

Global View on SMEFT interpretations and UV connection

Maeve Madigan

Heidelberg University



UNIVERSITÄT
HEIDELBERG
ZUKUNFT
SEIT 1386

LHC Reinterpretation Forum 31.08.23

Global SMEFT interpretations

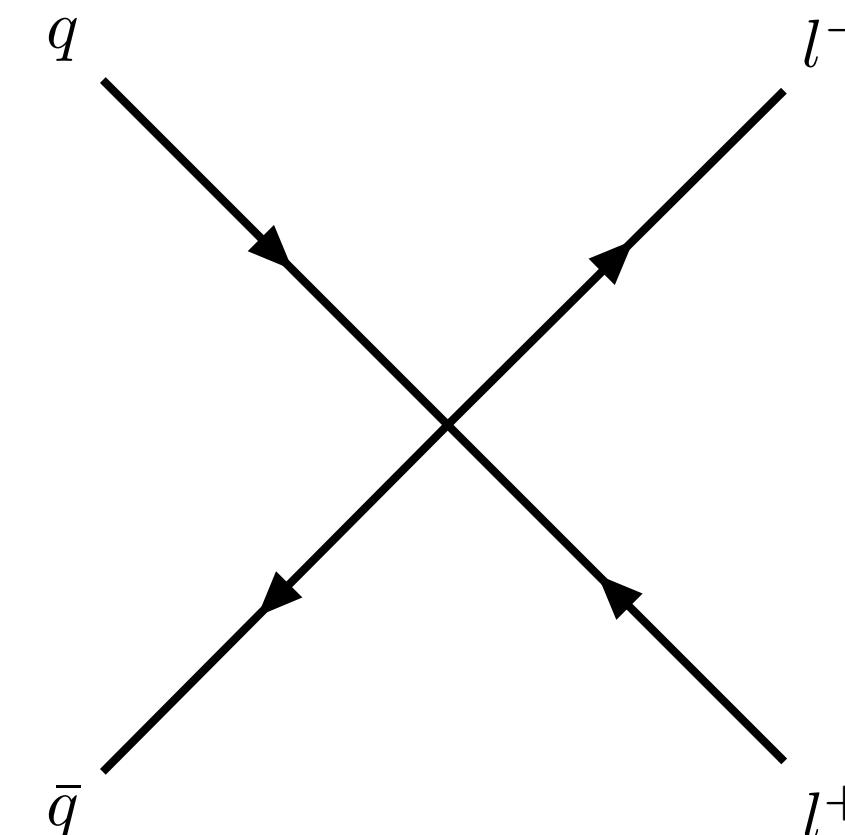
The SMEFT: a powerful framework for capturing deviations from the SM:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{C^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$

Global SMEFT interpretations

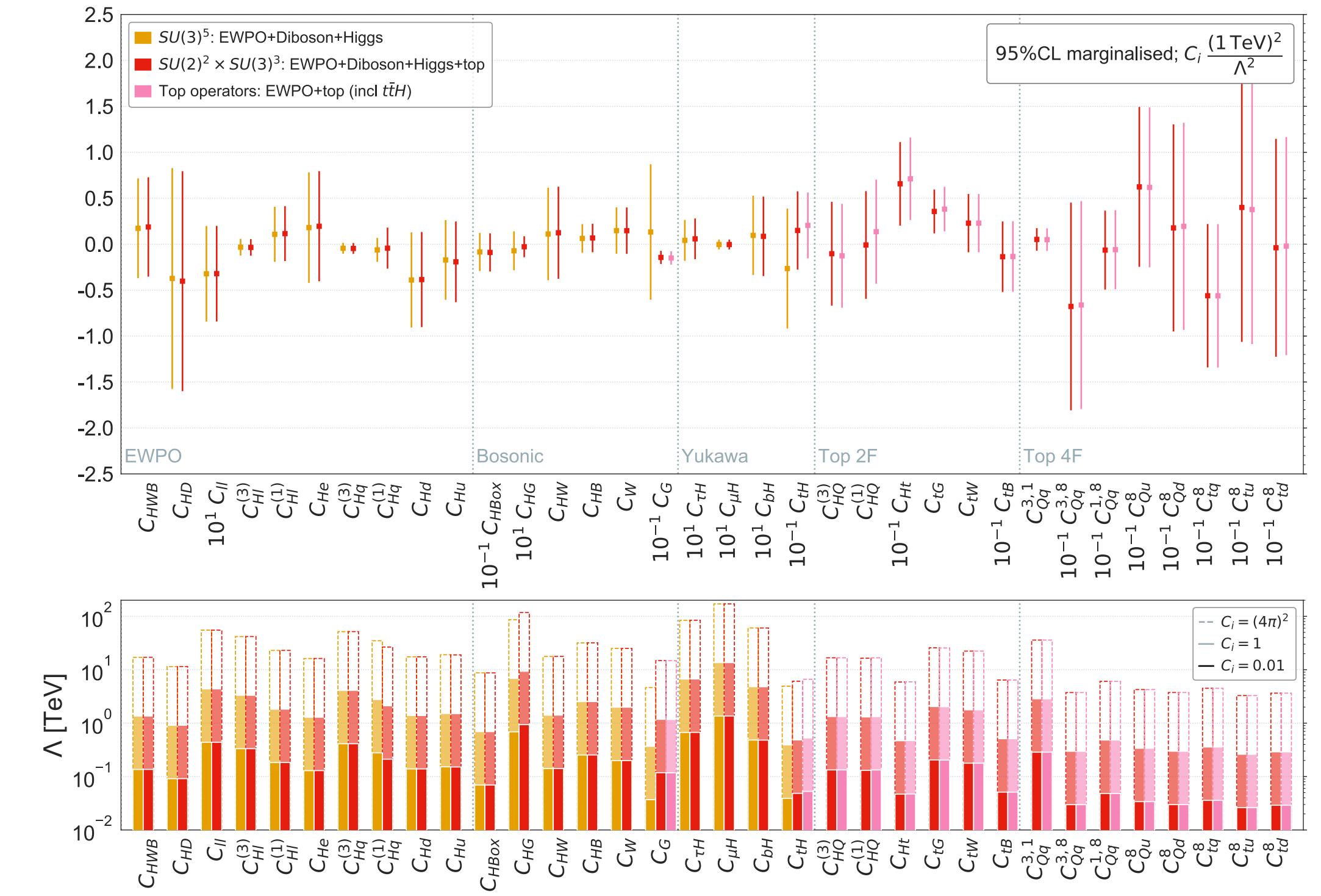
The SMEFT: a powerful framework for capturing deviations from the SM:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{C^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$



Talks by Jaco ter Hoeve, Danny van Dyk, Kirill Skovpen, Rahul Balasubramanian, ...

$$\Lambda \downarrow E$$

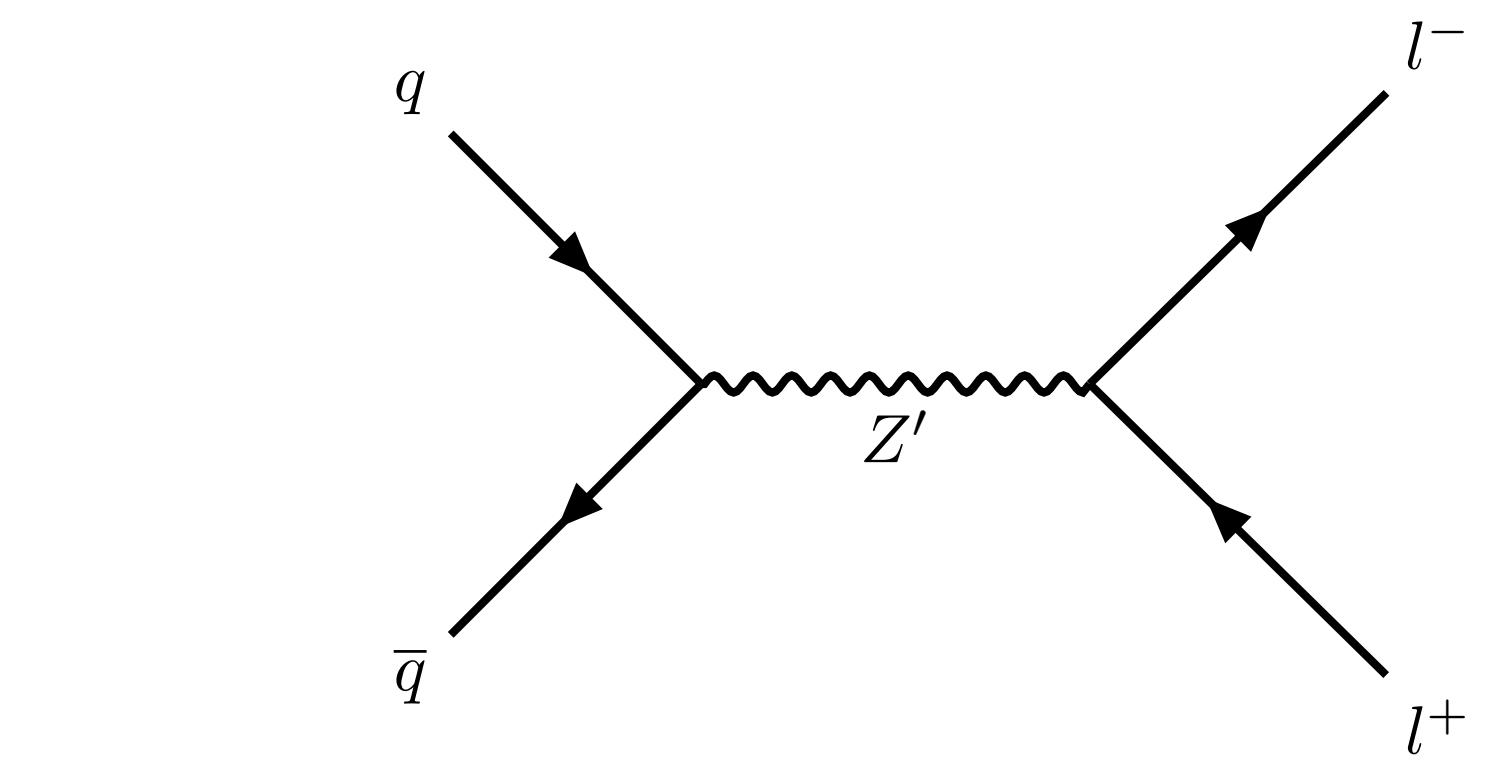


2012.02779, J. Ellis, MM, K. Mimasu, V. Sanz, T. You

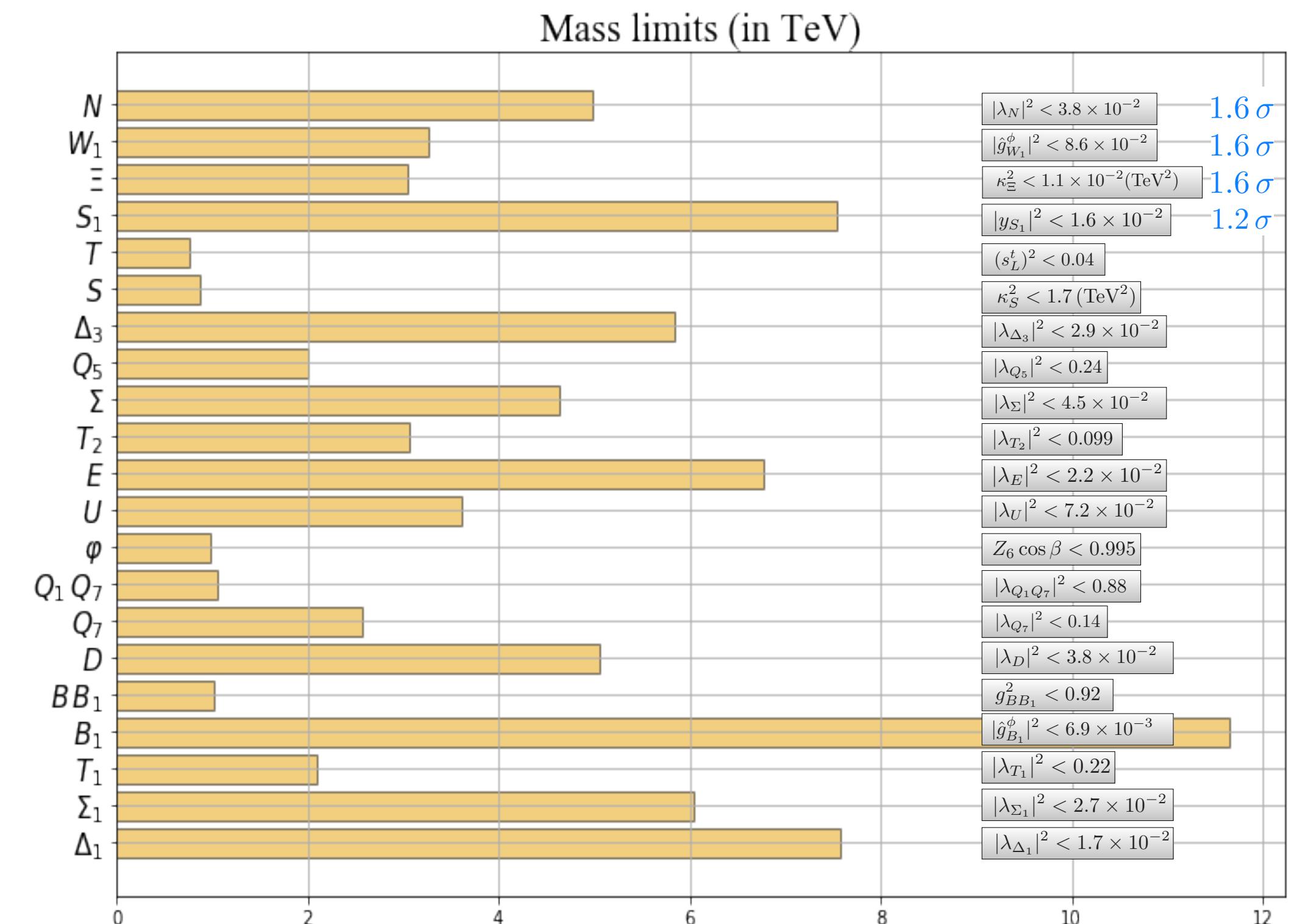
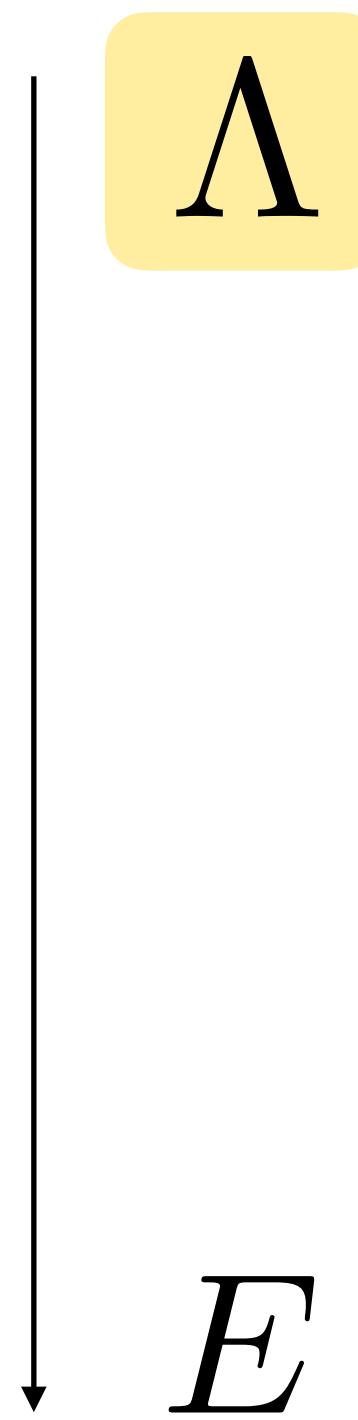
Global SMEFT interpretations

The SMEFT: a powerful framework for capturing deviations from the SM:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{C^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$

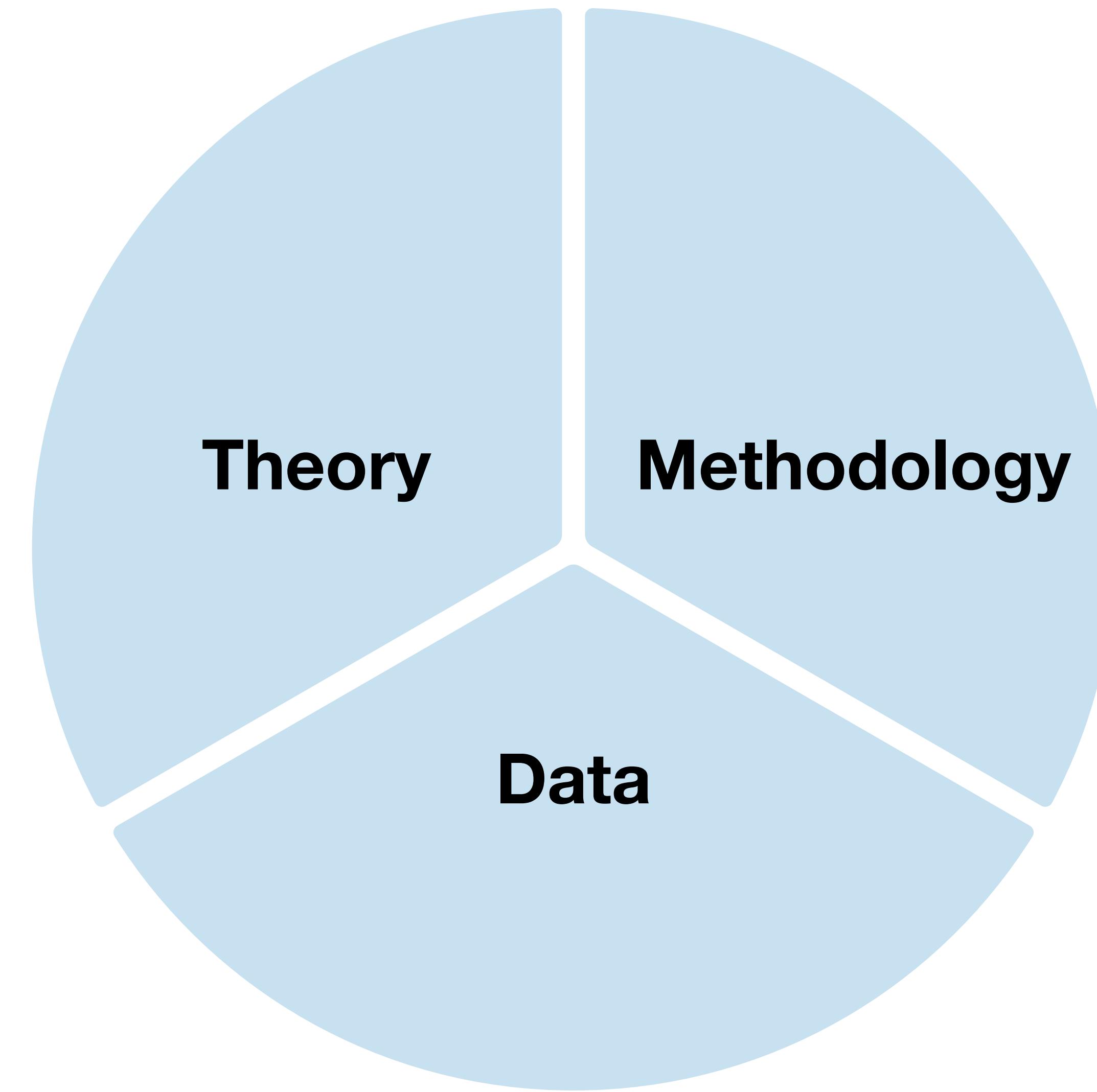


Talk by Shankha Banerjee, Jaco ter Hoeve, ...

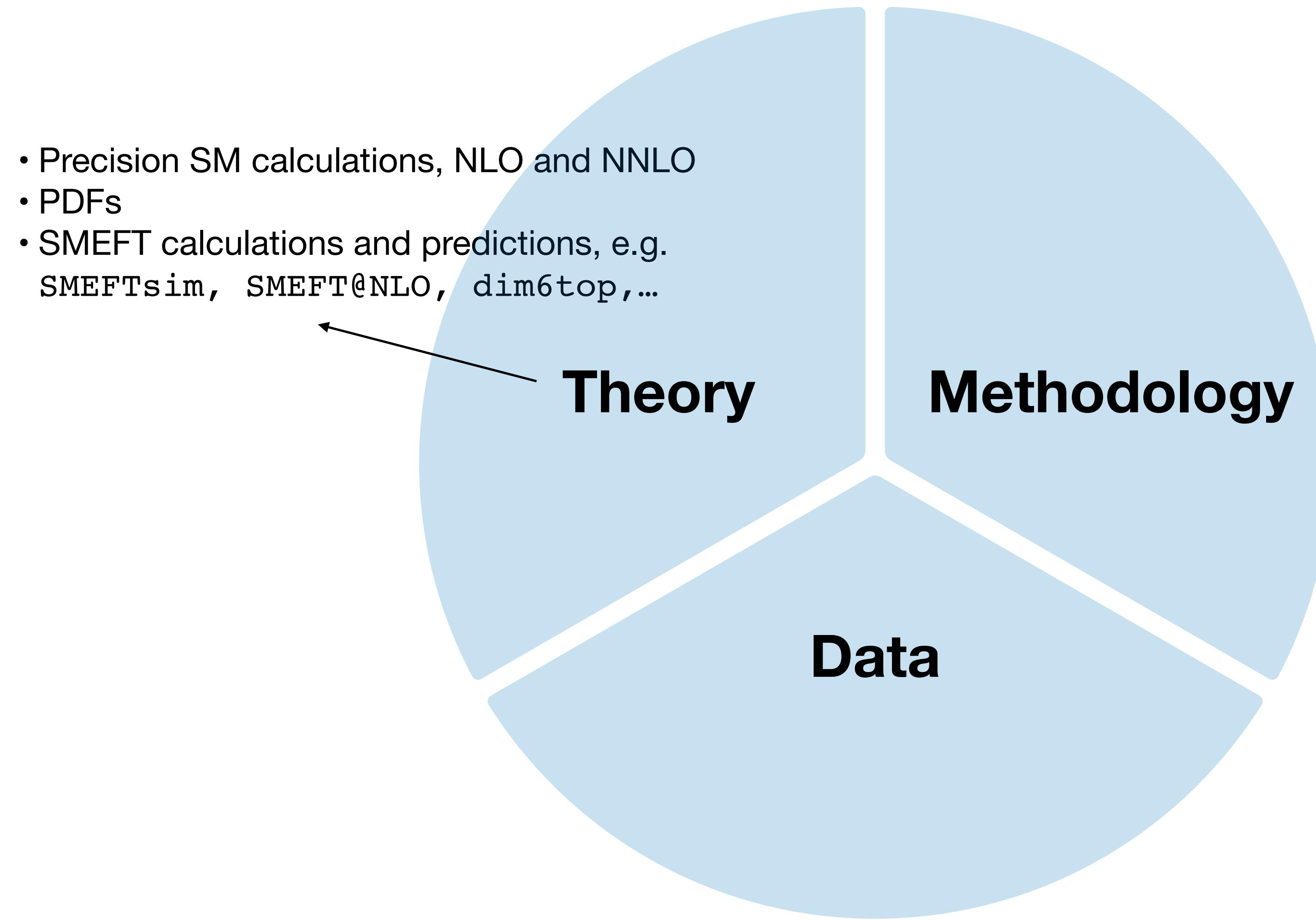


2012.02779, J. Ellis, MM, K. Mimasu, V. Sanz, T. You

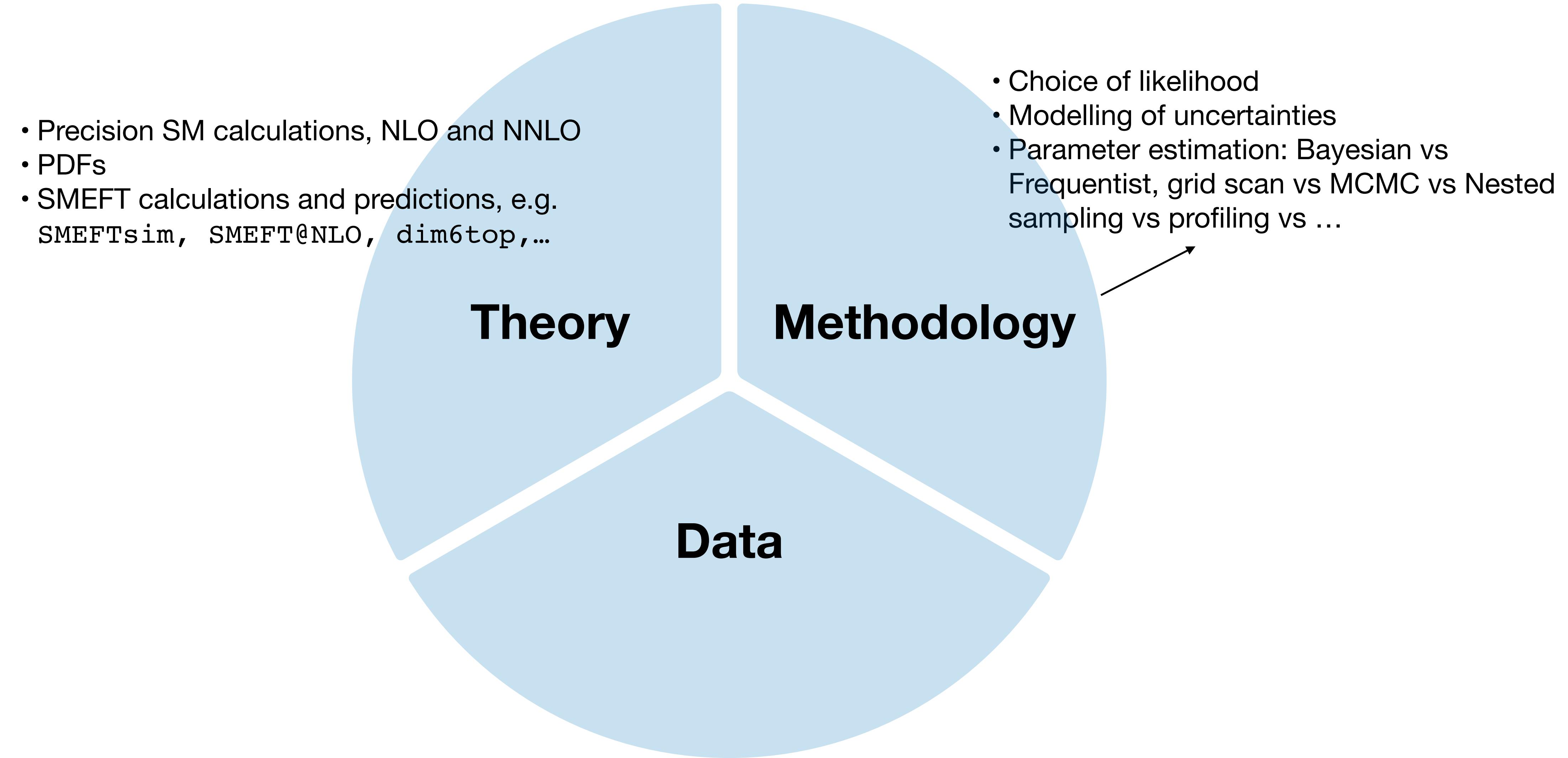
Global SMEFT interpretations



Global SMEFT interpretations



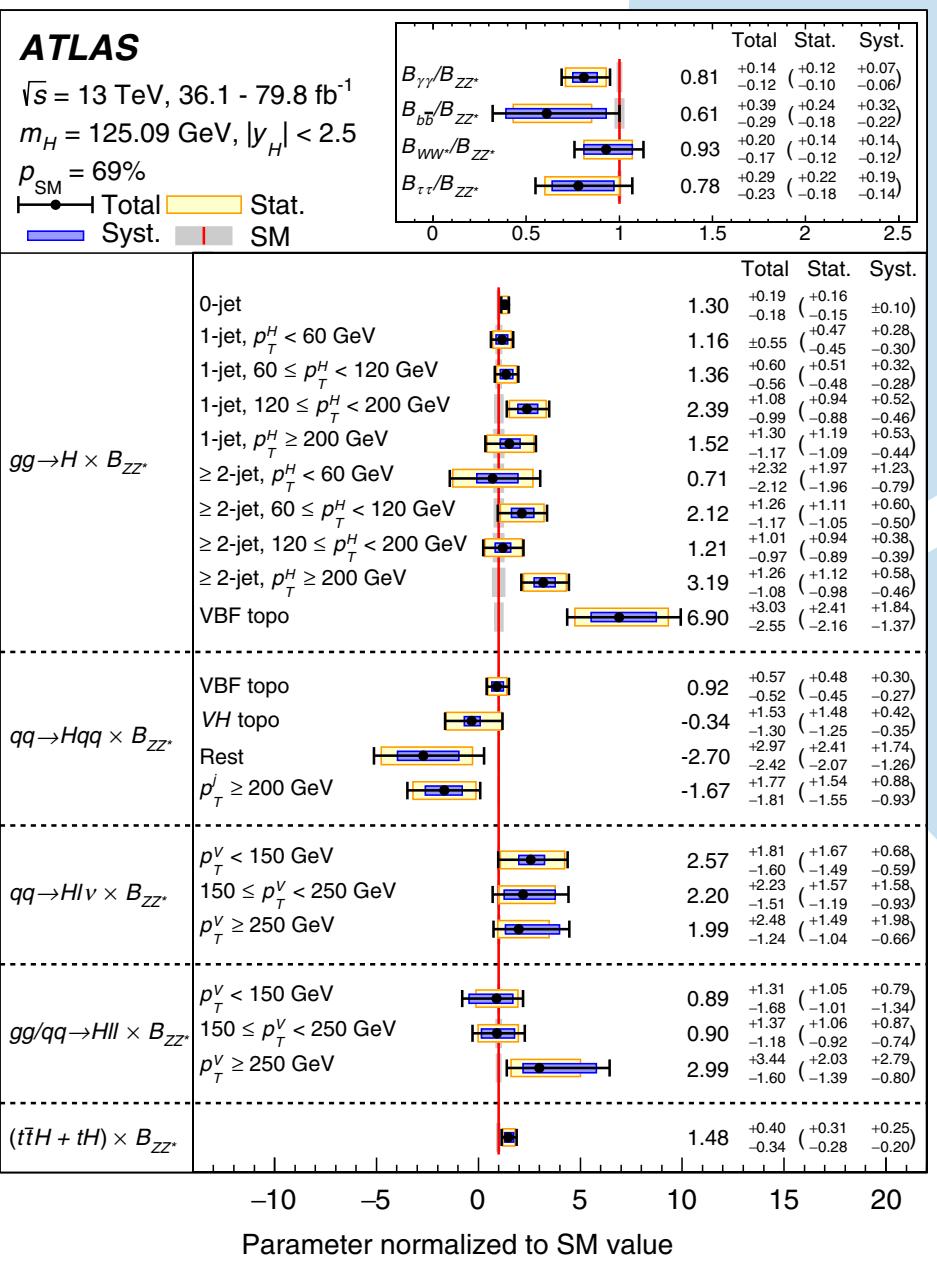
Global SMEFT interpretations



Global SMEFT interpretations

- Precision SM calculations, NLO and NNLO
- PDFs
- SMEFT calculations and predictions, e.g. SMEFTsim, SMEFT@NLO, dim6top, ...

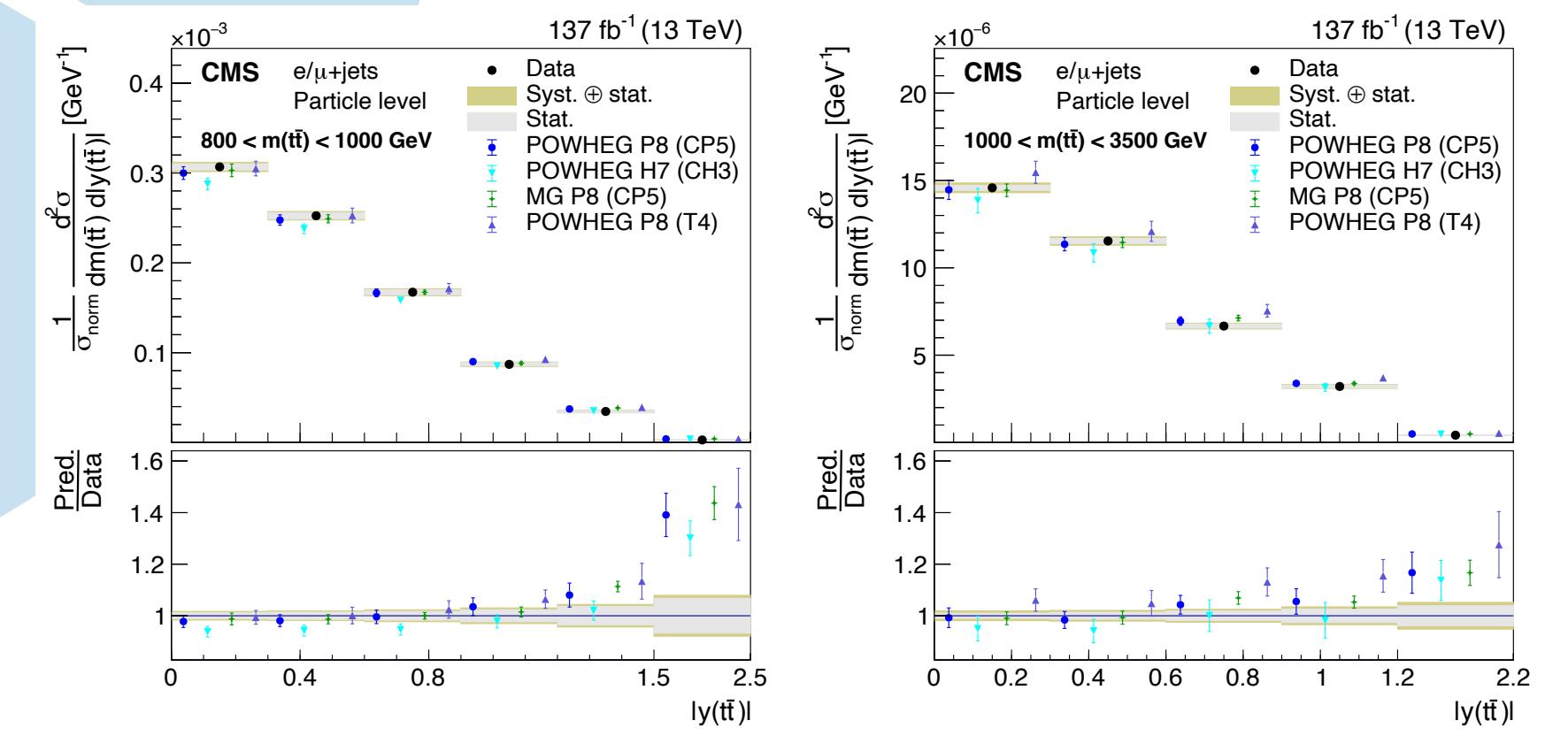
Theory



Data

- Choice of likelihood
- Modelling of uncertainties
- Parameter estimation: Bayesian vs Frequentist, grid scan vs MCMC vs Nested sampling vs profiling vs ...

Methodology



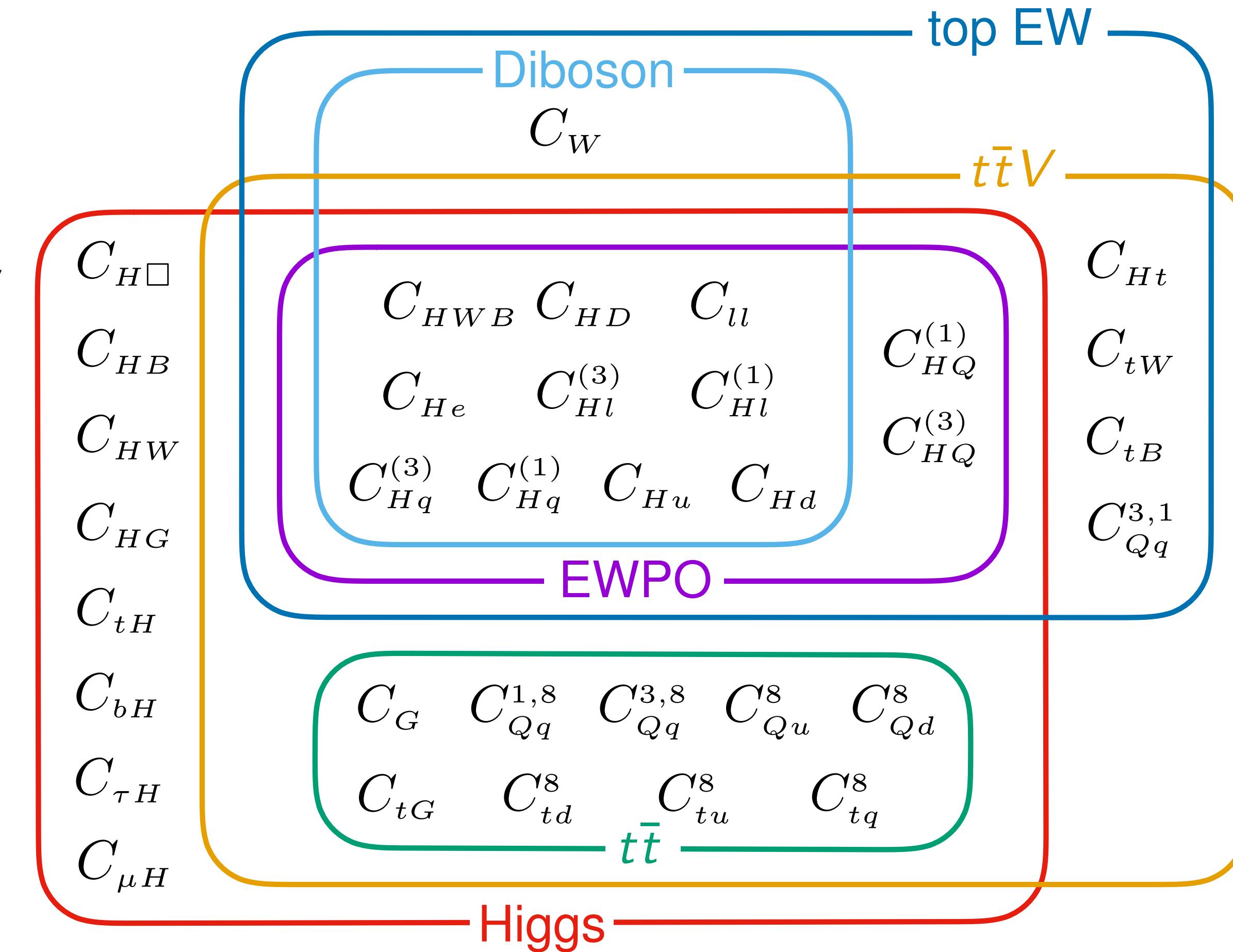
[ATLAS, Phys. Rev. D 101 (2020) 012002]

[CMS, Phys. Phys. Rev. D 104 (2021) 092013]

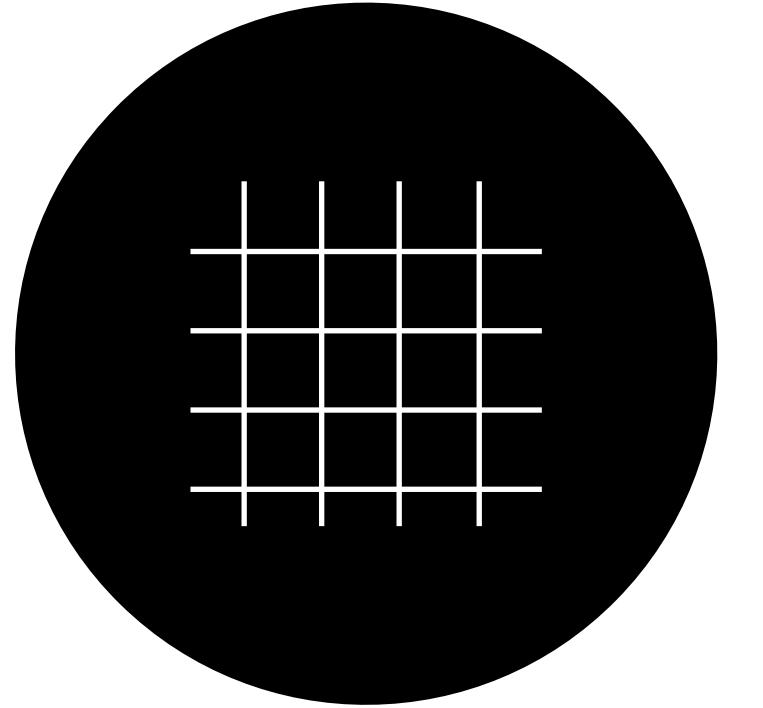
Global SMEFT interpretations: Data

Global: as much data constraining as many processes as possible

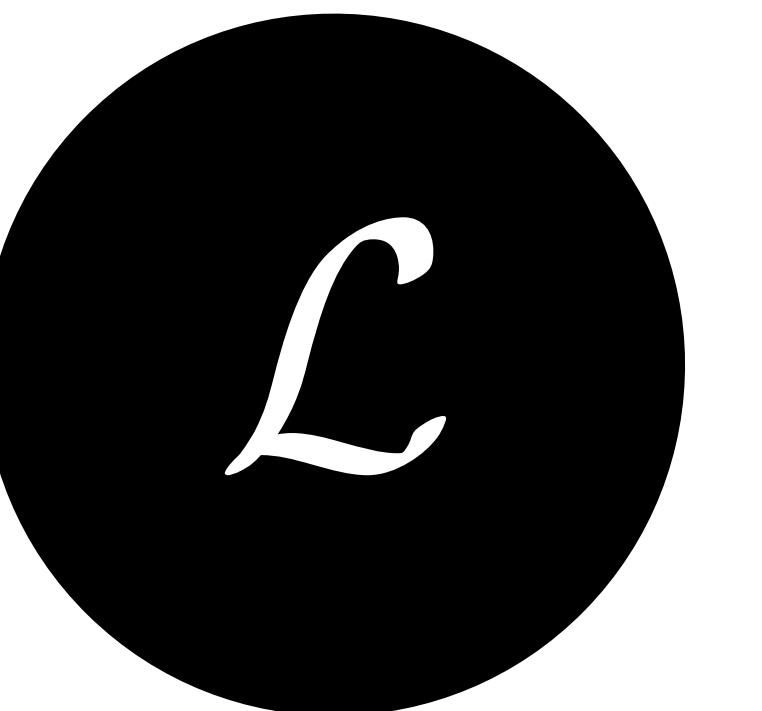
Data: the use of data in SMEFT fits depends on *how the data is presented*



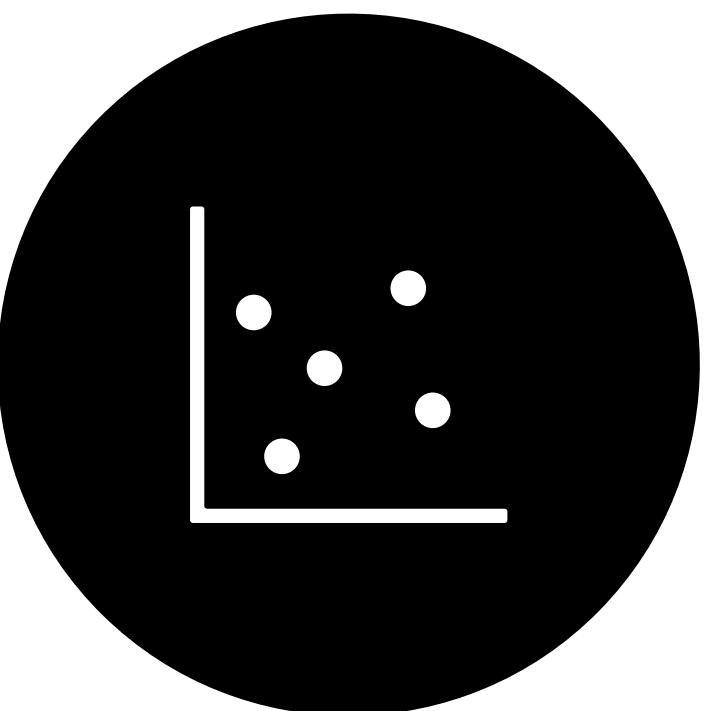
2012.02779, J. Ellis, MM, K. Mimasu, V. Sanz, T. You



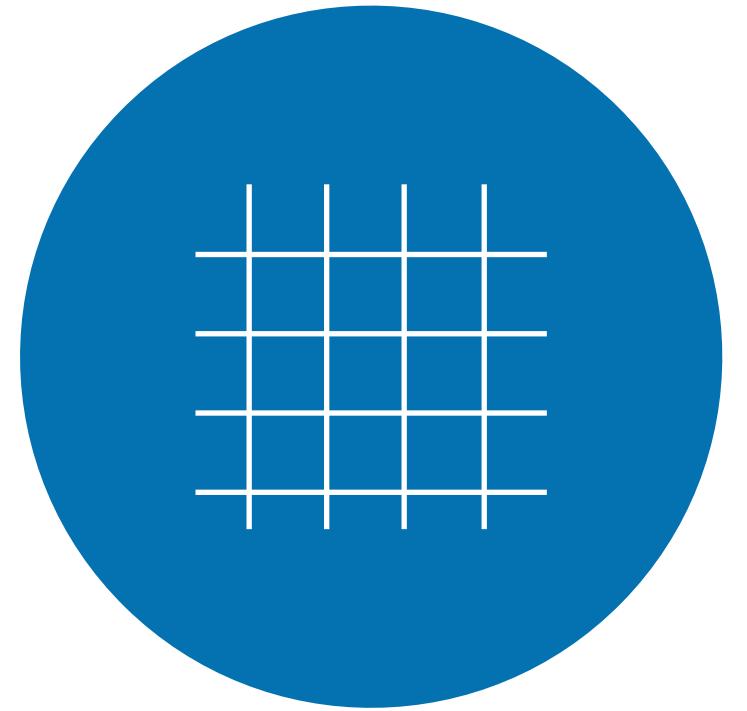
Datapoints & covariance
matrices



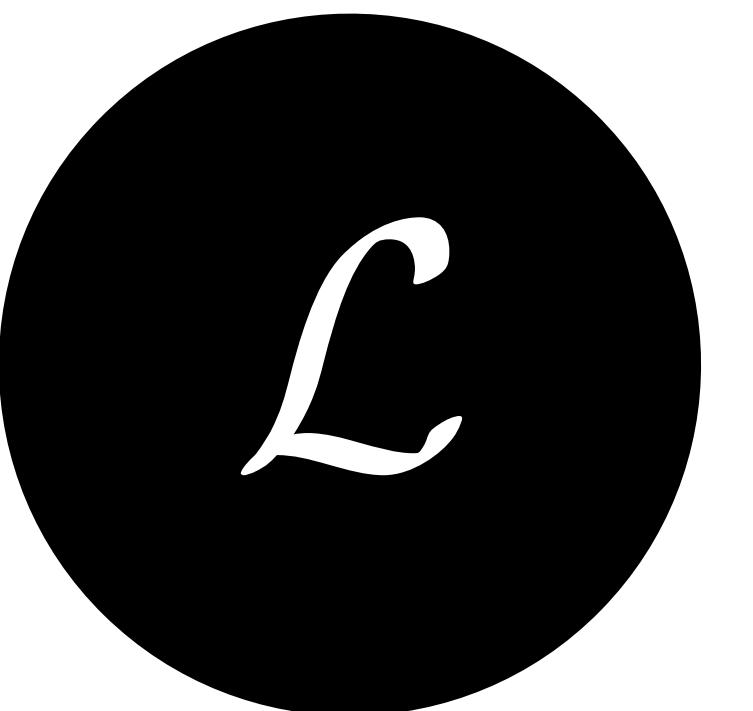
Likelihoods



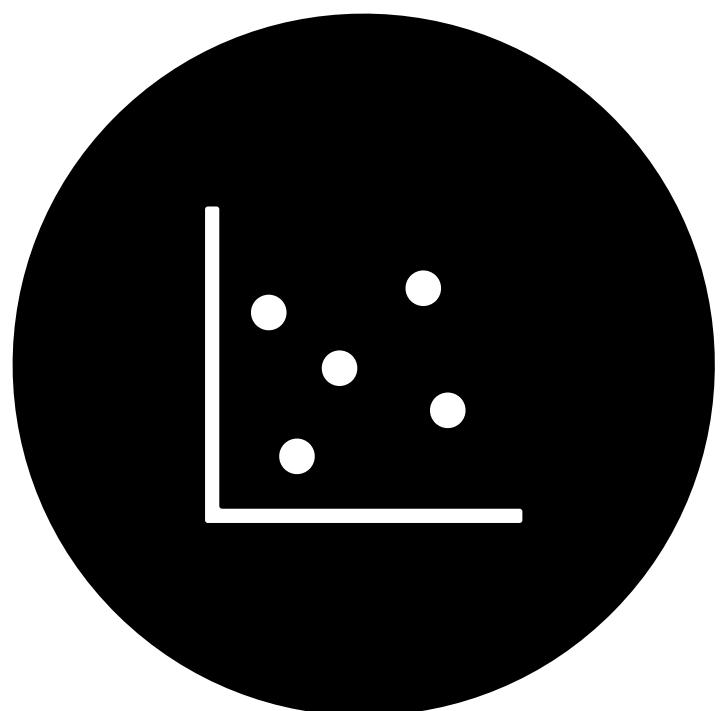
Unbinned
measurements



Datapoints & covariance
matrices



Likelihoods



Unbinned
measurements

Covariance matrices

Access to **datapoints** and their **correlated uncertainties (statistical and systematic)** allows us to calculate the Gaussian likelihood:

$$\mathcal{L}(c) \propto \exp\left(-\frac{1}{2}(D - T(c))^T V^{-1} (D - T(c))\right)$$

Covariance matrices

Access to **datapoints** and their **correlated uncertainties (statistical and systematic)** allows us to calculate the Gaussian likelihood:

$$\mathcal{L}(c) \propto \exp\left(-\frac{1}{2}(D - T(c))^T V^{-1} (D - T(c))\right)$$

Datapoints:

$p_T(t_h)$ [GeV]	$\frac{1}{\sigma_{\text{norm}}} \frac{d\sigma}{dp_T(t_h)}$ [GeV $^{-1}$]
0.0 - 40.0	0.002762 $\pm 1.883\text{e-}05$ stat $\pm 4.883\text{e-}05$ sys
40.0 - 80.0	0.006028 $\pm 2.874\text{e-}05$ stat $\pm 5.586\text{e-}05$ sys
80.0 - 120.0	0.005937 $\pm 2.685\text{e-}05$ stat $\pm 3.092\text{e-}05$ sys
120.0 - 160.0	0.004291 $\pm 2.126\text{e-}05$ stat $\pm 3.18\text{e-}05$ sys
160.0 - 200.0	0.002639 $\pm 1.502\text{e-}05$ stat $\pm 2.198\text{e-}05$ sys
200.0 - 250.0	0.001376 $\pm 7.9\text{e-}06$ stat $\pm 1.562\text{e-}05$ sys
250.0 - 300.0	0.0006568 $\pm 5.249\text{e-}06$ stat $\pm 9.987\text{e-}06$ sys
300.0 - 350.0	0.000311 $\pm 3.904\text{e-}06$ stat $\pm 6.535\text{e-}06$ sys
350.0 - 400.0	0.0001562 $\pm 3.109\text{e-}06$ stat $\pm 3.774\text{e-}06$ sys
400.0 - 450.0	7.761e-05 $\pm 2.569\text{e-}06$ stat $\pm 2.633\text{e-}06$ sys
450.0 - 500.0	4.085e-05 $\pm 2.015\text{e-}06$ stat $\pm 1.772\text{e-}06$ sys
500.0 - 600.0	1.791e-05 $\pm 6.153\text{e-}07$ stat $\pm 7.734\text{e-}07$ sys

CMS Phys.Rev.D 104 (2021) 092013, 2021.
<https://www.hepdata.net/record/ins1901295>

Availability of data on HEPData is growing 😊



Covariance matrices

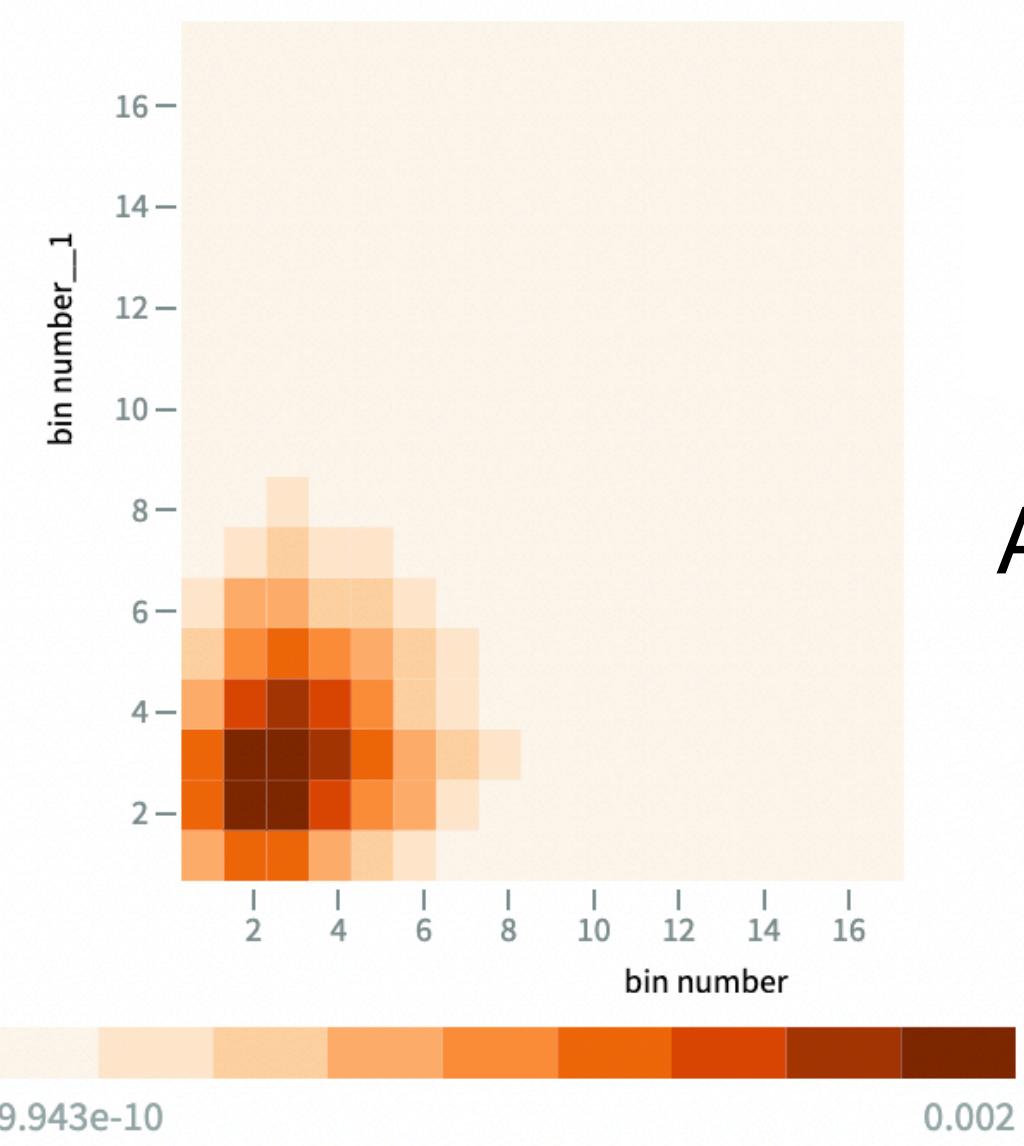
Access to **datapoints** and their **correlated uncertainties (statistical and systematic)** allows us to calculate the Gaussian likelihood:

$$\mathcal{L}(c) \propto \exp\left(-\frac{1}{2}(D - T(c))^T V^{-1} (D - T(c))\right)$$

Datapoints:

$p_T(t_h)$ [GeV]	$\frac{1}{\sigma_{\text{norm}}} \frac{d\sigma}{dp_T(t_h)}$ [GeV $^{-1}$]
0.0 - 40.0	0.002762 $\pm 1.883\text{e-}05$ stat $\pm 4.883\text{e-}05$ sys
40.0 - 80.0	0.006028 $\pm 2.874\text{e-}05$ stat $\pm 5.586\text{e-}05$ sys
80.0 - 120.0	0.005937 $\pm 2.685\text{e-}05$ stat $\pm 3.092\text{e-}05$ sys
120.0 - 160.0	0.004291 $\pm 2.126\text{e-}05$ stat $\pm 3.18\text{e-}05$ sys
160.0 - 200.0	0.002639 $\pm 1.502\text{e-}05$ stat $\pm 2.198\text{e-}05$ sys
200.0 - 250.0	0.001376 $\pm 7.9\text{e-}06$ stat $\pm 1.562\text{e-}05$ sys
250.0 - 300.0	0.0006568 $\pm 5.249\text{e-}06$ stat $\pm 9.987\text{e-}06$ sys
300.0 - 350.0	0.000311 $\pm 3.904\text{e-}06$ stat $\pm 6.535\text{e-}06$ sys
350.0 - 400.0	0.0001562 $\pm 3.109\text{e-}06$ stat $\pm 3.774\text{e-}06$ sys
400.0 - 450.0	7.761e-05 $\pm 2.569\text{e-}06$ stat $\pm 2.633\text{e-}06$ sys
450.0 - 500.0	4.085e-05 $\pm 2.015\text{e-}06$ stat $\pm 1.772\text{e-}06$ sys
500.0 - 600.0	1.791e-05 $\pm 6.153\text{e-}07$ stat $\pm 7.734\text{e-}07$ sys

Covariance matrix:



CMS Phys.Rev.D 104 (2021) 092013, 2021.

<https://www.hepdata.net/record/ins1901295>

Availability of data on HEPData is growing 😊



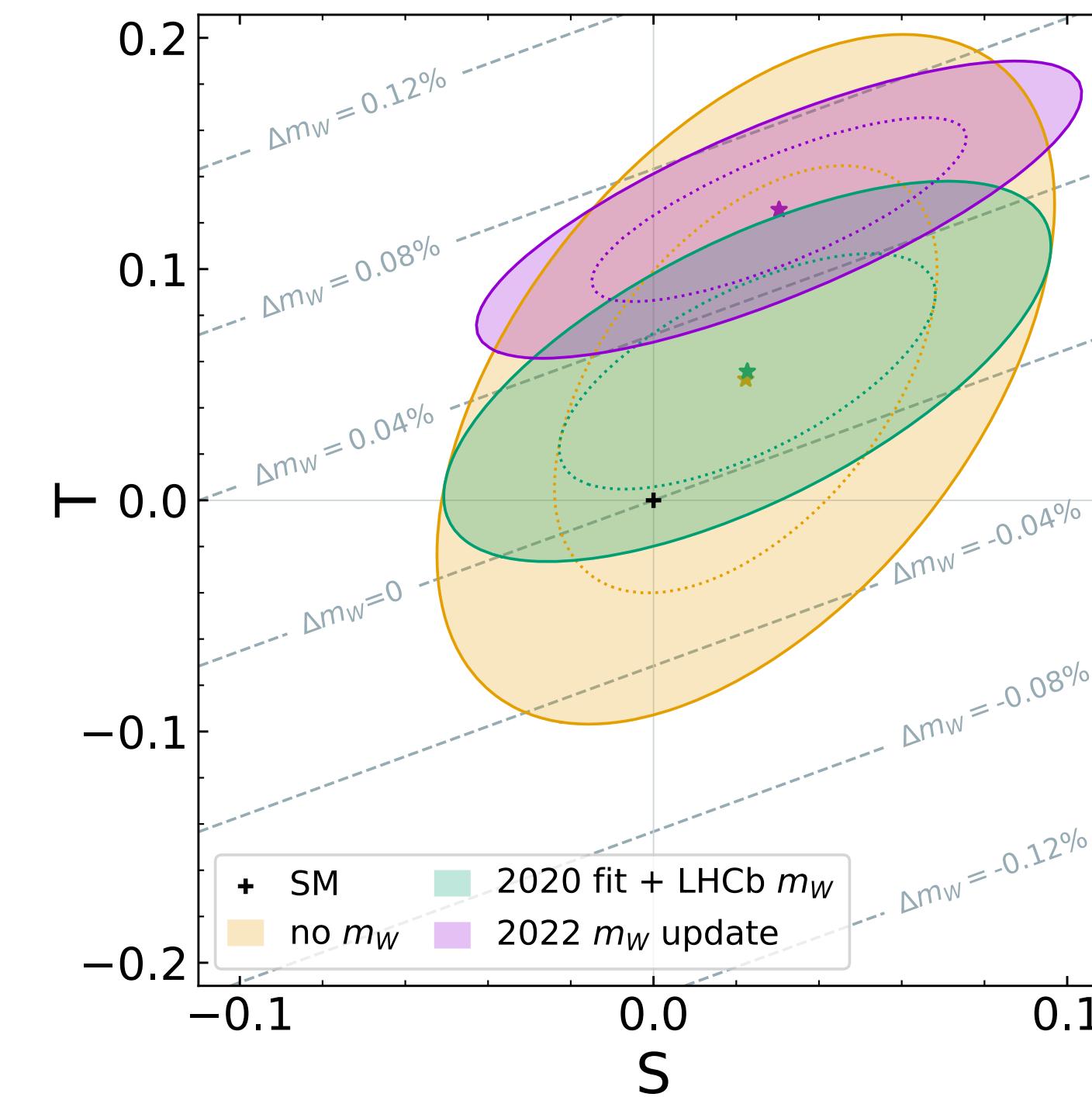
Covariance matrices

Access to **datapoints** and their **correlated uncertainties (statistical and systematic)** allows us to calculate the Gaussian likelihood:

$$\mathcal{L}(c) \propto \exp\left(-\frac{1}{2}(D - T(c))^T V^{-1} (D - T(c))\right)$$

Used in many fits, e.g.

- Fitmaker —————→
- SMEFIT - *see Jaco's talk this morning*
- PDF fits - *e.g. NNPDF, see also Zahari's talk*
- + *many others not mentioned here*



2204.05260, E. Bagnaschi, J. Ellis, MM, K. Mimasu, V. Sanz, T. You

Covariance matrices

Access to **datapoints** and their **correlated uncertainties (statistical and systematic)** allows us to calculate the Gaussian likelihood:

$$\mathcal{L}(c) \propto \exp\left(-\frac{1}{2}(D - T(c))^T V^{-1} (D - T(c))\right)$$

Be aware of assumptions, including:

- Relies on the assumption of symmetric uncertainties, see e.g. [Lilith, 1908.03952 Kraml et. al](#)
- Covariance matrices do not capture correlations **between different measurements** (more on this later)

Covariance matrices for double-differential distributions

Global SMEFT fits can benefit from **double differential distributions**

- more information  better constraints

Correlations between distributions are necessary

Covariance matrices for double-differential distributions

Global SMEFT fits can benefit from **double differential distributions**

- more information  better constraints

Correlations between distributions are necessary

Ideal for simultaneous **PDF & SMEFT** determinations;

- optimal choice of distribution may be different for PDFs and SMEFT

Covariance matrices for double-differential distributions

Global SMEFT fits can benefit from **double differential distributions**

- more information  better constraints

Correlations between distributions are necessary

Ideal for simultaneous **PDF & SMEFT** determinations;

- optimal choice of distribution may be different for PDFs and SMEFT

In arxiv 2303.06159, PDF & EFT fit in the top sector

see also Zahari's talk today

- ATLAS 2006.09274 13 TeV top pair production in the all-hadronic channel
- CMS 1703.01630 8 TeV top pair production in the dilepton channel

$$\frac{d^2\sigma}{dy_{t\bar{t}}dm_{t\bar{t}}}$$

Availability of covariance matrix made it possible to include both rapidity and top pair invariant mass

Covariance matrices and unfolded distributions



- Unfolding detector effects
- Measurements unfolded to the full phase space
- Measurements unfolded to e.g. stable top quarks

}

Efficient calculation of SMEFT and SM predictions for global SMEFT interpretations of O(100) datapoints

Covariance matrices and unfolded distributions

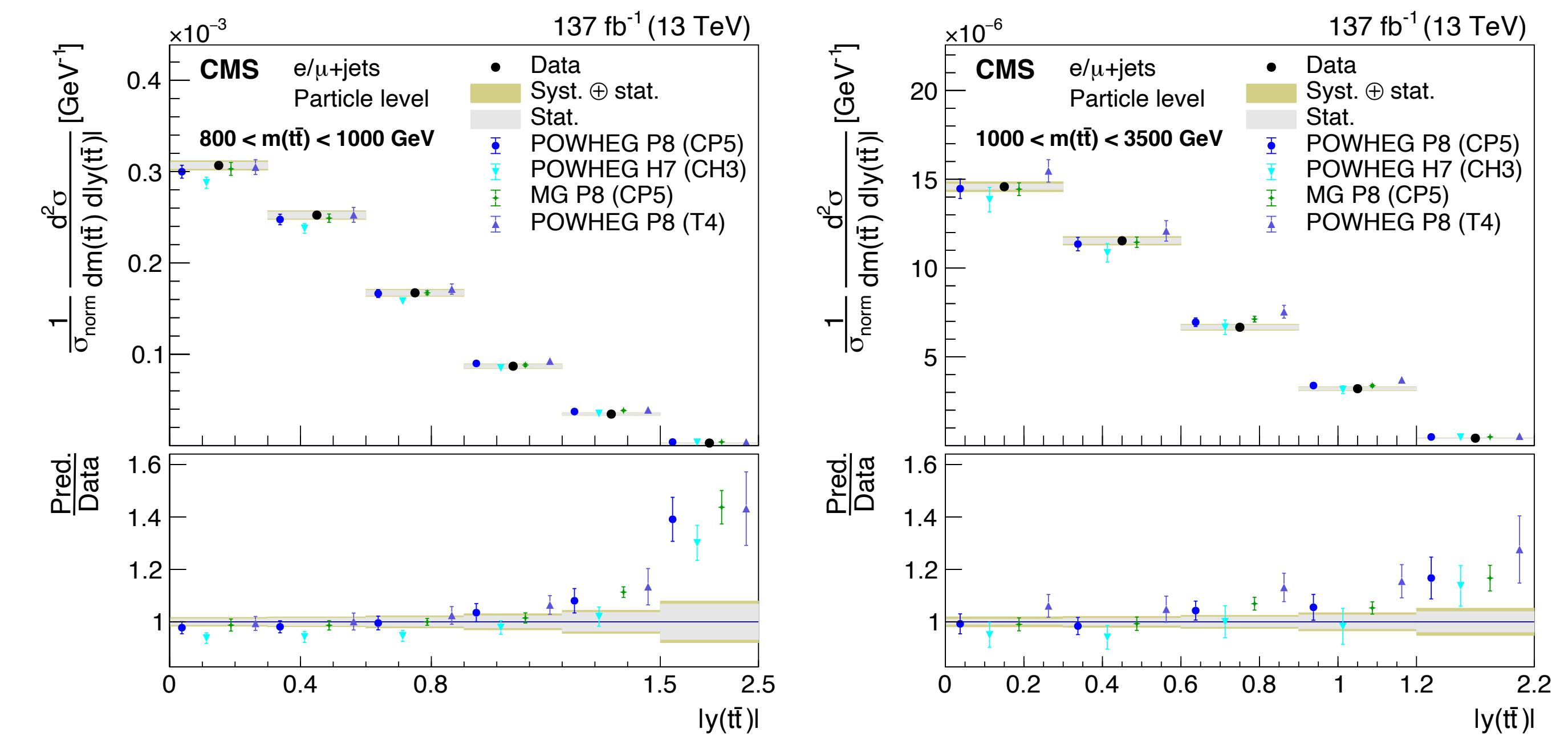


- Unfolding detector effects
- Measurements unfolded to the full phase space
- Measurements unfolded to e.g. stable top quarks



Efficient calculation of SMEFT and SM predictions for global SMEFT interpretations of O(100) datapoints

Many examples already exist, e.g. **in the top sector**: [\[CMS, Phys. Phys. Rev. D 104 \(2021\) 092013\]](#)



Covariance matrices and unfolded distributions



- Unfolding detector effects
- Measurements unfolded to the full phase space
- Measurements unfolded to e.g. stable top quarks



Efficient calculation of SMEFT and SM predictions for global SMEFT interpretations of O(100) datapoints

However: unfolding relies on SM assumptions

Covariance matrices and unfolded distributions



- Unfolding detector effects
- Measurements unfolded to the full phase space
- Measurements unfolded to e.g. stable top quarks

} Efficient calculation of SMEFT and SM predictions for global SMEFT interpretations of O(100) datapoints

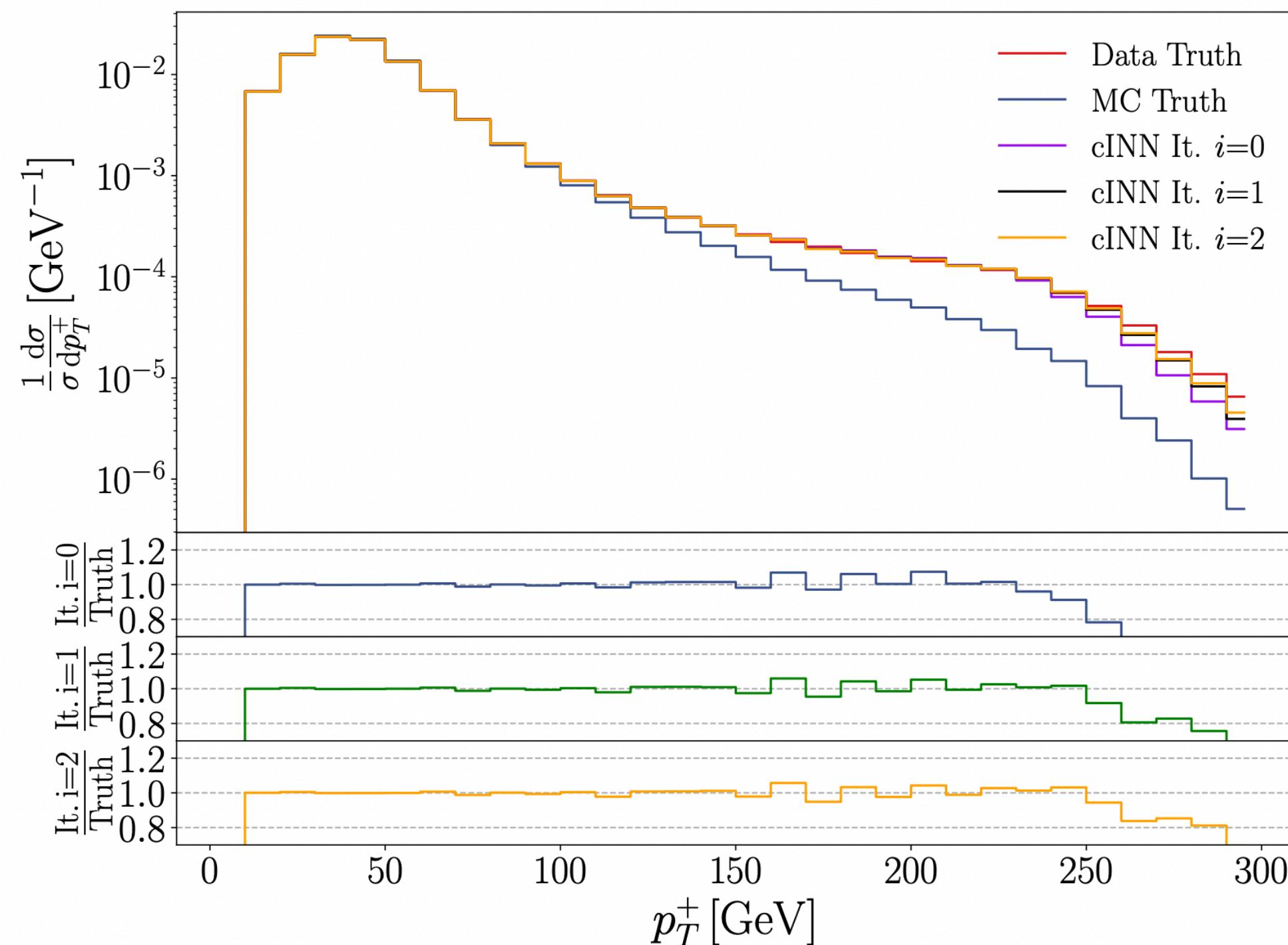
However: unfolding relies on SM assumptions

e.g. *Butter et. al, 2212.08674*

Unfolding using conditional invertible neural networks

- iteratively **remove model bias**
- dimension-8 SMEFT in $pp \rightarrow Z\gamma\gamma$

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{C_{T,8}}{\Lambda^4} B_{\mu\nu} B^{\mu\nu} B_{\alpha\beta} B^{\alpha\beta}$$



Covariance matrices and unfolded distributions



- Unfolding detector effects
- Measurements unfolded to the full phase space
- Measurements unfolded to e.g. stable top quarks

} Efficient calculation of SMEFT and SM predictions for global SMEFT interpretations of O(100) datapoints

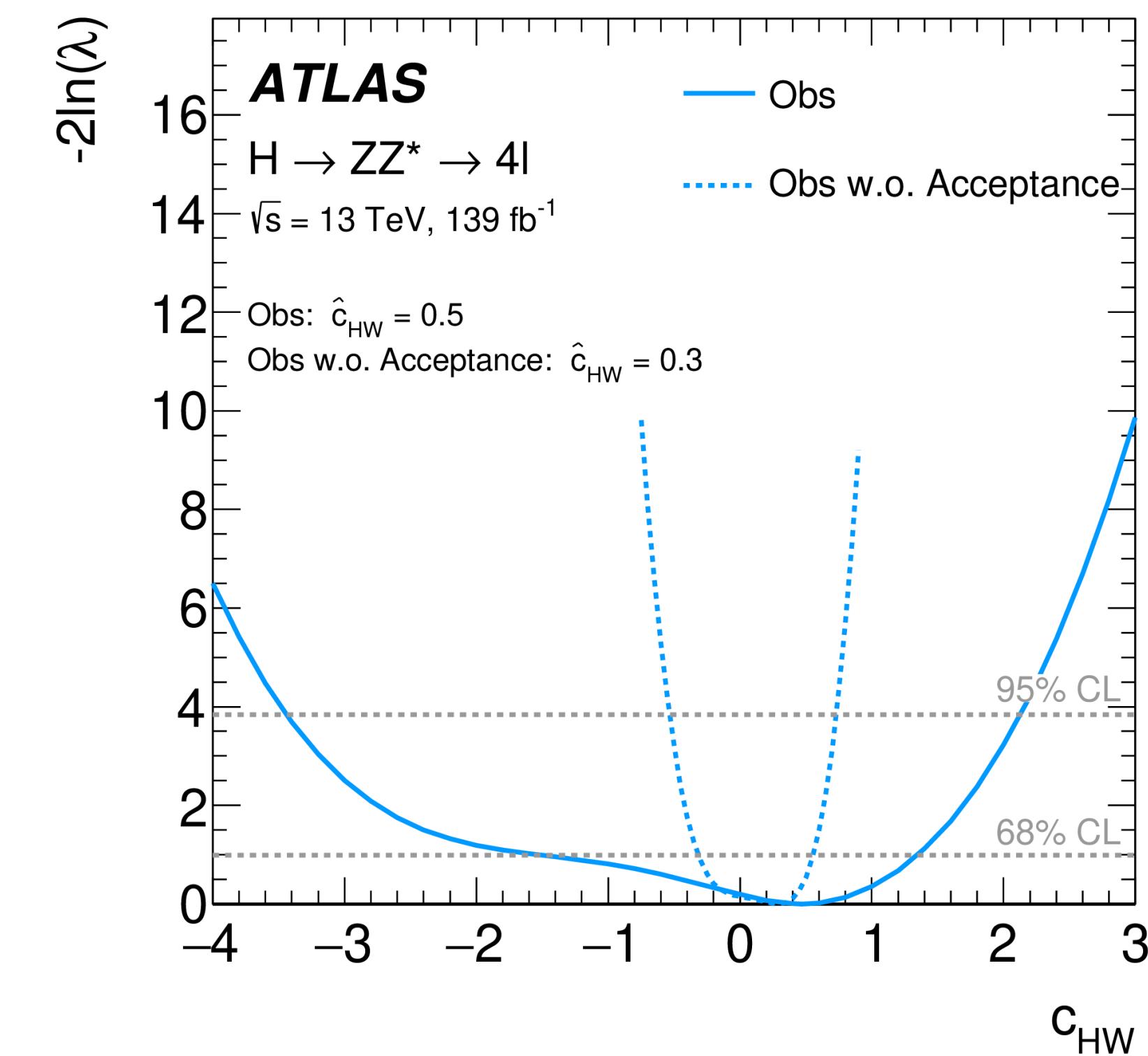
However: unfolding relies on SM assumptions

e.g. [ATLAS 2004.03447](#):

SM signal acceptance may not be a valid assumption for some SMEFT operators

- impact of including SMEFT dependence in signal acceptance is large
- effect reduced in a global fit, [2012.02779](#)

see also talk by Rahul

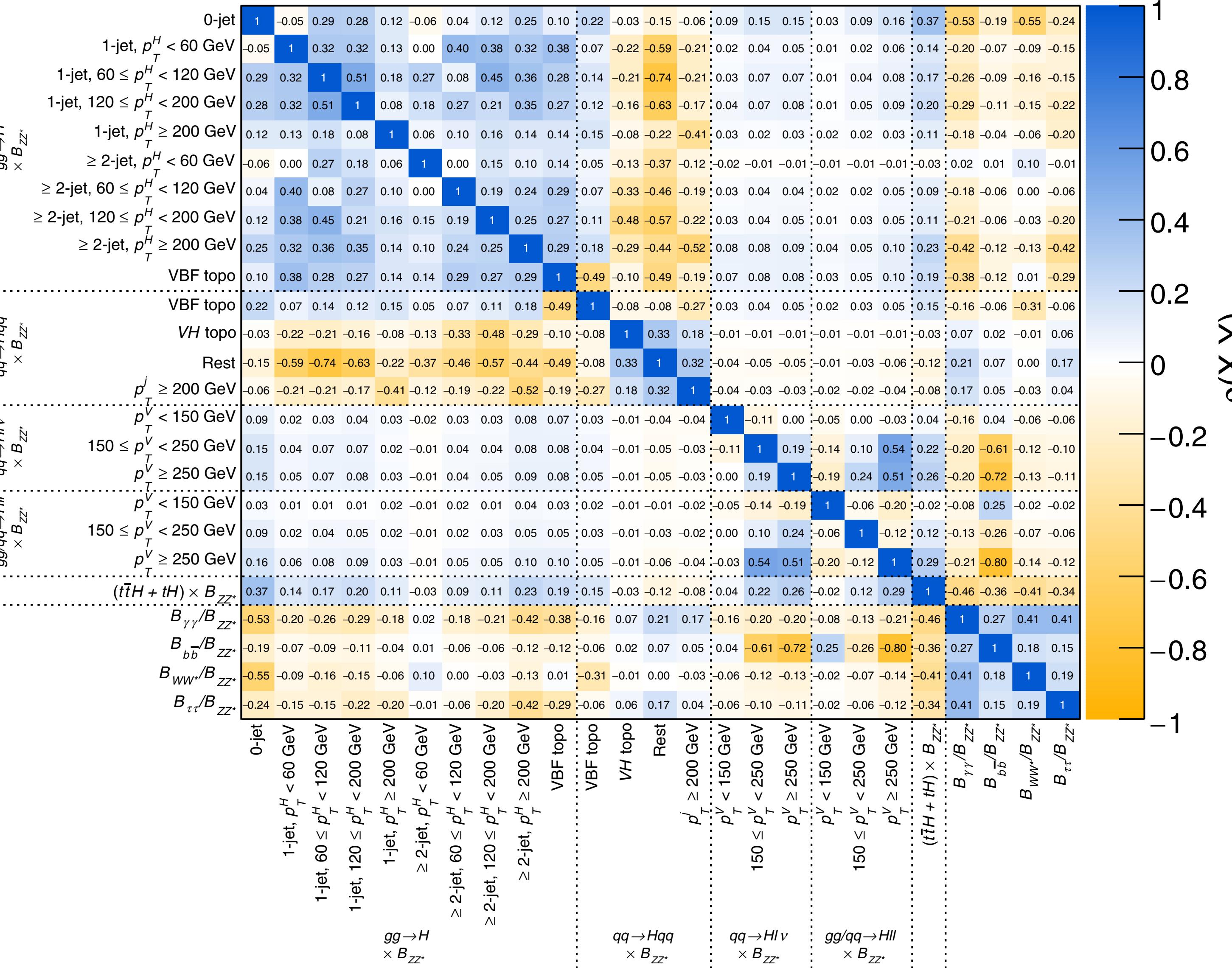


Combined fits

Combined Higgs fits of STXS/signal strengths



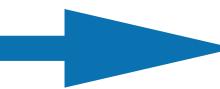
a single HEPData entry provides information on many different channels, *including correlations between channels*



Combined fits

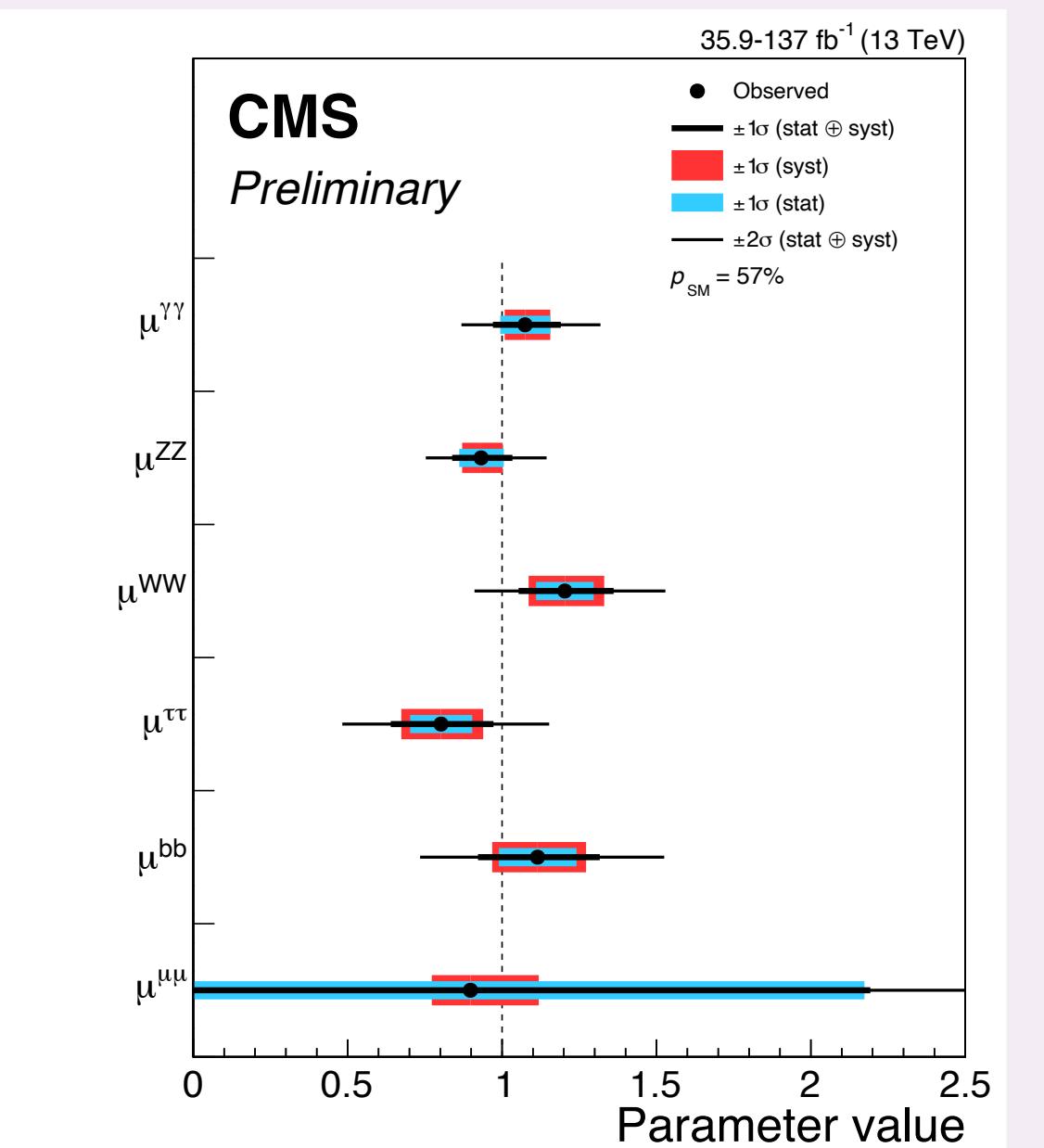
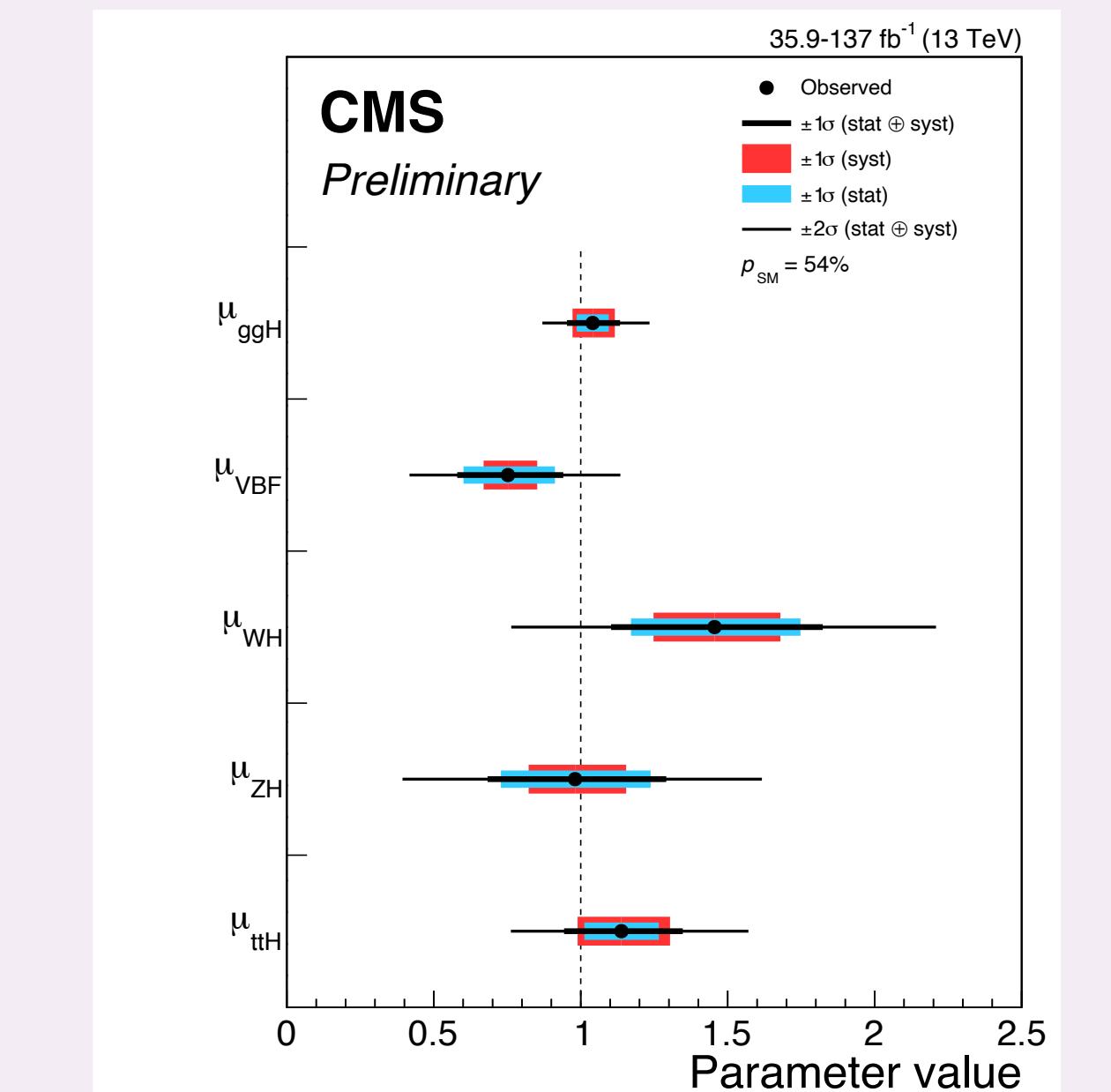
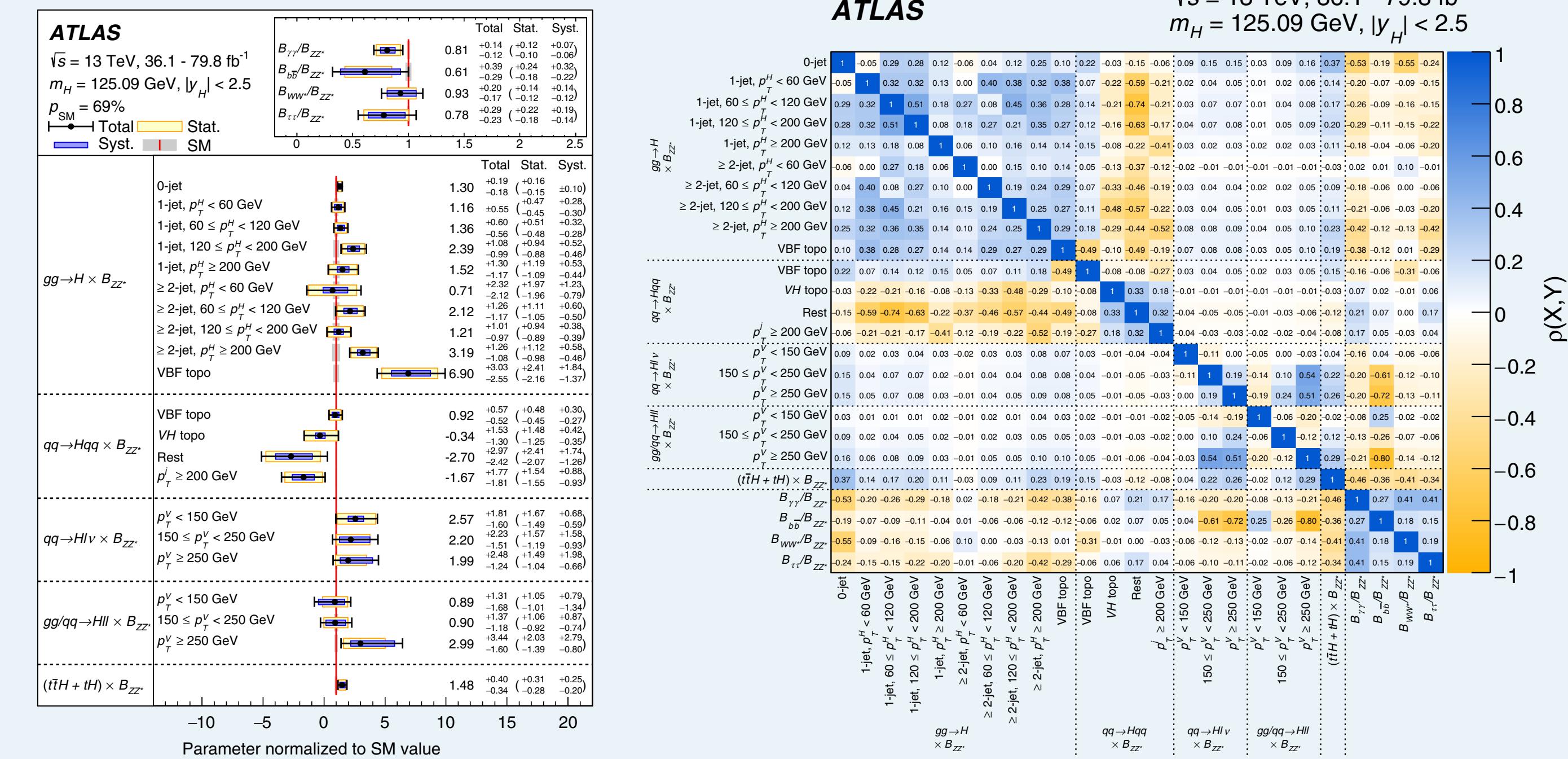
ATLAS Run II STXS combination:
[Phys. Rev. D 101 (2020) 012002]

Combined Higgs fits of STXS/signal strengths



a single HEPData entry provides information on many different channels, *including correlations between channels*

CMS Run II SS combination:
[CMS-PAS-HIG-19-005]



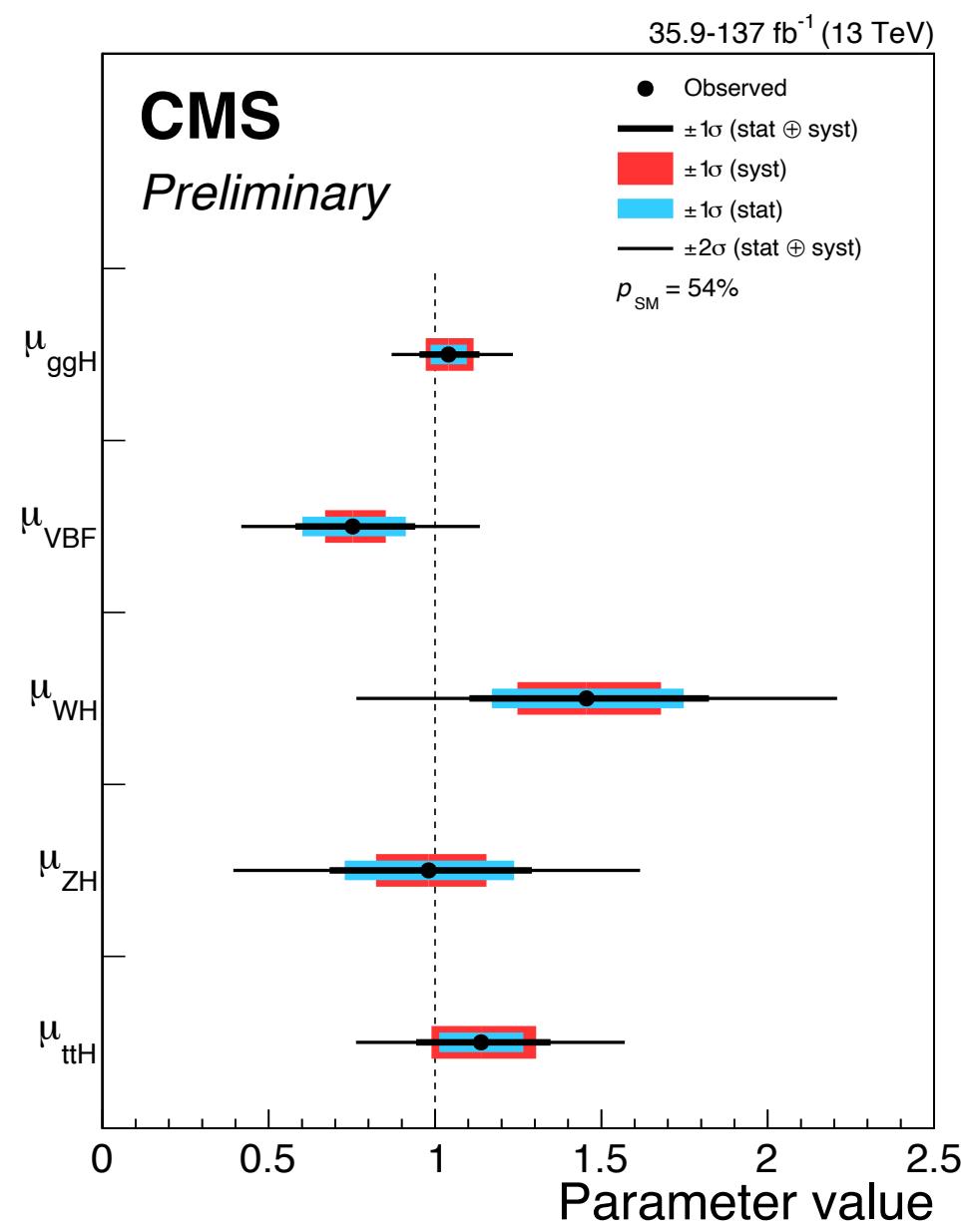
Combined fits

However: when new measurements become available, they cannot be easily swapped for measurements already in the combination

for example:

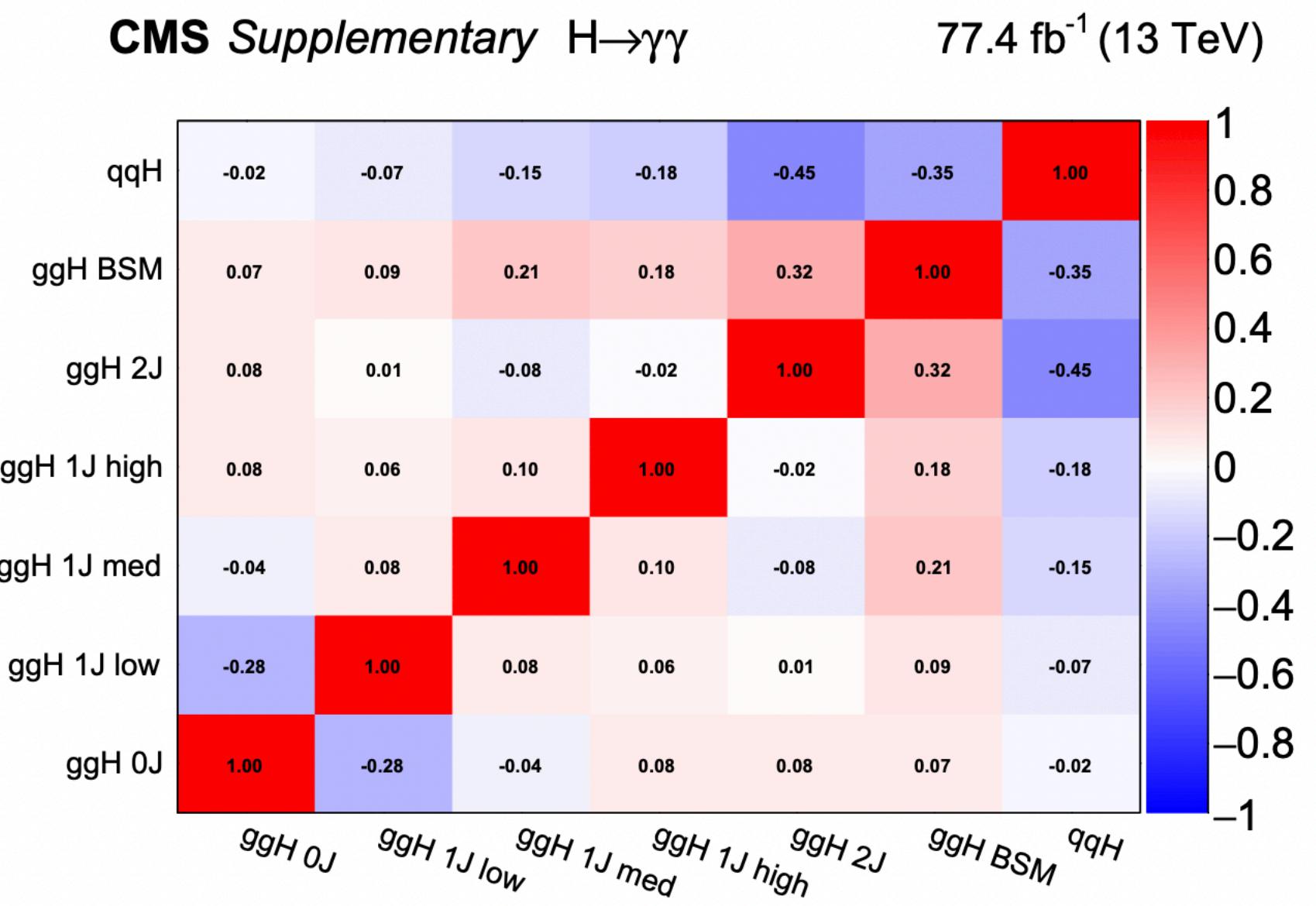
CMS Run II SS combination:
[CMS-PAS-HIG-19-005]

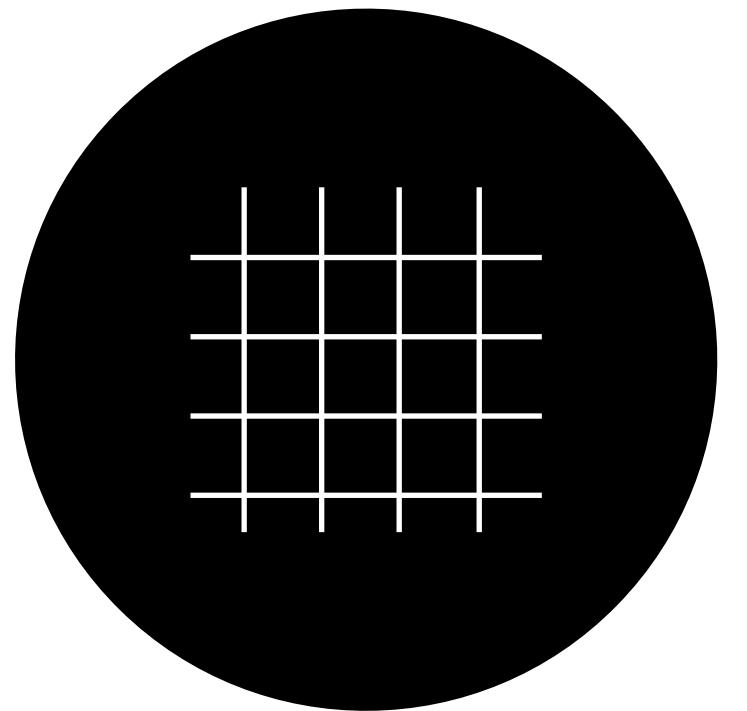
CMS Run II measurements of Higgs boson production via gluon fusion and vector boson fusion in the diphoton decay channel: [CMS-HIG-18-029-pas]



More channels - broader set of SMEFT coefficients constrained

Finer binning - better constraints on some energy-growing SMEFT coefficients

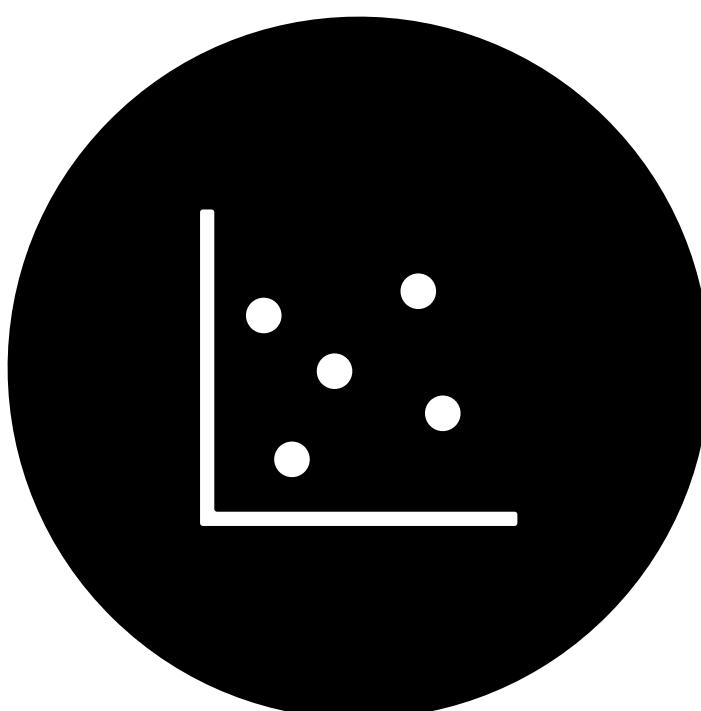




Datapoints & covariance
matrices



Likelihoods



Unbinned
measurements

Statistical Models

'Publishing statistical models: Getting the most out of particle physics experiments', 2109.04981, Cranmer et. al

$$p(n, x, y | \mu, \theta) = \prod_{i=1}^{N_c} \left[\text{Pois}(n_i | \nu_i(\mu, \theta)) \prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta) \right] p(y, \theta)$$

$\rightarrow \mathcal{L}(\mu, \theta)$ likelihoods

Statistical Models

'Publishing statistical models: Getting the most out of particle physics experiments', 2109.04981, Cranmer et. al

$$p(n, x, y | \mu, \theta) = \prod_{i=1}^{N_c} \left[\text{Pois}(n_i | \nu_i(\mu, \theta)) \prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta) \right] p(y, \theta)$$

Measured #events, features and auxiliary measurements

Predicted #events

Parameter of interest & nuisance parameters

Constraints on nuisance parameters from auxiliary measurements

Predictions for individual distributions

The diagram illustrates the components of the statistical model equation. The equation itself is:

$$p(n, x, y | \mu, \theta) = \prod_{i=1}^{N_c} \left[\text{Pois}(n_i | \nu_i(\mu, \theta)) \prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta) \right] p(y, \theta)$$

Annotations with arrows pointing to specific parts of the equation are as follows:

- Measured #events, features and auxiliary measurements: Points to the first term $p(n, x, y | \mu, \theta)$.
- Predicted #events: Points to the term $\text{Pois}(n_i | \nu_i(\mu, \theta))$.
- Parameter of interest & nuisance parameters: Points to the term $\prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta)$.
- Constraints on nuisance parameters from auxiliary measurements: Points to the term $p(y, \theta)$.
- Predictions for individual distributions: Points to the entire product term $\prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta)$.

Statistical Models

'Publishing statistical models: Getting the most out of particle physics experiments', 2109.04981, Cranmer et. al

$$p(n, x, y | \mu, \theta) = \prod_{i=1}^{N_c} \left[\text{Pois}(n_i | \nu_i(\mu, \theta)) \prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta) \right] p(y, \theta)$$

Measured #events, features and auxiliary measurements

Predicted #events

Parameter of interest & nuisance parameters

Constraints on nuisance parameters from auxiliary measurements

Predictions for individual distributions

Nuisance parameters are profiled/
marginalised away

Statistical Models

'Publishing statistical models: Getting the most out of particle physics experiments', 2109.04981, Cranmer et. al

$$p(n, x, y | \mu, \theta) = \prod_{i=1}^{N_c} \left[\text{Pois}(n_i | \nu_i(\mu, \theta)) \prod_{j=1}^{n_i} p_i(x_{ij} | \mu, \theta) \right] p(y, \theta)$$

Measured #events, features and auxiliary measurements

Predicted #events

Parameter of interest & nuisance parameters

Constraints on nuisance parameters from auxiliary measurements

Predictions for individual distributions

Nuisance parameters are profiled/
marginalised away

Access to full statistical model allows to e.g. reparametrise in terms
of new parameters of interest

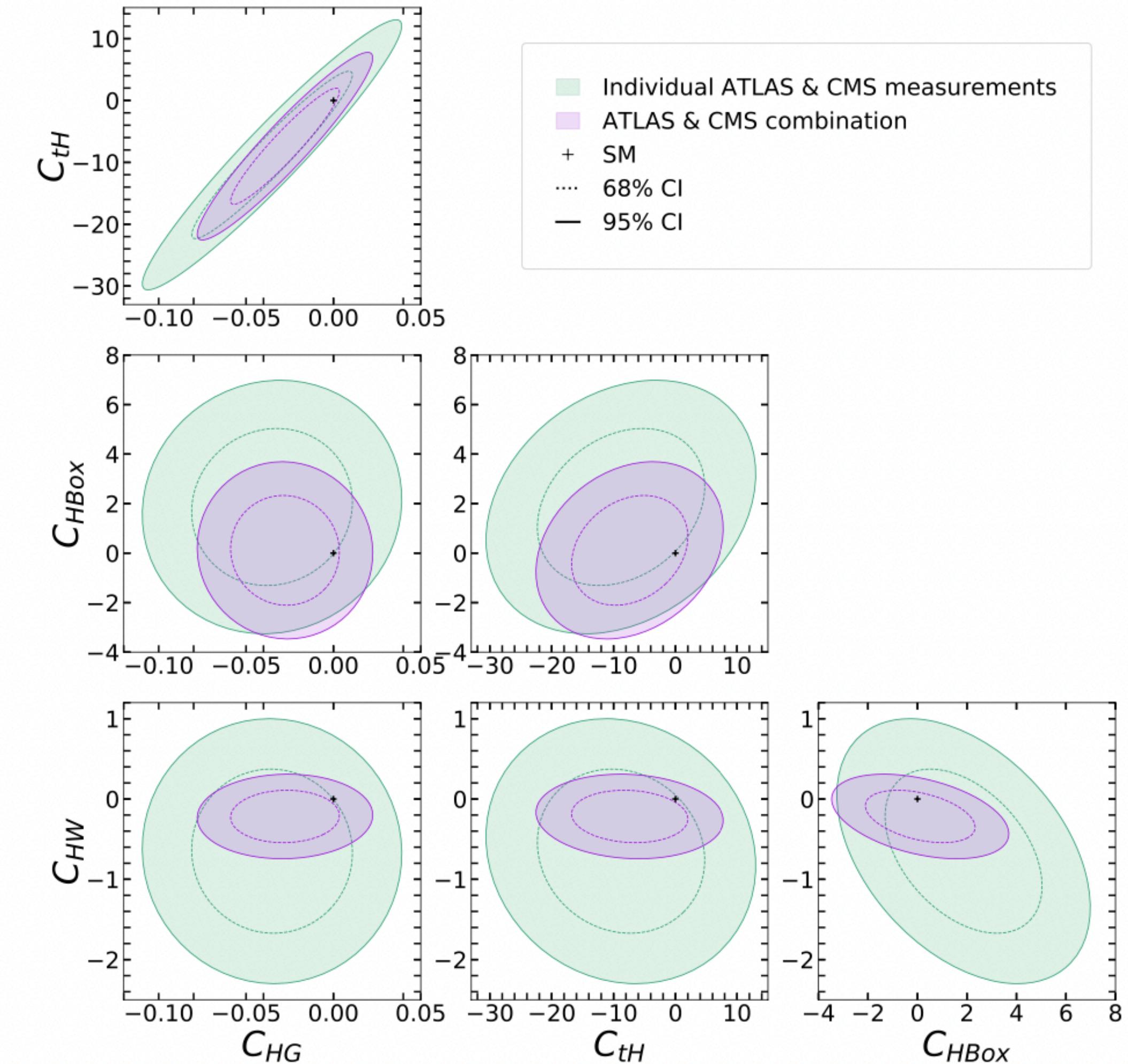
$$\mu \rightarrow \mu(\mu')$$

Example: statistical models allow for better combinations

Marginalised fit of

$$C_{tH}, C_{H\square}, C_{HW}, C_{HB}, C_{HG}$$

to Higgs Run I signal strengths.



2109.04981, Cranmer et. al

Example: statistical models allow for better combinations

Marginalised fit of

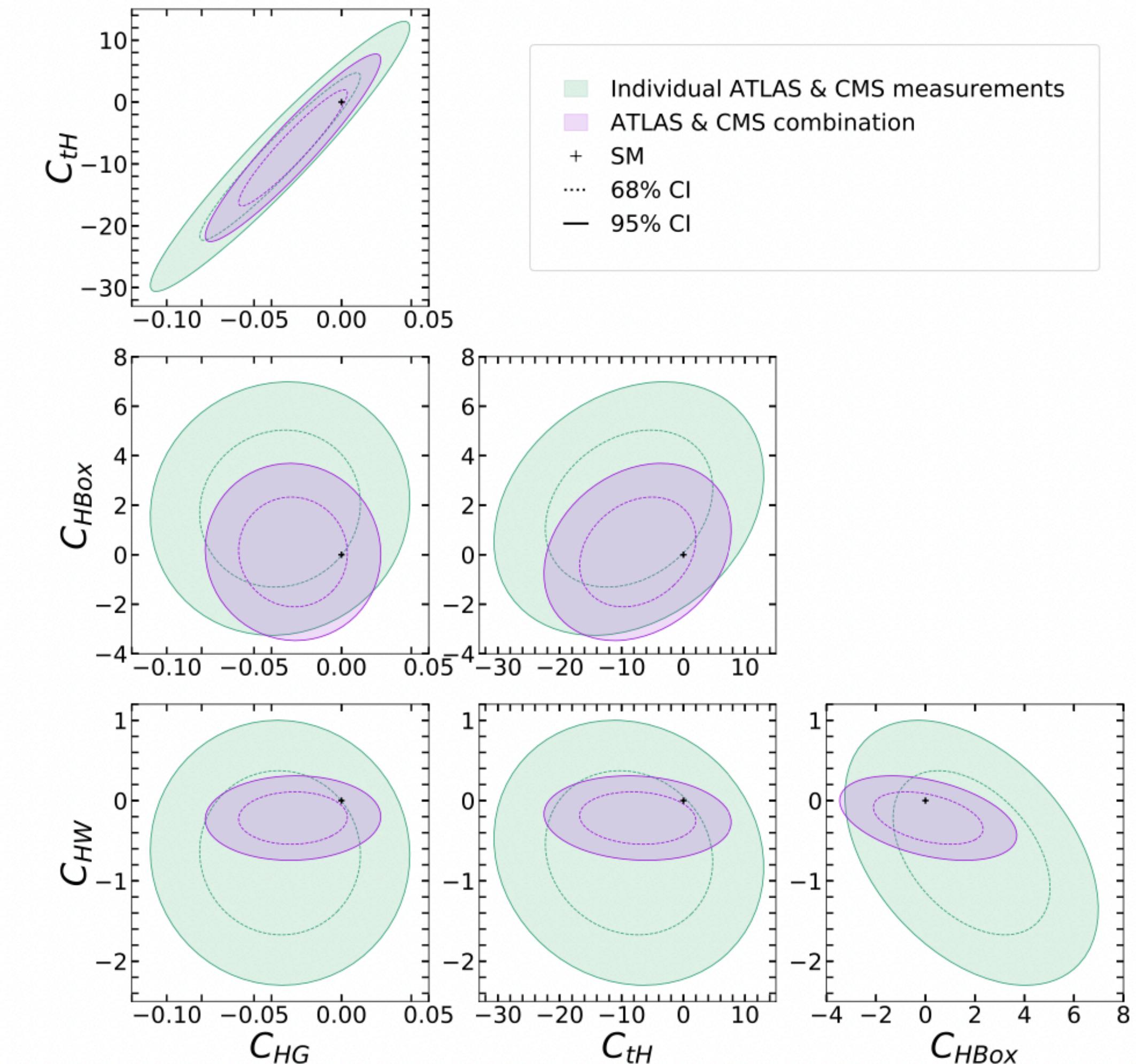
$$C_{tH}, C_{H\square}, C_{HW}, C_{HB}, C_{HG}$$

to Higgs Run I signal strengths.

Green:

ATLAS 1507.04548 (*including published covariance matrix*)
CMS 1412.8662 (*no published covariance matrix*)

Combined in Fitmaker code, neglecting correlations
between measurements



2109.04981, Cranmer et. al

Example: statistical models allow for better combinations

Marginalised fit of

$$C_{tH}, C_{H\square}, C_{HW}, C_{HB}, C_{HG}$$

to Higgs Run I signal strengths.

Green:

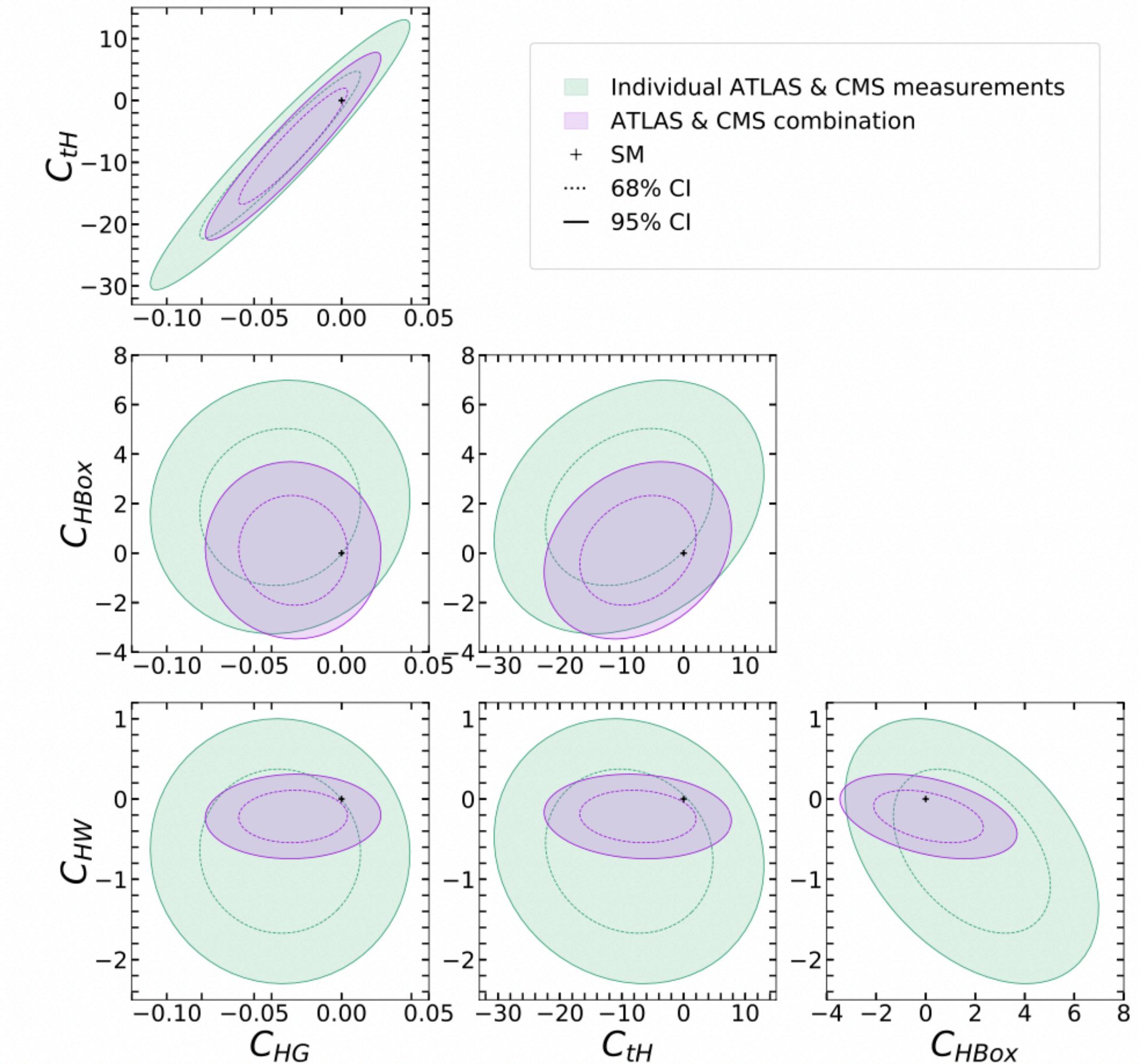
ATLAS 1507.04548 (*including published covariance matrix*)
CMS 1412.8662 (*no published covariance matrix*)

Combined in Fitmaker code, neglecting correlations
between measurements

Purple:

Combination by ATLAS & CMS 1606.02266

Combined datasets are presented with finer binning



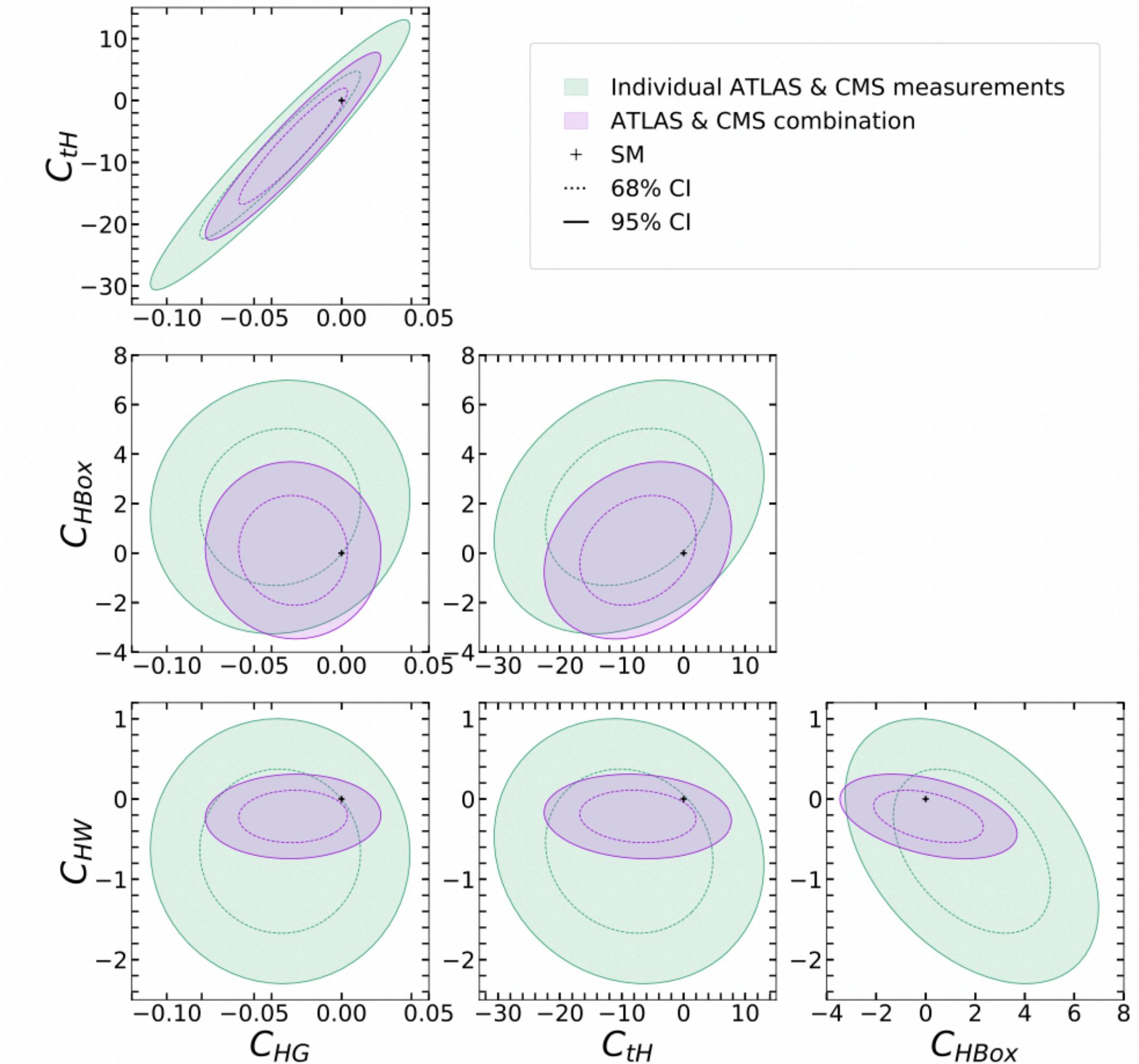
2109.04981, Cranmer et. al

Example: statistical models allow for better combinations

Access to the statistical model allows for proper combinations of data

- correlations can be taken into account
- better stats —————> finer binning can be used

→ Improved sensitivity to the SMEFT

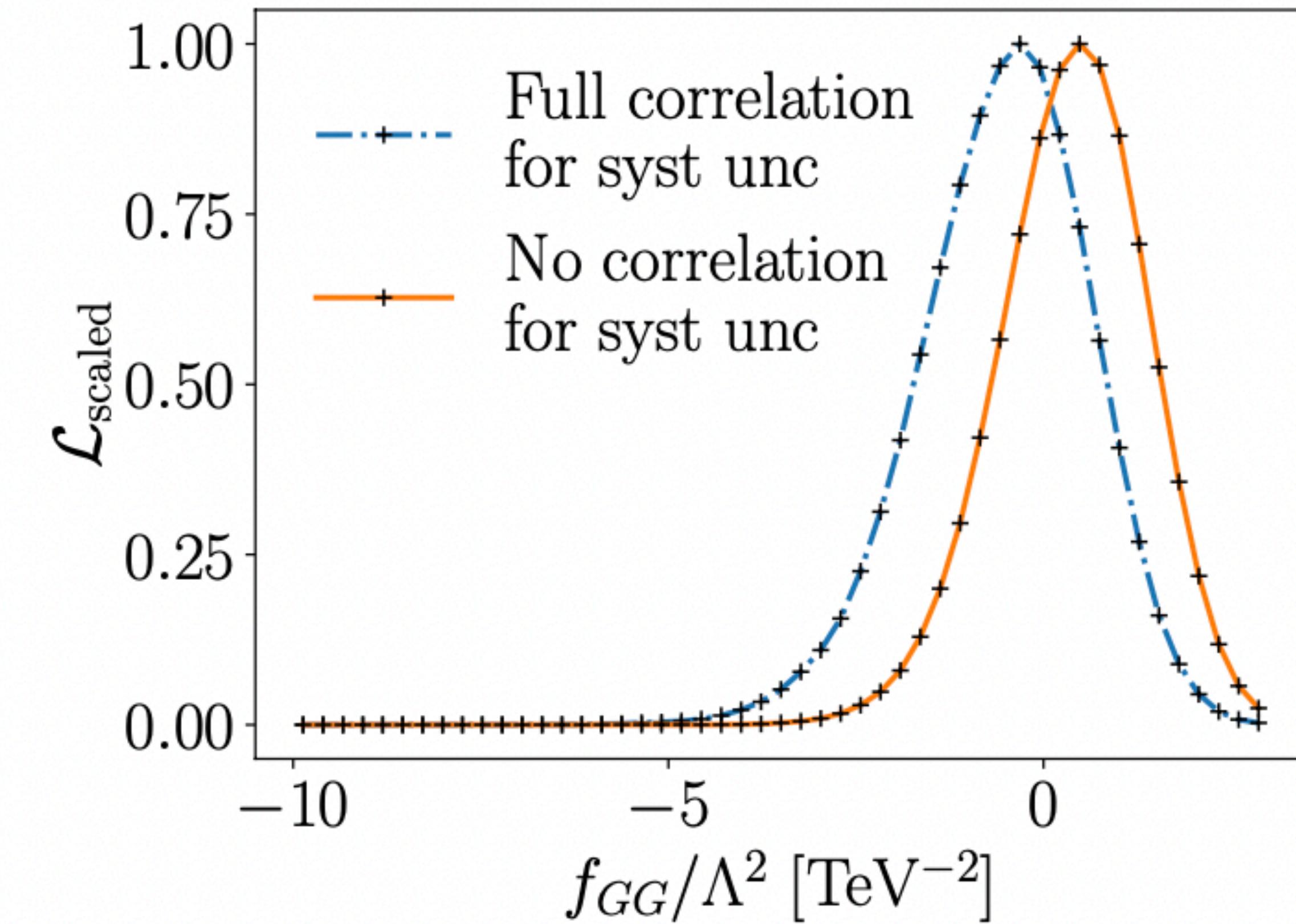


2109.04981, Cranmer et. al

Example: correlations between measurements

Published statistical models allow for a better understanding of correlated systematic uncertainties between measurements

Taking these into account has an impact on global SMEFT fits



I. Brivio et. al, 2208.08454

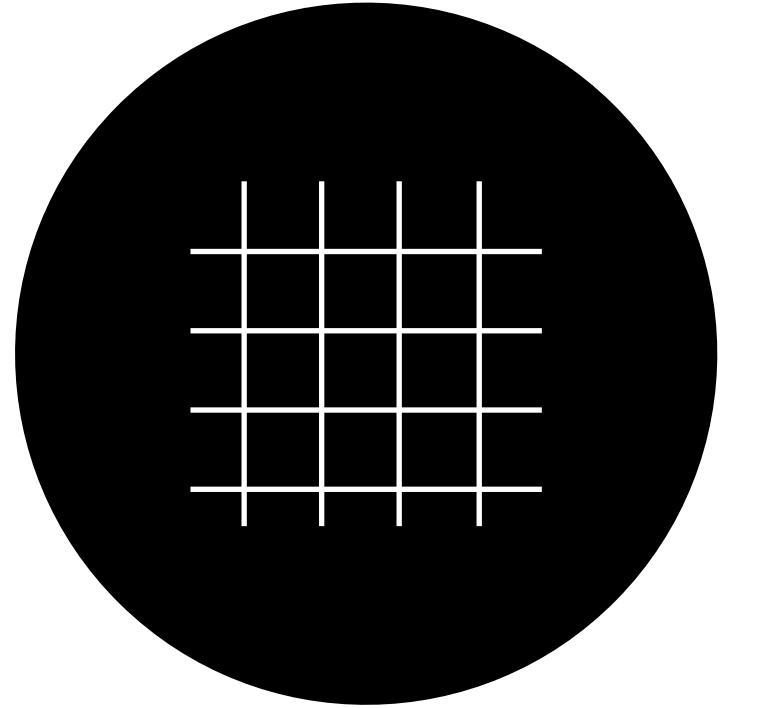
Published likelihoods

from <https://twiki.cern.ch/twiki/bin/view/AtlasPublic>

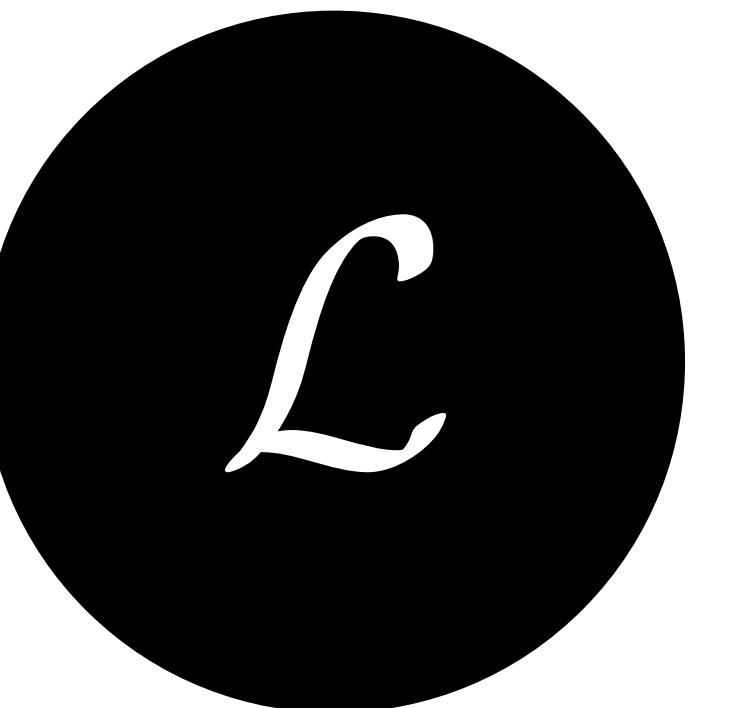
Observation of the tgamma production	TOPQ	Accepted by PRL	2023-02-02	13	140 fb ⁻¹	Documents 2302.01283 Inspire HepData Internal
Search for gluinos in multi-b final states	SUSY	Eur. Phys. J. C 83 (2023) 561	2022-11-15	13	139 fb ⁻¹	Documents 2211.08028 Inspire HepData Internal
Measurement of the s-channel single top cross-section at 13 TeV	TOPQ	JHEP 06 (2023) 191	2022-09-19	13	139 fb ⁻¹	Documents 2209.08990 Inspire HepData Internal
Search for flavor-changing neutral-current couplings between the top-quark and the photon at 13 TeV	TOPQ	Phys. Lett. B 842 (2023) 137379	2022-05-05	13	139 fb ⁻¹	Documents 2205.02537 Inspire HepData Internal
Search for SUSY in events with 2 leptons, jets and MET	SUSY	Eur. Phys. J. C 83 (2023) 515	2022-04-27	13	139 fb ⁻¹	Documents 2204.13072 Inspire HepData Internal
Search BSM H→hh→bb gamma gamma and hh→bb gamma gamma	HDBS	Phys. Rev. D 106 (2022) 052001	2021-12-22	13	139 fb ⁻¹	Documents 2112.11876 Inspire HepData Internal
Search for charginos and neutralinos in all-hadronic final states	SUSY	Phys. Rev. D 104 (2021) 112010	2021-08-17	13	139 fb ⁻¹	Documents 2108.07586 Inspire HepData Briefing Internal
4-top xsec measurement	TOPQ	JHEP 11 (2021) 118	2021-06-22	13	139 fb ⁻¹	Documents 2106.11683 Inspire HepData Internal
Search for gluinos, stops and electroweakinos in RPV models in final states with 1L and many jets	SUSY	Eur. Phys. J. C 81 (2021) 1023	2021-06-17	13	139 fb ⁻¹	Documents 2106.09609 Inspire HepData Briefing Internal

Many more likelihoods being published alongside measurements 😊

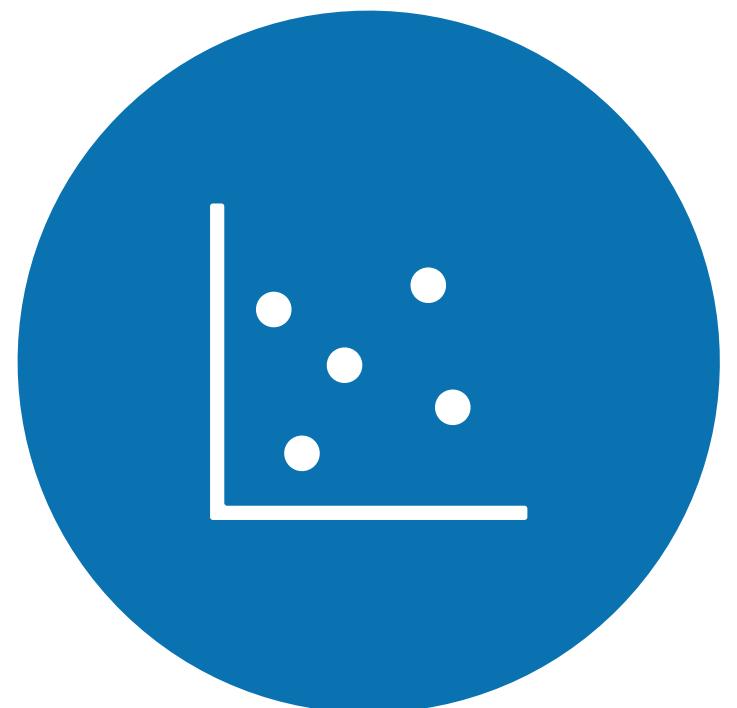
Work in progress will assess their impact on global SMEFT fits, *N. Elmer, MM, T. Plehn, N. Schmal*



Datapoints & covariance
matrices



Likelihoods



Unbinned
measurements

Unbinned Measurements

'Presenting Unbinned Differential Cross Section Results', Arratia et al, 2109.13243

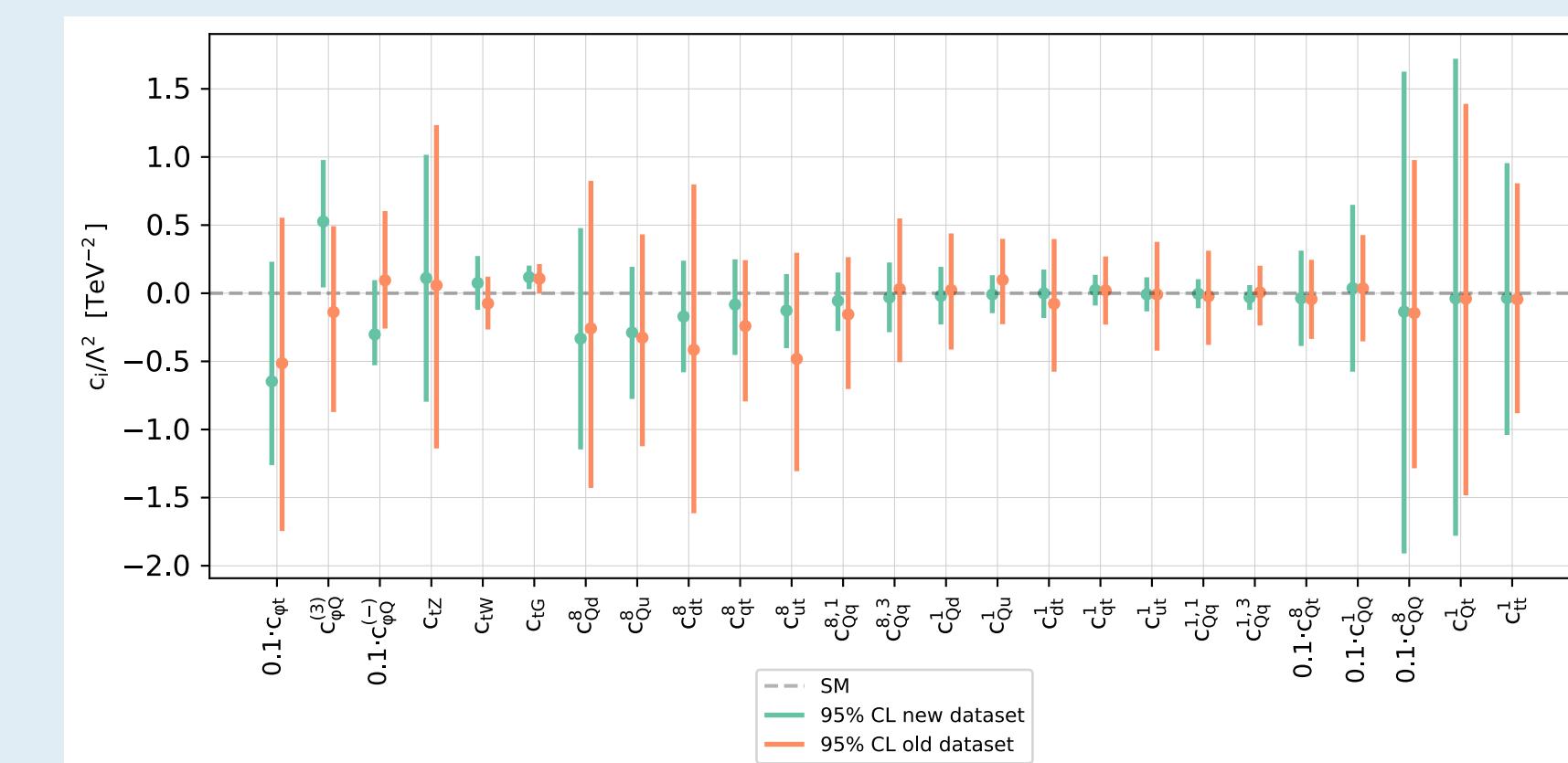
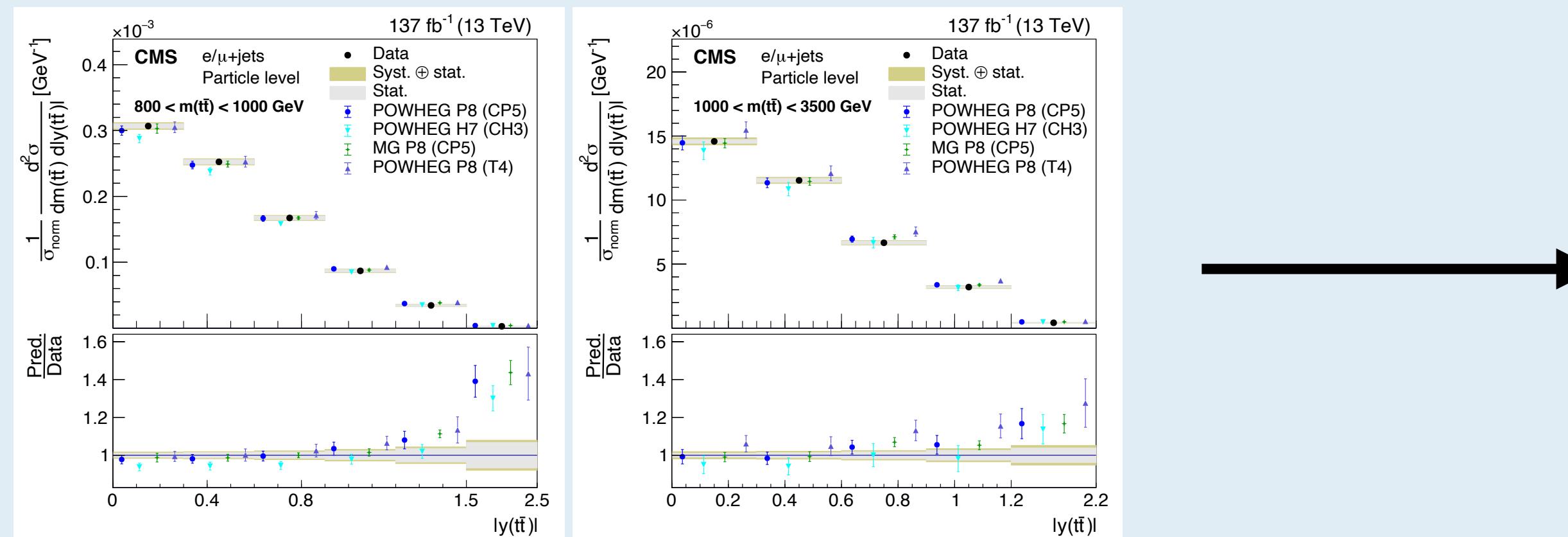
1. **Inference-aware binning:** ▶ optimal choice of binning can be made at the time of each statistical analysis or global fit

Unbinned Measurements

'Presenting Unbinned Differential Cross Section Results', Arratia et al, 2109.13243

- 1. Inference-aware binning:
 - optimal choice of binning can be made at the time of each statistical analysis or global fit

Typically we **reinterpret** measurements optimised for SM measurements or NP resonance searches



e.g. CMS measurement of top pair production in the $l+jets$ channel 2108.02803

Unbinned Measurements

'Presenting Unbinned Differential Cross Section Results', Arratia et al, 2109.13243

1. **Inference-aware binning:** ▶ optimal choice of binning can be made at the time of each statistical analysis or global fit
2. **Derivative measurements:** ▶ given measurements of features x_1, \dots, x_n , ‘post-hoc measurement’ of $f(x_1, \dots, x_n)$ possible

Unbinned Measurements

'Presenting Unbinned Differential Cross Section Results', Arratia et al, 2109.13243

1. **Inference-aware binning:** ▶ optimal choice of binning can be made at the time of each statistical analysis or global fit
2. **Derivative measurements:** ▶ given measurements of features x_1, \dots, x_n , ‘post-hoc measurement’ of $f(x_1, \dots, x_n)$ possible
3. **Extension to higher dimensions:** ▶ ML-based unbinned unfolding techniques better suited to multiple features

Open-source NN-based python framework for the integration of unbinned multivariate observables into global SMEFT interpretations.

Goal: to provide optimal constraints on the SMEFT:

Diagnostic tool:

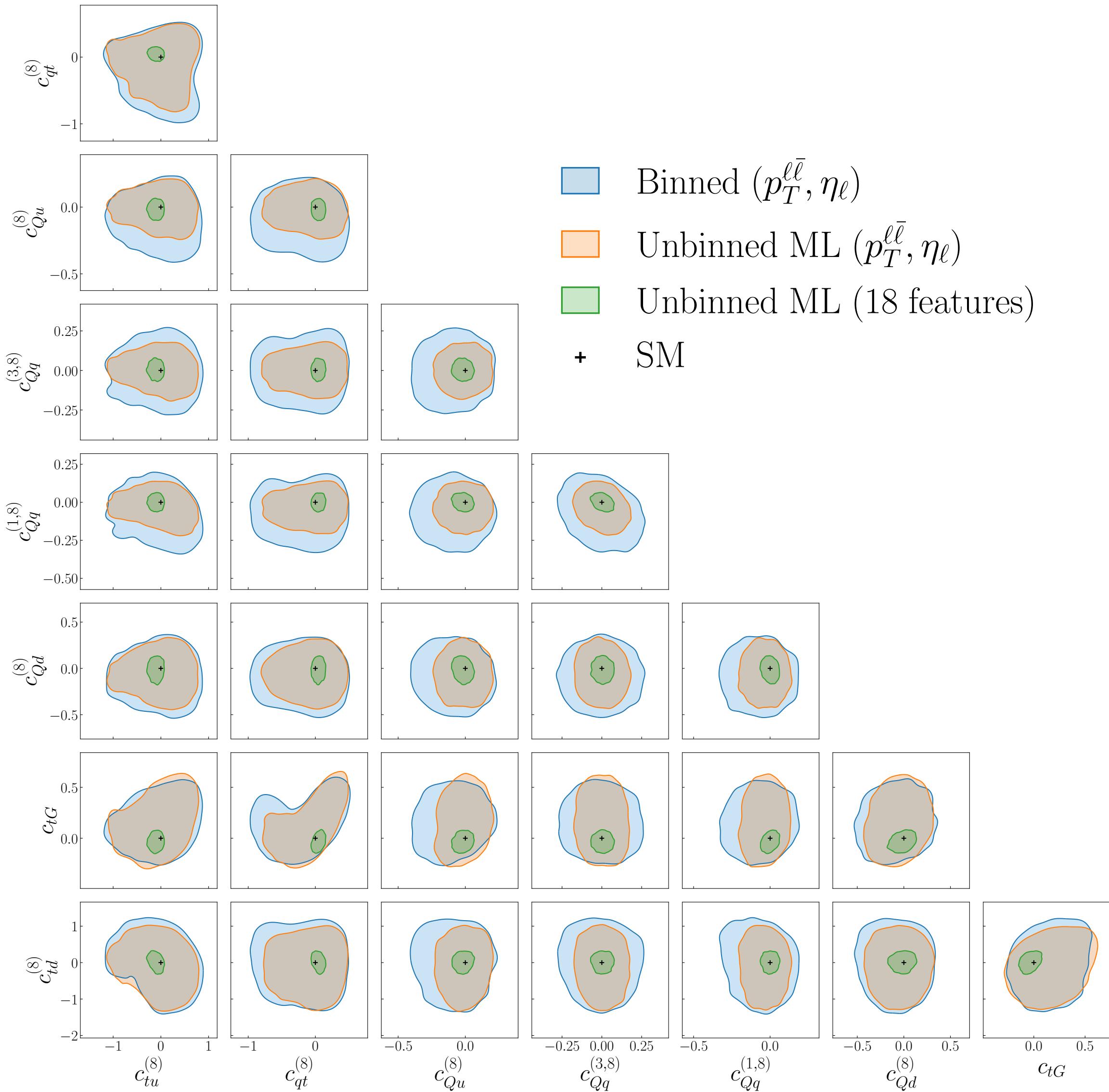
What is the information loss given a particular choice of bins?

Projections:

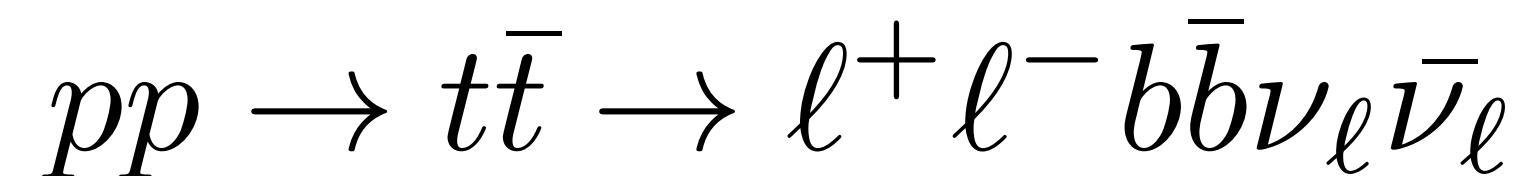
If unbinned data are made available, how will SMEFT constraints improve?

Unbinned observables in the top sector

Marginalised 95 % C.L. intervals, $\mathcal{O}(\Lambda^{-4})$ at $\mathcal{L} = 300 \text{ fb}^{-1}$



Particle-level top quark pair production in the dileptonic channel:

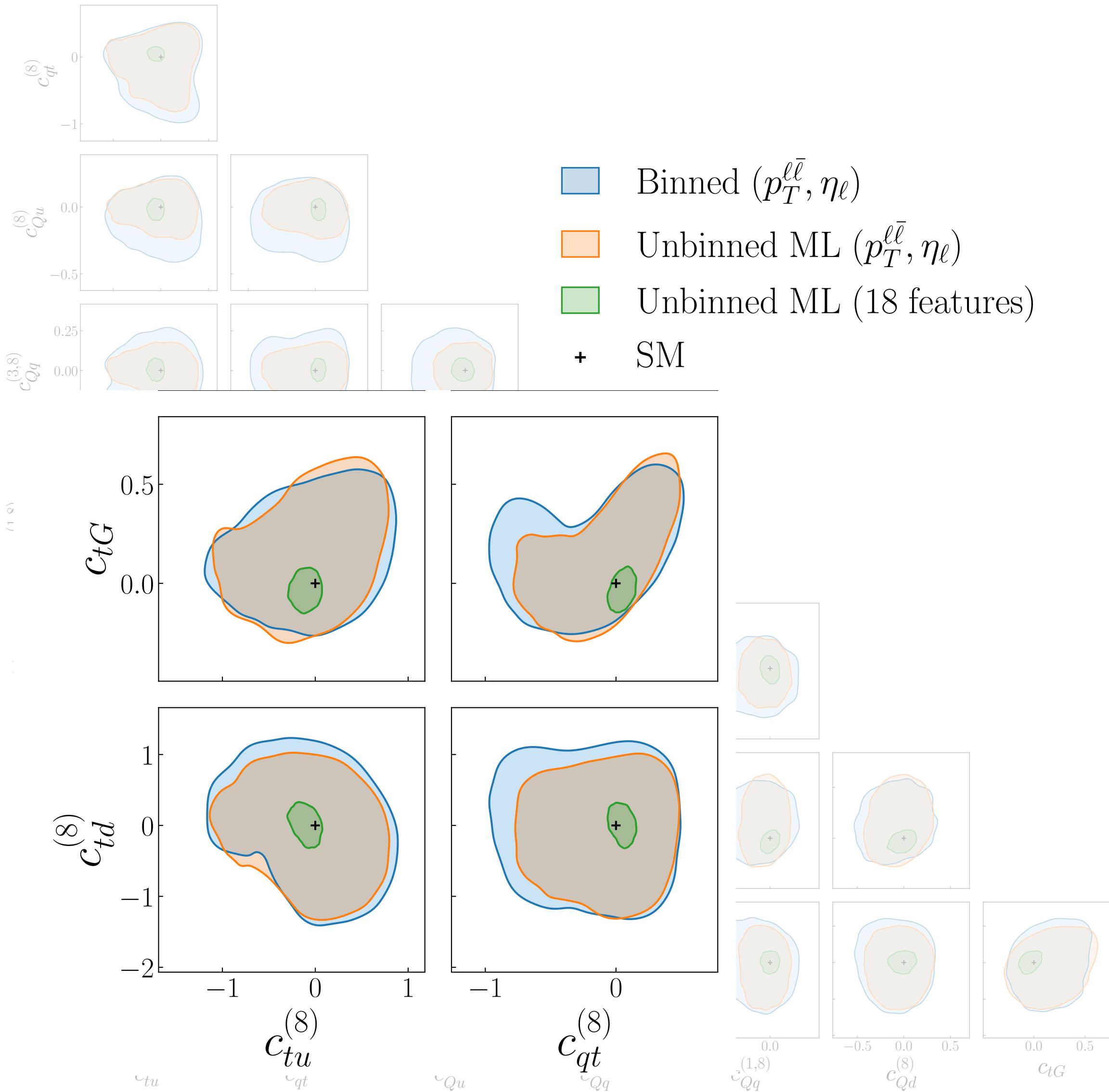


Constraints on 8 SMEFT operators:

O_{tG} + 4-fermion operators

Unbinned observables in the top sector

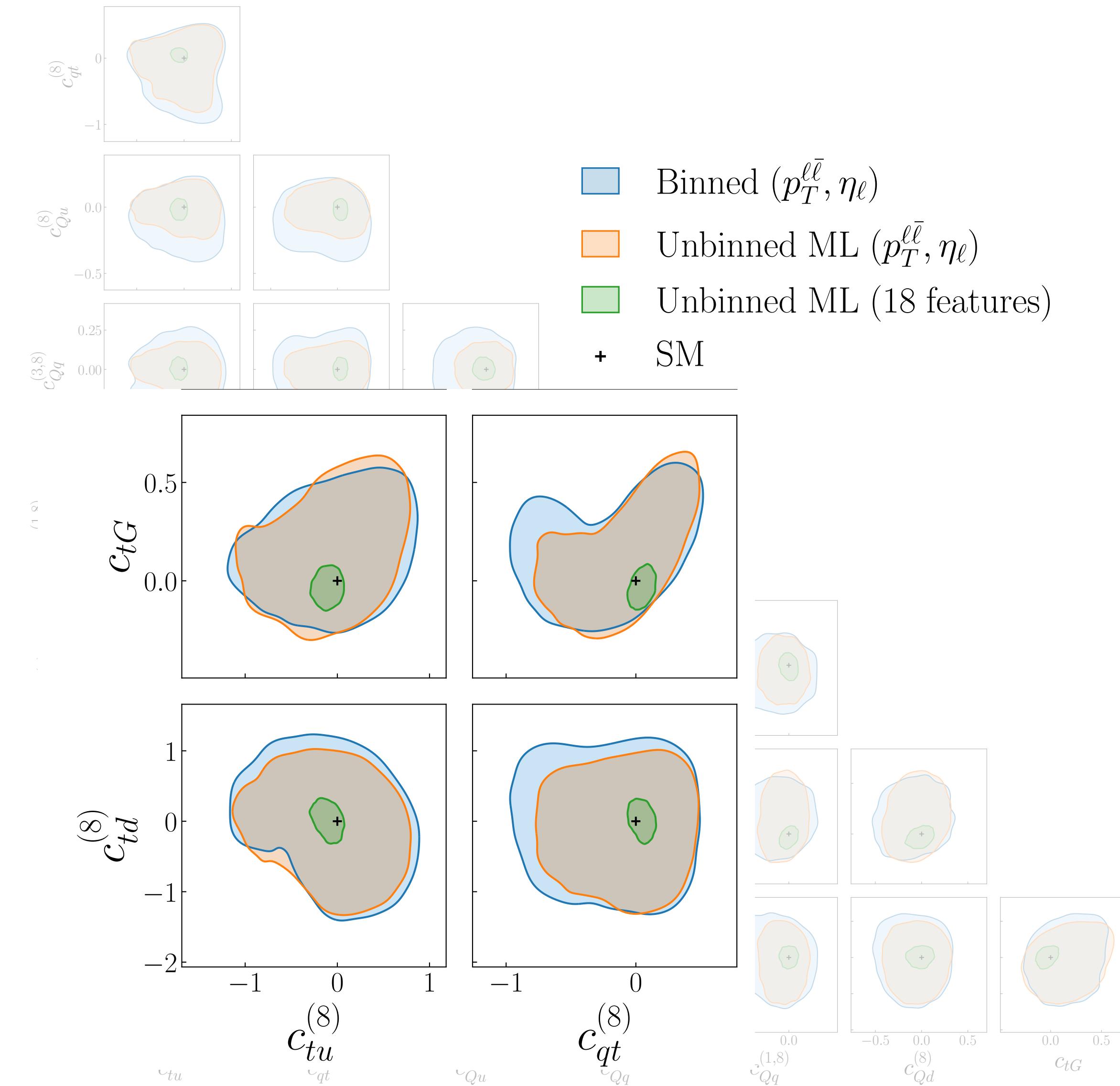
Marginalised 95 % C.L. intervals, $\mathcal{O}(\Lambda^{-4})$ at $\mathcal{L} = 300 \text{ fb}^{-1}$



Binned vs unbinned in $(p_T^{\ell\bar{\ell}}, \eta_\ell)$: small improvement from unbinned measurements, relative to nominal choice of bins

Unbinned observables in the top sector

Marginalised 95 % C.L. intervals, $\mathcal{O}(\Lambda^{-4})$ at $\mathcal{L} = 300 \text{ fb}^{-1}$



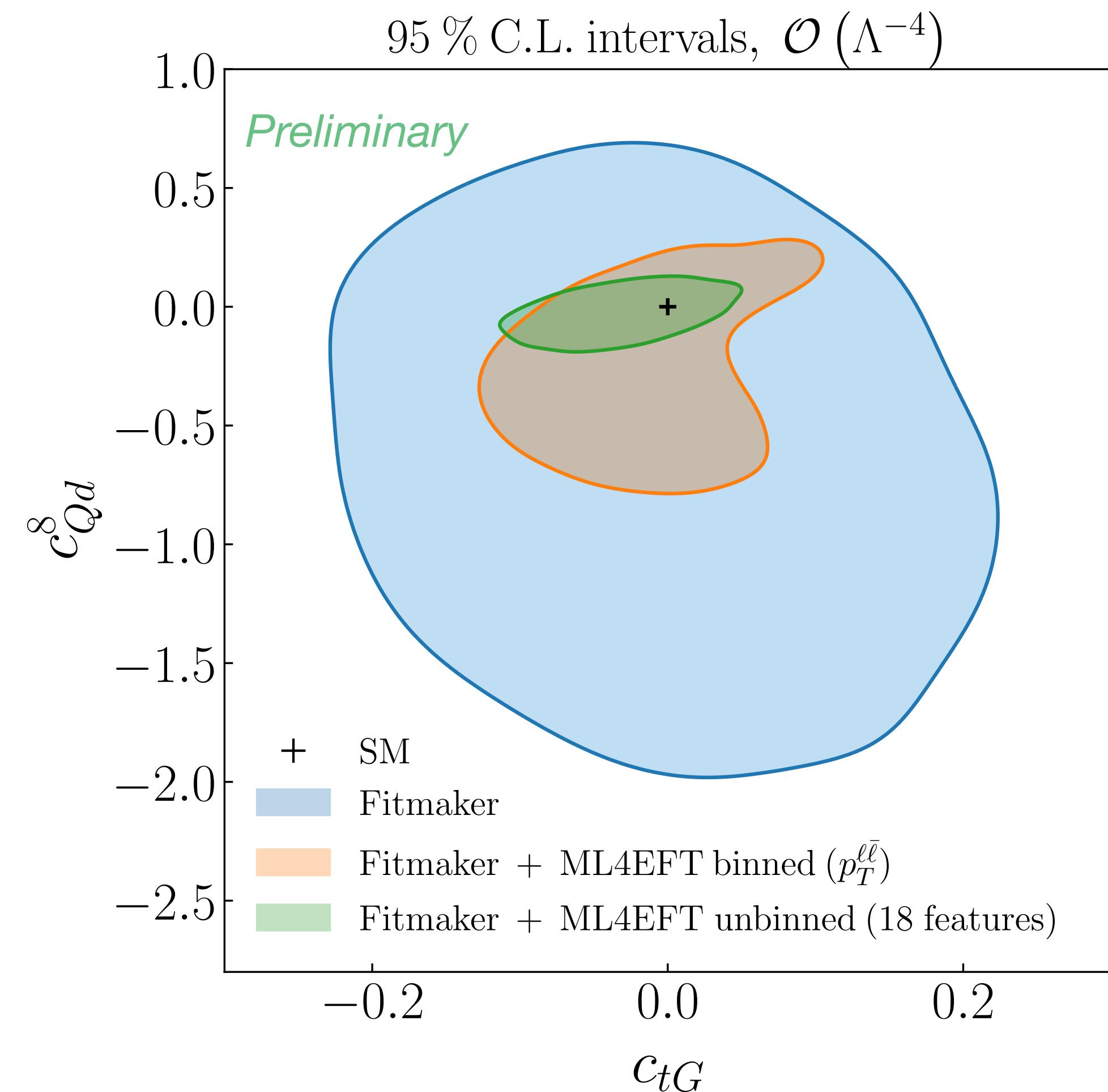
Binned vs unbinned in $(p_T^{\ell\bar{\ell}}, \eta_\ell)$: small improvement from unbinned measurements, relative to nominal choice of bins

2 features vs 18 features: vast improvement in constraining power

Integration into global SMEFT fits

New unbinned measurements can be combined alongside existing binned measurements:

$$\log\mathcal{L}(c) = \sum_{k=1}^{N_D^{(\text{unbinned})}} \log\mathcal{L}_k^{\text{unbinned}}(c) + \sum_{k=1}^{N_D^{(\text{binned})}} \log\mathcal{L}_k^{\text{binned}}(c)$$



Work in progress, Jaco ter Hoeve, MM

Conclusions

Reinterpretation of LHC data for global SMEFT fits
more information → **better SMEFT fits**

- 
- HEPData 😍
 - Likelihoods 🎉
 - Unbinned measurements 😁

Discussion points:

- Double and triple differential distributions with covariance matrices
- Unfolded data - careful of the unfolding assumptions and model dependence
- Preservation of combinations of measurements e.g. the Higgs sector
- Publication of likelihoods
- Unbinned measurements have the potential to better constrain the SMEFT

Conclusions

Thank you for listening!

Reinterpretation of LHC data for global SMEFT fits
more information → **better SMEFT fits**

- 
- HEPData 😍
 - Likelihoods 🎉
 - Unbinned measurements 😁

Discussion points:

- Double and triple differential distributions with covariance matrices
- Unfolded data - careful of the unfolding assumptions and model dependence
- Preservation of combinations of measurements e.g. the Higgs sector
- Publication of likelihoods
- Unbinned measurements have the potential to better constrain the SMEFT

Backup

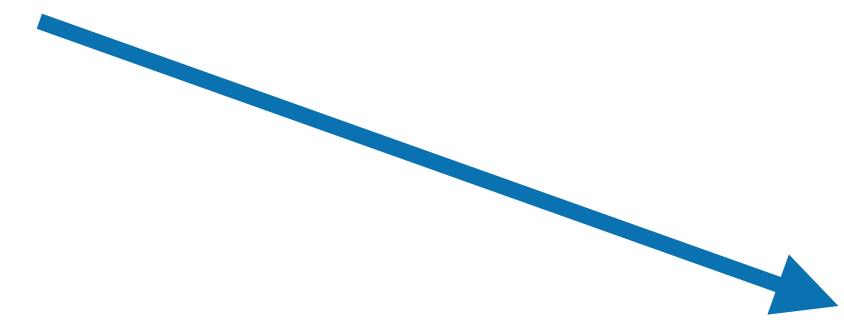
ML4EFT

For parameter estimation, we would like to be able to calculate the **likelihood**:

$$\mathcal{L}(D|\mathbf{c}) \propto \prod_{i=1}^{N_{ev}} f_\sigma(\mathbf{x}_i, \mathbf{c})$$

where $f_\sigma(\mathbf{x}, \mathbf{c}) = \frac{1}{\sigma(\mathbf{x}, \mathbf{c})} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}}$

$$D = \{\mathbf{x}_i\} \quad \mathbf{x}_i = \{m_{t\bar{t}}, p_T^{\ell_1}, p_T^{\ell_2}, \Delta\eta_{\ell_1, \ell_2}, \Delta\phi_{\ell_1, \ell_2}, \dots\}$$



multi-differential cross section in **all features**

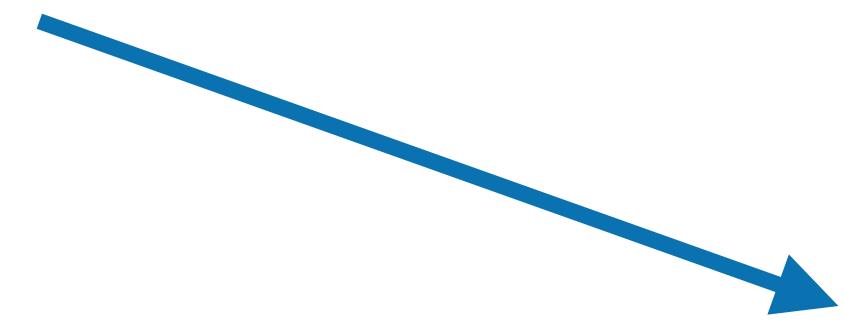
ML4EFT

For parameter estimation, we would like to be able to calculate the **likelihood**:

$$\mathcal{L}(D|\mathbf{c}) \propto \prod_{i=1}^{N_{ev}} f_\sigma(\mathbf{x}_i, \mathbf{c})$$

where $f_\sigma(\mathbf{x}, \mathbf{c}) = \frac{1}{\sigma(\mathbf{x}, \mathbf{c})} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}}$

$$D = \{\mathbf{x}_i\} \quad \mathbf{x}_i = \{m_{t\bar{t}}, p_T^{\ell_1}, p_T^{\ell_2}, \Delta\eta_{\ell_1, \ell_2}, \Delta\phi_{\ell_1, \ell_2}, \dots\}$$



multi-differential cross section in **all features**

However: analytical calculation of \mathcal{L} is intractable in most realistic cases.

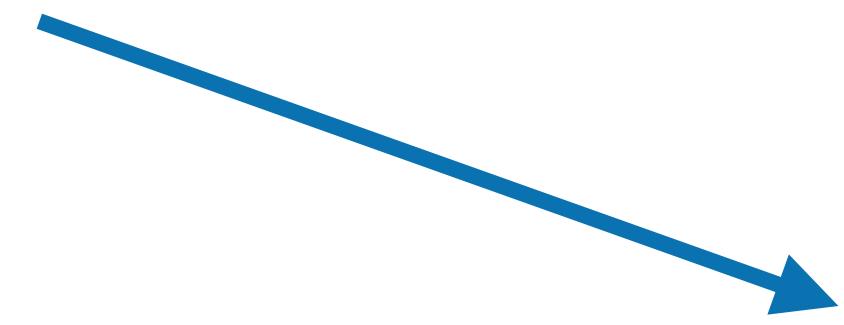
ML4EFT

For parameter estimation, we would like to be able to calculate the **likelihood**:

$$\mathcal{L}(D|\mathbf{c}) \propto \prod_{i=1}^{N_{ev}} f_\sigma(\mathbf{x}_i, \mathbf{c})$$

where $f_\sigma(\mathbf{x}, \mathbf{c}) = \frac{1}{\sigma(\mathbf{x}, \mathbf{c})} \frac{d\sigma(\mathbf{x}, \mathbf{c})}{d\mathbf{x}}$

$$D = \{\mathbf{x}_i\} \quad \mathbf{x}_i = \{m_{t\bar{t}}, p_T^{\ell_1}, p_T^{\ell_2}, \Delta\eta_{\ell_1, \ell_2}, \Delta\phi_{\ell_1, \ell_2}, \dots\}$$



multi-differential cross section in **all features**

However: analytical calculation of \mathcal{L} is intractable in most realistic cases.

Instead: approximate \mathcal{L} using neural networks

ML4EFT

Train a classifier \mathbf{g} to distinguish the SM from the SMEFT:

$$L[g(\mathbf{x}, \mathbf{c})] = - \sum_{i=1}^{N_{\text{ev}}^{\text{SMEFT}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{c})}{dx} \log(1 - g(\mathbf{x}_i, \mathbf{c})) - \sum_{i=1}^{N_{\text{ev}}^{\text{SM}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{0})}{dx} \log(g(\mathbf{x}_i, \mathbf{0}))$$

SMEFT training pseudodata sample *SM training pseudo data sample*



ML4EFT

Train a classifier \mathbf{g} to distinguish the SM from the SMEFT:

$$L[g(\mathbf{x}, \mathbf{c})] = - \sum_{i=1}^{N_{\text{ev}}^{\text{SMEFT}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{c})}{dx} \log(1 - g(\mathbf{x}_i, \mathbf{c})) - \sum_{i=1}^{N_{\text{ev}}^{\text{SM}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{0})}{dx} \log(g(\mathbf{x}_i, \mathbf{0}))$$

SMEFT training pseudodata sample *SM training pseudo data sample*

$$\frac{\delta L}{\delta g} = 0 \Rightarrow g(\mathbf{x}, \mathbf{c}) = \left(1 + \frac{d\sigma(\mathbf{x}, \mathbf{c})}{dx} / \frac{d\sigma(\mathbf{x}, \mathbf{0})}{dx} \right)^{-1} \equiv \frac{1}{1 + r_\sigma(\mathbf{x}, \mathbf{c})}$$

ML4EFT

Train a classifier \mathbf{g} to distinguish the SM from the SMEFT:

$$L[g(\mathbf{x}, \mathbf{c})] = - \sum_{i=1}^{N_{\text{ev}}^{\text{SMEFT}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{c})}{dx} \log(1 - g(\mathbf{x}_i, \mathbf{c})) - \sum_{i=1}^{N_{\text{ev}}^{\text{SM}}} \frac{d\sigma(\mathbf{x}_i, \mathbf{0})}{dx} \log(g(\mathbf{x}_i, \mathbf{0}))$$

SMEFT training pseudodata sample *SM training pseudo data sample*

$$\frac{\delta L}{\delta g} = 0 \Rightarrow g(\mathbf{x}, \mathbf{c}) = \left(1 + \frac{d\sigma(\mathbf{x}, \mathbf{c})}{dx} / \frac{d\sigma(\mathbf{x}, \mathbf{0})}{dx} \right)^{-1} \equiv \frac{1}{1 + r_\sigma(\mathbf{x}, \mathbf{c})}$$

In the limit of infinite training samples, the decision boundary is 1:1 with the likelihood

$$\hat{g} = \frac{1}{1 + \hat{r}_\sigma(x, c)}$$

$$r_\sigma(\mathbf{x}, \mathbf{c}) = \frac{d\sigma(\mathbf{x}, \mathbf{c})}{dx} / \frac{d\sigma(\mathbf{x}, 0)}{dx}$$

Exploit the polynomial structure of the SMEFT when defining the classifier \mathbf{g} :

$$\hat{r}_\sigma(\mathbf{x}, \mathbf{c}) = 1 + \sum_{j=1}^{n_{\text{eft}}} \text{NN}^{(j)}(\mathbf{x}) c_j + \sum_{j=1}^{n_{\text{eft}}} \sum_{k \geq j}^{n_{\text{eft}}} \text{NN}_\sigma^{(j,k)}(\mathbf{x}) c_j c_k$$

Parallelisable: generate a training sample with only c_i and learn only $\text{NN}^i(\mathbf{x})$

well-suited to global fits of many SMEFT coefficients