

Flavour-aware lepton dressing algorithms in the presence of photon splittings

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Introduction

- ▶ We are in an era of precision phenomenology in experiment and in theory
- ▶ Object definitions that are convenient both theoretically and experimentally are crucial
- ▶ Traditional cone-dressed leptons with photons are convenient until photons are further resolved

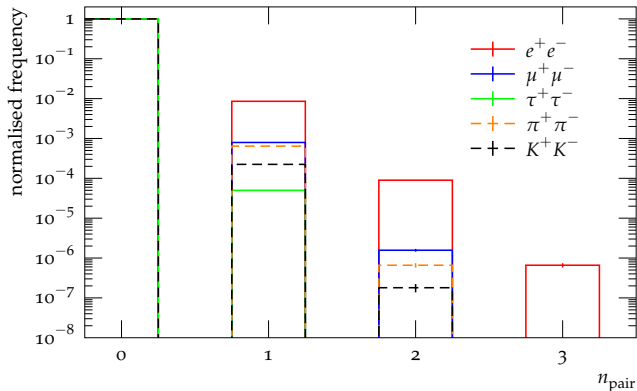
YFS and photon splittings

- ▶ Photons are usually produced as QED corrections to a process
- ▶ If these photons are further resolved, we should reevaluate lepton dressing procedures
- ▶ Here we use the YFS soft-photon resummation as implemented in SHERPA for photon production

Yennie, Frautschi, Suura '61

Krauss, Schönherr '08

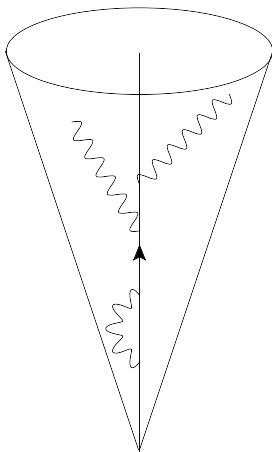
- ▶ We implemented a photon splitting algorithm which allows $\gamma \rightarrow f\bar{f}$ to occur, where $f = e, \mu, \tau, \pi, K$

Secondary flavour distribution for $Z \rightarrow e^+e^-$ 

Lepton definitions

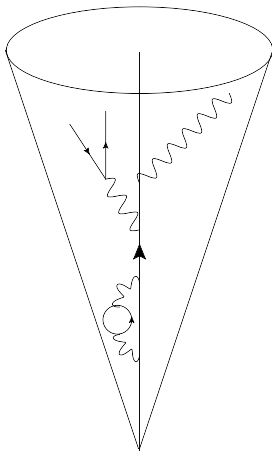
- ▶ Massive bare leptons are IR safe but experimentally difficult for electrons
- ▶ Electromagnetic calorimeter measures electrons and photons very similarly, though a magnetic field decouples their trajectories
- ▶ Usual practice: define a cone around a primary lepton, absorb all photons within its radius
- ▶ We will show this is very sensitive to higher-order corrections, especially $\gamma \rightarrow e^+ e^-$ which gives the largest mass enhancements

Lepton definitions cont'd



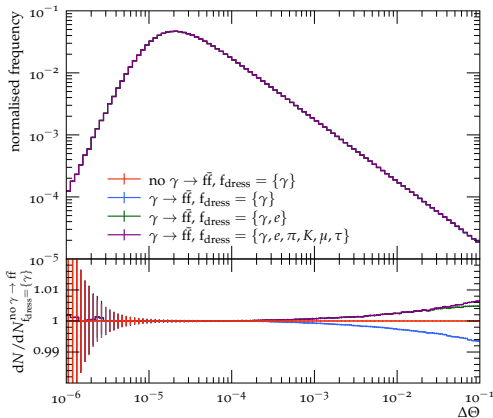
- ▶ For massless leptons, cone dressing with only photons is problematic
- ▶ Because we exclude real $\ell^+\ell^-$, there is nothing to cancel the virtual collinear singularity
- ▶ For massive leptons, there are contributions $\sim \log(m_\ell)$

Lepton definitions cont'd

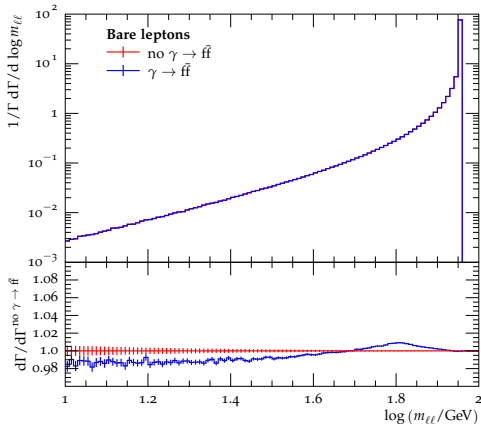


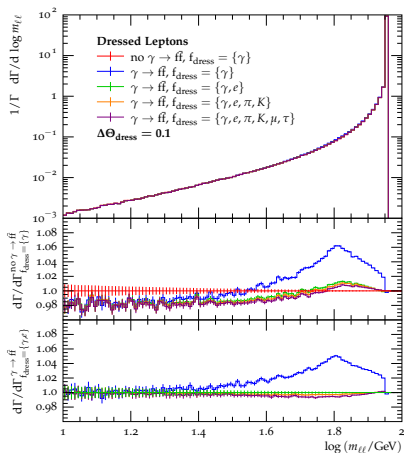
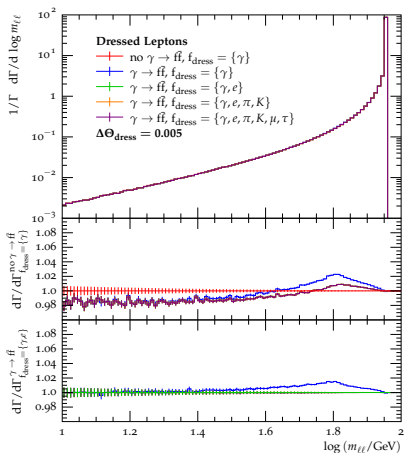
- ▶ For massless leptons, cone dressing with only photons is problematic
- ▶ Because we exclude real l^+l^- , there is nothing to cancel the virtual collinear singularity
- ▶ For massive leptons, there are contributions $\sim \log(m_\ell)$

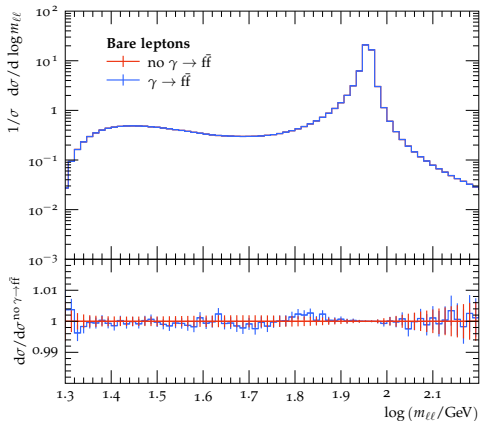
Dressed lepton structure

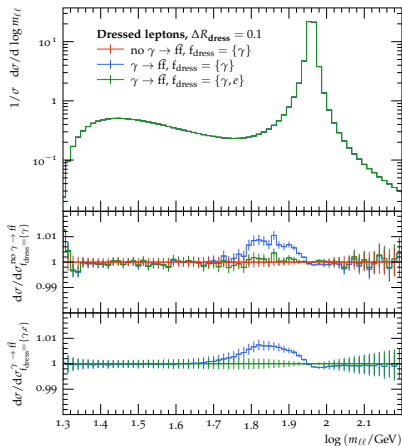
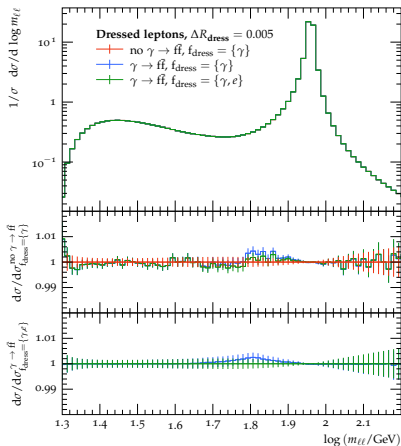


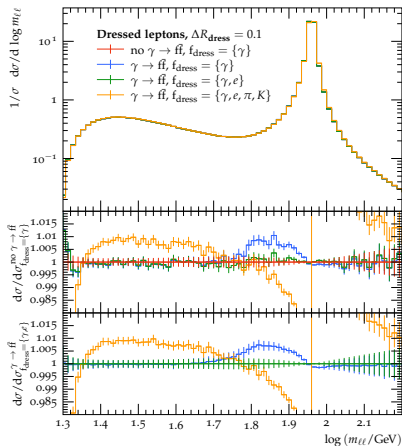
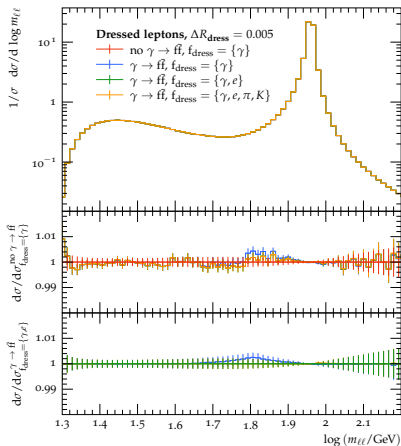
- ▶ On-shell $Z \rightarrow e^+e^-$
- ▶ For IR safety, $E_\gamma > 0.1\text{MeV}$
- ▶ At small $\Delta\Theta$, no difference in multiplicity
- ▶ At larger $\Delta\Theta$, more inclusive f_{dress} recoups loss of DL constituents due to photon splitting
- ▶ Including electrons as well as photons already reincorporates most constituents

Dilepton invariant mass for on-shell $Z \rightarrow e^+ e^-$ 

Dilepton invariant mass for on-shell $Z \rightarrow e^+e^-$ 

Dilepton invariant mass for $pp \rightarrow e^+ e^-$ (preliminary)

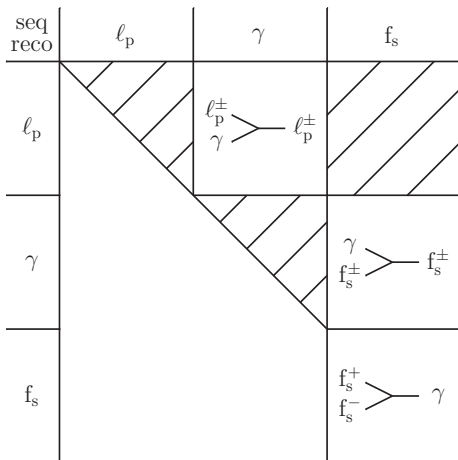
Dilepton invariant mass for $pp \rightarrow e^+e^-$ (preliminary)

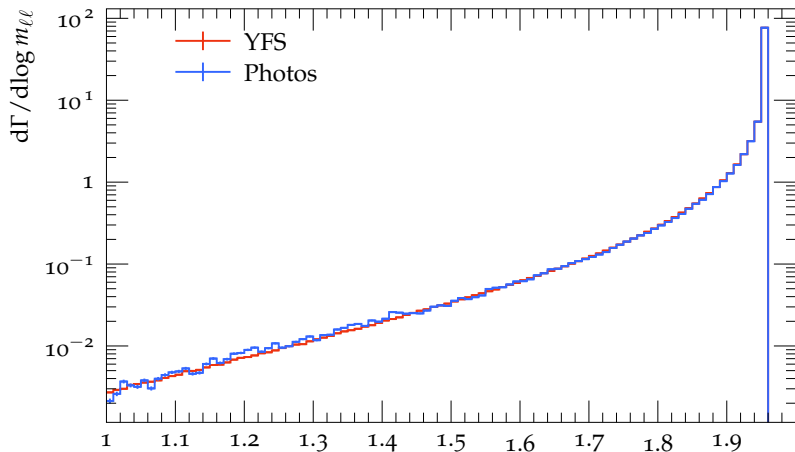
Dilepton invariant mass for $pp \rightarrow e^+e^-$ (preliminary)

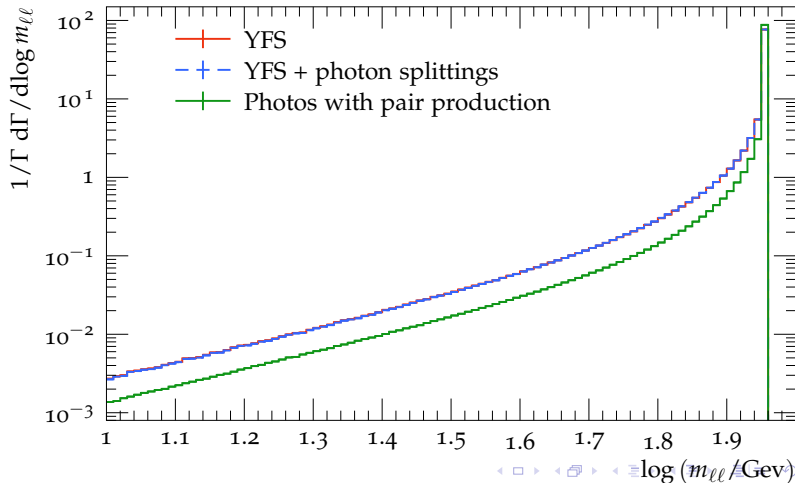
Conclusions and outlook

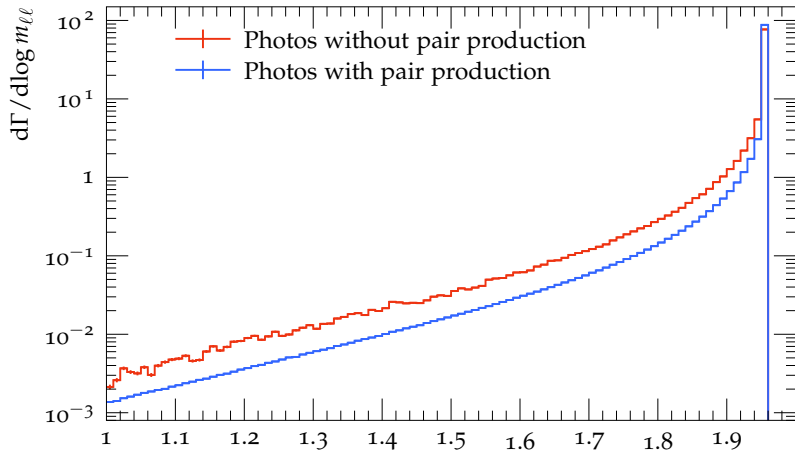
- ▶ We introduced a novel set of lepton dressing strategies which can be used theoretically in conjunction with QED corrections
- ▶ Results were presented for on-shell Z decay to leptons and Drell-Yan lepton production
- ▶ Flavour-aware dressing can be used when massless leptons are produced by photon splitting to ensure cancellation of collinear singularities
- ▶ What do we consider a reasonable definition of a dressed lepton?

Backup: Sequential recombination dressing

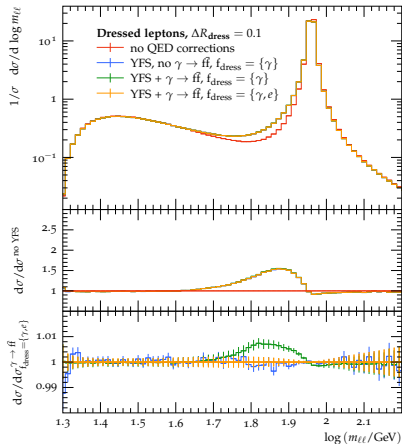


Backup: Comparison with Photos: $Z \rightarrow e^+e^-$ 

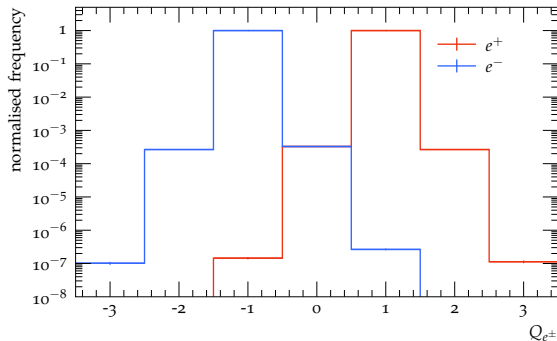
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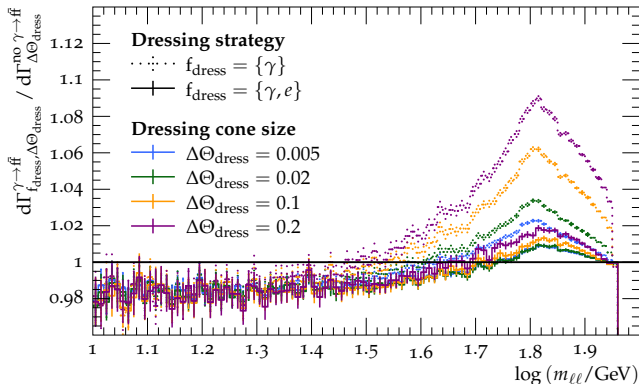
Backup: Dilepton invariant mass for $pp \rightarrow e^+e^-$



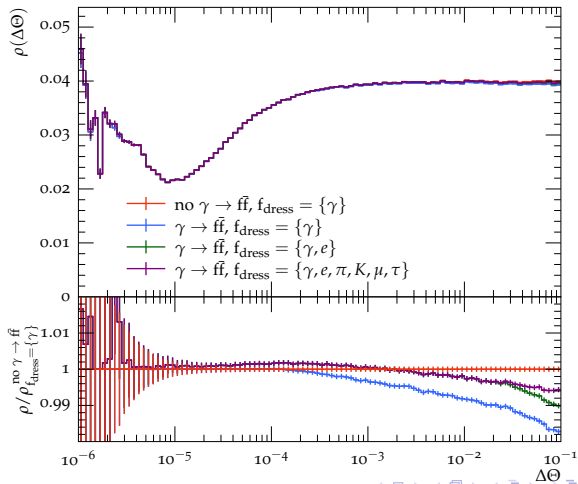
Backup: Dressed lepton charge



Backup: Dressing cone size dependence



Backup: Energy density of a dressed lepton



Backup: The YFS soft-photon resummation

$$d\Gamma^{\text{YFS}} = d\Gamma_0 \cdot e^{\alpha Y(\omega_{\text{cut}})} \cdot \sum_{n_\gamma=0}^{\infty} \frac{1}{n_\gamma!} \left[\prod_{i=1}^{n_\gamma} d\Phi_{k_i} \cdot \alpha \tilde{S}(k_i) \Theta(k_i^0 - \omega_{\text{cut}}) \cdot \mathcal{C} \right]$$

- ▶ $Y(\omega_{\text{cut}})$ is the YFS form factor containing soft-photon logarithms
- ▶ \tilde{S} is the eikonal (soft emission effects)
- ▶ \mathcal{C} corrects for hard-emission effects up to a given order
- ▶ YFS contains no description of charged particle production, $\gamma \rightarrow f\bar{f}$

Backup: Photon splitting algorithm

- ▶ One-step parton shower - subsequent emissions factorise when ordered in a scale variable t
- ▶ Input: primary charged particles and coherently emitted soft photons
- ▶ We reconstruct the scale t_{start} (GeV^2) from the input
- ▶ Then calculate probabilities for each photon to split and let all possible splittings compete

Backup: Splitting functions

$$S_{s_{\tilde{ij}}(\tilde{k}) \rightarrow s_i \gamma_j(k)} = -Q_{\tilde{ij}\tilde{k}}^2 \alpha \left[\frac{2}{1-z+zy} - \frac{\tilde{V}_{\tilde{ij},\tilde{k}}}{v_{ij,k}} \left(2 + \frac{m_i^2}{p_i p_j} \right) \right]$$

$$S_{f_{\tilde{ij}}(\tilde{k}) \rightarrow f_i \gamma_j(k)} = -Q_{\tilde{ij}\tilde{k}}^2 \alpha \left[\frac{2}{1-z+zy} - \frac{\tilde{V}_{\tilde{ij},\tilde{k}}}{v_{ij,k}} \left(1 + z + \frac{m_i^2}{p_i p_j} \right) \right]$$

$$S_{\gamma_{\tilde{ij}}(\tilde{k}) \rightarrow s_i \bar{s}_j(k)} = S_{\gamma_{\tilde{ij}}(\tilde{k}) \rightarrow f_i \bar{f}_j(k)} = -Q_{\tilde{ij}\tilde{k}}^2 \alpha \left[1 - 2z(1-z) - z_+ z_- \right]$$

Catani et al. '02, Dittmaier et al. '08, Schumann, Krauss '08

Backup: Leptonic W decay

- ▶ The W is charged and the neutrino is not, so instead of an FF dipole we have an FI dipole $W - \ell$
- ▶ Large W mass suppresses photon emissions, so neglect it as an emitter
- ▶ Modify kinematic variables and splitting functions – we keep the W eikonal term

Basso et al. '16