



YTF 2022



High Energy Jets Applied to Higgs plus jet(s) Processes

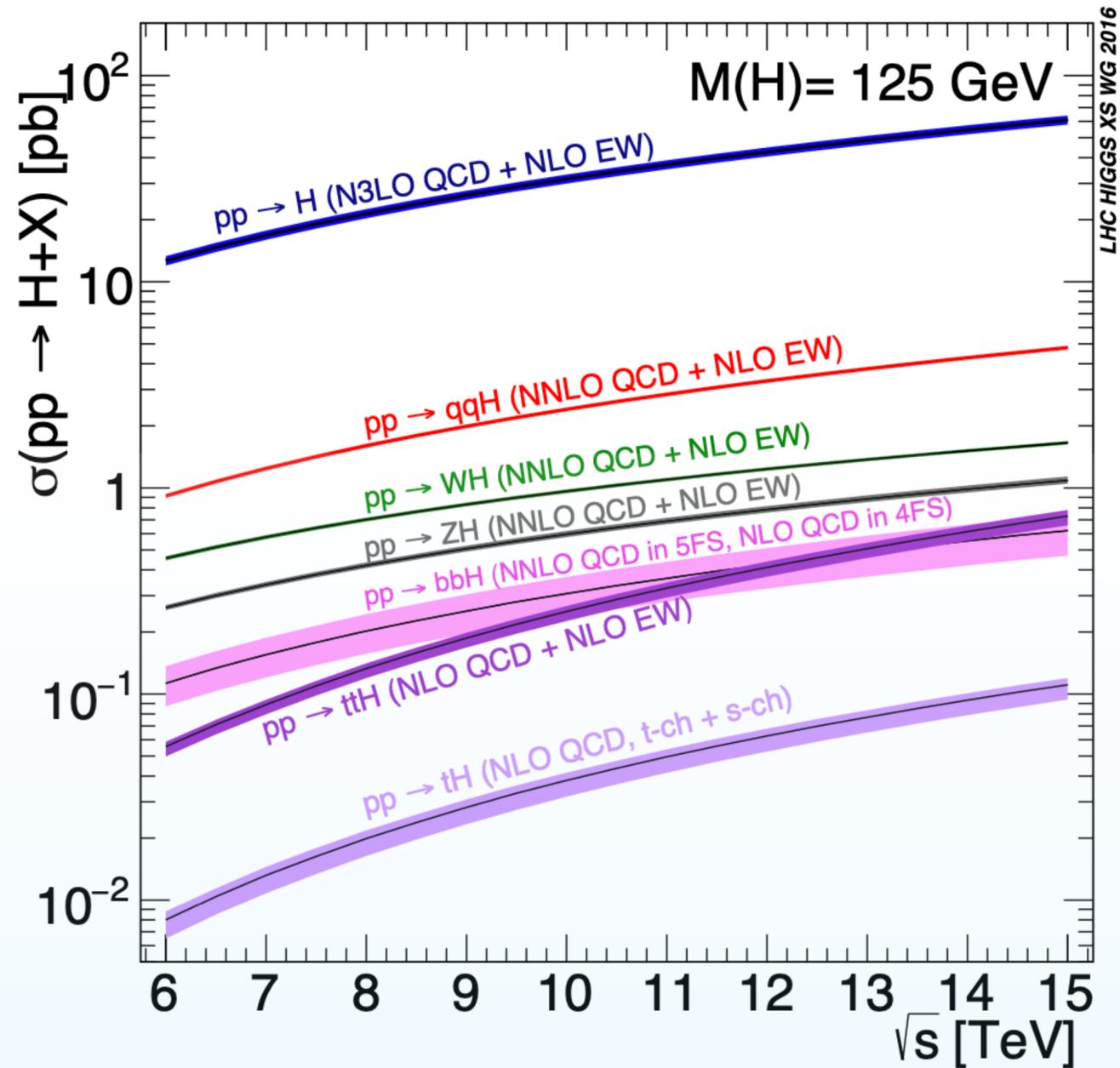
Jérémy Paltrinieri
University of Edinburgh

High Energy Jets (HEJ) Collaboration

16 December 2022 - Durham



Higgs production at the LHC



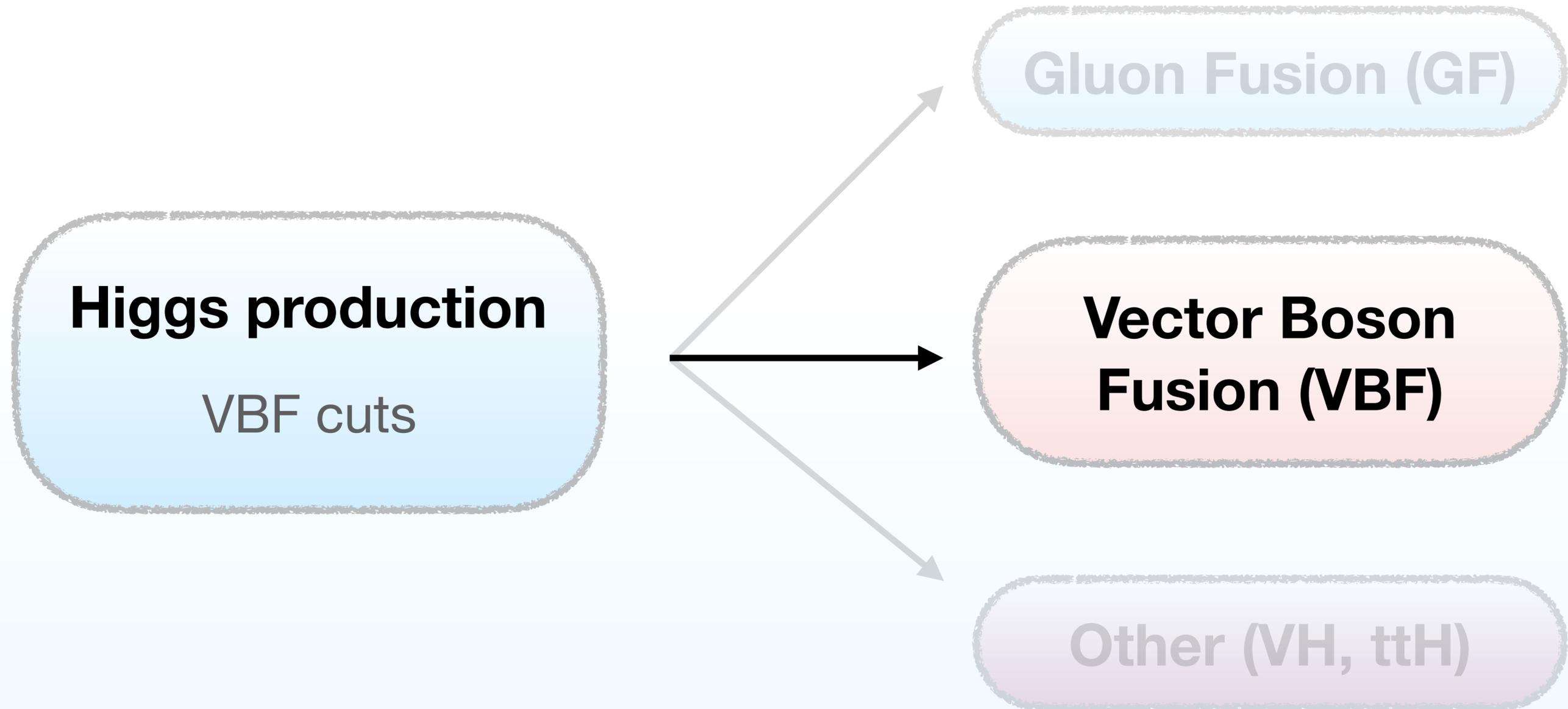
Gluon Fusion (GF)

Vector Boson Fusion (VBF)

Other (VH, ttH)



Higgs production at the LHC





Vector Boson Fusion (VBF) Cuts



Higgs production

VBF cuts

$$|\Delta y_{j_1 j_2}| > y_{\text{cut}}$$

$$m_{j_1 j_2} > m_{\text{cut}}$$



Vector Boson Fusion (VBF) Cuts



Higgs production

VBF cuts

$$|\Delta y_{j_1 j_2}| > y_{\text{cut}}$$

$$m_{j_1 j_2} > m_{\text{cut}}$$



Gluon Fusion (GF)

Enhance large logarithms
Damage perturbative expansion

VBF cuts make it difficult to get reliable QCD background predictions!

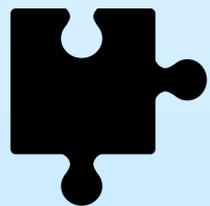


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HEJ Formalism

High Energy Limit
Resummation
Building Blocks
All-order results



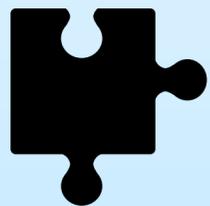


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Higgs + dijet

Theory
Finite quark masses
Comparisons to FO
VBF cuts



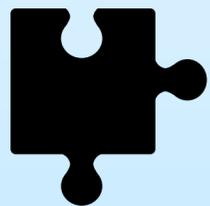


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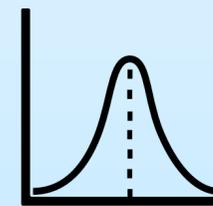
Higgs + dijet

Theory
Finite quark masses
Comparisons to FO
VBF cuts



Higgs + one jet

Theory
Comparisons to
experimental data
Comparisons to FO



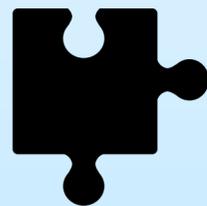


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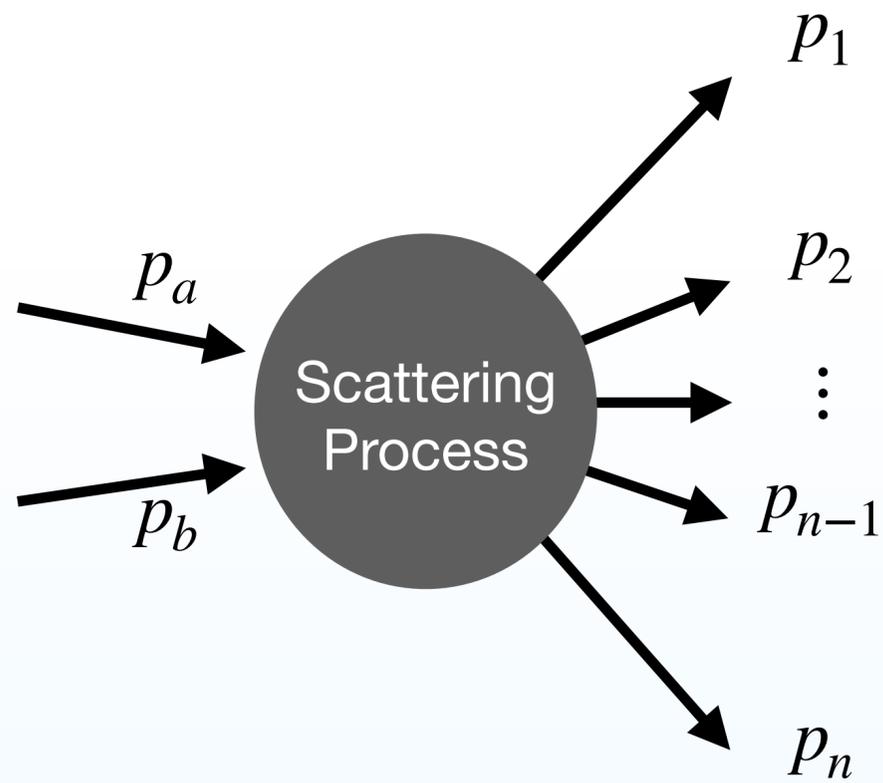
HEJ References

Constructing paper
[\[0908.2786\]](#)

Factorisation in qg
[\[0910.5113\]](#)

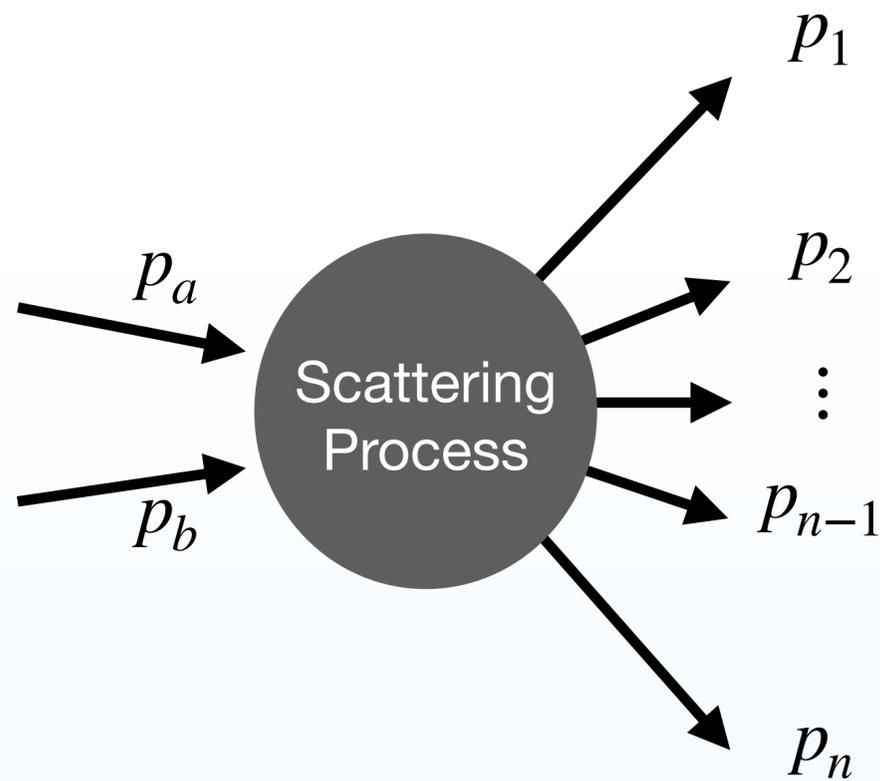


High Energy (HE) limit





High Energy (HE) limit



$$|p_{1\perp}| \approx |p_{2\perp}| \approx \dots \approx |p_{n\perp}| \text{ finite}$$
$$y_1 \gg y_2 \gg \dots \gg y_n$$

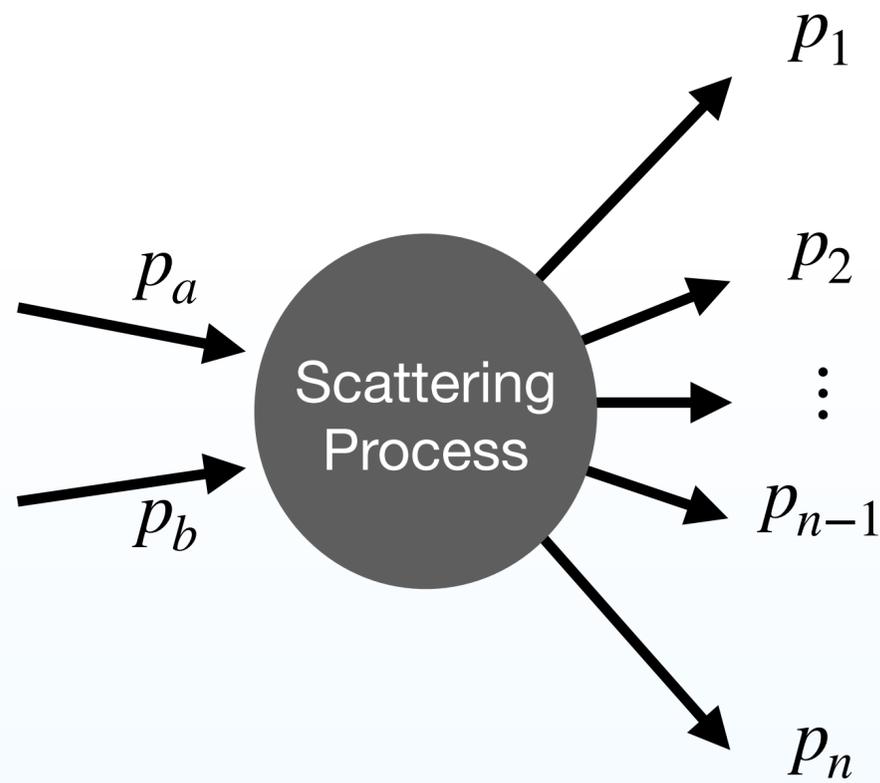
High Energy Limit

or equivalently

t – channel momenta squared finite
large invariant dijet masses $s_{i,i+1}$



High Energy (HE) limit

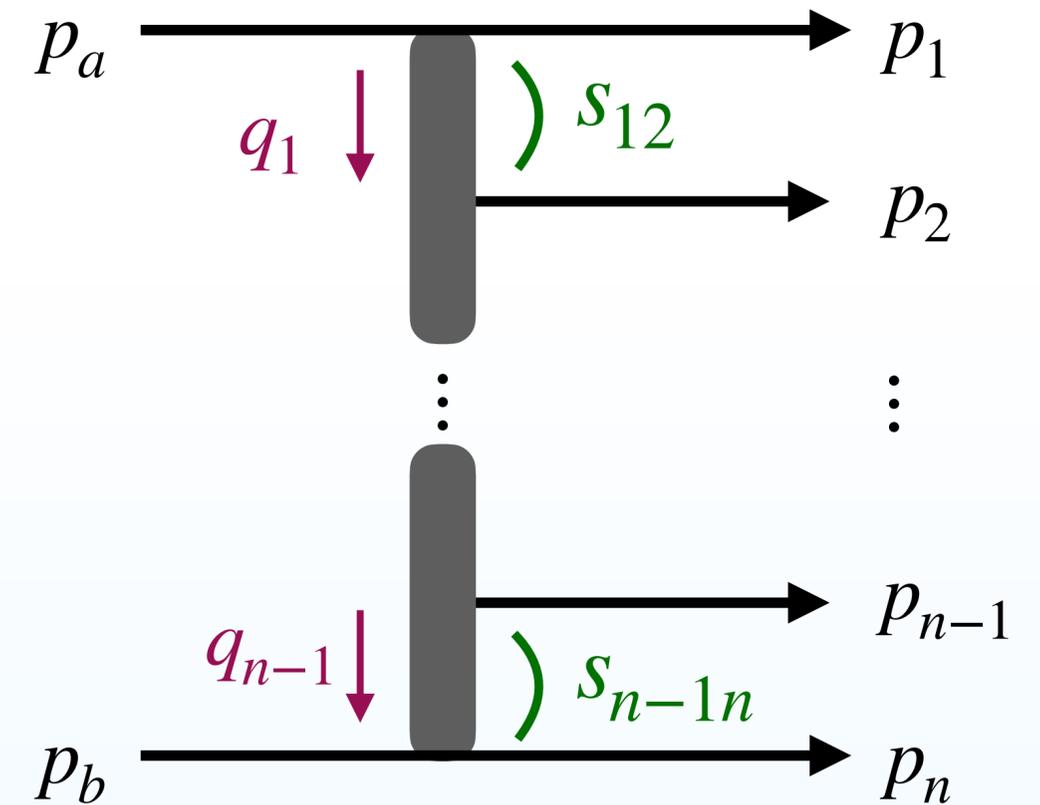


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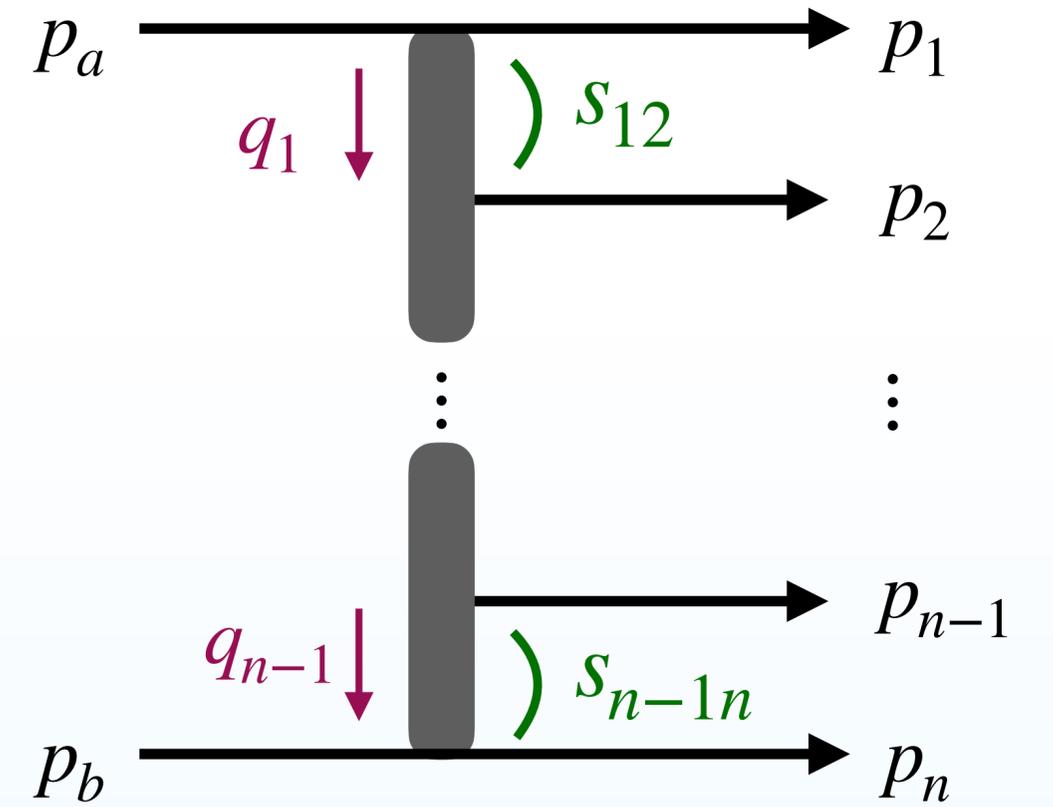
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 t – channel momenta squared finite
 large invariant dijet masses $s_{i,i+1}$





Regge scaling

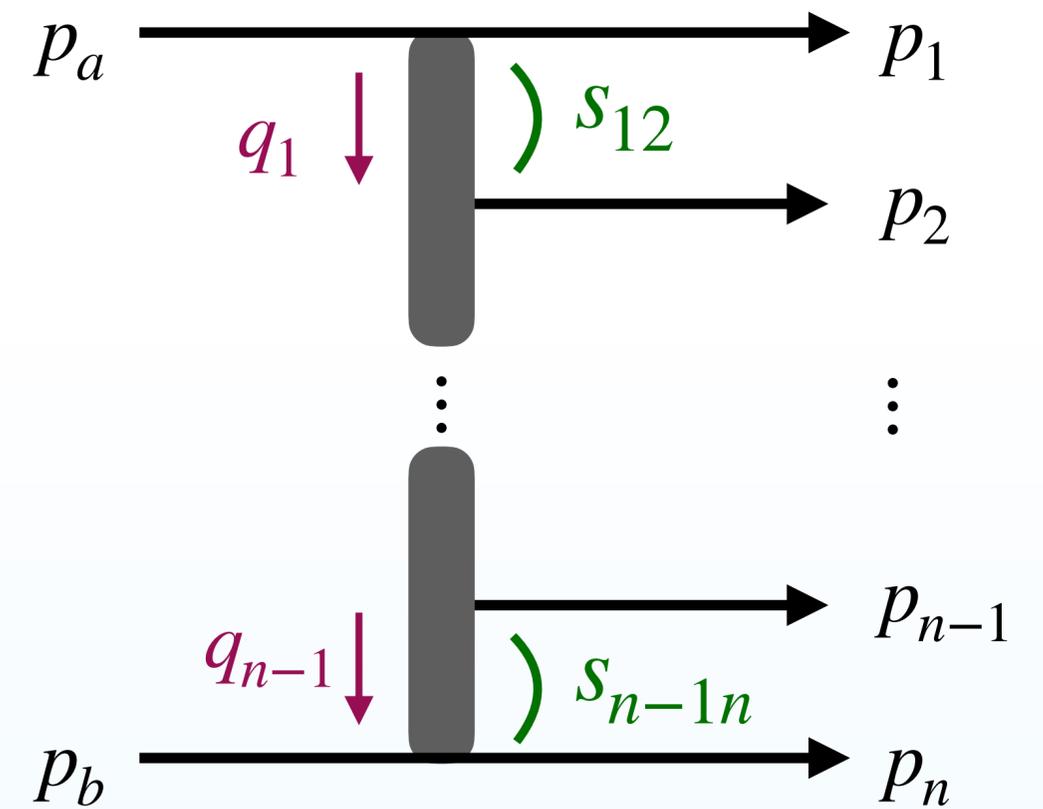




Regge scaling



Regge scaling: amplitudes = product of pieces
Get leading configurations in the HE limit:





Regge scaling

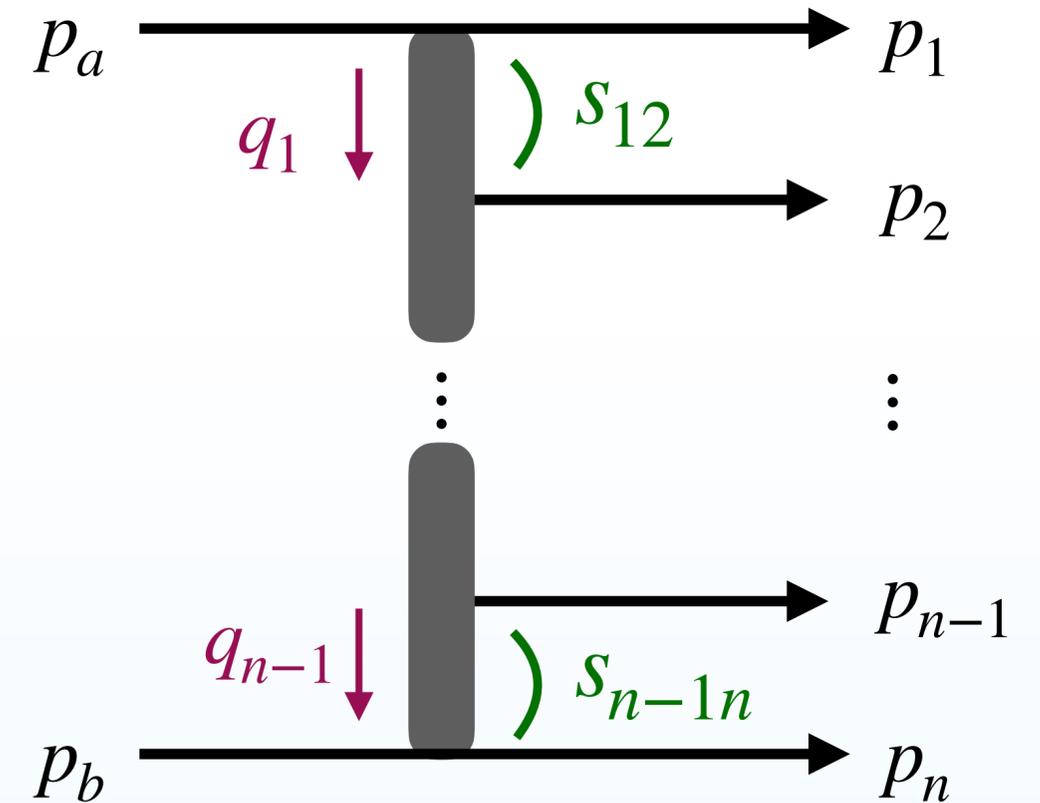


Regge scaling: amplitudes = product of pieces
Get leading configurations in the HE limit:

$$\mathcal{M} = s_{12}^{\alpha_1(q_1)} \cdots s_{n-1n}^{\alpha_{n-1}(q_{n-1})} \times \Gamma(q_1^2, \cdots, q_{n-1}^2)$$

Spin of particle q_1

Finite factor in the HE limit





Regge scaling

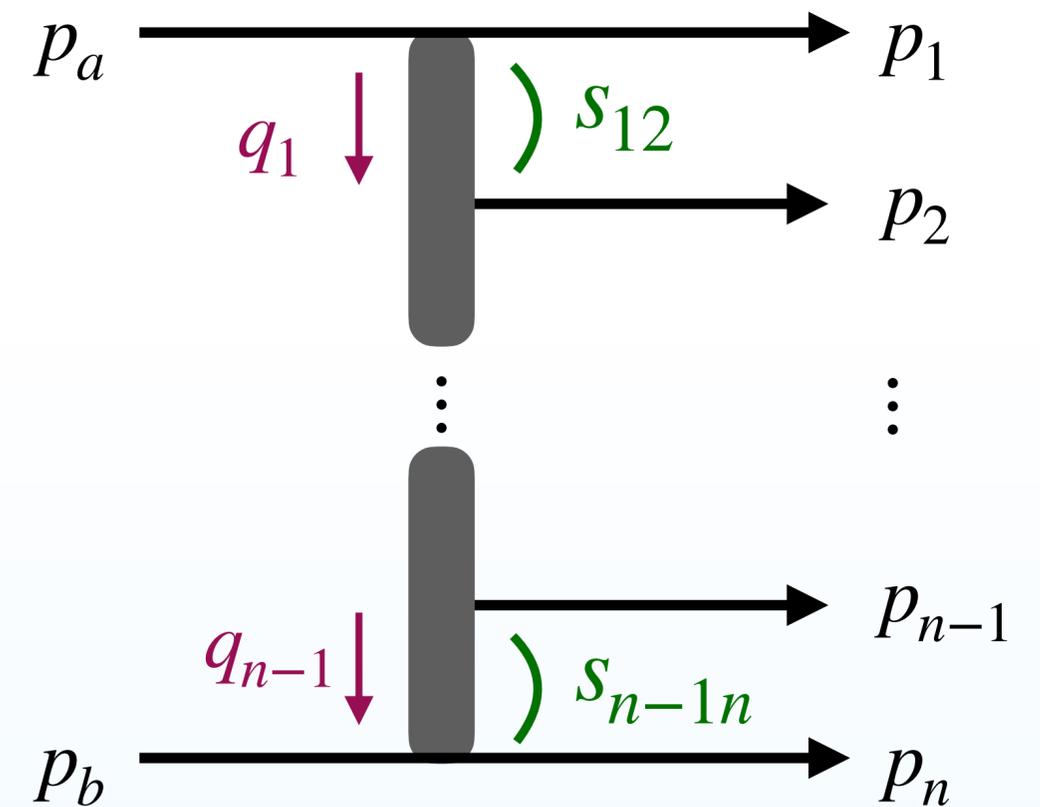


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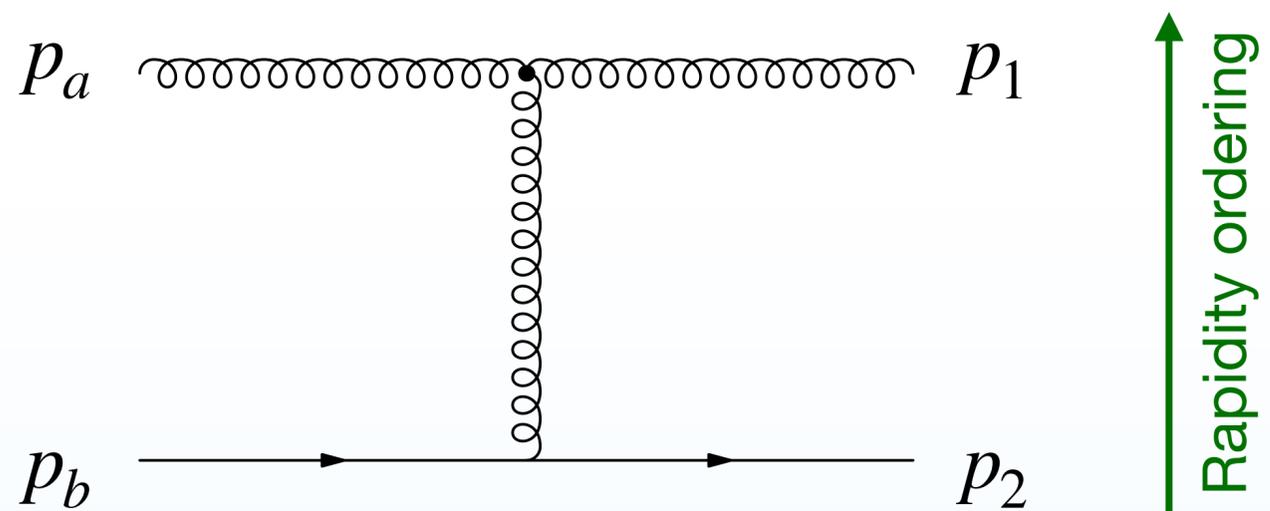
Finite factor in the HE limit



Leading configurations: maximise number of t-channel gluons exchanges

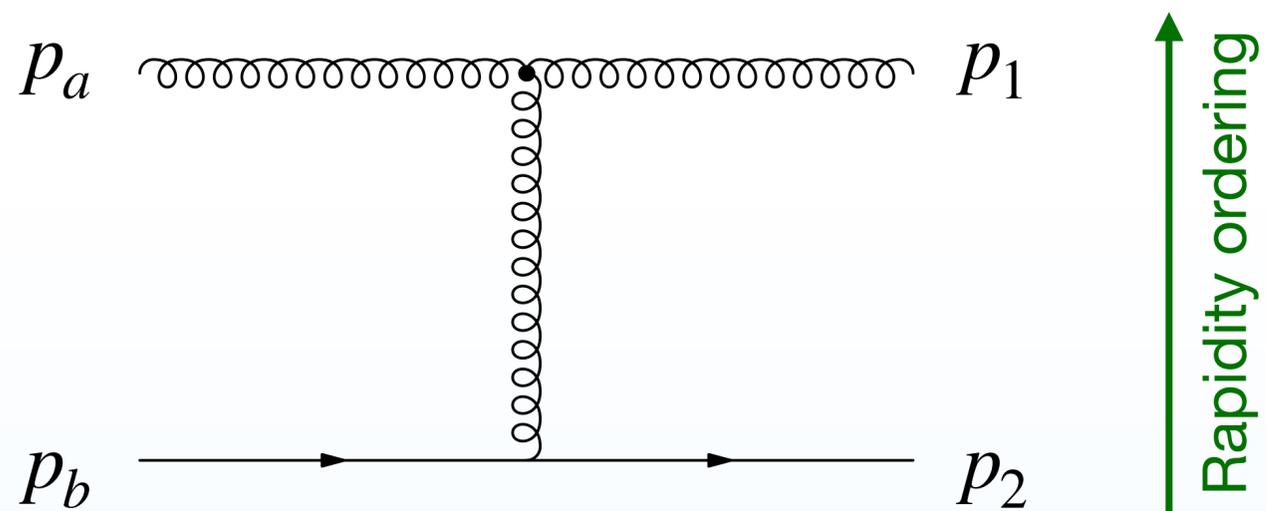


QCD scattering: quark-gluon





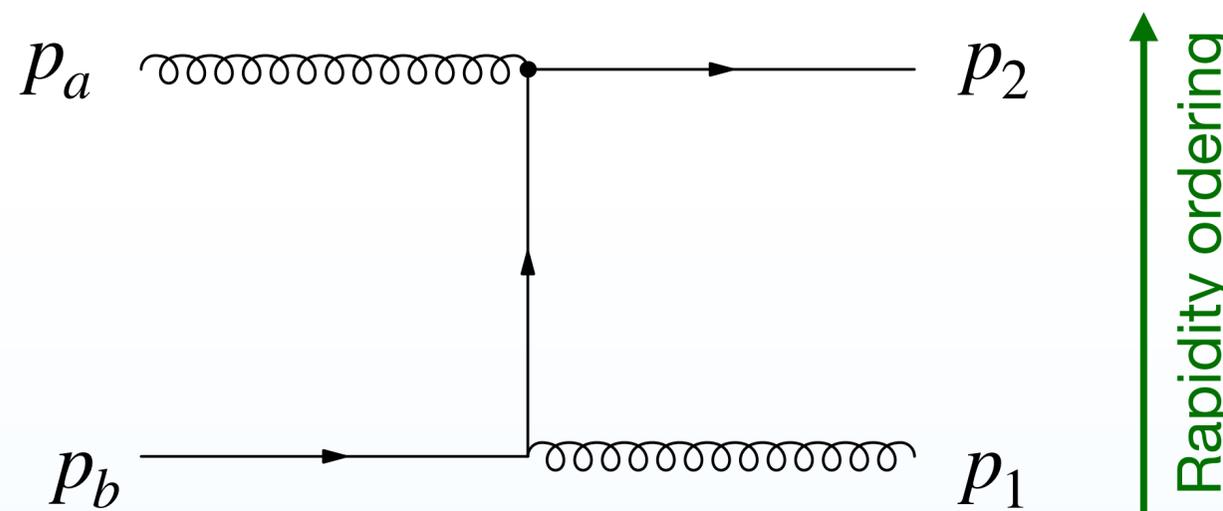
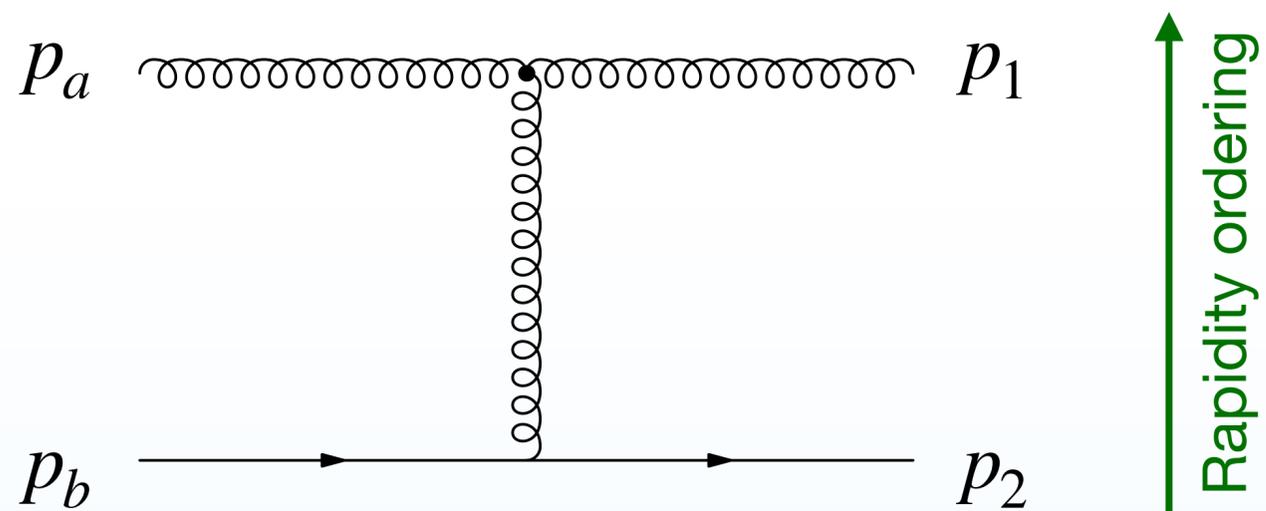
QCD scattering: quark-gluon



$$\mathcal{M} \propto s_{12}$$



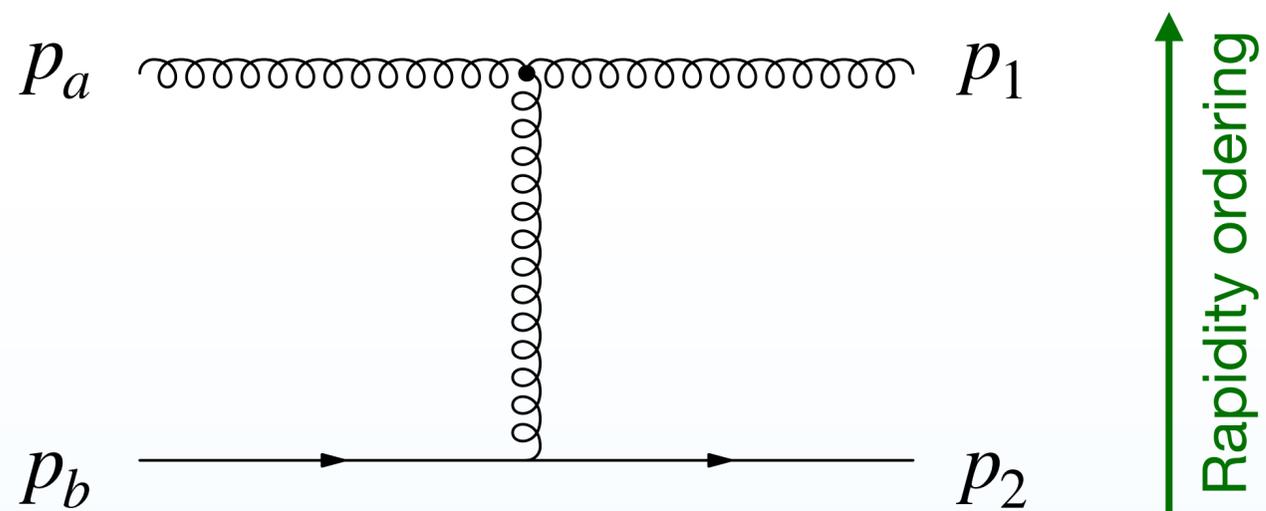
QCD scattering: quark-gluon



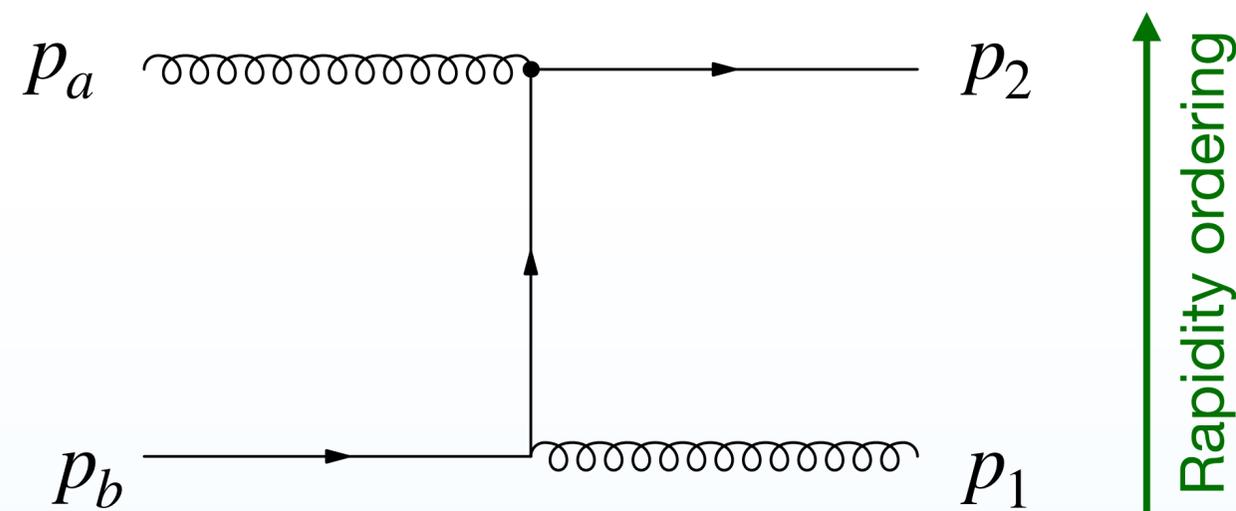
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QCD scattering: quark-gluon



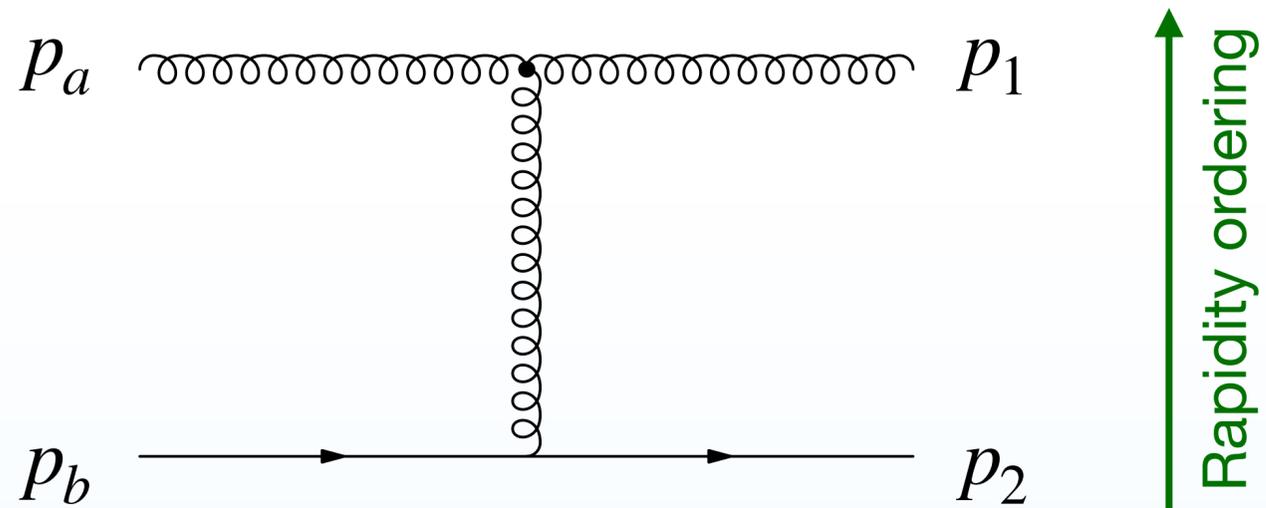
$$\mathcal{M} \propto s_{12}$$



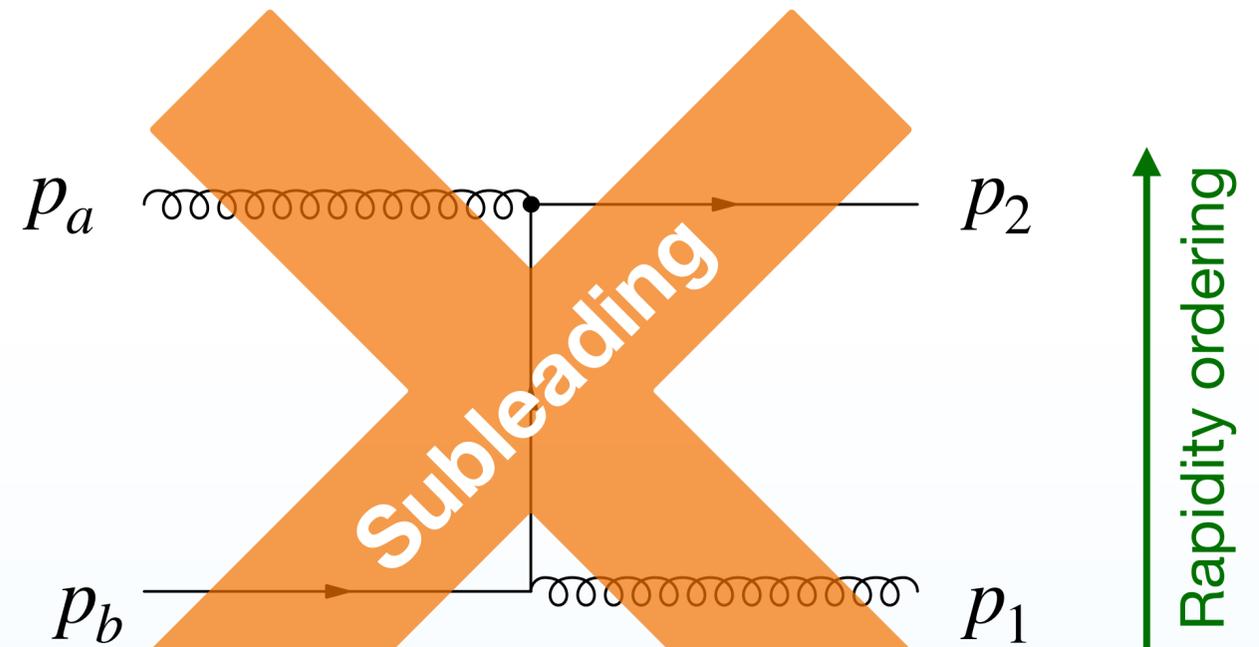
$$\mathcal{M} \propto \sqrt{s_{12}}$$



QCD scattering: quark-gluon



$$\mathcal{M} \propto s_{12}$$



$$\mathcal{M} \propto \sqrt{s_{12}}$$



The need for resummation



In the High Energy Phase-Space, the perturbative cross-section expansion contains factors of numerically significant logarithms:

$$\begin{aligned} |\mathcal{M}_{2j \text{ inc.}}|^2 &= \alpha_s^2 c_{\text{LO}} \\ &+ \alpha_s^3 c_{\text{NLO}} \\ &+ \alpha_s^4 c_{\text{NNLO}} \\ &+ \dots \end{aligned}$$

Perturbative expansion valid
as long as coefficients do
not grow too much



The need for resummation



In the High Energy Phase-Space, the perturbative cross-section expansion contains factors of numerically significant logarithms:

$$\begin{aligned} |\mathcal{M}_{2j \text{ inc.}}|^2 &= \alpha_s^2 c_{\text{LO}} \\ &+ \alpha_s^3 (c_{11} \log(s/|t|) + c_{12}) \\ &+ \alpha_s^4 (c_{21} \log^2(s/|t|) + c_{22} \log(s/|t|) + c_{23}) \\ &+ \dots \end{aligned}$$



The need for resummation



In the High Energy Phase-Space, the perturbative cross-section expansion contains factors of numerically significant logarithms:

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Leading Logarithm (LL)



The need for resummation



In the High Energy Phase-Space, the perturbative cross-section expansion contains factors of numerically significant logarithms:

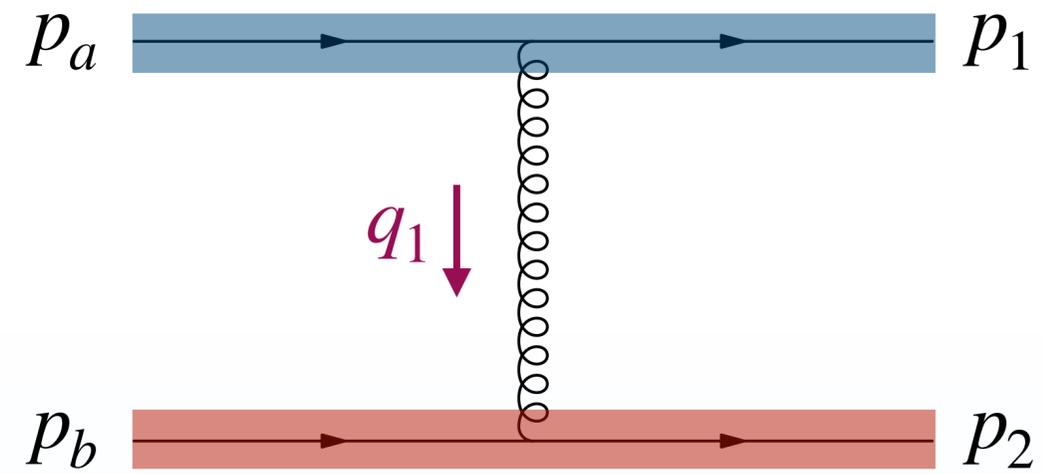
$$\begin{aligned} |\mathcal{M}_{2j \text{ inc.}}|^2 &= \alpha_s^2 c_{\text{LO}} \\ &+ \alpha_s^3 (c_{11} \log(s/|t|) + c_{12}) \\ &+ \alpha_s^4 (c_{21} \log^2(s/|t|) + c_{22} \log(s/|t|) + c_{23}) \\ &+ \dots \end{aligned}$$

Leading Logarithm (LL)

**Next-to-Leading
Logarithm (NLL)**

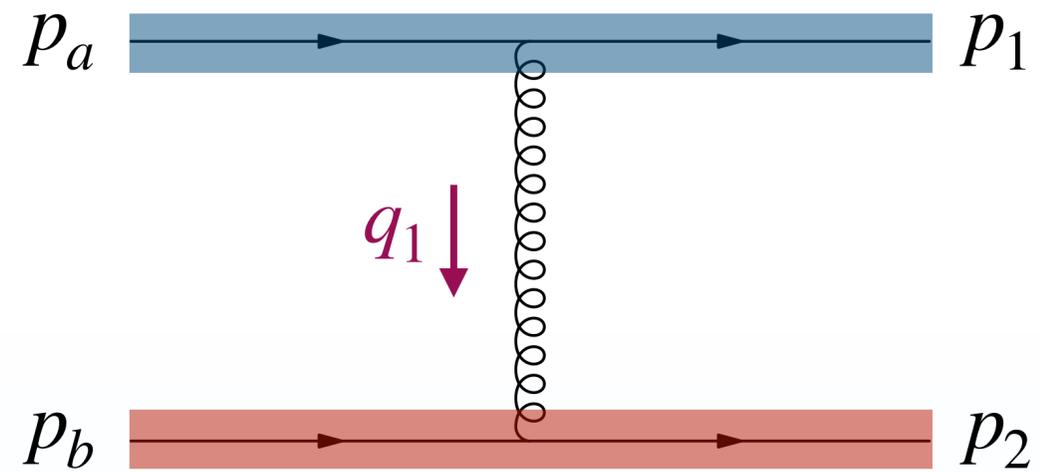


Building blocks of HEJ





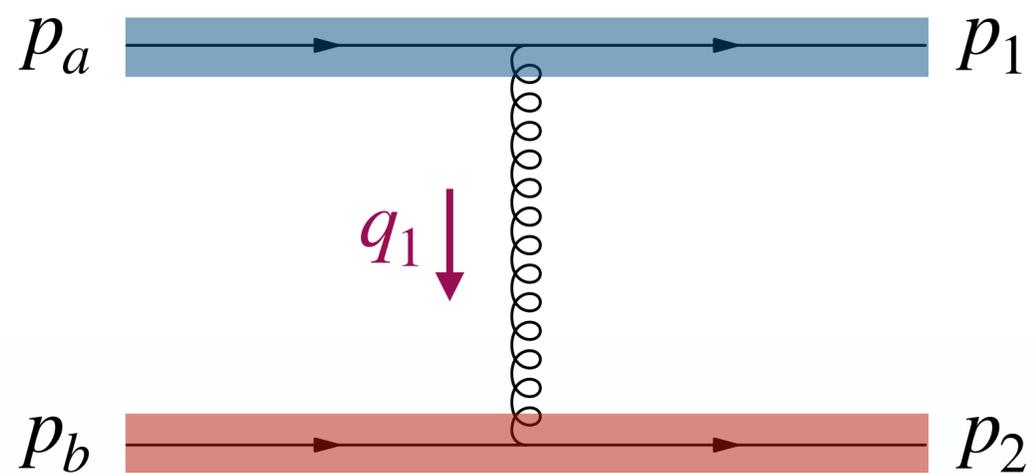
Building blocks of HEJ



$$|\mathcal{M}|^2 \propto C_F^2 \left(\frac{1}{q_1^2} \right)^2 |j^\mu(p_1, p_a) j_\mu(p_2, p_b)|^2$$



Building blocks of HEJ

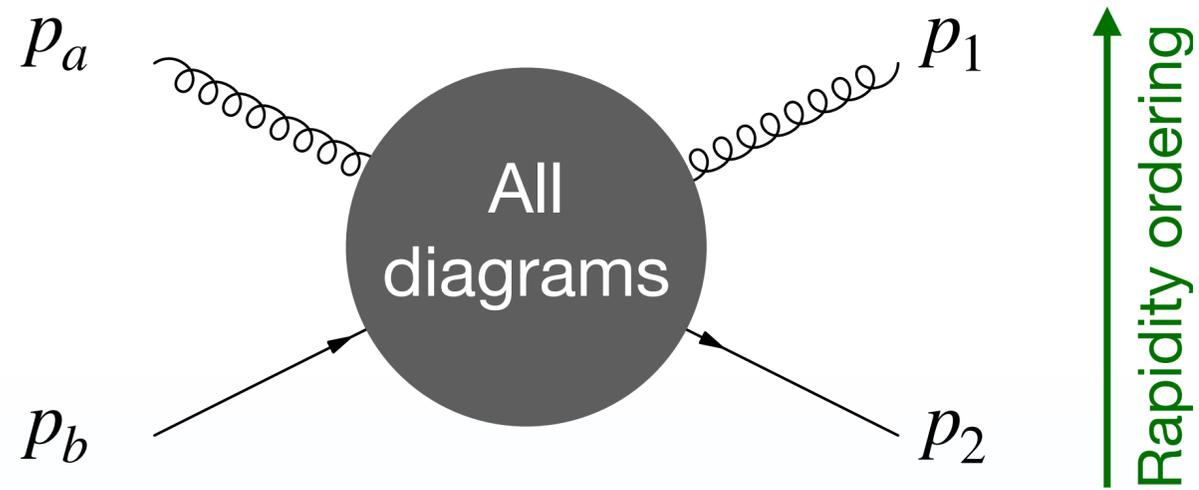


- ◆ Contraction of currents over a t-channel pole
- ◆ Looks natural: only one diagram contributes at tree-level!
- ◆ What about gluon-induced processes?

$$|\mathcal{M}|^2 \propto C_F^2 \left(\frac{1}{q_1^2} \right)^2 |j^\mu(p_1, p_a) j_\mu(p_2, p_b)|^2$$

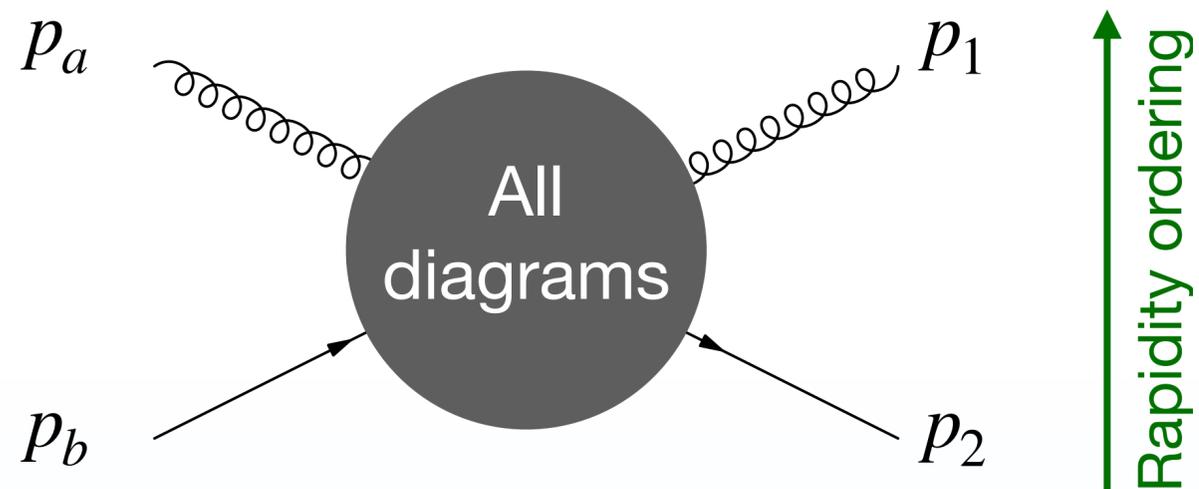


Building blocks of HEJ

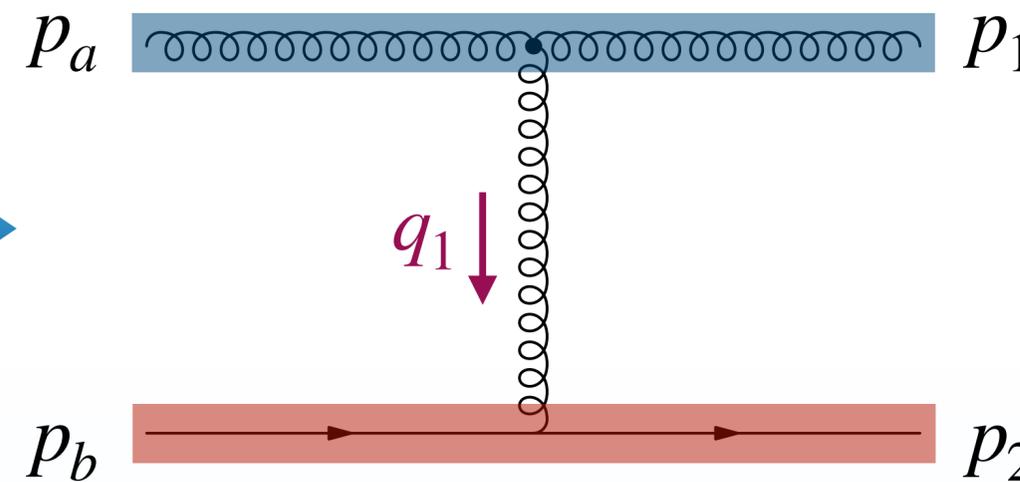




Building blocks of HEJ



High Energy Limit





Building blocks of HEJ

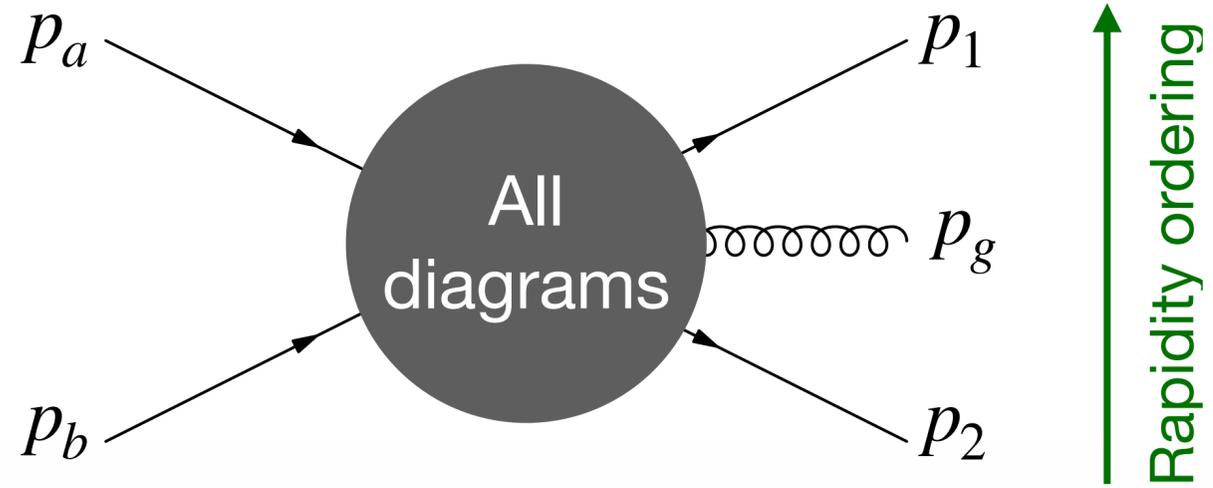


$$|\mathcal{M}|^2 \propto C_F \times \text{CAM} \times \left(\frac{1}{q_1^2} \right)^2 |j^\mu(p_1, p_a) j_\mu(p_2, p_b)|^2$$

CAM \rightarrow C_A in the High Energy Limit

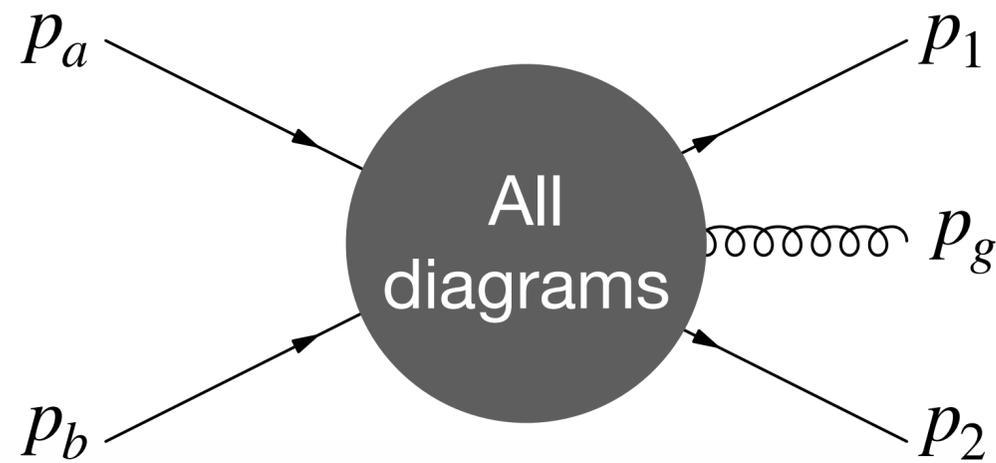


Real Corrections



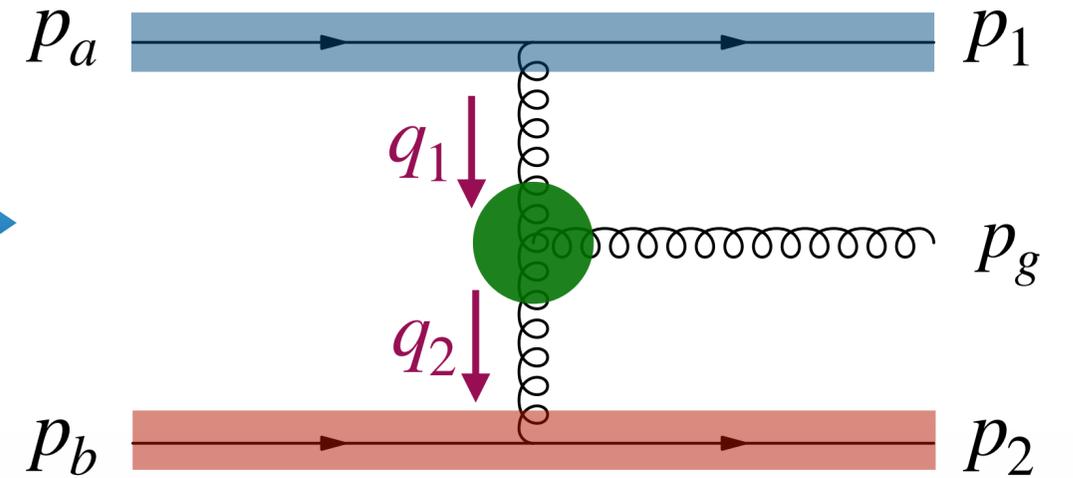


Real Corrections



Rapidity ordering

High Energy Limit





Real Corrections

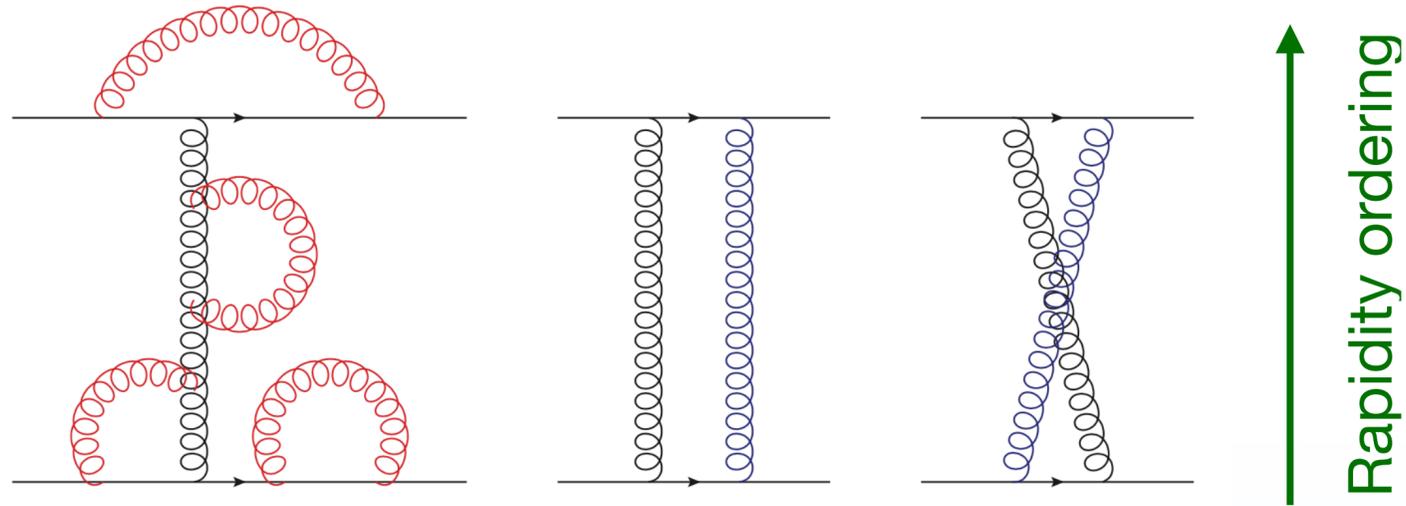


$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} j^\mu(p_1, p_a) j_\mu(p_2, p_b) V^\rho \epsilon_\rho^*(p_g)$$

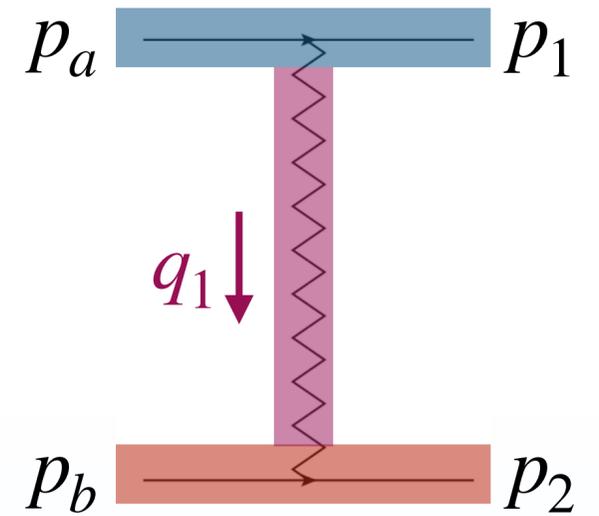
Effective Lipatov vertex, gauge invariant



Virtual Corrections

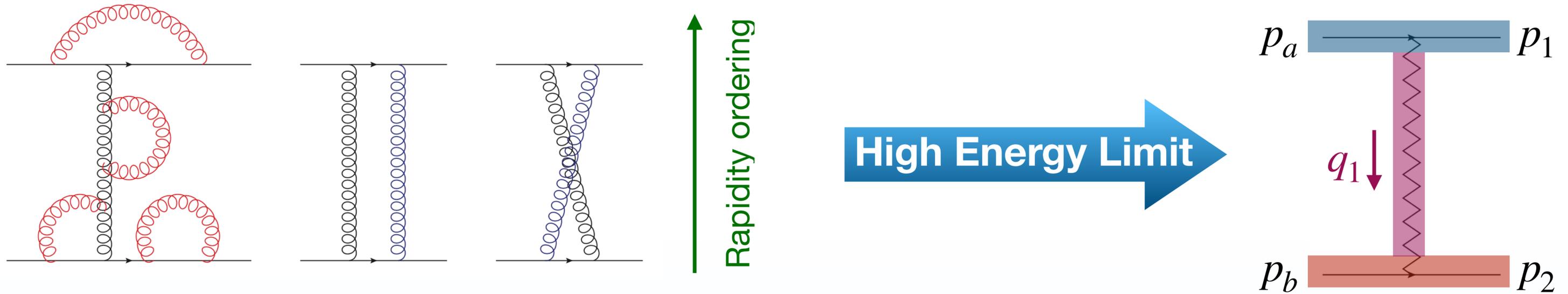


High Energy Limit





Virtual Corrections



Finite expression: soft divergences cancel with real corrections

$$\mathcal{M} \propto \frac{\exp(\alpha(q_1)\Delta y_{12})}{q_1^2} j^\mu(p_1, p_a) j_\mu(p_2, p_b)$$

The gluon in the t-channel reggeizes: Lipatov Ansatz (valid even at NLL)



All-order Corrections





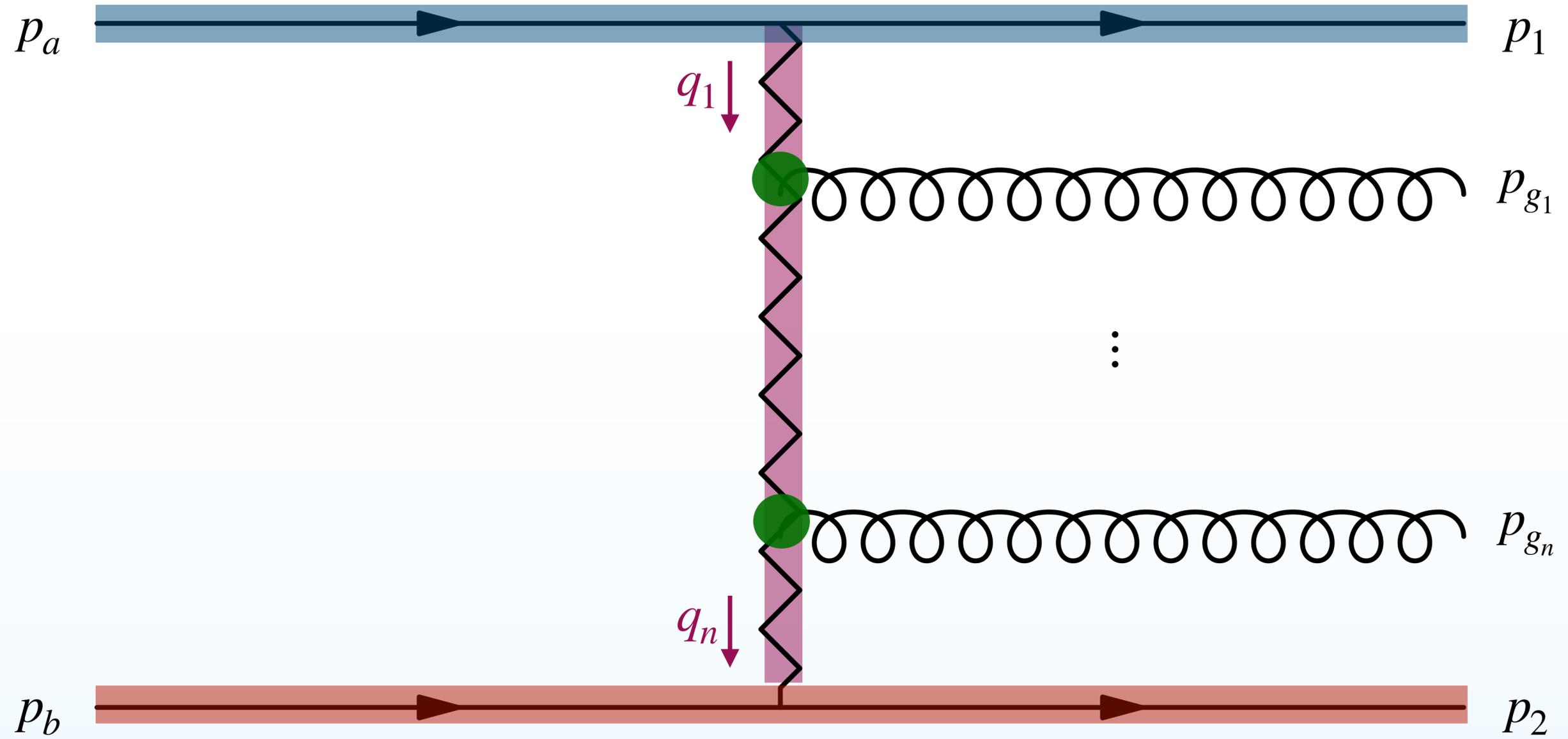
All-order Corrections



$$\begin{aligned}
 \mathcal{M} &\propto j^\mu(p_1, p_a) j_\mu(p_2, p_b) \\
 &\times V_1^{\rho_1} \epsilon_{\rho_1}^*(p_{g_1}) V_1^{\rho_{n-1}} \epsilon_{\rho_{n-1}}^*(p_{g_{n-1}}) \\
 &\times \frac{\exp(\alpha(q_1) \Delta y_{12})}{q_1^2} \dots \frac{\exp(\alpha(q_n) \Delta y_{n-1n})}{q_n^2}
 \end{aligned}$$



All-order Corrections

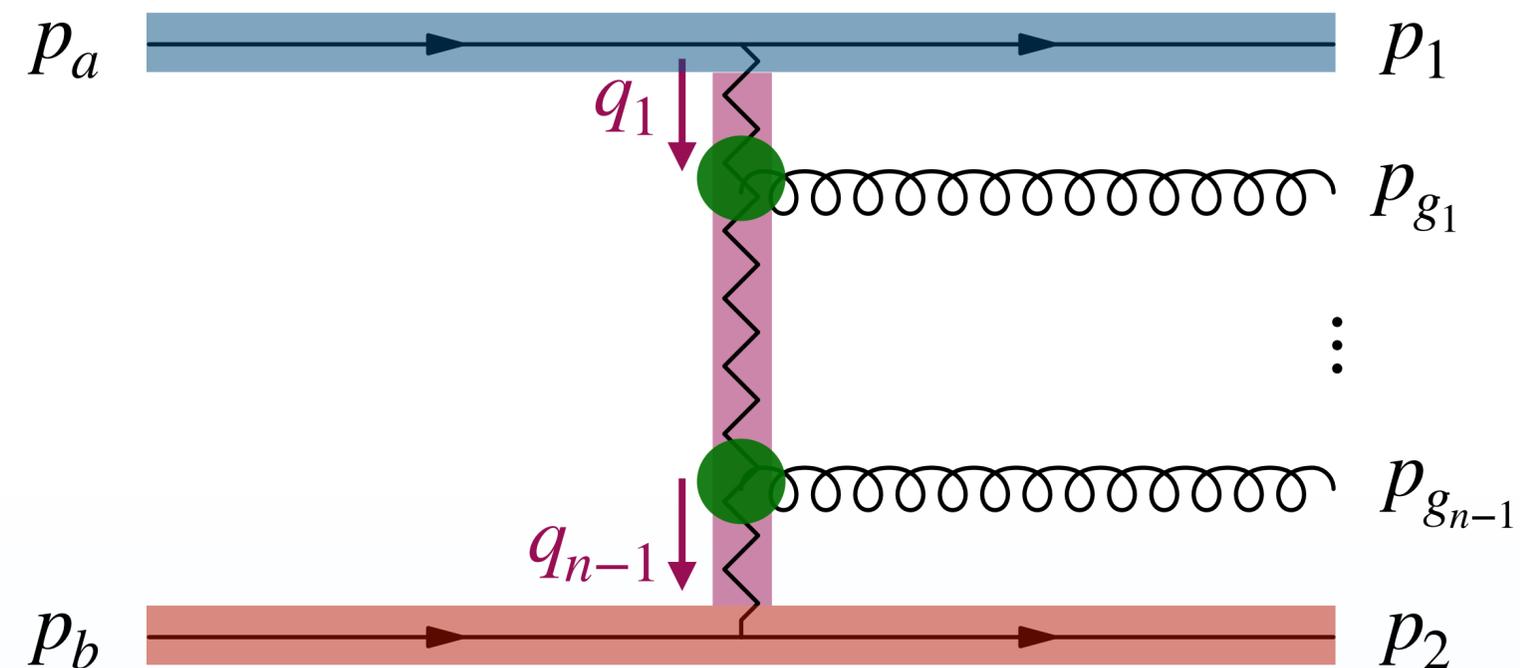




All-order Corrections



- ◆ All-order leading-log results
- ◆ Gauge-invariant in all phase-space
- ◆ Phase-space not approximated
- ◆ Monte-Carlo integration
- ◆ IR divergences cancel



$$\begin{aligned}
 \mathcal{M} &\propto j^\mu(p_1, p_a) j_\mu(p_2, p_b) \\
 &\times V_1^{\rho_1} \epsilon_{\rho_1}^*(p_{g_1}) V_1^{\rho_{n-1}} \epsilon_{\rho_{n-1}}^*(p_{g_{n-1}}) \\
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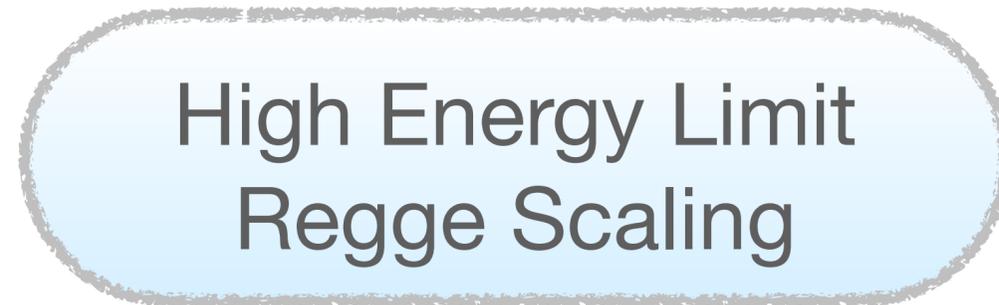
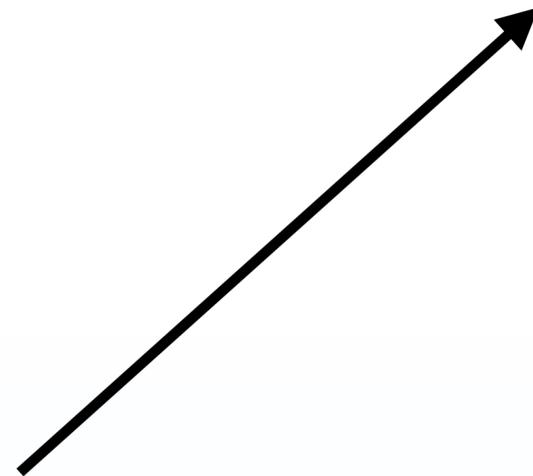
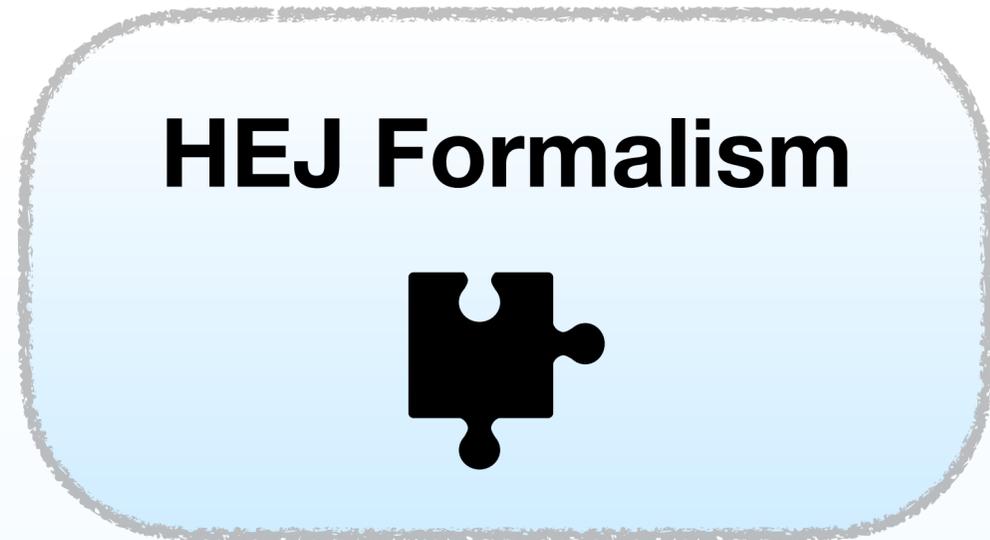


HEJ Formalism



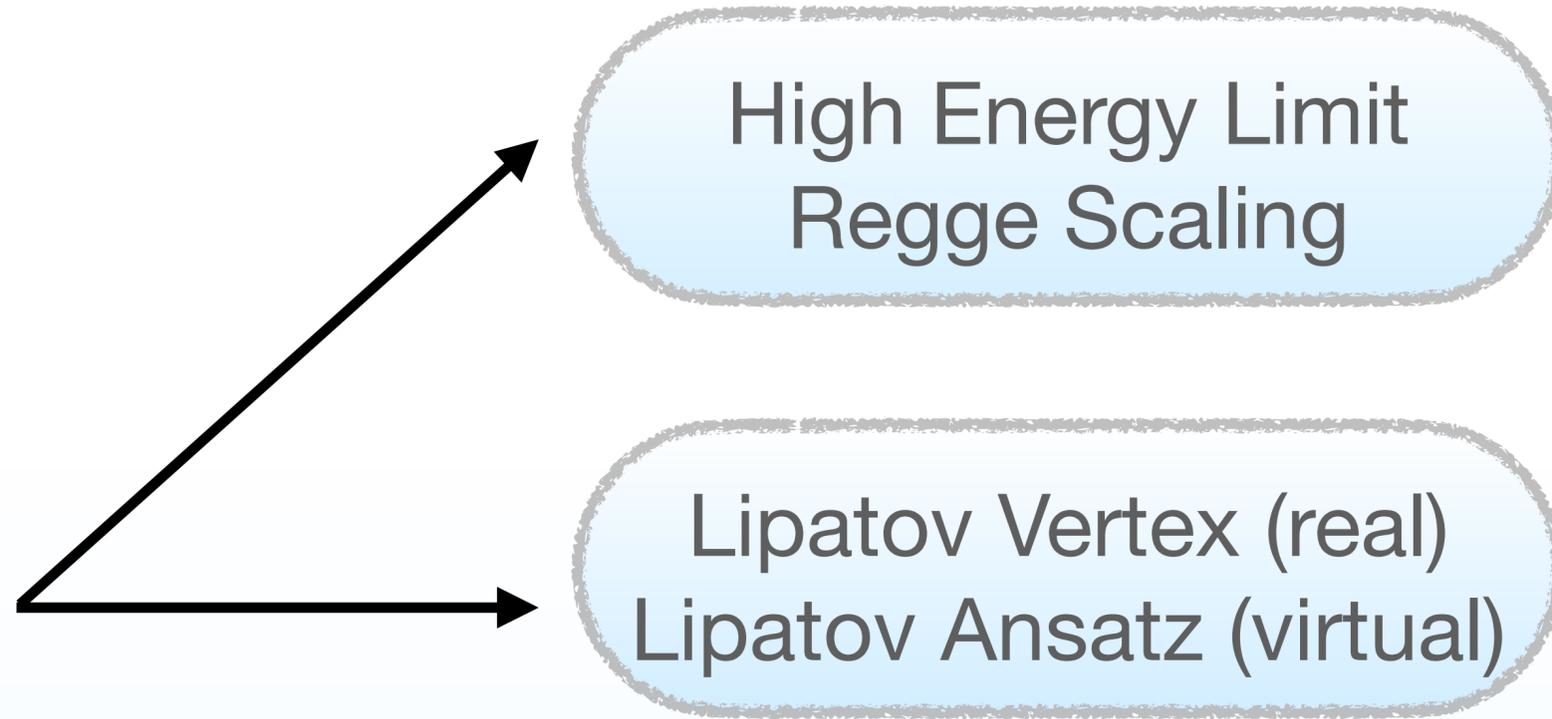
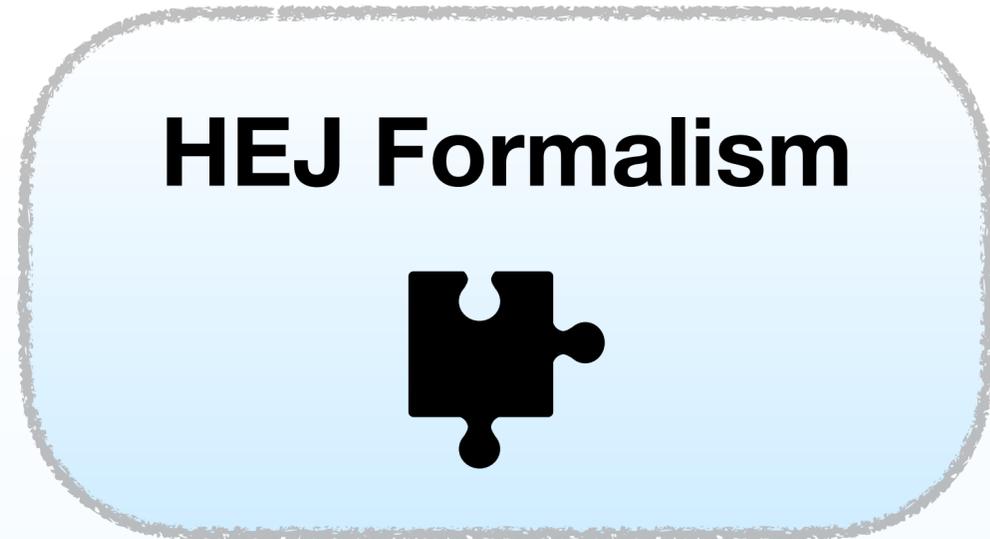


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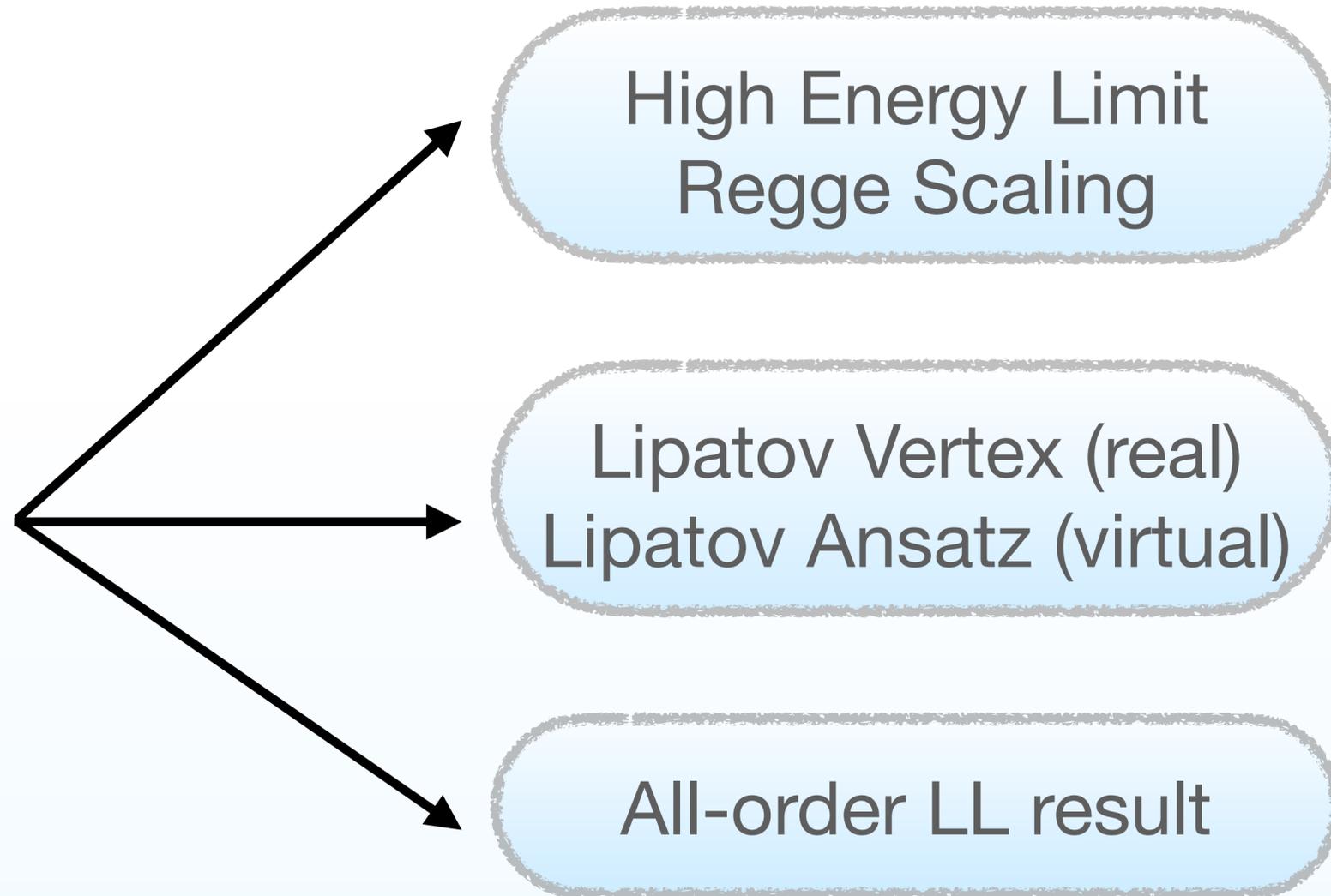
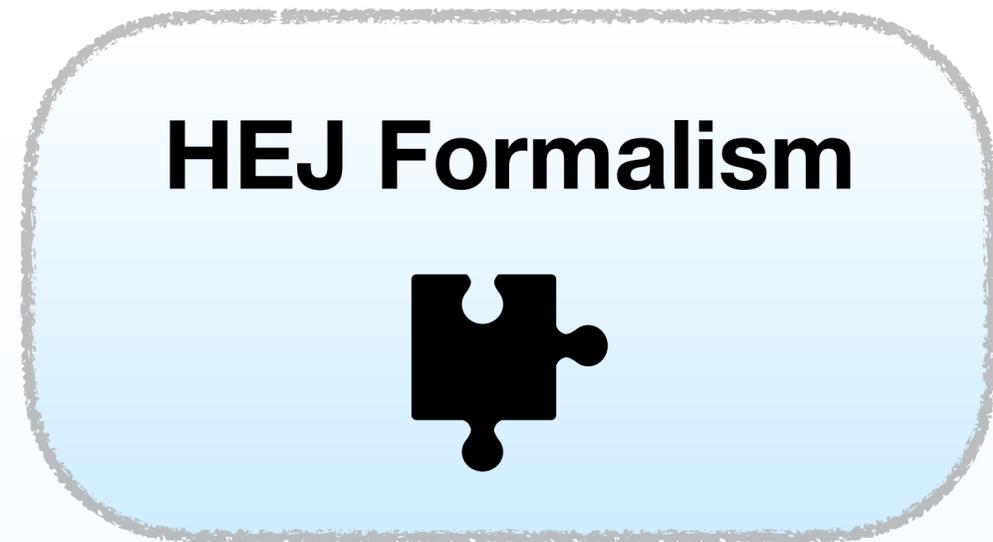


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YTF 2022





YTF 2022



Higgs + dijet

Theory

Finite quark masses

Comparisons to FO

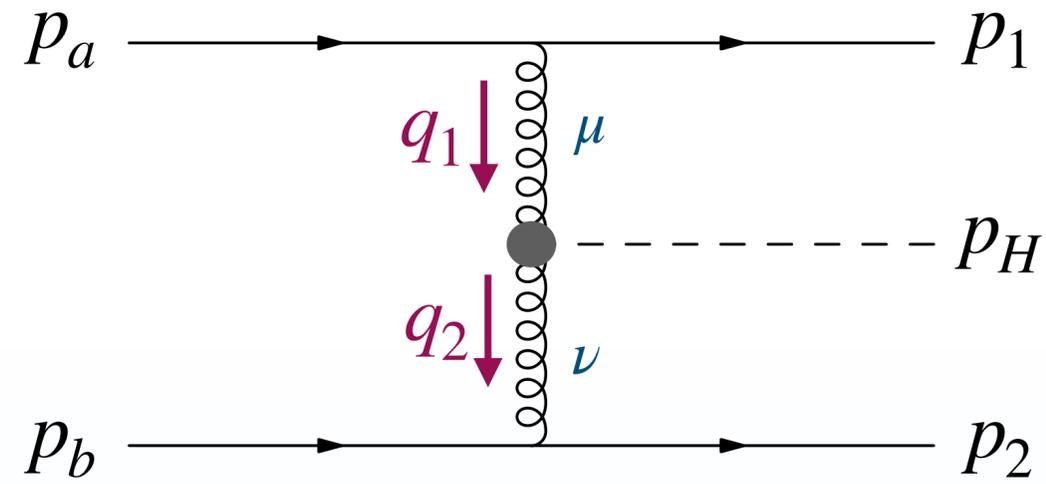
VBF cuts

[1812.08072]



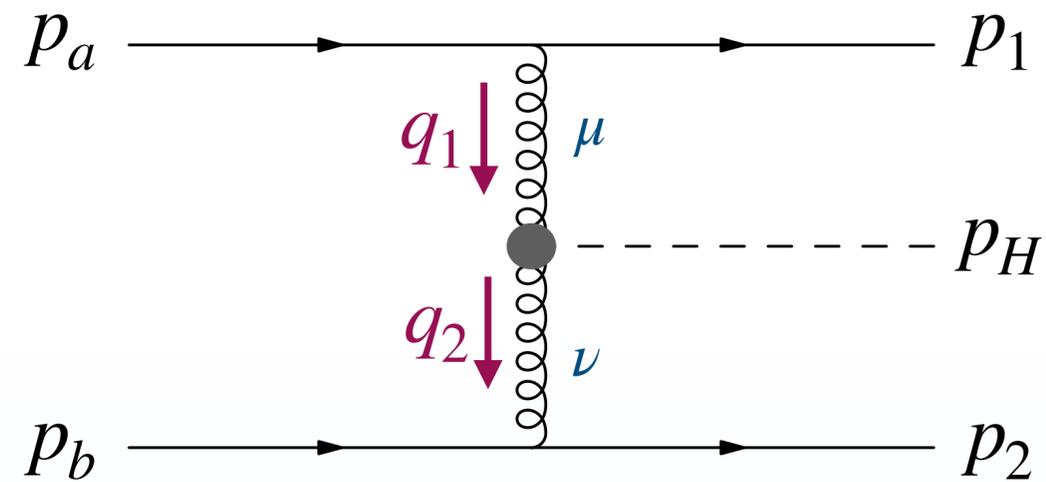


Central Higgs Production





Central Higgs Production

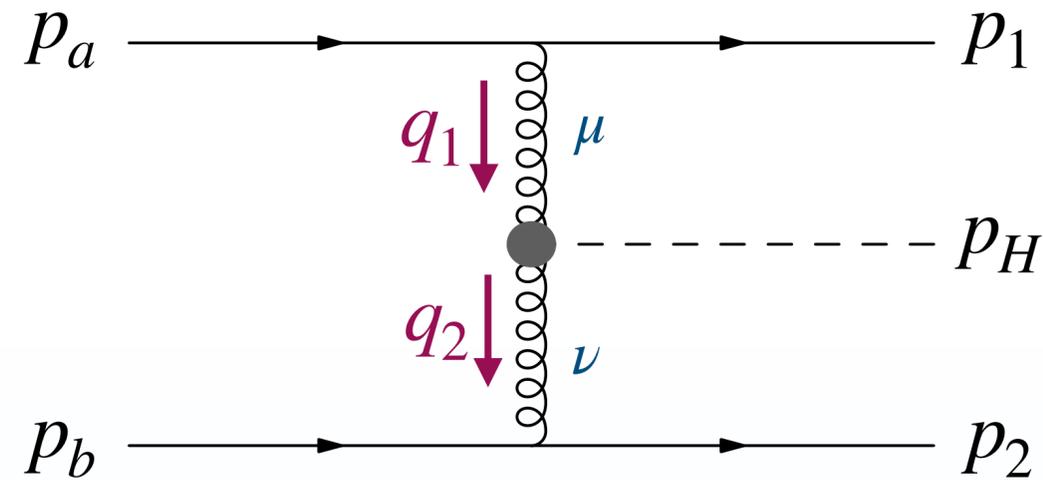


ggH vertex has the following tensor form:

$$V_H^{\mu\nu}(q_1, q_2) = \frac{\alpha_s m^2}{\pi v} (g^{\mu\nu} T_1(q_1, q_2) - q_2^\mu q_1^\nu T_2(q_1, q_2))$$
$$\xrightarrow{m \rightarrow \infty} \frac{\alpha_s}{3\pi v} (g^{\mu\nu} q_1 \cdot q_2 - q_2^\mu q_1^\nu)$$



Central Higgs Production



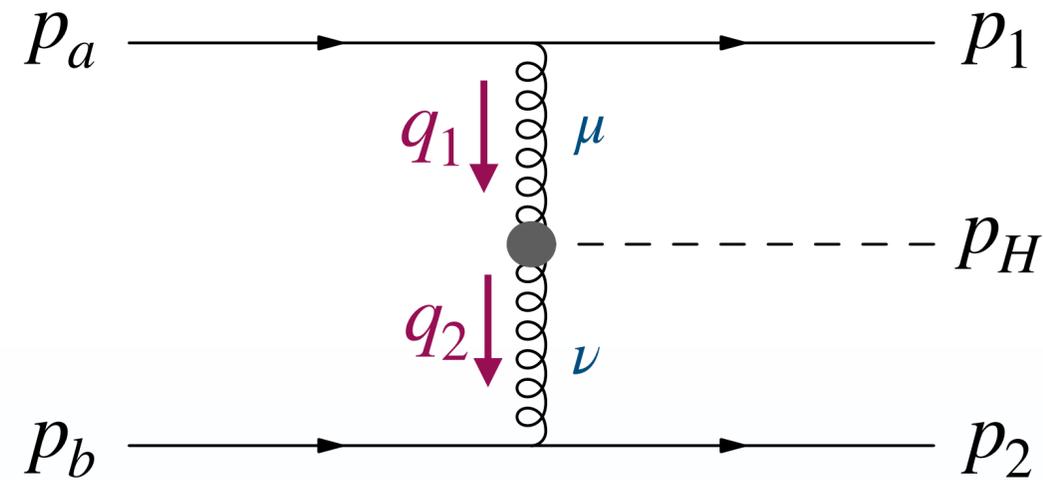
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~~$$\frac{\alpha_s}{3\pi v} (g^{\mu\nu} q_1 \cdot q_2 - q_2^\mu q_1^\nu)$$~~



Central Higgs Production



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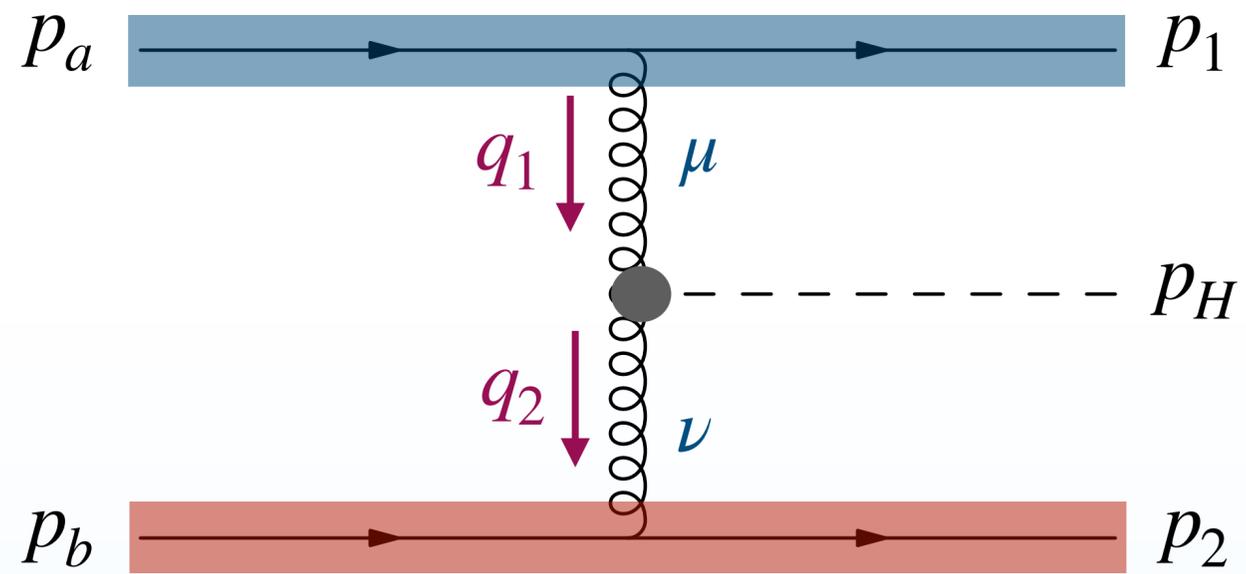
$$V_H^{\mu\nu}(q_1, q_2) = \frac{\alpha_s m^2}{\pi v} (g^{\mu\nu} T_1(q_1, q_2) - q_2^\mu q_1^\nu T_2(q_1, q_2))$$

~~$$\frac{\alpha_s}{3\pi v} (g^{\mu\nu} q_1 \cdot q_2 - q_2^\mu q_1^\nu)$$~~

t-channel factorised form allows us to keep the full quark masses dependence

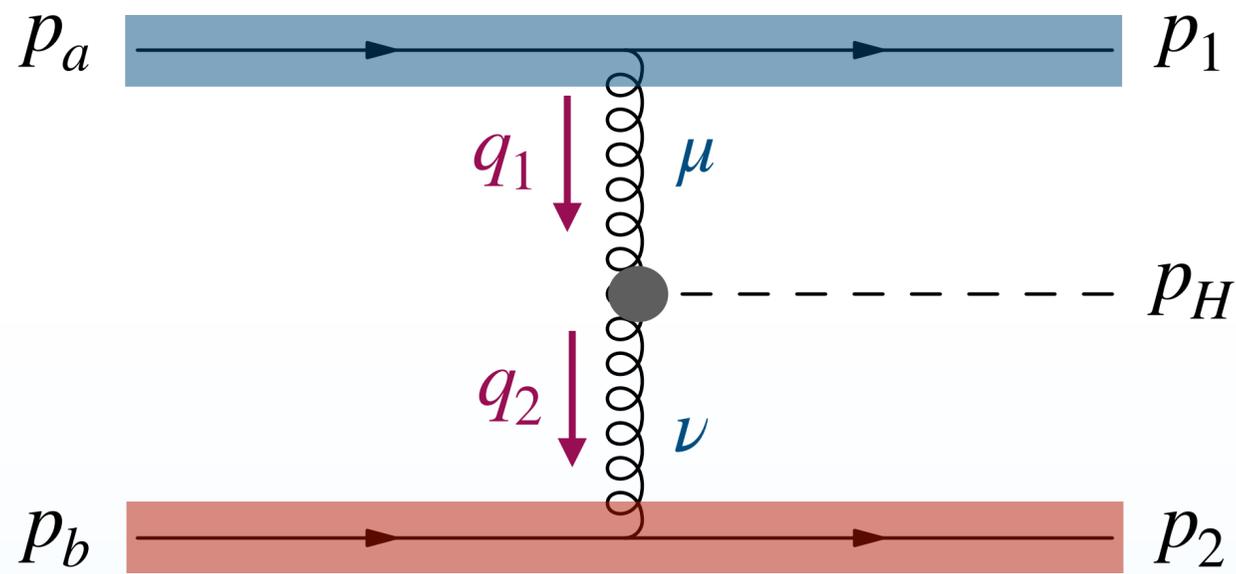


Central Higgs Production





Central Higgs Production

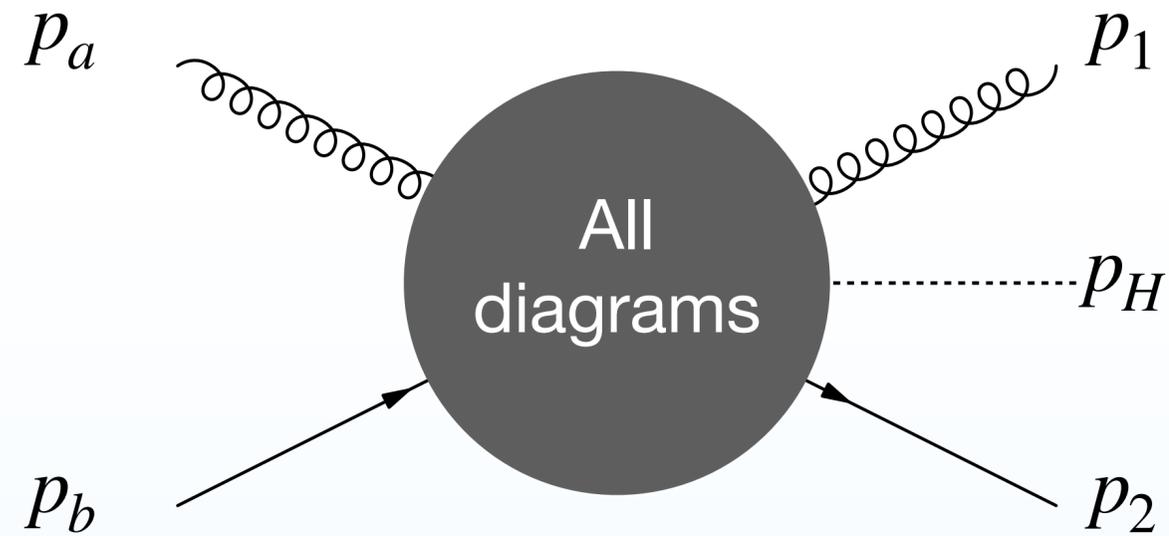


$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} j^\mu(p_1, p_a) j_\mu(p_2, p_b) V_H^{\mu\nu}(q_1, q_2)$$

Simple factorised expression, no approximations here!

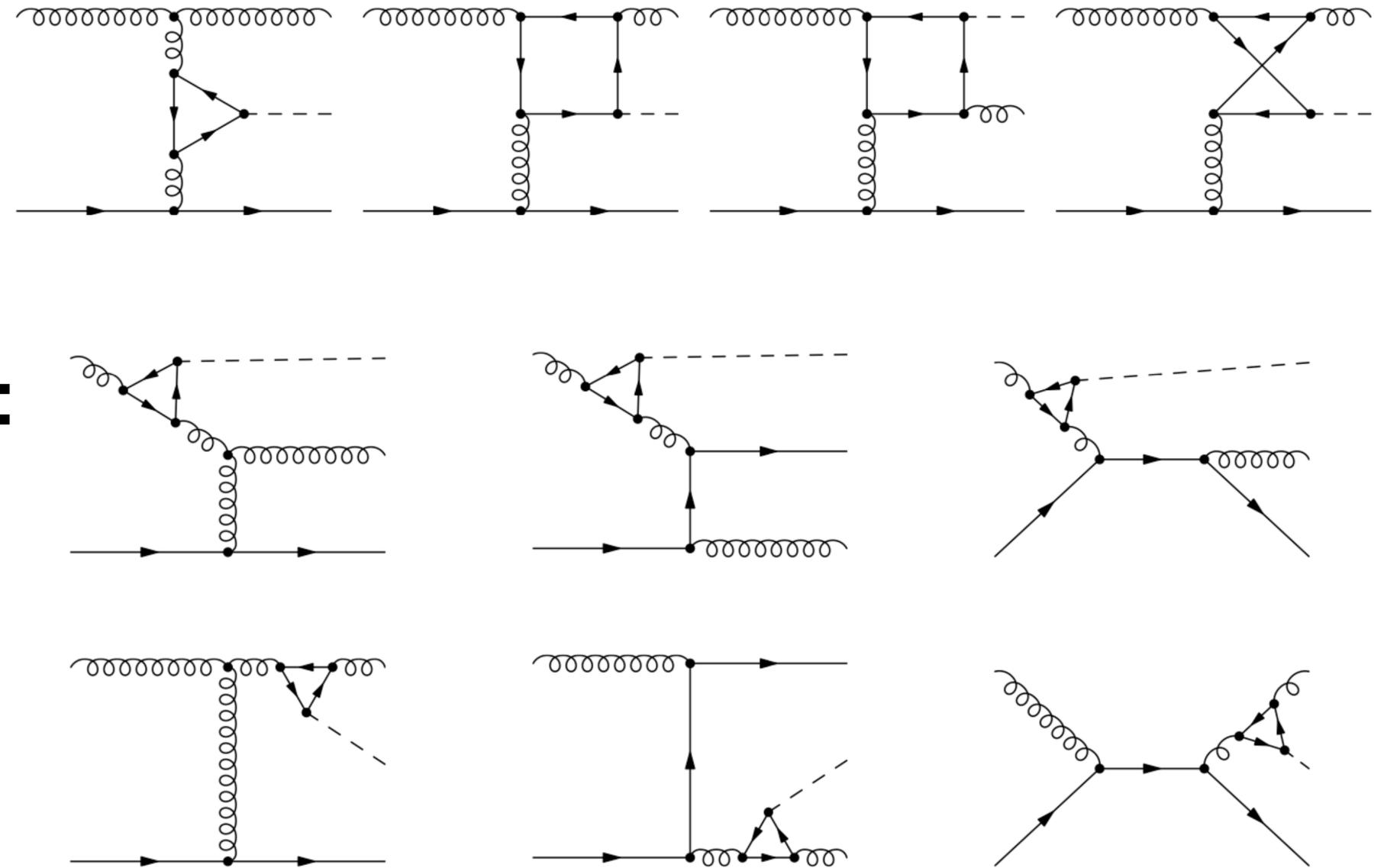
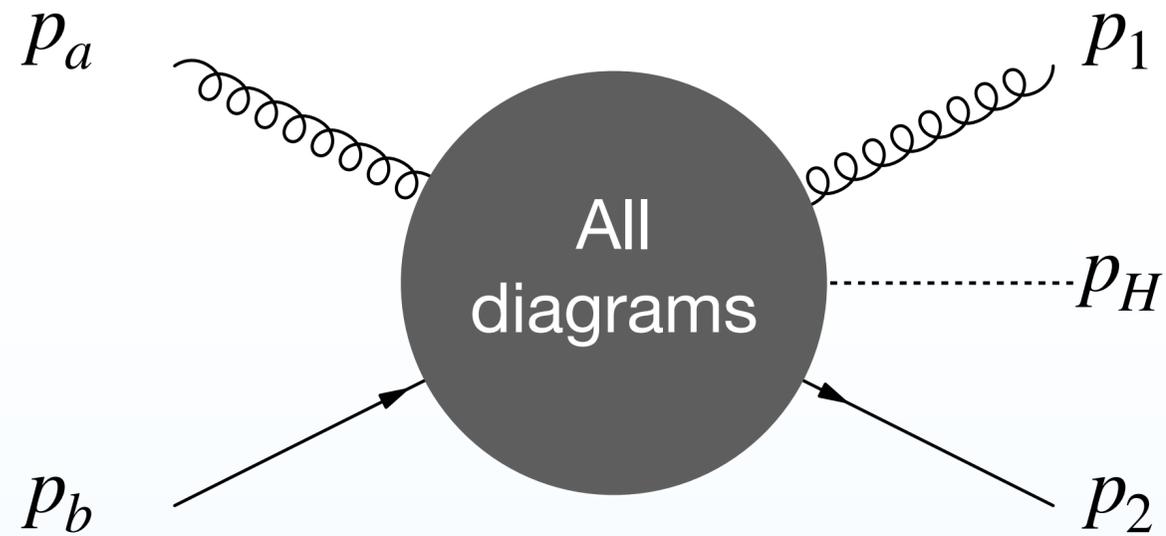


Central Higgs Production



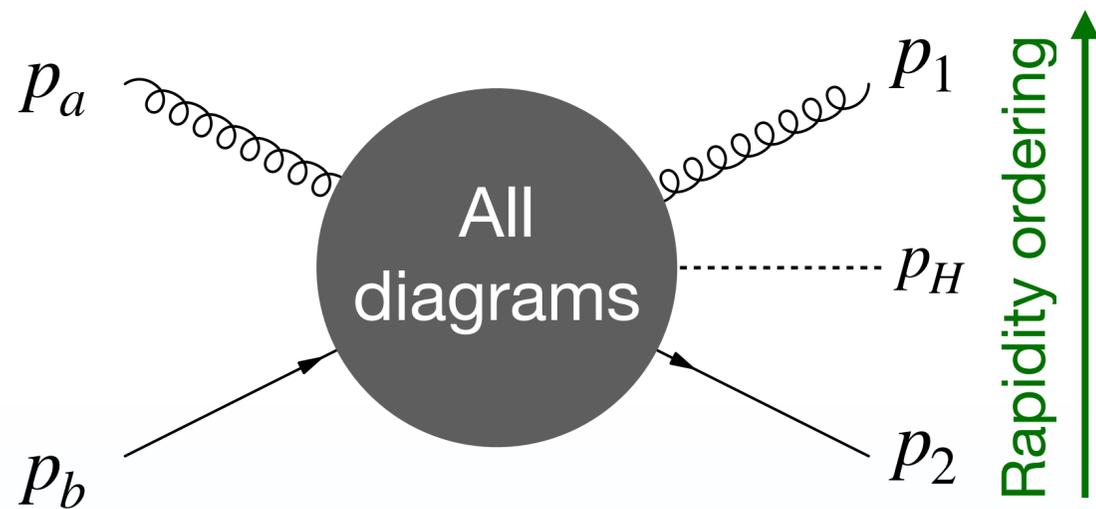


Central Higgs Production





Central Higgs Production



$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} j^\mu(p_1, p_a) j_\mu(p_2, p_b) V_H^{\mu\nu}(q_1, q_2)$$



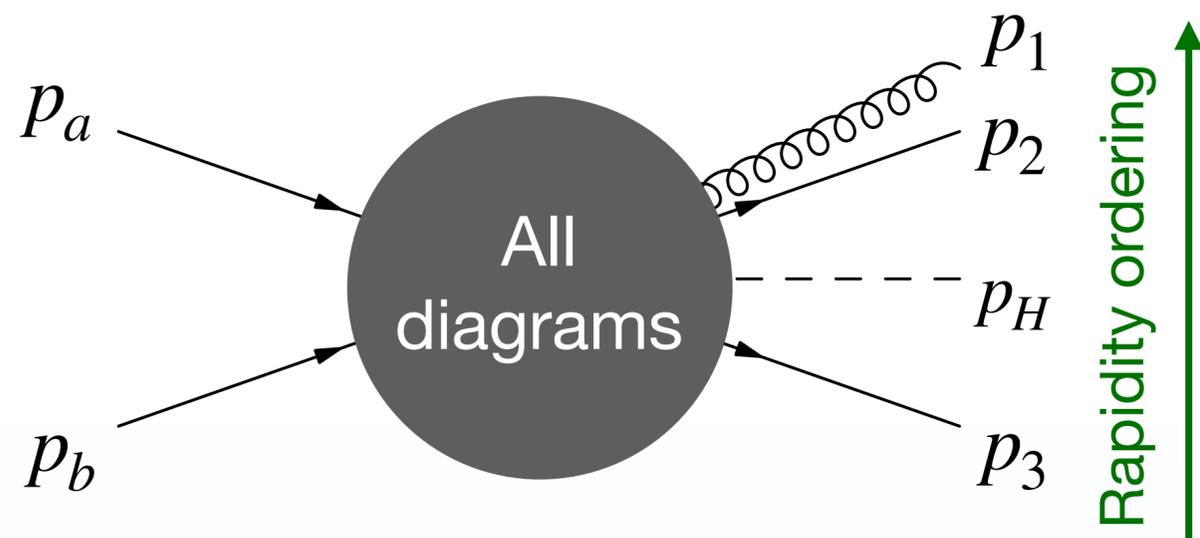
Central Higgs Production



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Unordered gluon emission



Set of gauge invariant NLL corrections

$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} j_\mu^{\text{uno}}(p_a, p_1, p_2) j^\mu(p_3, p_b) V_H^{\mu\nu}(q_1, q_2)$$



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Higgs plus dijet

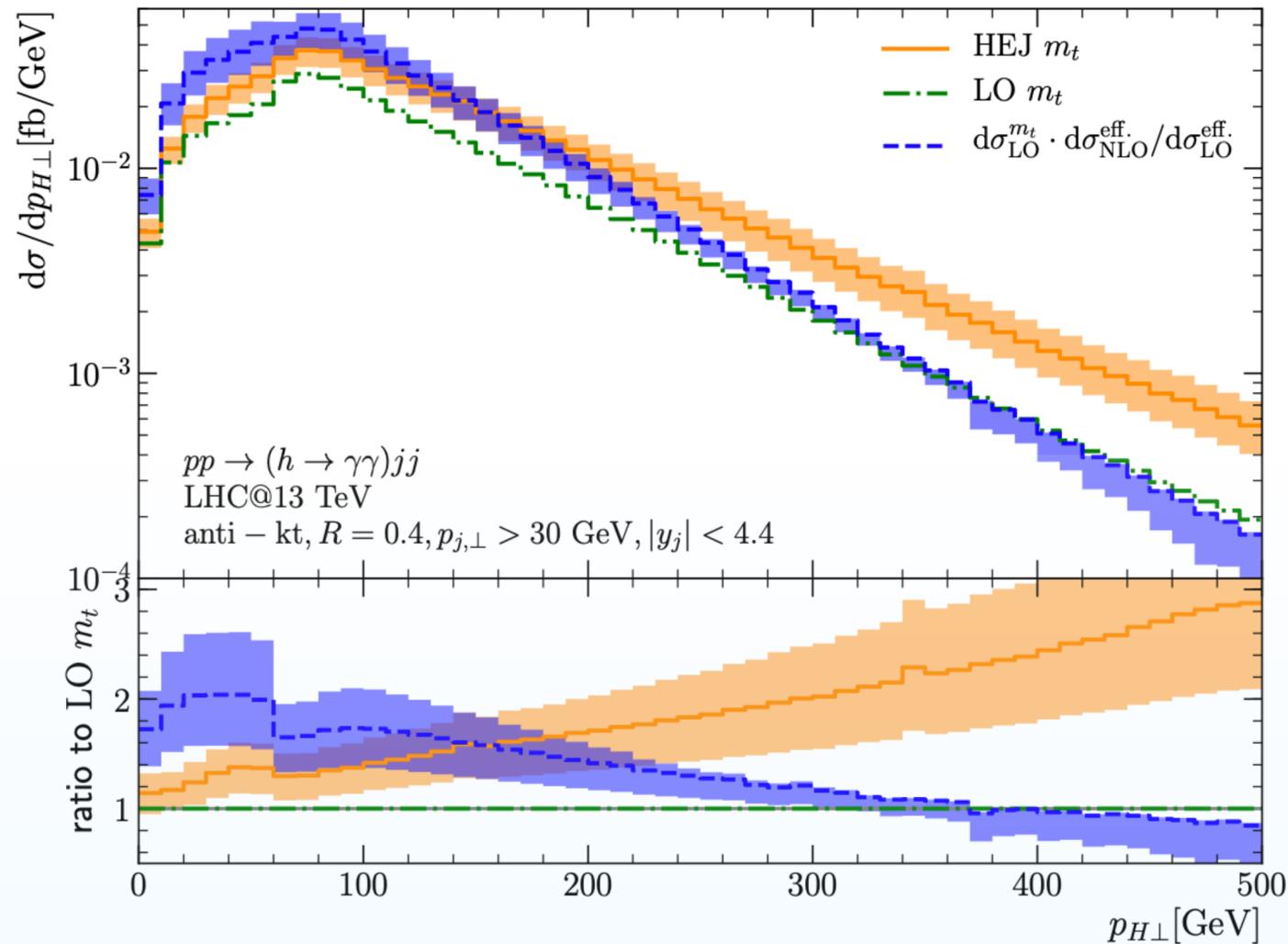


Adapted from ATLAS analysis [[1407.4222](#)]

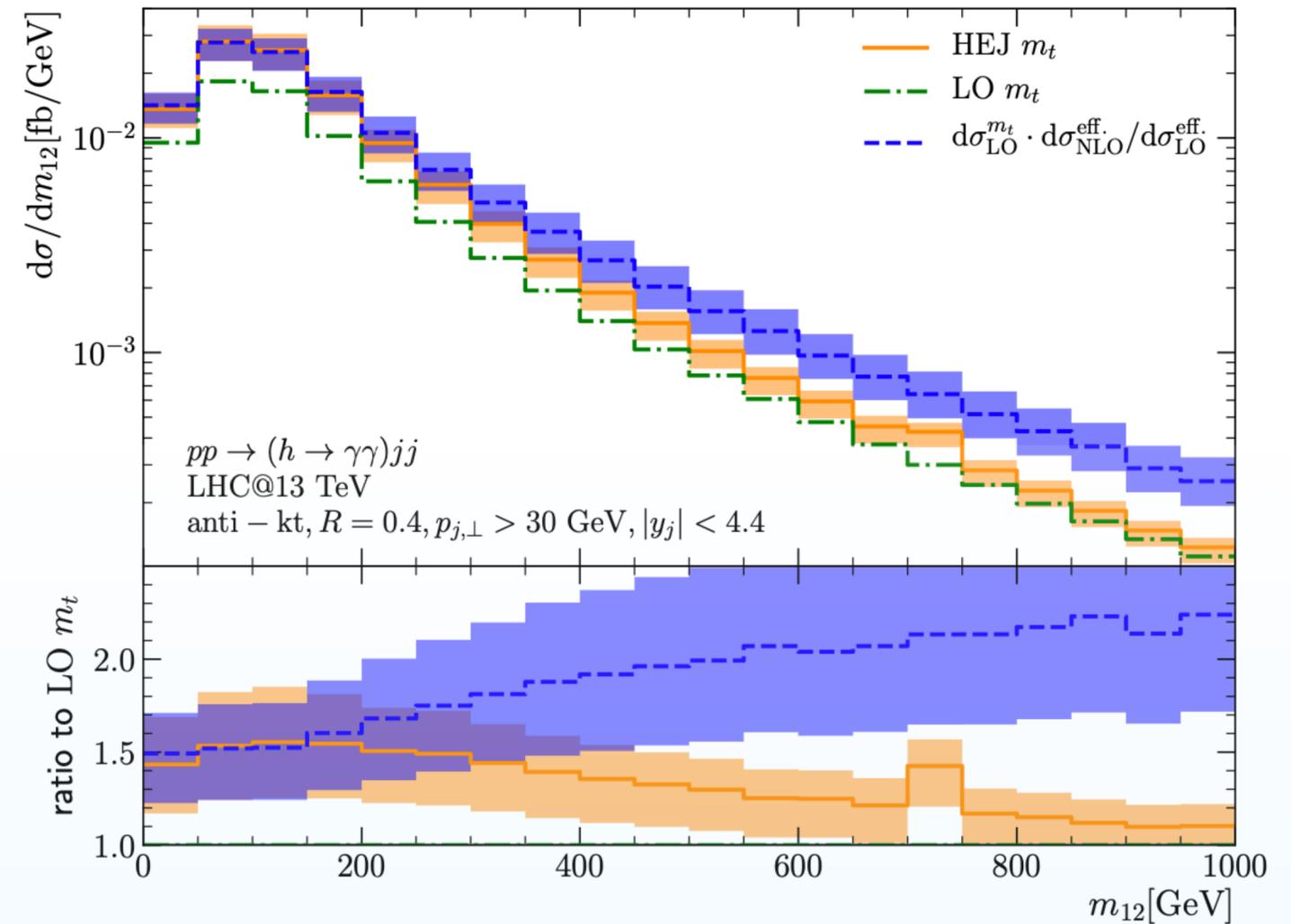
Baseline Jet Cuts	
Rapidity	$ y_j < 4.4$
Transverse momentum	$p_{\perp,j} > 30 \text{ GeV}$
Baseline Photon Cuts	
Rapidity	$ y_\gamma < 2.37$
Diphoton invariant mass	$105 \text{ GeV} < m_{\gamma_1\gamma_2} < 160 \text{ GeV}$
Transverse momentum hardest photon	$p_{\perp,\gamma_1} > 0.35 m_{\gamma_1\gamma_2}$
Transverse momentum other photon	$p_{\perp,\gamma_2} > 0.25 m_{\gamma_1\gamma_2}$
Vector Boson Fusion (VBF) Cuts	
Rapidity jet difference	$ y_{j_1} - y_{j_2} < 2.8$
Invariant dijet mass	$m_{j_1j_2} > 400 \text{ GeV}$



Resummation Effects



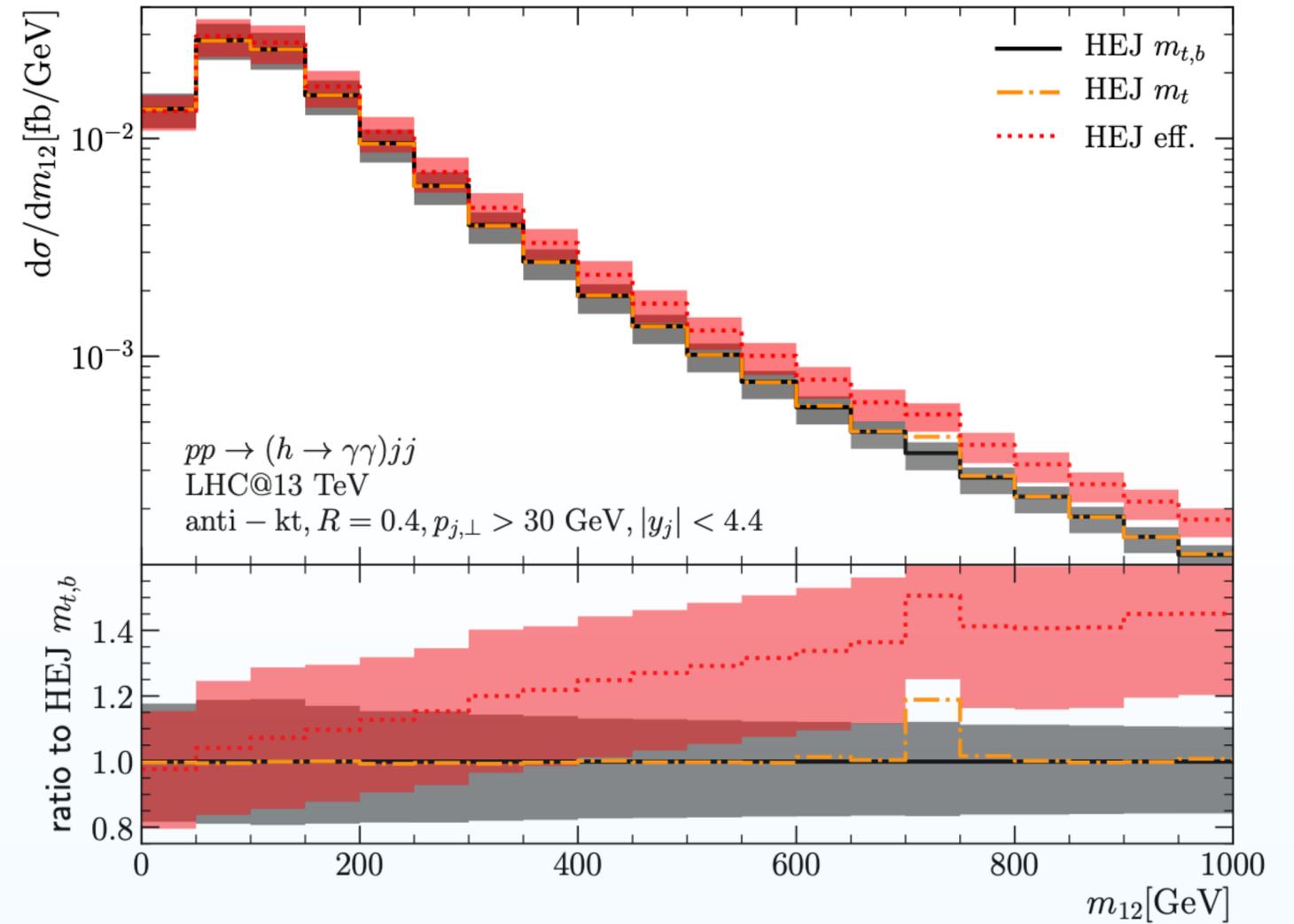
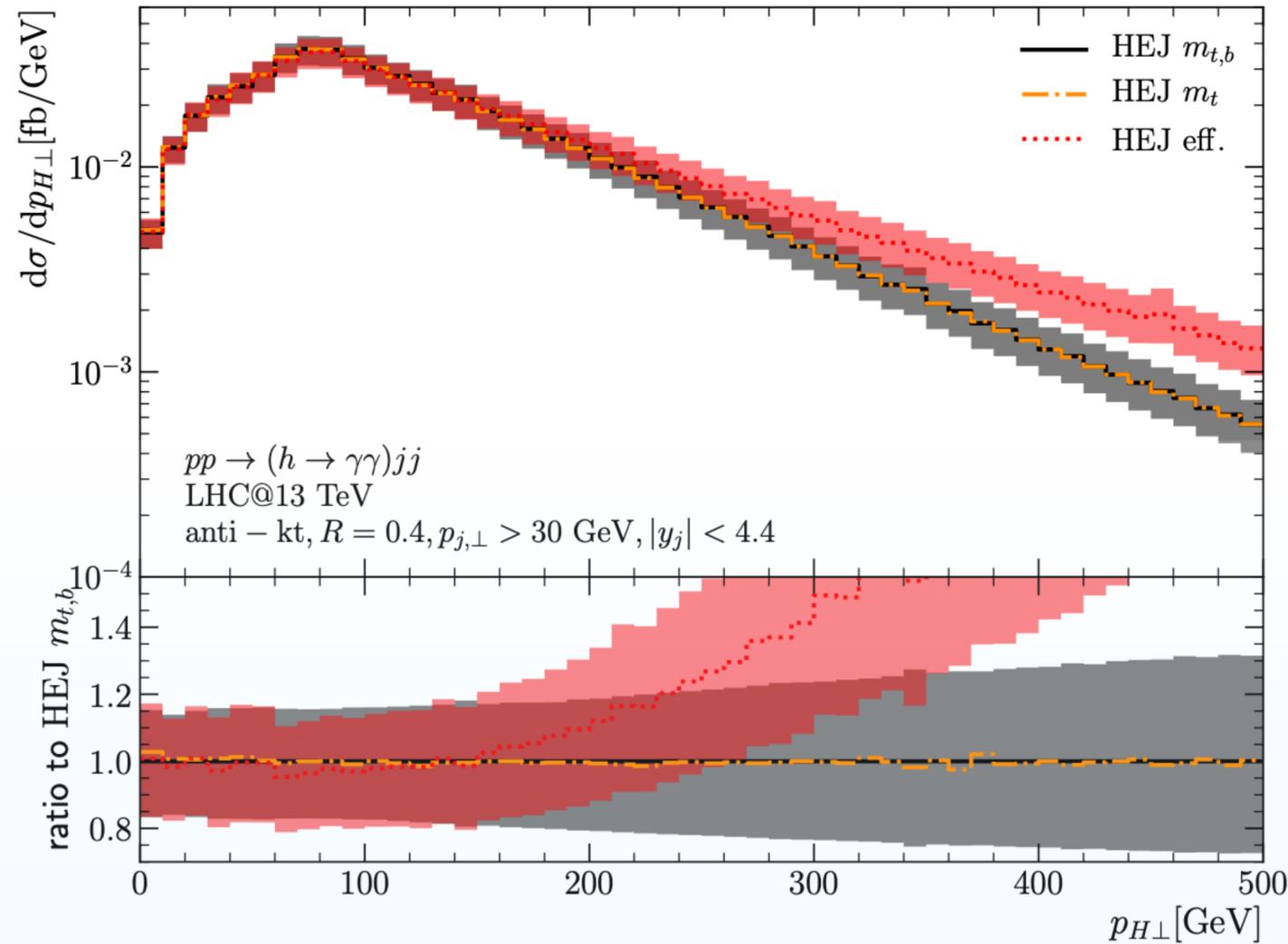
Transverse momentum sensitive to resummation
Resummation hardens the tail



Important observable for VBF cuts!
Resummation softens the tail



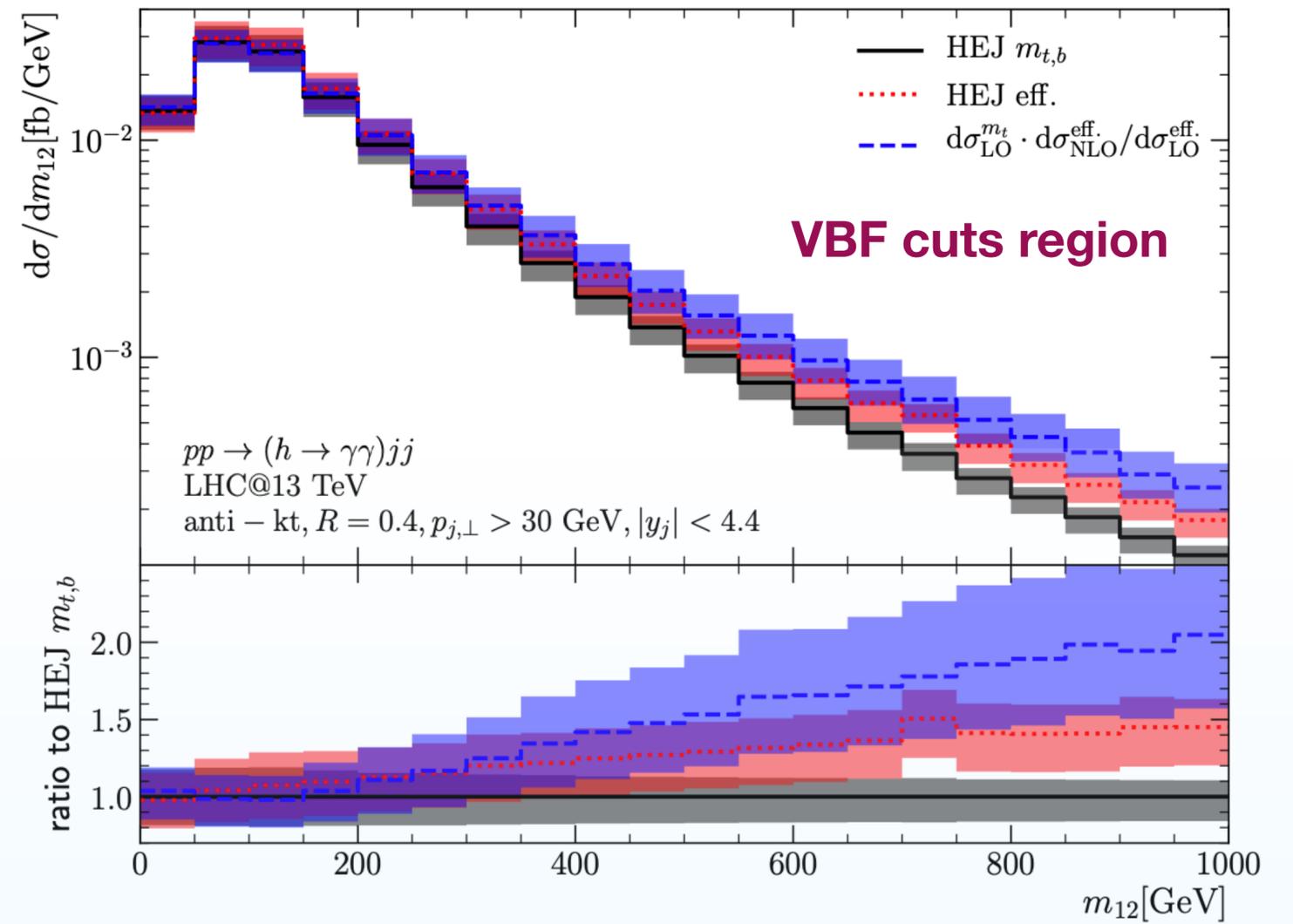
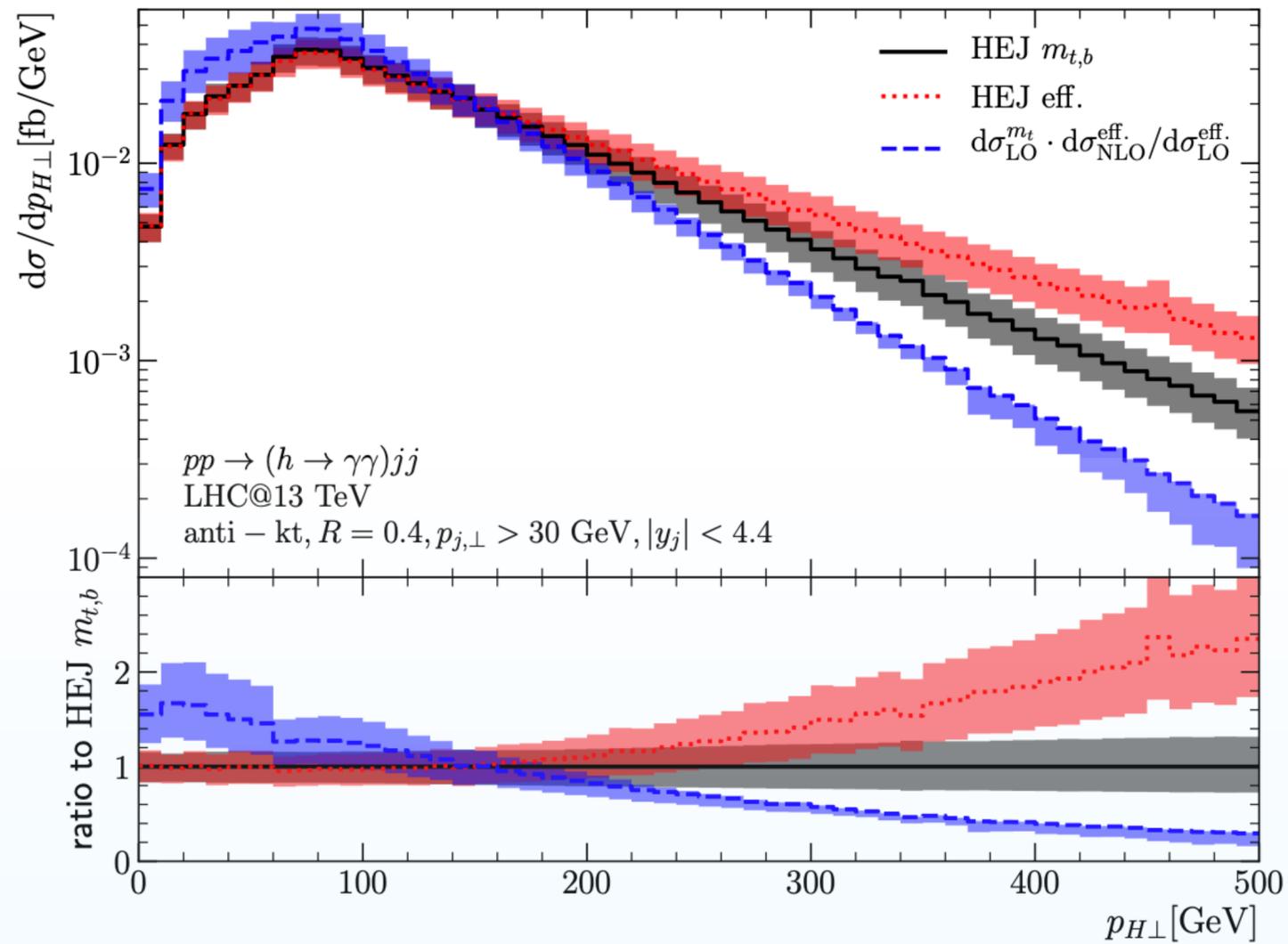
Finite Quark Masses Effects



Finite top mass effects is sizeable (bottom mass is not)

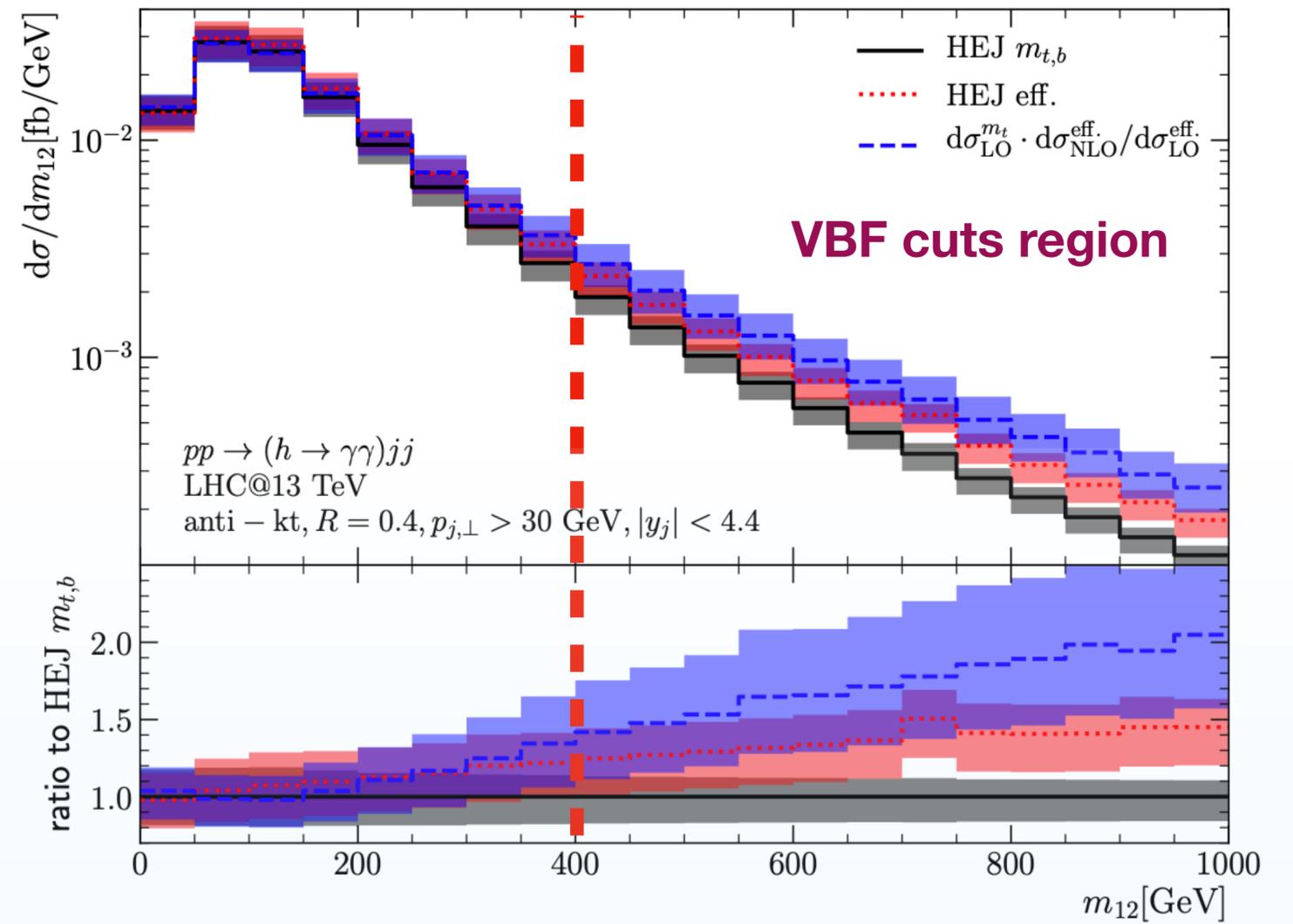
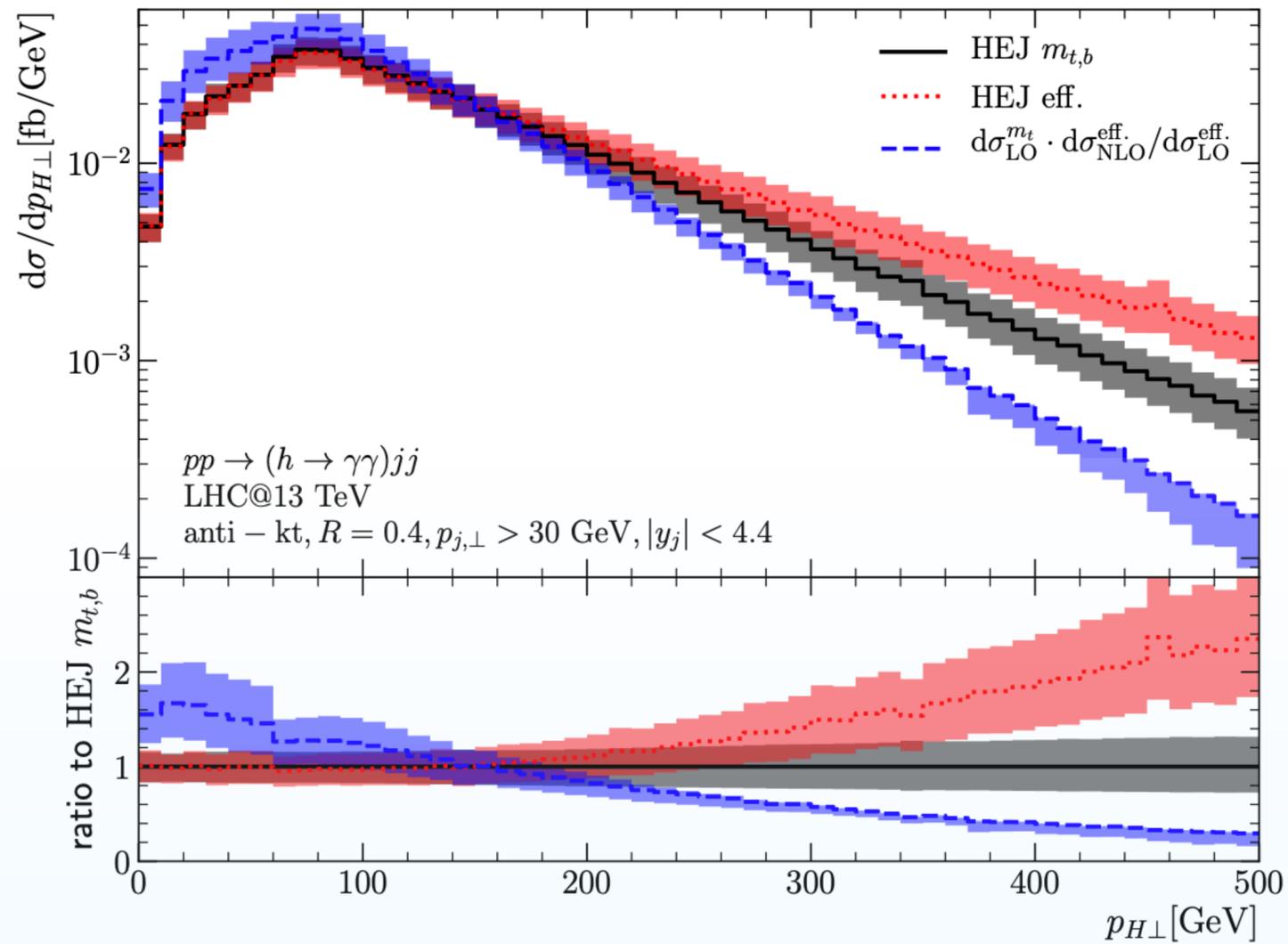


Combined effects



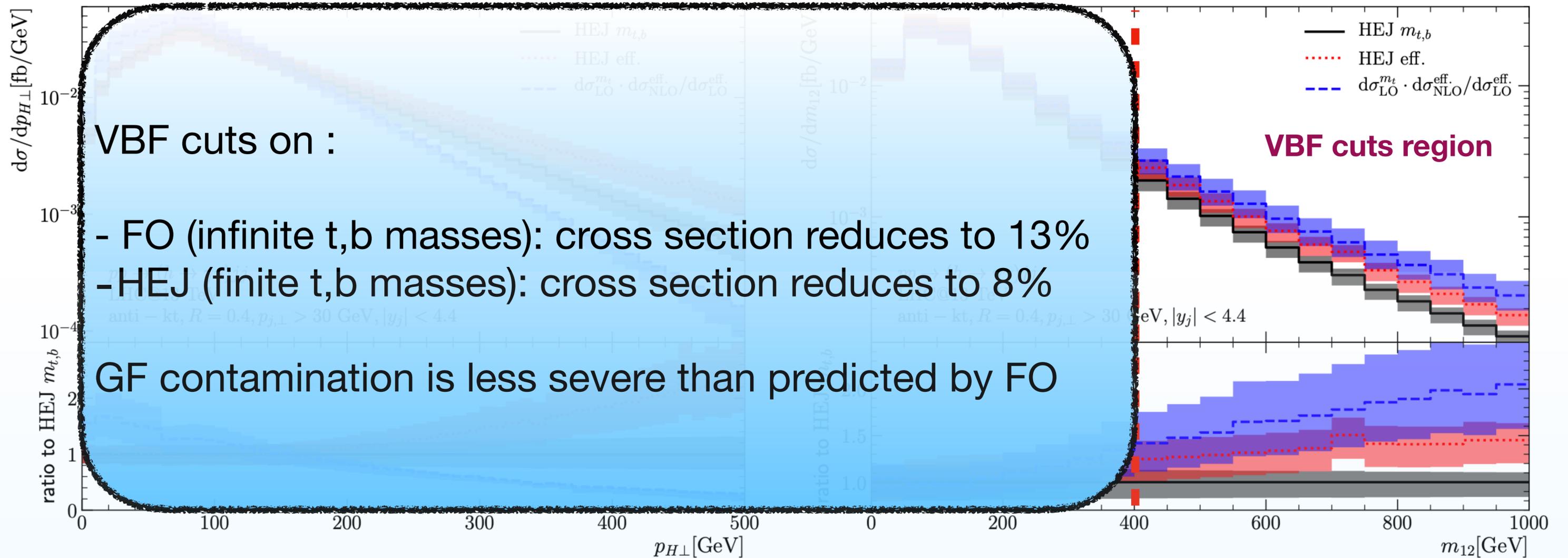


Combined effects

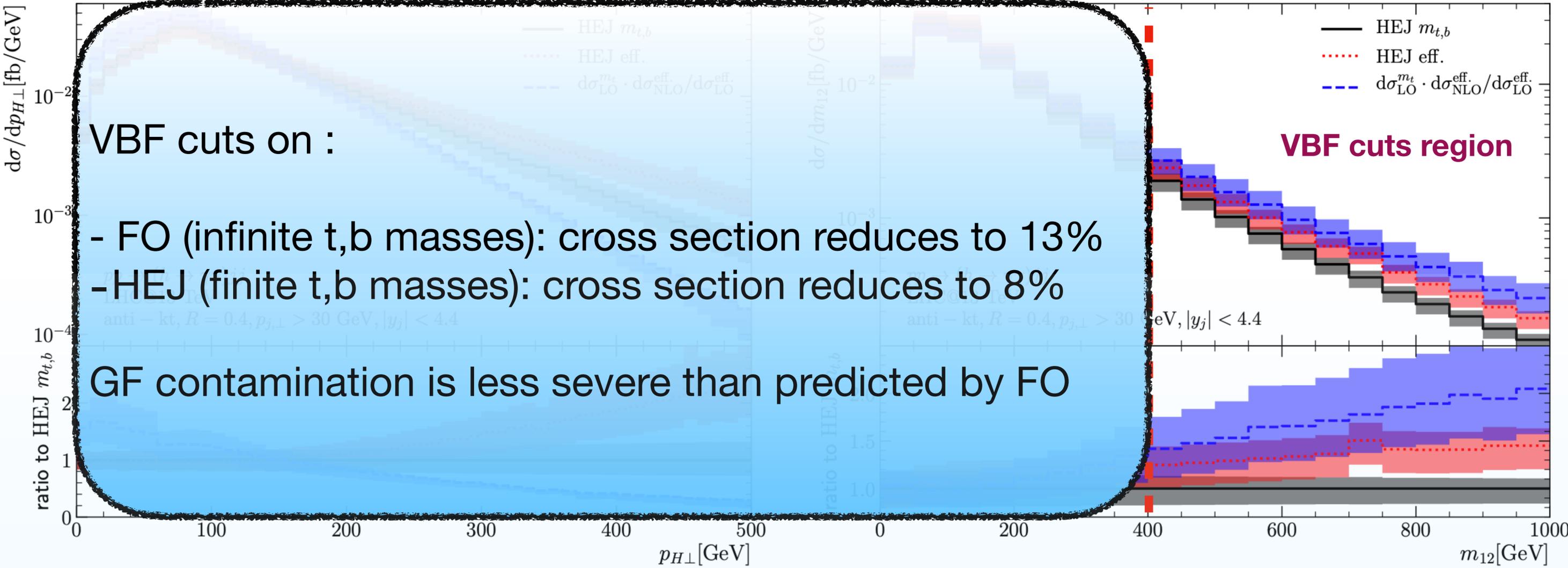




Combined effects



Combined effects



Effects add up: VBF cuts are more efficient than predicted by FO



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Higgs + dijet

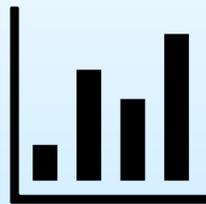




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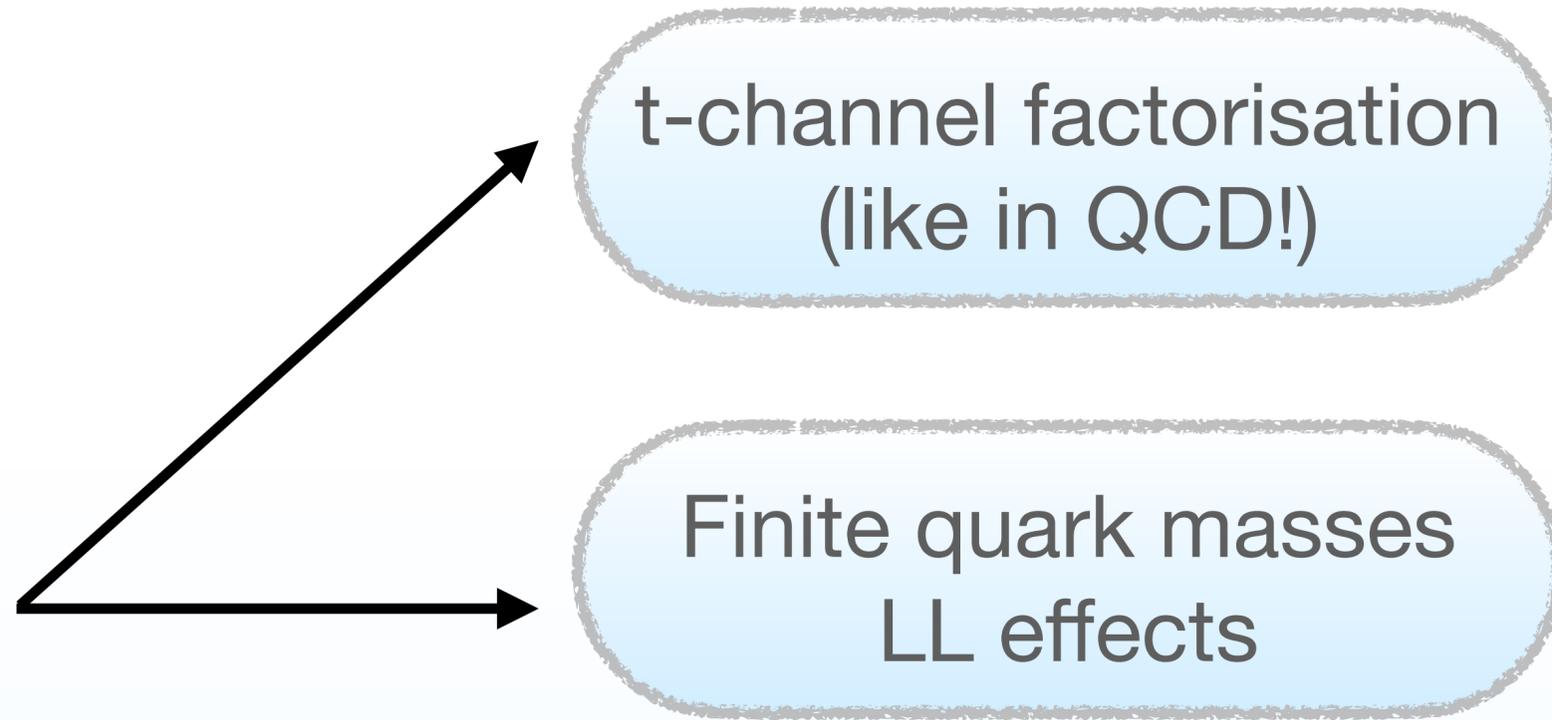
Higgs + dijet



t-channel factorisation
(like in QCD!)

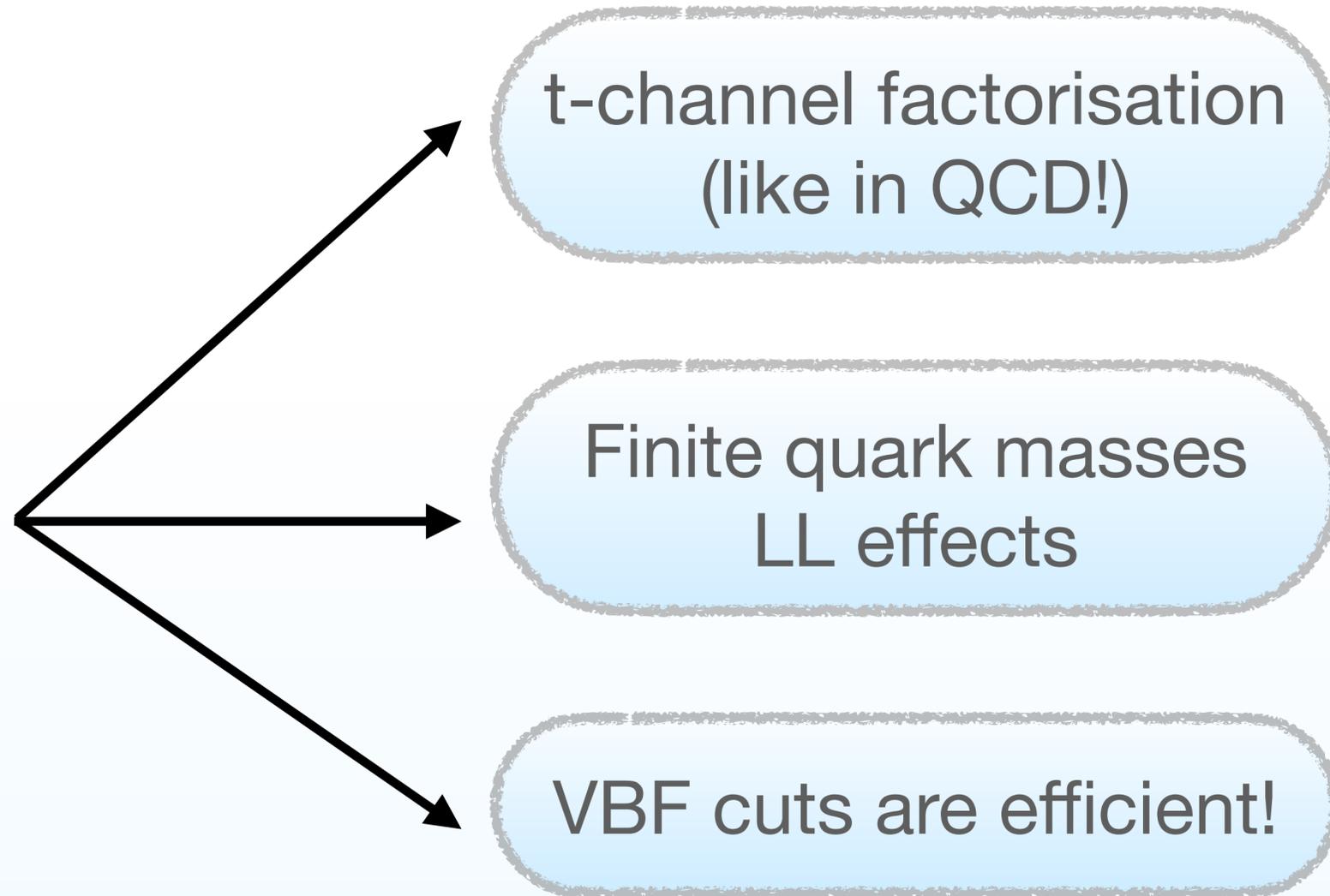


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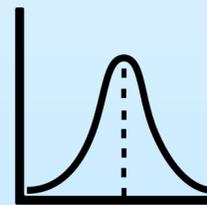
Higgs + one jet

Theory

Comparisons to
experimental data

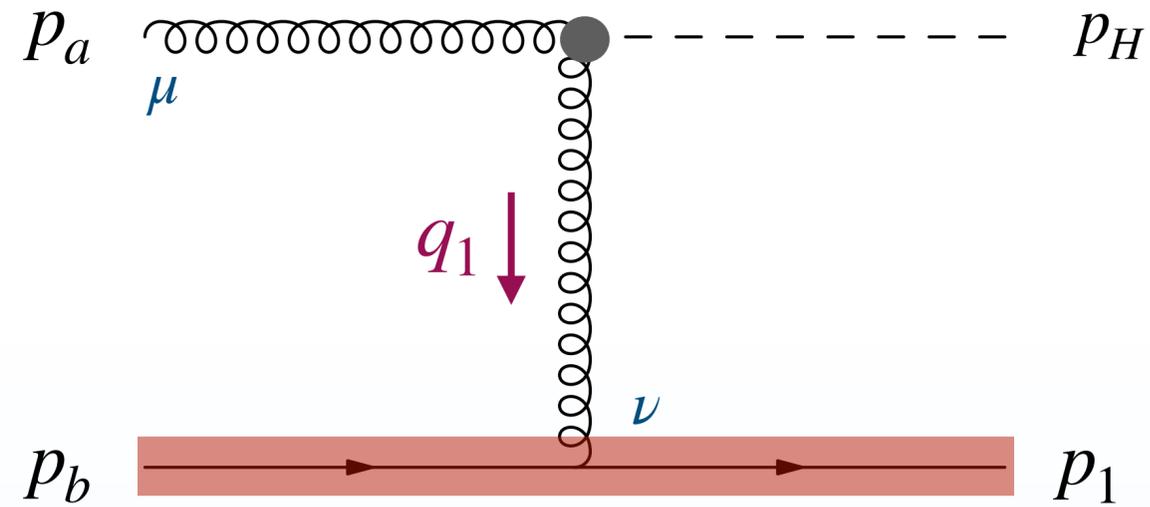
Comparisons to FO

[\[2210.10671\]](#)





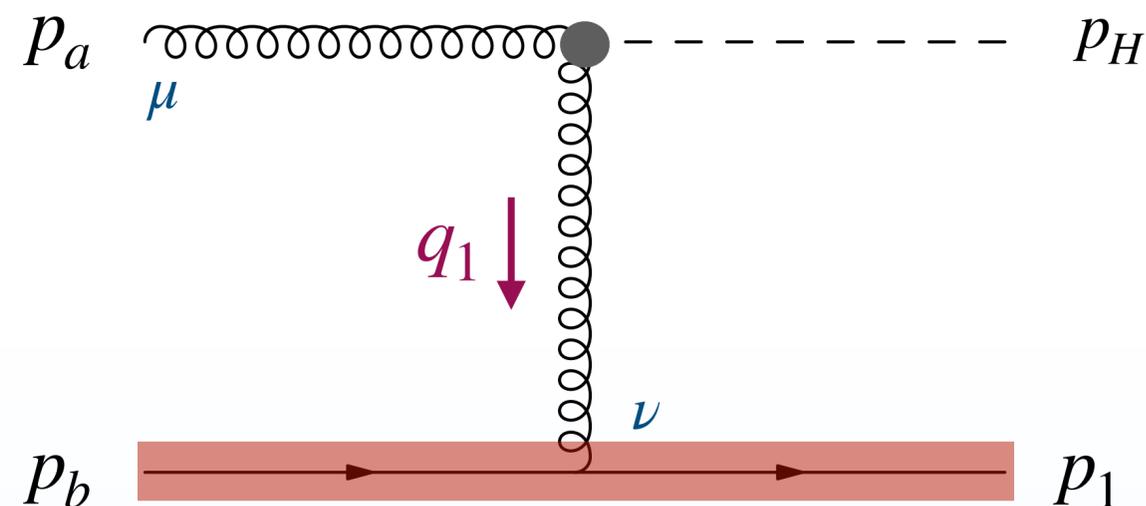
New Components for H+1j



$$\mathcal{M} \propto \frac{1}{q_1^2} \epsilon_\mu(p_a) j^\mu(p_1, p_b) V_H^{\mu\nu}(p_a, q_1)$$



New Components for H+1j



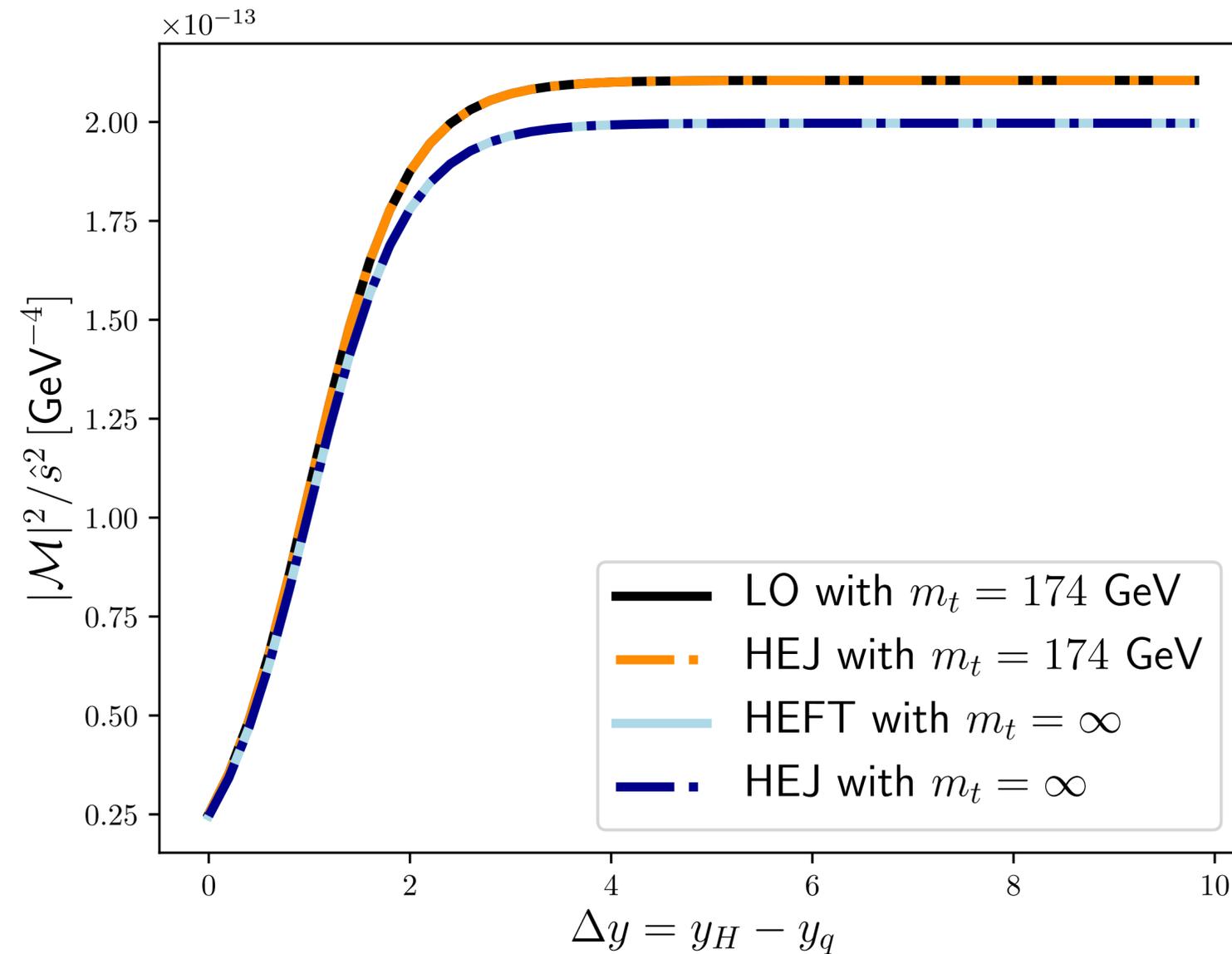
$$\mathcal{M} \propto \frac{1}{q_1^2} \epsilon_\mu(p_a) j^\mu(p_1, p_b) V_H^{\mu\nu}(p_a, q_1)$$

Simple factorised expression, no approximations here!

Before moving on, check Regge scaling

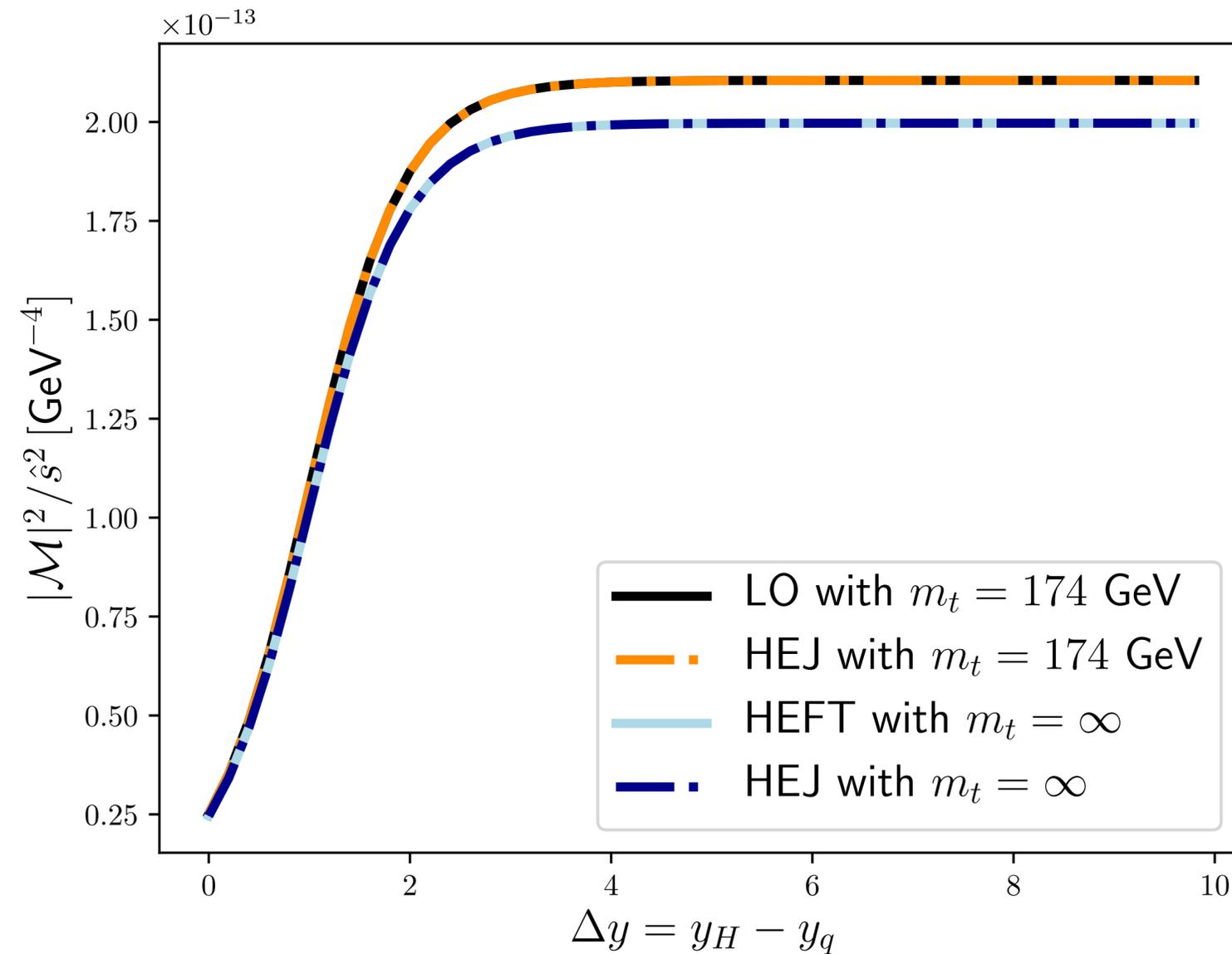


New Components for H+1j





New Components for H+1j



Indeed, no approximations,
exact LO description **with**
or **without** finite top mass

Regge scaling is verified:

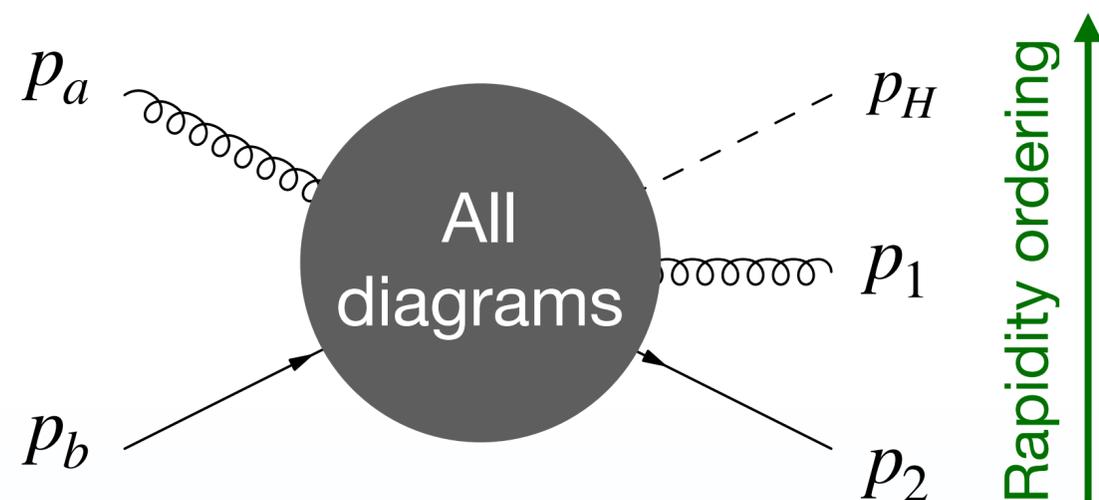
$$\mathcal{M} \propto s_{Hq}^1$$



Even though the Higgs is
not a coloured particle



New Components for H+1j



$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} \epsilon_\rho(p_a) j_\mu(p_2, p_b) V_H^{\rho\mu}(p_a, q_1) V^\lambda \epsilon_\lambda^*(p_1)$$



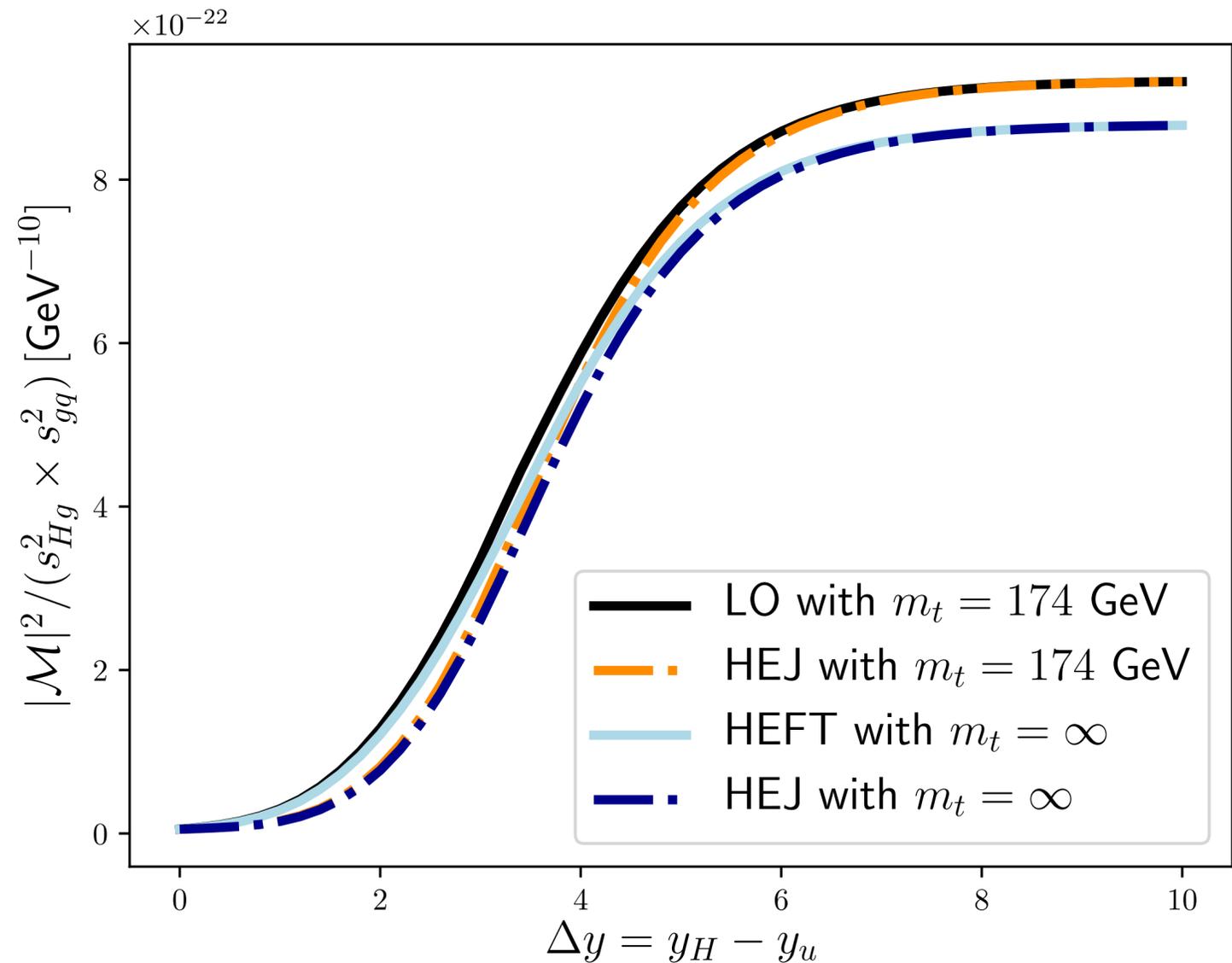
New Components for H+1j



$$\mathcal{M} \propto \frac{1}{q_1^2} \frac{1}{q_2^2} \epsilon_\rho(p_a) j_\mu(p_2, p_b) V_H^{\rho\mu}(p_a, q_1) V^\lambda \epsilon_\lambda^*(p_1)$$



New Components for H+1j



Approximations remain decent at low rapidity difference **with** or **without** finite top mass.

Regge theory is verified!

$$\mathcal{M} \propto s_{Hg}^1 \times s_{gu}^1$$





Higgs plus one jet results

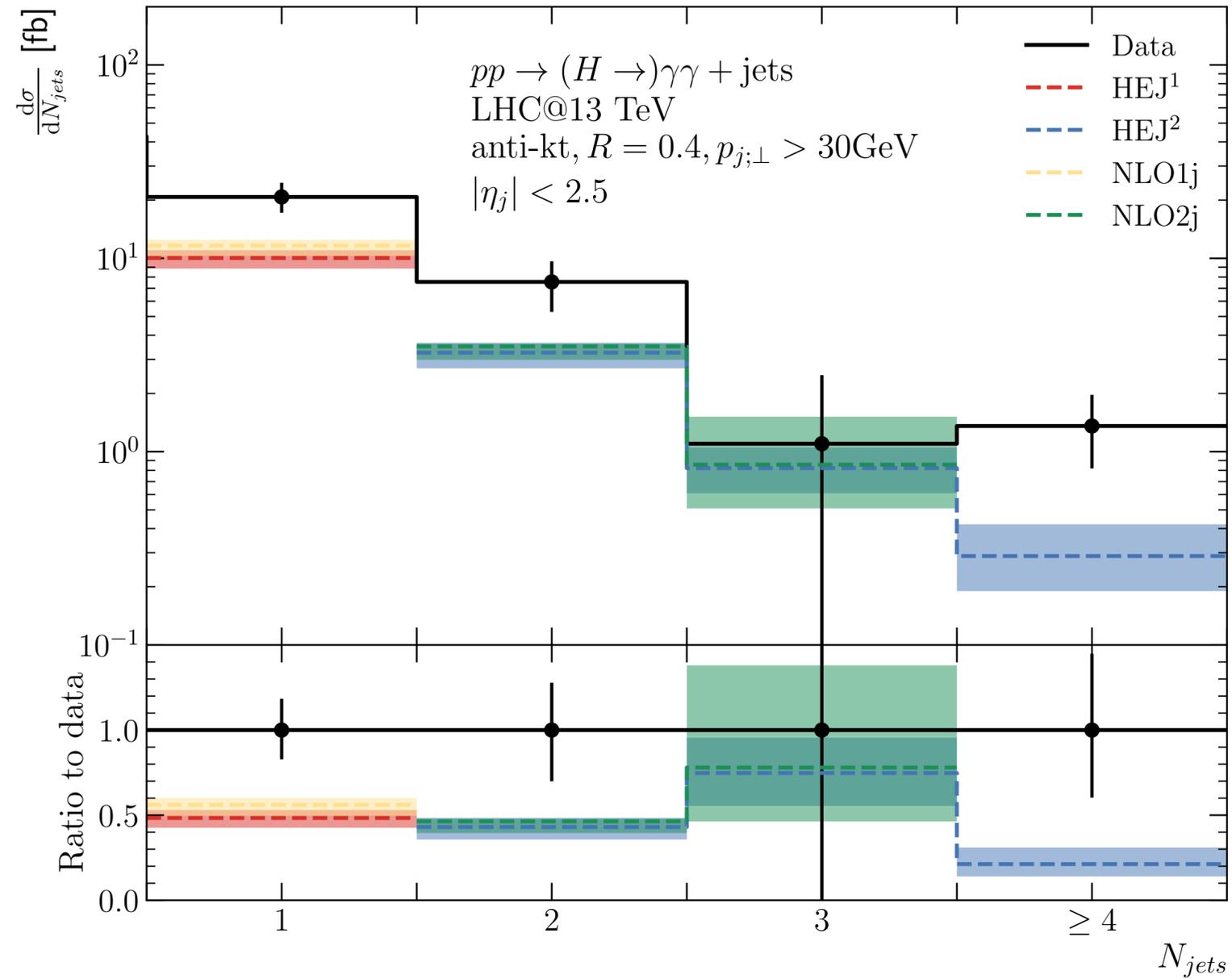


CMS analysis [[1807.03825](#), [2208.12279](#)]

Baseline Jet Cuts	
Pseudo-Rapidity	$ \eta_j < 4.7$
Transverse momentum	$p_{\perp,j} > 30 \text{ GeV}$
Baseline Photon Cuts	
Pseudo-Rapidity	$ \eta_\gamma < 2.5$
Diphoton invariant mass	$m_{\gamma_1\gamma_2} > 90 \text{ GeV}$
Transverse momentum hardest photon	$p_{\perp,\gamma_1} > \max(1/3 m_{\gamma_1\gamma_2}, 30 \text{ GeV})$
Transverse momentum other photon	$p_{\perp,\gamma_2} > 0.25 m_{\gamma_1\gamma_2}$



Higgs plus one jet results



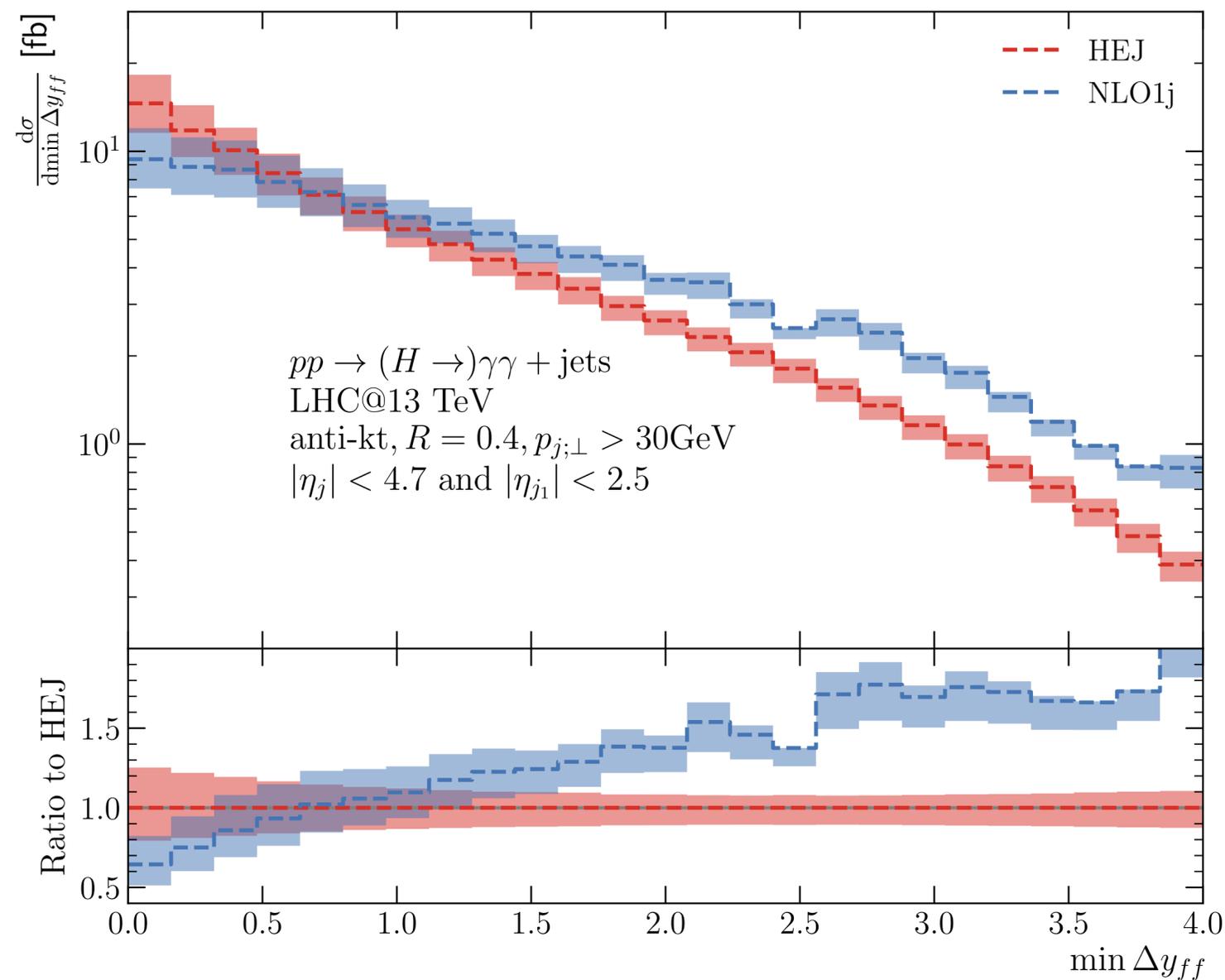
Resummed HEJ with finite m_t matched to NLO 1j
Resummed HEJ with finite m_t matched to NLO 2j
NLO 1j with infinite m_t
NLO 2j with infinite m_t

Resummed predictions for 4 jets and more

Predictions undershoot the data, but adding electroweak processes bring the predictions within the uncertainty of the data



Higgs plus one jet results



Resummed HEJ with finite m_t matched to NLO 1j
NLO 1j with infinite m_t

NLO effects harden the tail of large dijet rapidity separation

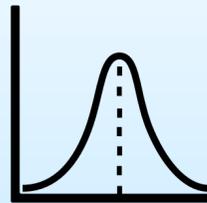
HEJ resummation soften the tail: the logarithms are numerically significant!



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Higgs + one jet

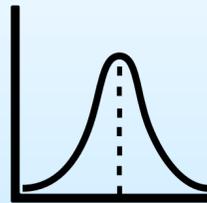




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Higgs + one jet



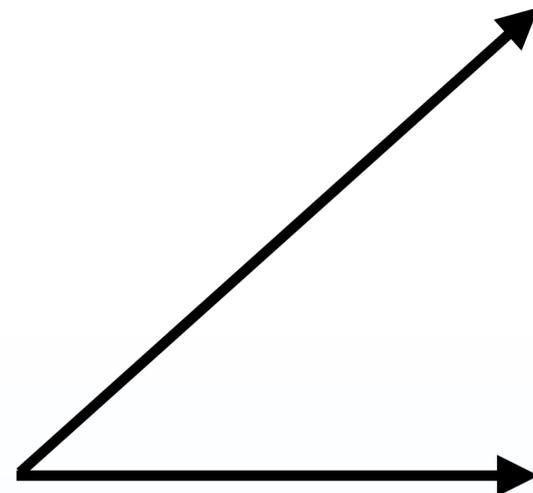
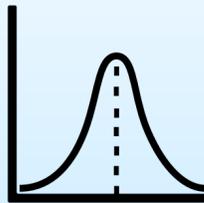
t-channel factorisation
still valid



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Higgs + one jet



t-channel factorisation
still valid

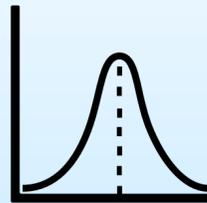
Finite quark masses
still important



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Higgs + one jet



t-channel factorisation
still valid

Finite quark masses
still important

logs are significant in
tails of distributions

Conclusion



HEJ publicly available on [here](#)

Resums High-Energy large logarithms to LL accuracy, work is ongoing towards NLL accuracy.



Conclusion



HEJ publicly available on [here](#)

Resums High-Energy large logarithms to LL accuracy, work is ongoing towards NLL accuracy.

H+1j inclusive: new paper!

Takes into account finite quark masses unlike FO approaches, more observables are compared in paper



Conclusion



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Resums High-Energy large logarithms to LL accuracy, work is ongoing towards NLL accuracy.

H+2j inclusive

VBF cuts are more efficient than what FO predicts, finite quark masses effects and High-Energy Logarithms work together

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H+1j inclusive: new paper!

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Ongoing work

Merging with parton shower, reaching NLL accuracy, amplitudes for Vector Boson production



Thank you for listening



Any questions?

