

Studying the Higgs boson in the SMEFT: gluon-induced ZH production

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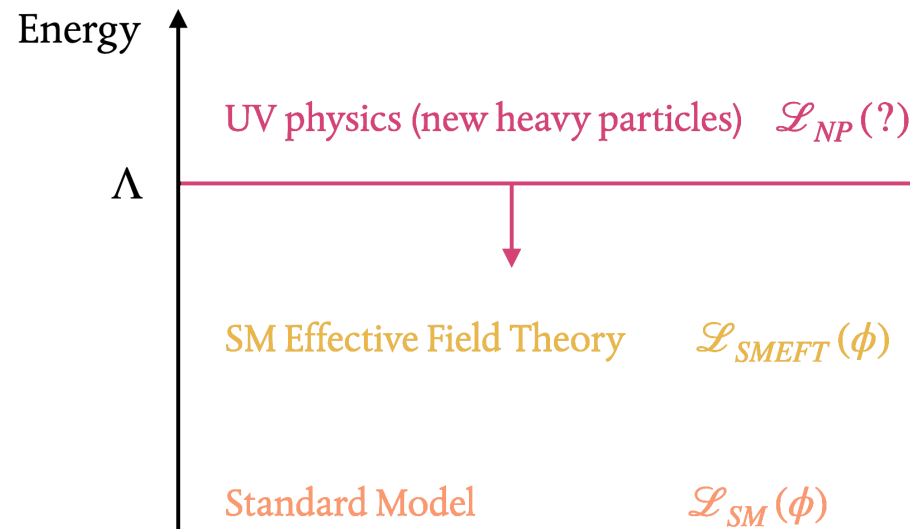
Abingdon, 21/11/23

Based on arXiv: 2306.09963 (accepted in JHEP)

In collaboration with A. Rossia and E. Vryonidou



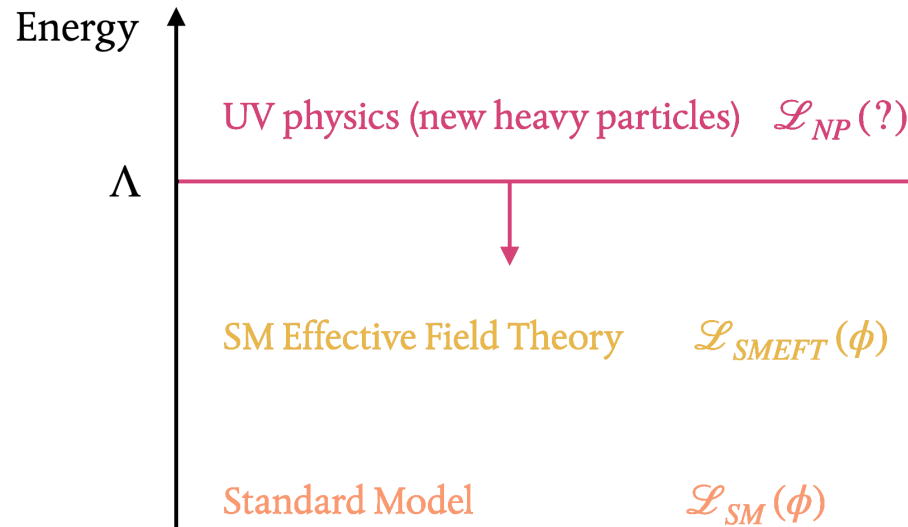
About the SMEFT framework



$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}}^{(4)} + \frac{1}{\Lambda} c^{(5)} \mathcal{O}^{(5)} + \frac{1}{\Lambda^2} \sum_k c_k^{(6)} \mathcal{O}_k^{(6)} + \mathcal{O}\left(\frac{1}{\Lambda^3}\right)$$



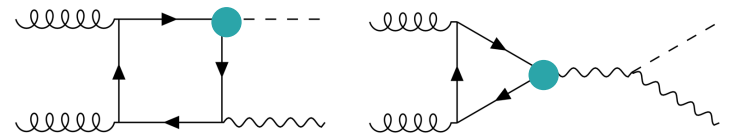
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Focus on dimension-6 operators with Warsaw basis.

$gg \rightarrow ZH$: top quark loop \rightarrow probes Higgs and top operators



Growing amplitudes in $gg \rightarrow ZH$

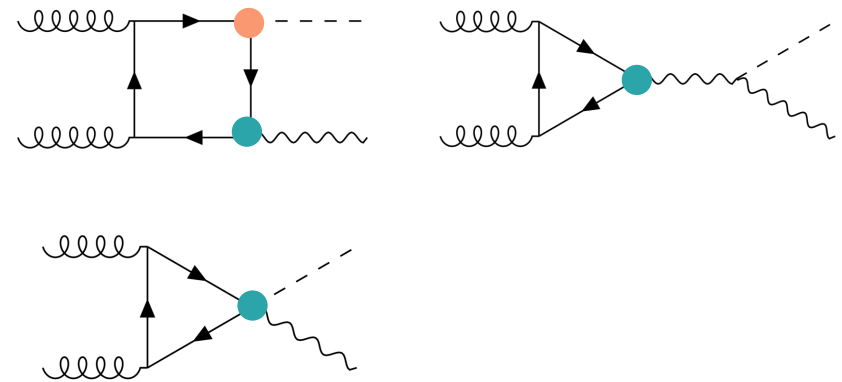
- Calculated analytical helicity amplitudes with 1 insertion of dim-6 SMEFT operators.
- Studied high-energy behaviour of amplitudes \rightarrow Which operators grow with energy?



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$\lambda_{g_1}, \lambda_{g_2}, \lambda_H, \lambda_Z$	$\mathcal{O}_{\varphi t}$	$\mathcal{O}_{\varphi Q}^{(-)}$	$\mathcal{O}_{t\varphi}$
$+, +, 0, 0$	$\frac{m_t^2 v e g_s^2}{32\pi^2 m_Z c_w s_w} \left[\log\left(\frac{s}{m_t^2}\right) - i\pi \right]^2$	$\frac{m_t^2 v e g_s^2}{32\pi^2 m_Z c_w s_w} \left[\log\left(\frac{s}{m_t^2}\right) - i\pi \right]^2$	$\frac{m_t v^2 e g_s^2}{32\sqrt{2}\pi^2 m_Z c_w s_w} \left[\log\left(\frac{s}{m_t^2}\right) - i\pi \right]^2$

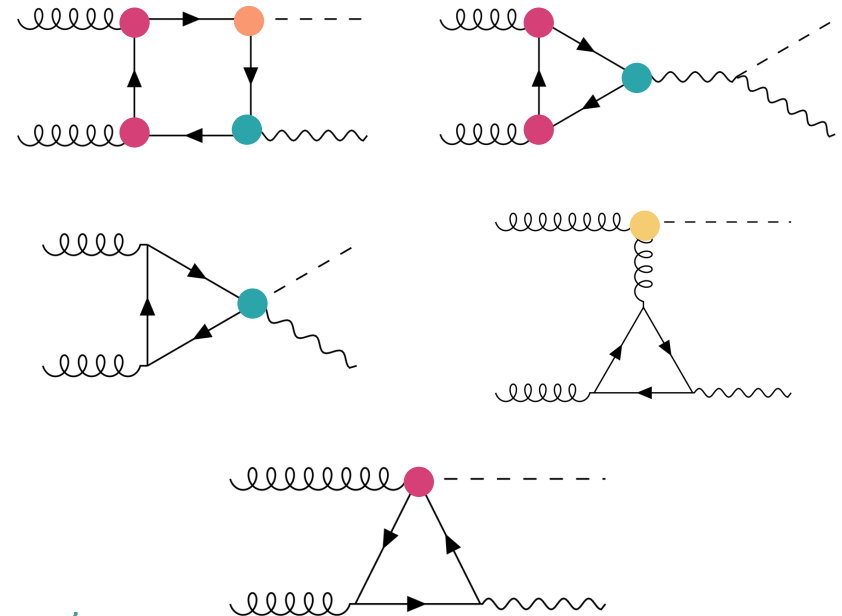


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$\lambda_{g_1}, \lambda_{g_2}, \lambda_H, \lambda_Z$	\mathcal{O}_{tG}	$\mathcal{O}_{\varphi G}$
$+, +, 0, +$	$\sqrt{s} m_t \log\left(\frac{s}{m_t^2}\right)$	—
$+, +, 0, -$	$\sqrt{s} m_t \log^2\left(\frac{s}{m_t^2}\right)$	—
$+, +, 0, 0$	$\frac{m_t v^2}{m_Z} \log^2\left(\frac{s}{m_t^2}\right)$	$\frac{m_t^2 v}{m_Z} \log^2\left(\frac{s}{m_t^2}\right)$
$+, -, 0, +$	$\sqrt{s} m_t$	—
$+, -, 0, 0$	$s \frac{m_t}{m_Z}$	$\frac{m_t^2 v}{m_Z} \log^2\left(\frac{s}{m_t^2}\right)$



Can measuring $pp \rightarrow ZH$ improve the bounds on Higgs and top operators?

Third generation operators

Quark and gluon channels interplay

$$\mathcal{O}_{\varphi Q}^{(1)} \quad c_{\varphi Q}^{(1)}$$

$$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{Q} \gamma^\mu Q)$$

$$\mathcal{O}_{\varphi Q}^{(3)} \quad c_{\varphi Q}^{(3)}$$

$$i(\varphi^\dagger \overleftrightarrow{D}_\mu \tau_I \varphi) (\bar{Q} \gamma^\mu \tau^I Q)$$

$$\mathcal{O}_{\varphi Q}^{(-)} \quad c_{\varphi Q}^{(-)}$$

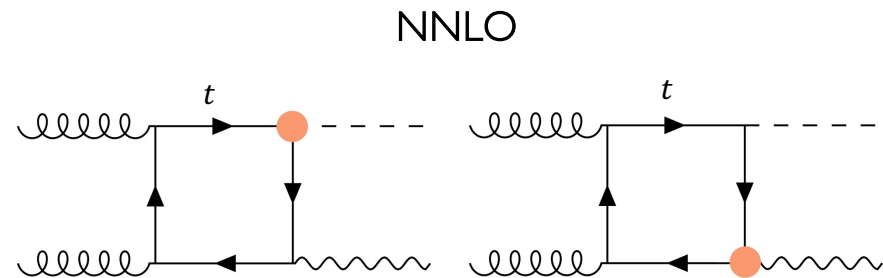
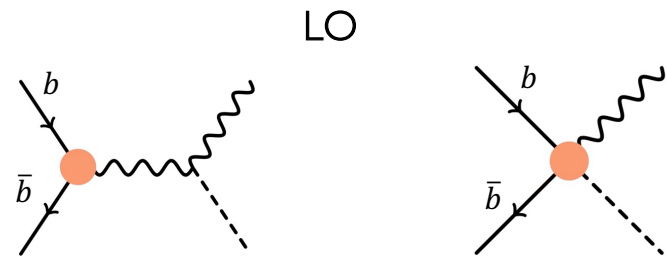
$$c_{\varphi Q}^{(1)} - c_{\varphi Q}^{(3)}$$

$$\mathcal{O}_{\varphi t} \quad c_{\varphi t}$$

$$i(\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{t} \gamma^\mu t)$$

$$\mathcal{O}_{t\varphi} \quad c_{t\varphi}$$

$$\left(\varphi^\dagger \varphi - \frac{v^2}{2}\right) \bar{Q} t \tilde{\varphi} + \text{h.c.}$$



Probed by $gg \rightarrow ZH$

Probed by $qq \rightarrow ZH$



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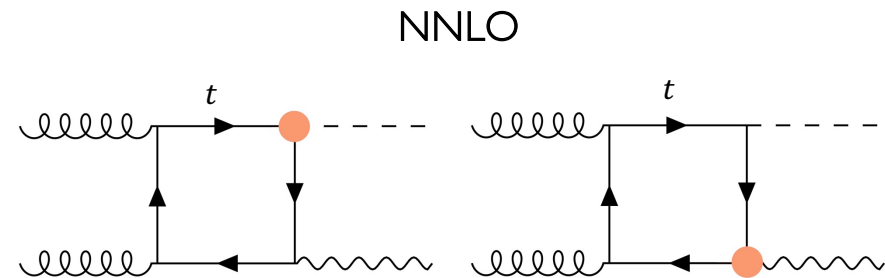
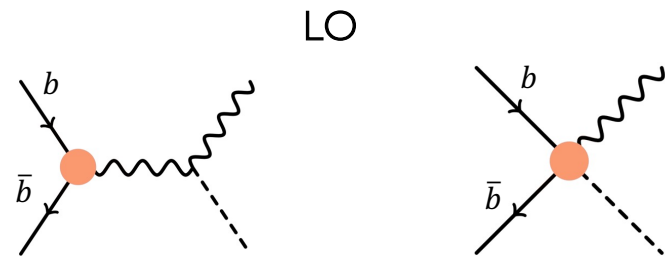
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Probed by $gg \rightarrow ZH$

Probed by $qq \rightarrow ZH$

- Used $qq \rightarrow ZH$ analysis by Bishara, Englert, Grojean, Panico and Rossia, arXiv:2208.11134.
- Predictions obtained with Madgraph in the presence of one operator at a time.



HL-LHC projected bounds from $pp \rightarrow ZH$

WC [TeV^{-2}]	95% C.L. Bound (5% syst.)
$c_{\varphi Q}^{(3)}$	$[-0.72, 0.57]$
$c_{\varphi Q}^{(-)}$	$[-1.5, 1.1]$
$c_{\varphi t}$	$[-8.1, 19.6]$
$c_{t\varphi}$	$[-19.4, 8.0]$

Probed by $gg \rightarrow ZH$ (loop induced)

Probed by $qq \rightarrow ZH$ (tree-level)



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Compare to global fits of LHC data:

$$|c_{\varphi Q}^{(3)}| \lesssim 0.6 \text{ TeV}^{-2} \quad c_{\varphi t} \in [-13.3, 4.0] \text{ TeV}^{-2}$$

$$|c_{\varphi Q}^{(-)}| \lesssim 2.9 \text{ TeV}^{-2} \quad c_{t\varphi} \in [-2.3, 2.8] \text{ TeV}^{-2}$$

SMEFIT Collaboration, arXiv:2105.00006

Probed by $gg \rightarrow ZH$ (loop induced)
 Probed by $qq \rightarrow ZH$ (tree-level)

Competitive against current bounds

→ Motivates precision measurements and inclusion in global fits.



Thank you!

