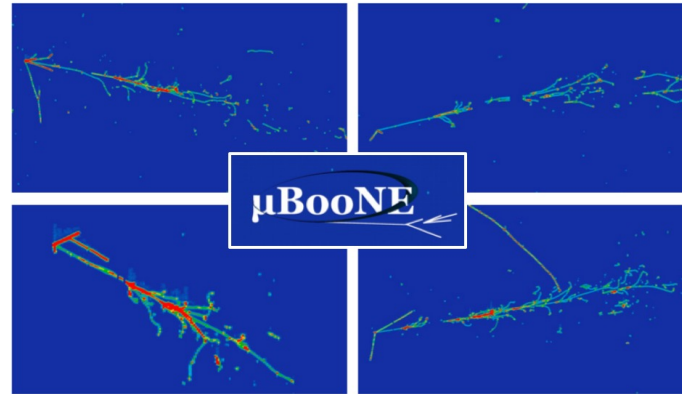


Sterile neutrinos



Justin Evans

Sterile neutrinos

Singlet states with no weak interactions

Dark sector fermion–portal models very naturally produce sterile neutrinos

$$\mathcal{L} \supset y (i\sigma^2 H^*) L N$$

Neutrino mass needs right-handed states

- Whether that's through Dirac or Majorana terms

$$\mathcal{L}_{\text{mass}} \supset -M_D^{\alpha\beta} \nu^\alpha N^\beta + h.c.$$

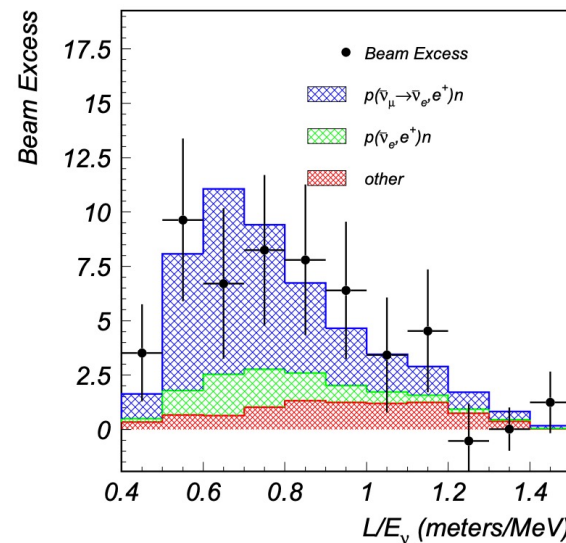
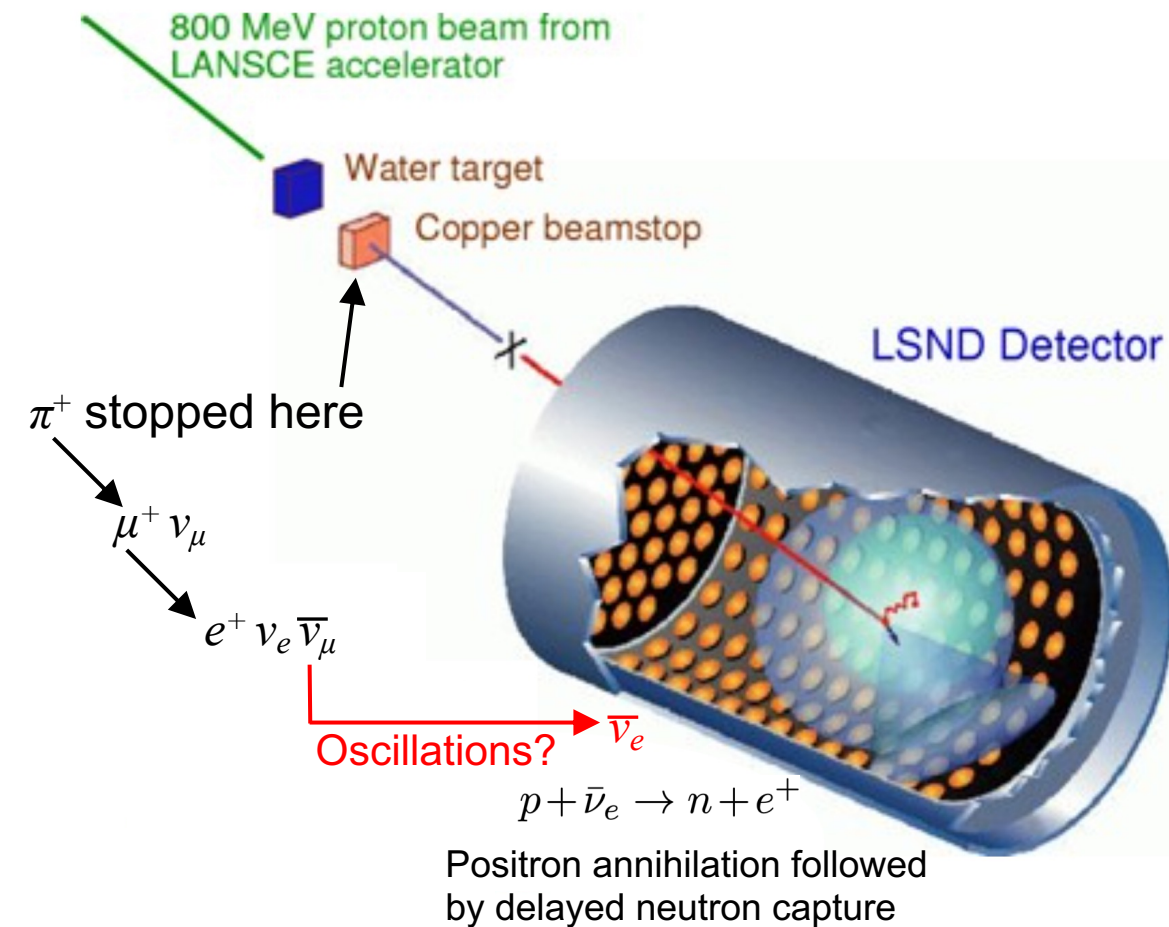
Sterile neutrinos are theoretically very favourable

- 'Just a neutral fermion'
- But tend to be heavy – MeV scale or above

The experimental hints are for light sterile neutrinos

- Around the eV scale

LSND



$87.9 \pm 22.4 \pm 6.0$ excess
 $\bar{\nu}_e$ -like events observed

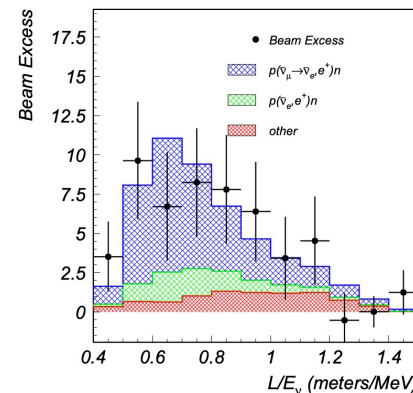
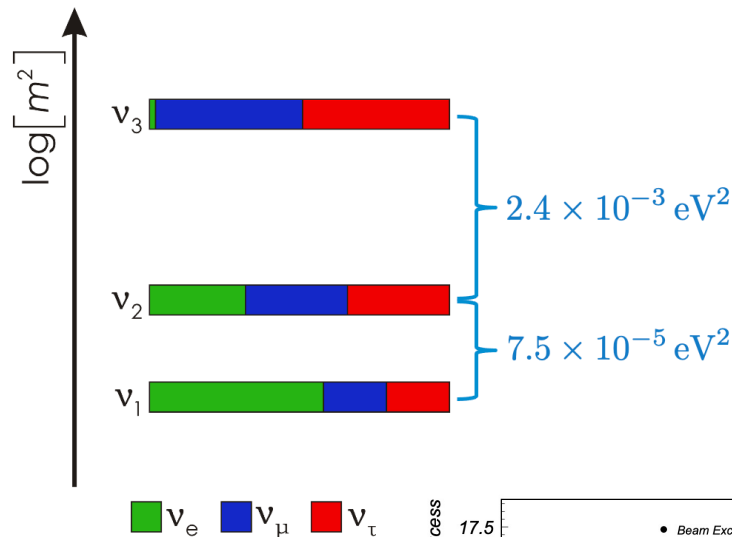
3.8σ significance

The oscillation interpretation

The neutrinos have not traveled far enough for standard oscillations to occur

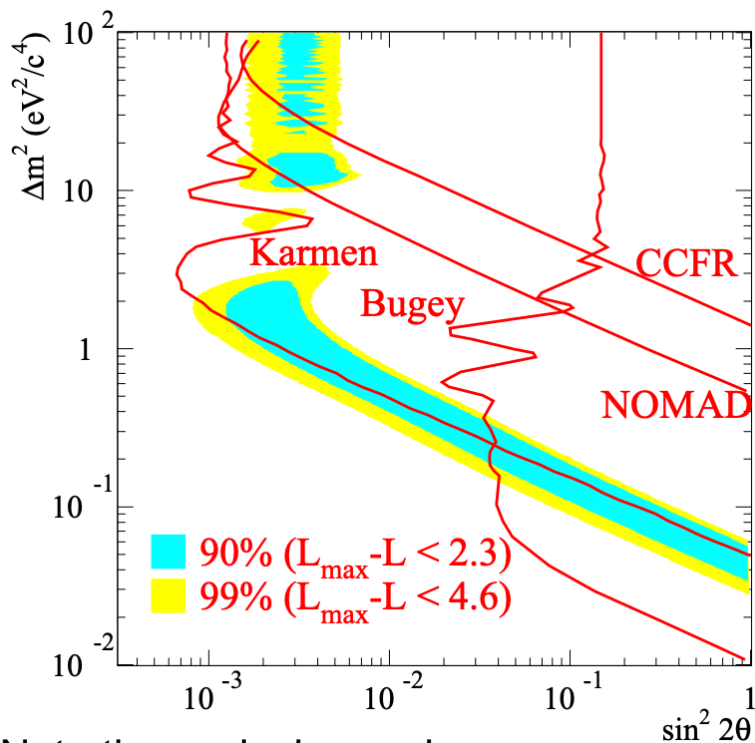
Requires a mass splitting $O(1 \text{ eV}^2)$

➤ Fourth neutrino state

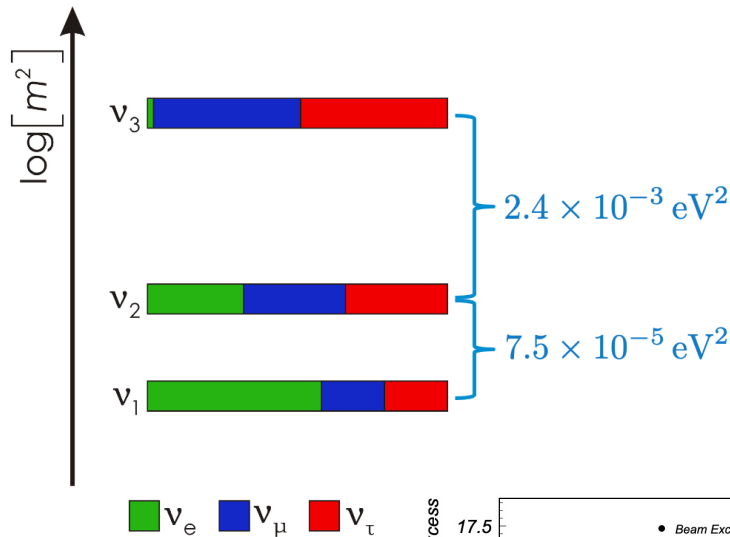


$$P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu) = \sin^2(2\theta) \sin^2 \left(1.27 \Delta m^2 \frac{L}{E} \right)$$

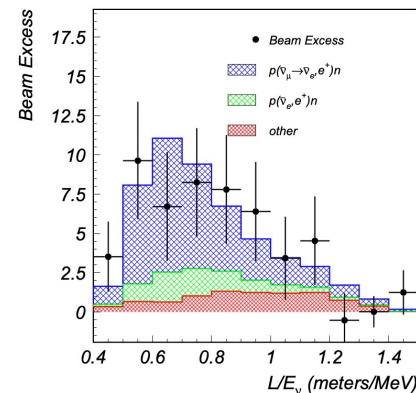
The oscillation interpretation



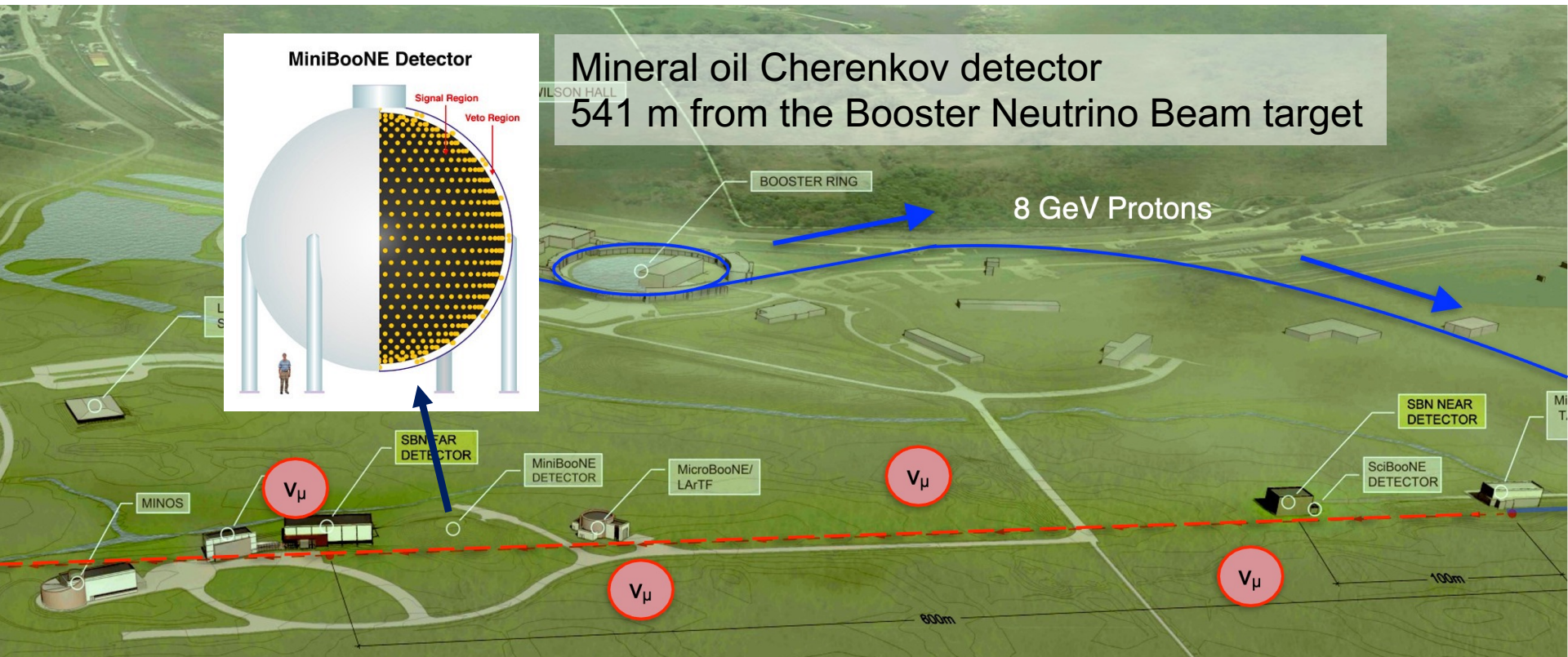
Note the exclusion regions,
particularly Karmen



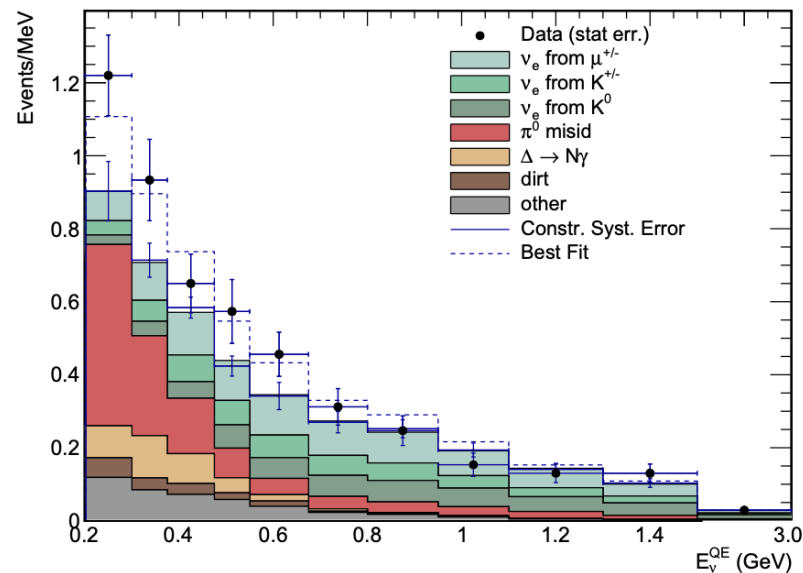
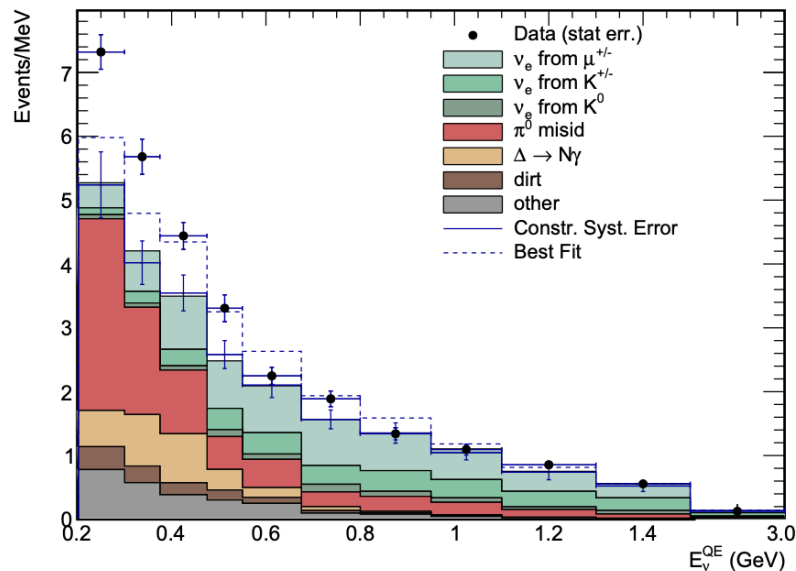
$$P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu) = \sin^2(2\theta) \sin^2 \left(1.27 \Delta m^2 \frac{L}{E} \right)$$



MiniBooNE



MiniBooNE



Excesses of electron-like events observed

➤ In both neutrino and antineutrino runs

MiniBooNE

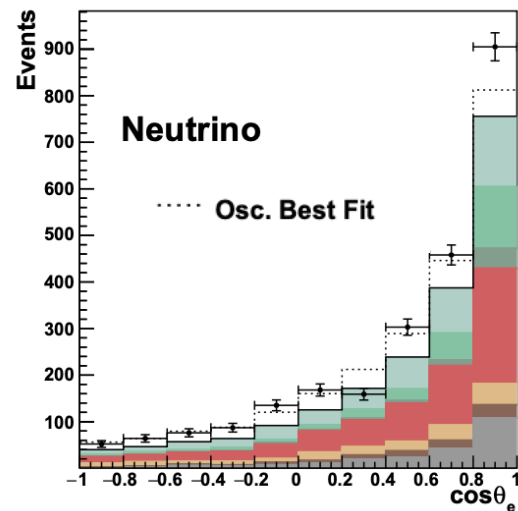
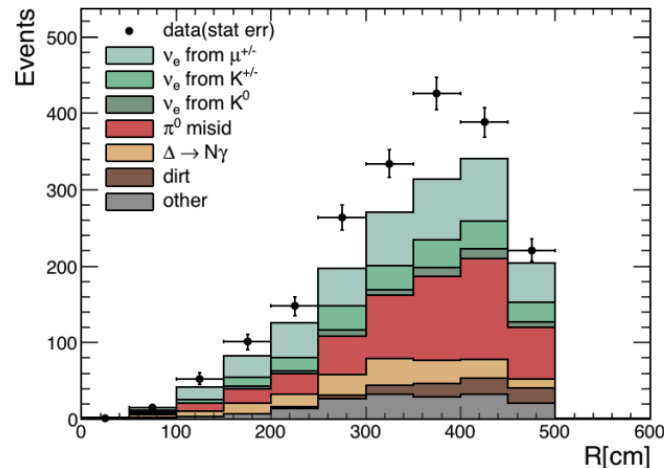
Excess is distributed throughout the detector

- Not consistent with escaping photons from π^0 decay

π^0 rate is measured in data

Excess enhanced in forward lepton angles

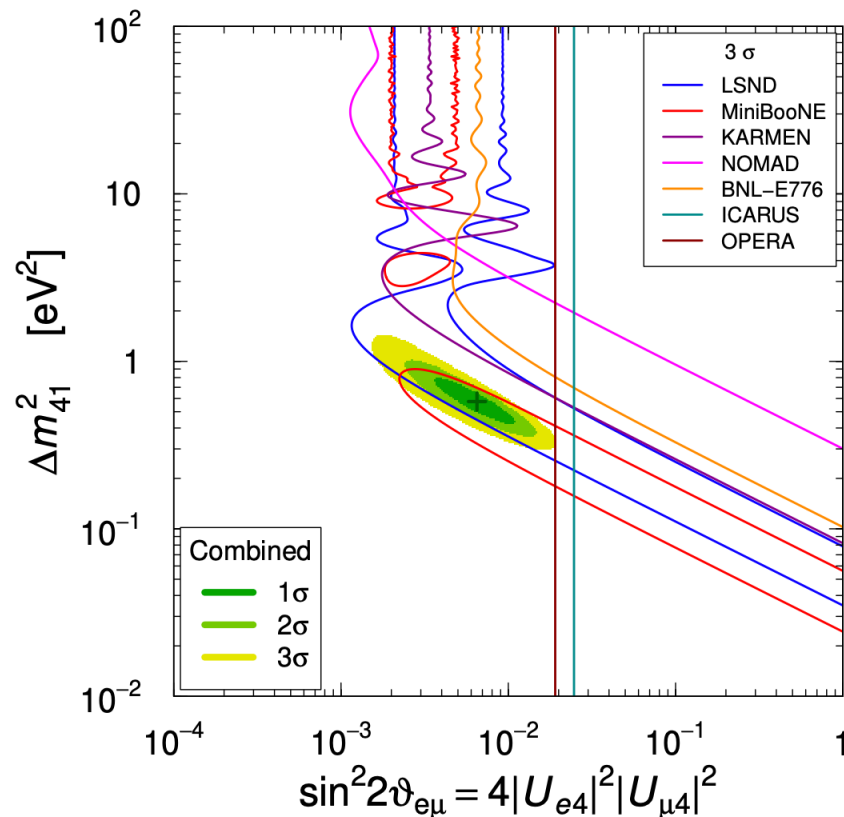
- Less consistent with oscillations



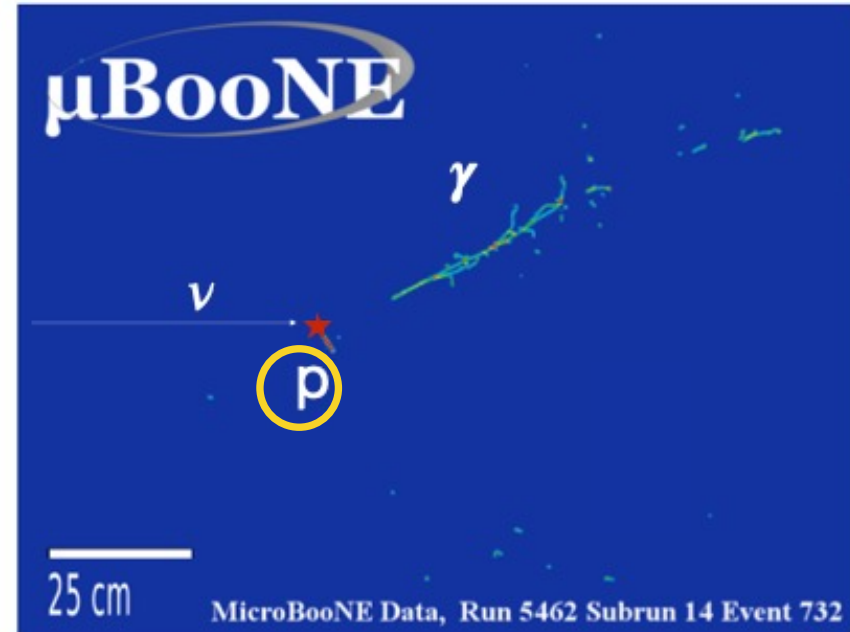
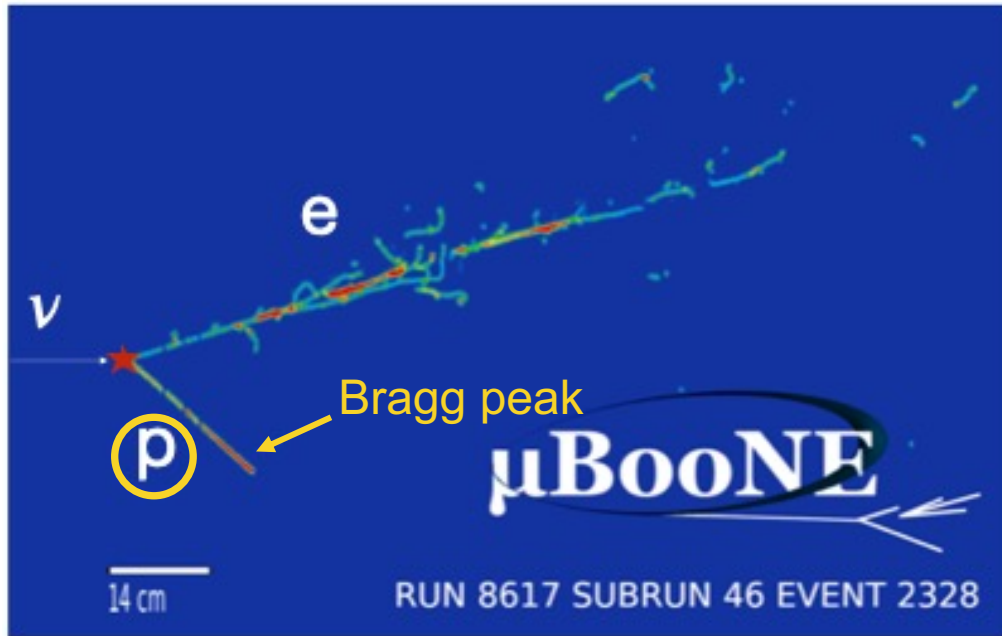
MiniBooNE & LSND

When combined together,
along with exclusion
regions, a mass splitting a
little below 1 eV^2 can fit the
data reasonably well

- e.g. Giunti & Lasserre,
Ann. Rev. Nucl. Part. Sci.
69, 163 (2019)

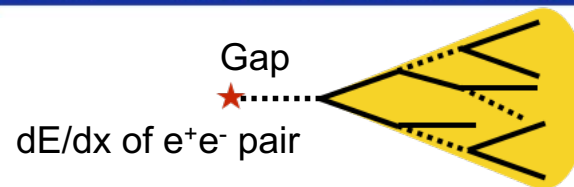
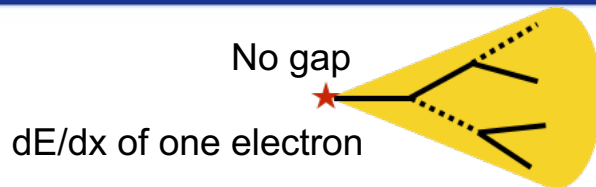
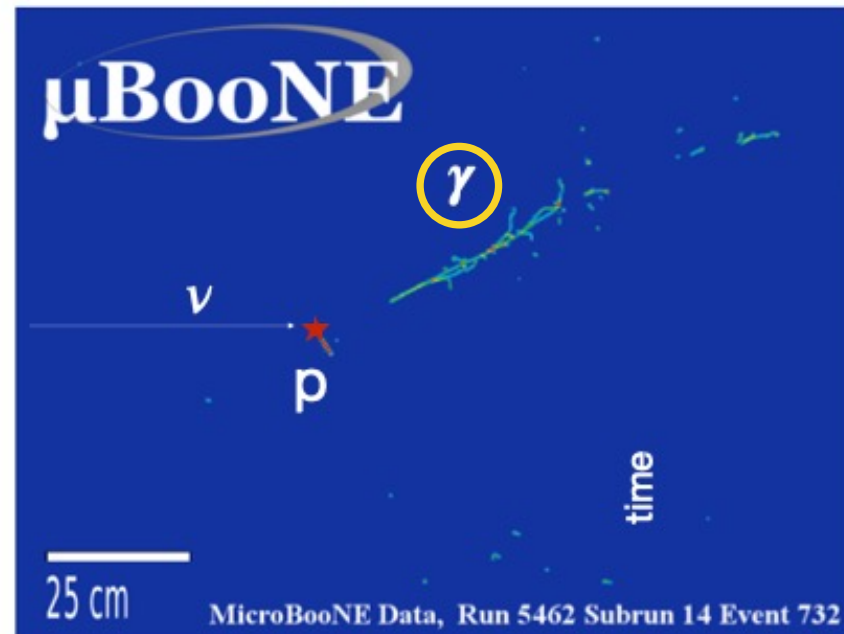


Liquid argon: MicroBooNE

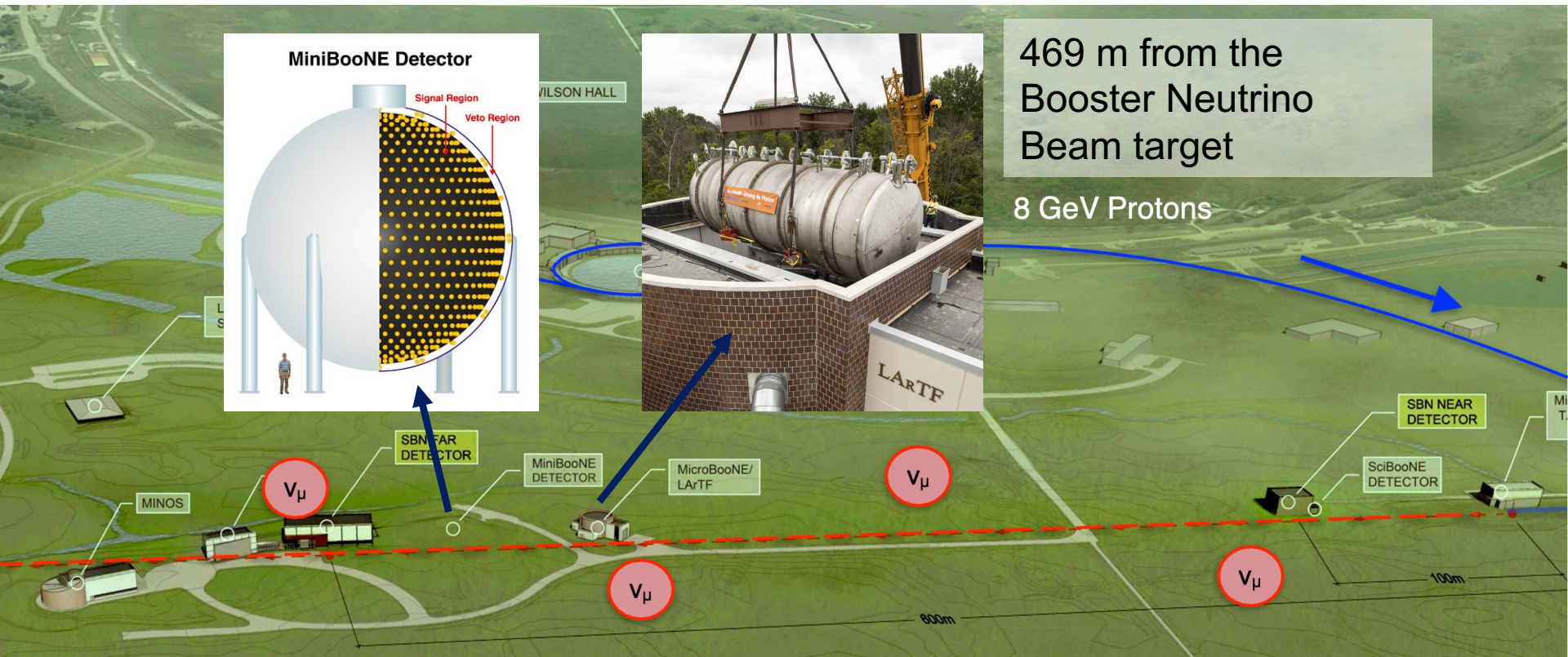


We can see the hadronic final state

Liquid argon: MicroBooNE



Liquid argon: MicroBooNE



First low-energy excess search

Four independent analyses

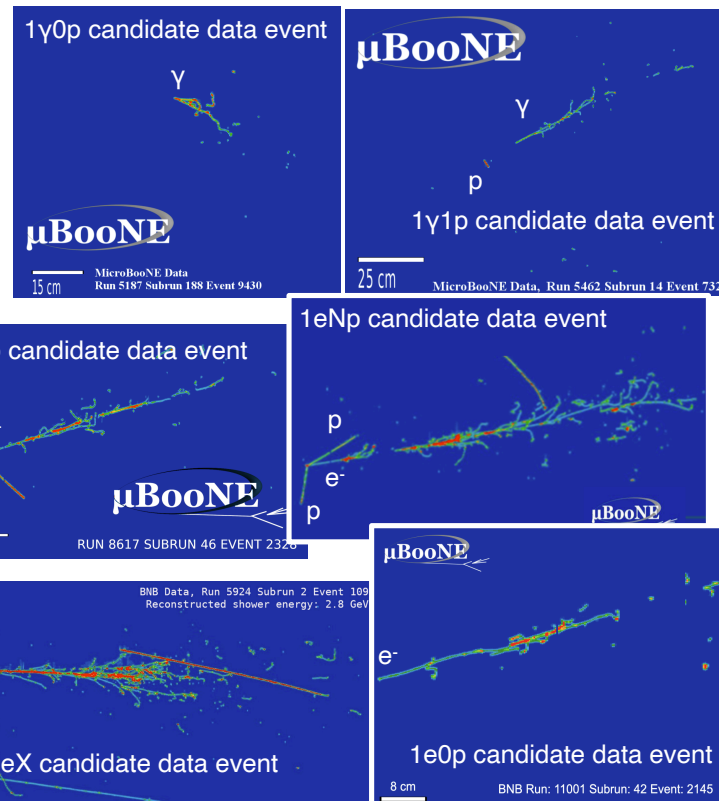
- Targeting six different final states

Single-photon analysis

- NC $\Delta \rightarrow N\gamma$ hypothesis
- $1\gamma 0p$, $1\gamma 1p$

Searches for a ν_e excess

- Quasi-elastic kinematics ($1e1p$)
- MiniBooNE-like final states ($1eNp$, $1e0p$)
- All ν_e final states ($1eX$)



Single photons: $\Delta \rightarrow N\gamma$

Several photon sources in MiniBooNE

NC π^0 misidentification

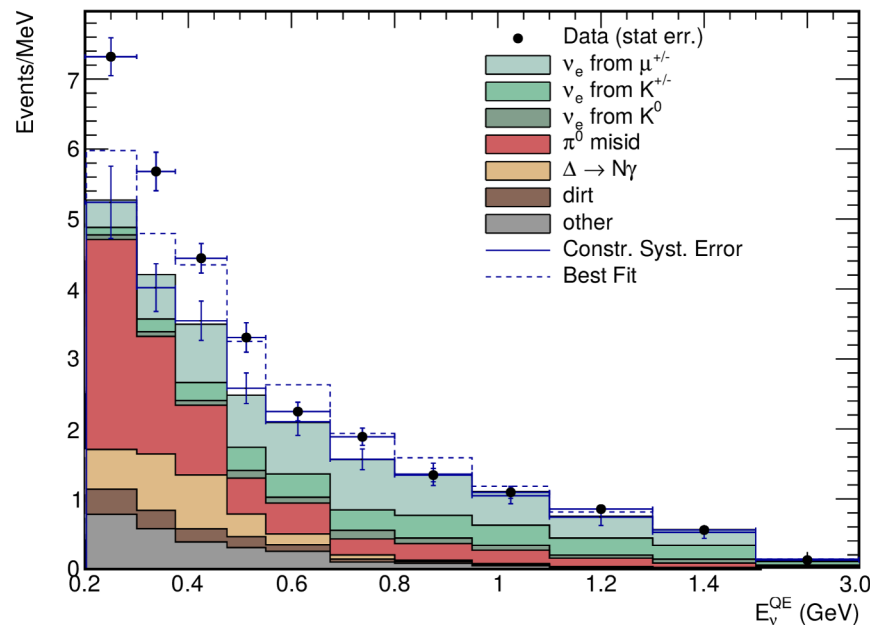
- Measured in MiniBooNE with sidebands

Interactions outside the detector

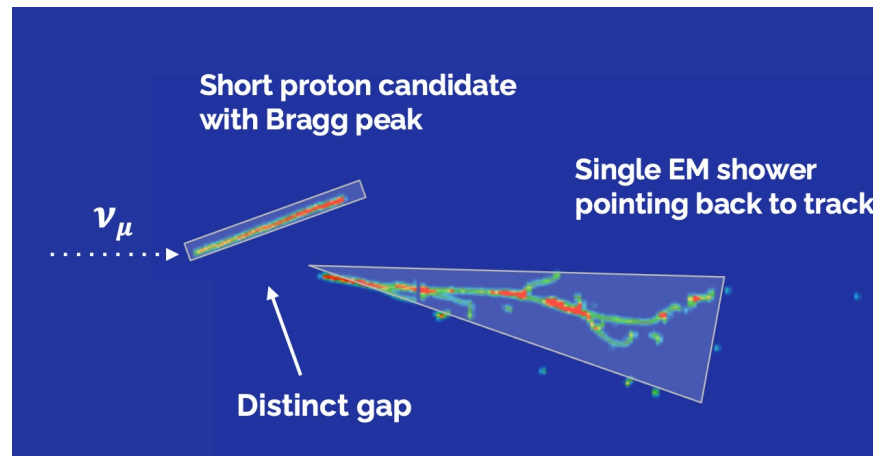
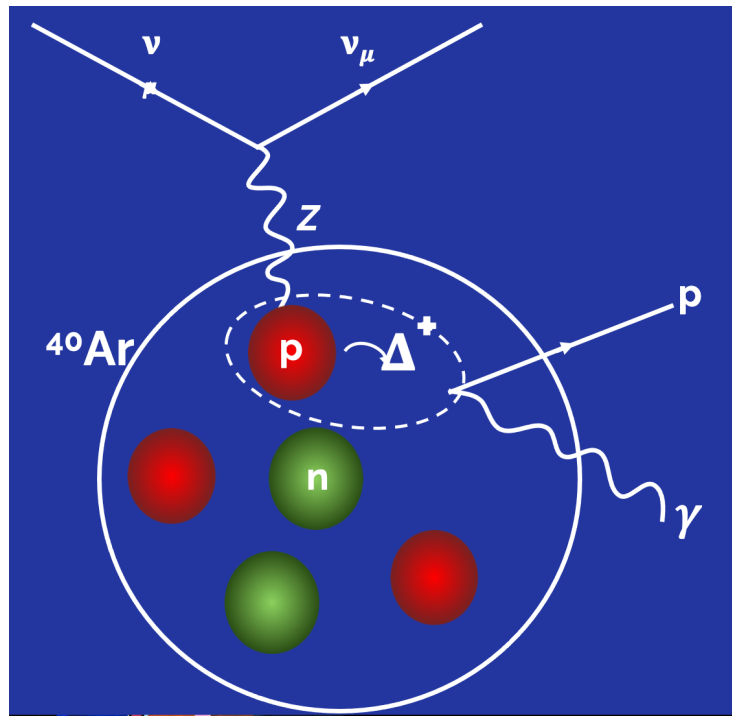
- Eliminated using beam timing and radial cuts

NC $\Delta \rightarrow N\gamma$

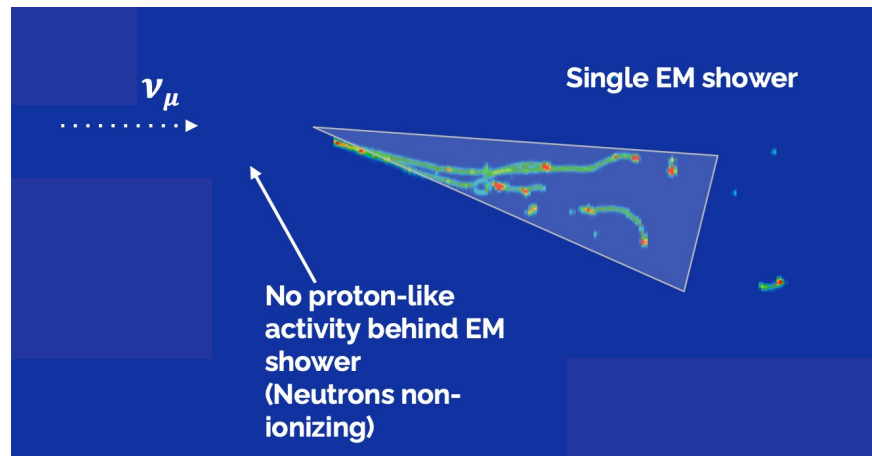
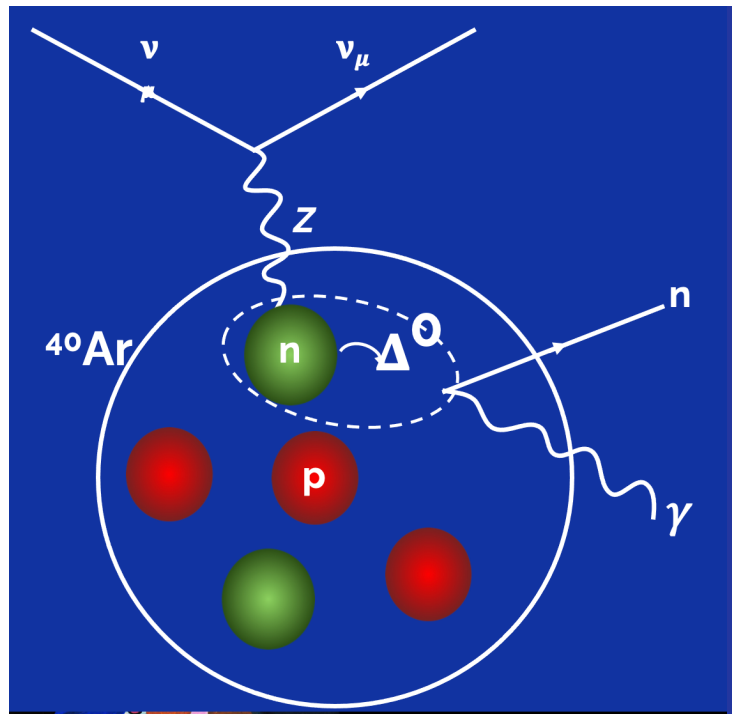
- **NC delta radiative decay**
- Not constrained directly by MiniBooNE
- Used π^0 measurements and a theoretical branching ratio for the delta radiative decay



$\Delta \rightarrow N\gamma$: $1\gamma 1p$ topology



$\Delta \rightarrow N\gamma$: $1\gamma 0p$ topology

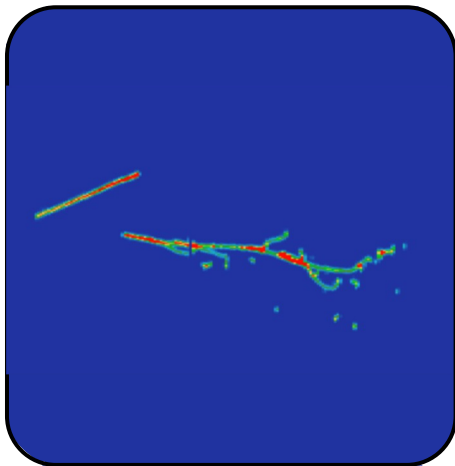


Single photon analysis

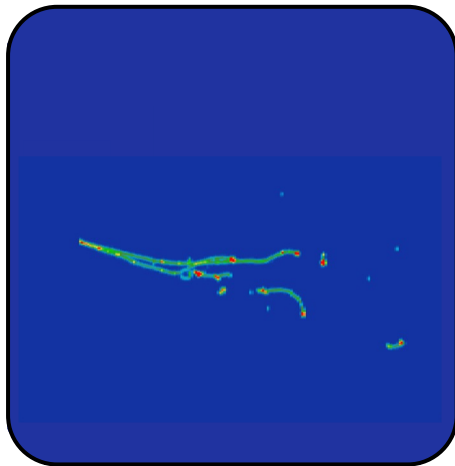
Simultaneous fit of four topologically distinct samples

Two **NC $\Delta \rightarrow N\gamma$** rich **single-photon** selections

1 γ 1p

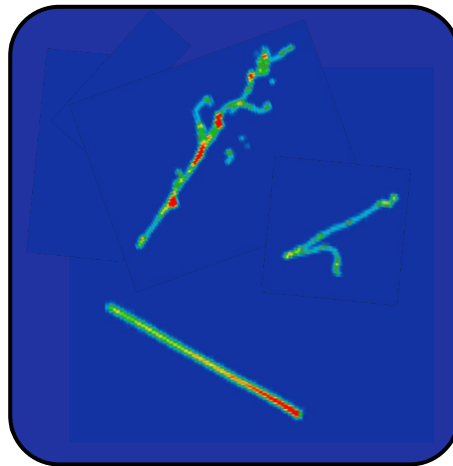


1 γ 0p

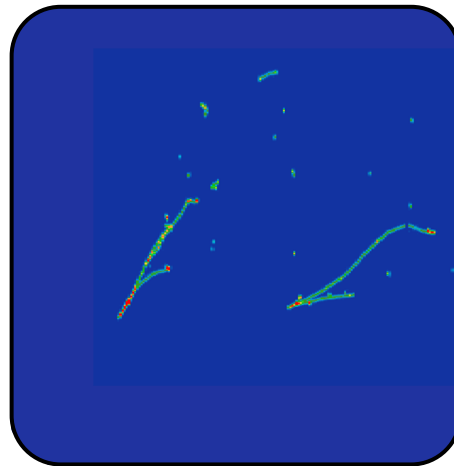


Two high-statistics **NC π^0** **rich two-photon** selections to constrain backgrounds

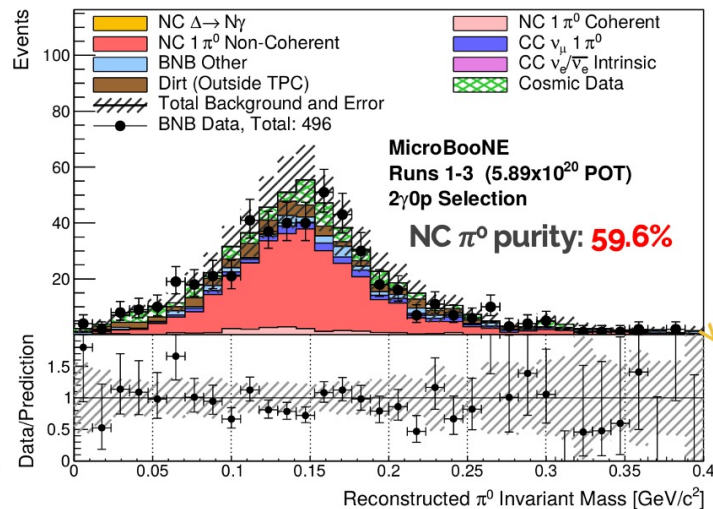
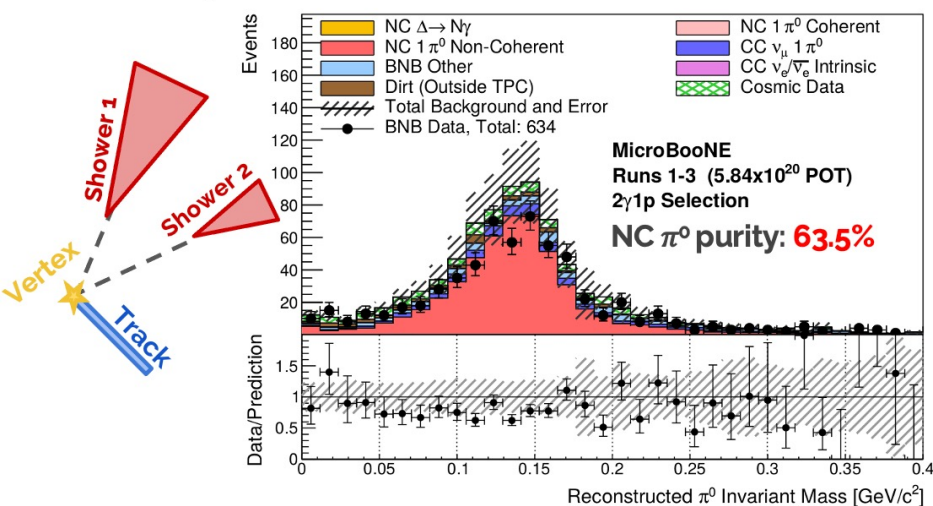
2 γ 1p



2 γ 0p



1p2 γ and 0p2 γ samples



High statistics: 1130 candidate π^0 interactions

Used to constrain the π^0 backgrounds in the $\Delta \rightarrow N\gamma$ signal samples

➤ And validate shower reconstruction and energy measurement

Single-photon results

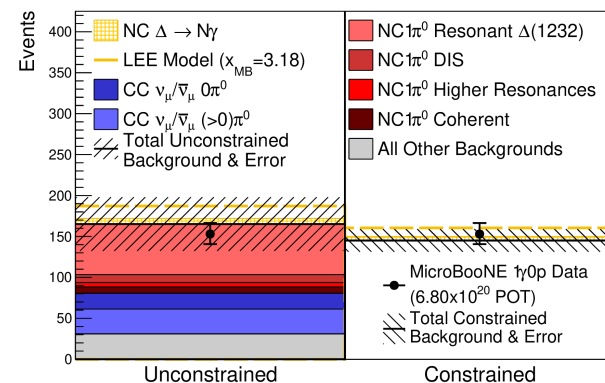
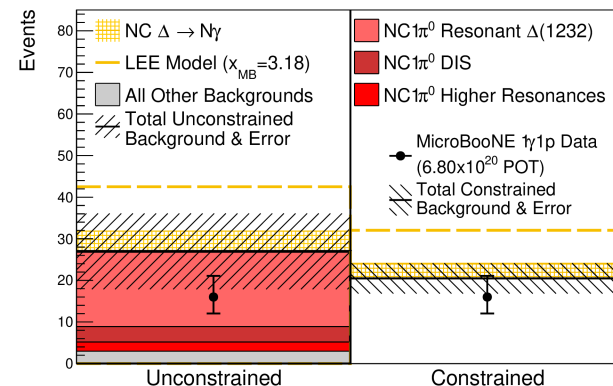
No evidence for an enhanced rate of single photons from NC $\Delta \rightarrow N\gamma$ decay above nominal GENIE expectations

➤ x3.18 scaling disfavoured at 94.8% C.L.

One-sided bound on the normalisation of NC $\Delta \rightarrow N\gamma$ events of $x_{\Delta} < 2.3$ (90% C.L.)

$$\mathcal{B}_{\text{eff}}(\Delta \rightarrow N\gamma) < 1.38\% \text{ (90\% C.L.)}$$

More than 50 times better than the world's previous limit



Electron search

Three independent searches across multiple final states

Two-body CC quasi-elastic: $1e1p$

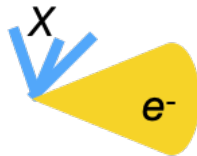


CC ν_e interactions without final-state pions: $1eNp0\pi$ & $1e0p0\pi$

- Matches the topology of MiniBooNE events



Fully inclusive charged-current ν_e : $1eX$



Start with muon neutrinos

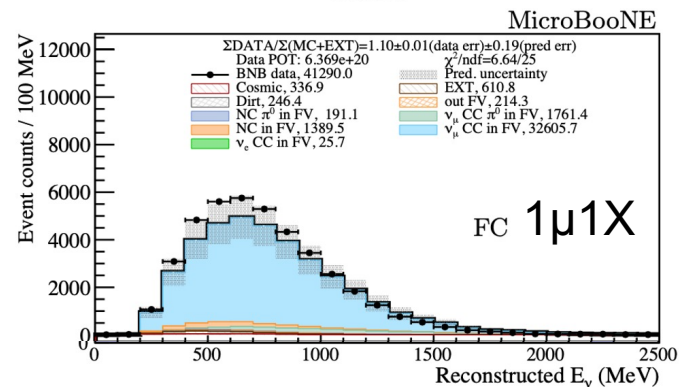
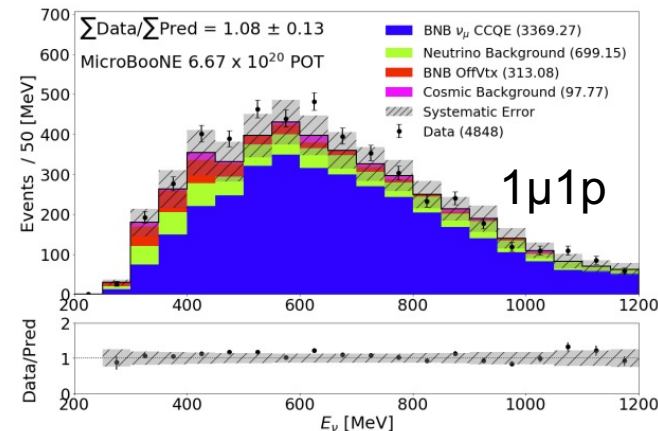
High-statistics CC ν_μ samples

Leverage ν_μ and ν_e correlations

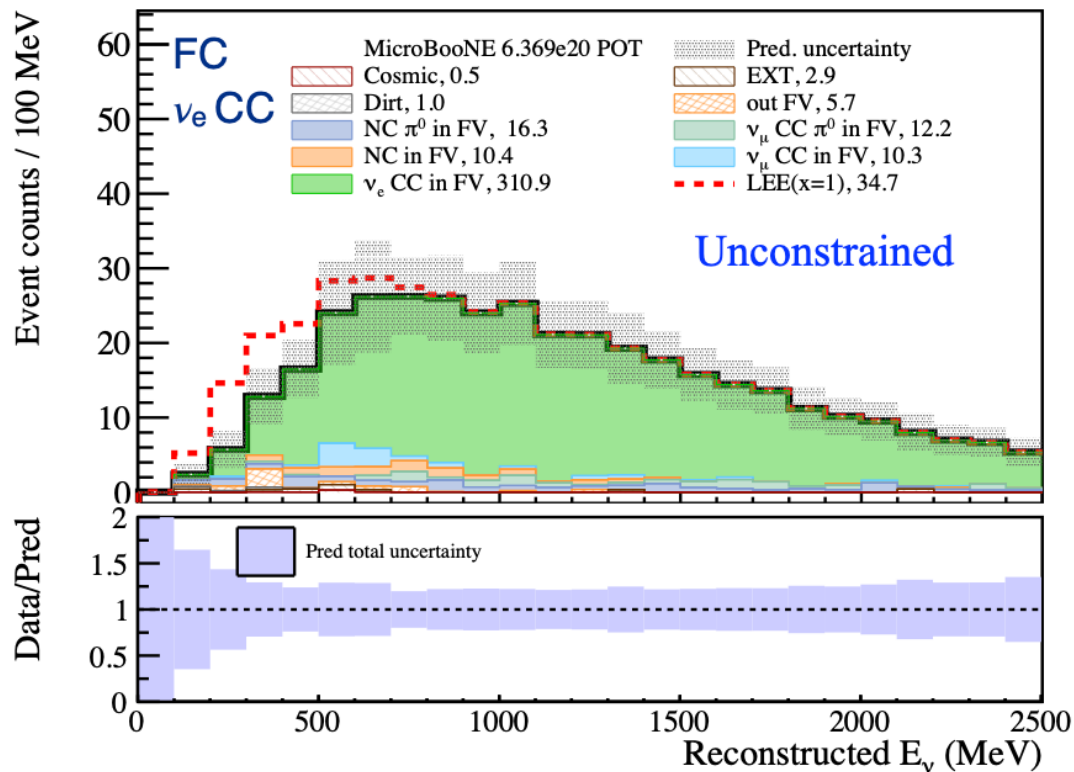
- Common flux parentage
- Lepton universality

Use our ν_μ sample to create a data-driven ν_e prediction

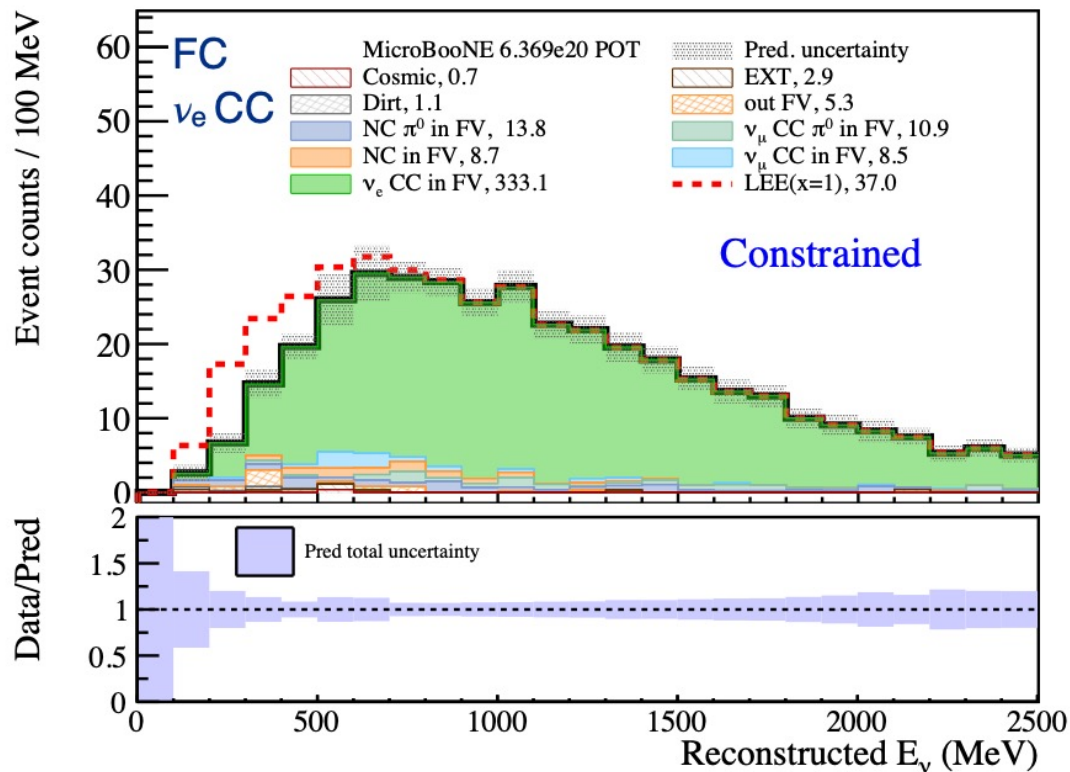
- Systematic uncertainties incorporated through a covariance matrix
- This process reduces the uncertainty on the ν_e prediction



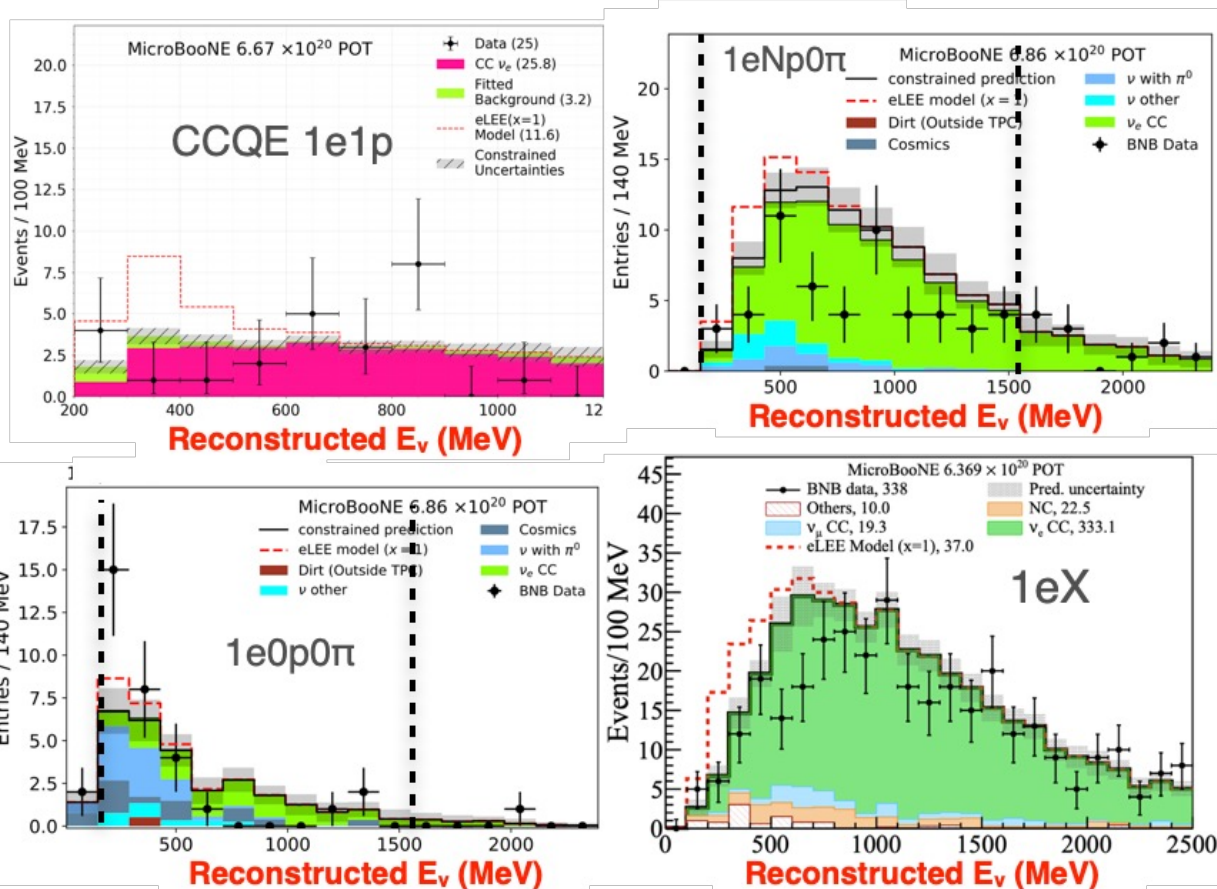
1eX prediction before constraint



1eX prediction after constraint

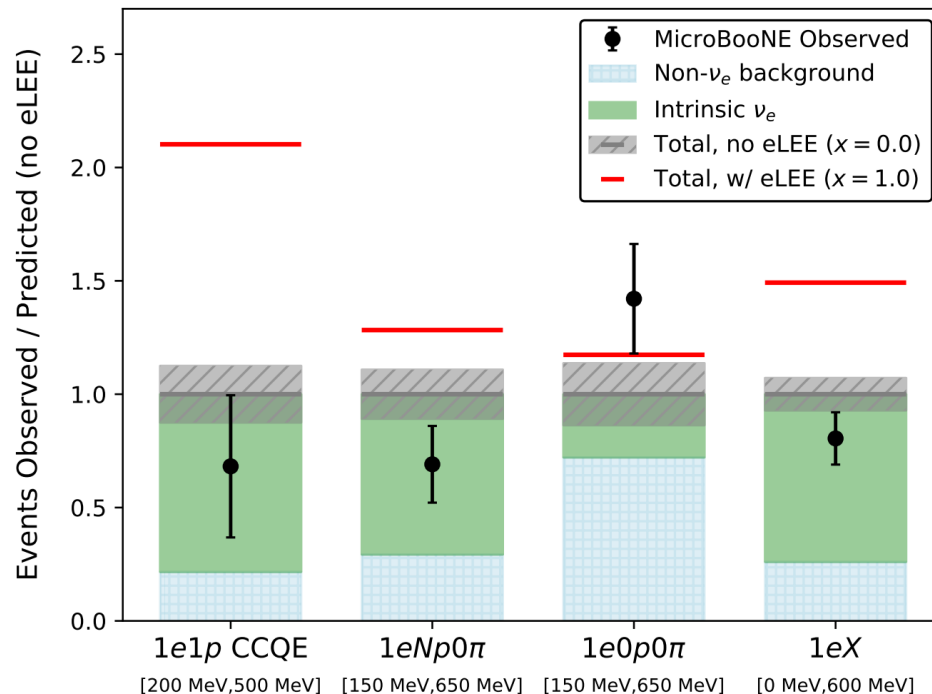


Electron results

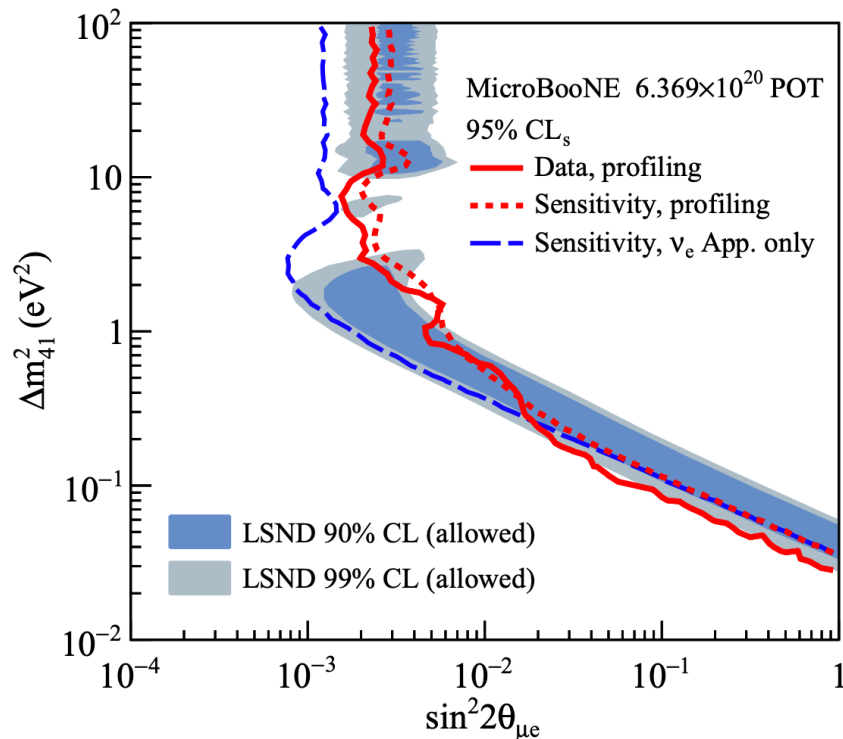


Electron results

- Observe ν_e candidate event rates in agreement with, or below, the predicted rates
- Reject the hypothesis that ν_e CC interactions are fully responsible for the MiniBooNE excess at $>97\%$ C.L. in all analyses
- Inclusive analysis rejects our median MiniBooNE electron-excess model at 3.75σ



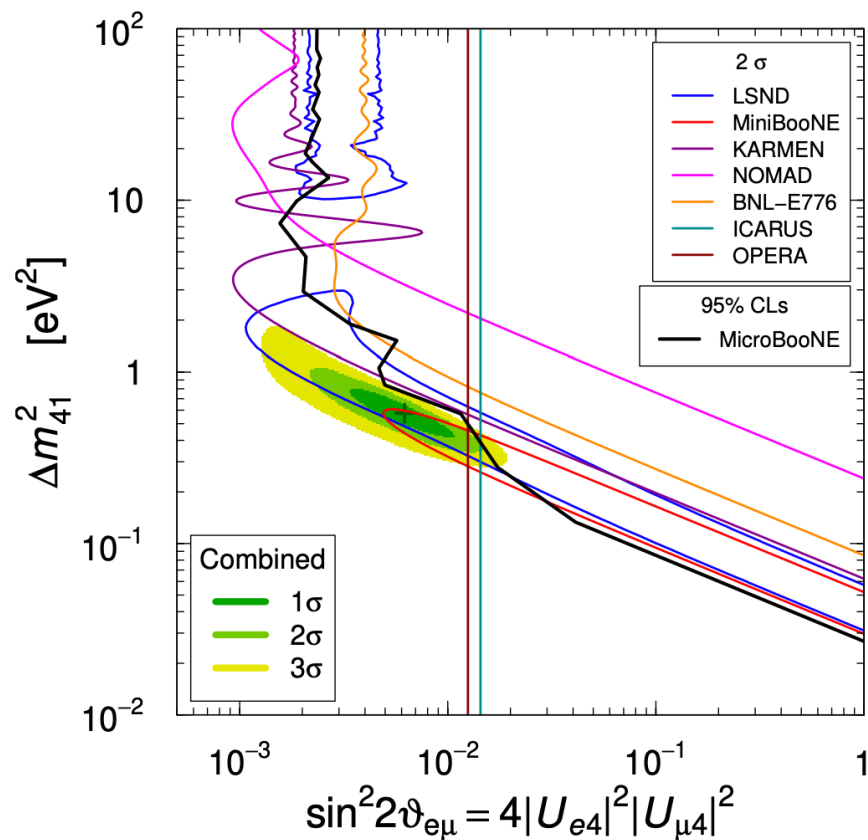
MicroBooNE oscillation limits



With the first half of our data, we are excluding part of the LSND allowed region

- Note that a degeneracy between ν_e disappearance and appearance reduces our sensitivity
- But we will use the NuMI beam – a second baseline – to overcome this in a future analysis

MicroBooNE and the global fits



$\nu_\mu \rightarrow \nu_e$ is not the only observable

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

θ_{14} : ν_e disappearance

θ_{24} : ν_μ disappearance

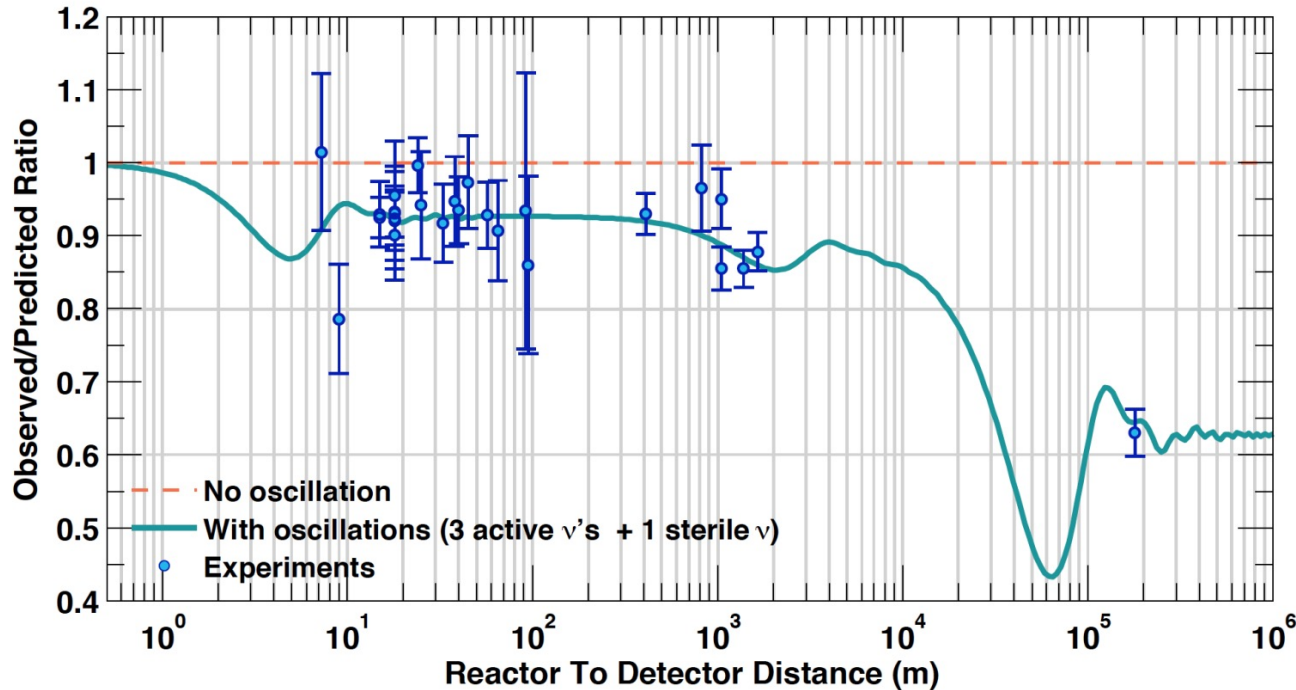
θ_{34} : ν_τ disappearance

$\sin^2(2\theta_{14}) \sin^2 \theta_{24}$: $\nu_\mu \rightarrow \nu_e$

The simplest sterile neutrino oscillation model has a single sterile flavour+mass state

- Introduces three new mixing angles

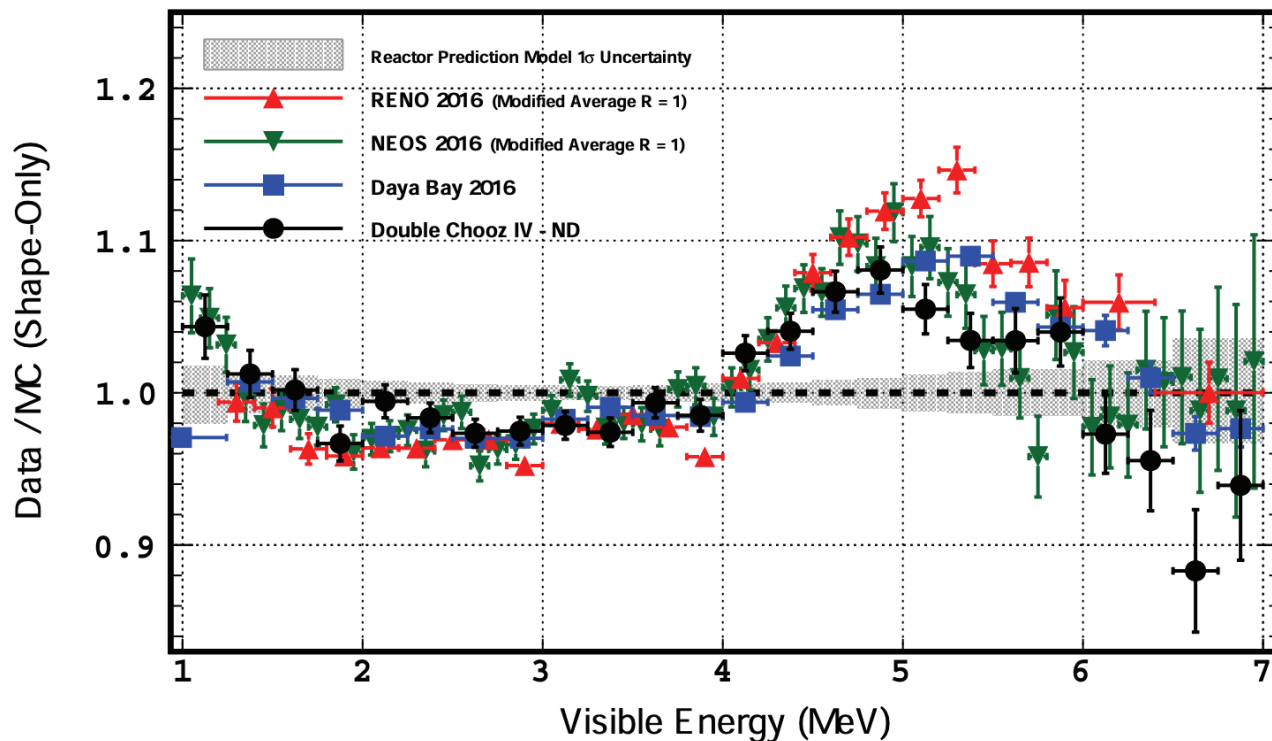
Reactor neutrinos



In 2011 recalculations of reactor neutrino fluxes left historic experiments now seeing few-percent deficits at $\sim 3\sigma$ significance

➤ Non-zero θ_{14} ?

Reactor neutrinos



➤ But can we trust reactor flux predictions at the few-percent level?

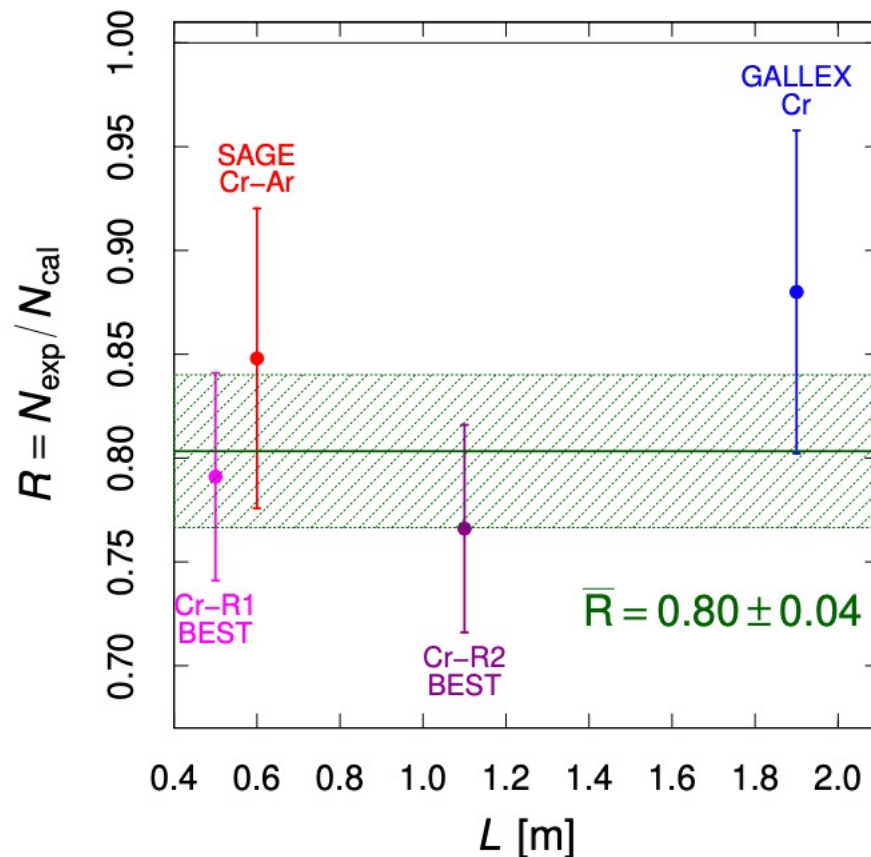
Gallium detectors

Calibration runs with ^{51}Cr and ^{37}Ar sources at the SAGE and GALLEX solar-neutrino experiments

- Recently confirmed by the BEST experiment

Deficits of ν_e interactions seen with respect to the expectation

- Requires dead-reckoning of cross-sections, but all cross-section models confirm the deficit
- No clear oscillatory pattern seen as a function of distance



Electron-neutrino disappearance

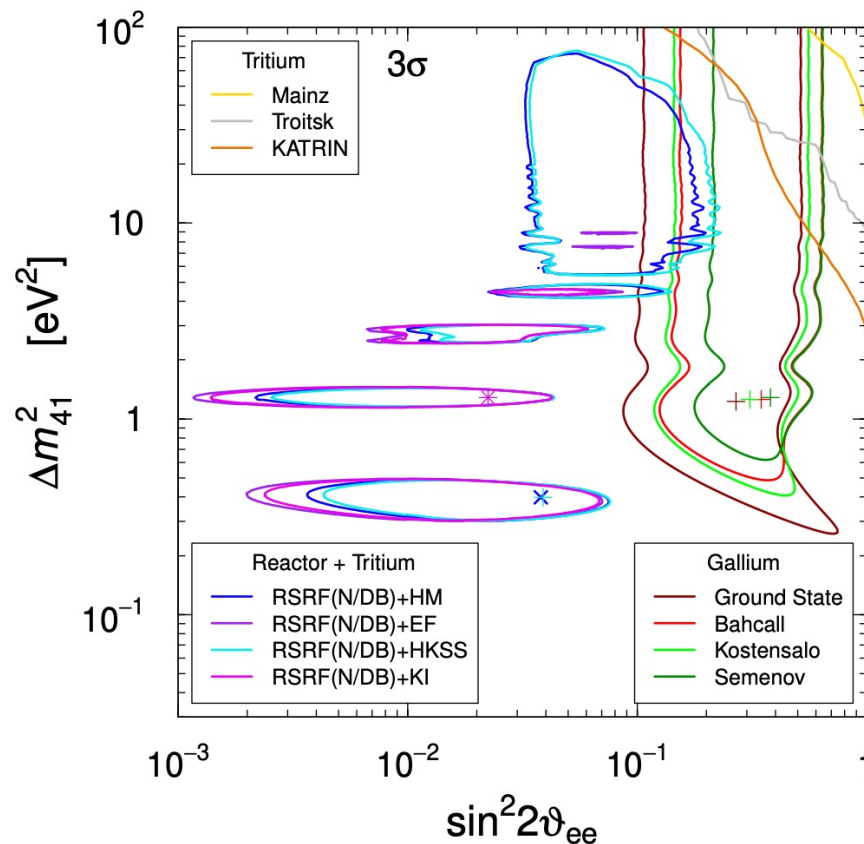
Tritium experiments (such as KATRIN) also place limits at high Δm^2

$$m_\beta^2 = \sum_{i=1}^3 |U_{ei}|^2 m_i^2.$$

No flux or cross-section model choices can get the reactor and gallium anomalies to agree on a sterile-neutrino cause

- And the gallium anomaly is in tension with bounds from solar neutrinos

N.B. the Neutrino-4 anomaly has received some criticism, and is largely ruled out by more recent measurements

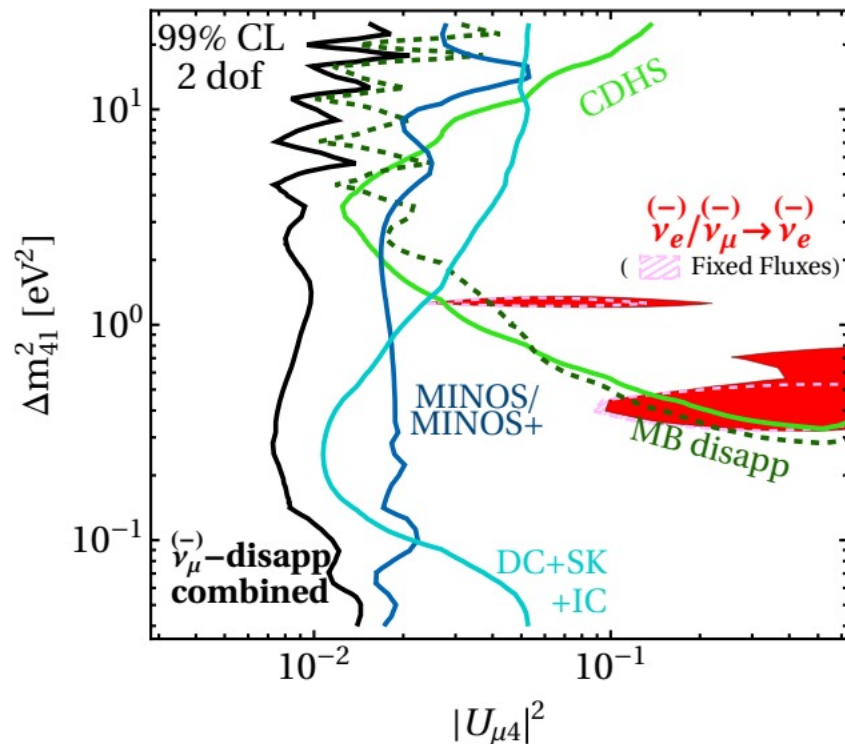


Muon neutrino disappearance

If $\nu_\mu \rightarrow \nu_e$ happens, then muon neutrinos must disappear

Here, the red area shows the level of ν_μ disappearance required to explain the $\nu_\mu \rightarrow \nu_e$ appearance in LSND, given the ν_e -disappearance allowed regions

- The ν_μ disappearance results strongly exclude this red region
- Adding more neutrinos (beyond 3+1) does not help

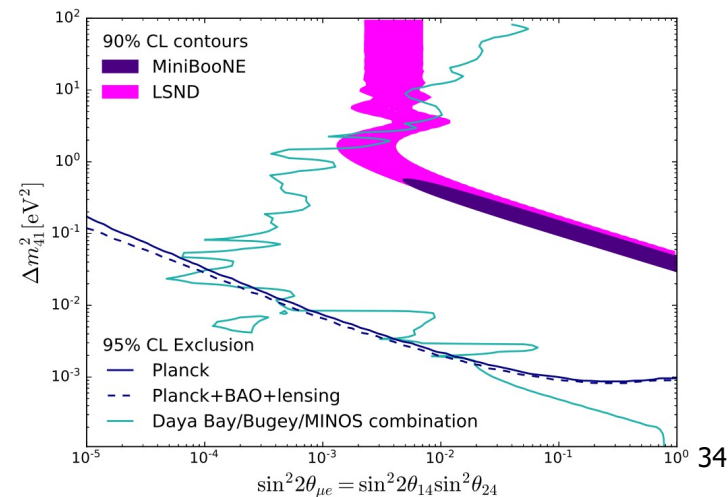
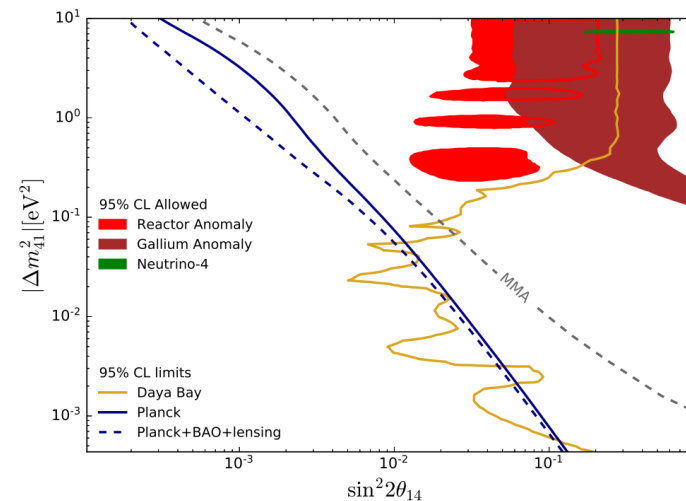


Cosmological constraints

Planck data sets strong limits on the effective neutrino mass

- This mass would increase if active flavour states mixed into heavier mass states

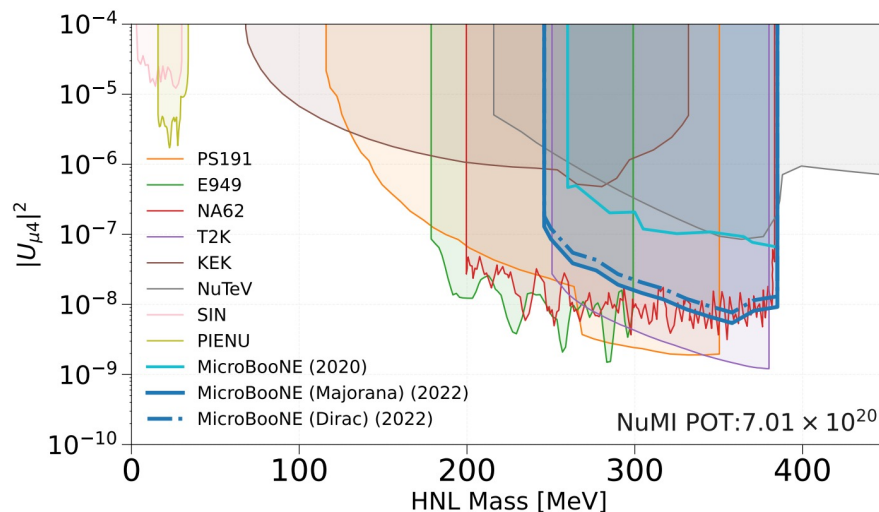
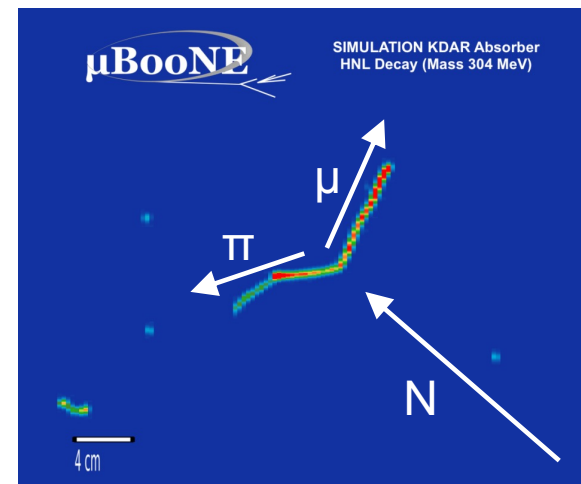
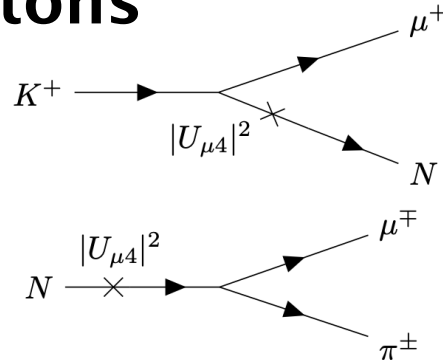
These strong limits also exclude the parameter space needed to explain the $\nu_\mu \rightarrow \nu_e$ and ν_e -disappearance anomalies with sterile neutrinos



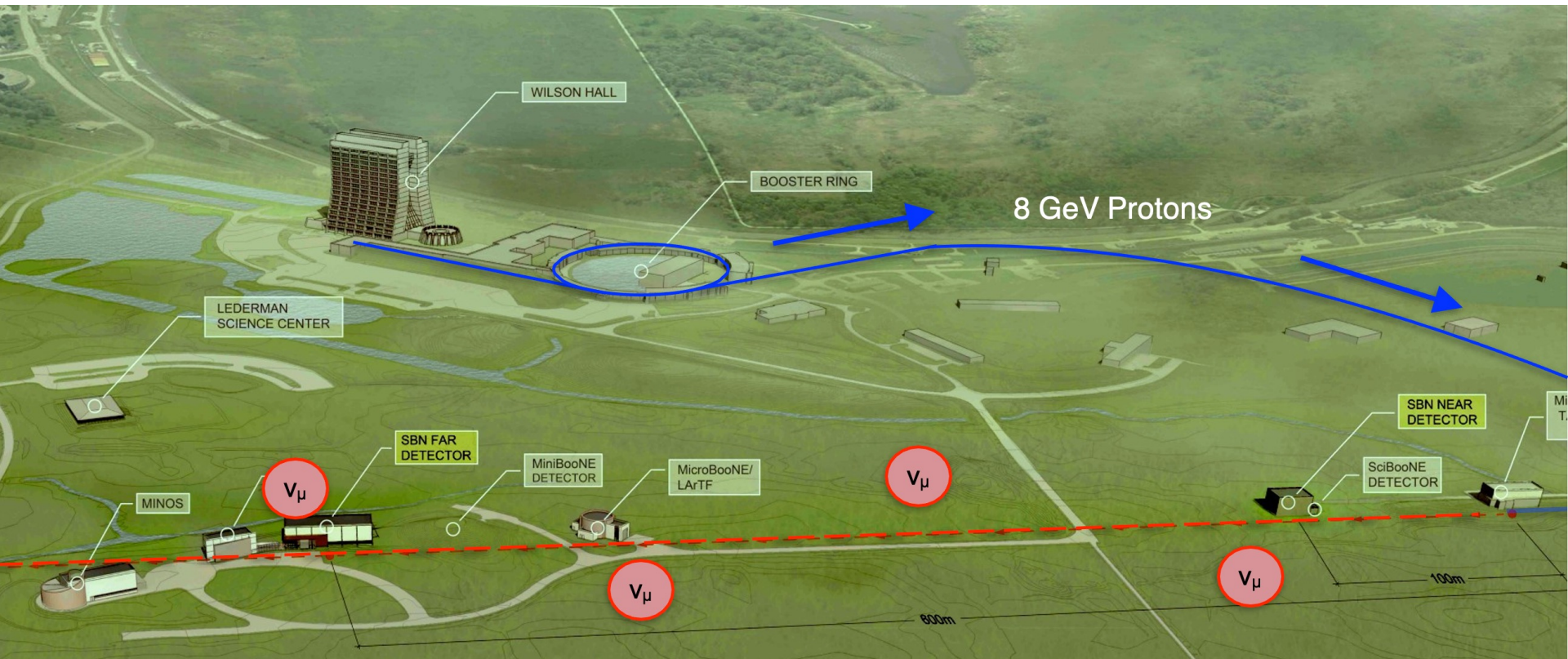
Heavy neutral leptons

Many theoretically-favoured sterile neutrinos are MeV-scale or higher in mass

- We can look for these in short-baseline experiments as they would be produced in the beam and decay in the detector



Onwards to the SBN programme



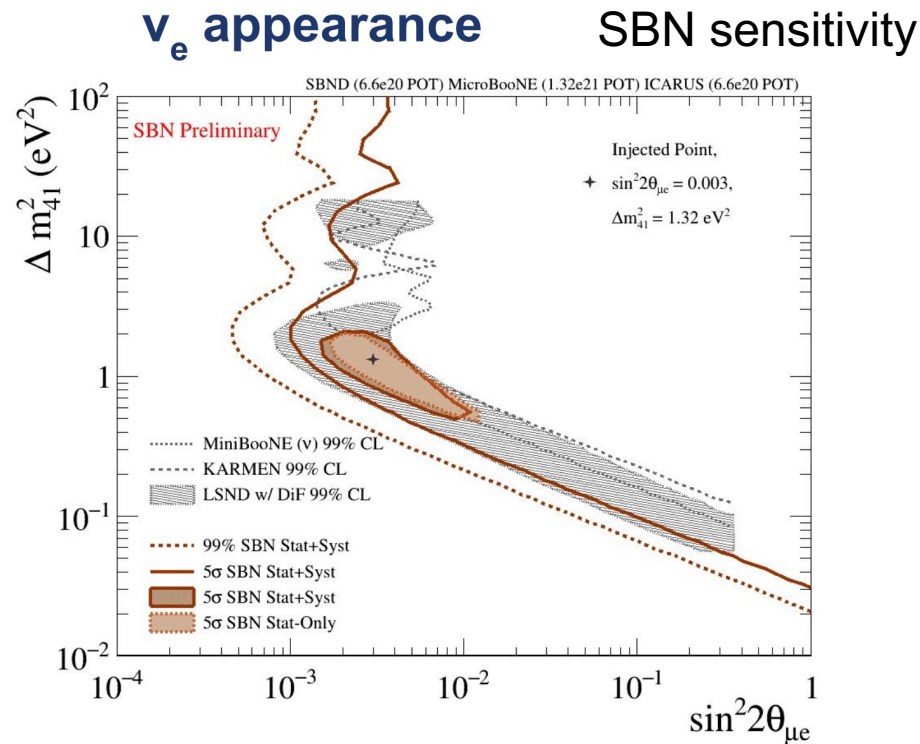
The SBN programme

MicroBooNE alone still has half its data to analyse

- And by including data from the NuMI can resolve degeneracies between ν_e appearance and disappearance
- Expect new results, with significantly improved sensitivity, next year

The SBN programme will take a three-detector approach to the sterile-neutrino problem

- With multiple baselines, the need to dead-reckon cross-sections and fluxes cancels out



Summary

Neutrino physics is hard!

- Dead-reckoning fluxes and cross sections is fraught with challenges
- Extensive neutrino cross-section programmes (e.g. MicroBooNE, SBN) are essential
- Multiple baselines (i.e. near detectors) are essential to cancel model uncertainties

There are numerous appearance and disappearance anomalies that can individually be explained by a light sterile neutrinos

- But looked at together, no sterile neutrino can fit all the pieces of the puzzle

Heavy sterile neutrinos are theoretically favourable

- An active experimental programme is ongoing here

We cannot yet explain away many of the anomalies

- But future programmes, such as the SBN programme, will continue to shed light on the short-baseline behaviour of neutrinos

