Non-Linear Solutions and the Double Copy

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Within the last 2 decades, a duality has been (known as the Double Copy) written down for amplitudes between Quantum Gravity and Yang-Mills theory.

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- At a conceptual level, the double copy challenges our current understanding of the foundations for both QFT and GR.

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- Within the last decade, this duality has been extended to classical relativistic theories (YM and General Relativity).
- At a conceptual level, the double copy challenges our current understanding of the foundations for both QFT and GR.
- The double copy has also found great success as a tool for research; having been utilised in Cosmological (Gravitational Waves) and Optical settings.

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- So far little attention has been paid to studying the duality between solutions existing within Euclidean Spacetimes.
- In particular, non-linear solutions have remained frustratingly elusive.

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The Kerr-Schild Double Copy



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The Kerr-Schild Double Copy



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Biadjoint Scalar Theory (BAST)

In $\mathbb{R}^{(1,d-1)}$, the Lagrangian density for BAST with Gauge structure $\mathfrak{su}(2) \times \mathfrak{su}(2)$ is given by:

$$\mathscr{L} = \frac{1}{2} \partial^{\mu} \Phi^{aa'} \partial_{\mu} \Phi^{aa'} + \frac{g}{3} \epsilon^{abc} \epsilon^{a'b'c'} \Phi^{aa'} \Phi^{bb'} \Phi^{cc'} \tag{1}$$

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Performing the Wick rotation $(\tau \rightarrow it)$ on (1), we derive the E.O.M for BAST in \mathbb{R}^d to be:

$$\triangle \Phi^{aa'} + g \epsilon^{abc} \epsilon^{a'b'c'} \Phi^{bb'} \Phi^{cc'} = 0$$
⁽²⁾

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Considering only spherical symmetric solutions, set:

$$\Phi^{aa'} = \frac{\delta^{aa'}}{gT_A} f(r) \tag{3}$$

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Inputting (3) into (2), we derive:

$$\frac{1}{r^{d-1}}\frac{d}{dr}\left(r^{d-1}\frac{df(r)}{dr}\right) + f^{2}(r) = 0,$$
(4)

Choosing a power-like ansatz:

$$f(r) = Ar^{\alpha} \tag{5}$$

We find a general solution for $\Phi^{aa'}$ in *d* dimensions:

$$\Phi^{aa'} = \frac{2\delta^{aa'}}{gT_A} \frac{d-4}{r^2}.$$
 (6)

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Where for d = 4, we find that there exists no power-like instanton solutions.

The Mystery of $d \neq 4$

To see why this is the case, choose a more general ansatz (in d dimensions):

$$f(r) = \frac{K(r) - 1}{r^2},$$
 (7)

Where taking $K(r) \rightarrow 1$ yields the trivial solution.

To see why this is the case, choose a more general ansatz (in d dimensions):

$$f(r) = \frac{K(r) - 1}{r^2},$$
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Where taking $K(r) \rightarrow 1$ yields the trivial solution. Substituting (7) into (4) and performing the substitution $r = e^{-\xi}$ ($-\infty < \xi < \infty$), we find:

$$\frac{\partial^2 K}{\partial \xi^2} - (d-6)\frac{\partial K}{\partial \xi} + (K-1)(K-2d+7) = 0.$$
 (8)

We must employ numerical methods in order to solve this (*critical points and integral curves*).

Numerical Solutions: Part 1

Setting $\psi = \frac{\partial K}{\partial \xi}$, we receive a pair of coupled differential equations:

$$\psi = \frac{\partial K}{\partial \xi} \tag{9}$$

$$\frac{\partial \psi}{\partial \xi} = (d-6)\psi - (K-1)(K-2d+7)$$
(10)

Now writing both (9) and (10) into a vector field :

$$\left(\frac{\partial K}{\partial \xi}, \frac{\partial \psi}{\partial \xi}\right) = \left(\psi, (d-6)\psi - (K-1)(K-2d+7)\right)$$
(11)

Solving (11) yields a vector field (K, ψ) , where in *d* dimensions, we have two solutions:

$$K = 1, \psi = 0 \tag{12}$$

$$K = 2d - 7, \psi = 0$$
 (13)

Notice, for d = 4, the only solution for $\Phi^{aa'}$ corresponds to the trivial solution. Thus, there are **no non-linear power-like** solutions for d = 4.

D = 4 Solution



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- In ℝ⁴, there exist no radial power-like solution to the full (non-linear) EOM for BAST.
- However, non-linear solutions do exist for $d \neq 4$.
- ► There exists solutions for the *linearised* EOM for BAST in ℝ⁴. Through the Double Copy, we know that these solutions are related to the **Eguchi Hanson gravitation instanton**.
- Investigation into such solutions would require the use of the Operator Formalism of the Kerr-Schild Double Copy. (10.1007/JHEP08(2022)160)

The Kerr-Schild Double Copy: Operator Formalism



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