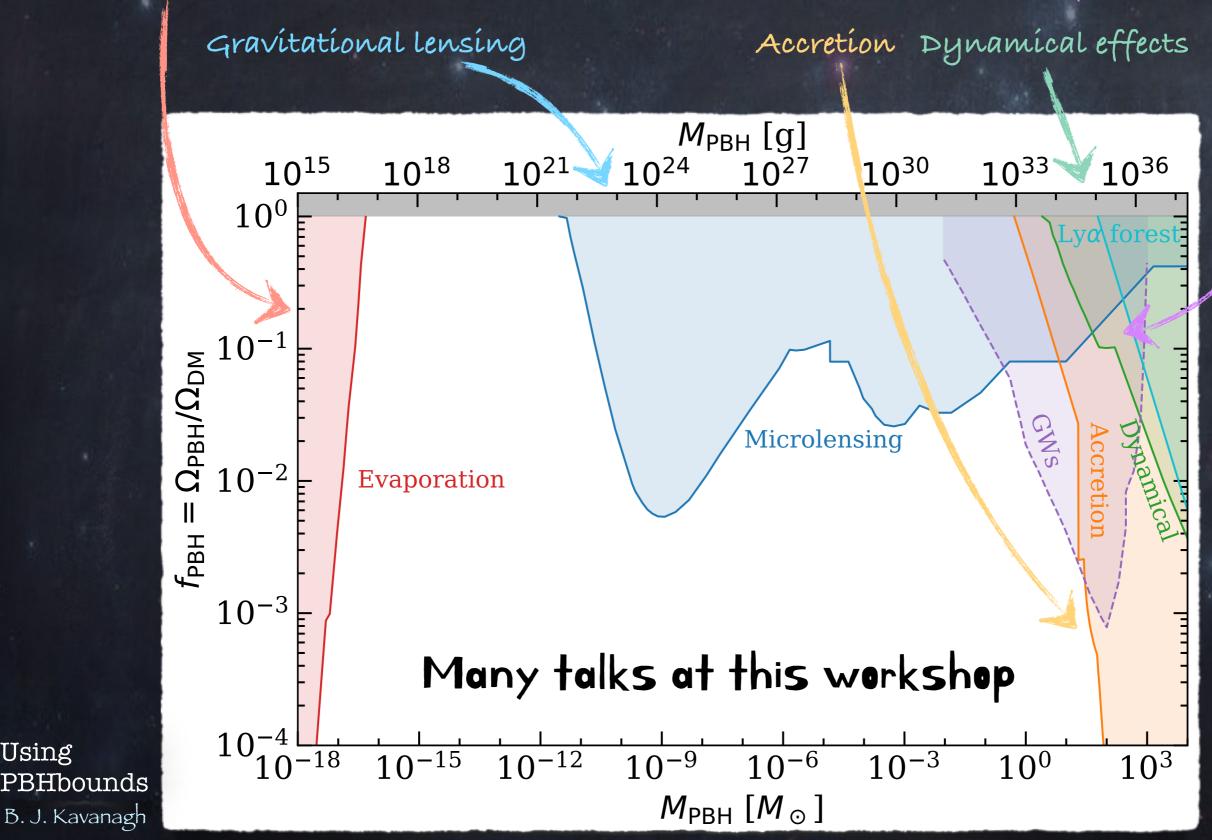




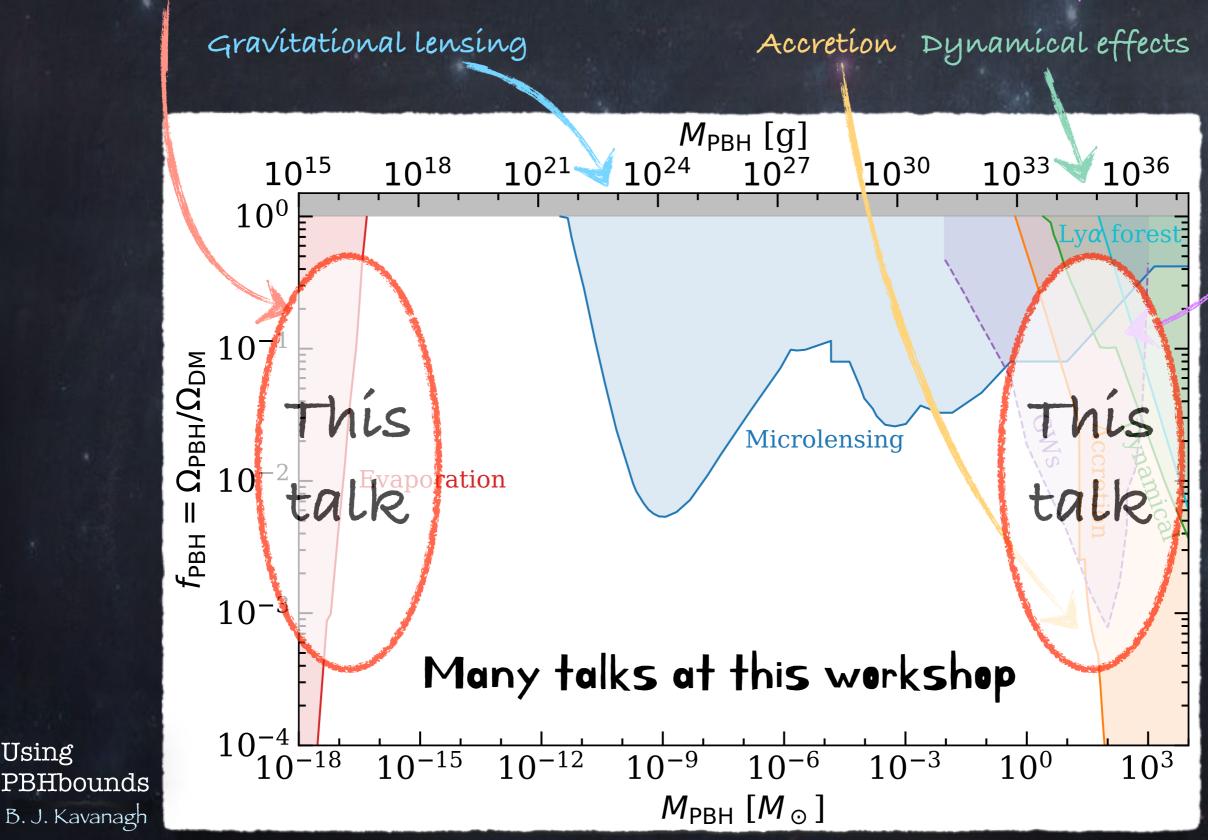












Even if they cannot form all the dark matter... still of great interest

see talks by B. Carr, L. Visinelli

Recent detection of black hole mergers with gravitational waves

B. P. Abbott et al. [LVC], Phys. Rev. Lett. 116:061102, 2016; Phys. Rev. Lett. 116:241103, 2016; Phys. Rev. Lett. 116:131102, 2016; Phys. Rev. X6:041015, 2016; Phys. Rev. Lett. 118:221101, 2017; Astrophys. J. 851:L35, 2017; Phys. Rev. Lett. 119:141101, 2017

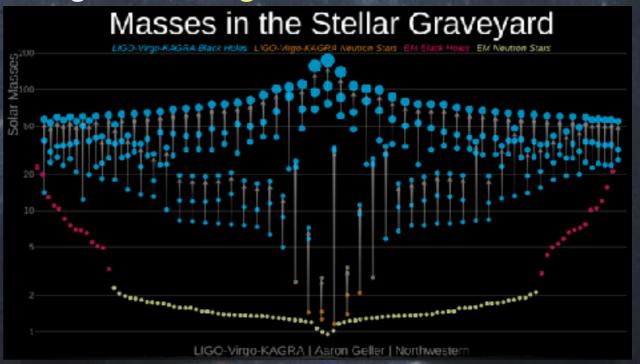
Did LIGO detect dark matter?

S. Bird et al., Phys. Rev. Lett. 116:201301, 2016

see talks by A. Riotto, A. Miller

Insight into early universe physics (inflation, phase transitions...)

many talks at this workshop



WIMPs and PBHs relation: no go

B. Lacki and J. F. Beacom, Astrophys. J. 720:L67, 2010

R. Saito and S. Shirai, Phys. Lett. B697:95, 2011

D. Zhang, Mon. Not. R. Astron. Soc. 418:1850, 2011

PBHs as DM (or other exotics) generators

see talks by I. Masina, N. Bernal, A. Cheek, Y. Perez, D. Montanino

Timing problem: Could PBHs be connected to the origin of SMBHs?

e.g., A. Smith and V. Bromm, Contemp. Phys. 60:111, 2019

see A. Ireland's talk

SOLAR MASS PBHS ABUNDANCE

VIRGO/LIGO:

Merger rates and masses related to BHs abundance

M. Sasakí, T. Suyama, T. Tanaka and S. Yokoyama, Phys. Rev. Lett. 117:061101, 2016 (E: Phys. Rev. Lett. 121:059901, 2018)

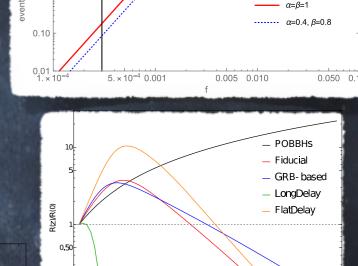
see A. Riotto's talk... but see also B. J. Kavanagh's talk

Gravitational waves at large redshifts (z>40) by O(10) solar mass PBH mergers: Einstein Telescope

S. M. Koushiappas and A. Loeb, Phys. Rev. Lett. 119:22 104, 2017

Z.-C. Chen and Q.-G. Huang, JCAP 08:039, 2020

See A. Riotto's talk



Accretion of gas onto PBHs: Emission of broad band spectrum observed excesses (X-rays, infrared, radio) see F. Kuhnel's talk heating and ionization of the IGM: cosmological implications

B. J. Carr, Mon. Not. R. Astron. Soc. 194:639, 1981

PBHs clustering modifies small scale structure: shot noise > isocurvature perturbations

N. Afshordí, P. McDonald and D. N. Spergel, Astrophys. J. 594, L71, 2003



Cosmic Infrared and X-ray unresolved backgrounds?

A. Kashlinsky, Astrophys. J. 823:L25, 2016

Primordial black holes and the 21cm line

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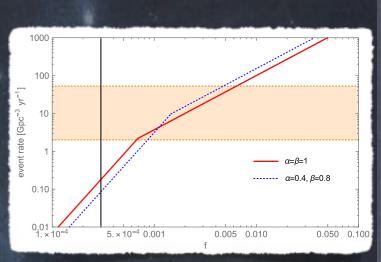
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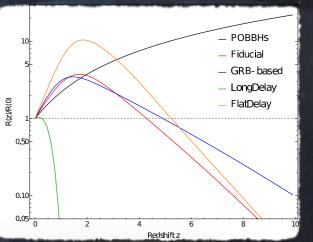
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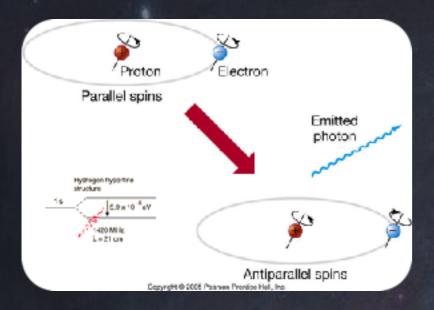
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Primordial black holes and the 21cm line

THE 21CM LINE

Predicted by H. van de Hulst in 1944 and first observed by H. I. Ewen and E. M. Purcell in 1951

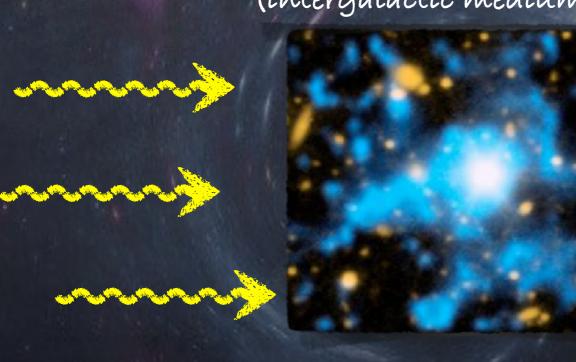


Hyperfine transition: v = 1420 Mhz

21cm photon from HI clouds during cosmic dawn: v ~ 100 Mhz

$$v = \frac{v_0}{(1+z)}$$

neutral hydrogen gas (intergalactic medium: IGM)



observer



emission/absorption

CMB photons as backlight

z~1000

Population of ground and excited states controlled by: absorption and stimulated emission of background radiation zeo collisions of neutral hydrogen excitation/de-excitation by Lyman-a photons



THE 21CM LINE

Probing Dark Ages and Cosmic Dawn Cosmic tomography









Interferometers LOFAR, MWA, PAPER, GMRT, LEDA, HERA, SKA

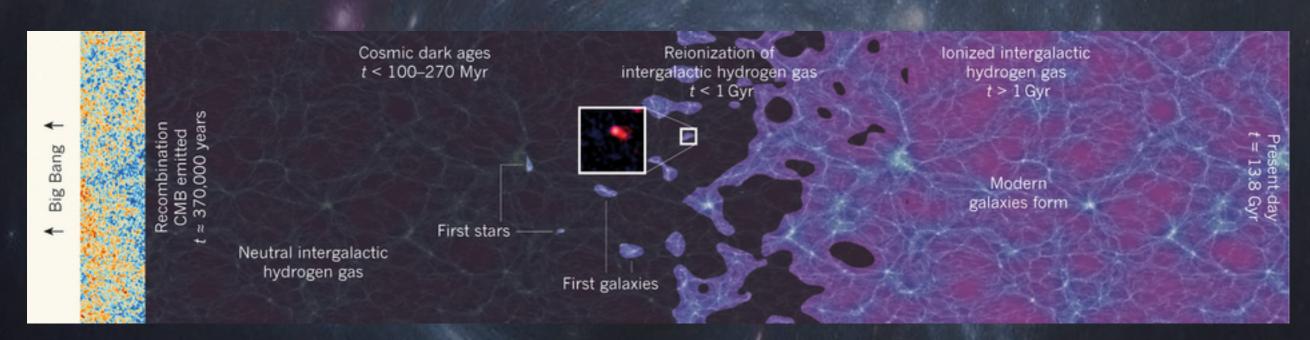




Galaxy Surveys







z~1000

Dark ages

z~30

Cosmic dawn

Galaxy formation



First stars

Reionization



THE 21CM SIGNAL

Differential brightness temperature

$$\delta T_b(v) \simeq 27 x_{HI} (1+\delta) \left(1 - \frac{T_{CMB}}{T_S}\right) \left(\frac{1+z}{10}\right)^{1/2}$$

Fraction of neutral H Reionization suppresses the signal

Baryon overdensity

$$\frac{n_1}{n_0} = 3e^{-T_{21}/T_5}$$

Spin temperature: occupation of the two states

signal in emission, can saturate

signal in absorption, limited by gas temperature

$$\delta T_b \approx 0$$
 if $T_S \sim T_{CMB}$
 $\delta T_b > 0$ if $T_S > T_{CMB}$
 $\delta T_b < 0$ if $T_S < T_{CMB}$

Astrophysical processes decouple T_{S} from T_{CMB}

Reionization: Dense medium: no neutral Spin temperature coupled hydrogen to gas via collisions and gas coupled to CMB via X-ray heating: Compton scattering $T_{
m CMB}$ from absorption CMB decouples: to emission gas cools faster density decreases: collisions not effective first stars: Lyman a coupling

10 20 30

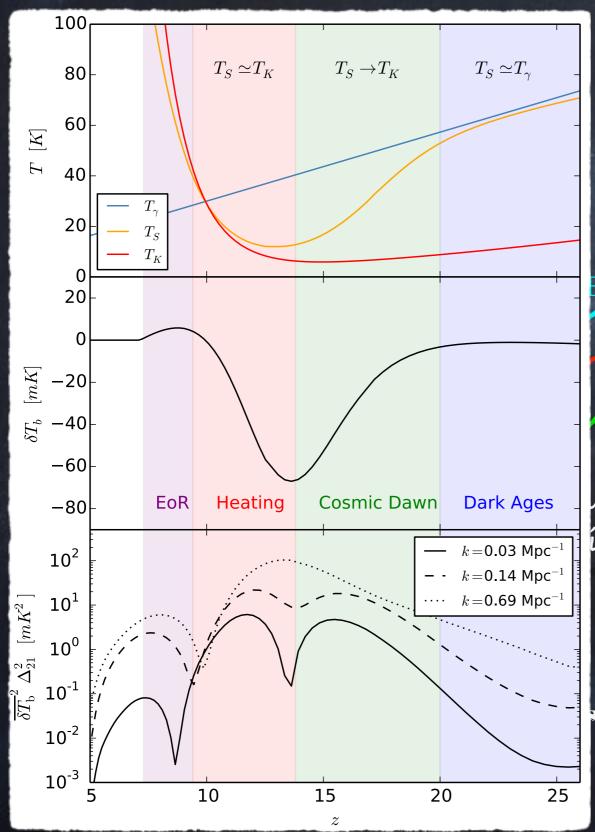
100

300



emberalare

Redshift



Dense medium:
Spin temperature coupled to gas via collisions and gas coupled to CMB via Compton scattering

CMB decouples: gas cools faster

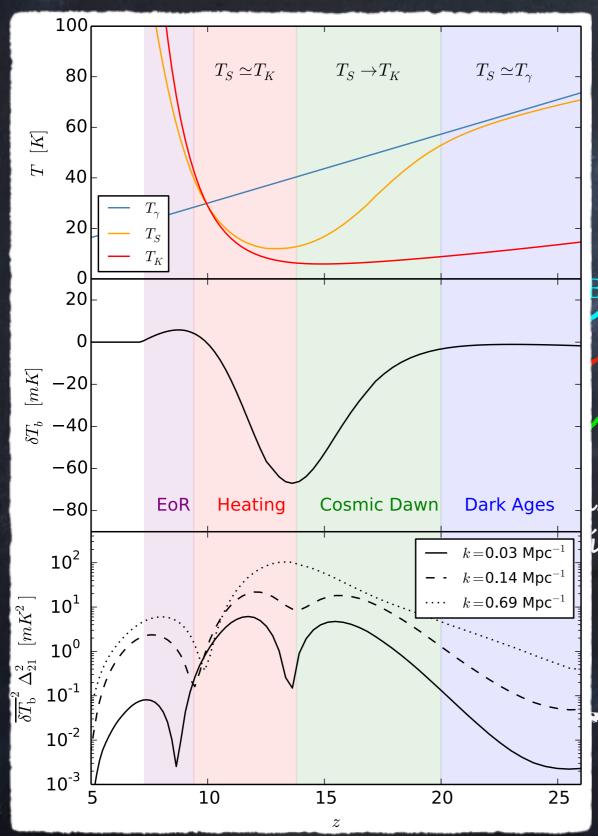
sity decreases: ions not effective

100

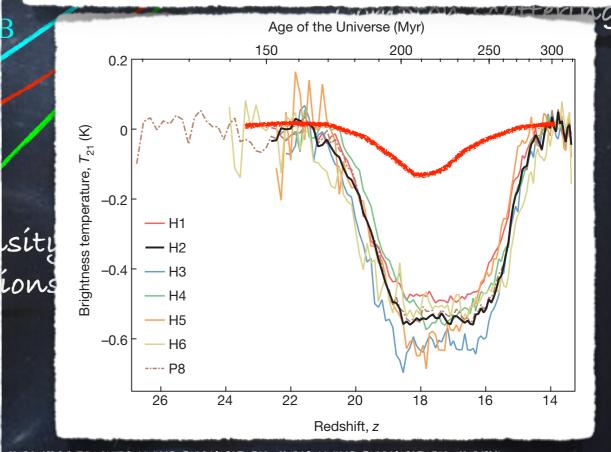
300



P. Villanueva-Domingo, PhD Thesis, 2021



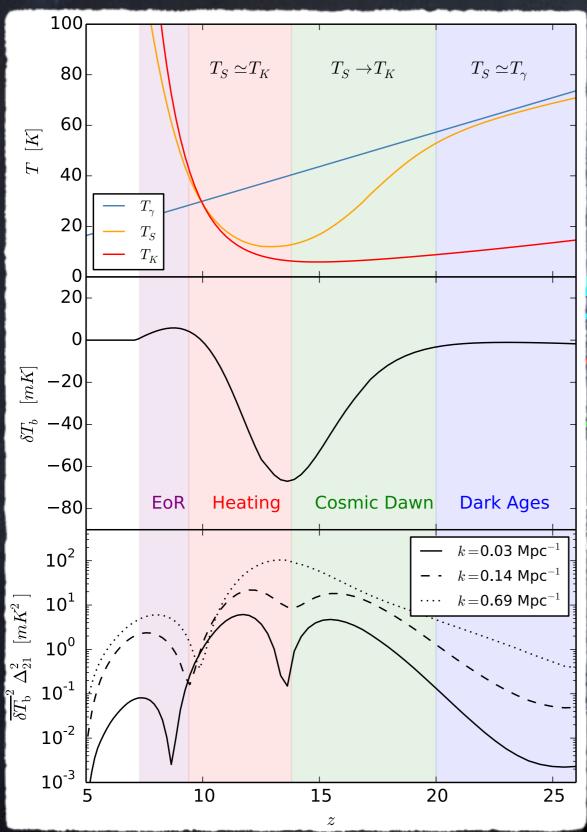
EDGES: new physics? astrophysics? systematics?



J. D. Bowman et al., Nature *555*:67, 2018



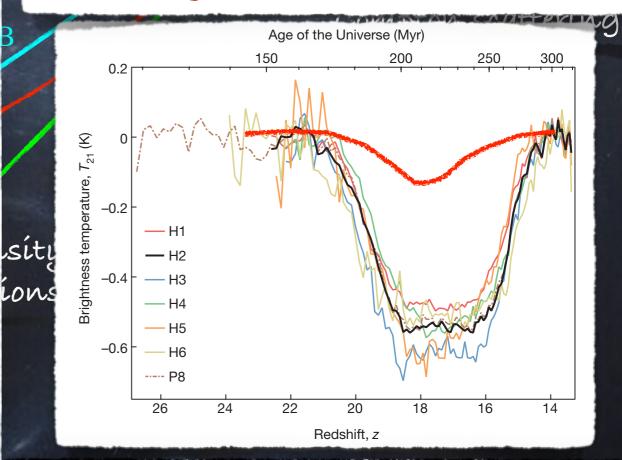
P. Villanueva-Domingo, PhD Thesis, 2021



EDGES:

new physics? astrophysics?

systematics?



Fully incompatible (amplitude, location AND shape) with standard scenarios

S. Witte, P. Villanueva-Domingo, S. Garíazzo, O. Mena and SPR, Phys. Rev. D97:103533, 2018



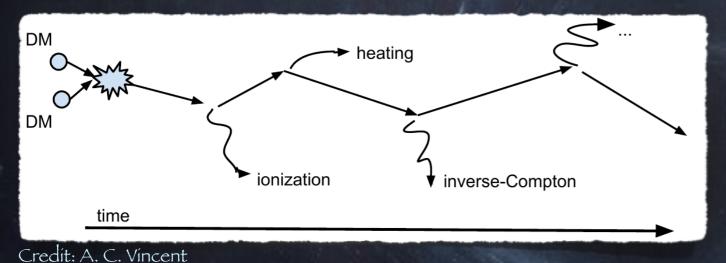
P. Villanueva-Domingo, PhD Thesis, 2021

10

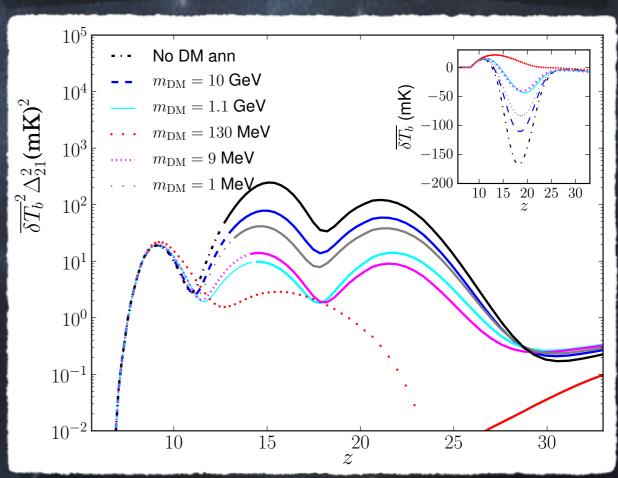
ENERGY INJECTION: EFFECT ON THE 21CM SIGNAL

Dark matter annihilations: inject energy into the IGM -> excite, ionize and heat

Chen'03, Hansen'03, Pierpaoli'03, Padmanabhan'05, Shchenikov'06, Furlanetto'06, Valdes'07, Chuzhoy'07, Cumberbatch'08, Natarajan'09, Yuan'09, Valdes'12, Evoli'14, Poulin'16, Liu'18...



DM annihilations: suppress power, but effects are degenerate with astrophysics



L. Lopez-Honorez, O. Mena, A. Moliné, SPR and A. C. Vincent, JCAP 1608:004, 2016

21CM SIGNAL FROM PBHS

ACCRETION

 $M_{\mathrm{PBH}} \gtrsim 0.1 M_{\odot}$

 $M_{\rm PBH} \lesssim 10^{-16} M_{\odot}$ EVAPORATION

K. J. Mack and D. H. Wesley, arXiv:0805.1531

uniform heating and ionization of the IGM

In the context of EDGES (accretion):

A. Hektor et al., Phys. Rev. D98:023503, 2018

Y. Yang, Phys. Rev. D104:063528, 2021

Forecasts (accretion):

O. Mena, SPR, P. Villanueva-Domingo and S. J. Witte, Phys. Rev. D100:043540, 2019

In the context of EDGES (evaporation):

S. Clark et al., Phys. Rev. D98:043006, 2018

Y. Yang, Phys. Rev. D102:083538, 2020

A. Halder and M. Pandey, MNRAS 508:3446, 2021

A. Halder and S. Banerjee, Phys. Rev. D103:0530044, 2021

S. Míttal et al., JCAP 03:030, 2022

U. Mukhopadhyay, D. Majumdar and A. Halder, JCAP 10:099, 2022

A. K. Saha and R. Laha, Phys. Rev. D105:103026, 2022

Forecasts (evaporation):

P. K. Natwariya, A. C. Nayak and T. Srivastava, MNRAS 510, 4236, 2021

J. Cang, Y. Gao and Y.-Z. Ma, JCAP 03:012, 2022

Y. Yang, Phys. Rev. D106:123508, 2022





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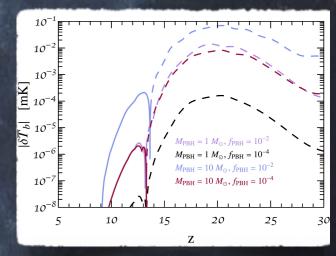
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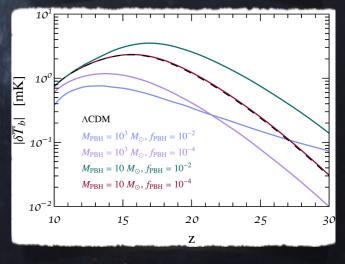
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J. Cang, Y. Gao and Y.-Z. Ma, JCAP 03:012, 2022 Y. Yang, Phys. Rev. D106:123508, 2022

Heating and ionization of the local environment of isolated PBHs

Enhanced signal from míní-halos by Poisson noise









21CM SIGNAL FROM PBHS

K. J. Mack and D. H. Wesley, arXiv:0805.1531

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A. Halder and M. Pandey, MNRAS 508:3446, 2021

A. Halder and S. Banerjee, Phys. Rev. D103:0530044, 2021 S. Míttal et al., JCAP 03:030, 2022

<u>-</u> 10^{−5}

U. Mukhopadhyay, D. Majumdar and A. Halder, JCAP 10:099, 2022 A. K. Saha and R. Laha, Phys. Rev. D105:103026, 2022

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Heating and ionization of the local Small contributions environment of isolate

mini-halos by Poisson noise

See, however, effects on the 21cm forest:

P. Villanueva-Domíngo and K. Ichíkí, Publ. Astron. Soc. Jpn. 75, S33, 2023

K. Kadota et al., JCAP 03:017, 2023

 $M_{\text{PBH}} = 10 \ M_{\odot}, \ f_{\text{PBH}} = 10^{-2}$ $M_{\rm PBH} = 10 \ M_{\odot}, f_{\rm PBH} = 10^{-6}$

O. Mena, SPR, P. Villanueva-Domingo and S. J. Witte, Phys. Rev. D100:043540, 2019

Sergio Palomares-Ruiz

ENERGY INJECTION DUE TO ACCRETION BY PBHS

Injected energy by PBHs:
$$\left(\frac{dE}{dVdt}\right)_{inj} = L_{acc} n_{PBH} = L_{acc} \frac{f_{PBH} \rho_{DM}}{M_{PBH}}$$

luminosity due to accretion:

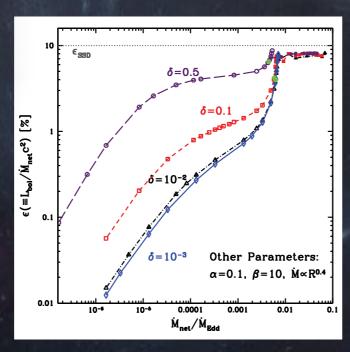
$$L_{acc} = \epsilon M_{PBH}$$

if cooling is inefficient, accretion could form a thick disk

V. Poulín, P. D. Serpico, F. Calore, S. Clesse and K. Kohrí, Phys. Rev. D96:083524, 2017

Radiative efficiency (ADAF)

$$\epsilon = \epsilon_0 \left(\frac{10 \, \dot{M}_{PBH}}{L_{Edd}} \right)^a$$



F.-G. Xie and F. Yuan, Mon. Not. R. Astron. Soc. 427:1580, 2012

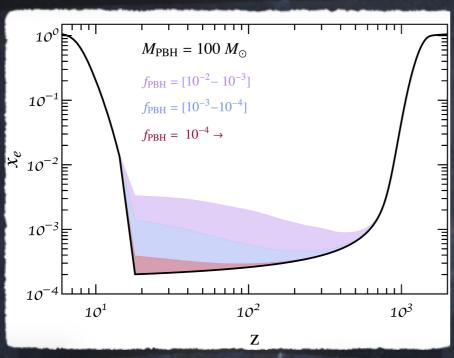
Bondí-Hoyle-Lyttleton accretion

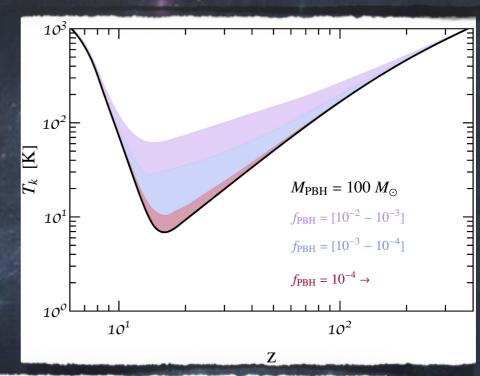
$$\dot{M}_{PBH} = 4\pi \lambda \rho_{\infty} \frac{(GM_{PBH})^2}{\left(c_s^2 + v_{rel}^2\right)^{3/2}}$$

F. Hoyle and R. A. Lyttleton, Mon. Not. R. Astron. Soc. 101:227, 1941 H. Bondi and F. Hoyle, Mon. Not. R. Astron. Soc. 104:273, 1944 H. Bondí, Mon. Not. R. Astron. Soc. 112:195, 1952

PBHS: BRIGHTNESS TEMPERATURE

Accretion: Injected energy goes into ionizing and heating the IGM





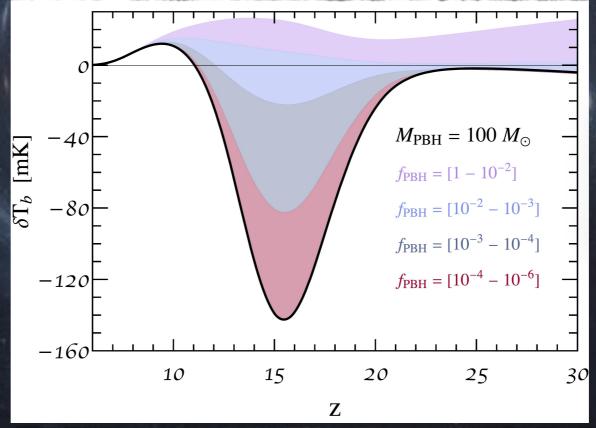
We use CosmoRec

J. Chluba and R. M. Thomas, Mon. Not. R. Astron. Soc. 412:478, 2011

Impact on the brightness temperature: suppression of signal

We use 21cmFAST

A. Mesinger, S. Furlanetto and R. Cen, Mon. Not. R. Astron. Soc. 411:955, 2011



From absorption to emission



PBHs: 21CM POWER SPECTRUM

We use 21cmSense

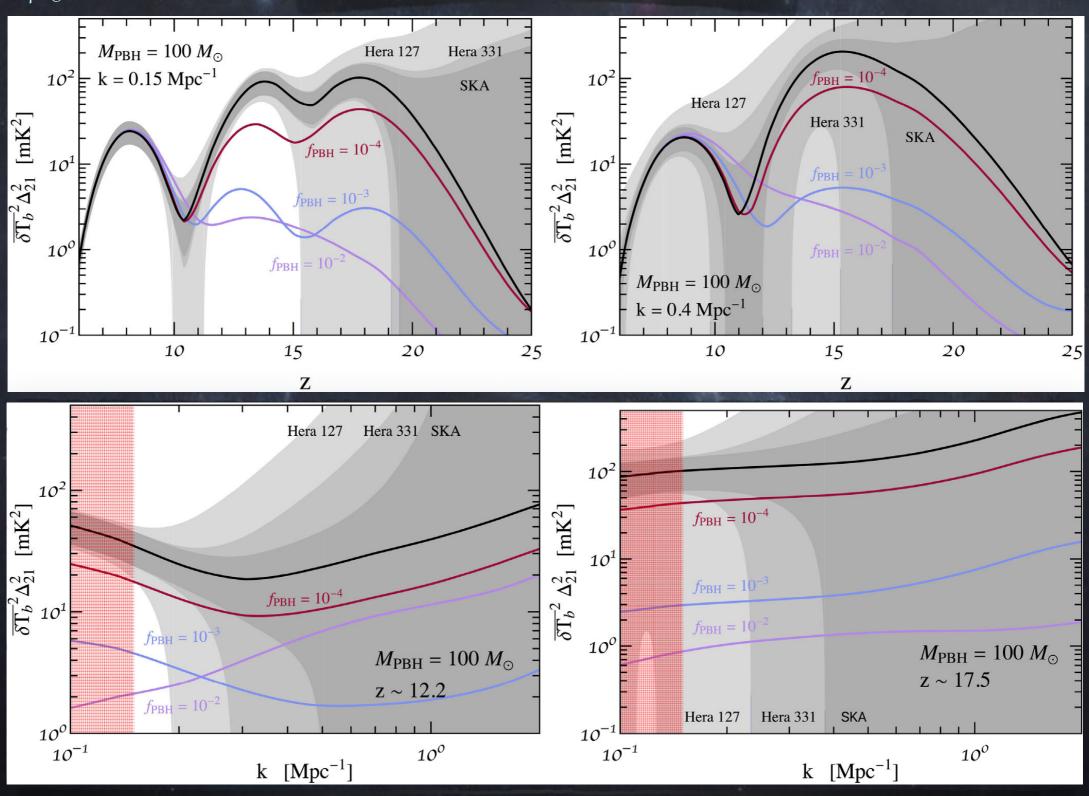
J. C. Pober et al., Astrophys. J. 145:65, 2013

J. C. Pober et al., Astrophys. J. 782:66, 2014

Four-parameter astrophysical model

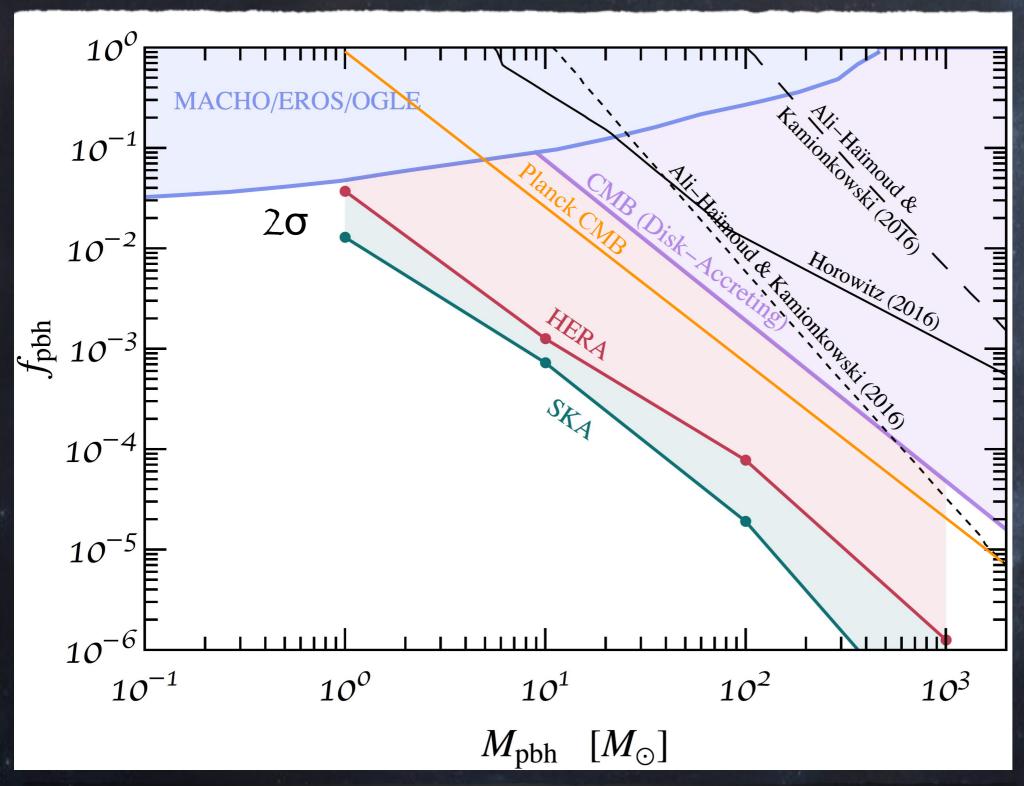








PBHs abundance: Sensitivity



Y. Alí-Haïmoud and M. Kamíonkowskí, Phys. Rev. D95:043534, 2017

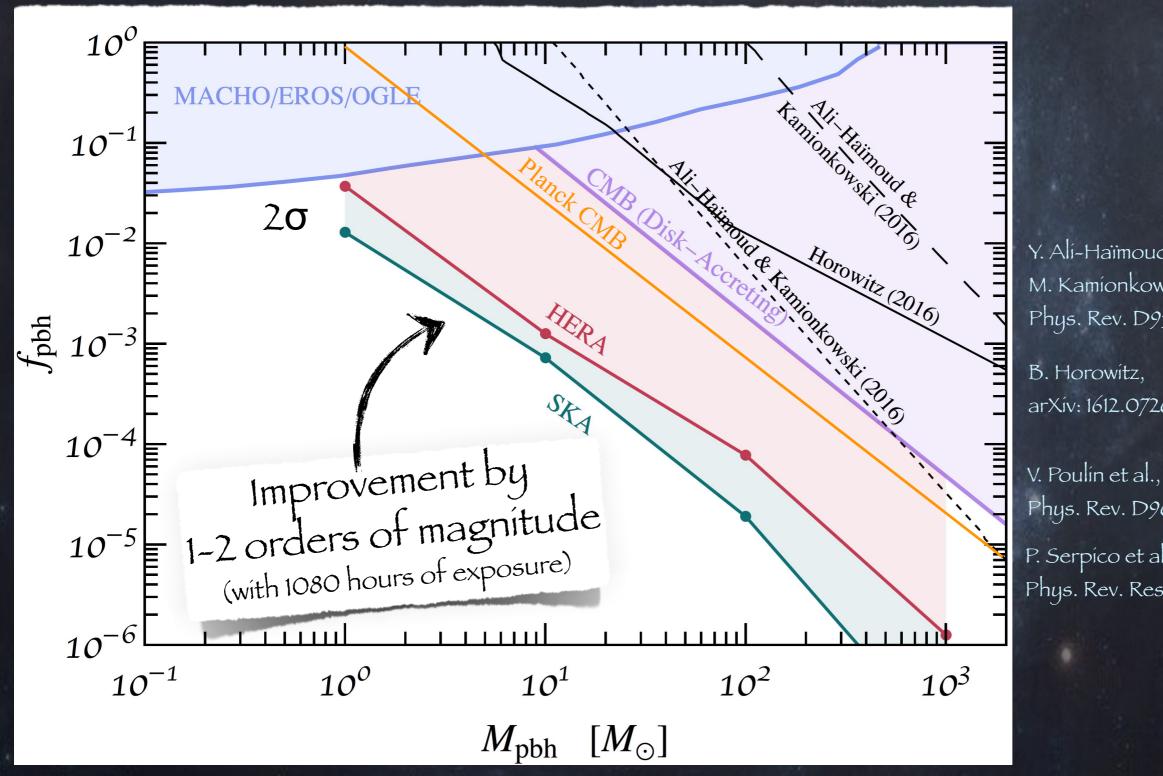
B. Horowitz, arXiv: 1612.07264

V. Poulín et al., Phys. Rev. D96:083524, 2017

P. Serpico et al., Phys. Rev. Res. 2:023204, 2020



PBHs abundance: Sensitivity



Y. Alí-Haimoud and M. Kamionkowski, Phys. Rev. D95:043534, 2017

arXiv: 1612.07264

Phys. Rev. D96:083524, 2017

P. Serpico et al., Phys. Rev. Res. 2:023204, 2020



PBH EVAPORATION EFFECTS ON THE 21CM SIGNAL

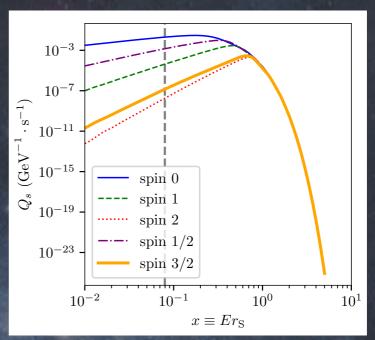
Emission of a quasi-thermal black-body spectrum

$$\frac{d^{2}N_{i}(E,t)}{dtdE}\bigg|_{prim} = \frac{g_{i}}{2\pi} \frac{\Gamma_{i}(E,M_{BH})}{e^{E/T_{BH}} \pm 1}$$

S. W. Hawking, Mon. Nature 248:30, 1974

S. W. Hawking, Commun. Math, Phys. 43:199, 1975

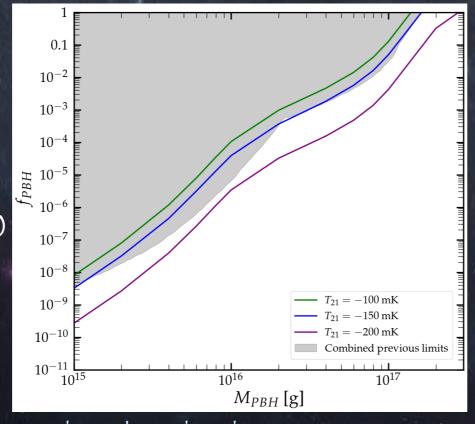
BlackHawk



A. Arbey and J. Auffinger, Eur. Phys. J. C79:693, 2019 A. Arbey and J. Auffinger, Eur. Phys. J. C81:10, 2021

PBH abundance: Sensitivity (only using the global signal)

> for other constraints, see talks by F. Calore, V. De Romeri



$$T_{\rm BH} \sim \frac{1}{8 \,\pi \, G M_{\rm BH}} \sim 10 \left(\frac{10^{15} \,\rm g}{M_{\rm BH}}\right) \,\rm MeV$$

$$\tau(M_{\rm BH}) \sim G^2 M_{\rm BH}^3 \sim 100 \left(\frac{M_{\rm BH}}{10^{15} \, \rm g}\right)^3 \, {\rm Gyr}$$



A. K. Saha and R. Laha, Phys. Rev. D105:103026, 2022

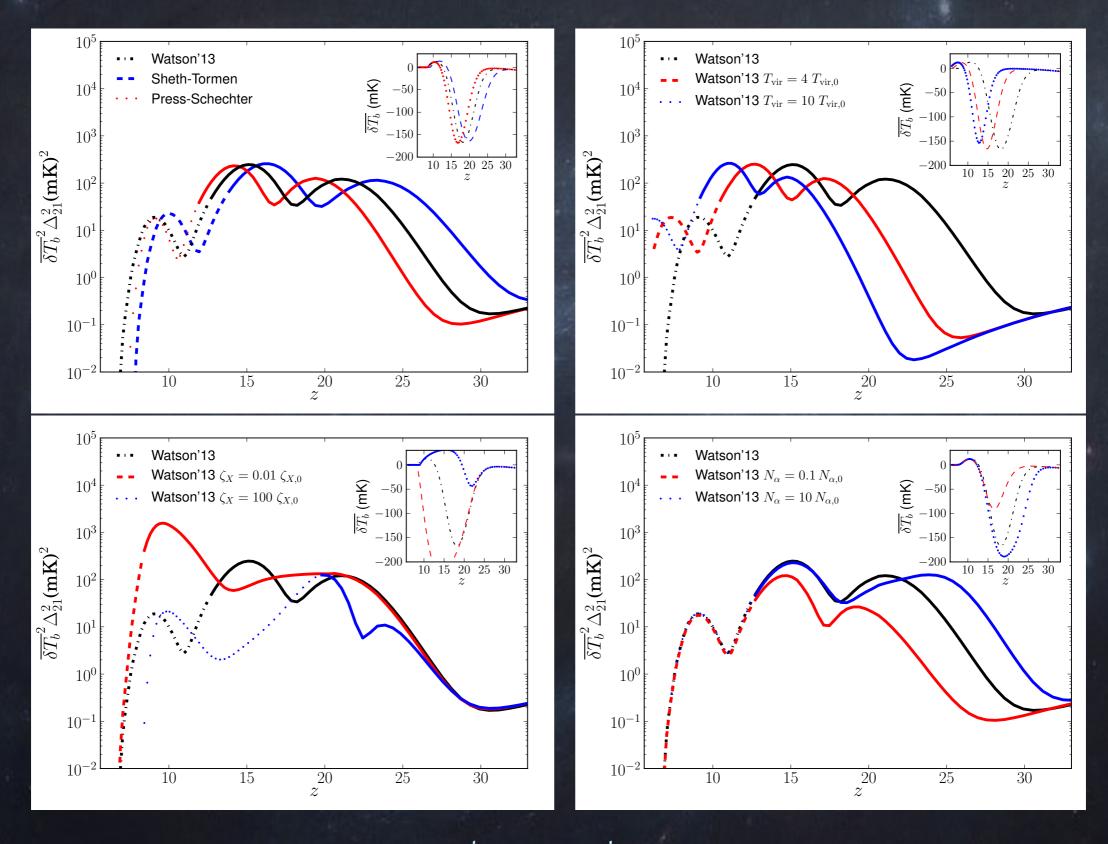
CONCLUSIONS

21cm radio observatories will be a powerful tool to learn about exotic energy injection mechanisms into the IGM (PBHs, particle DM...) during dark ages and cosmic dawn

Sensitivity to solar mass PBHs will be improved by up to 2 orders of magnitude with future interferometers (HERA, SKA)... and also great sensitivity to comet-size PBHs



DEPENDENCE ON ASTROPHYSICAL PARAMETERS





L. Lopez-Honorez, O. Mena, A. Molíné, SPR and A. C. Vincent, JCAP 1608:004, 2016