

# Higher Order Corrections in Pure Spinor SYM

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1 Introduction

2 Non-Linear Super Yang-Mills

3 Ansatz Building

4  $\alpha'^3$  Corrections

- Pure Spinor Superstring OPEs inspired a recursive approach [1012.3987]  
 $\lambda^\alpha$  where  $\lambda \gamma^m \lambda = 0$ .

# Non Linear Fields

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- Inspired Non-linear fields that generate Multiparticle fields  
 $A_\alpha(x, \theta)$ ,  $A^m(x, \theta)$ ,  $F^{mn}(x, \theta)$ ,  $W^\alpha(x, \theta)$

# Non Linear Fields

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- Inspired Non-linear fields that generate Multiparticle fields  
 $A_\alpha(x, \theta)$ ,  $A^m(x, \theta)$ ,  $F^{mn}(x, \theta)$ ,  $W^\alpha(x, \theta)$
- Capture the  $N$ -point expansion of the field and Kinematic factors!

$$\text{Tree-Level} \rightarrow \frac{1}{3} \text{Tr} \langle V V V \rangle$$

← [1510.08843]

$$\langle \lambda^3 \theta^5 \rangle = 2880, \quad V = \lambda^\alpha A_\alpha$$

# Higher Order Corrections & Methods

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- BRST operator in PS = Simple!

$$Q = \chi^\alpha D_\alpha \quad [\text{hep-th/0001035}]$$

Spinor Derivative

$$D_\alpha \equiv \partial_\alpha + \frac{1}{2} (\gamma^m \theta)_\alpha \partial_m$$

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- Find BRST Closed (NOT EXACT) expression at Correct Mass dimension

↳ Obtain the Perturbative expansion



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# Non-Linear Equation of Motion

- Non-Linear SYM is built from the Spinor and Vector Covariant derivatives,

$$\nabla_\alpha = D_\alpha - A_\alpha, \quad \nabla_m = \partial_m - A_m$$

- We impose the following constraint,

$$\{ \nabla_\alpha, \nabla_\beta \} = \gamma_{\alpha\beta}^m \nabla_m$$

[Witten '86]

# Non-Linear Equation of Motion

- Using the Bianchi identity and the above constraint yields the equations of motion:

$$\{D_{(\alpha}, \mathbb{A}_{\beta)}\} = \gamma_{\alpha\beta}^m \mathbb{A}_m + \{\mathbb{A}_\alpha, \mathbb{A}_\beta\},$$

$$\{D_\alpha, \mathbb{W}^\beta\} = \frac{1}{4} (\gamma^{mn})_\alpha{}^\beta \mathbb{F}_{mn} + \{\mathbb{A}_\alpha, \mathbb{W}^\beta\},$$

$$[D_\alpha, \mathbb{A}_m] = [\partial_m, \mathbb{A}_\alpha] + (\gamma_m \mathbb{W})_\alpha + [\mathbb{A}_\alpha, \mathbb{A}_m],$$

$$[D_\alpha, \mathbb{F}^{mn}] = (\mathbb{W}^{[m} \gamma^{n]})_\alpha + [\mathbb{A}_\alpha, \mathbb{F}^{mn}],$$

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$$[D_{\alpha}, \mathbb{A}_m] = [\partial_m, \mathbb{A}_{\alpha}] + (\gamma_m \mathbb{W})_{\alpha} + [\mathbb{A}_{\alpha}, \mathbb{A}_m],$$

$$[D_{\alpha}, \mathbb{F}^{mn}] = (\mathbb{W}^{[m} \gamma^{n]})_{\alpha} + [\mathbb{A}_{\alpha}, \mathbb{F}^{mn}],$$

We can use these for

BRST variations  $\rightarrow [Q, \mathbb{A}_m] = [\partial_m, \lambda \mathbb{A}] + (\lambda \gamma_m \mathbb{W}) + [\lambda \mathbb{A}, \mathbb{A}_m].$

# Non-Linear Equation of Motion

- Dirac and Yang-Mills :

$$\gamma^m_{\alpha\beta} [\nabla_m, W^\beta] = 0, \quad \partial_m F^{mn} = \gamma^n_{\alpha\beta} \{W^\alpha, W^\beta\}$$

# Non-Linear Equation of Motion

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$$\gamma^m_{\alpha\beta} [\nabla_m, W^\beta] = 0, \quad \partial_m F^{mn} = \gamma^n_{\alpha\beta} \{W^\alpha, W^\beta\}$$

- We can treat **higher mass** fields like  $[\nabla_m, W^\alpha]$  as 'fundamental',

$$K^{m_1 \dots m_k | N} = [\nabla^{m_1}, K^{m_2 \dots m_k | N}]$$

↳ They also obey 'Dirac + Yang-Mills' eqns. [2210.1424].

Contains Deconcatenation.

PS Review

# Generating Series

- Non-linear fields can be expanded as a Lie algebra valued generating series with BG coefficients.

[1501.05562]. 
$$K = \sum_{p=1}^{\infty} \sum_{i_1 \dots i_p} K_{i_1 \dots i_p} t^{i_1} \dots t^{i_p}$$

BG Currents

Multiparticle Labels.

Lie algebra generators

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- Non-linear fields can be expanded as a Lie algebra valued generating series with BG coefficients.

$$K = \sum_{p=1}^{\infty} \sum_{i_1 \dots i_p} K_{i_1 \dots i_p} t^{i_1} \dots t^{i_p}$$

- Multiparticle BG currents have correct pole structure,

$$K_p = \frac{1}{s_p} \sum_{xy=p} \underbrace{K_{[x,y]}}_{\text{From BG EoM}}$$

$p = i_1 i_2 \dots i_p$

[1510.08843]



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# Dynkin Labels of $\mathfrak{so}(10)$

○ Consider all terms of correct mass dimension

↳ At  $\alpha^3$  this is dimension 7.5.

$$\nabla_m \leftrightarrow 1$$

$$A_\alpha \leftrightarrow 0.5$$

$$A_m \leftrightarrow 1$$

$$F_{mn} \leftrightarrow 2$$

$$W^\alpha \leftrightarrow 1.5$$

## Dynkin Labels of $\mathfrak{so}(10)$

- Consider all terms of correct mass dimension
  - ↳ At  $\alpha^3$  this is dimension 7.5.
- Use Dynkin labels to determine if in  $SO(10)$  group structure.
  - Also check Constraints + Scalars

## Dynkin Labels of $\mathfrak{so}(10)$

- Consider all terms of correct mass dimension  
↳ At  $\alpha'^3$  this is dimension 7.5.
- Use Dynkin labels to determine if in  $SO(10)$  group structure.
- Find all allowed permutations + Contractions  
→ This gives the ansatz.

# Dynkin Labels of $\mathfrak{so}(10)$

- Basis labels :

$$\text{Scalar } A = (0, 0, 0, 0, 0),$$

$$\text{Vector } A^m = (1, 0, 0, 0, 0),$$

$$\text{2-Form } A^{[mn]} = (0, 1, 0, 0, 0),$$

$$\text{3-Form } A^{[mnp]} = (0, 0, 1, 0, 0),$$

$$\text{Anti-Weyl } \xi_\alpha = (0, 0, 0, 1, 0),$$

$$\text{Weyl } \xi^\alpha = (0, 0, 0, 0, 1),$$

[hep-th/0205165]

# Ansatz Terms

- Example :  $\mathbb{V}(\lambda \gamma \mathbb{W})(\lambda \gamma \mathbb{W}) \mathbb{D}^2 \mathbb{F}$

- Decomposition  $\Rightarrow 1(0, 0, 0, 0, 0) + \dots$

- Only a single scalar :

$$\mathbb{V}(\lambda \gamma^m \mathbb{W})(\lambda \gamma^p \mathbb{W}) \mathbb{F}^{m \times n \times p \times n}$$

# Ansatz Terms

- Example :  $\mathbb{V}(\lambda \gamma W)(\lambda \gamma W) D^2 F$ 
  - Decomposition  $\Rightarrow 1(0,0,0,0,0) + \dots$
  - Only a single scalar :  
 $\mathbb{V}(\lambda \gamma^m W)(\lambda \gamma^p W) F^{mnlpn}$

○ Repeating this for all terms we can think of gives 162 terms  $\rightarrow$  After PS + EoM.

# Higher Mass Identities

- Higher order terms can often be re-written as lower order



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- Higher order terms can often be re-written as lower order
- E.g.  $D^N F$  and  $D^N W$ ,  $N \geq 2$  can be decomposed into sym. Traceless part + other fields.

$$F^{ab|mn} = \hat{F}^{(ab)|mn} + \frac{\delta^{ab}}{5} ([F^{mp}, F^{pn}] + \{(W^{[m} \gamma^{n]}, W)\}) + \frac{1}{2} [F^{mn}, F^{ab}]$$

 sym. + Traceless in a, b.

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- After applying BRST Closure  $\rightarrow$  Fix last unknown  
 $\rightarrow$  With 4-point amp.

BRST fixes the Correction  
up to a Mandelstam invariant  
 $\rightarrow$  We fix this

- After applying BRST Closure  $\rightarrow$  Fix last unknown with 4-point amp.
- The Final result gives

$$\begin{aligned}
 & \mathbb{V}(\mathbb{F}^{mn}(\lambda\gamma^n\mathbb{W}^p)(\lambda\gamma^p\mathbb{W}^m) + (\lambda\gamma^m\mathbb{W}^n)\mathbb{F}^{np}(\lambda\gamma^p\mathbb{W}^m)) \rightarrow \text{2-Loop 4-point Kin. Factor.} \\
 & \left[ \begin{aligned}
 & + \frac{1}{2} [\mathbb{F}^{mn}\mathbb{F}^{np}, \{(\lambda\gamma^m\mathbb{W}), (\lambda\gamma^p\mathbb{W})\}] \\
 & + \frac{1}{4} \mathbb{F}^{mn}\mathbb{F}^{pq} \{(\lambda\gamma^m\mathbb{W}), (\lambda\gamma^{npq}\mathbb{W})\} \\
 & + \frac{1}{4} \mathbb{F}^{mn}(\lambda\gamma^{npq}\mathbb{W}) [(\lambda\gamma^m\mathbb{W}), \mathbb{F}^{pq}] \\
 & - \frac{1}{4} \mathbb{F}^{pq}(\lambda\gamma^m\mathbb{W}) [(\lambda\gamma^{npq}\mathbb{W}), \mathbb{F}^{mn}] \\
 & - \frac{1}{4} (\lambda\gamma^{npq}\mathbb{W})(\lambda\gamma^m\mathbb{W}) [\mathbb{F}^{mn}, \mathbb{F}^{pq}]
 \end{aligned} \right] \leftarrow \begin{array}{l} \text{'Stringy'} \\ \text{SYM Corrections} \end{array}
 \end{aligned}$$

- Connection between  $\mathbb{W}$  and NO- $\mathbb{W}$  Corrections,

$$\mathbb{F}^{m|p|q} (\lambda \gamma^p \mathbb{W}) (\lambda \gamma^q \mathbb{W}) (\lambda \gamma^m \mathbb{W}).$$

Find ghost n<sup>o</sup> 2 expression  
+ Match in Superspace.

- Connection between  $\mathbb{W}$  and  $\text{NO-}\mathbb{W}$  Corrections,  
$$F^{m1pq} (\lambda \gamma^p \mathbb{W}) (\lambda \gamma^q \mathbb{W}) (\lambda \gamma^m \mathbb{W}).$$
- $(\alpha')$ <sup>4</sup> Corrections in Pure Spinor Superspace
  - Ansatz Over 12,000 terms
  - Canonicalization difficult.
  - ↳ Would yield Perturbative Expansion!  
n-point.

# Thank You