

ALESSANDRO BACCHETTA, PAVIA U. AND INFN

**MULTIDIMENSIONAL STRUCTURE OF THE
PROTON AND OPPORTUNITIES AT A NEW
ELECTRON ION COLLIDER**

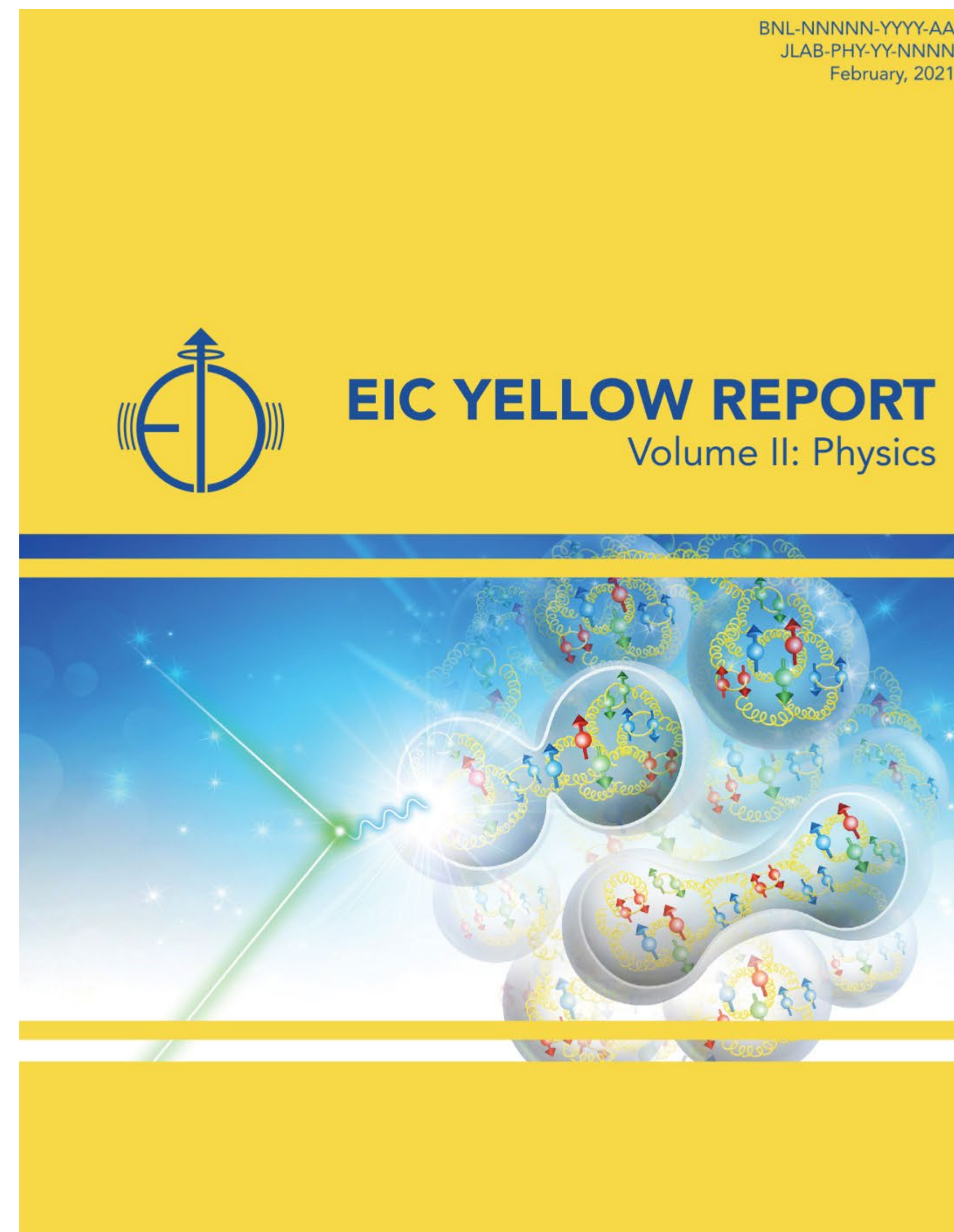
THE EIC PROJECT

The EIC is a new electron-ion collider to be built by 2035 at Brookhaven National Laboratory.

Its main goal is to study the structure of nucleons and nuclei.

<https://www.eicug.org>

[arXiv:2103.05419](https://arxiv.org/abs/2103.05419)



A NEW ERA OF DISCOVERY

THE 2023 LONG RANGE PLAN FOR NUCLEAR SCIENCE

2023 | VERSION 1.2



RECOMMENDATION 3

We recommend the expeditious completion of the EIC as the highest priority for facility construction.

A NEW ERA OF DISCOVERY

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“

To achieve the scientific goals of the EIC, a parallel investment in quantum chromodynamics (QCD) theory is essential,

Progress in theory and computing has already helped to drive and refine the physics program of the EIC.

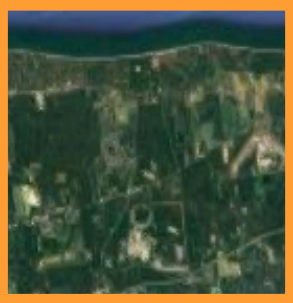
To maximize the scientific impact of the facility and to prepare for the precision expected at the EIC, theory must advance on multiple fronts, and new collaborative efforts are required.



New York



New York

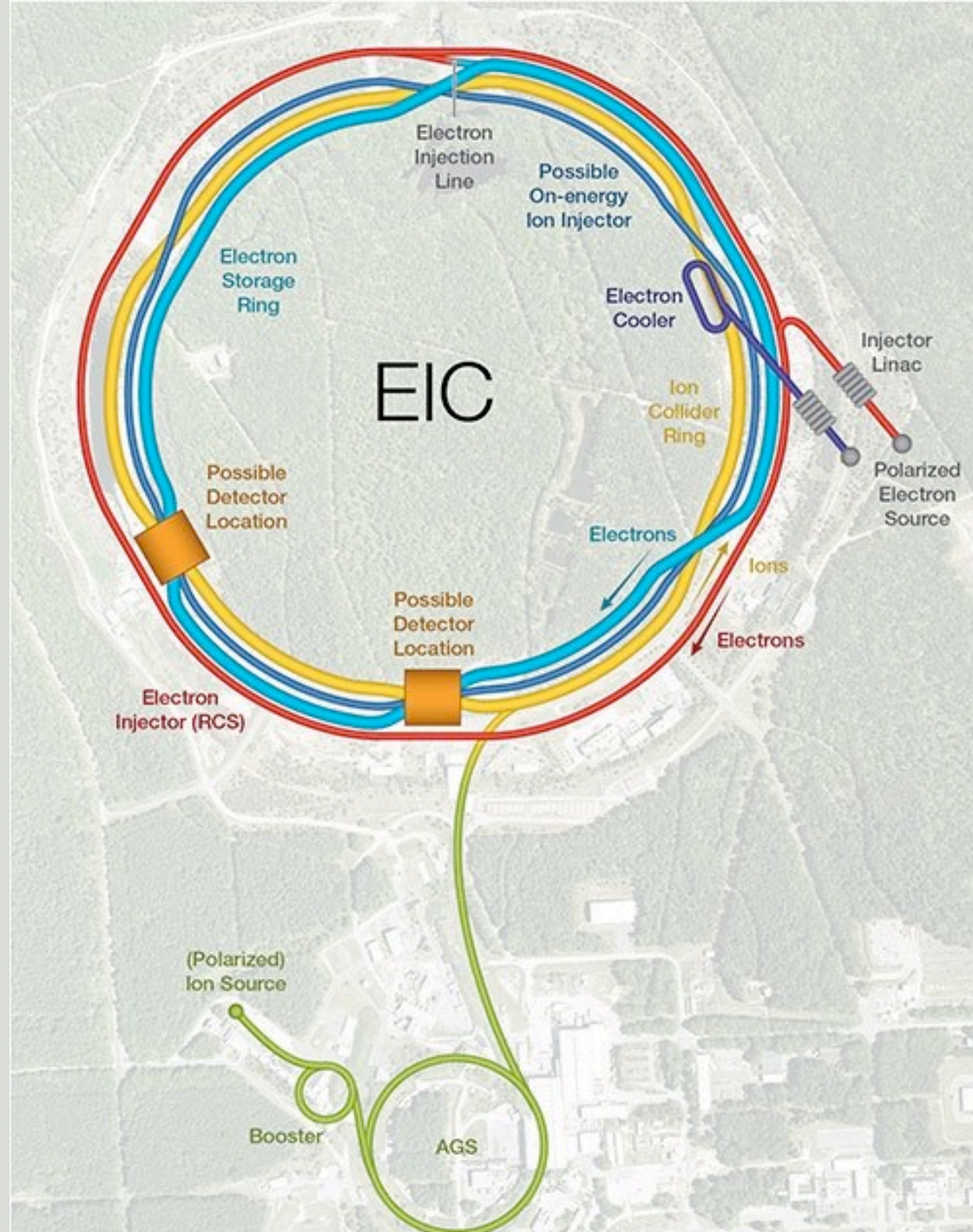


Brookhaven National Lab

Partnership:

BROOKHAVEN
NATIONAL LABORATORY

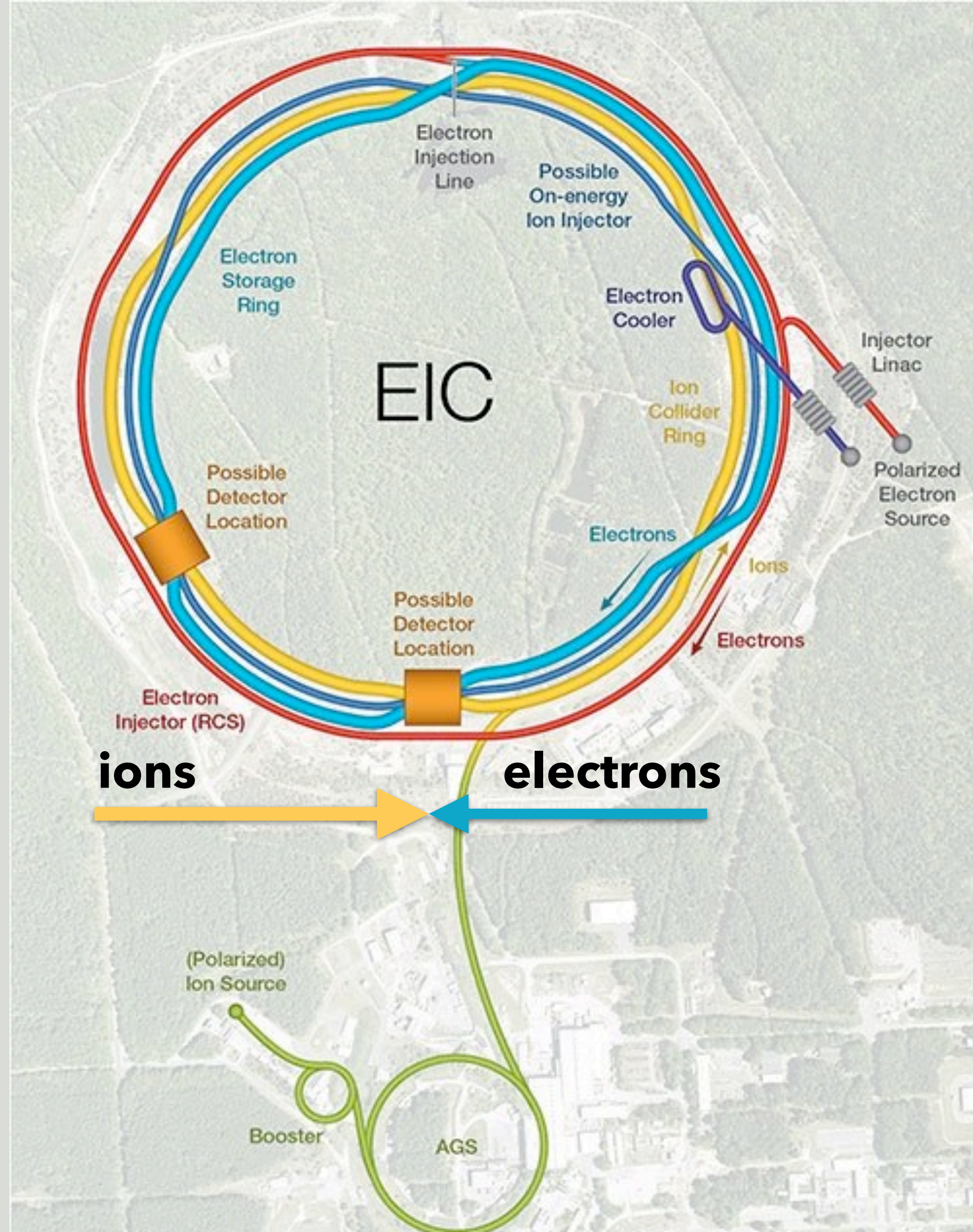
Jefferson Lab



Partnership:

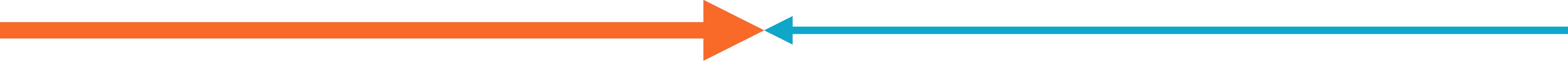
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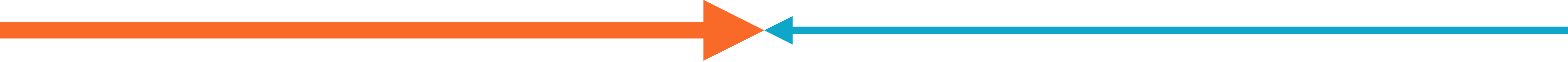
protons to uranium

electrons



protons to uranium

electrons

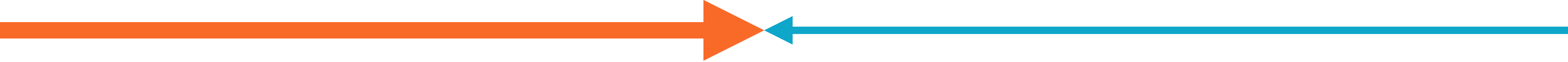


70% polarization

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protons to uranium

electrons



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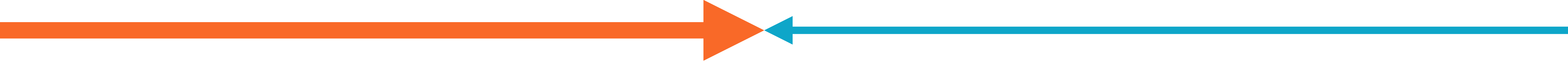
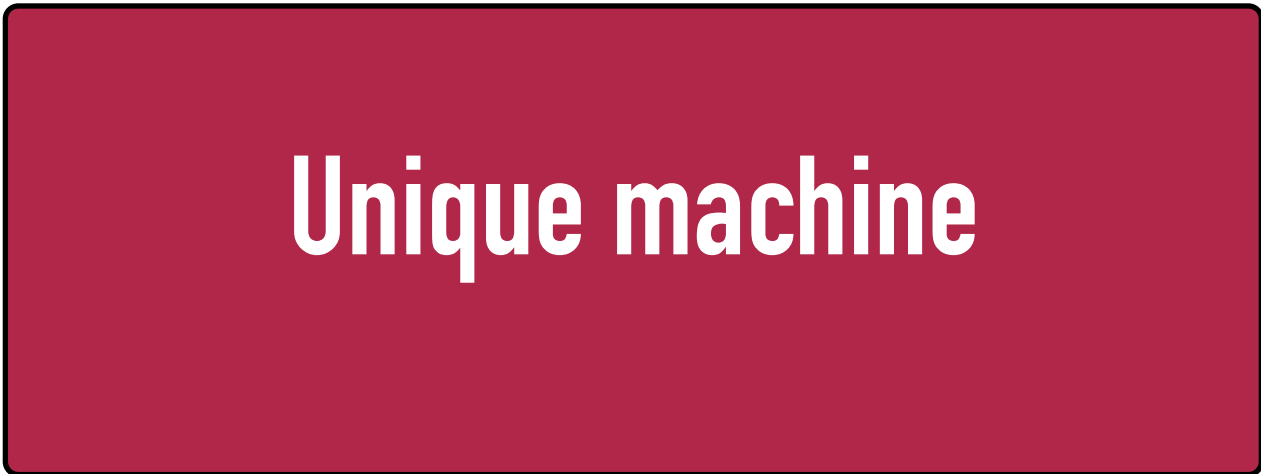
70% polarization

41-275 GeV

5-18 GeV

protons to uranium

electrons

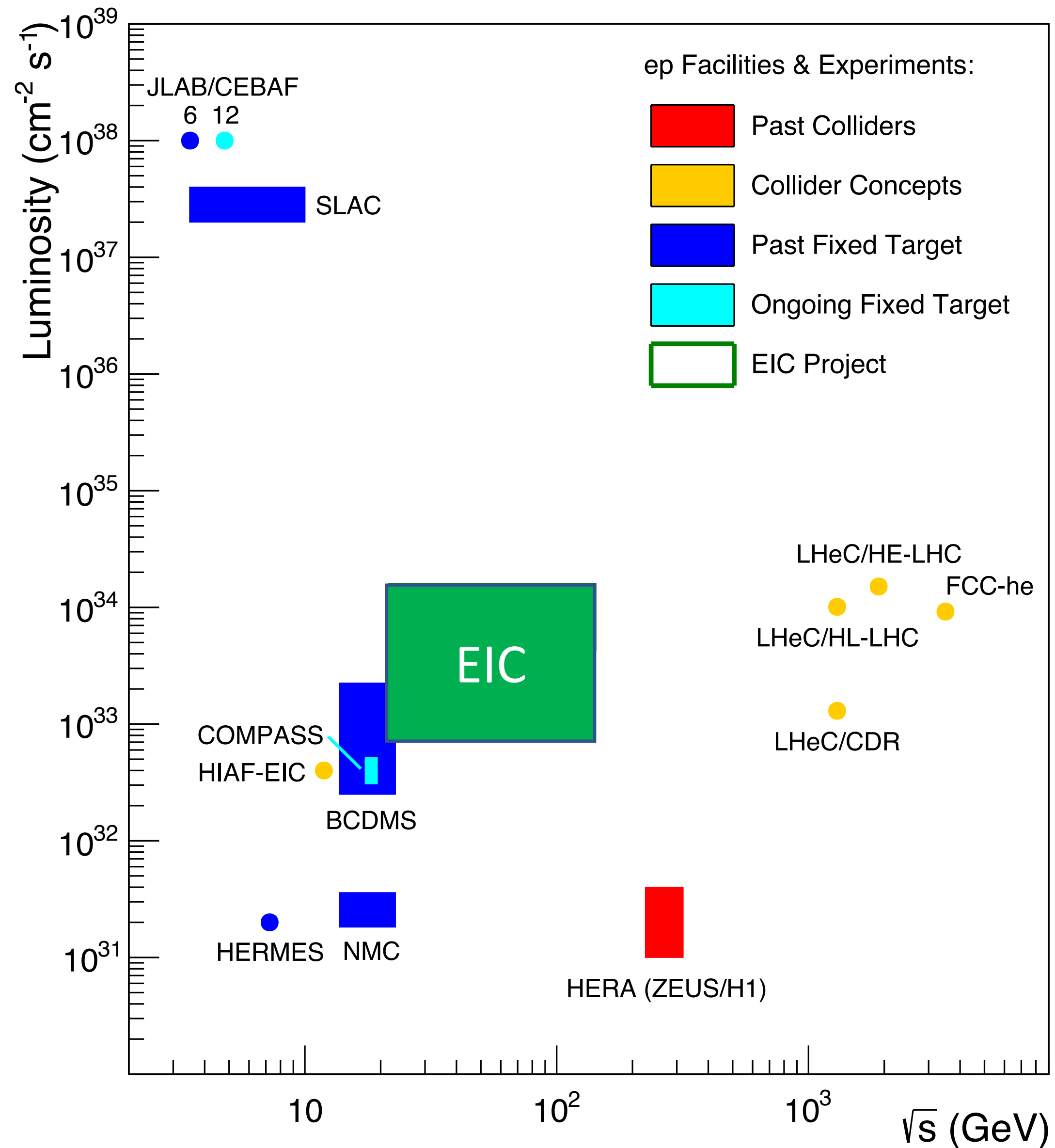


70% polarization

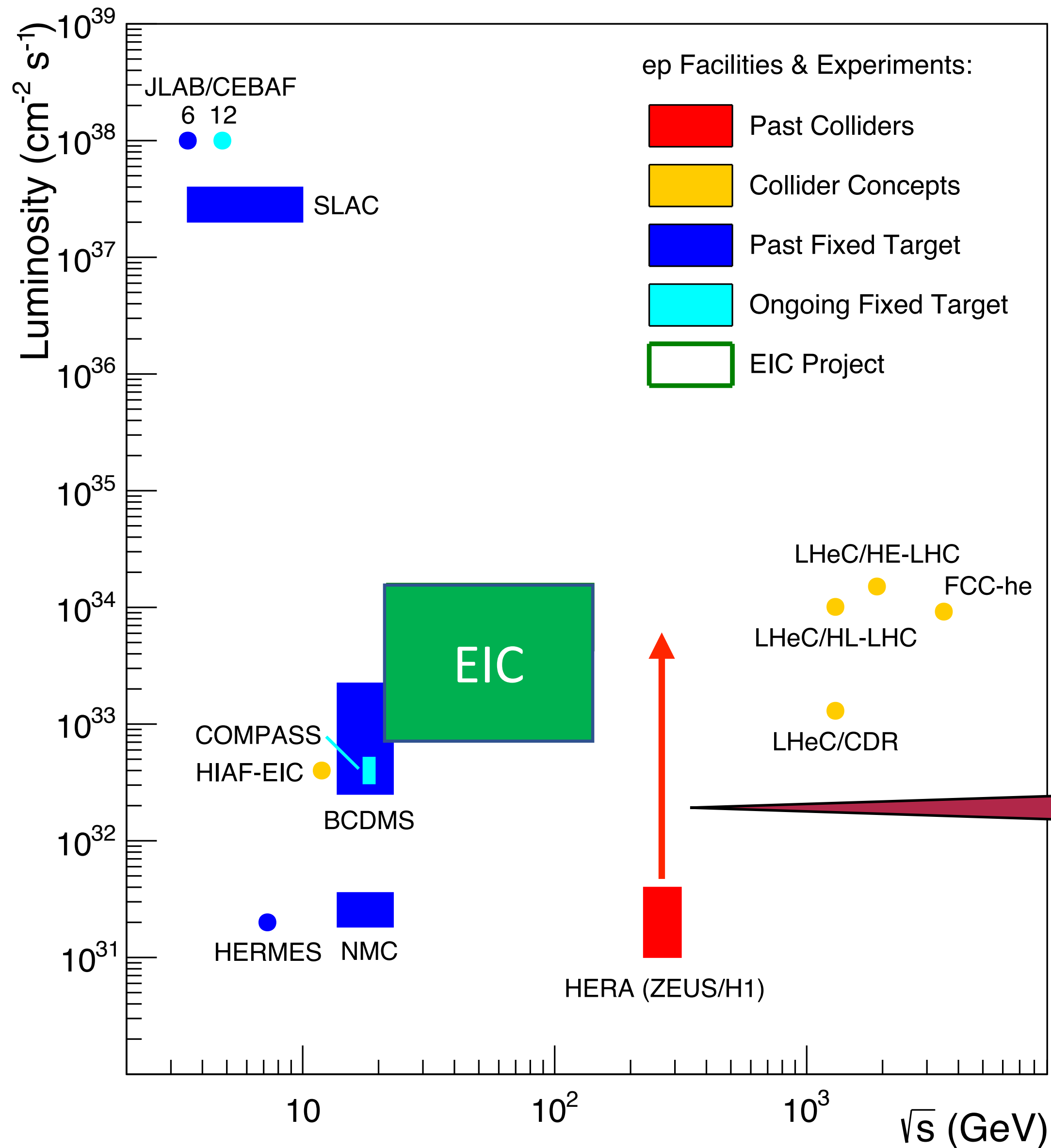
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- **high luminosity** $\sim 10^{33-34} \text{ cm}^{-2} \text{ sec}^{-1}$
- **wide energy range** $\sqrt{s} \sim 29 - 140 \text{ GeV}$

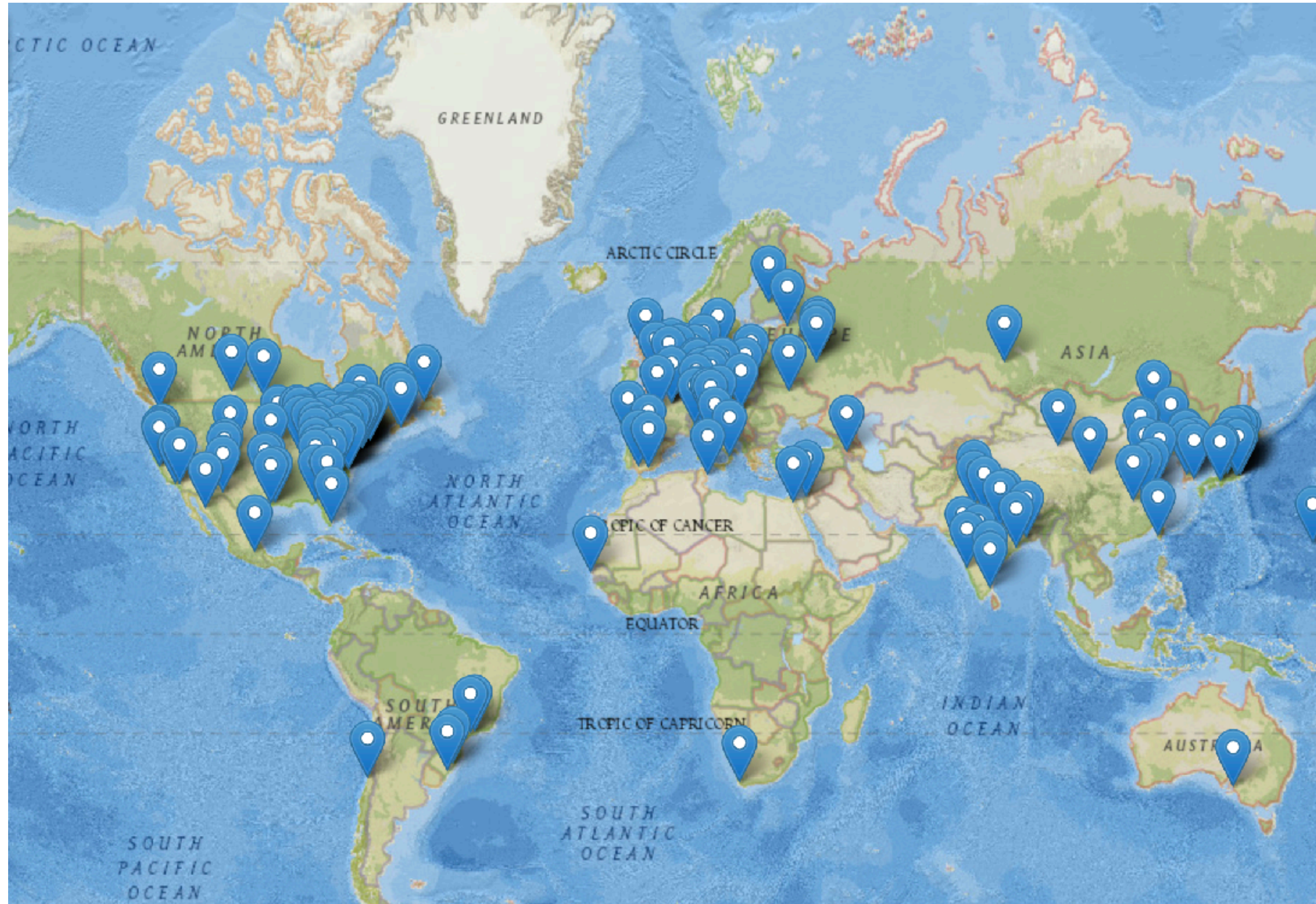


- **high luminosity** $\sim 10^{33-34} \text{ cm}^{-2} \text{ sec}^{-1}$
- **wide energy range** $\sqrt{s} \sim 29 - 140 \text{ GeV}$

100 to 1000 times
the luminosity of HERA!
In principle, in a couple of months
can get the same statistics as HERA

INTERNATIONAL COMMUNITY (EIC USER GROUP)

<https://www.eicug.org>



<https://www.eicug.org>

Phonebook statistics



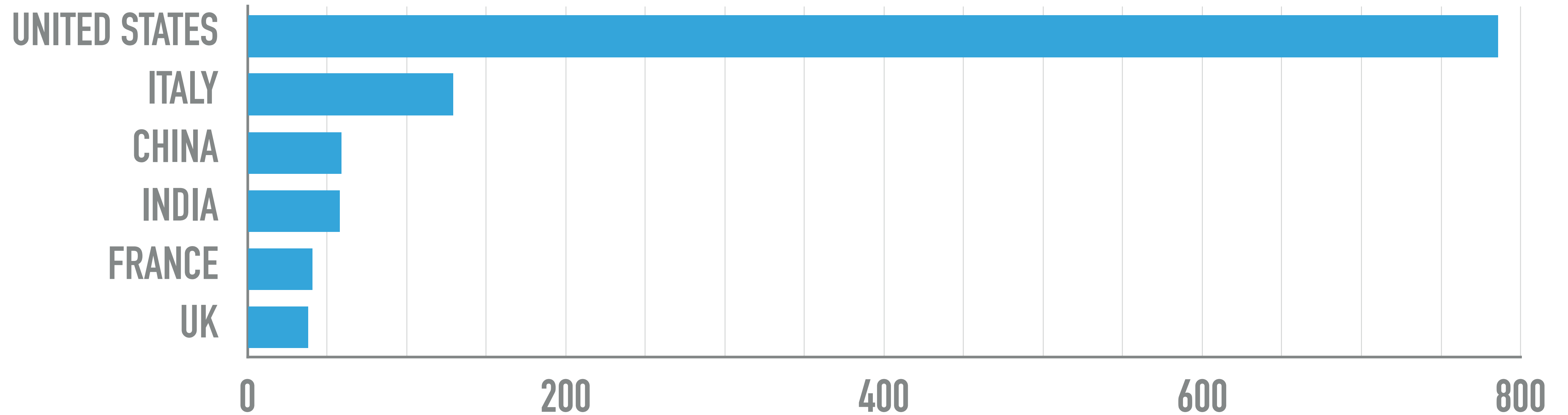
I. EIC User Group:

- **1435** members
- **295** institutions
- **40** countries (**6** world regions)

Experiment Scientists: **905**, Theory Scientists: **363**, Accelerator Scientists: **151**,
Computer Scientists: **10**, Support: **3**, Other: **3**

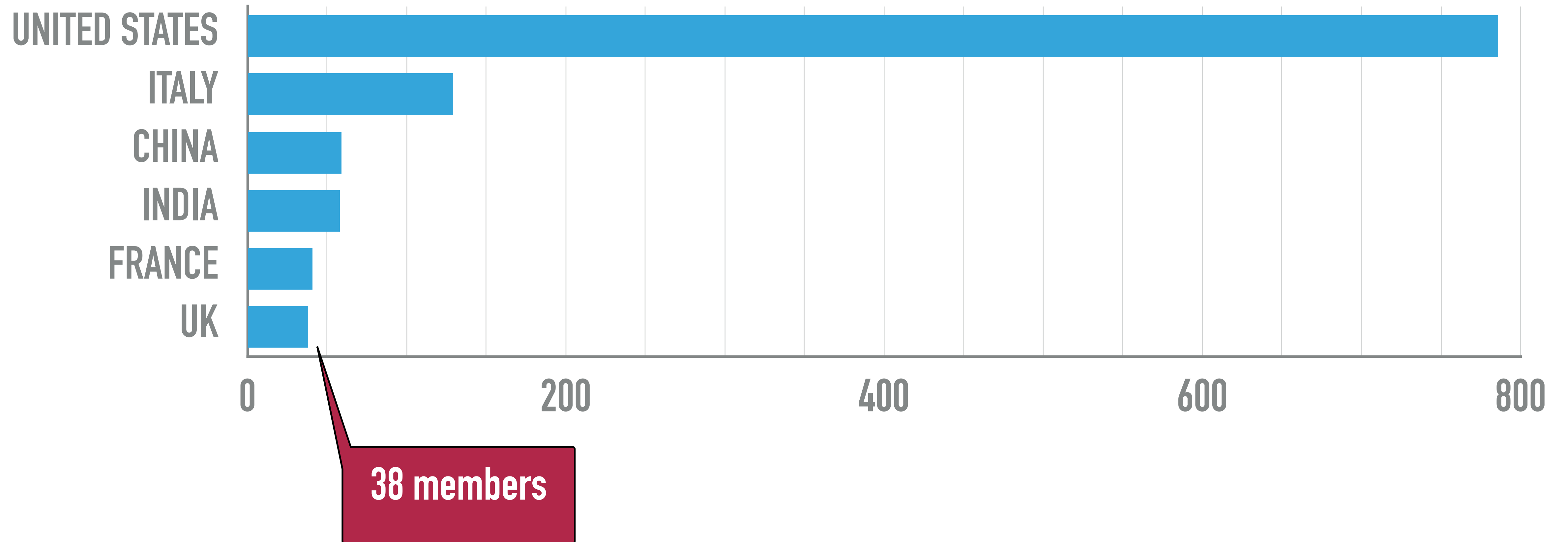
EIC USER GROUP MEMBERS

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PROJECT TIMELINE

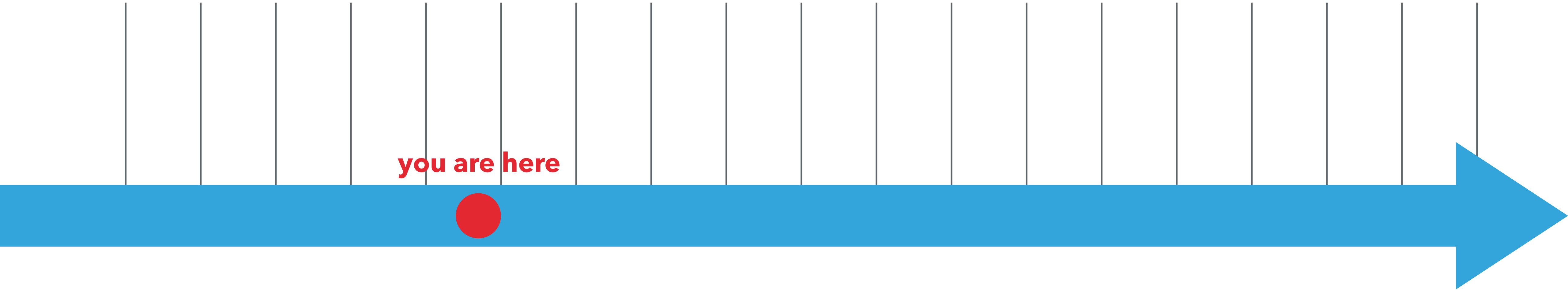
2020

2025

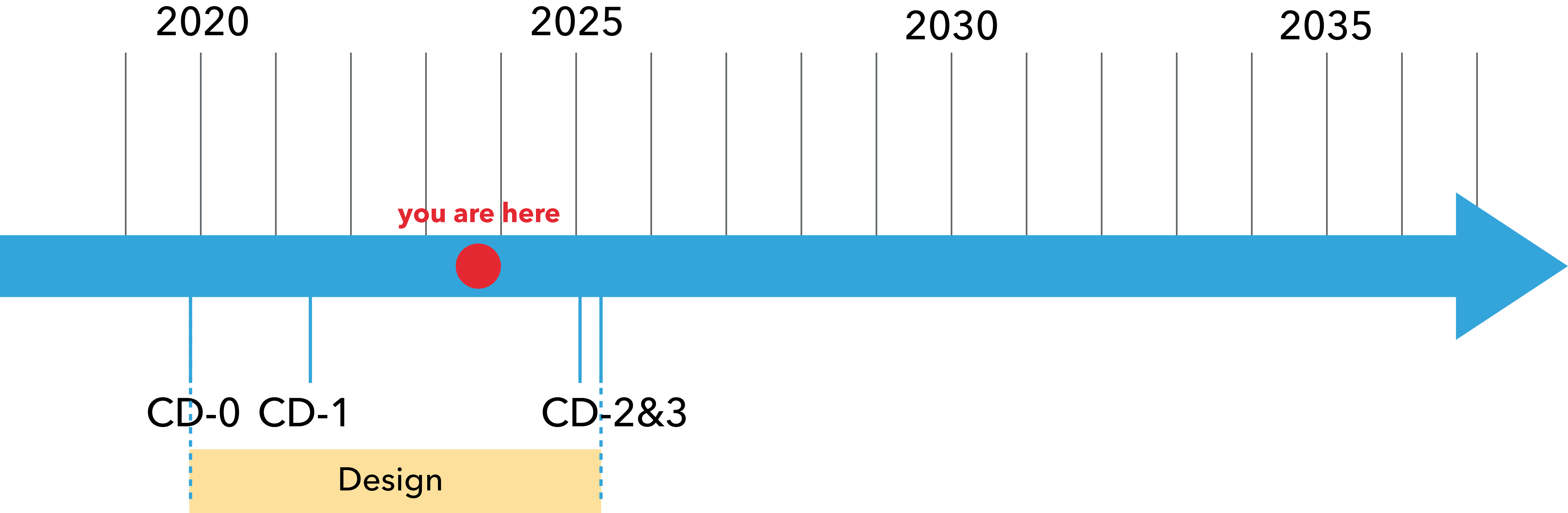
2030

2035

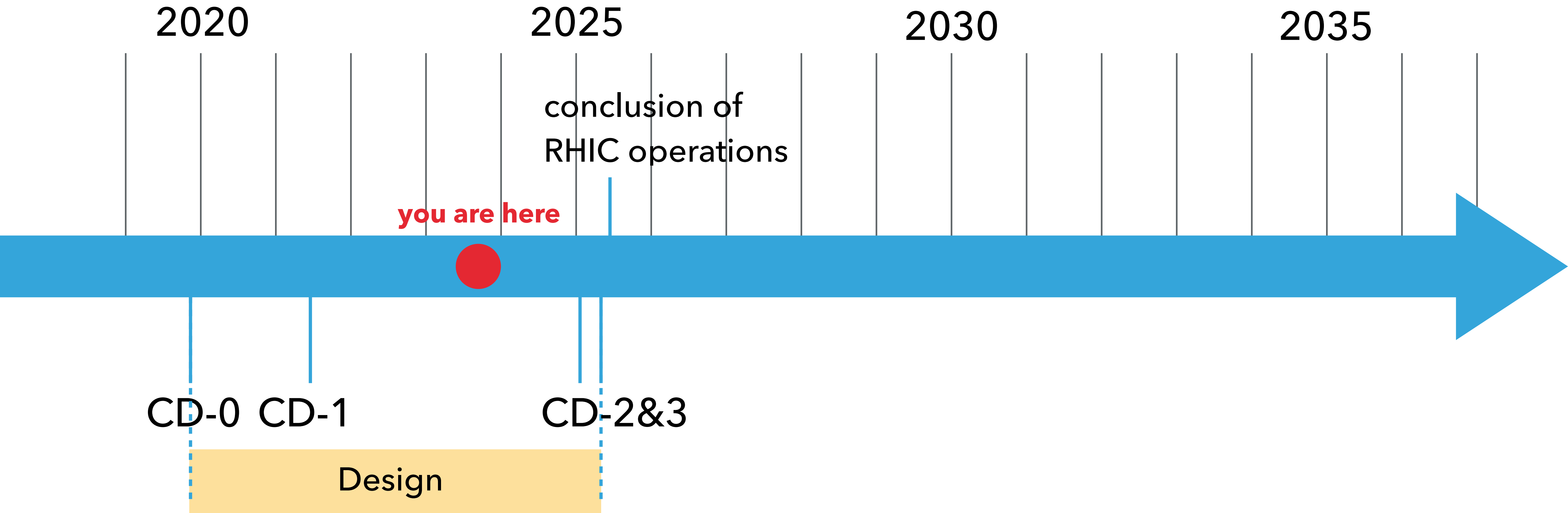
you are here



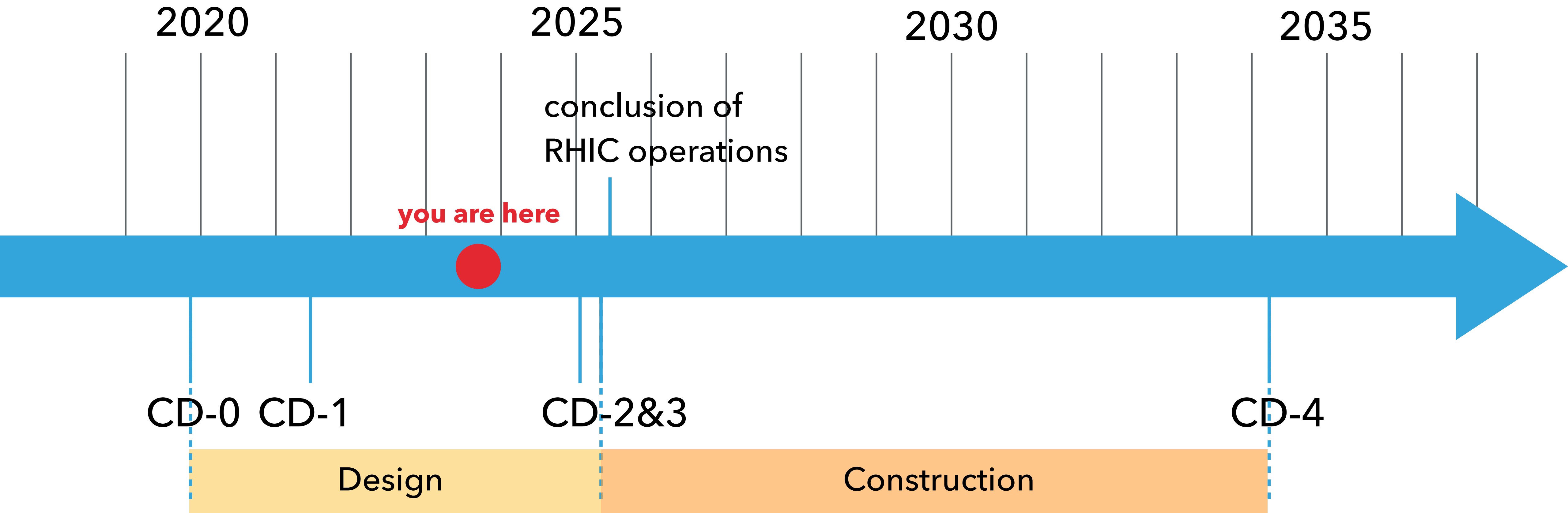
PROJECT TIMELINE



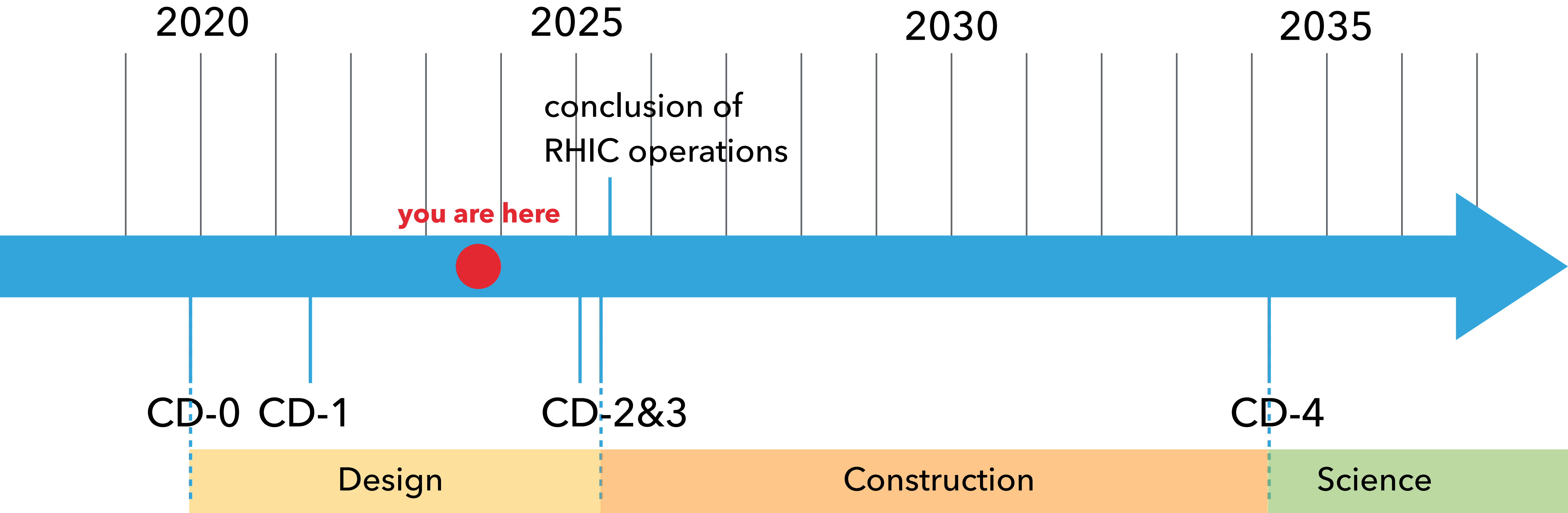
PROJECT TIMELINE



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PROJECT TIMELINE



1.7-2.8 billion US dollars (source: DOE)



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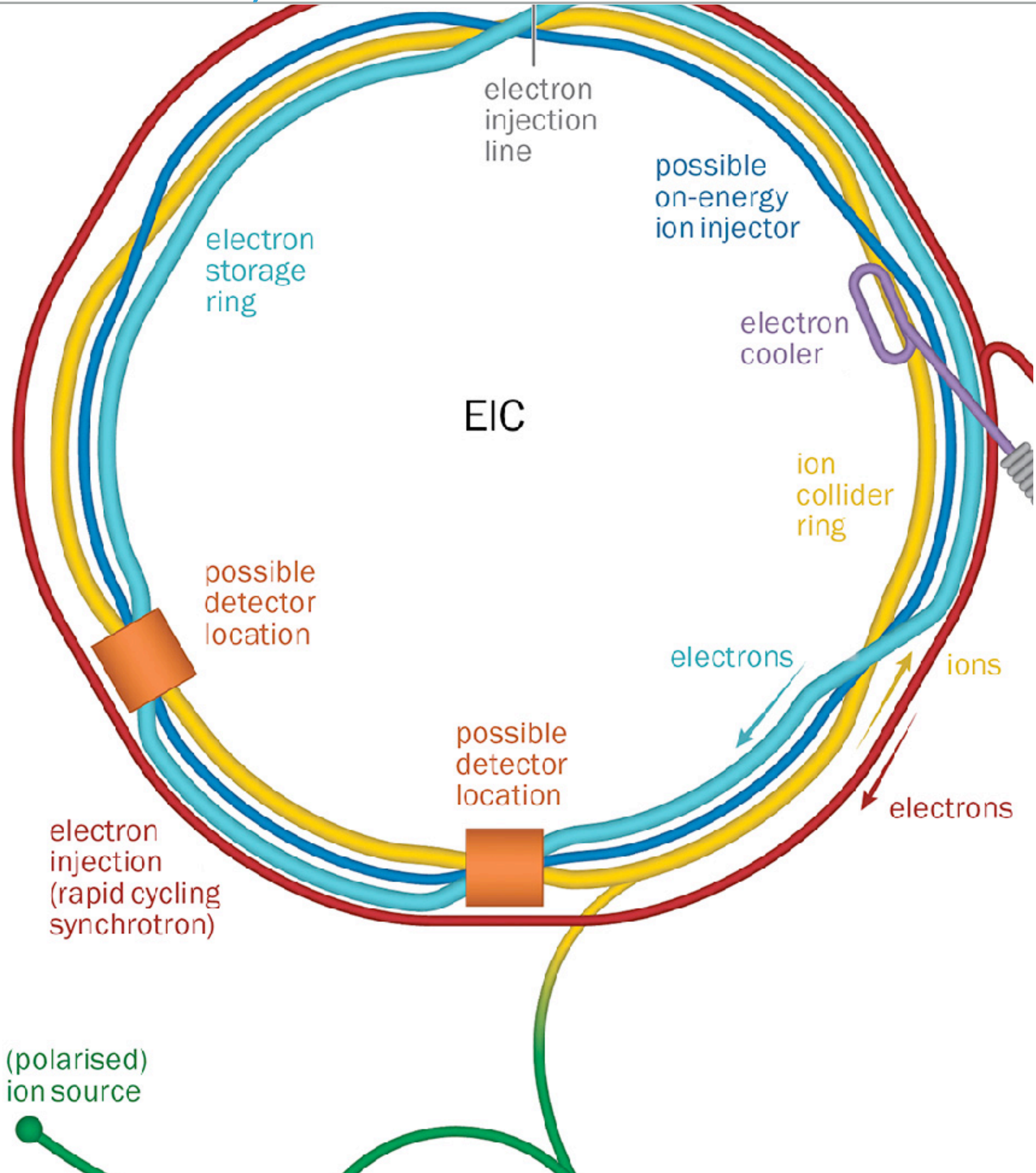
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- James Webb Telescope: 10 billion US dollars (source: Wikipedia)

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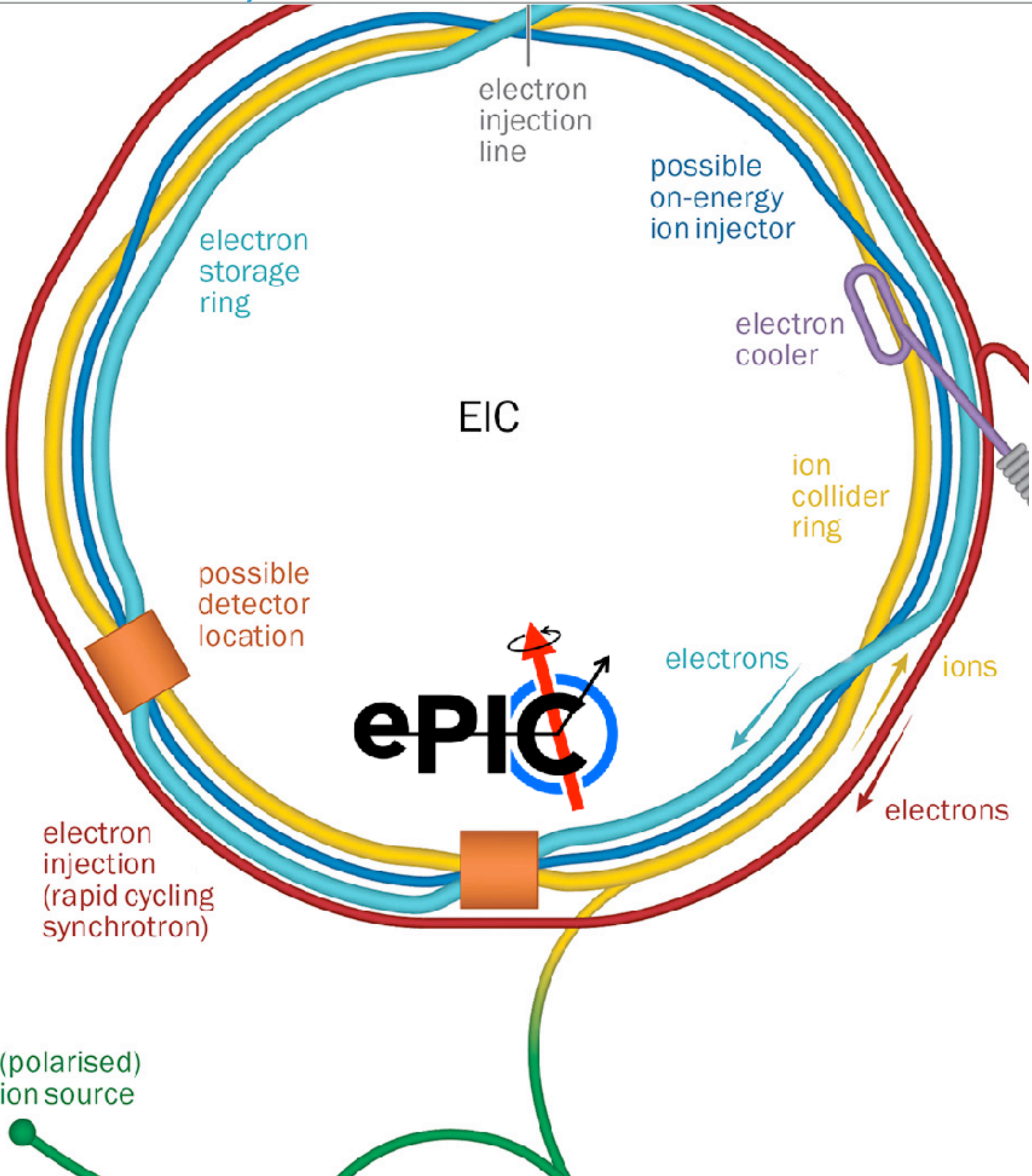


- FAIR: 2.5 billion US dollars (source: Wikipedia)
- Einstein Telescope: 2 billion US dollars (source: Scientific American)
- James Webb Telescope: 10 billion US dollars (source: Wikipedia)
- FCC: 20 billion US dollars (source: Wikipedia)

DETECTOR (OR DETECTORS)

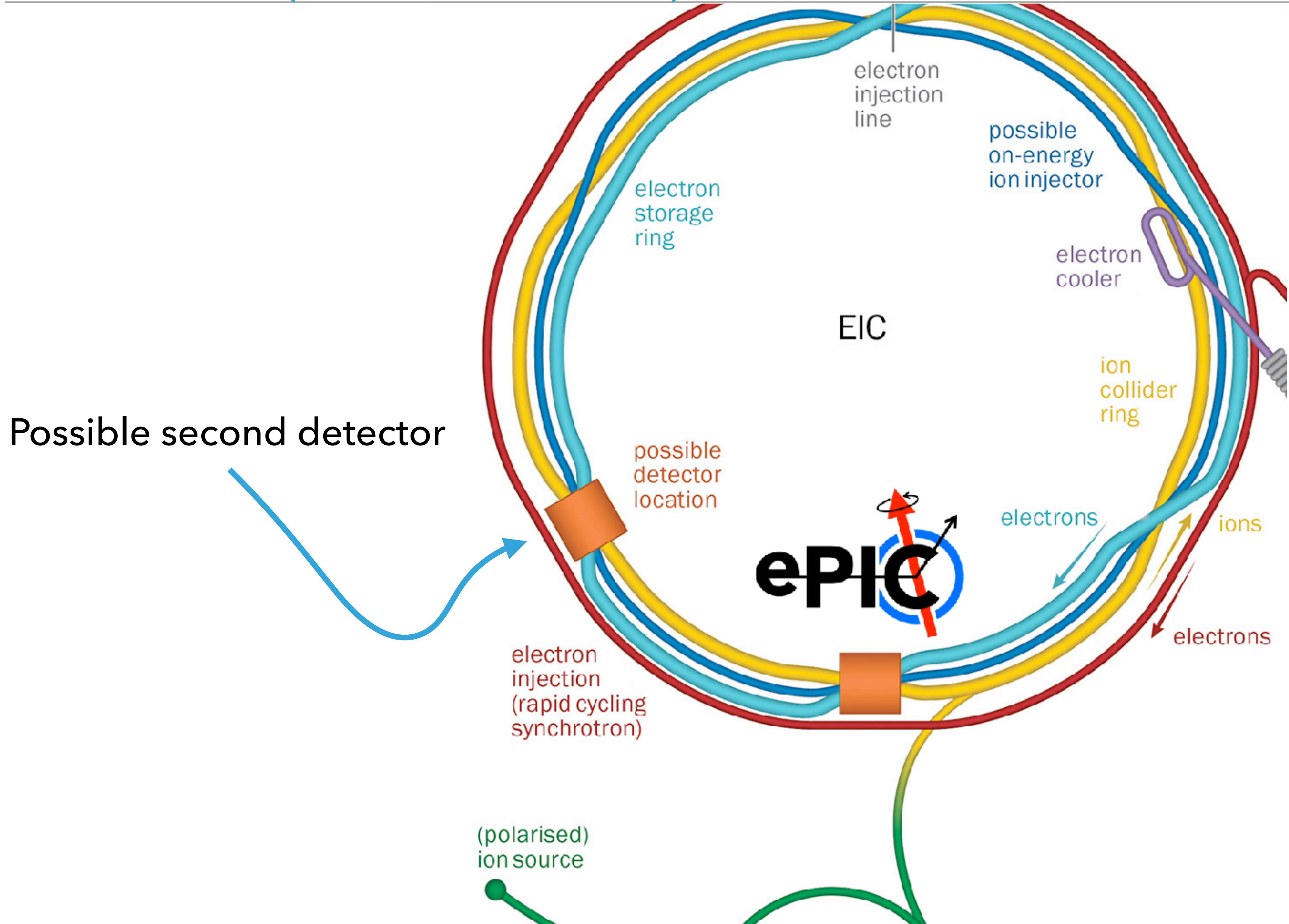


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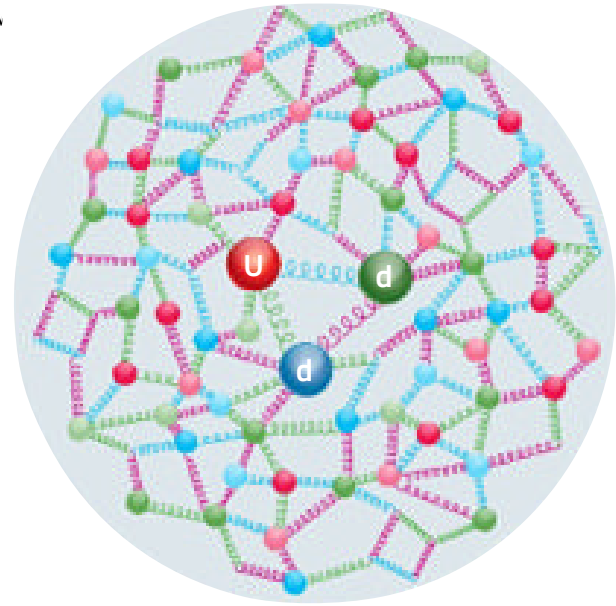
Collaboration
24 countries
171 Institutions
500+ members

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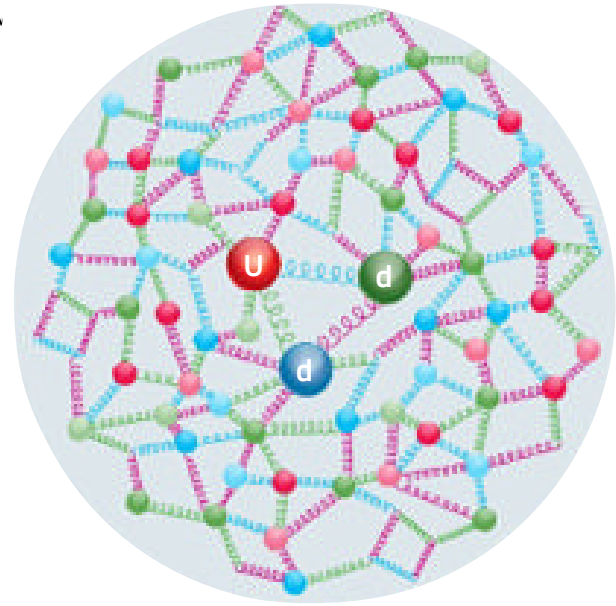


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WHAT DO WE WANT TO DO WITH THE EIC?

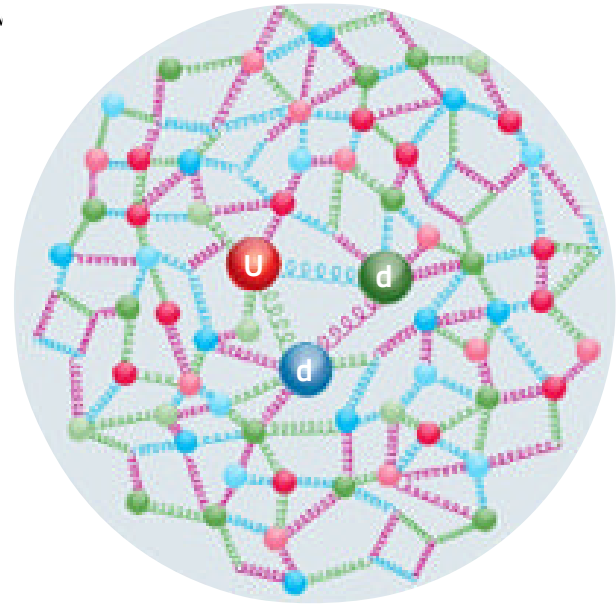


1) How are partons with their spins distributed in space and momentum inside the nucleon, such that its properties emerge from their interactions?



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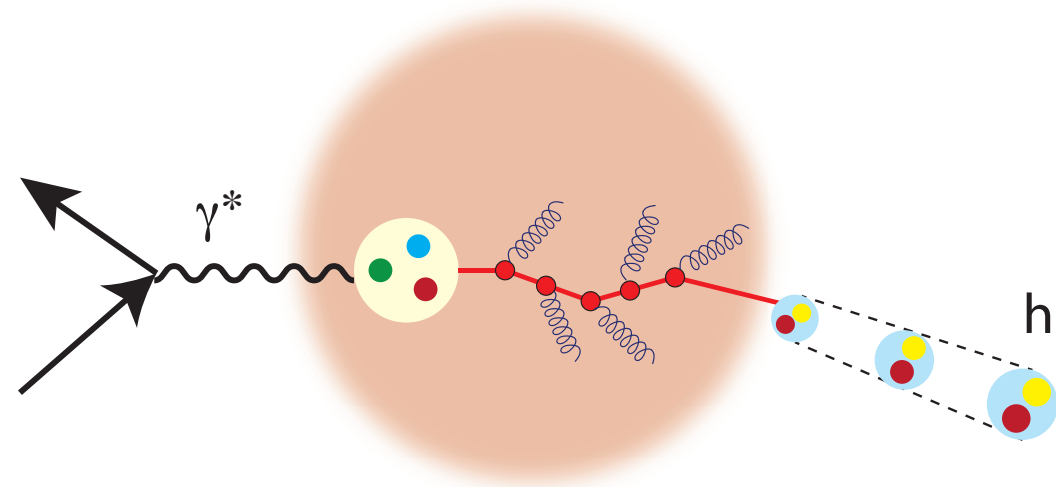
Nucleon “femtography”

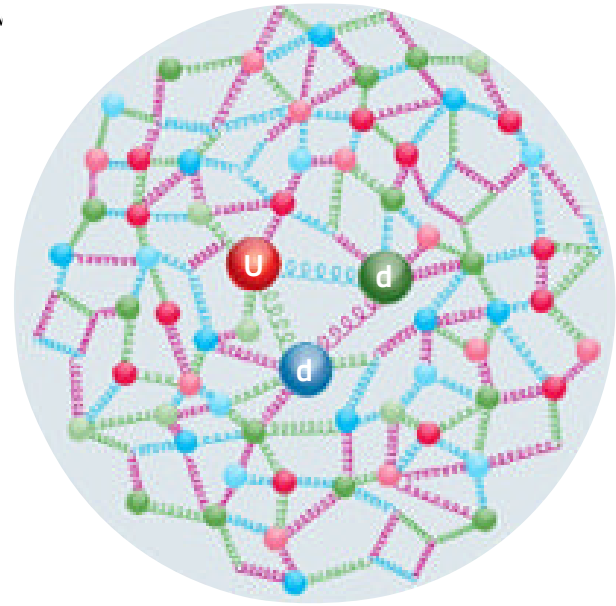


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Nucleon “femtography”

2) How do colored partons propagate and interact with nuclear medium such that eventually colorless hadrons emerge?

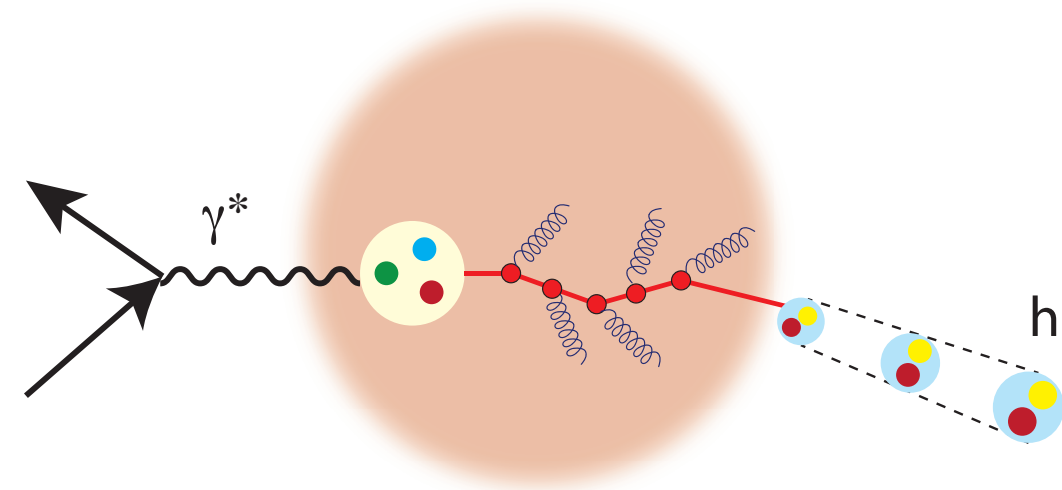




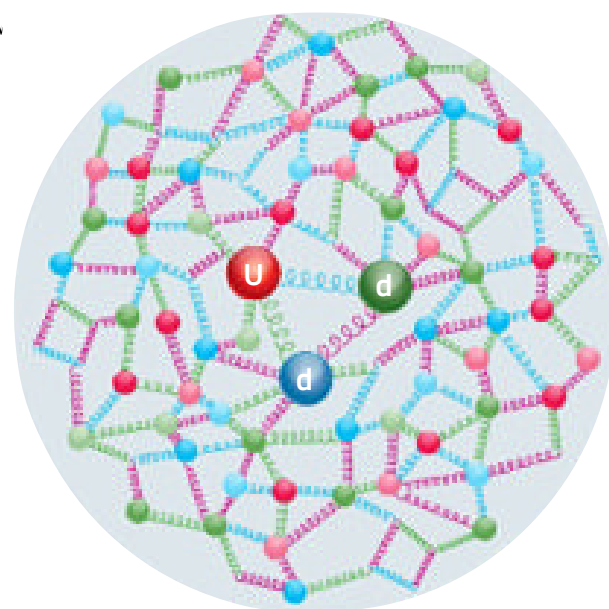
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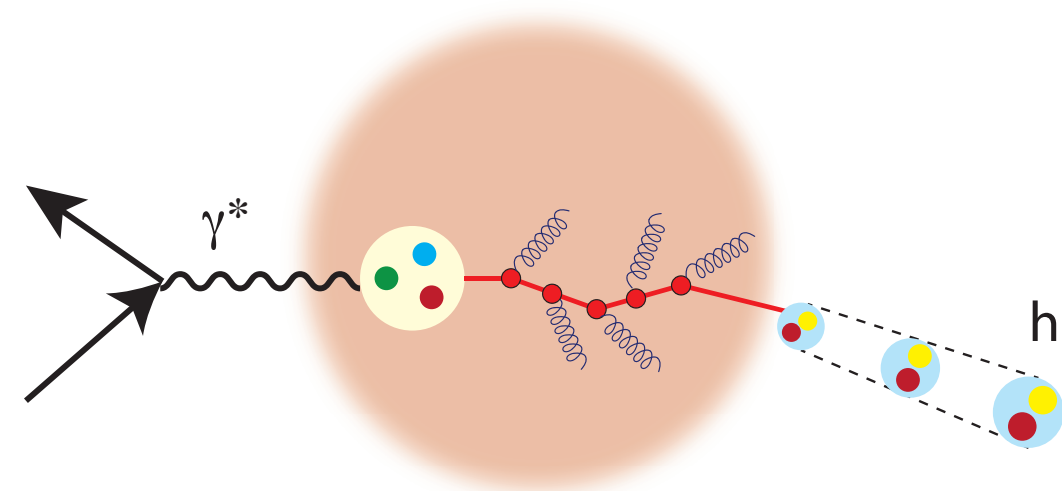


Mechanisms of color confinement and nuclear binding



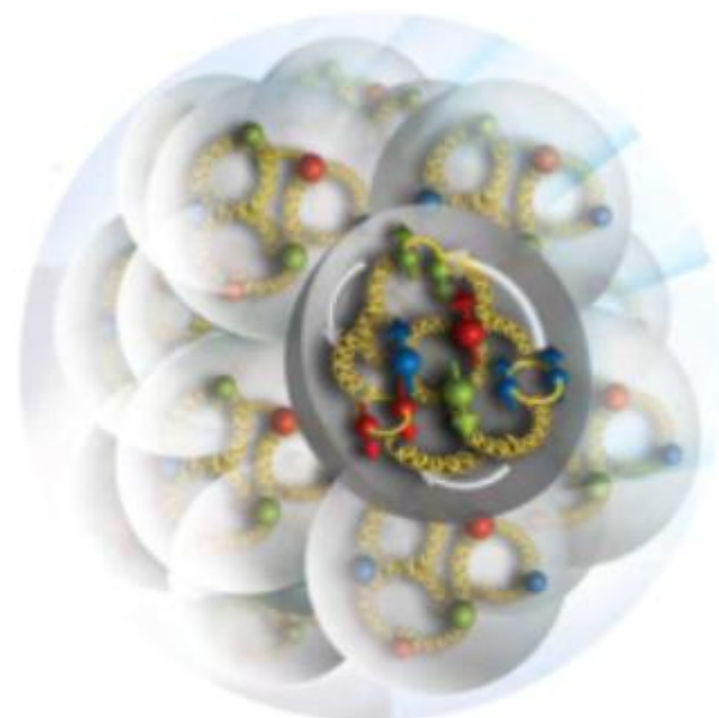
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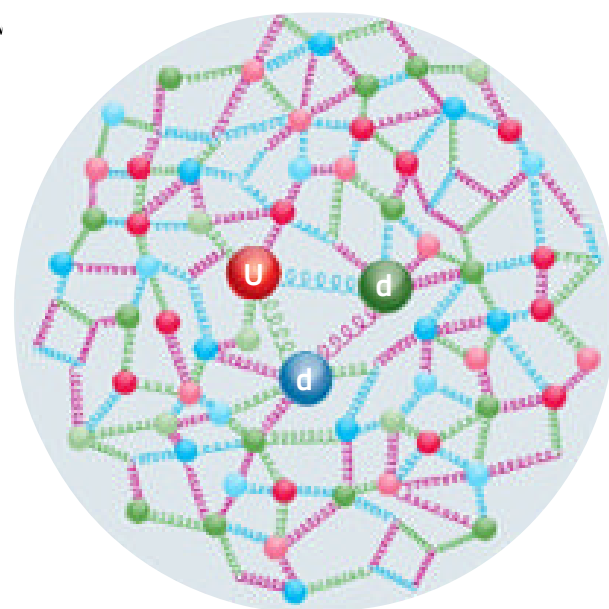


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Mechanisms of color confinement and nuclear binding

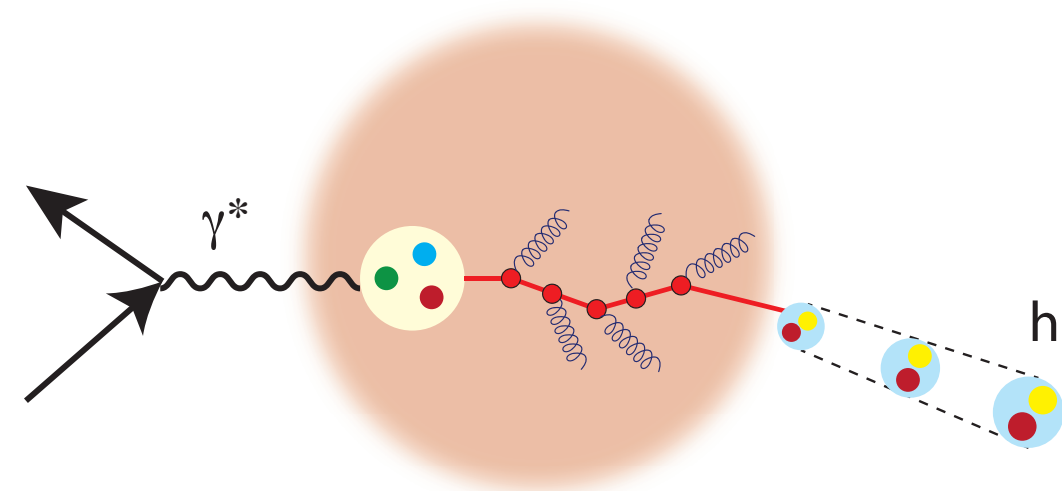


3) Does gluon density saturate at high energy, giving rise to a universal gluonic matter?



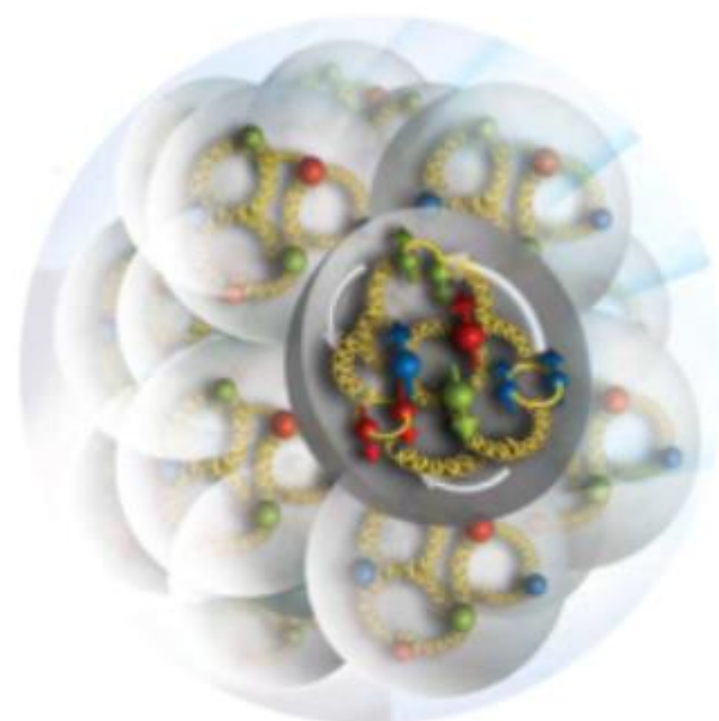
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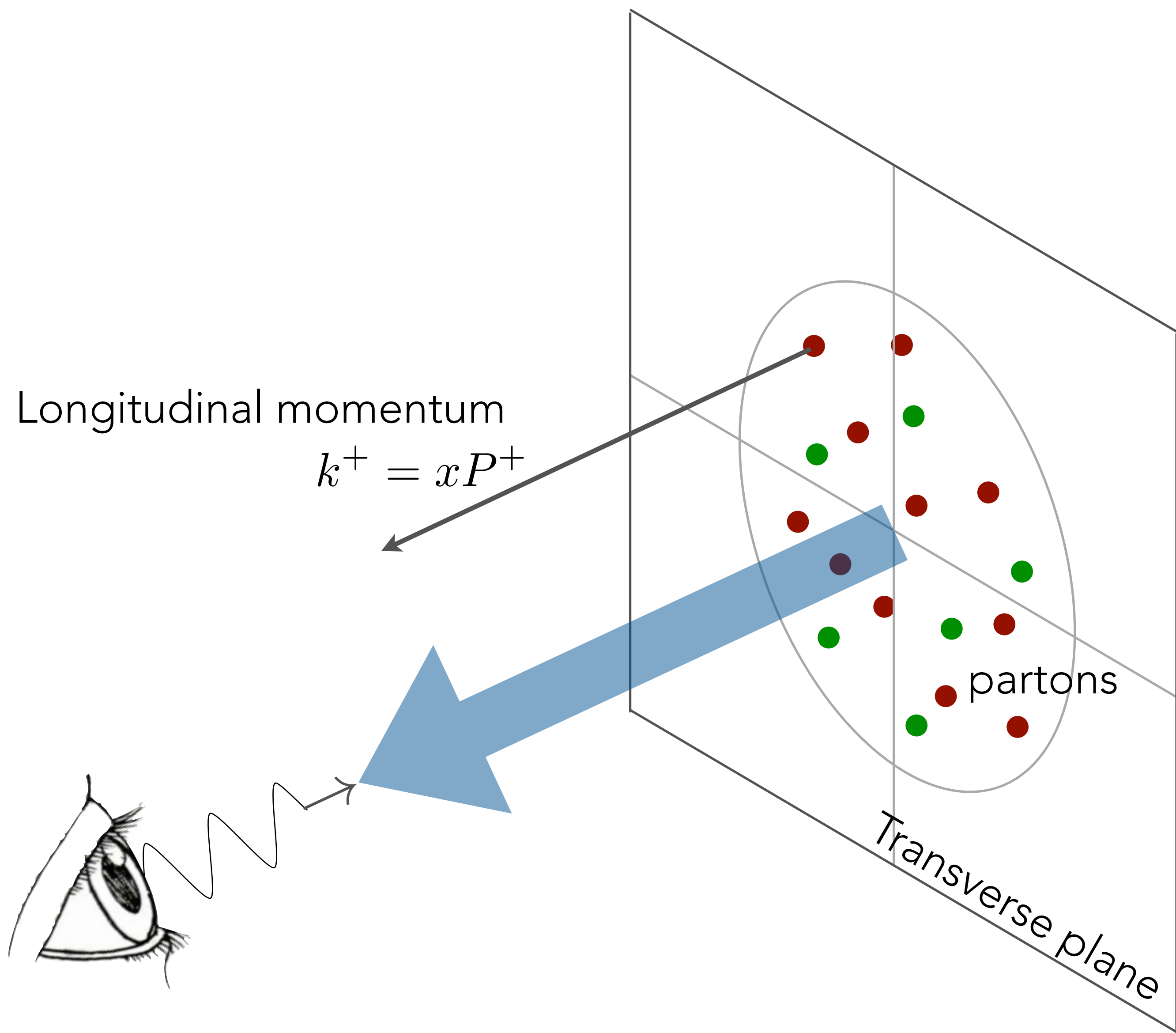
Mechanisms of color confinement and nuclear binding



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Gluon saturation

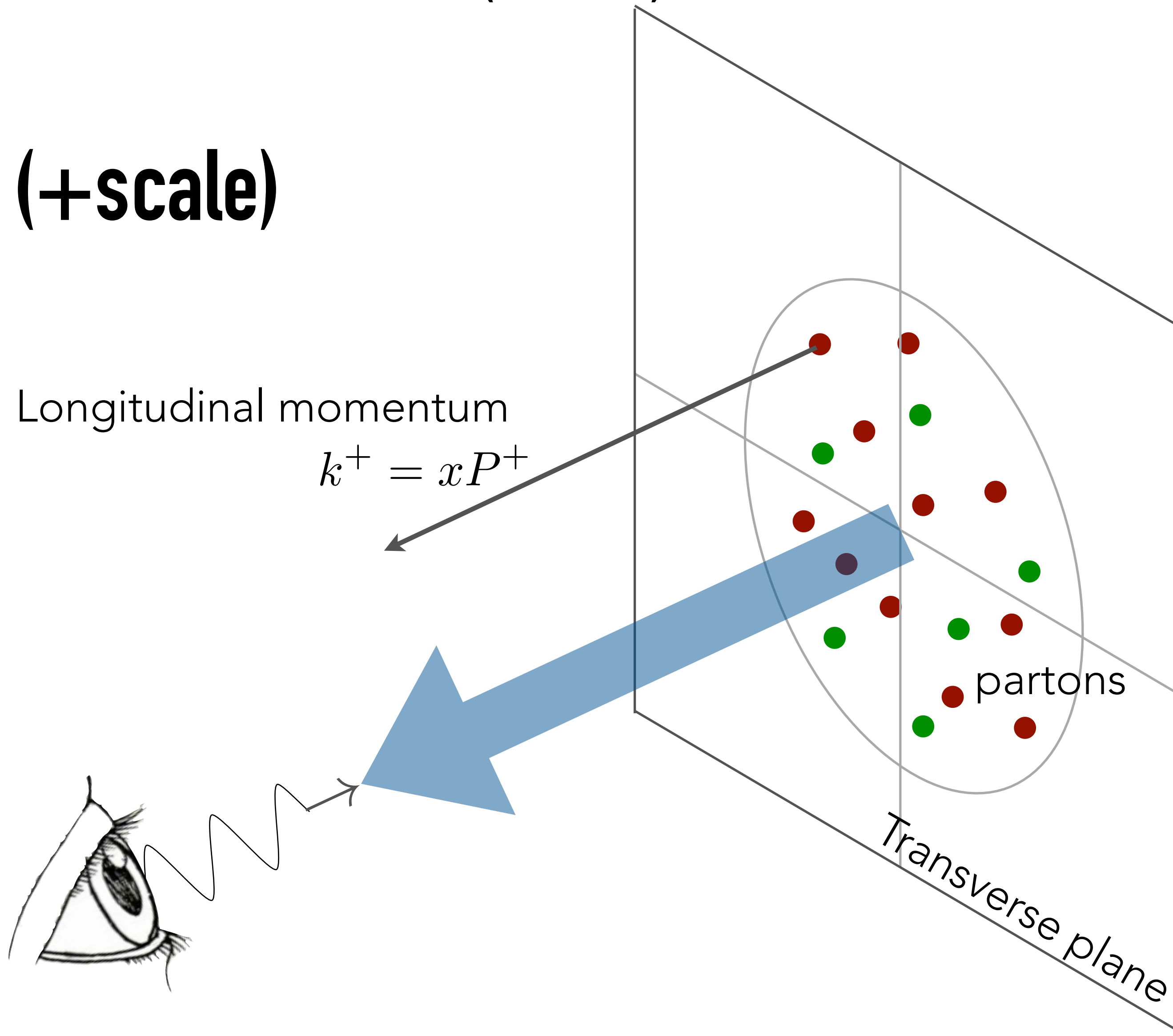
NUCLEON FEMTOGRAPHY

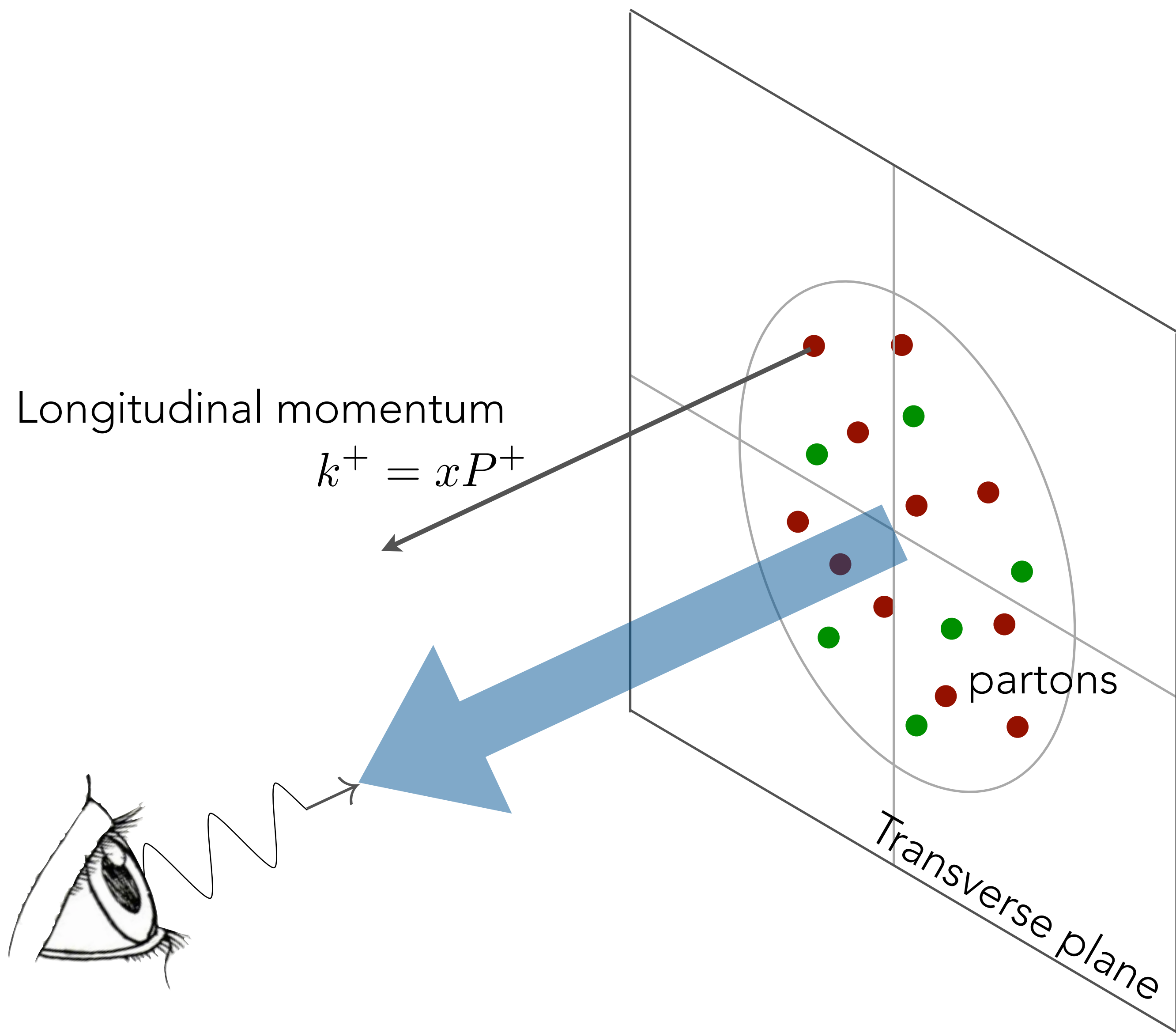


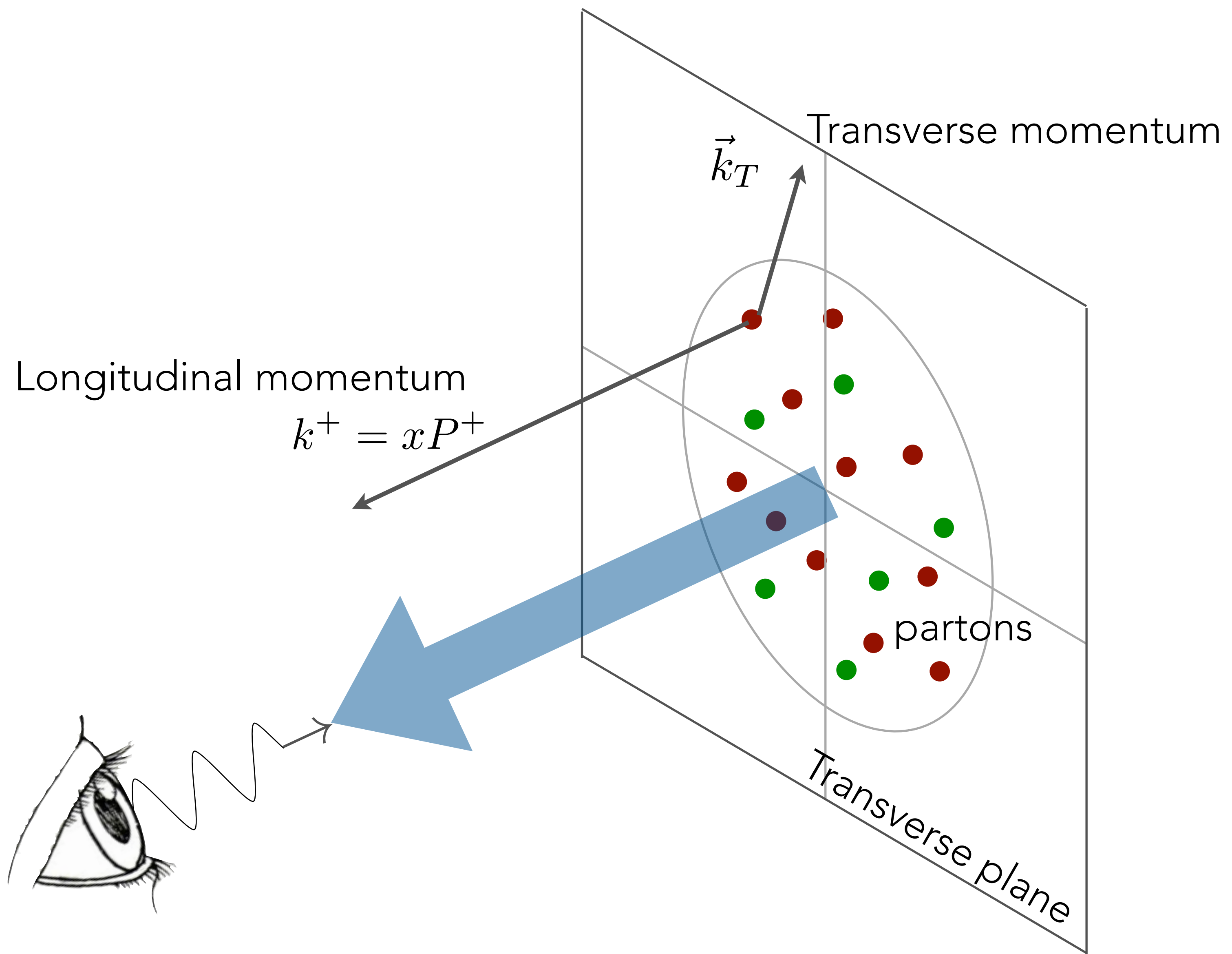
Parton Distribution Functions (PDFs)

$$f(x)$$

1 dimensional (+scale)



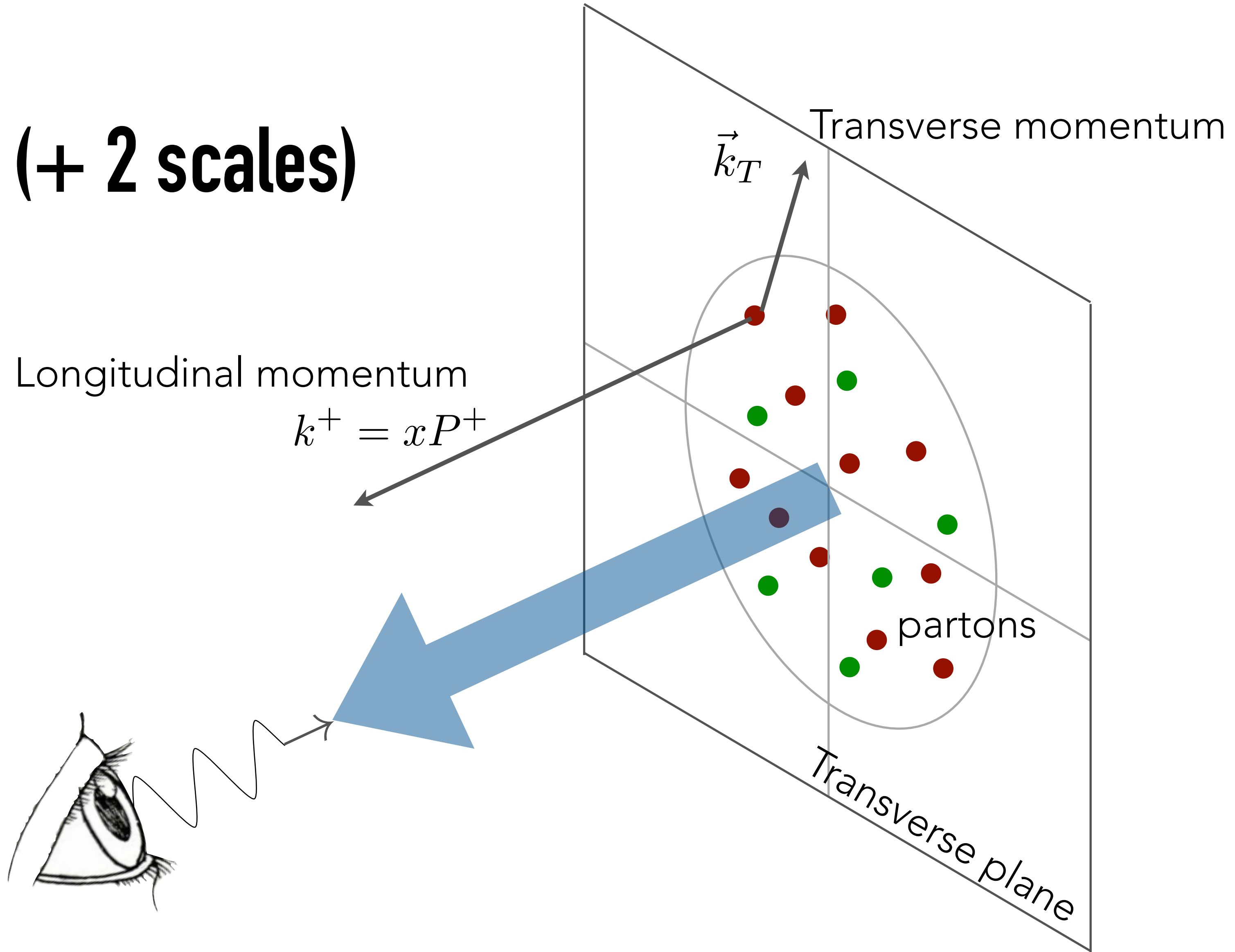


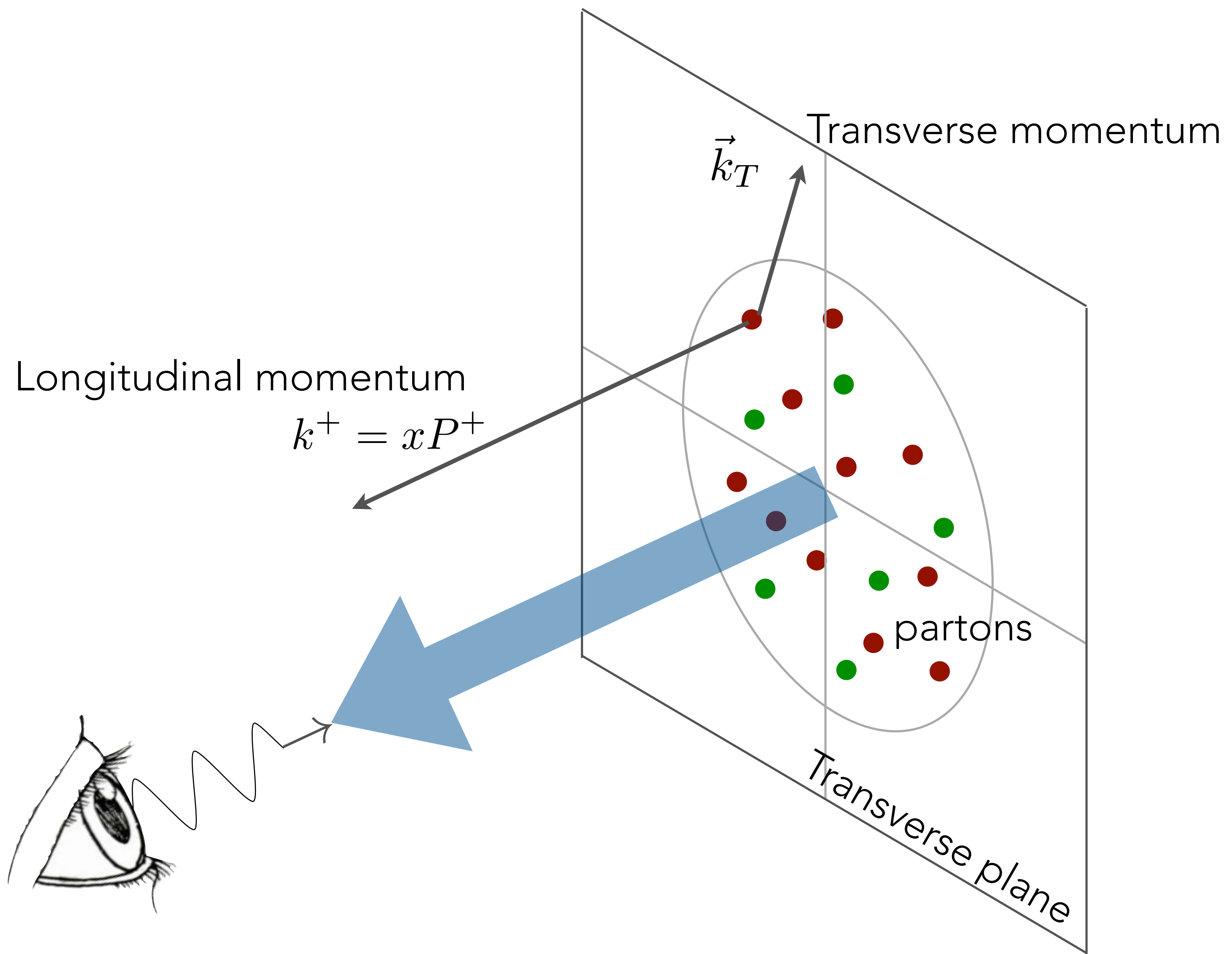


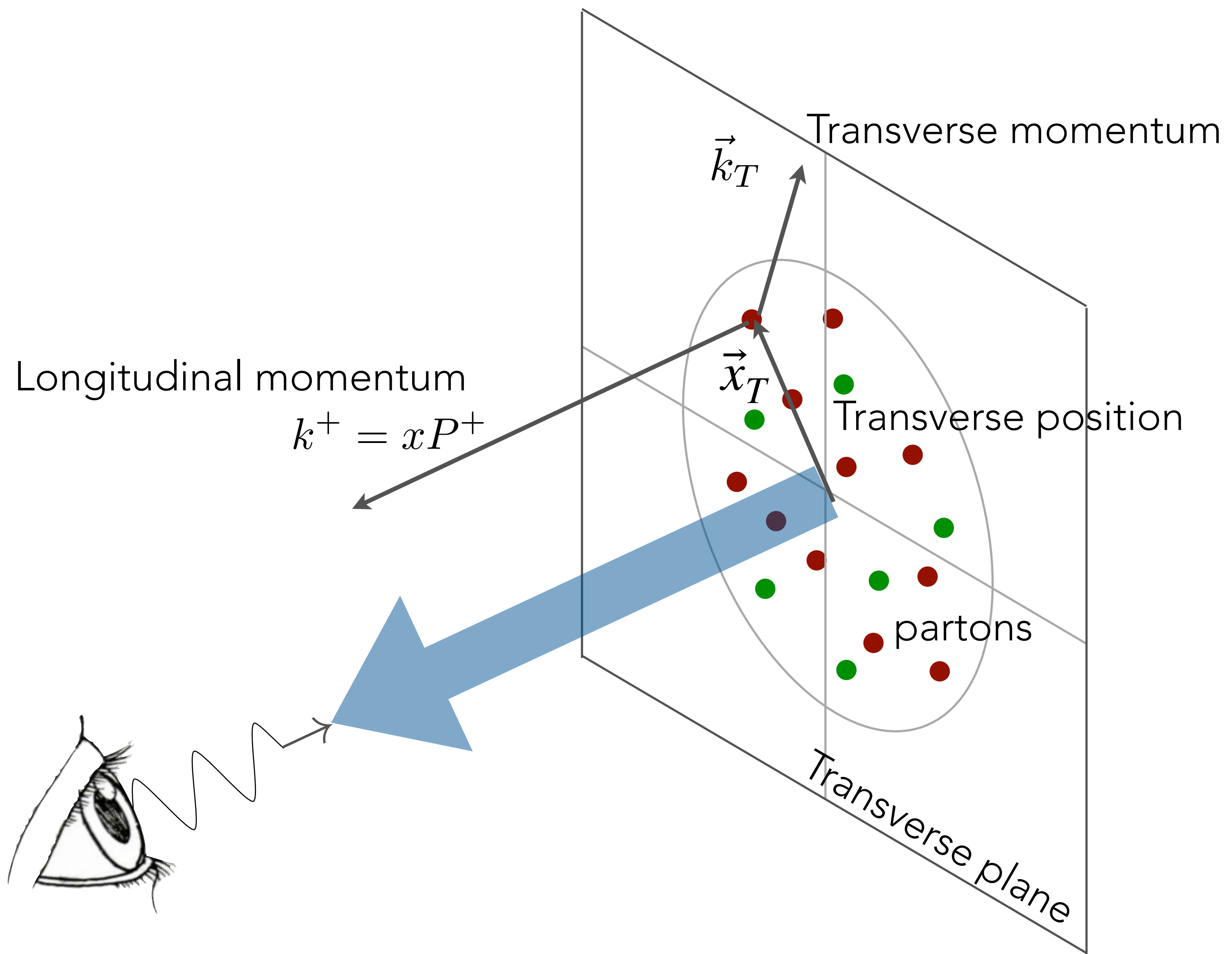
Transverse-Momentum Distributions (TMDs)

$$f(x, \vec{k}_T)$$

3 dimensional (+ 2 scales)



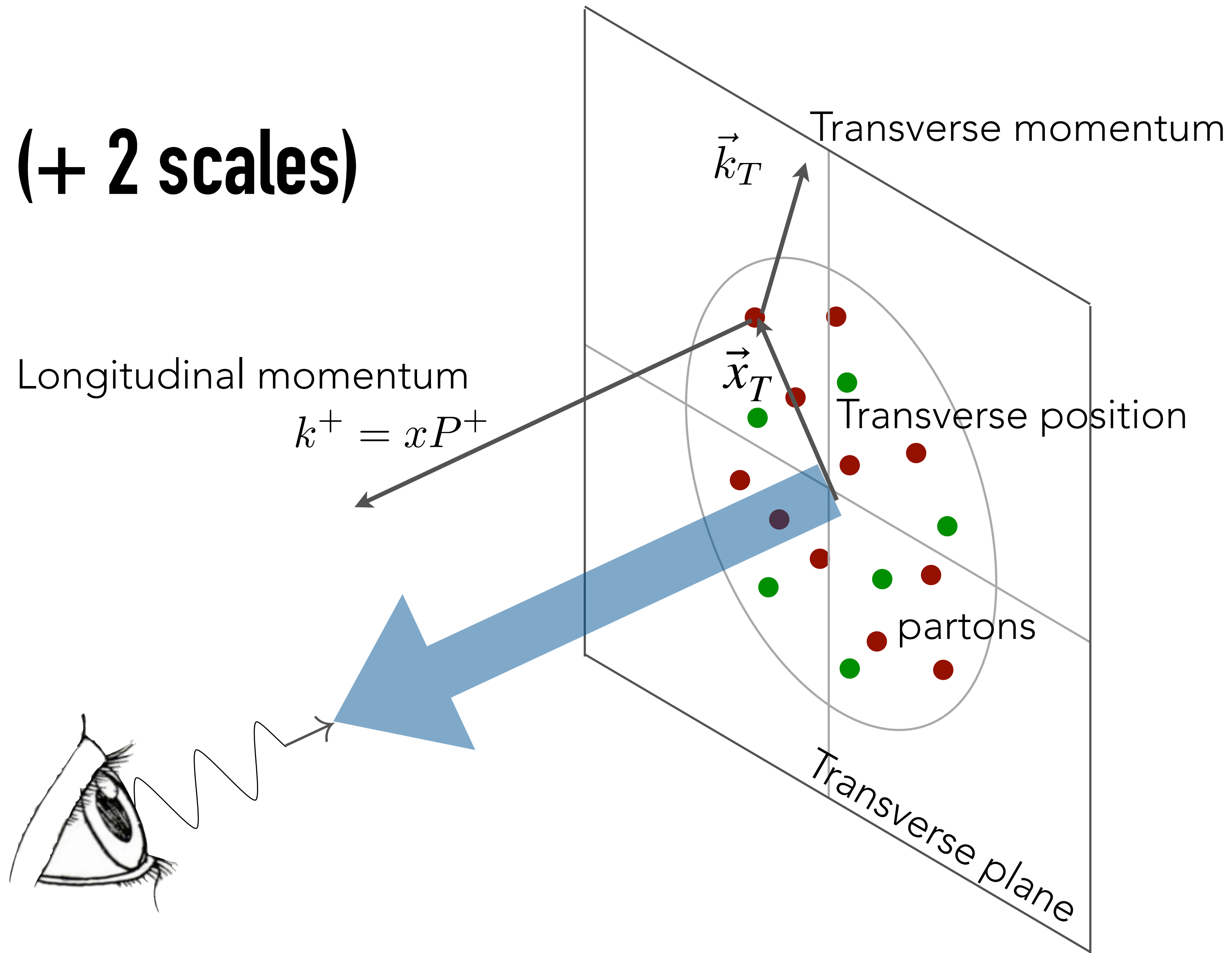




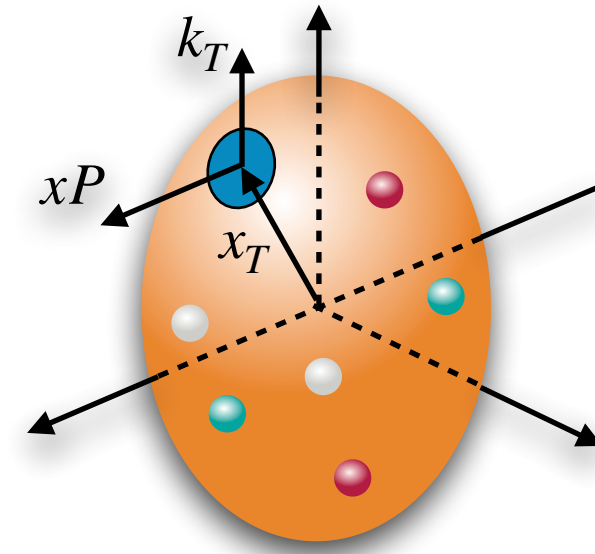
Wigner Distributions

$$f(x, \vec{k}_T, \vec{x}_T)$$

5 dimensional (+ 2 scales)

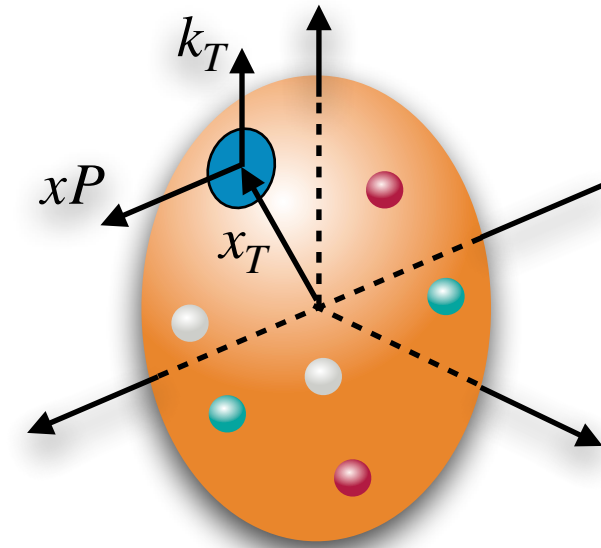
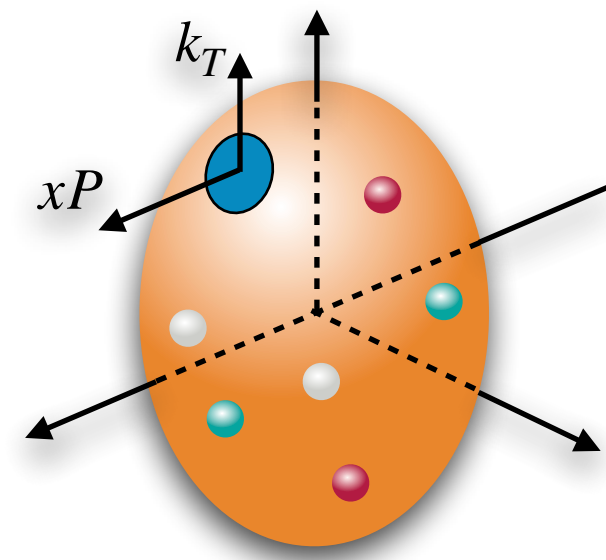


Wigner distributions
(Fourier transform of
GTMDs = Generalized
Transverse Momentum Distributions)



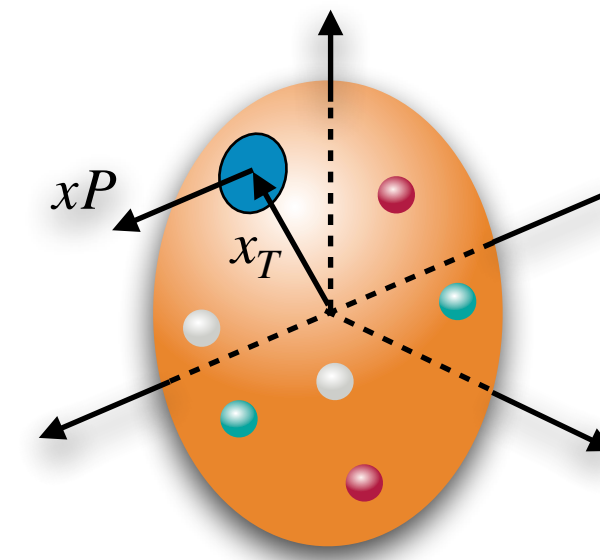
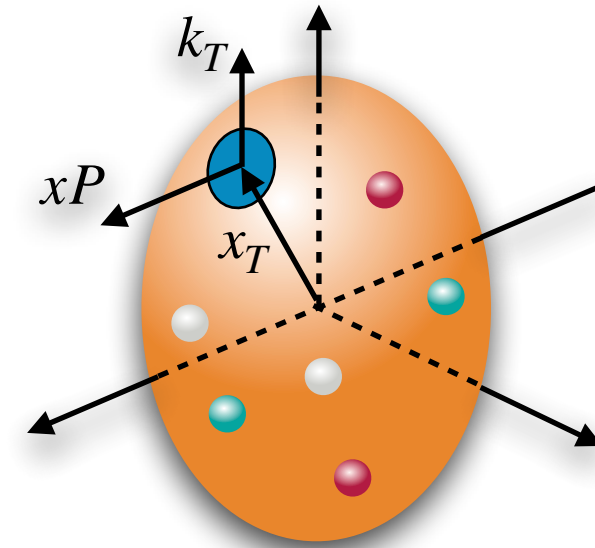
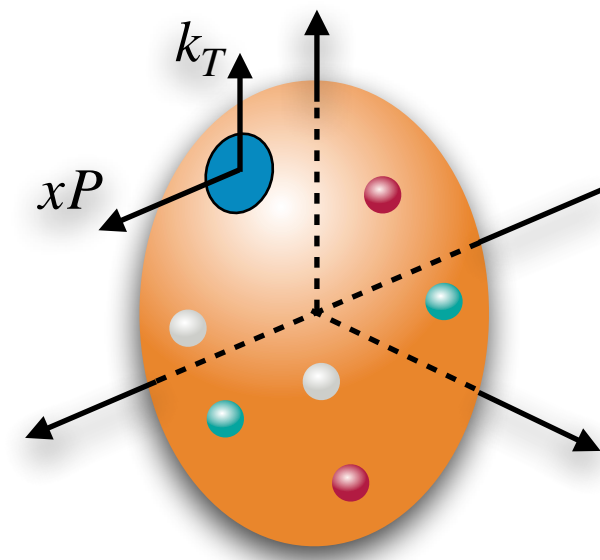
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TMDs



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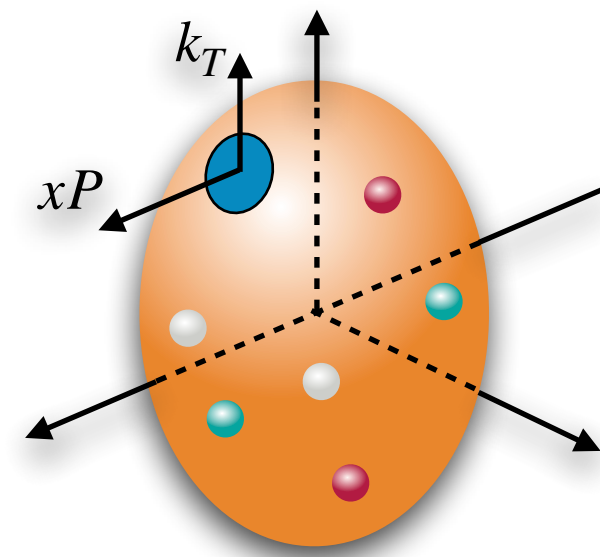
TMDs



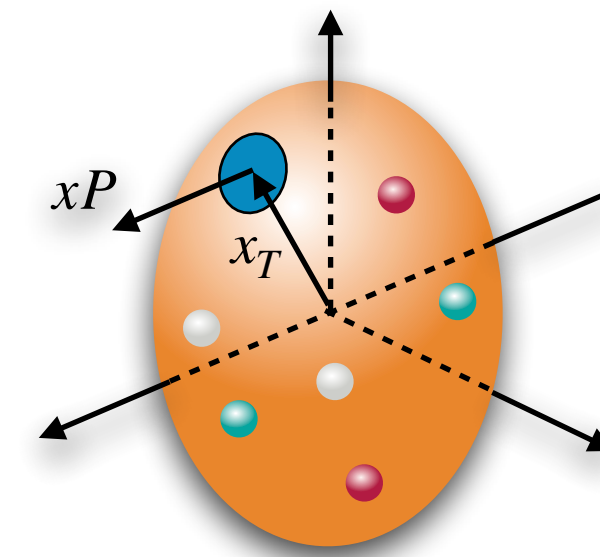
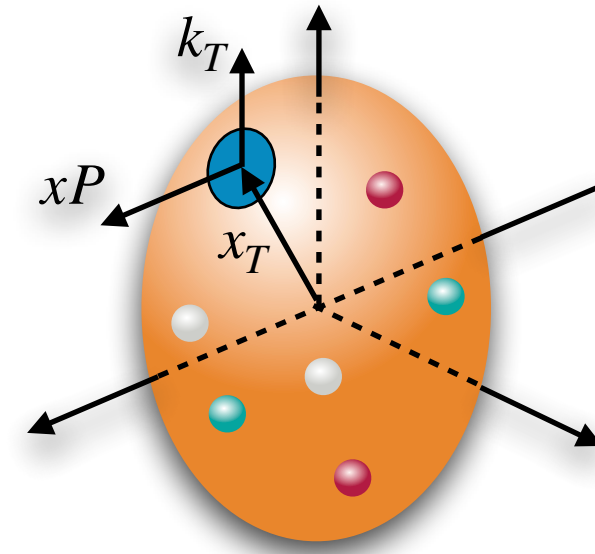
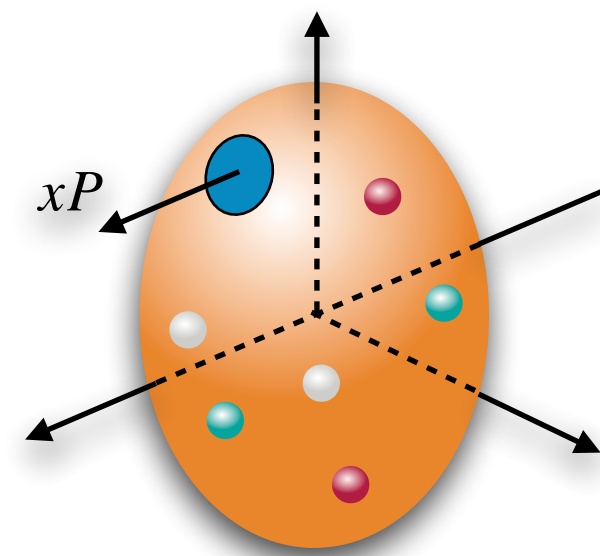
Fourier transform
of Generalized Parton
Distributions

Wigner distributions
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TMDs



PDFs

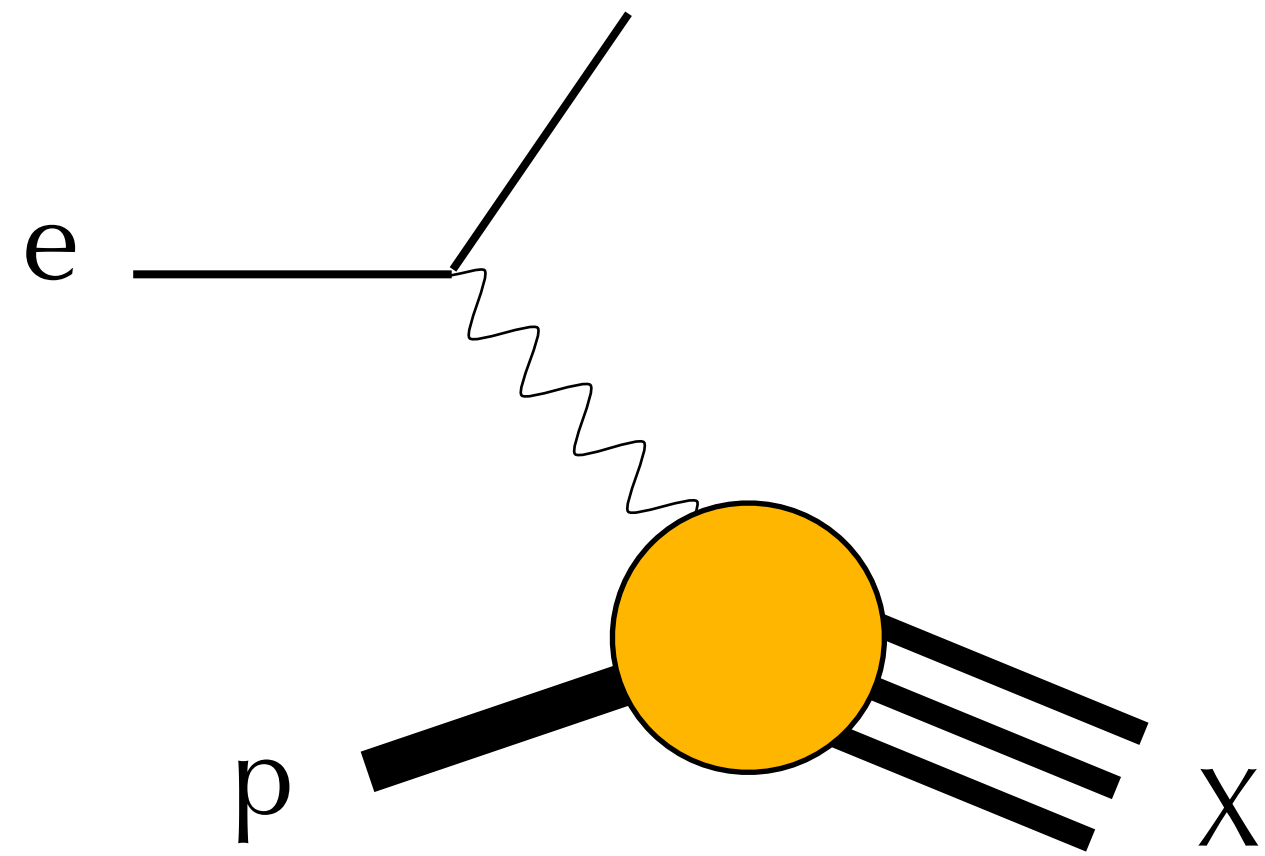


Fourier transform
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Distributions

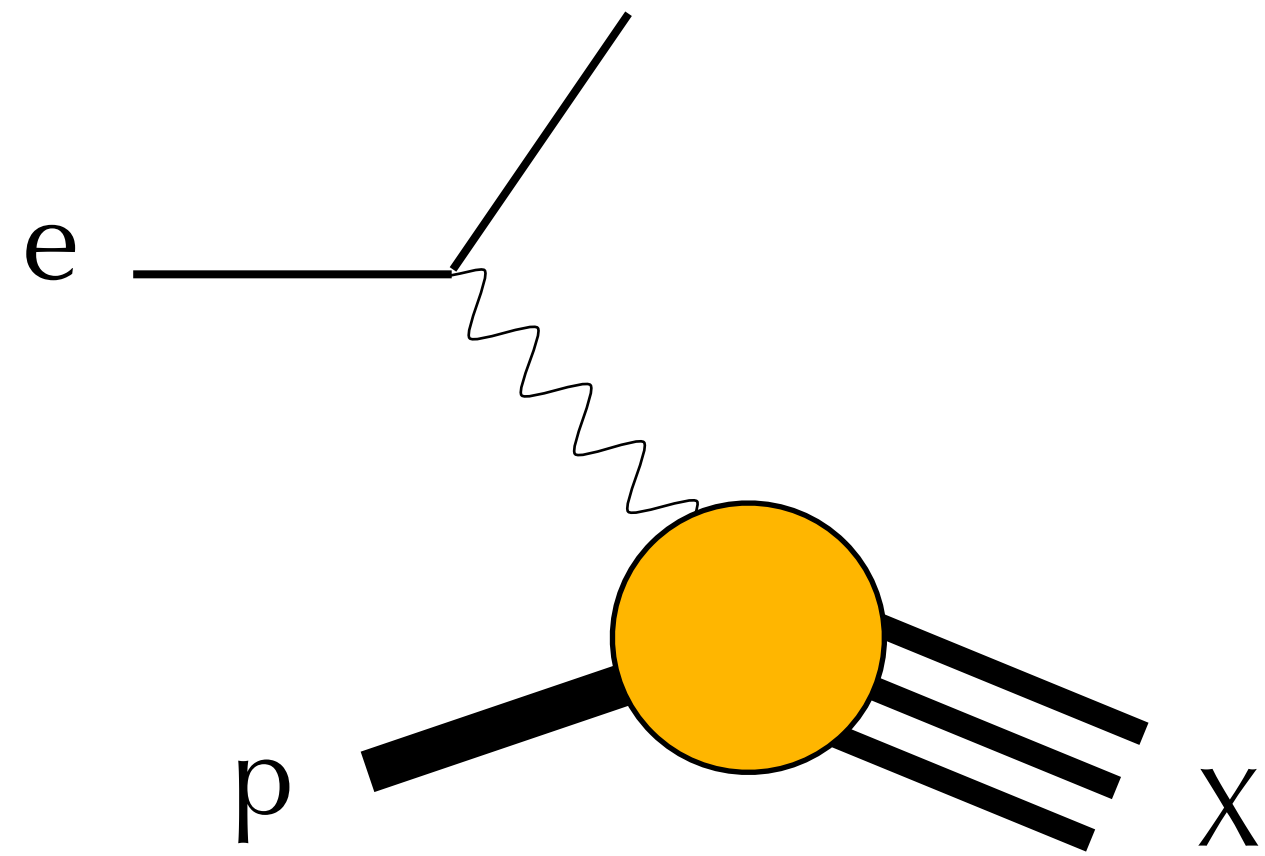
Fourier transform
of Form Factors

HOW TO RECONSTRUCT THESE MAPS?

Inclusive DIS

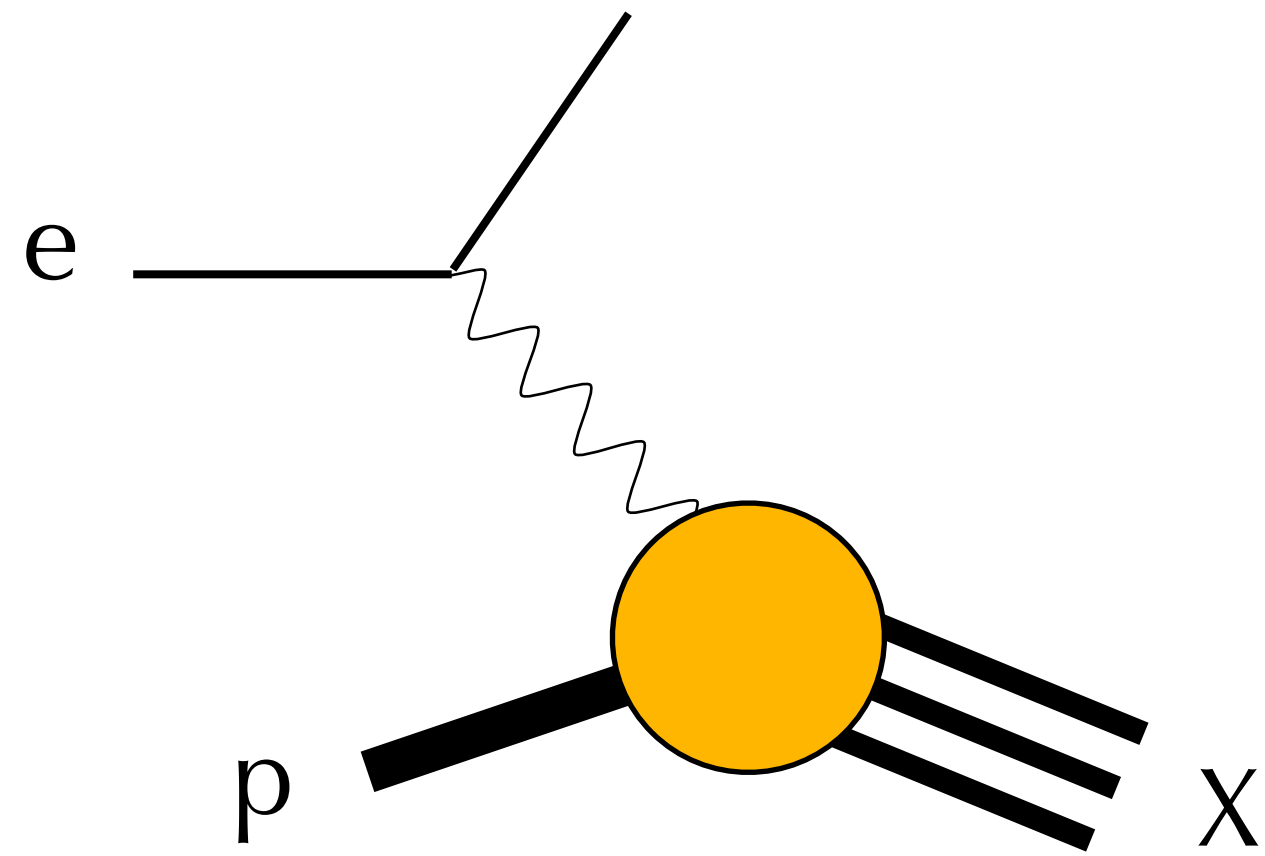


Inclusive DIS



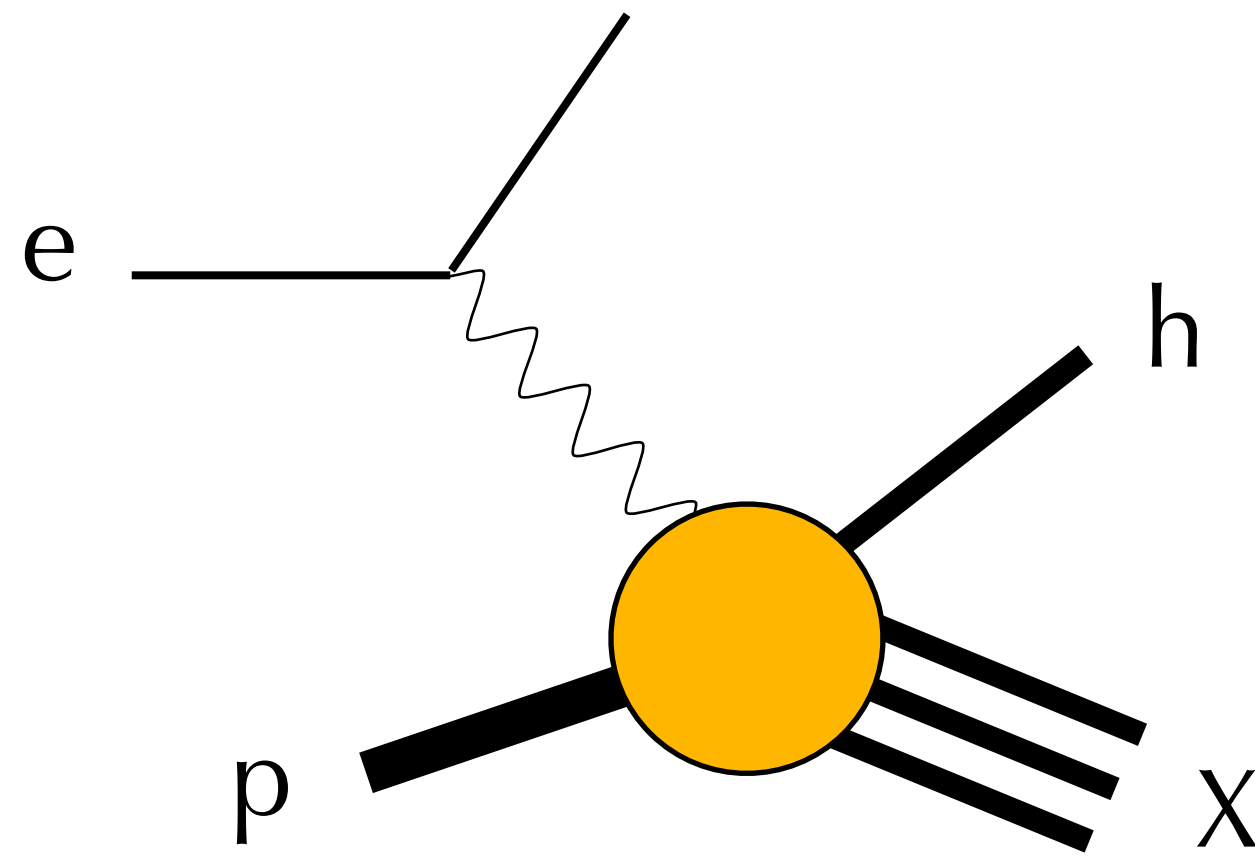
access to
Parton Distribution Functions

Inclusive DIS

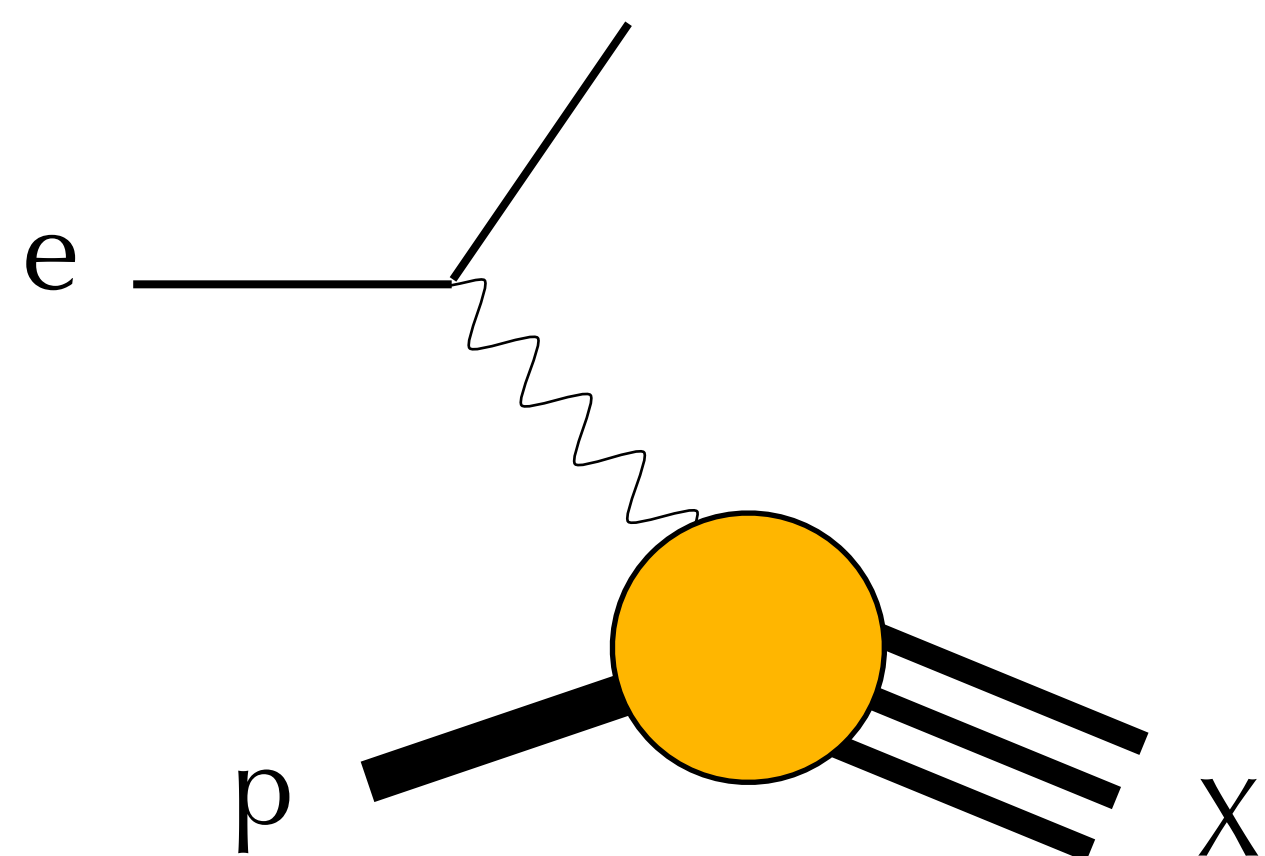


access to
Parton Distribution Functions

Semi-Inclusive DIS

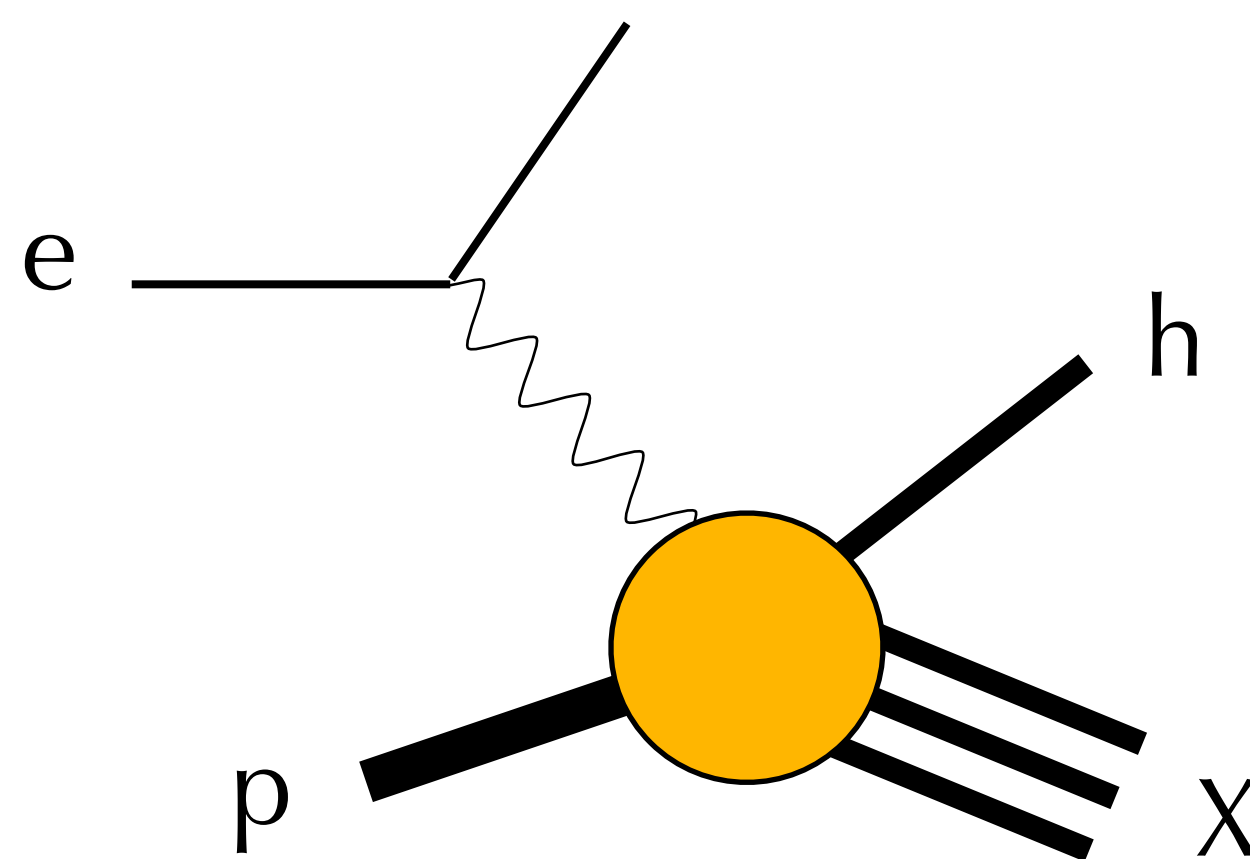


Inclusive DIS



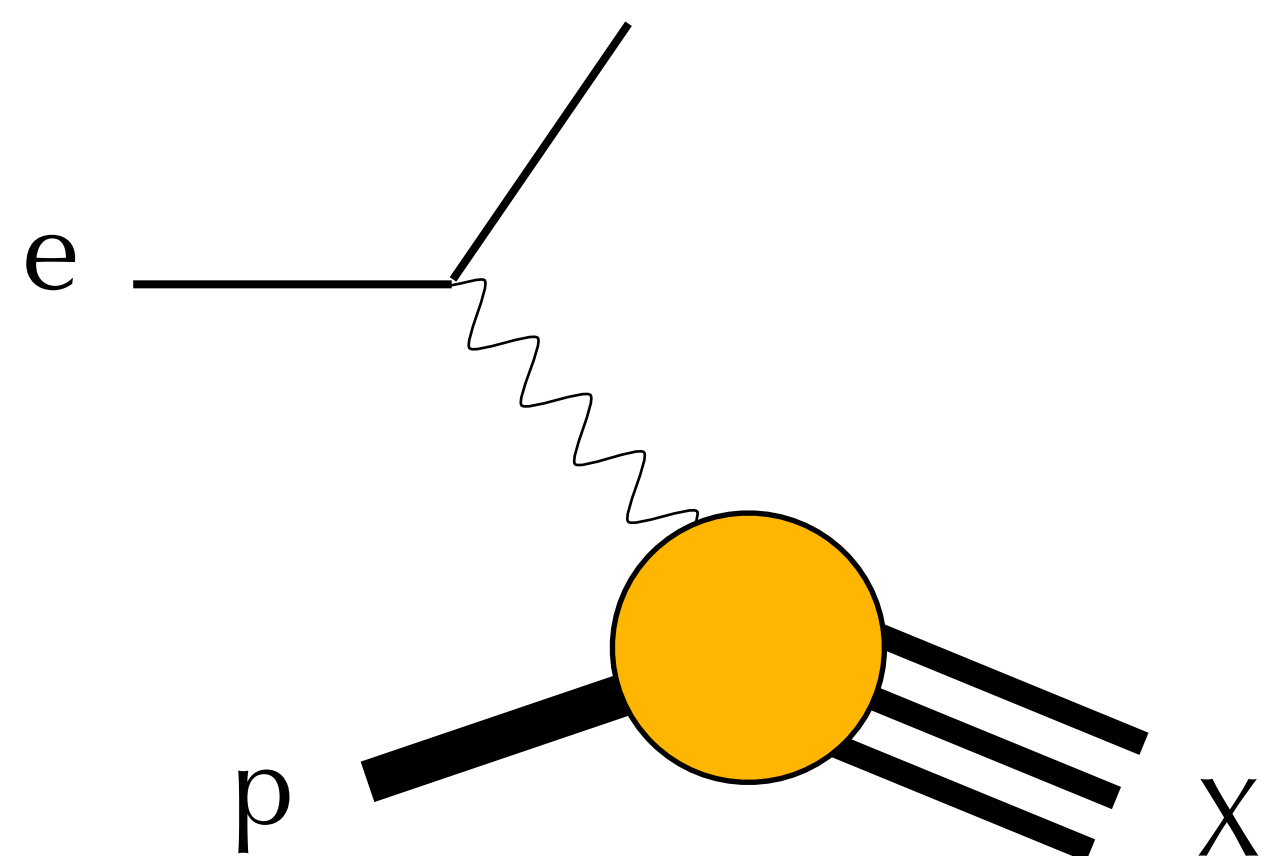
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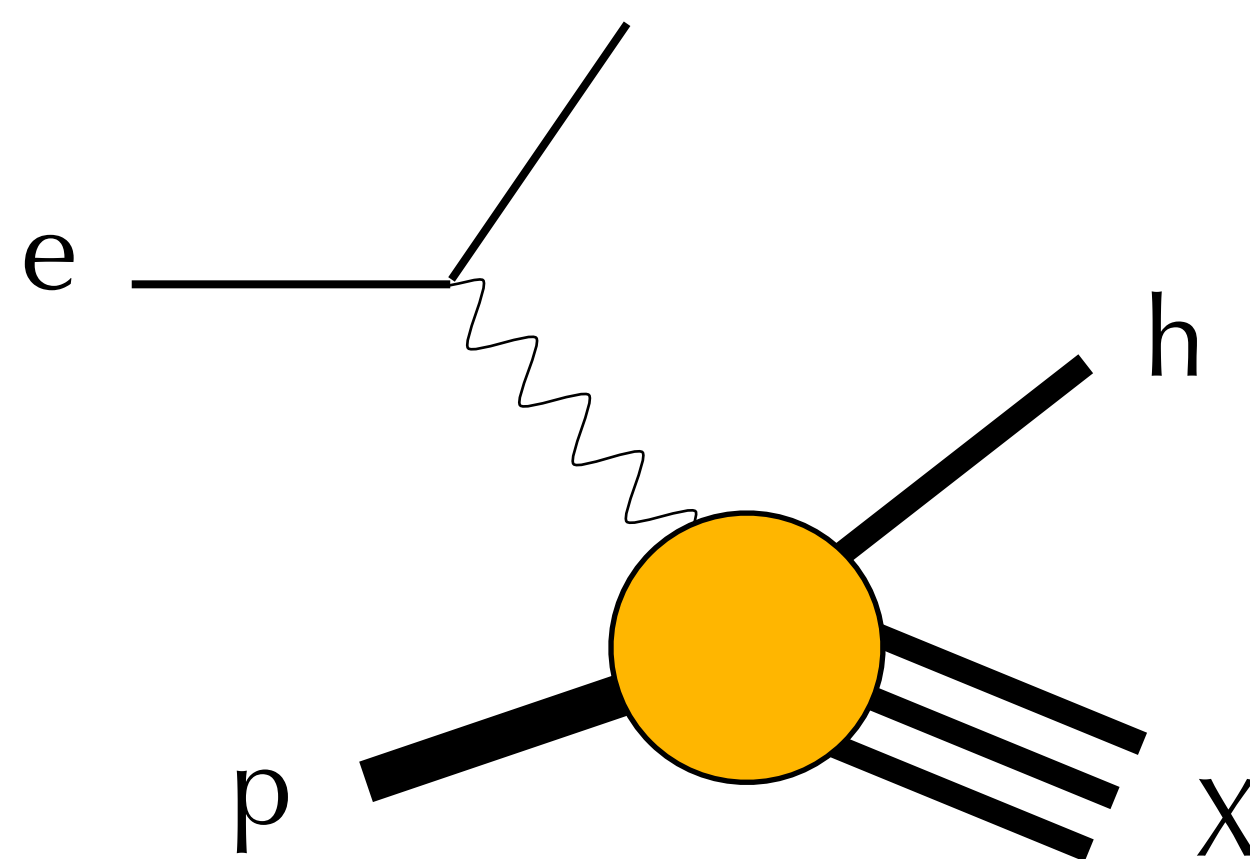
access to
Transverse Momentum Distributions

Inclusive DIS



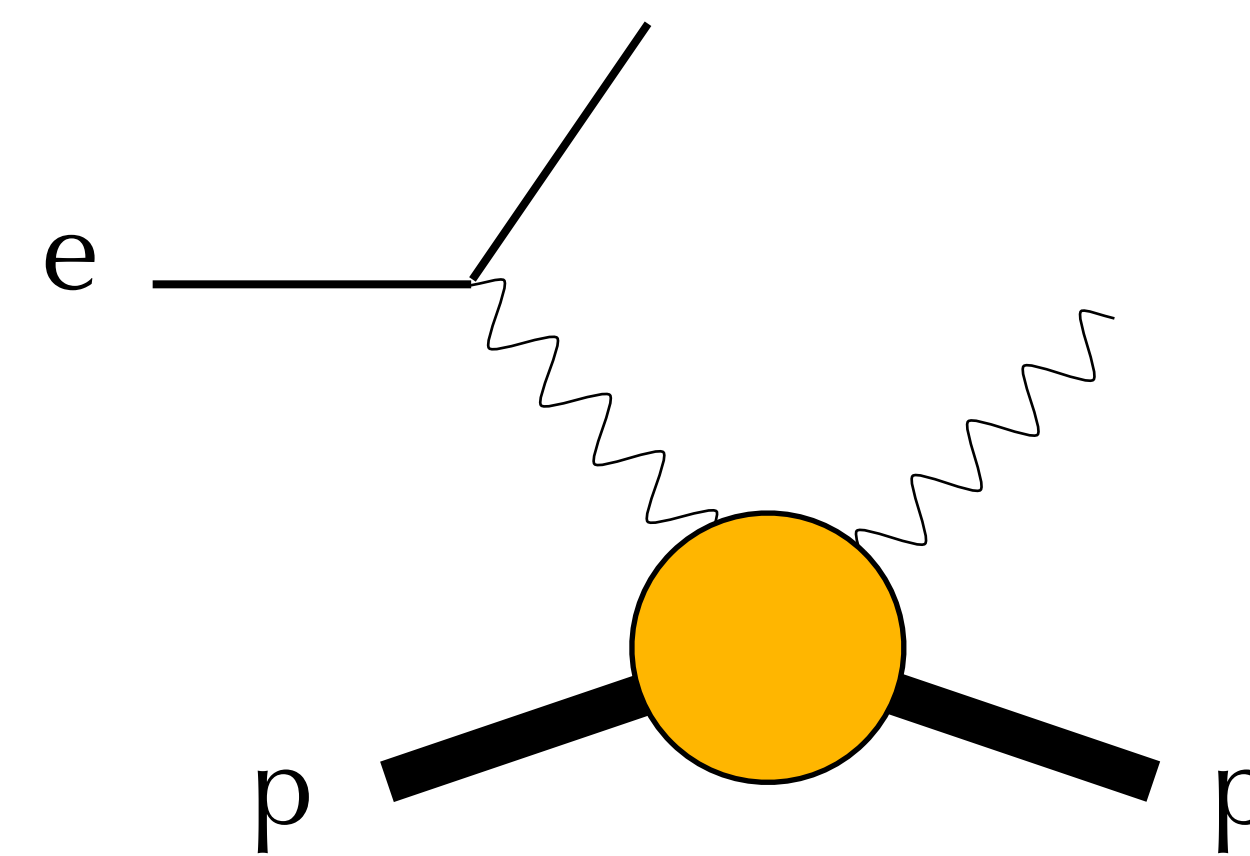
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Parton Distribution Functions

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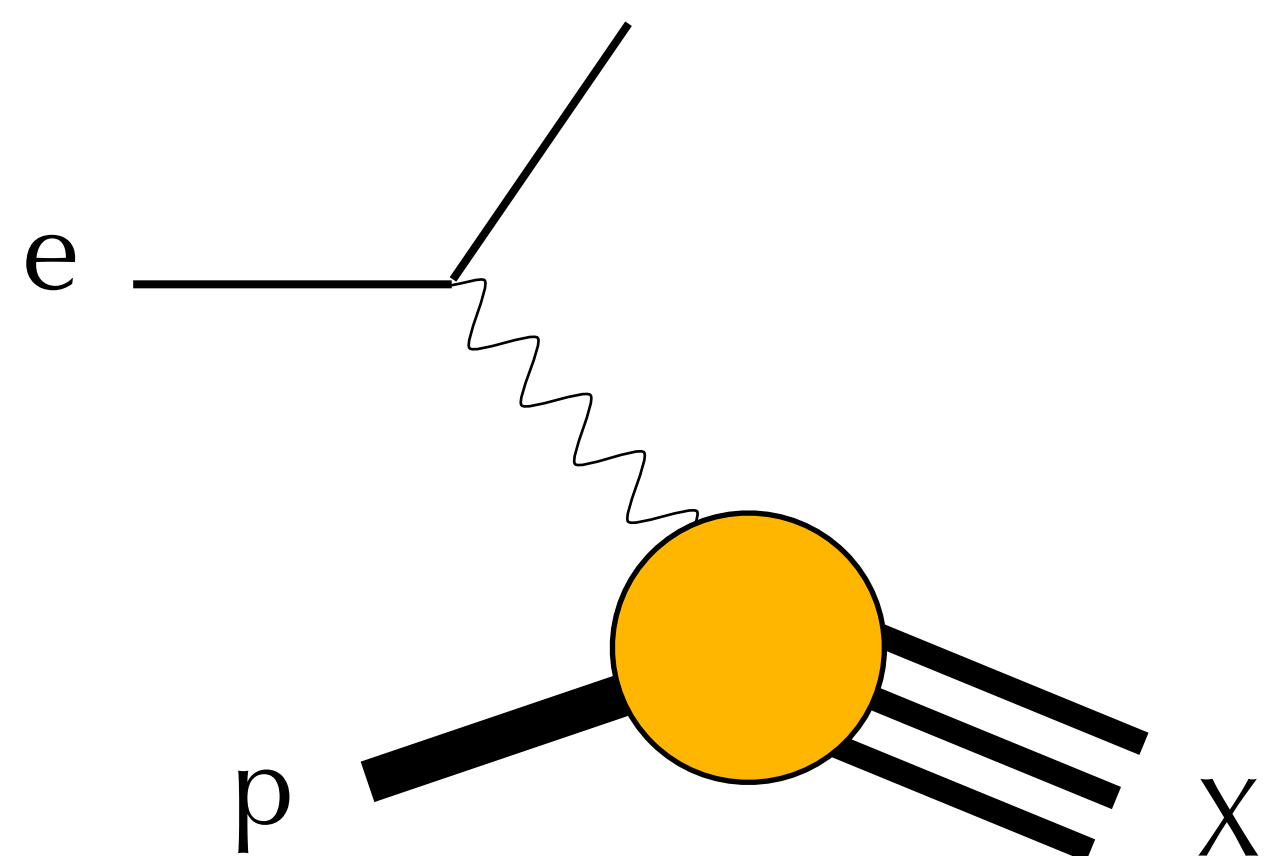


access to
Transverse Momentum Distributions

Exclusive processes

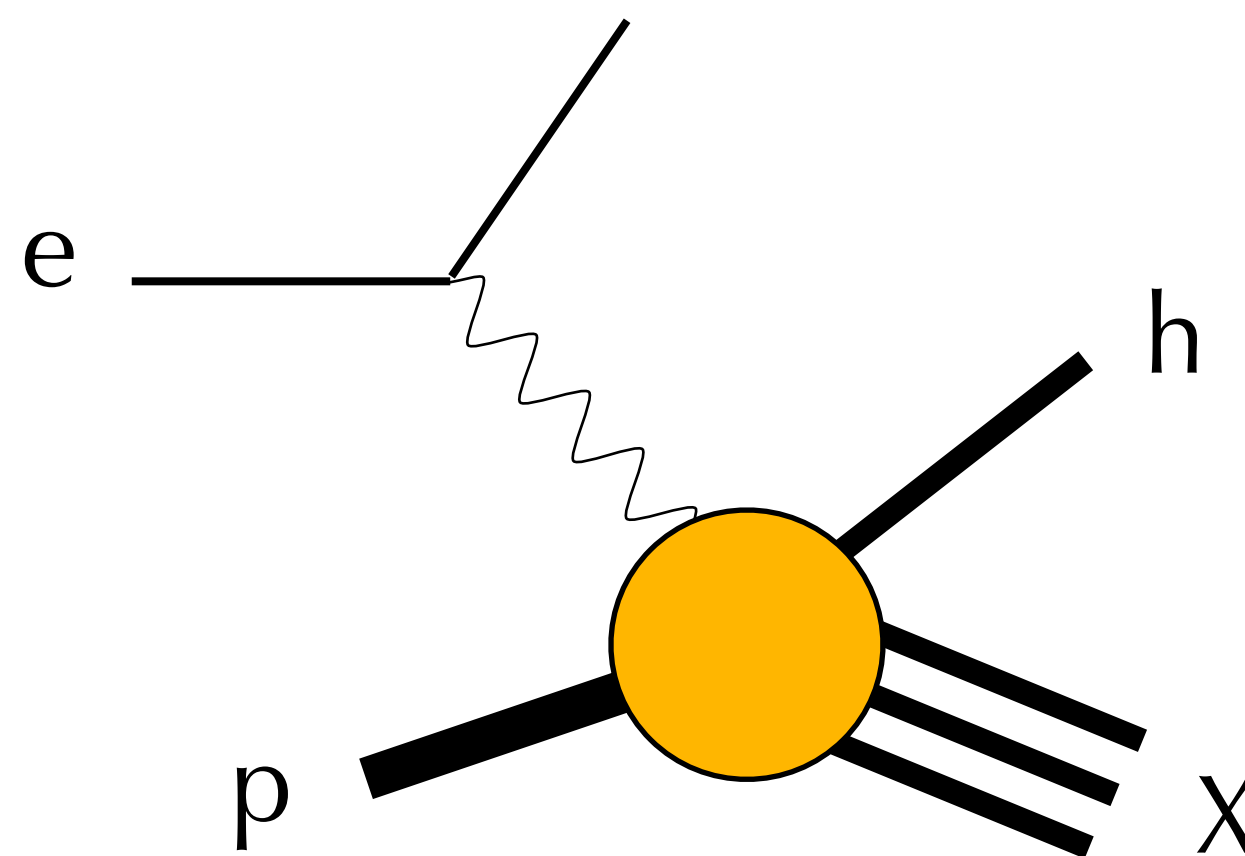


Inclusive DIS



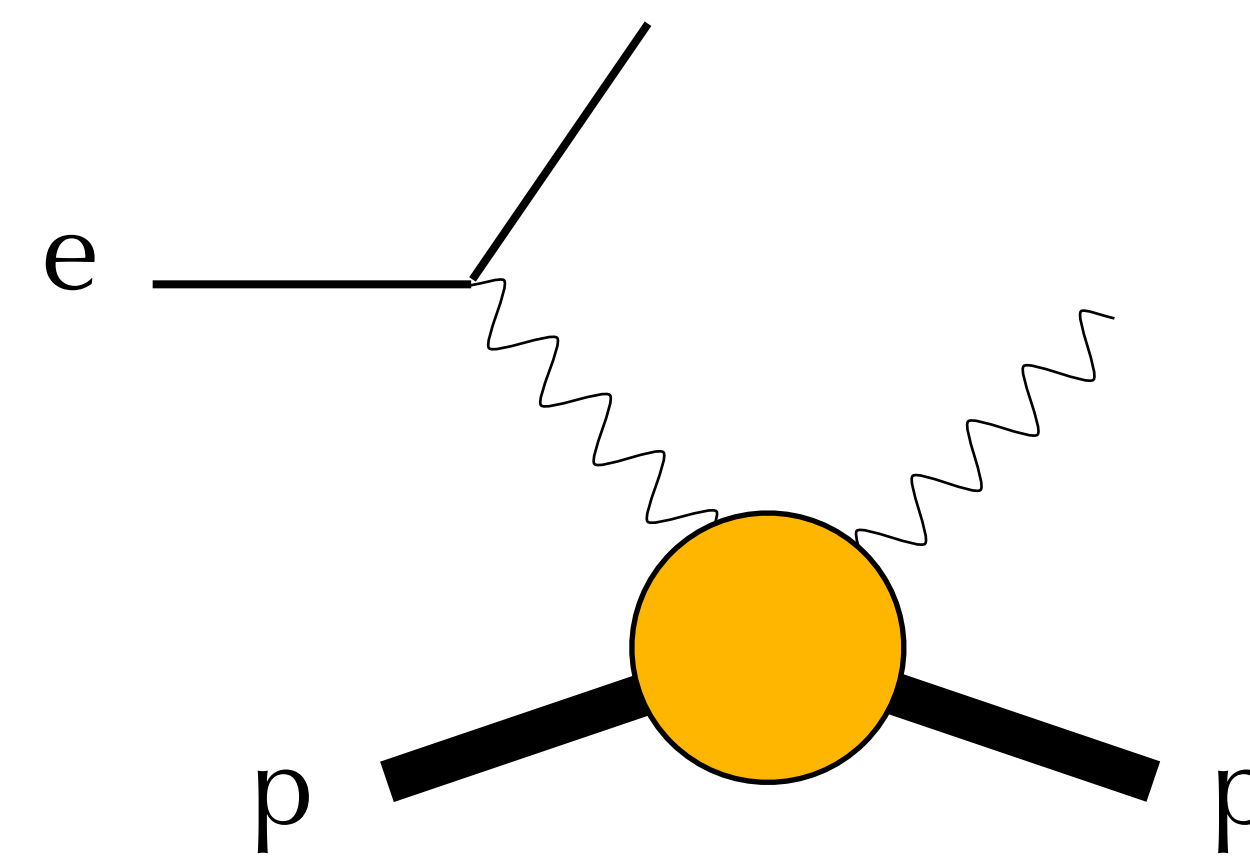
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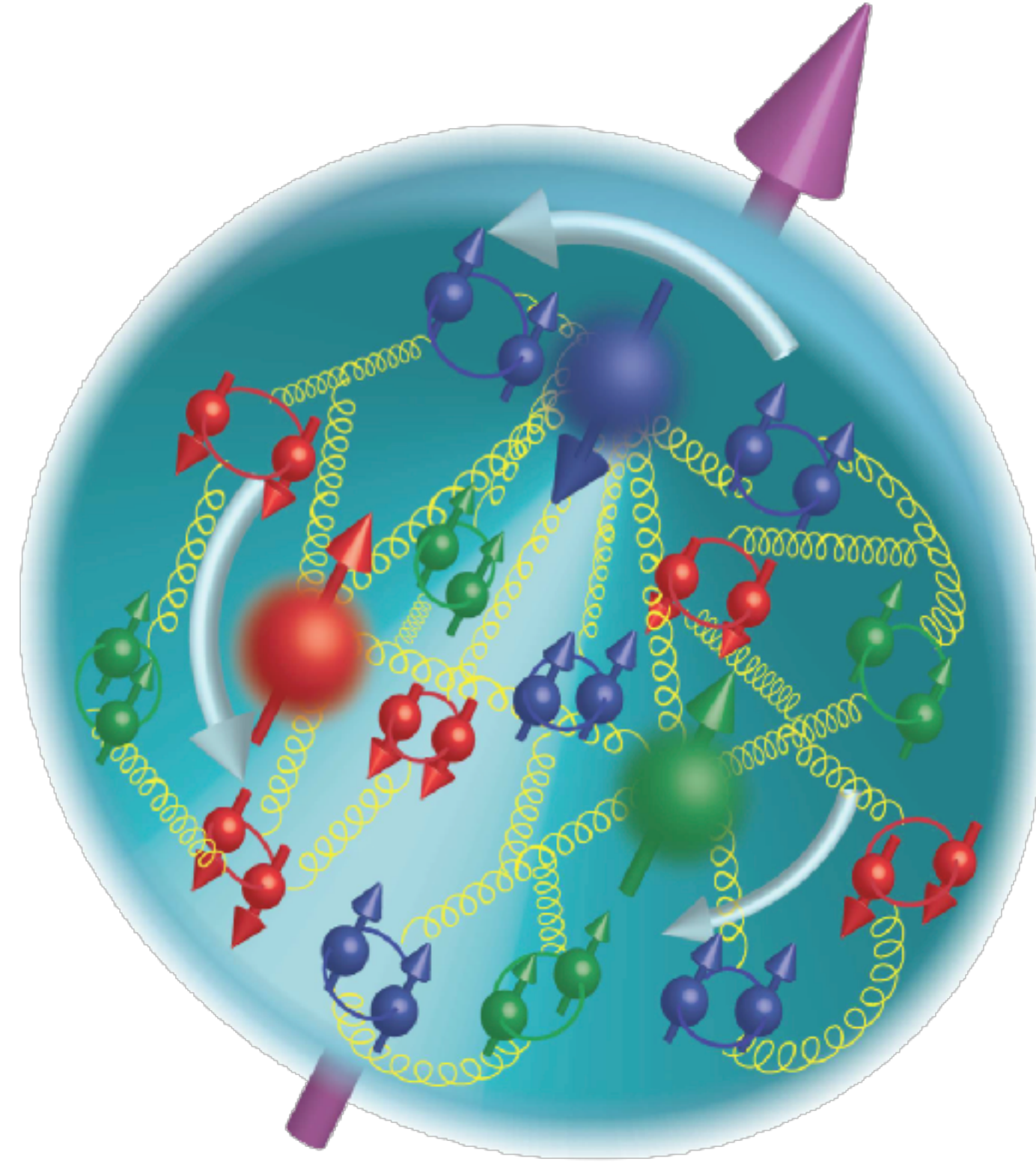
access to
Transverse Momentum Distributions

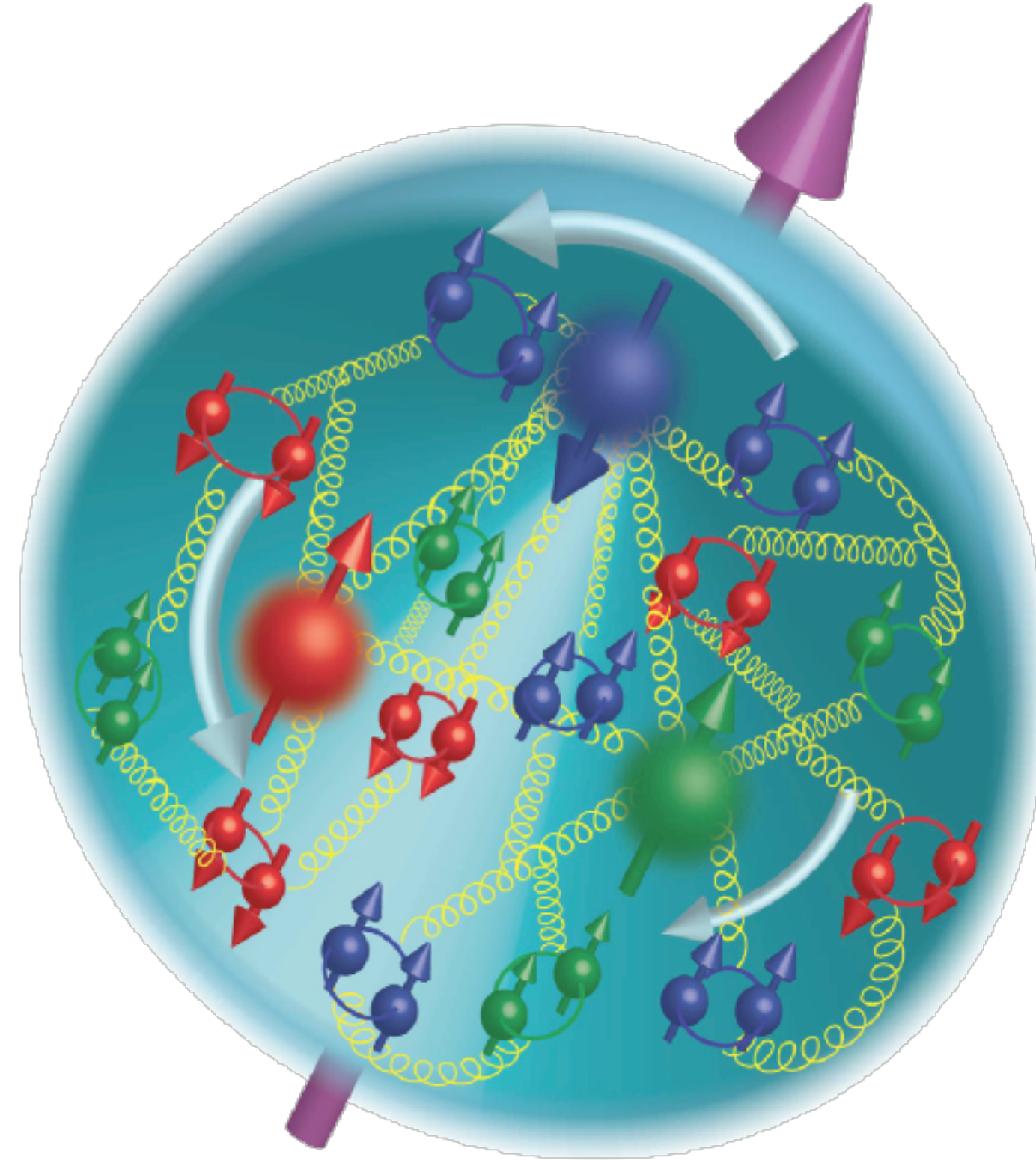
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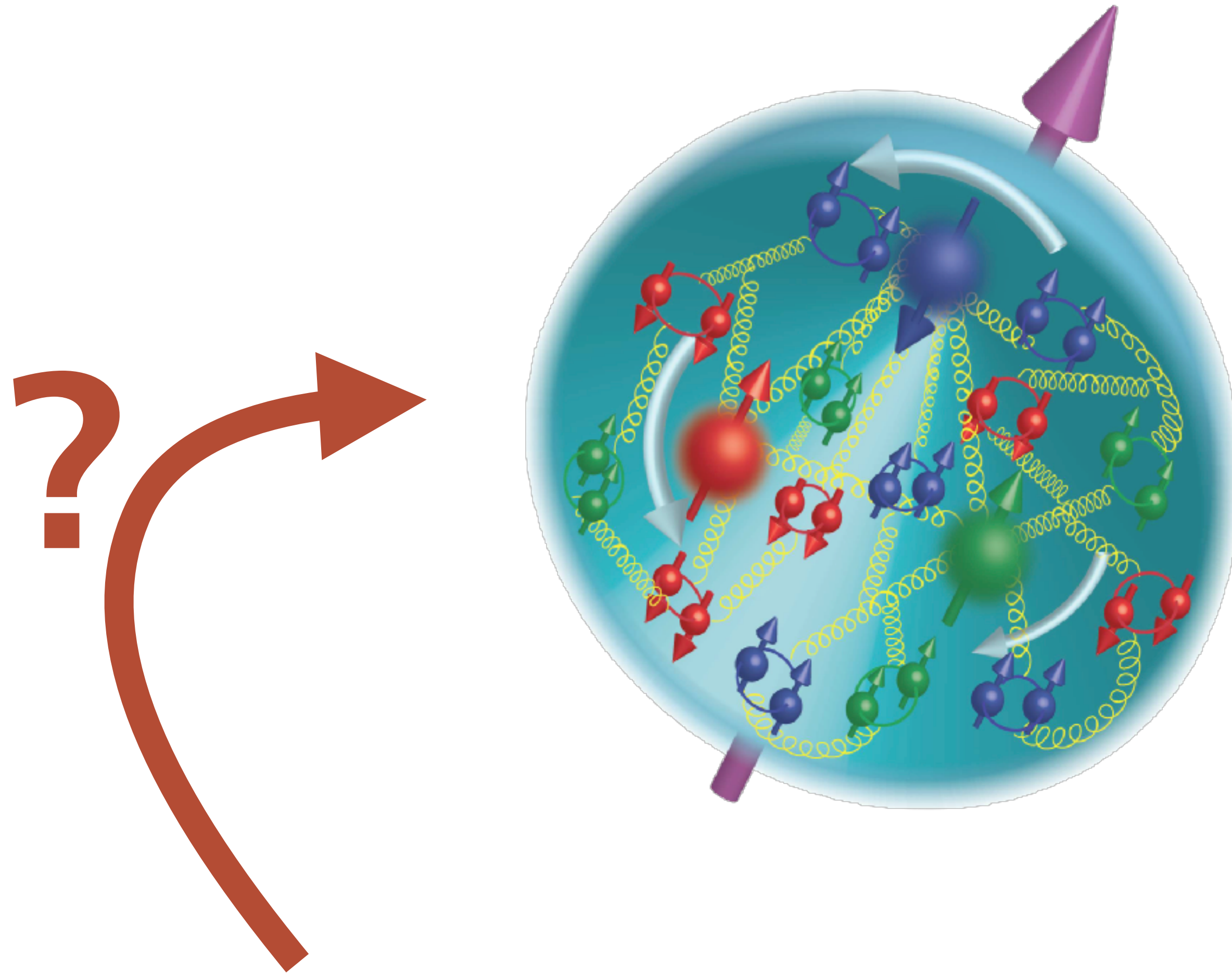
access to
Generalized Parton Distributions

WHY IS IT INTERESTING TO MAP THE NUCLEON?





$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{\partial} - g \not{A} + m) \psi_q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$



$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{\partial} - g \not{A} + m) \psi_q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

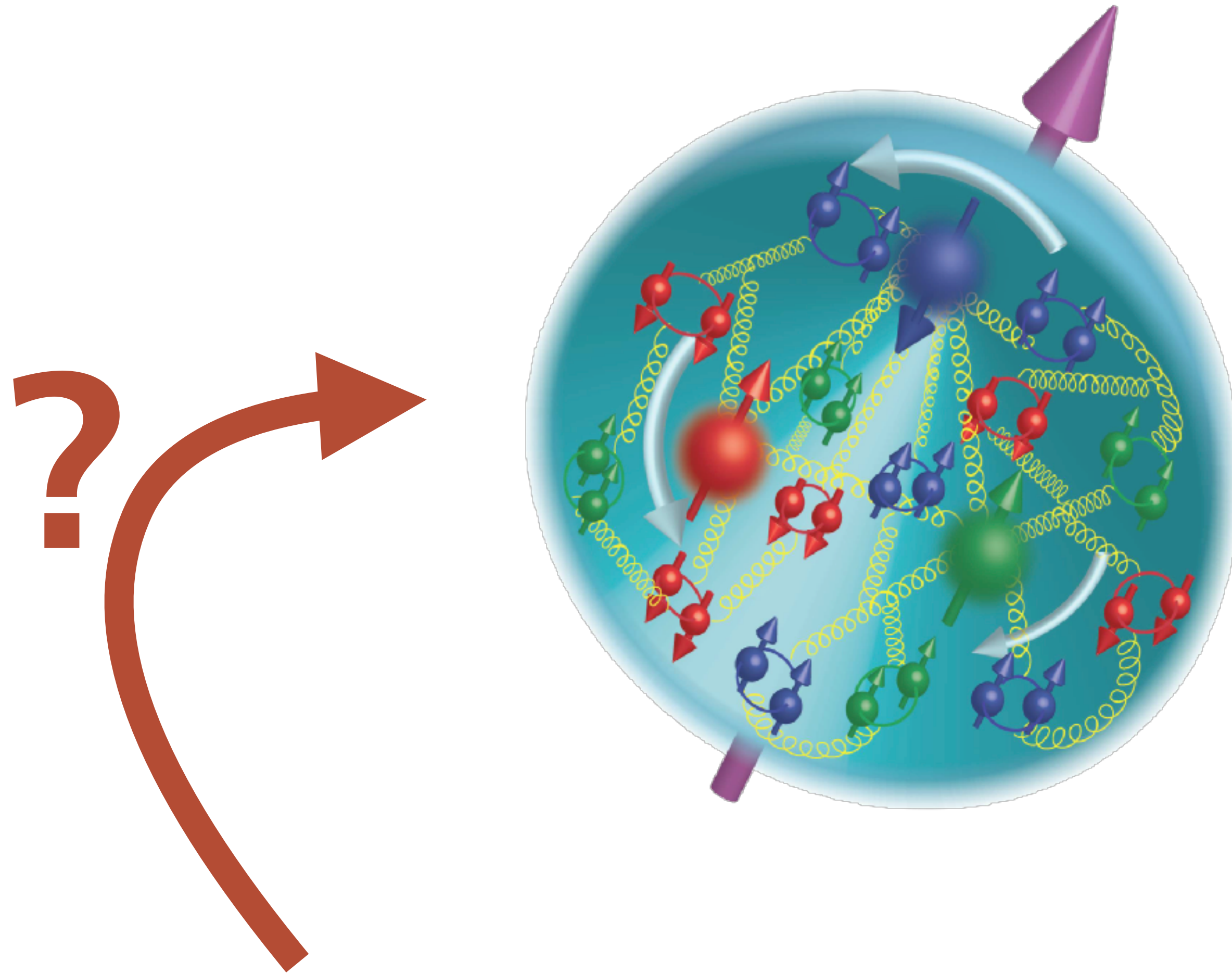


QCD: the **WILD SIDE** of the Standard Model



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there are more things that we cannot explain than we can explain...

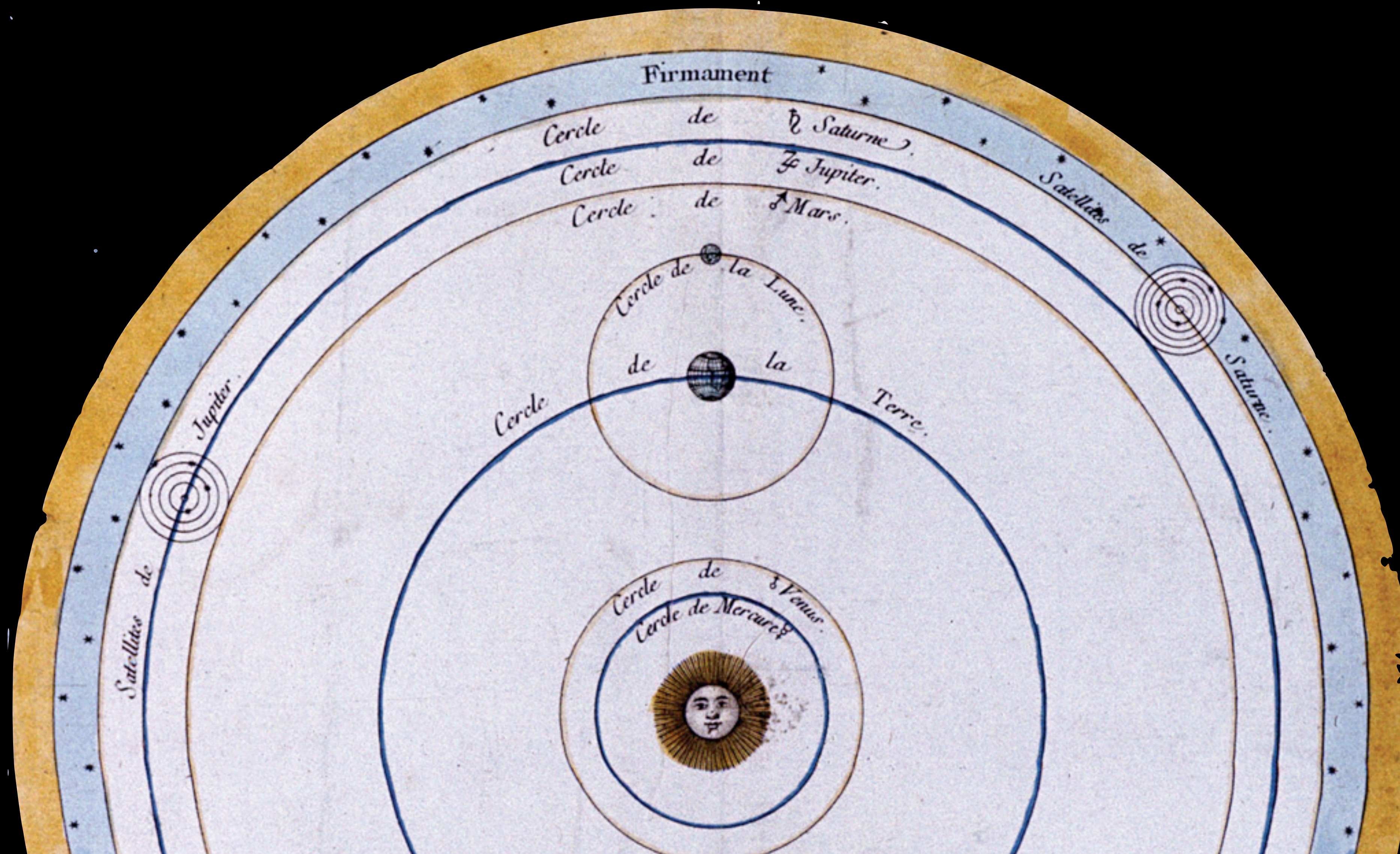


$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{\partial} - g \not{A} + m) \psi_q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

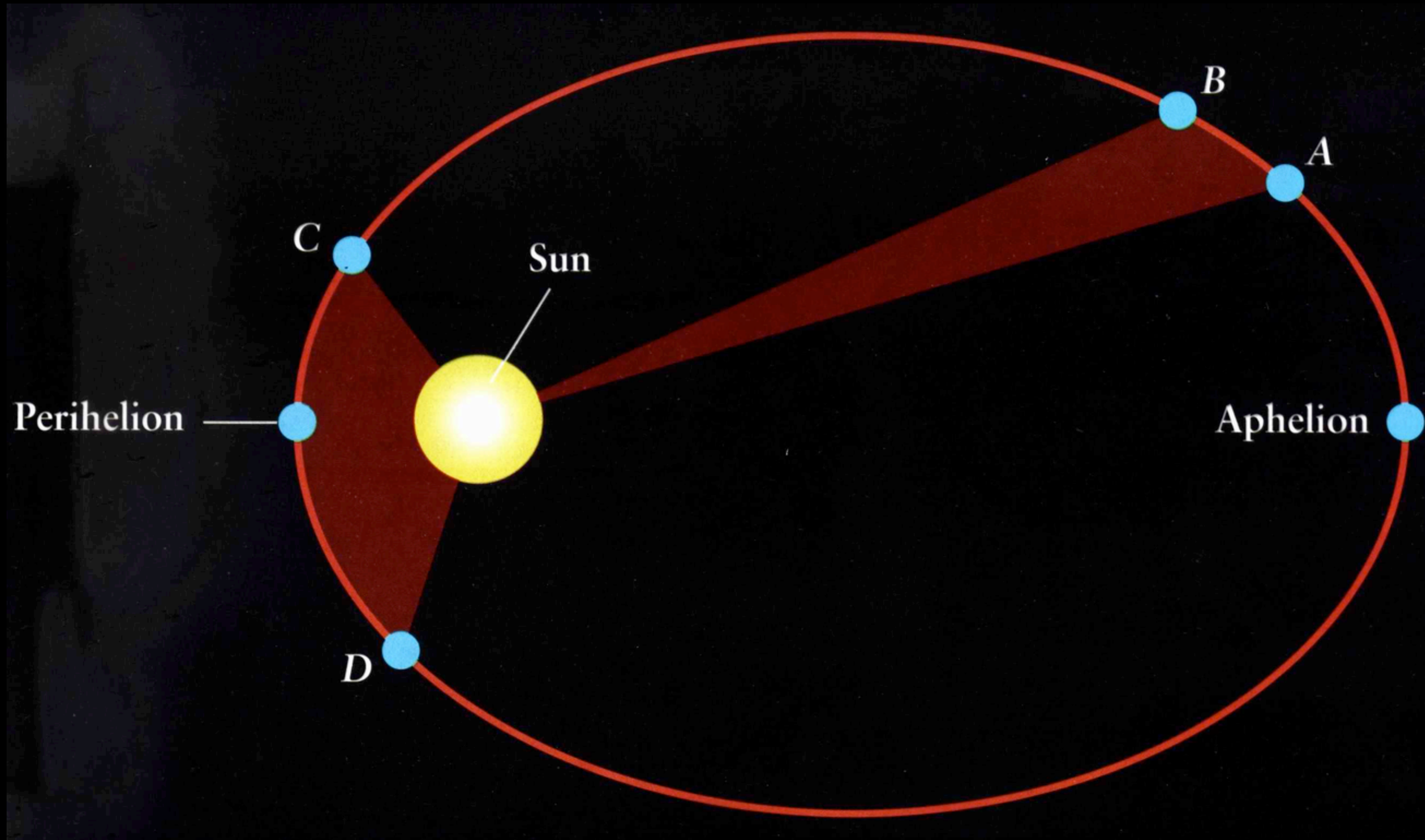
WHY IS IT INTERESTING TO MAP THE SKY?



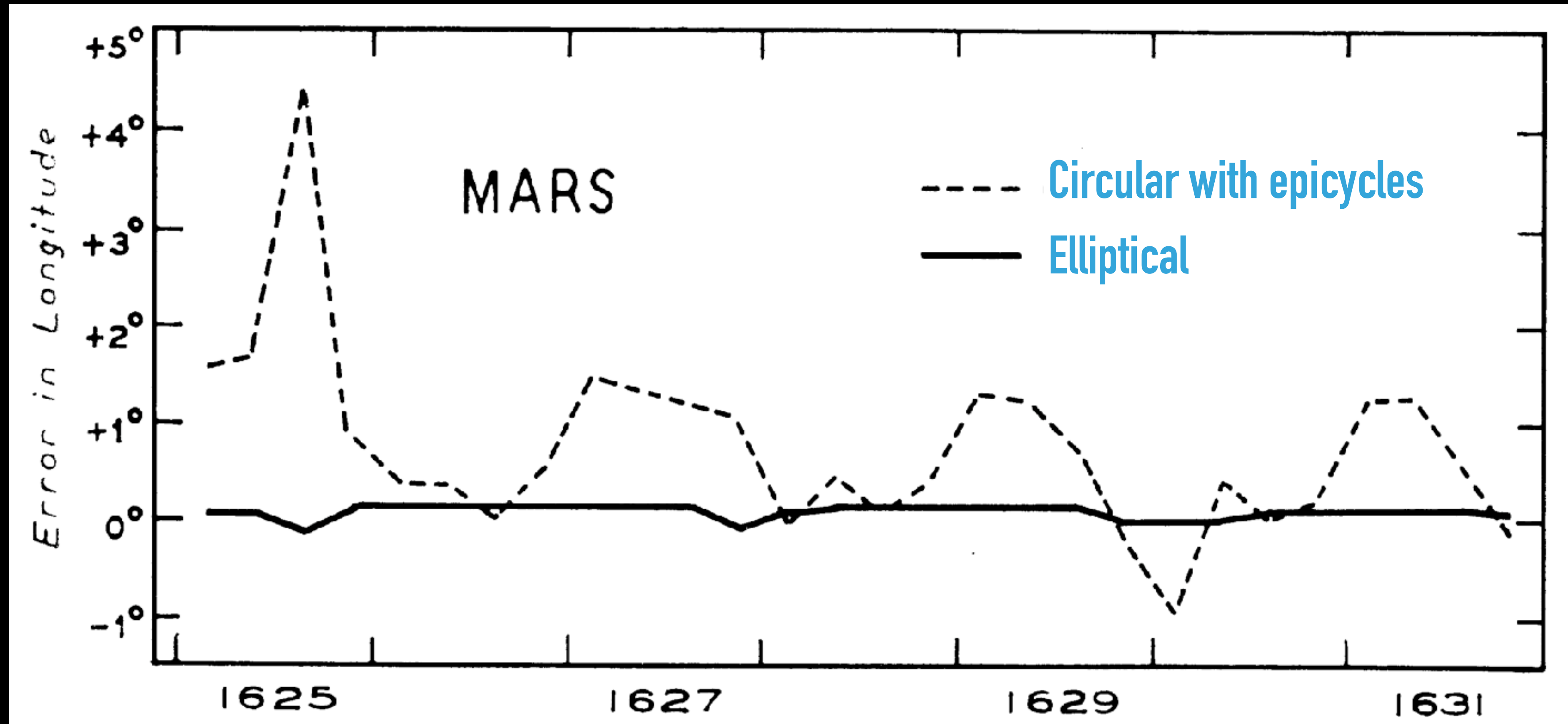
COPERNICAN MODEL



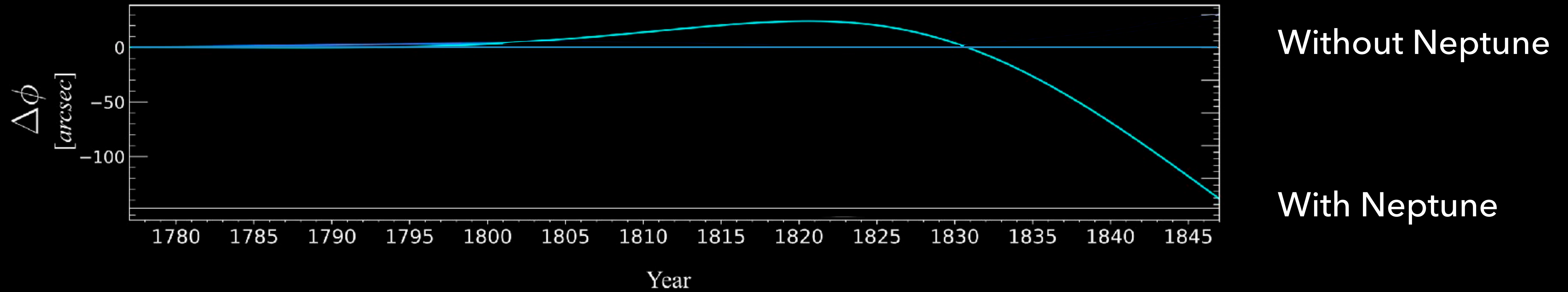
KEPLER'S MODEL



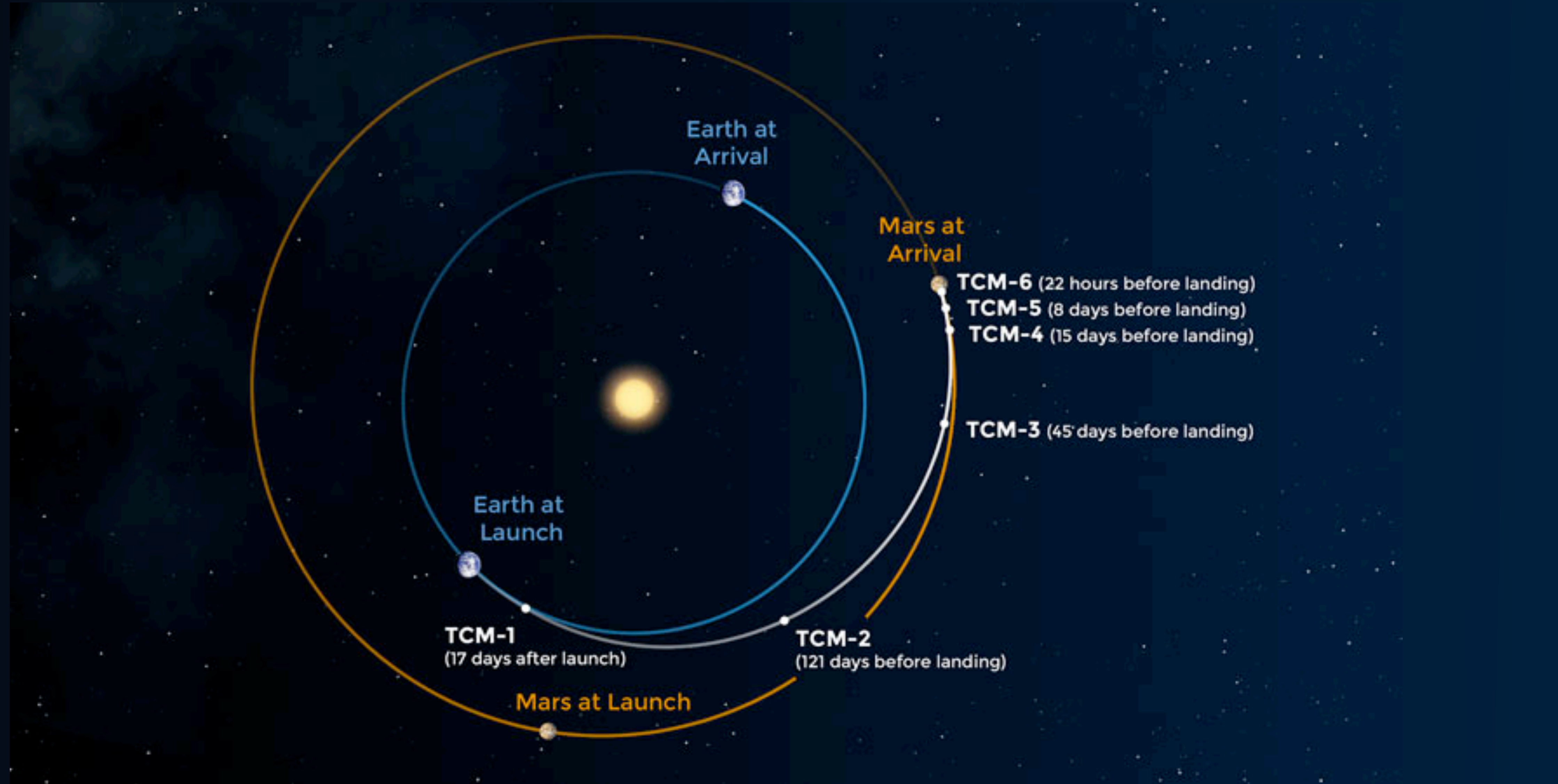
CHECK PREDICTIONS

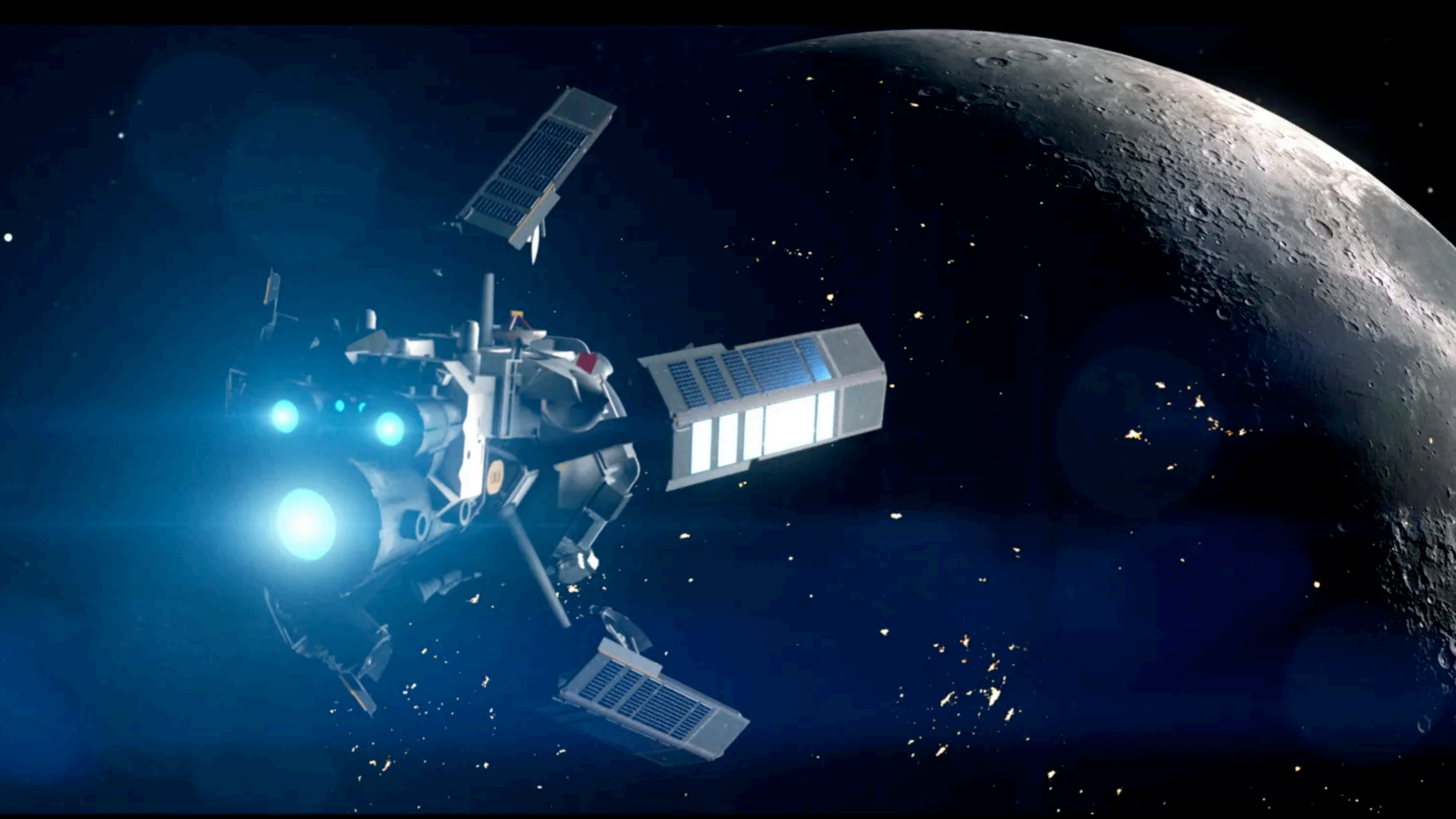


Uranus's longitude predictions

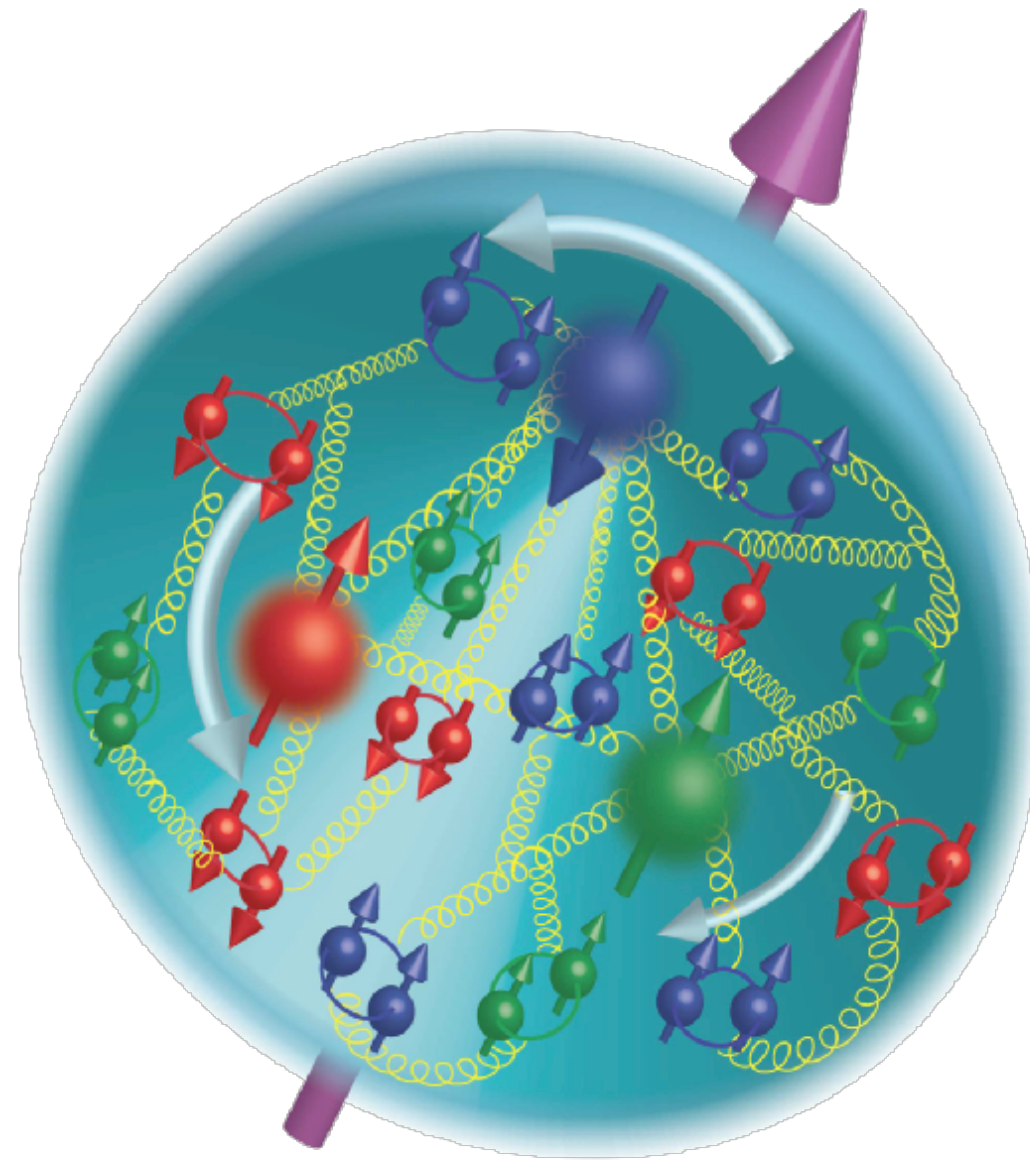


MAKE PREDICTIONS

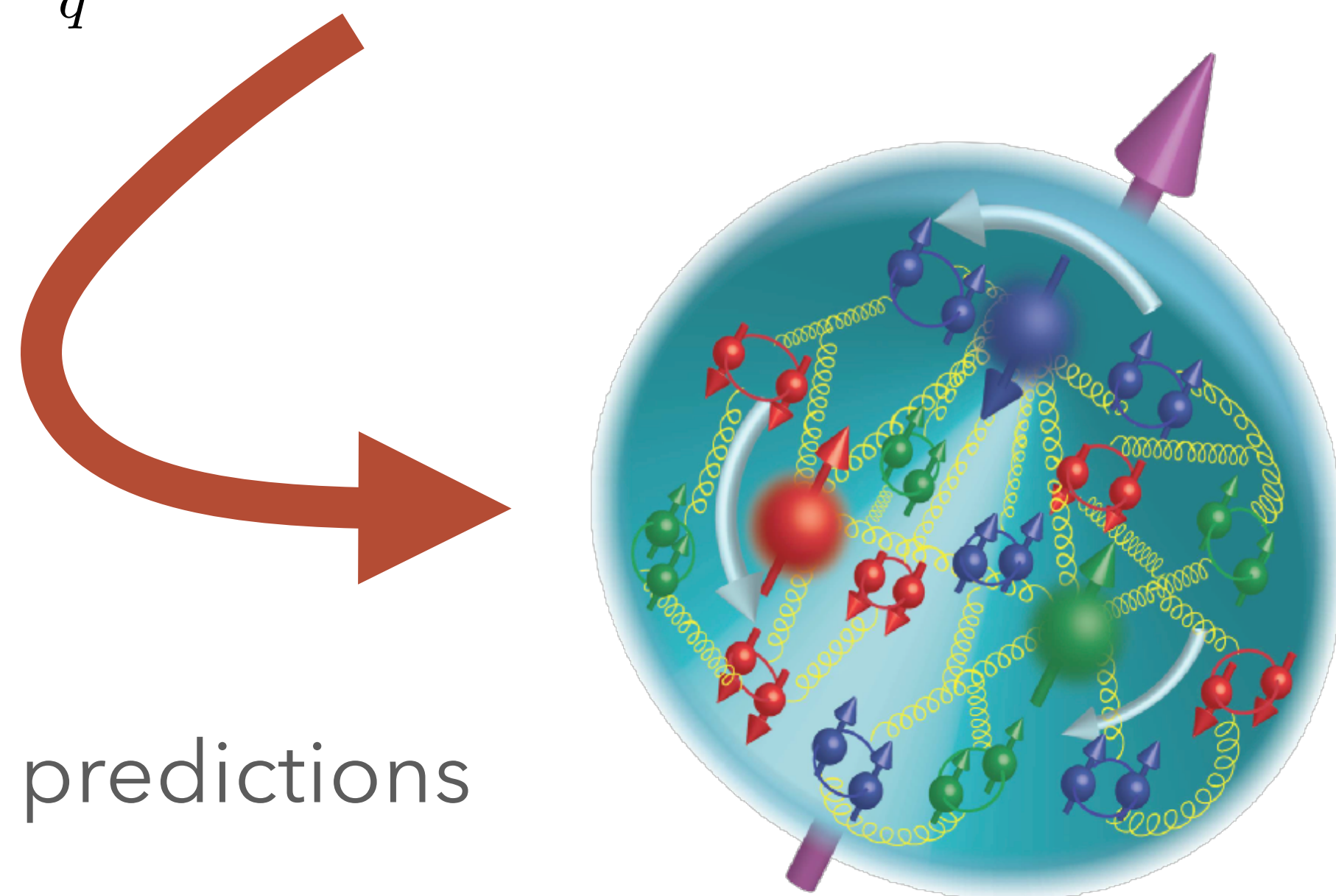




WHY IS IT INTERESTING TO MAP THE NUCLEON?



$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{\partial} - g \not{A} + m) \psi_q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

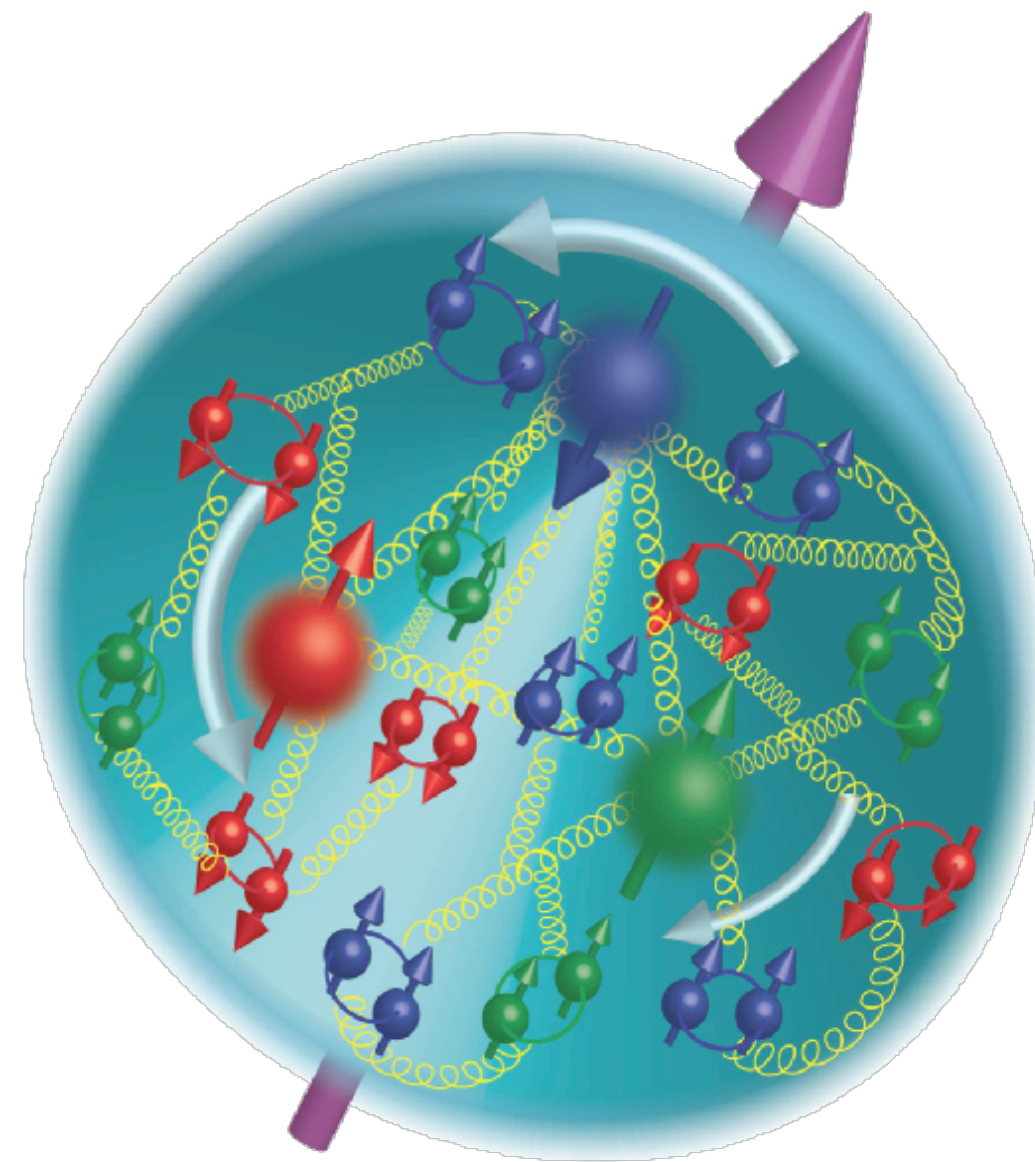


Check predictions

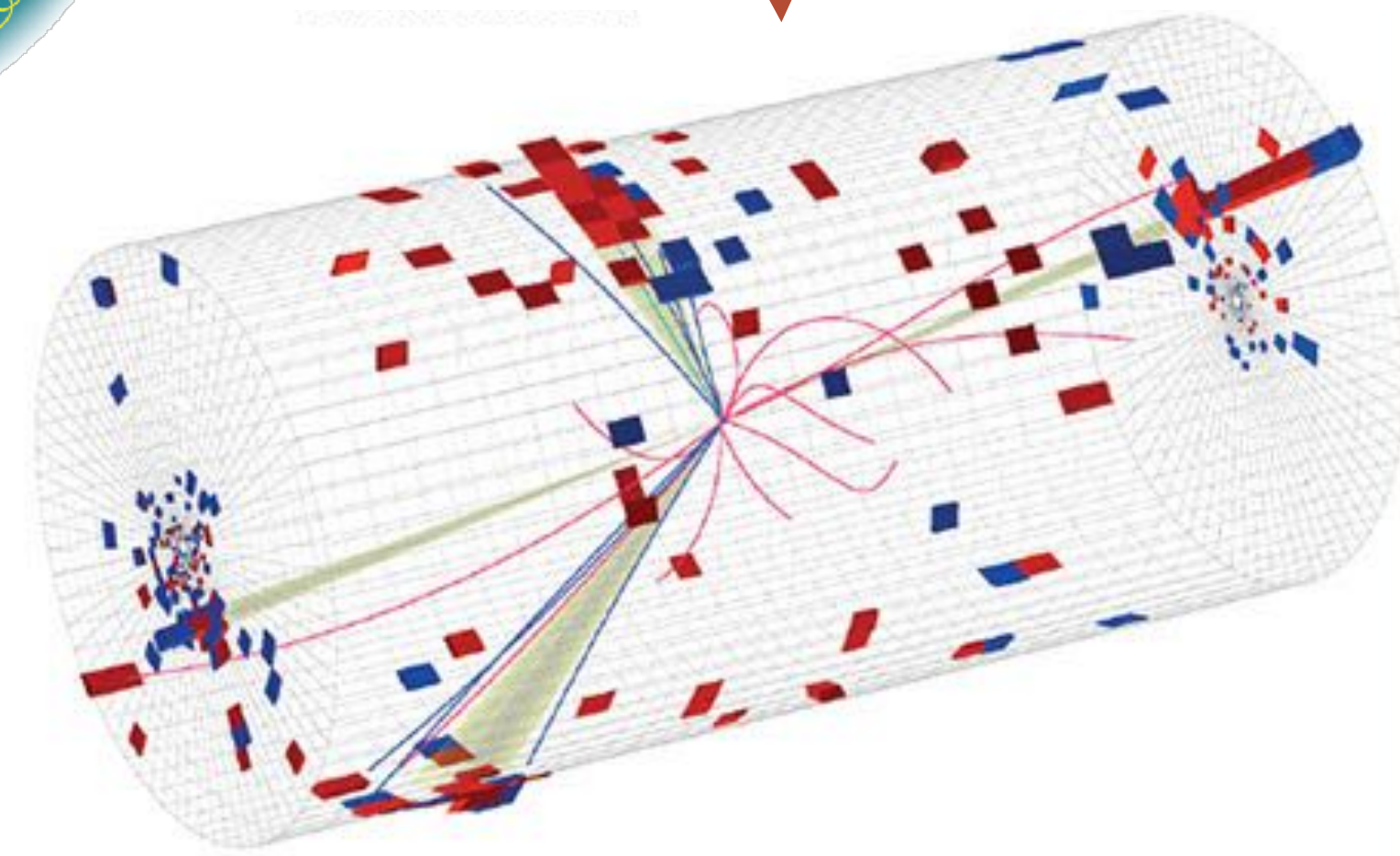
WHY IS IT INTERESTING TO MAP THE NUCLEON?

$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{\partial} - g \not{A} + m) \psi_q - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$

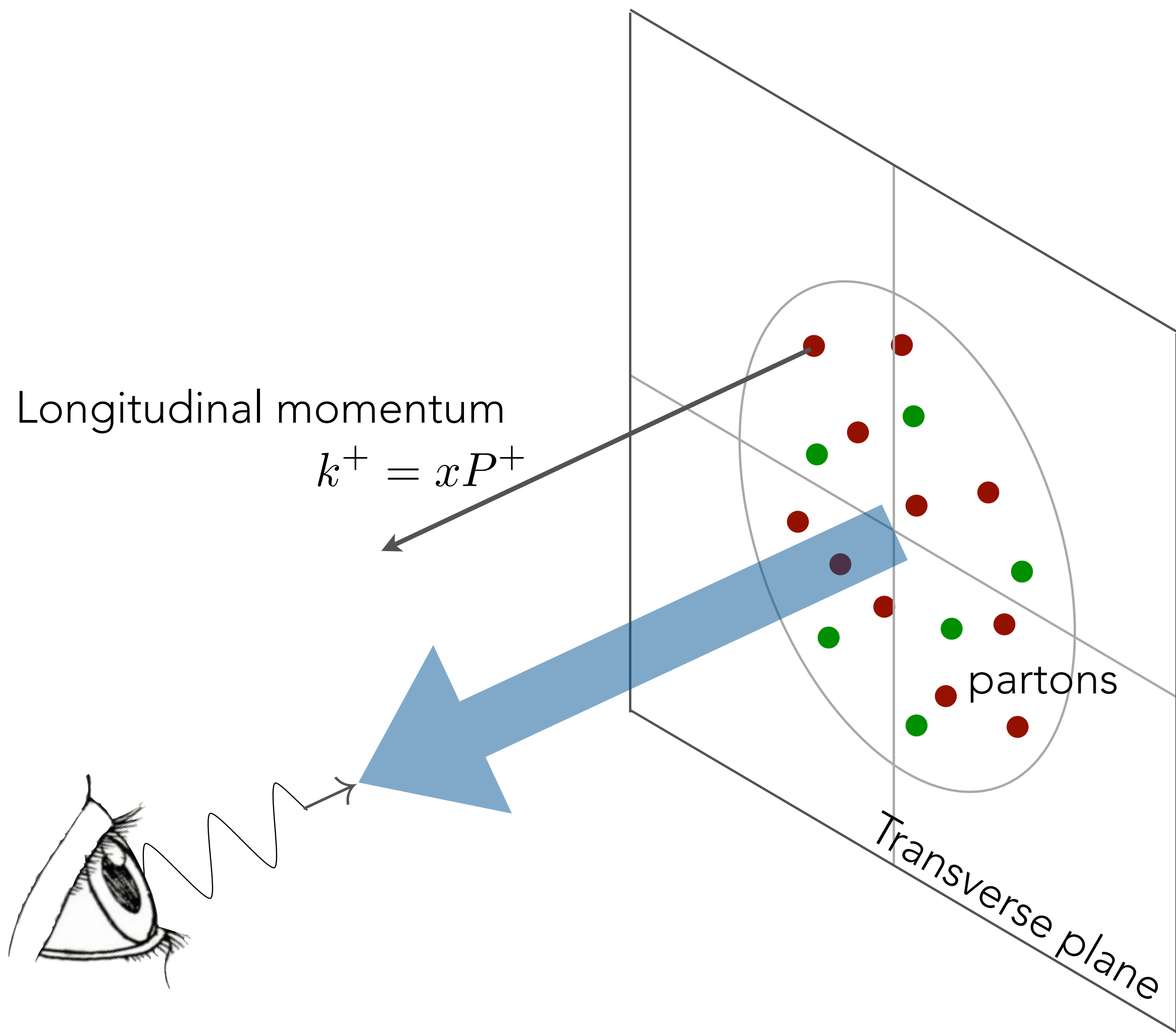
Check predictions



Make predictions



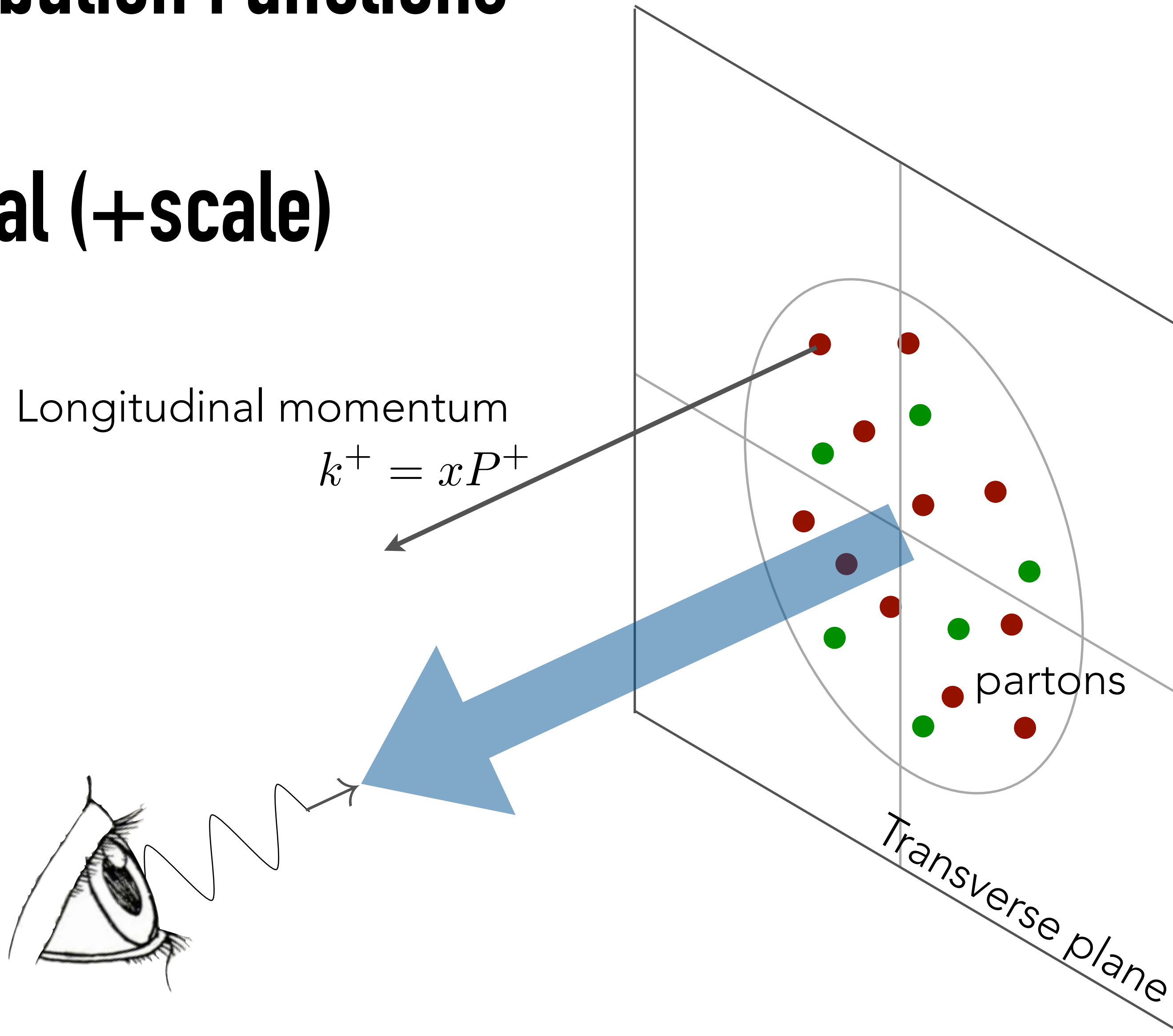
PRESENT KNOWLEDGE

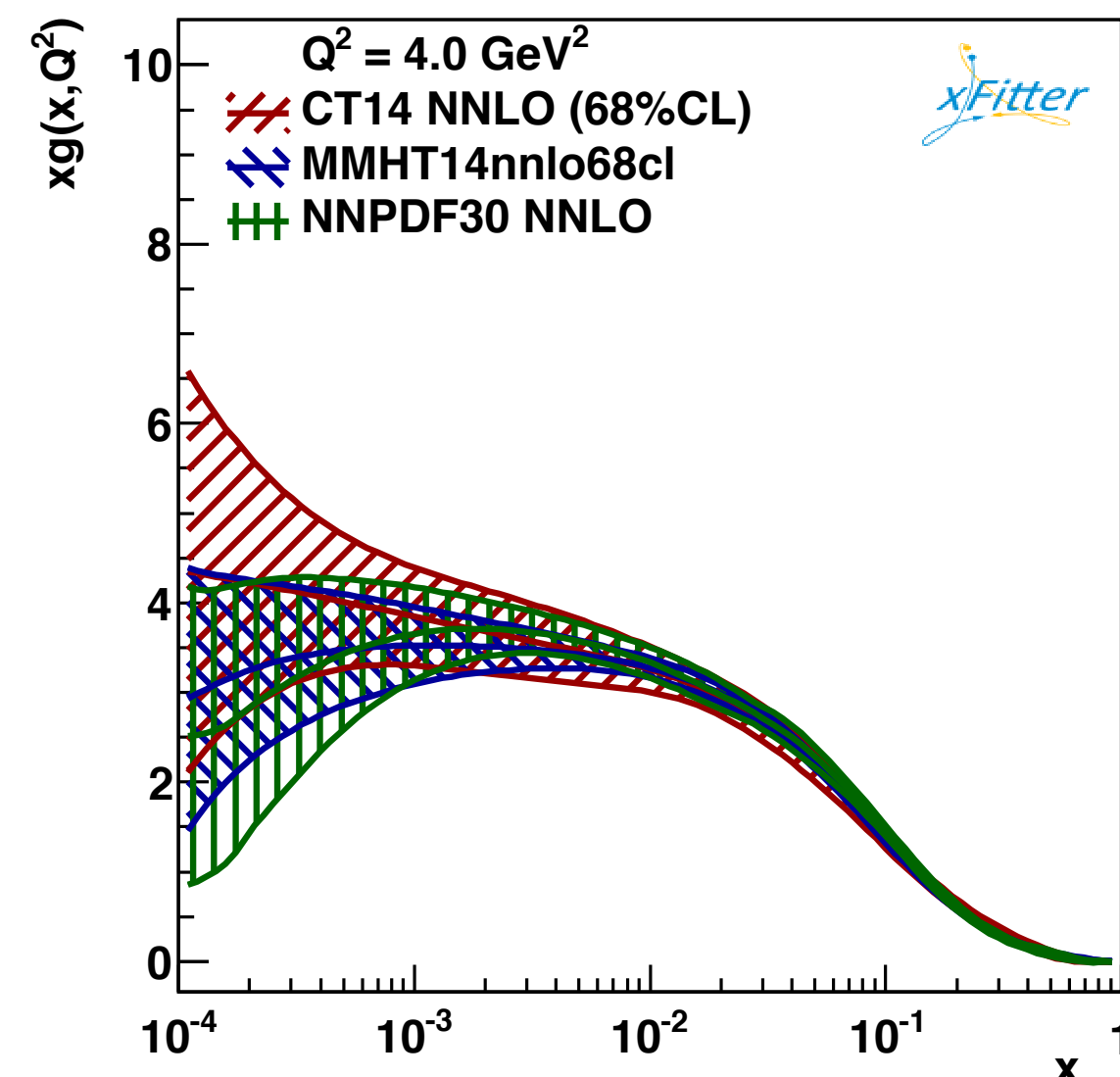
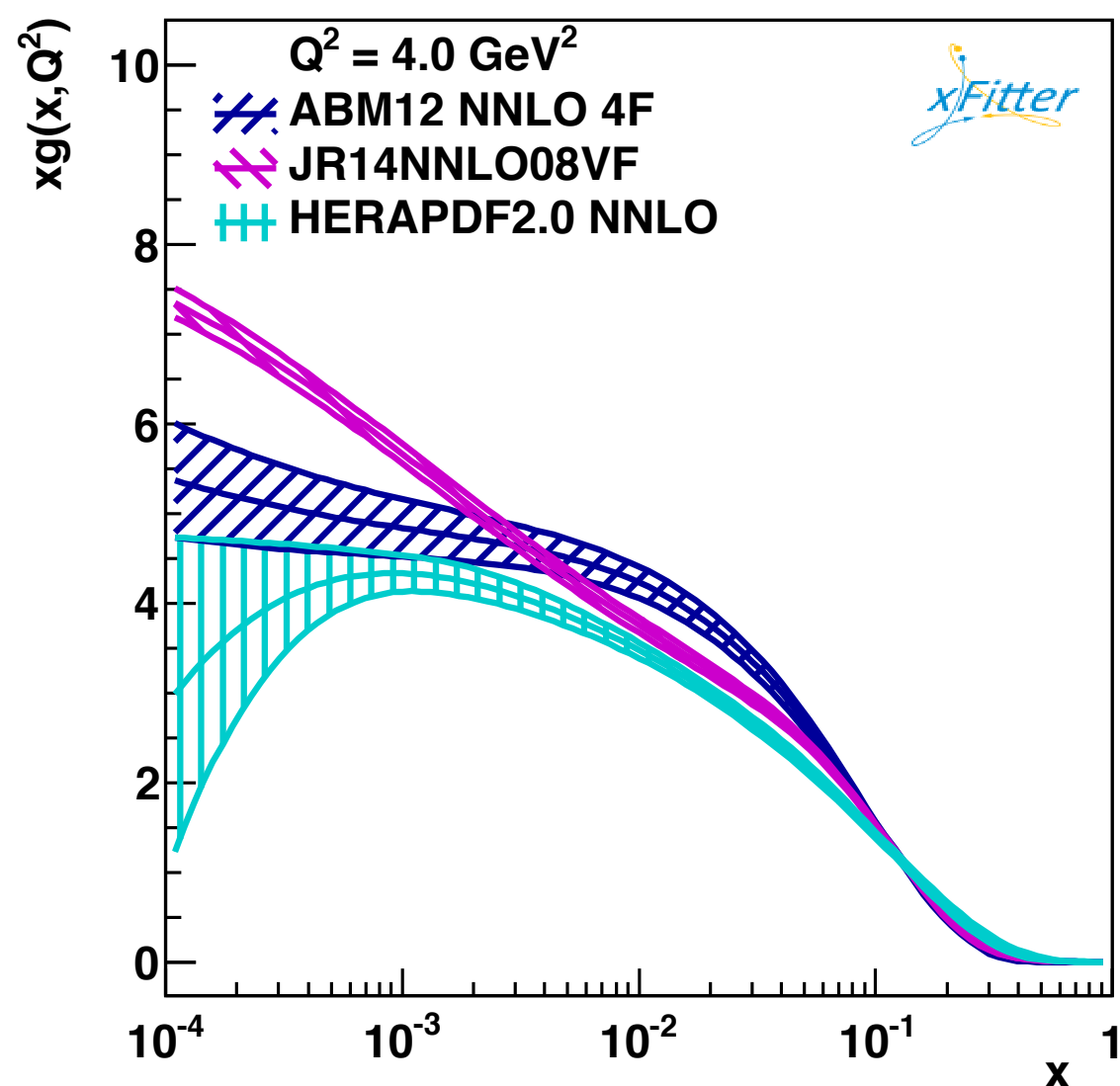
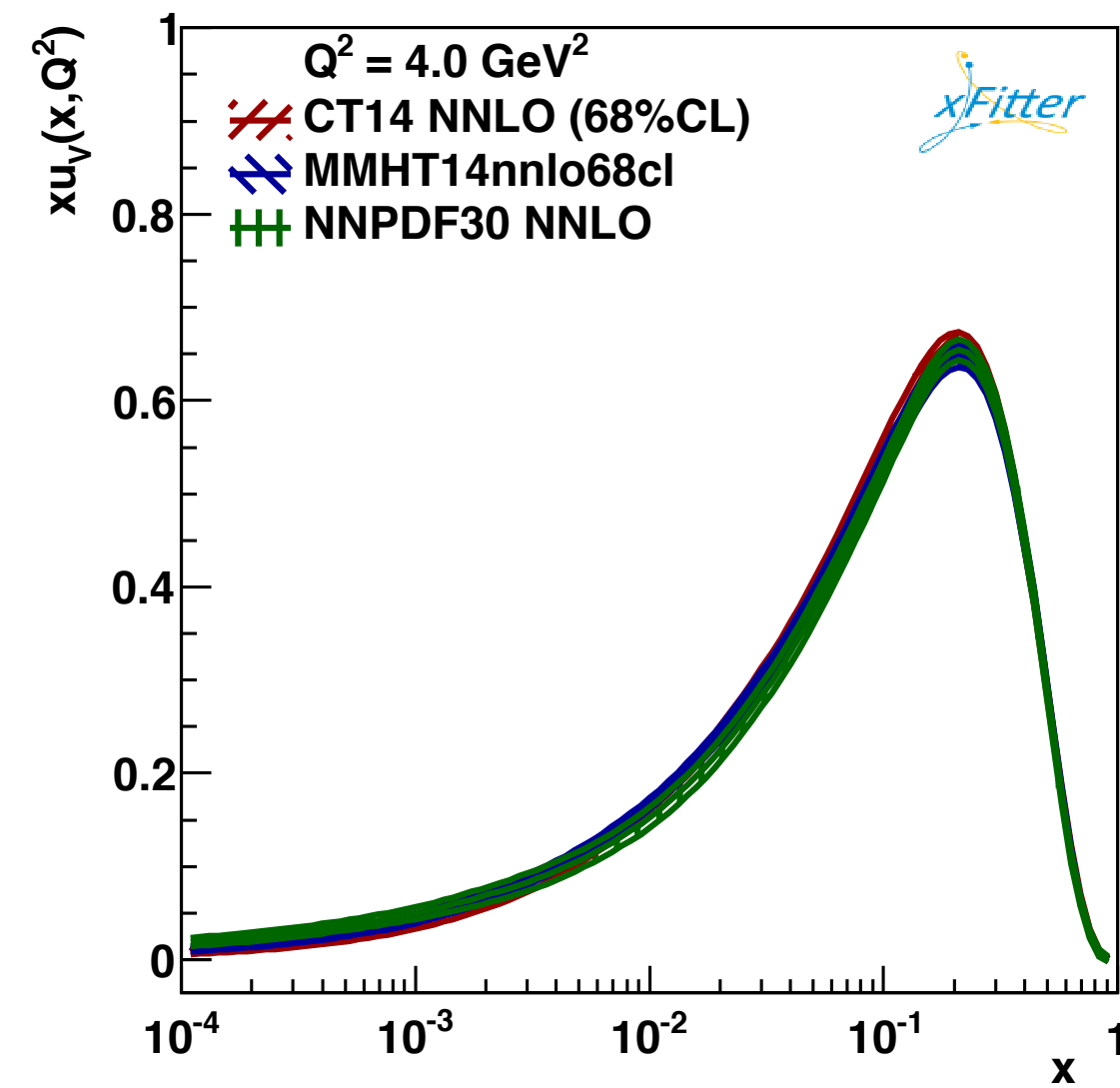
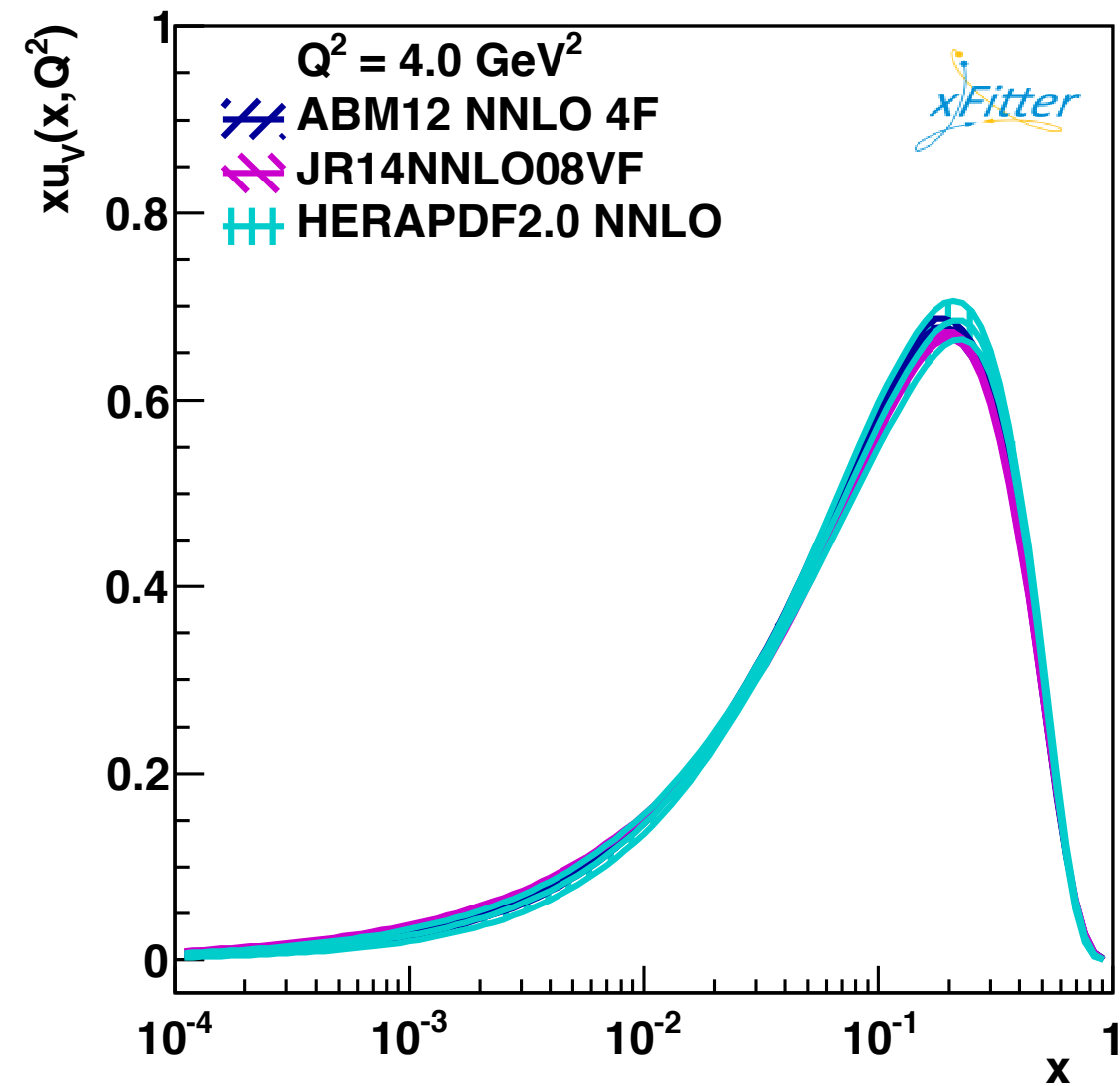


Parton Distribution Functions

$$f(x)$$

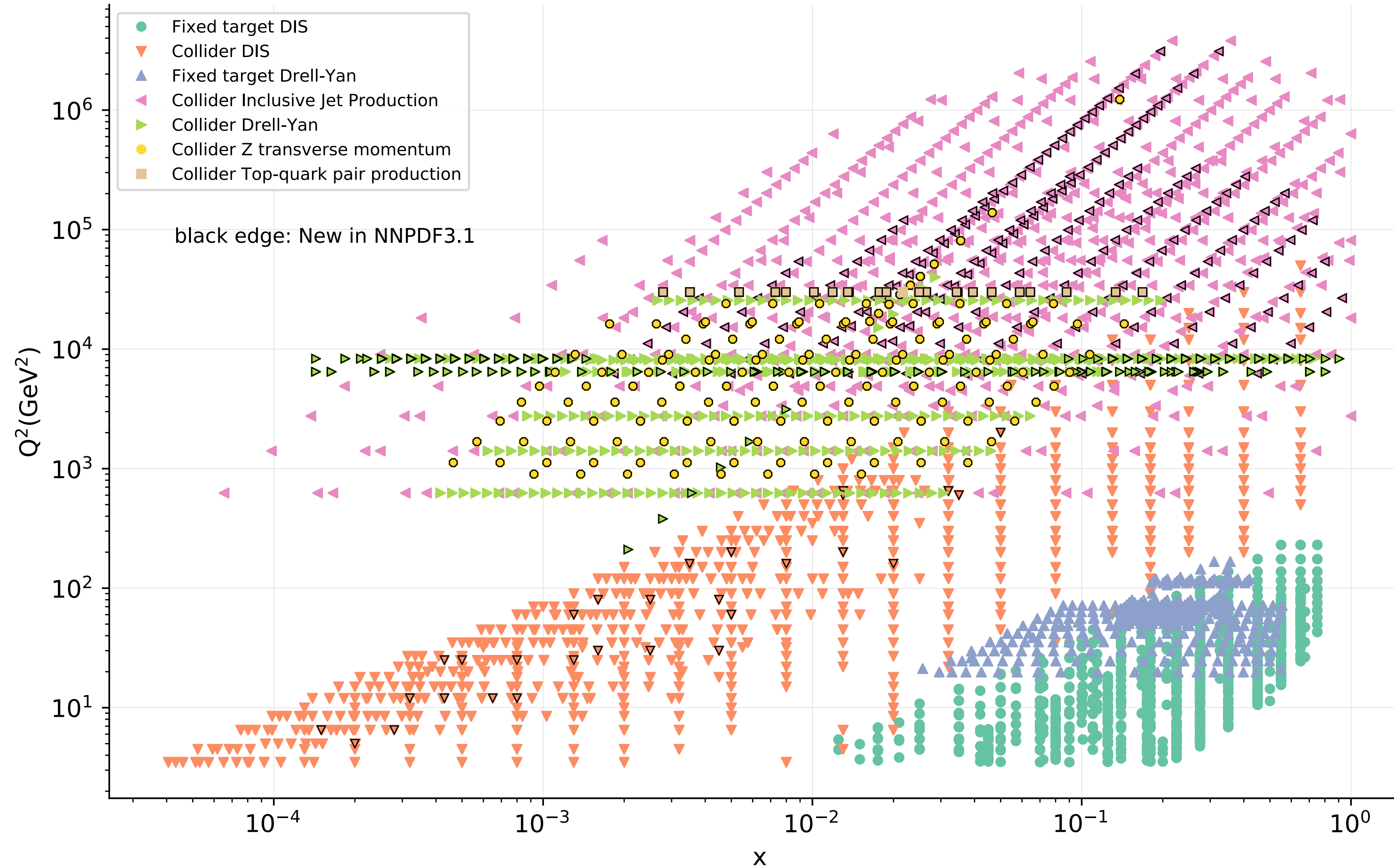
1 dimensional (+scale)

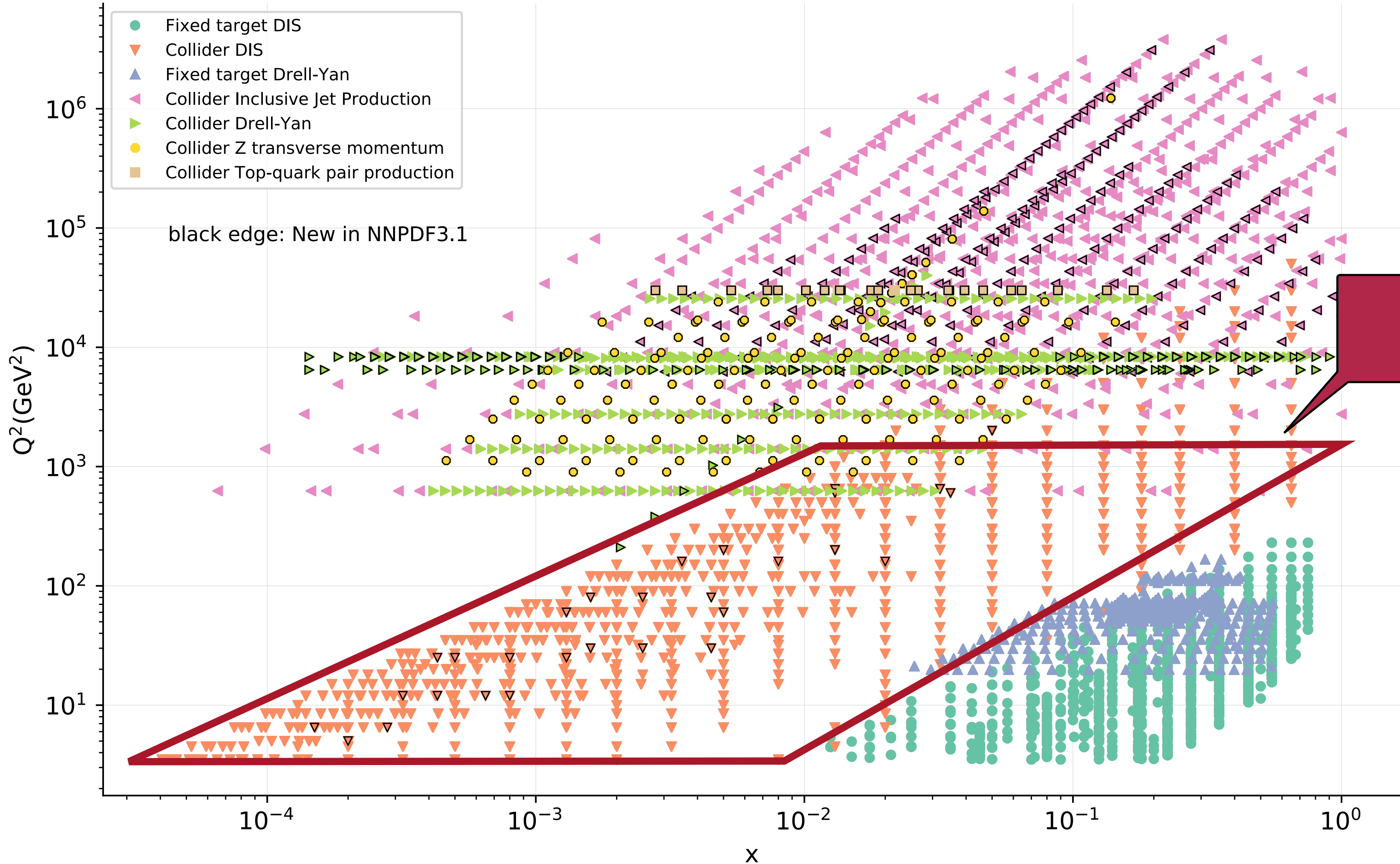




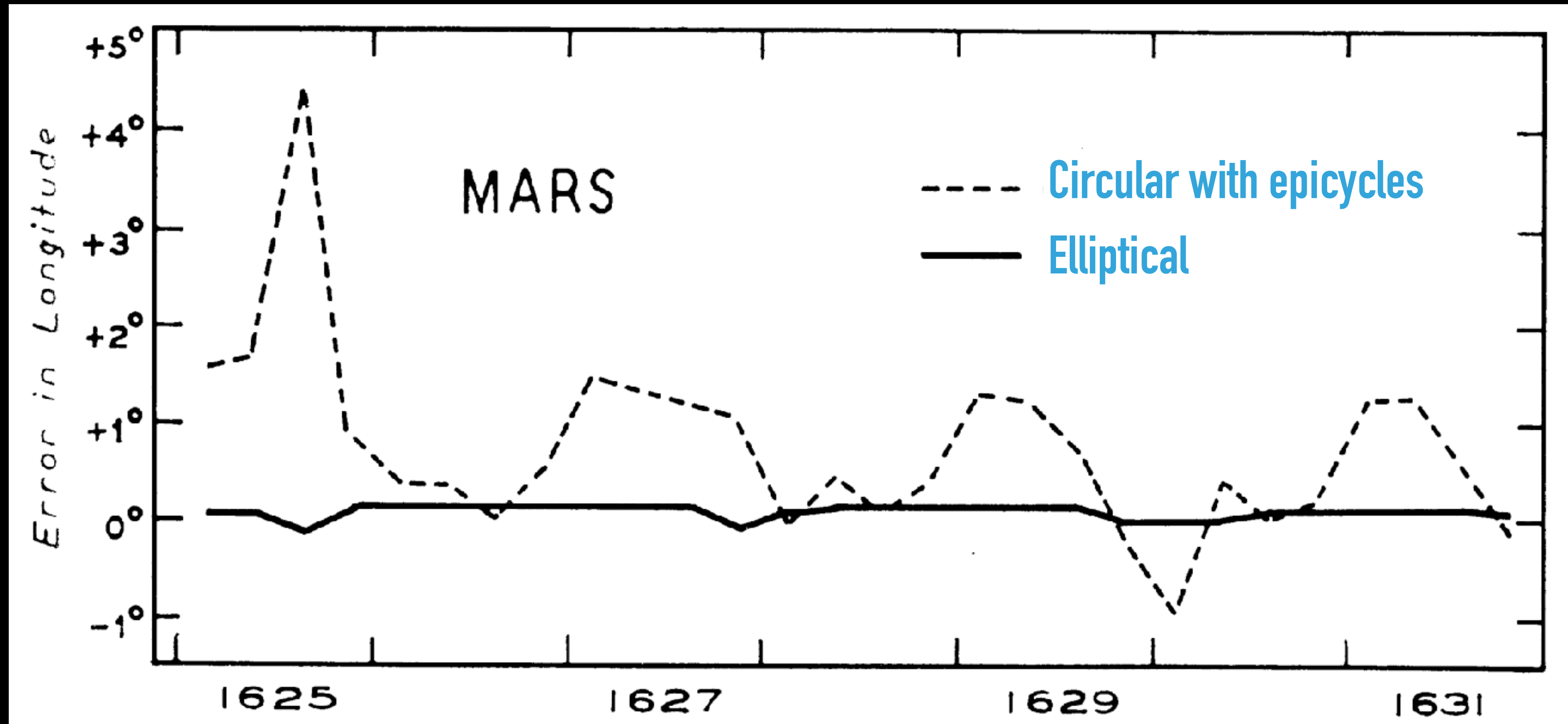
Standard collinear PDFs describe the distribution of partons in one dimension in momentum space.

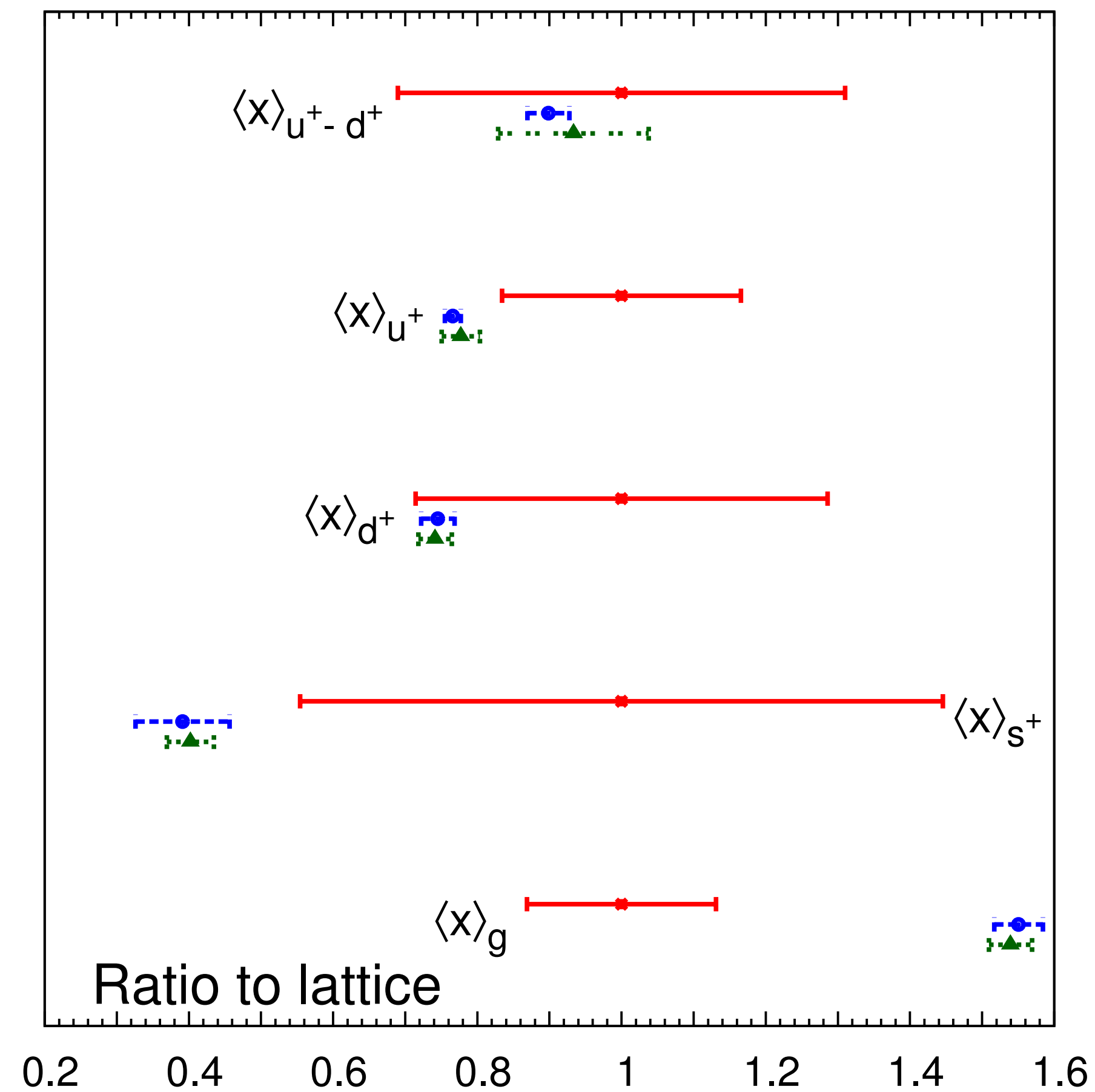
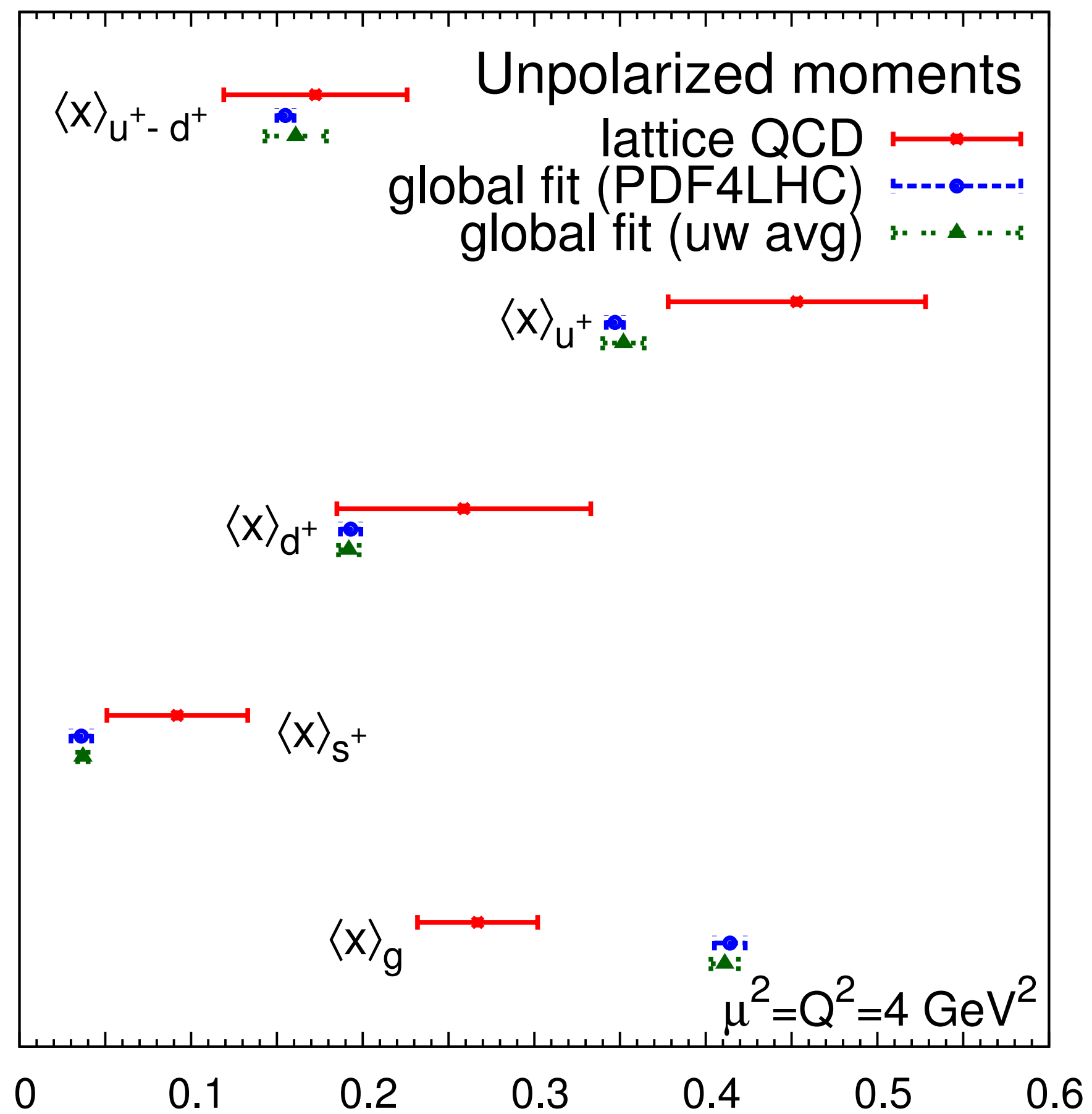
They are extracted through global fits

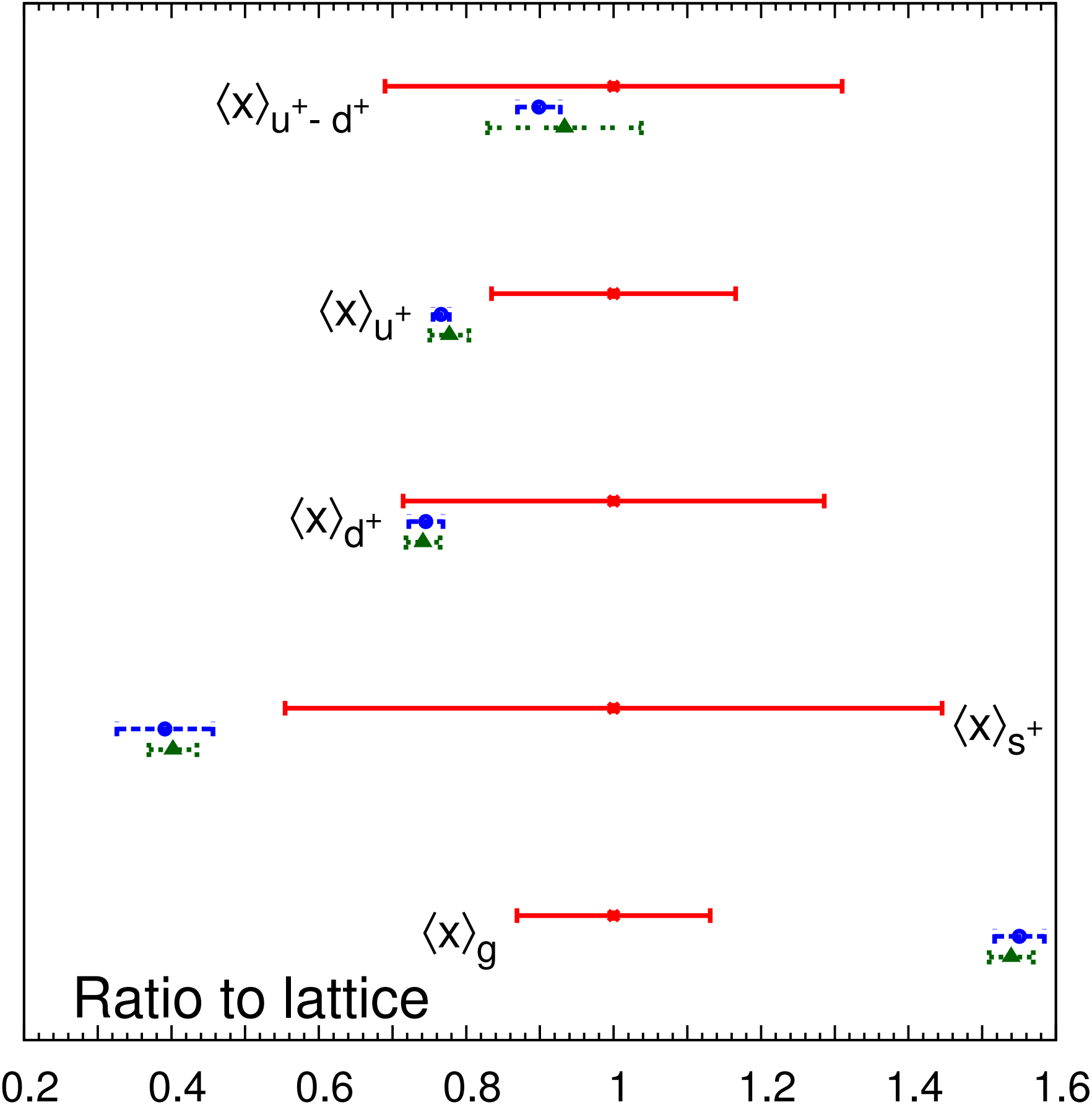
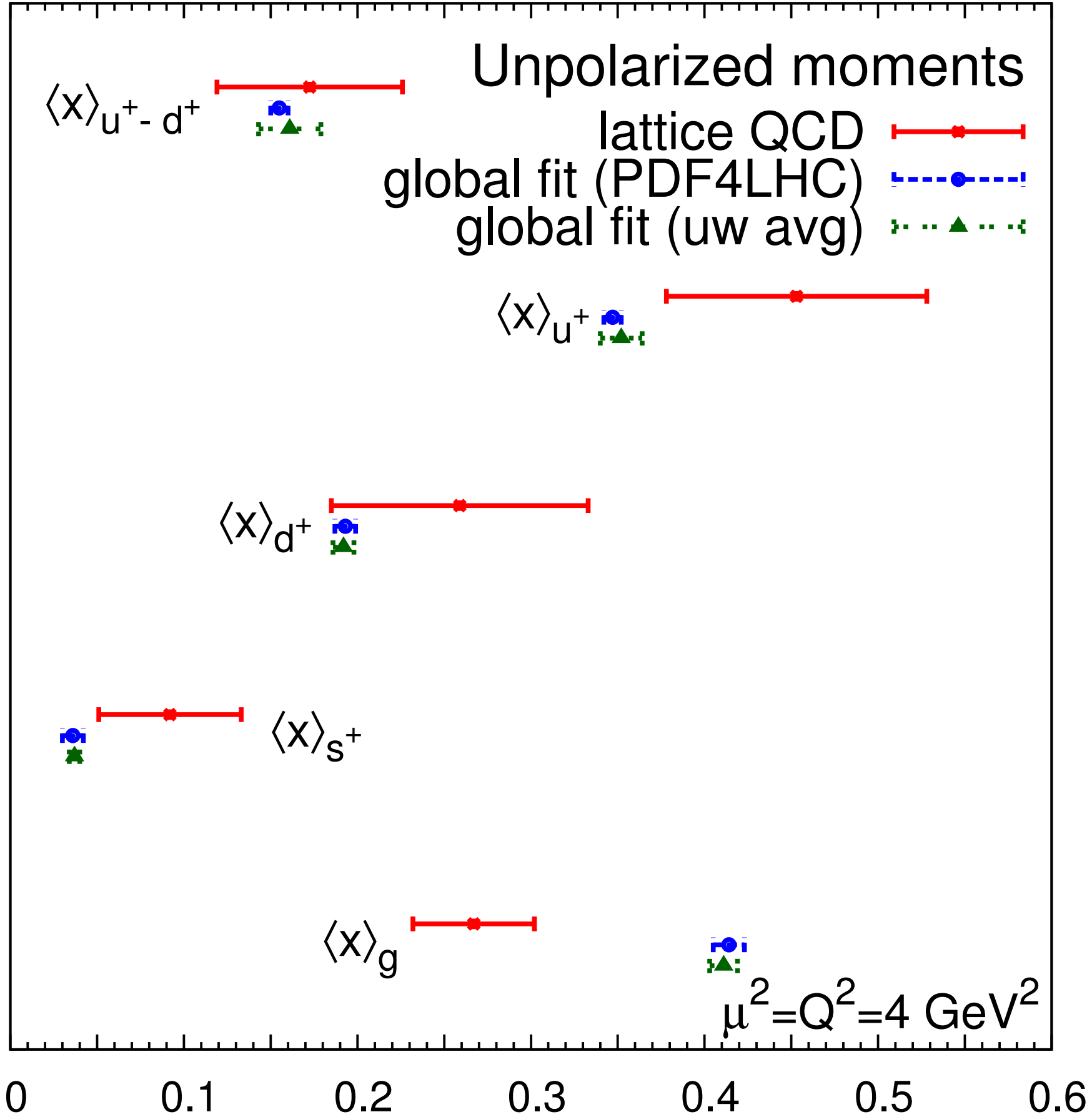




CHECK PREDICTIONS



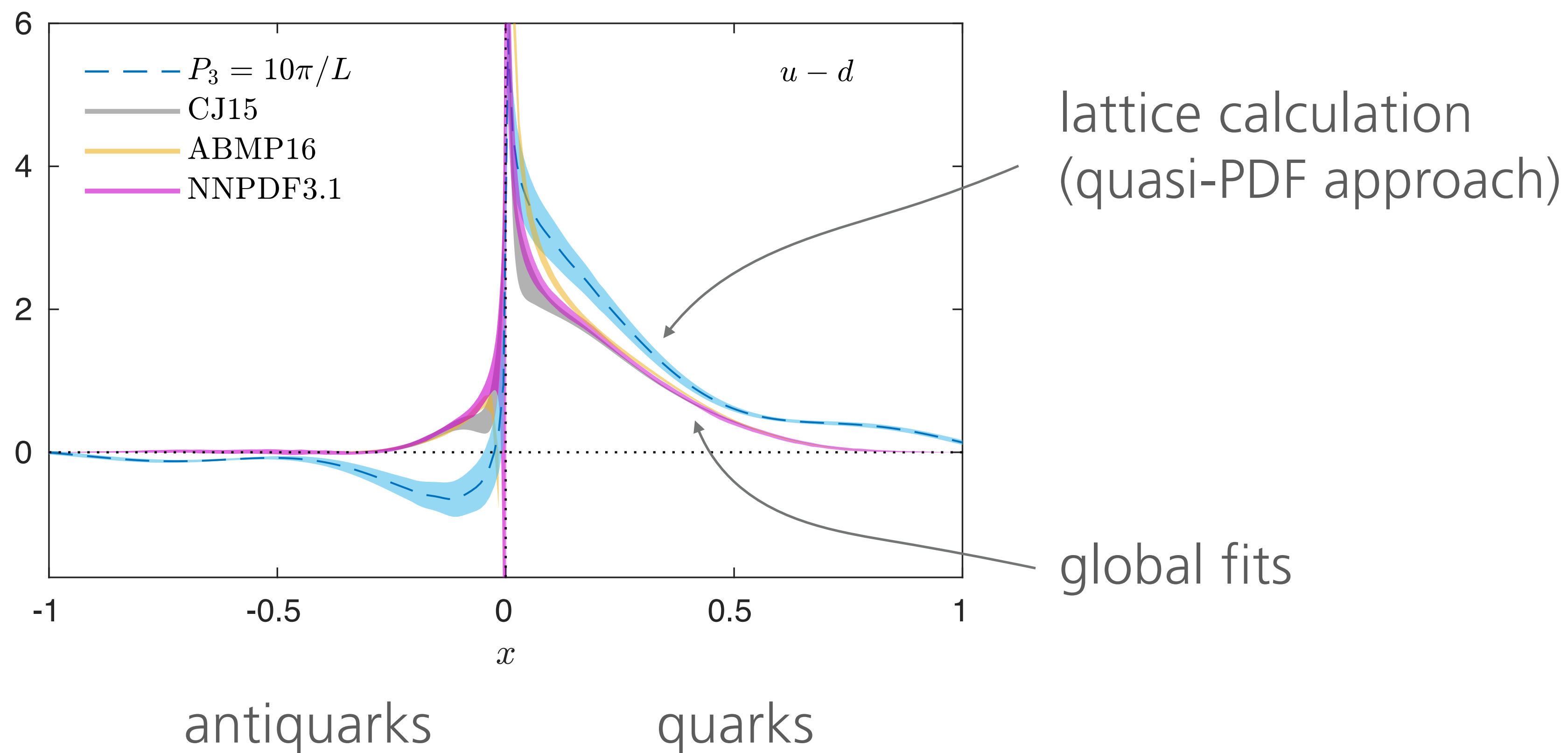


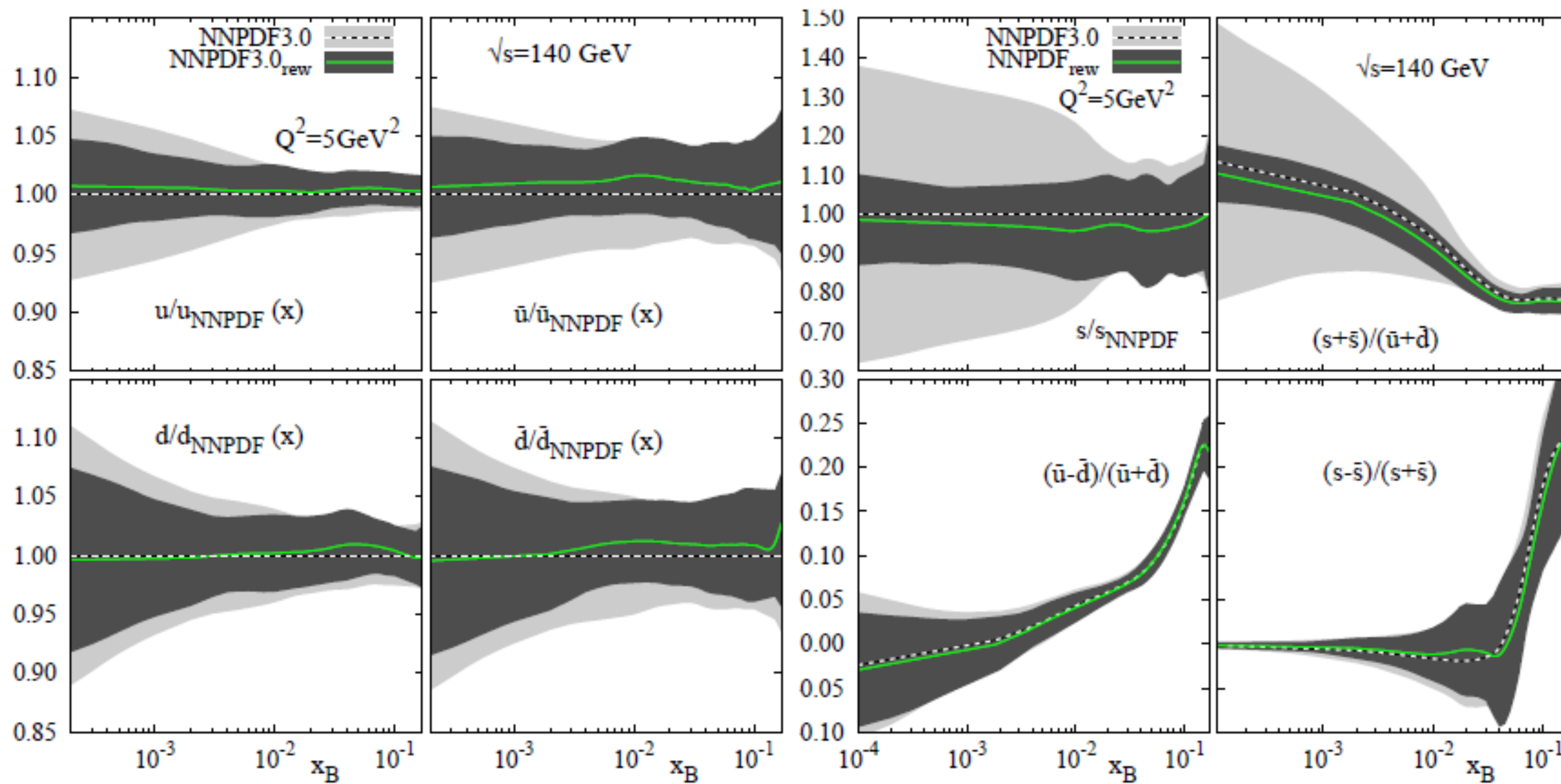


PDFLattice White Paper, arXiv:1711.07916

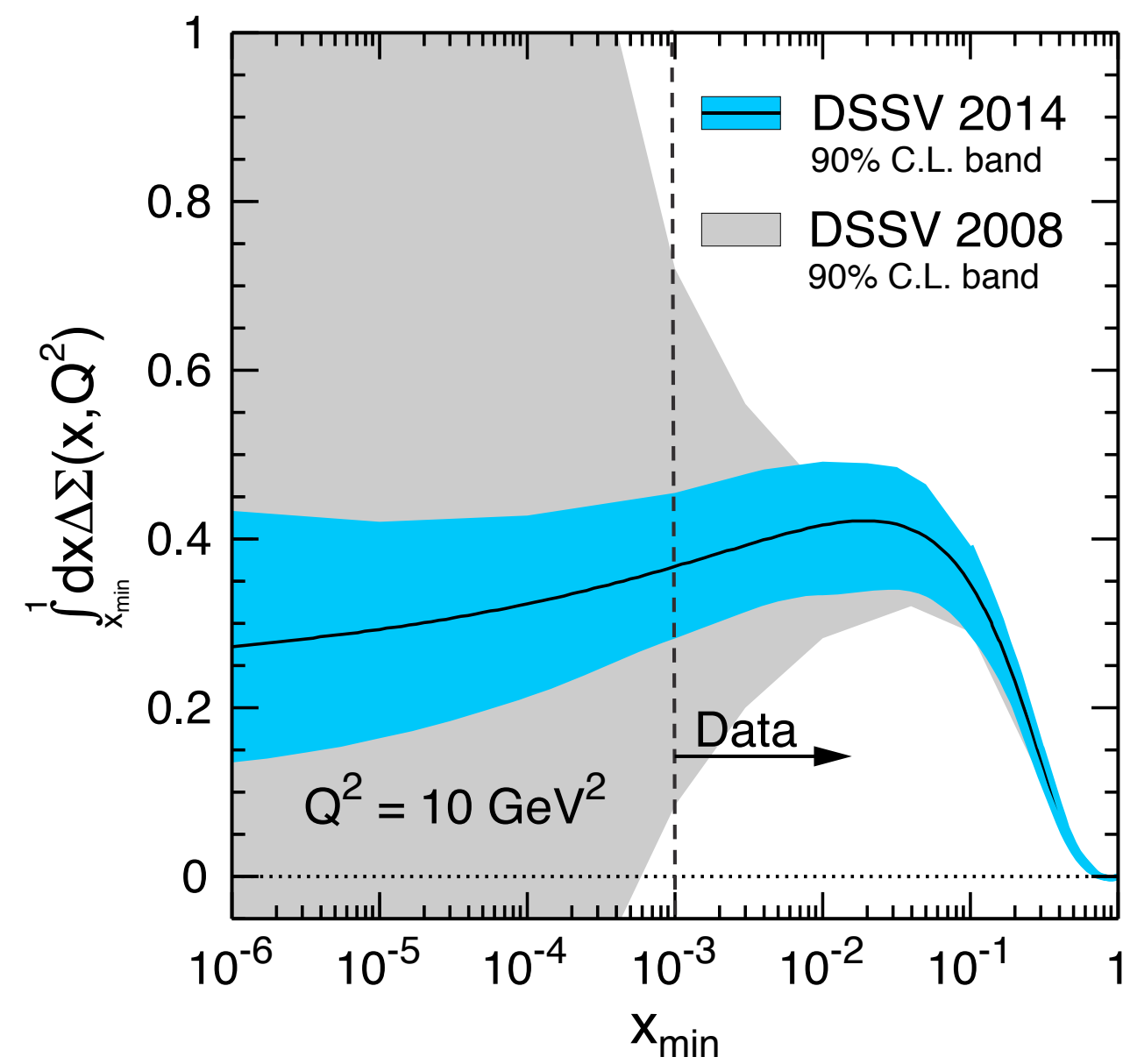
Fair agreement, but far from perfect

Alexandrou, Cichy, Constantinou, Hadjiyiannakou, Jansen, Scapellato, Steffens, arXiv:1902.00587

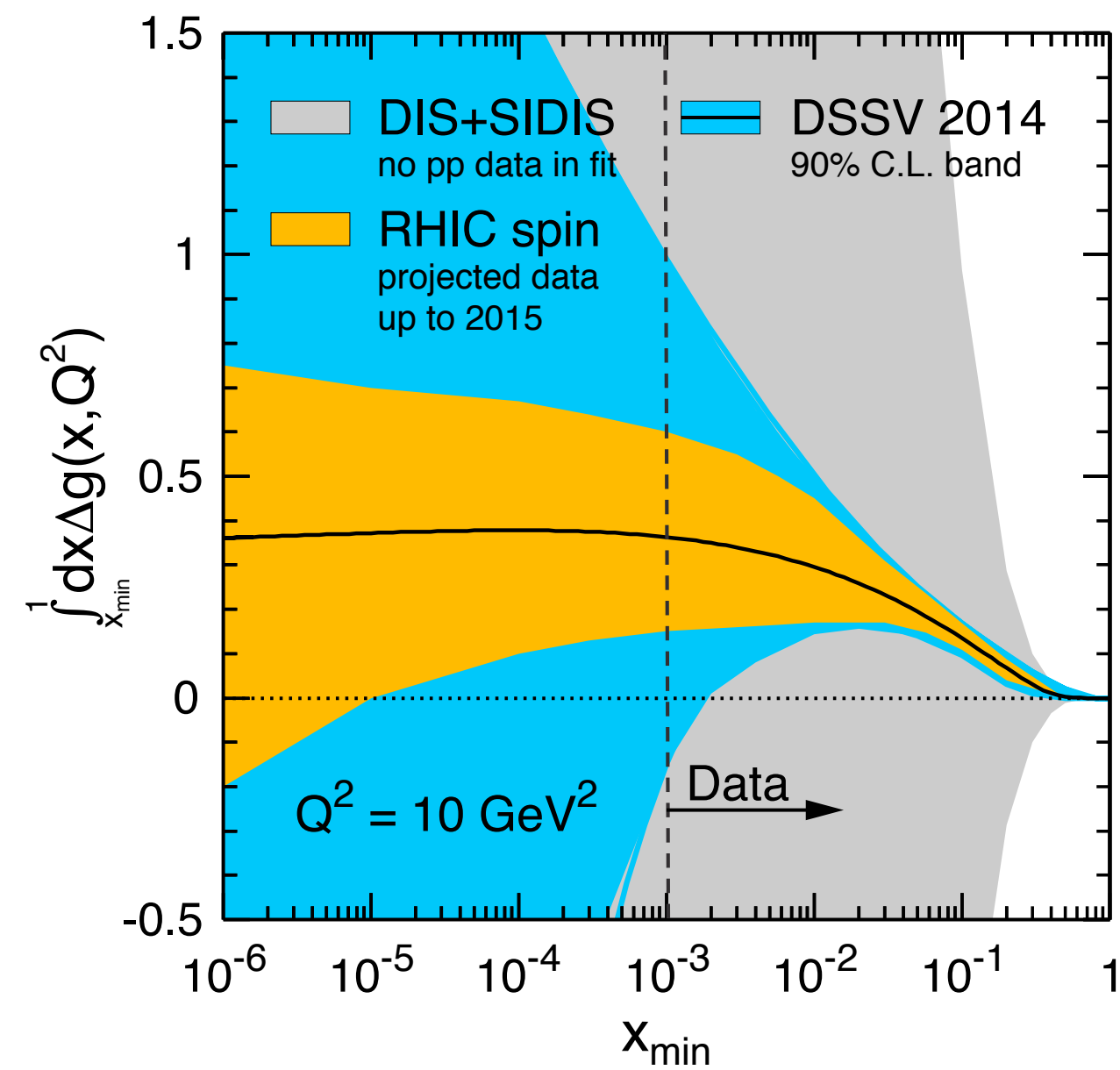




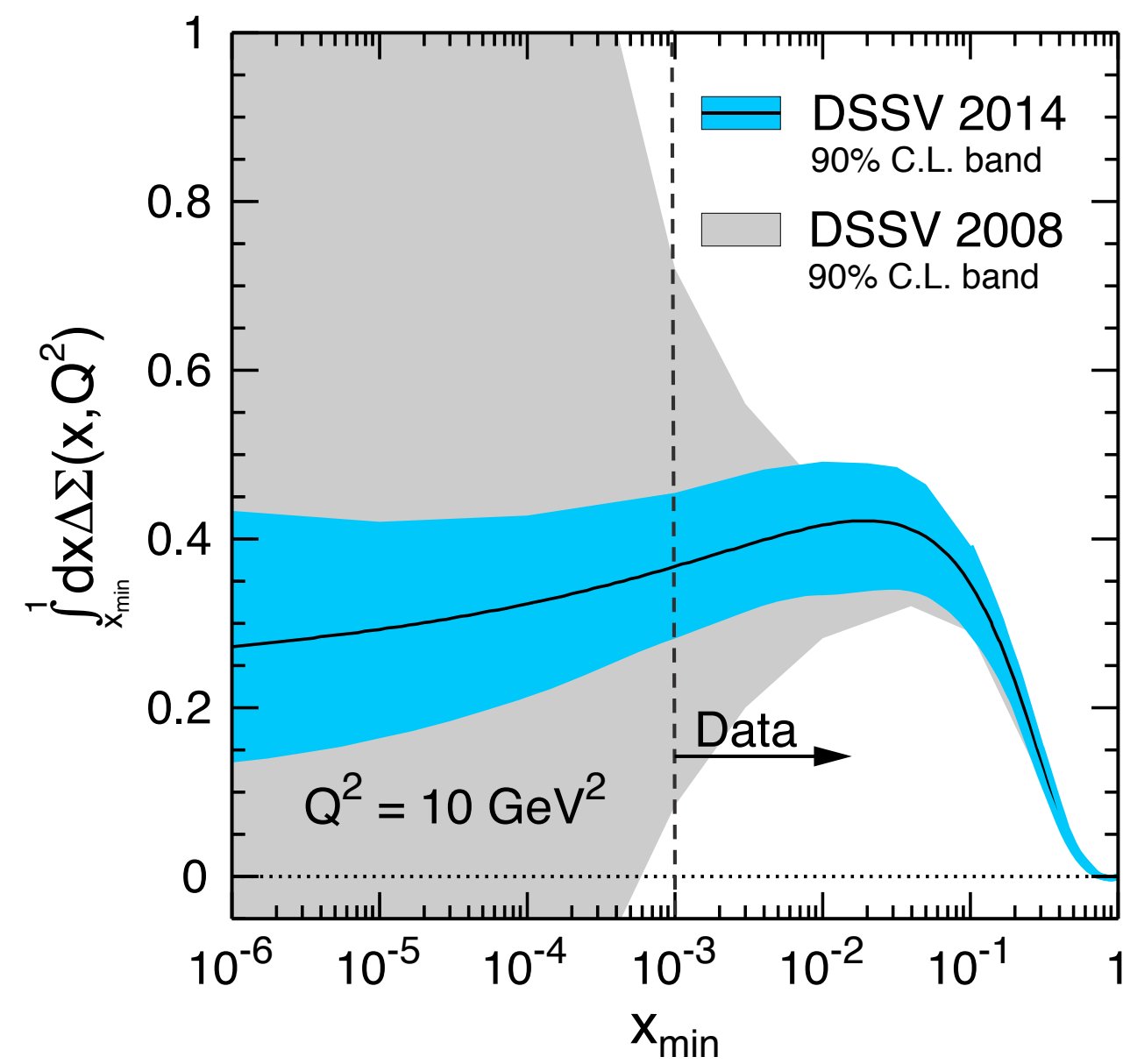
quark spin



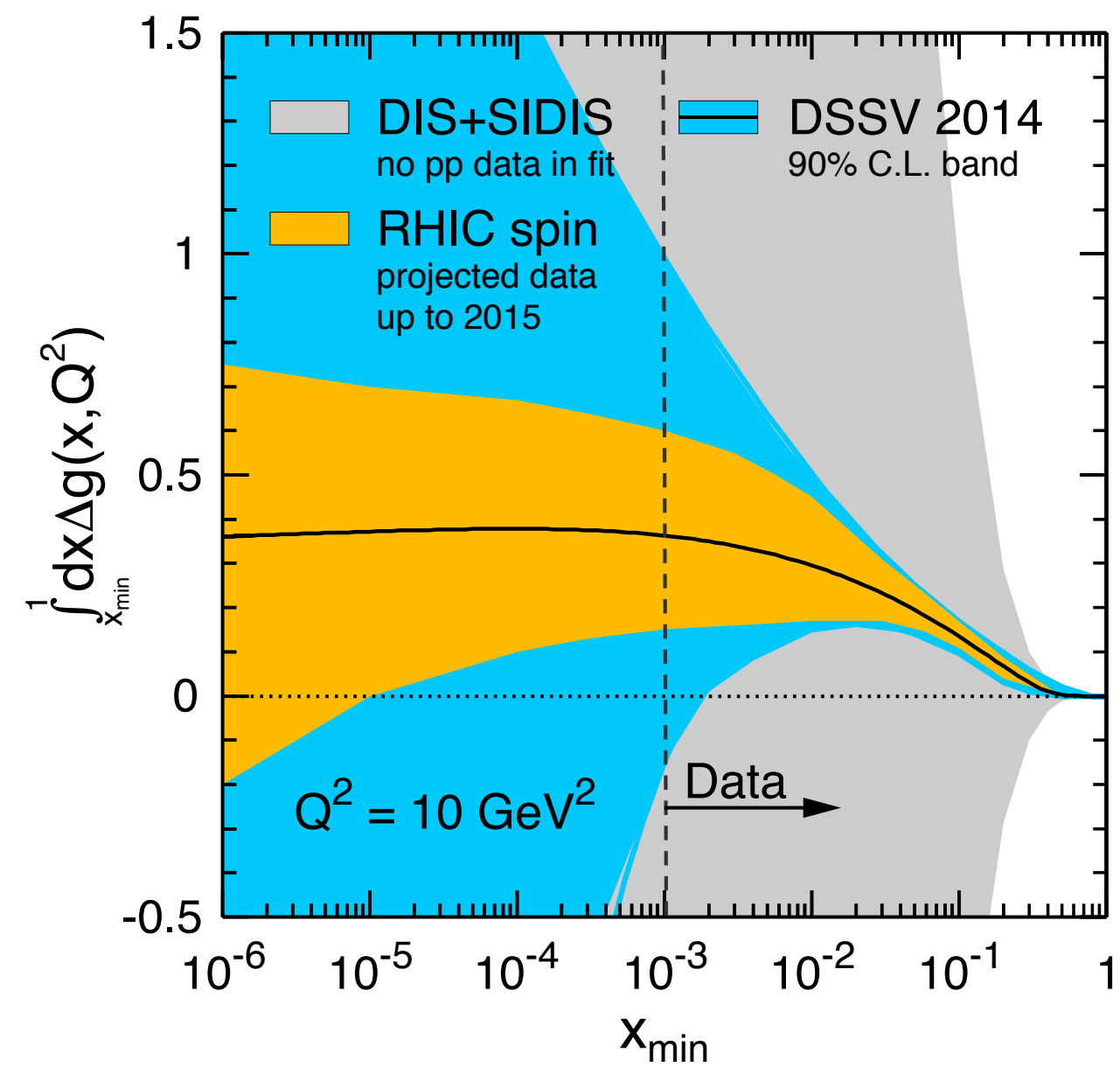
gluon spin



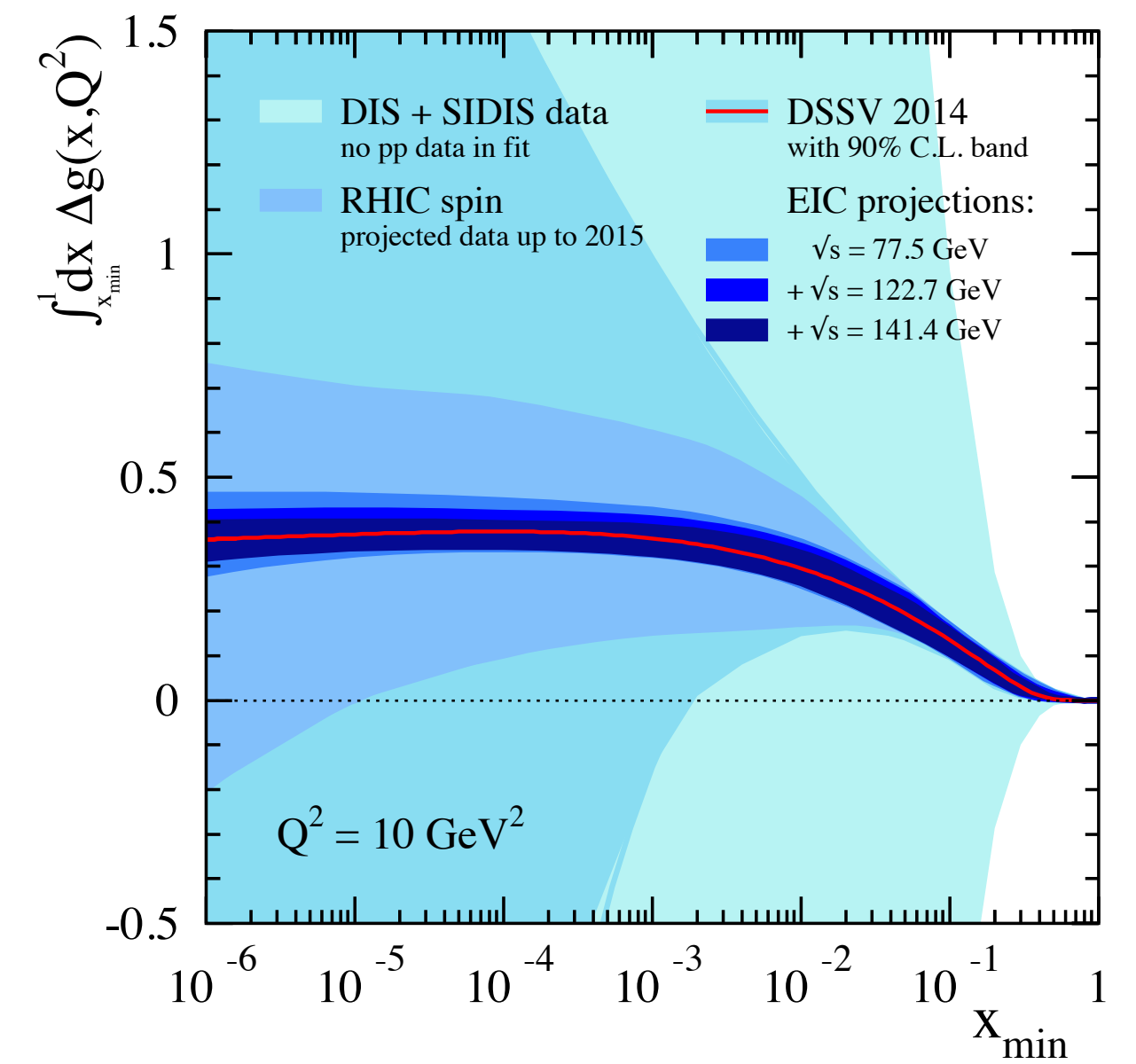
quark spin



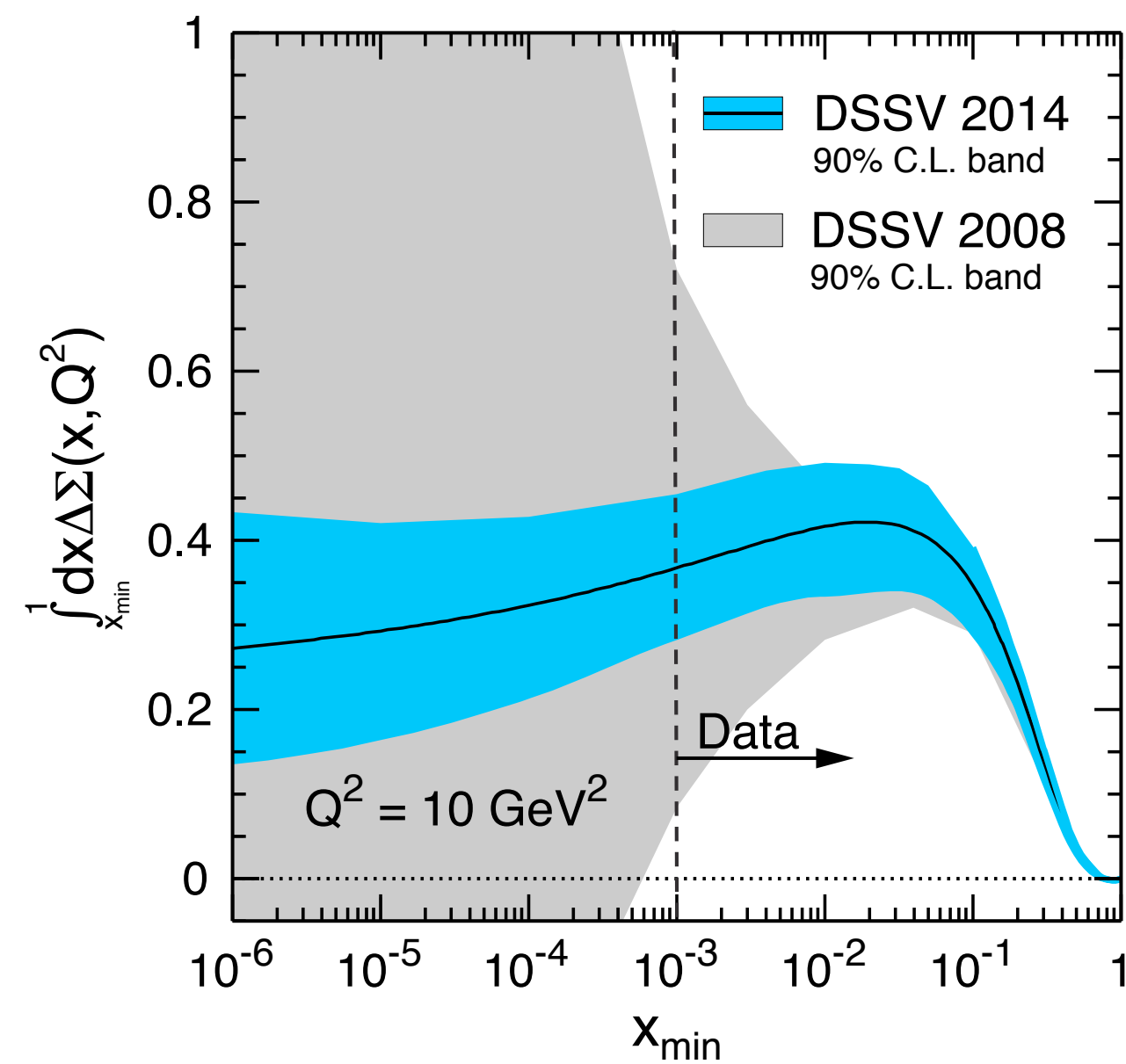
gluon spin



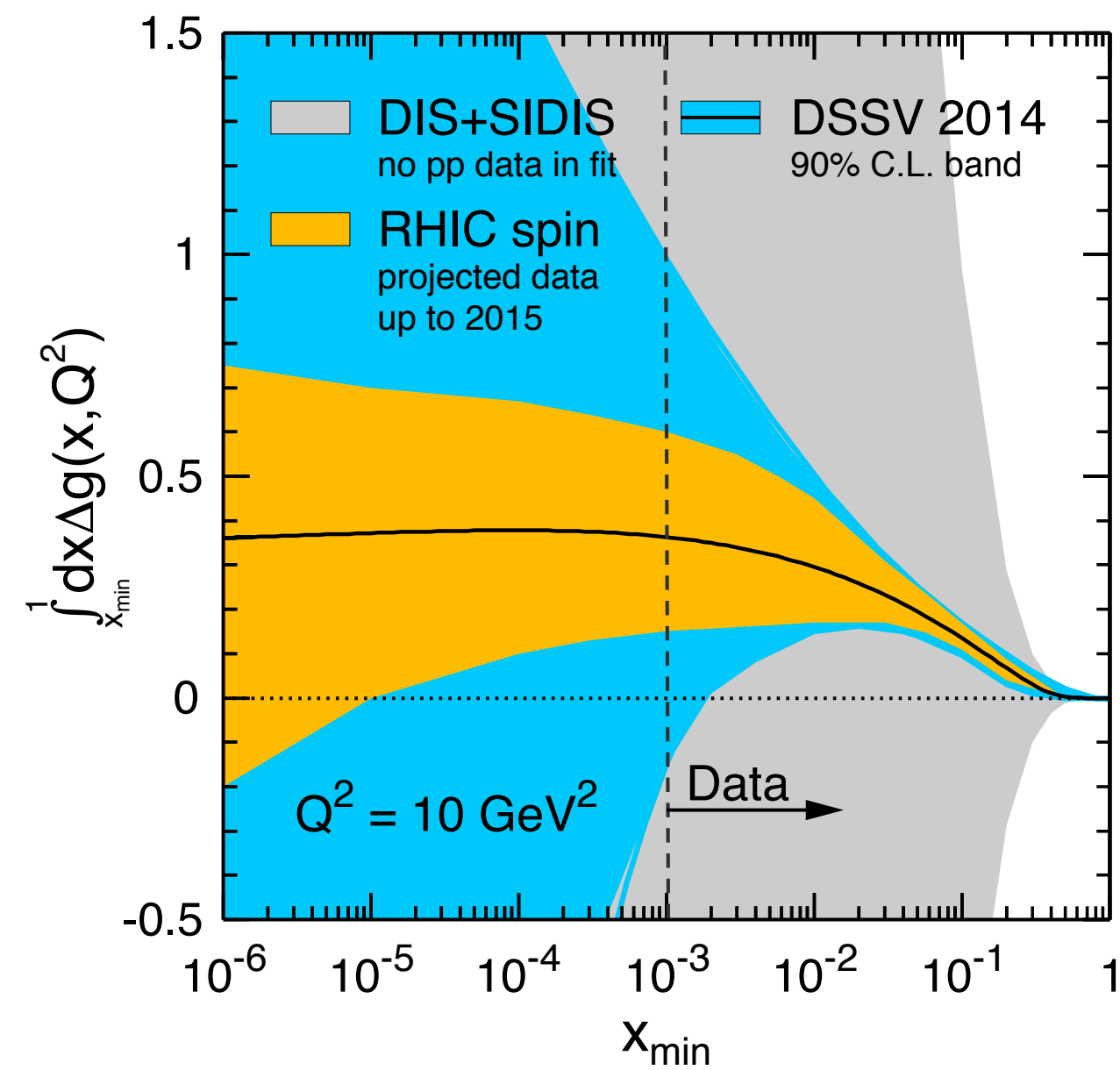
gluon spin w/ EIC



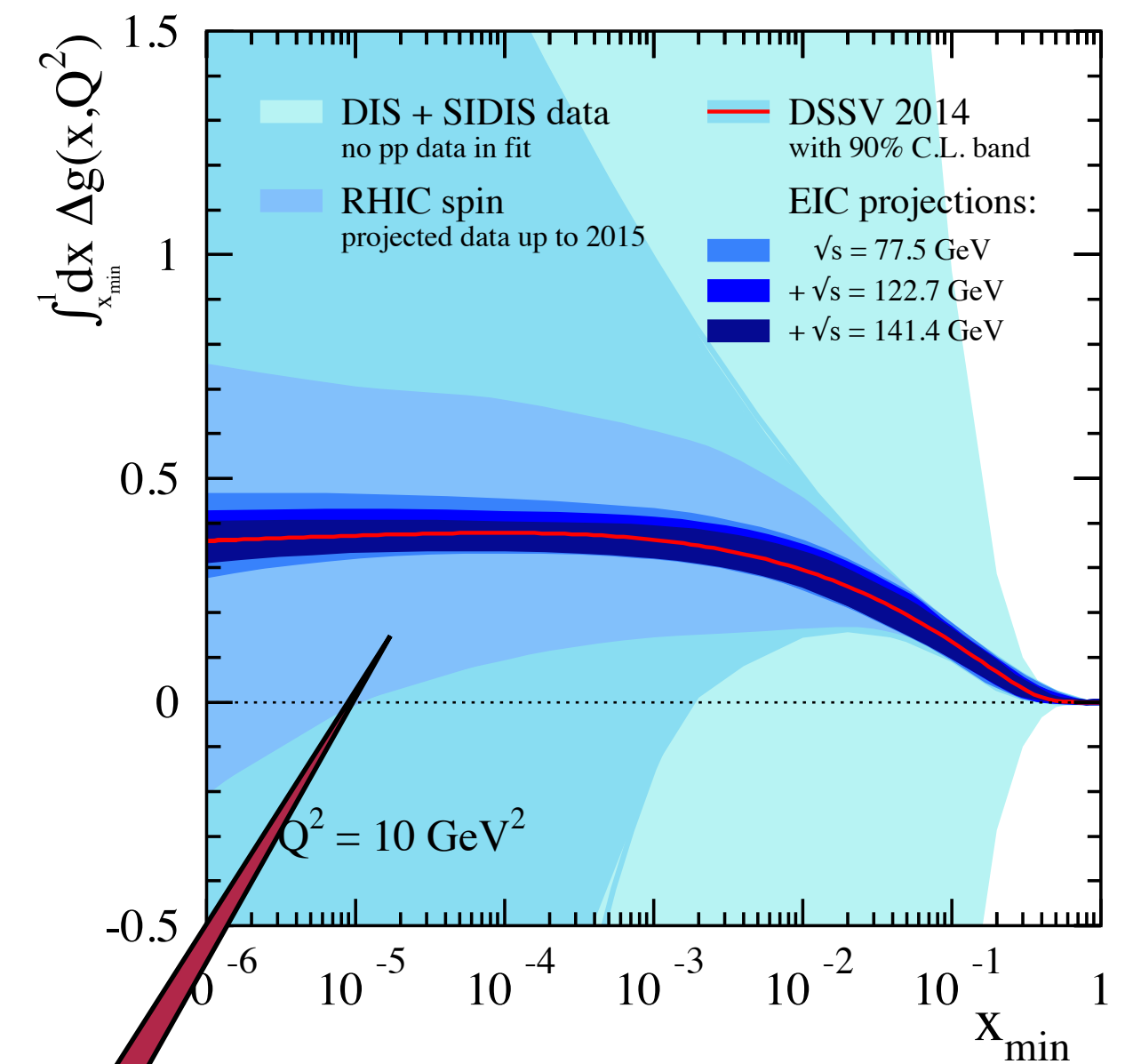
quark spin



gluon spin



gluon spin w/ EIC




Enormous impact

Aschenauer et al., [arXiv:1708.01527](https://arxiv.org/abs/1708.01527) and [arXiv:1509.06489](https://arxiv.org/abs/1509.06489)

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta g + L$$

Aschenauer et al., [arXiv:1708.01527](https://arxiv.org/abs/1708.01527) and [arXiv:1509.06489](https://arxiv.org/abs/1509.06489)

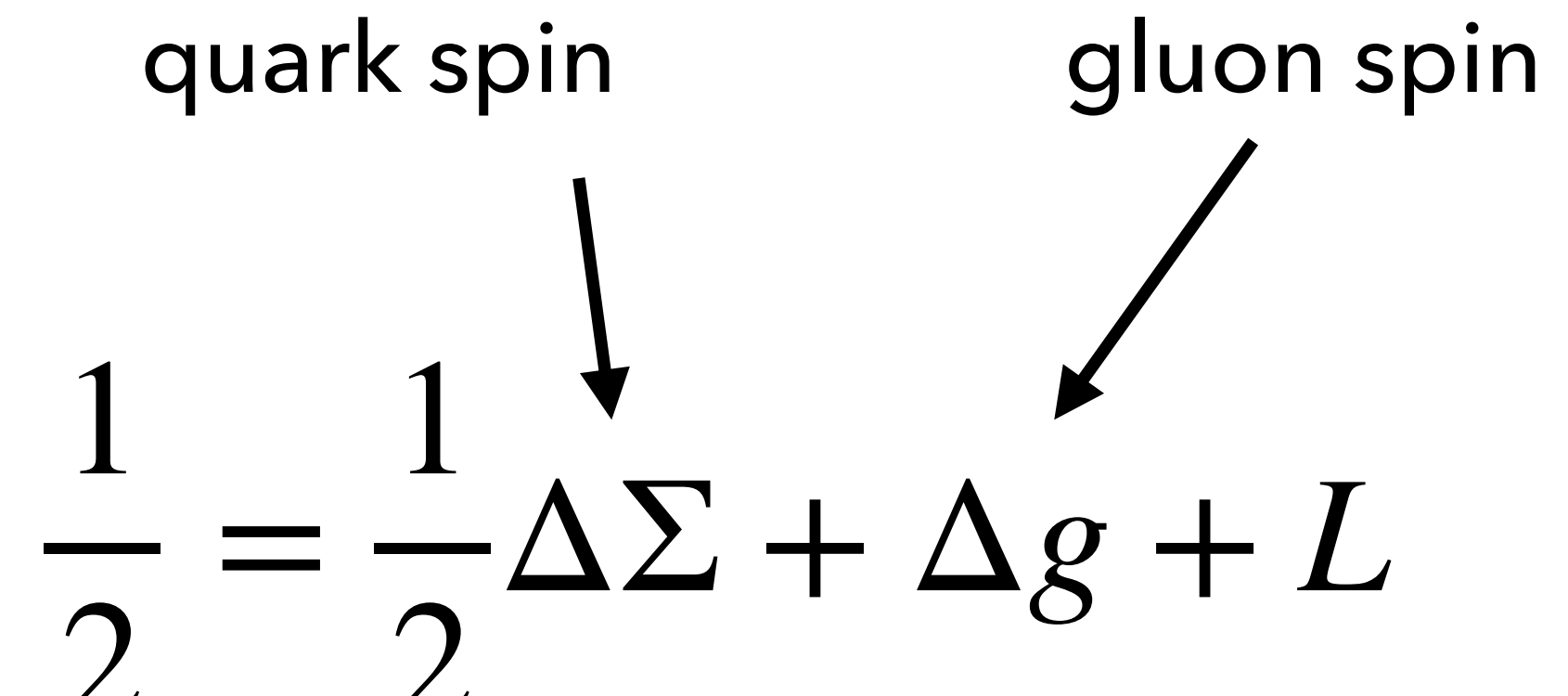
quark spin

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta g + L$$


Aschenauer et al., [arXiv:1708.01527](https://arxiv.org/abs/1708.01527) and [arXiv:1509.06489](https://arxiv.org/abs/1509.06489)

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta g + L$$

quark spin gluon spin



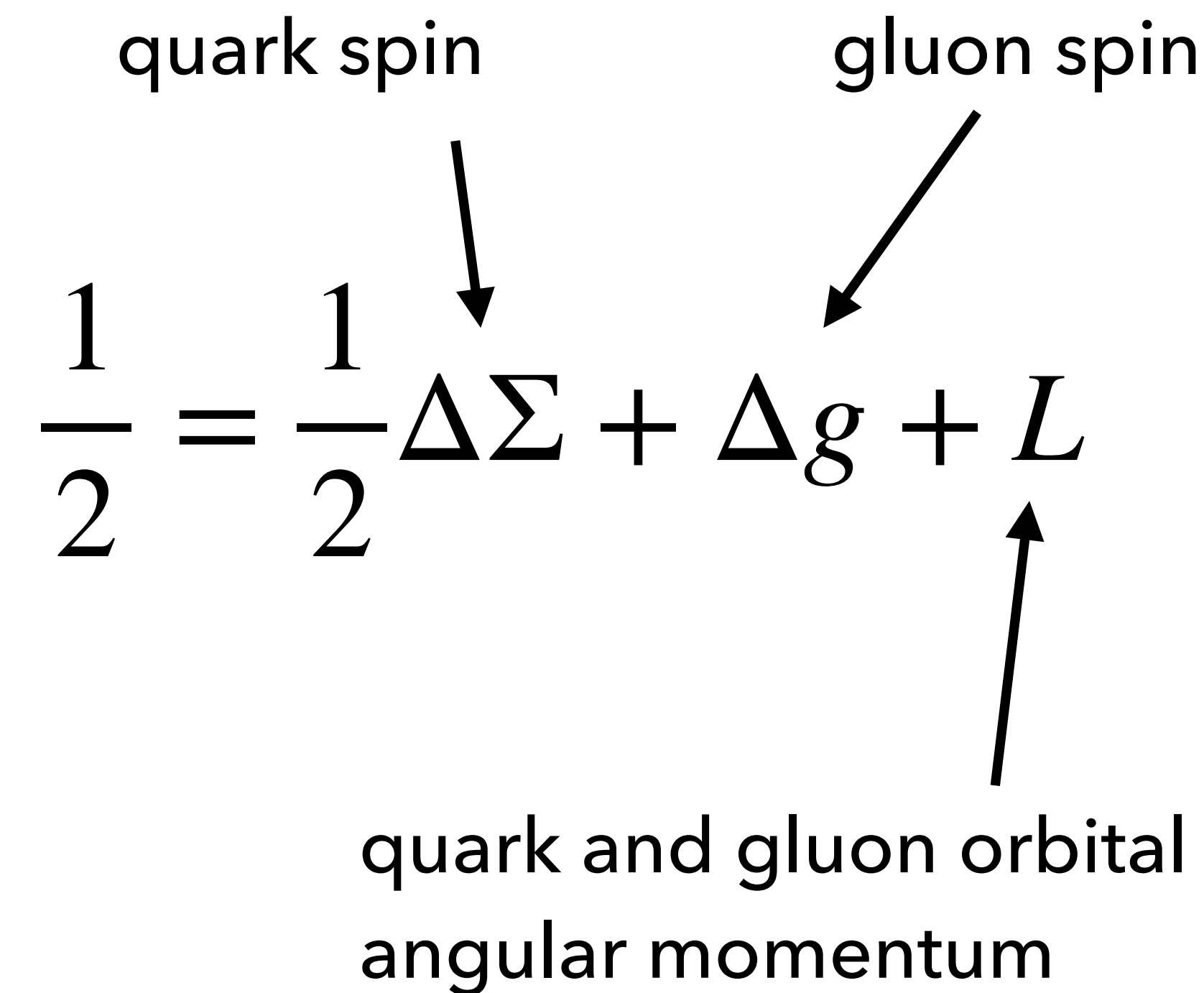
Aschenauer et al., [arXiv:1708.01527](https://arxiv.org/abs/1708.01527) and [arXiv:1509.06489](https://arxiv.org/abs/1509.06489)

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta g + L$$

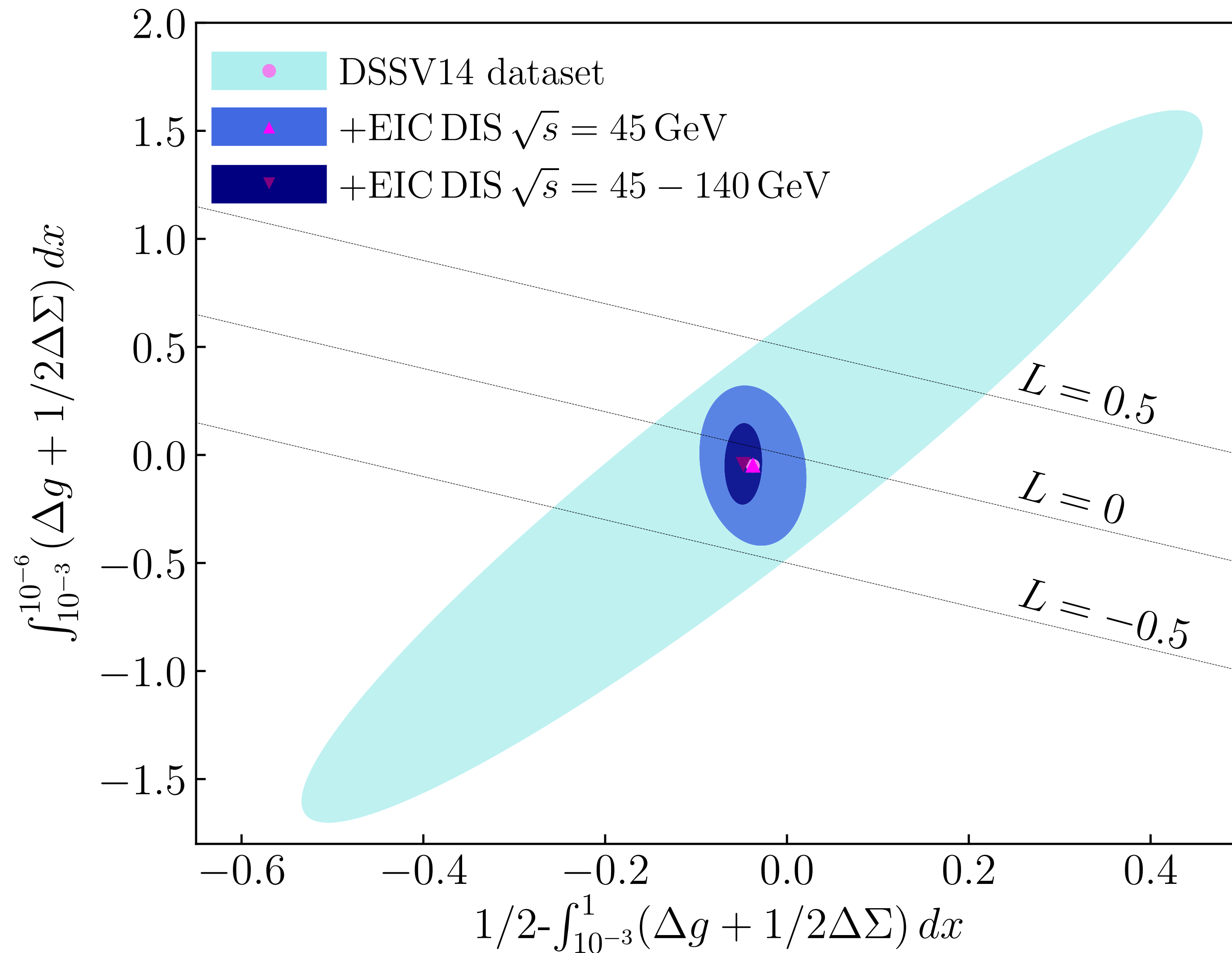
quark spin

gluon spin

quark and gluon orbital angular momentum



Aschenauer et al., [arXiv:1708.01527](https://arxiv.org/abs/1708.01527) and [arXiv:1509.06489](https://arxiv.org/abs/1509.06489)



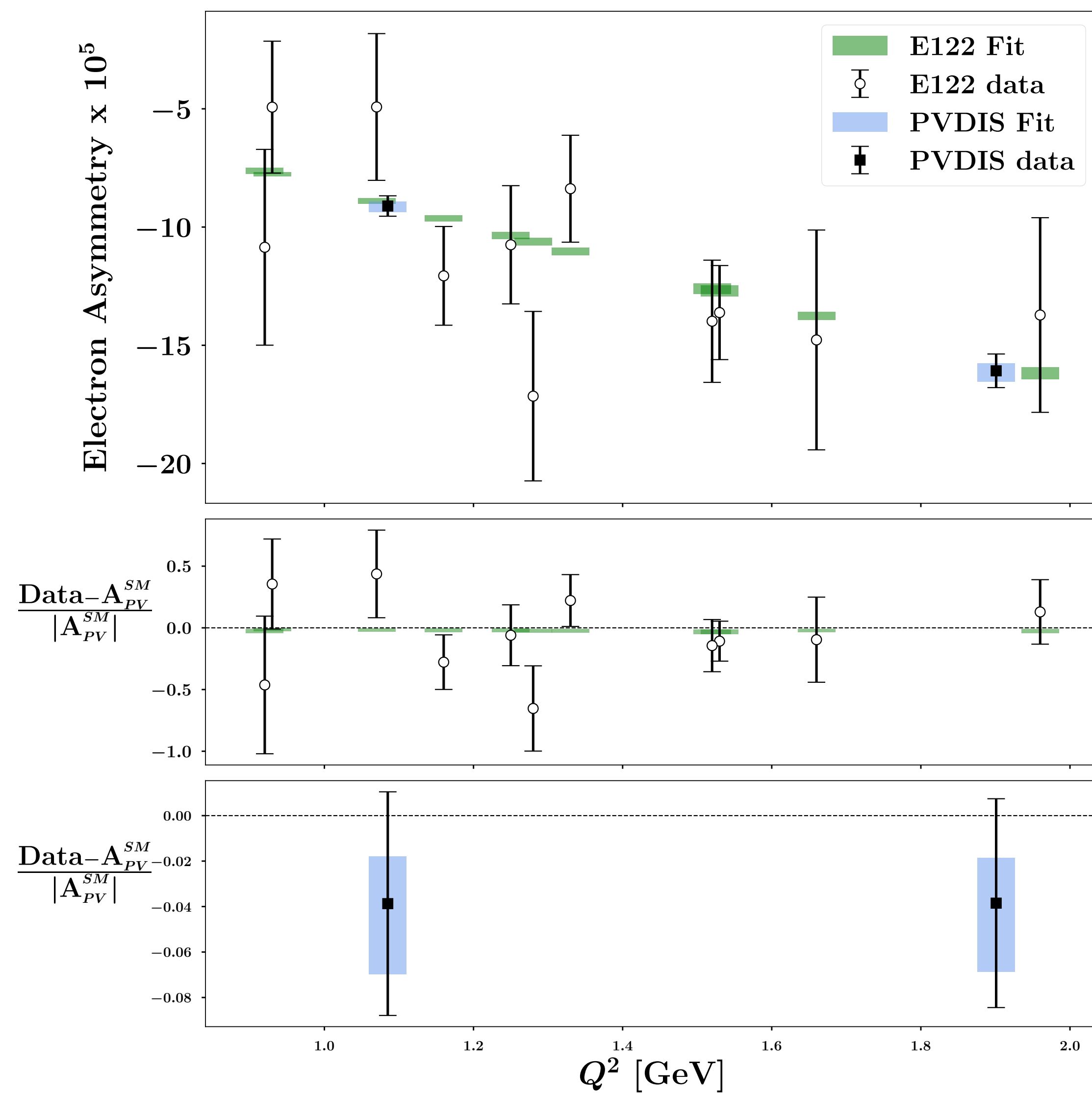
quark spin

gluon spin

$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta g + L$$

quark and gluon orbital angular momentum

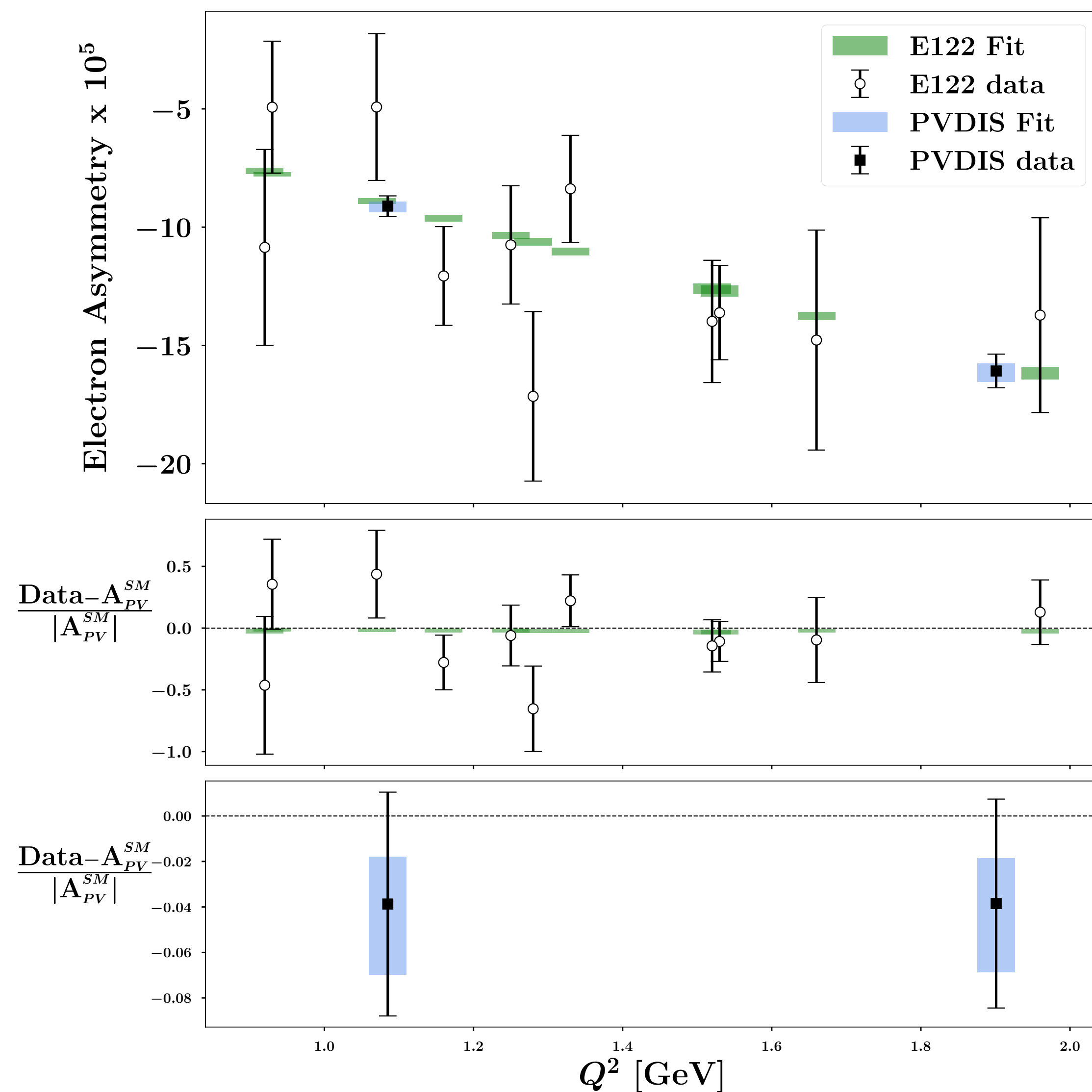
CAN THERE STILL BE SURPRISES?



[Bacchetta, Cerutti, Manna, Radici, Zheng, arxiv:2306.04704](#)

A_{PV} asymmetry (with polarized leptons)

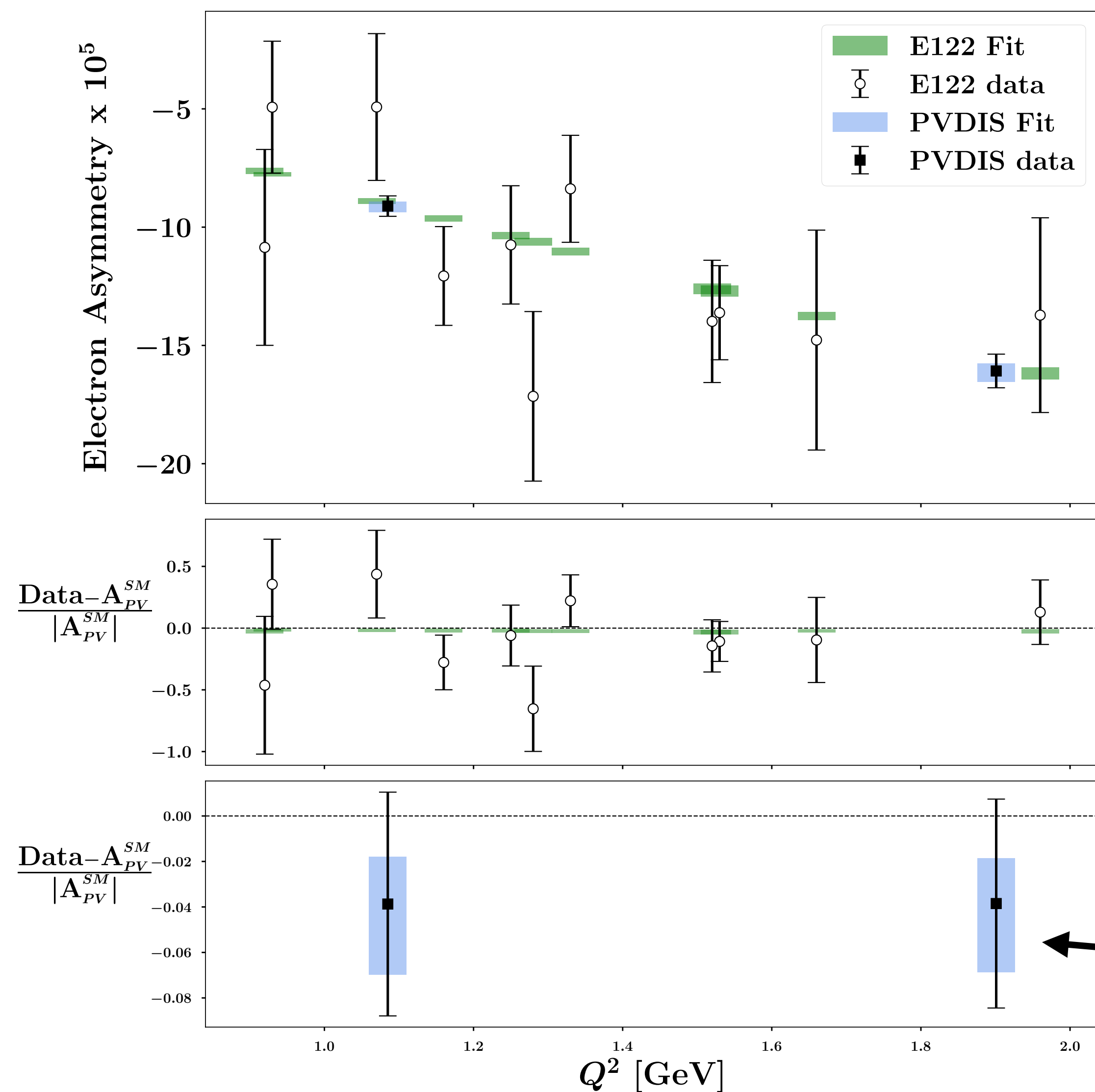
[Bacchetta, Cerutti, Manna, Radici, Zheng, arxiv:2306.04704](#)



A_{PV} asymmetry (with polarized leptons)

Standard Model prediction

[Bacchetta, Cerutti, Manna, Radici, Zheng, arxiv:2306.04704](#)

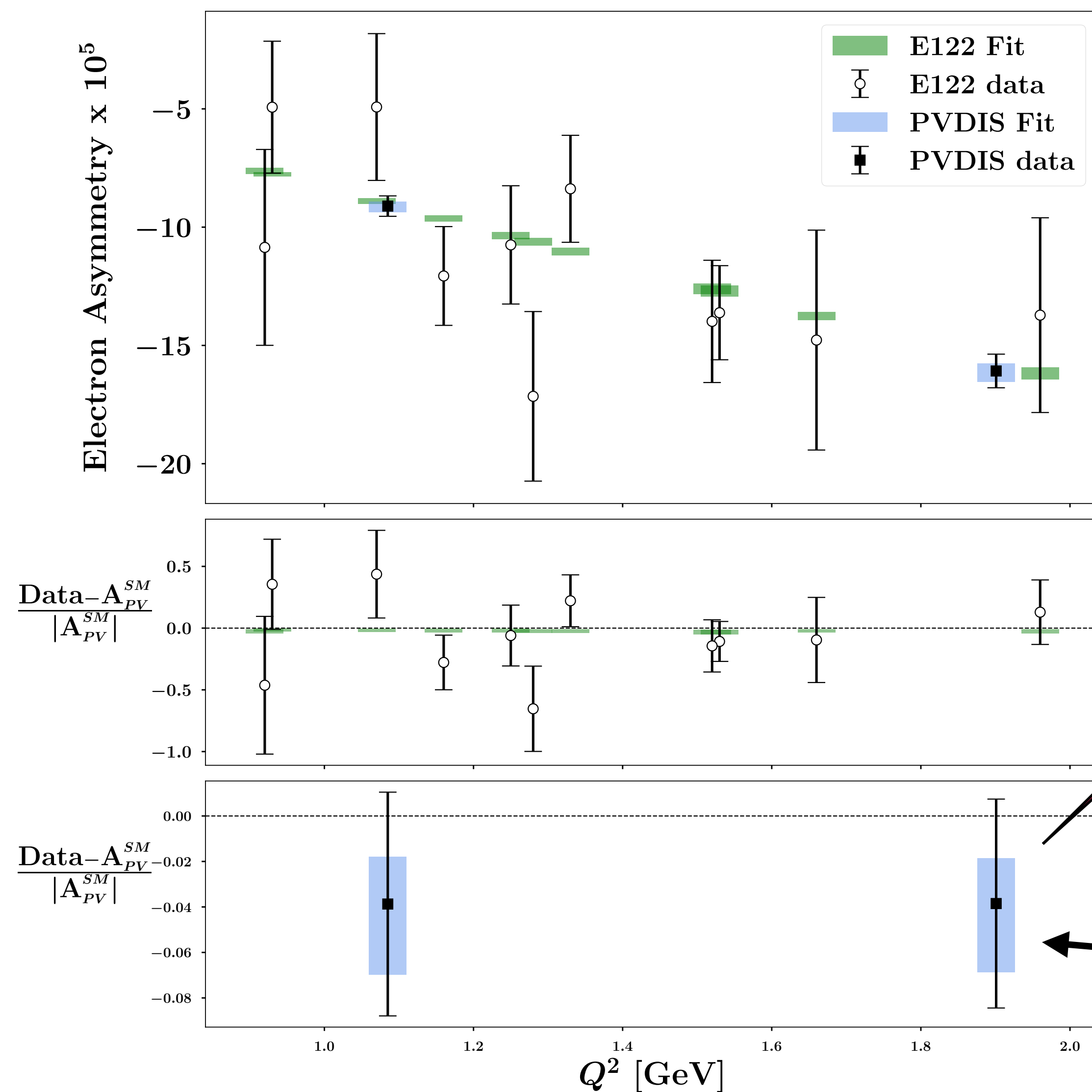


A_{PV} asymmetry (with polarized leptons)

Standard Model prediction

Fit with the inclusion of a strong parity violating parton distribution function

[Bacchetta, Cerutti, Manna, Radici, Zheng, arxiv:2306.04704](#)



A_{PV} asymmetry (with polarized leptons)

Precise DIS data may expose signals of strong parity violation

Standard Model prediction

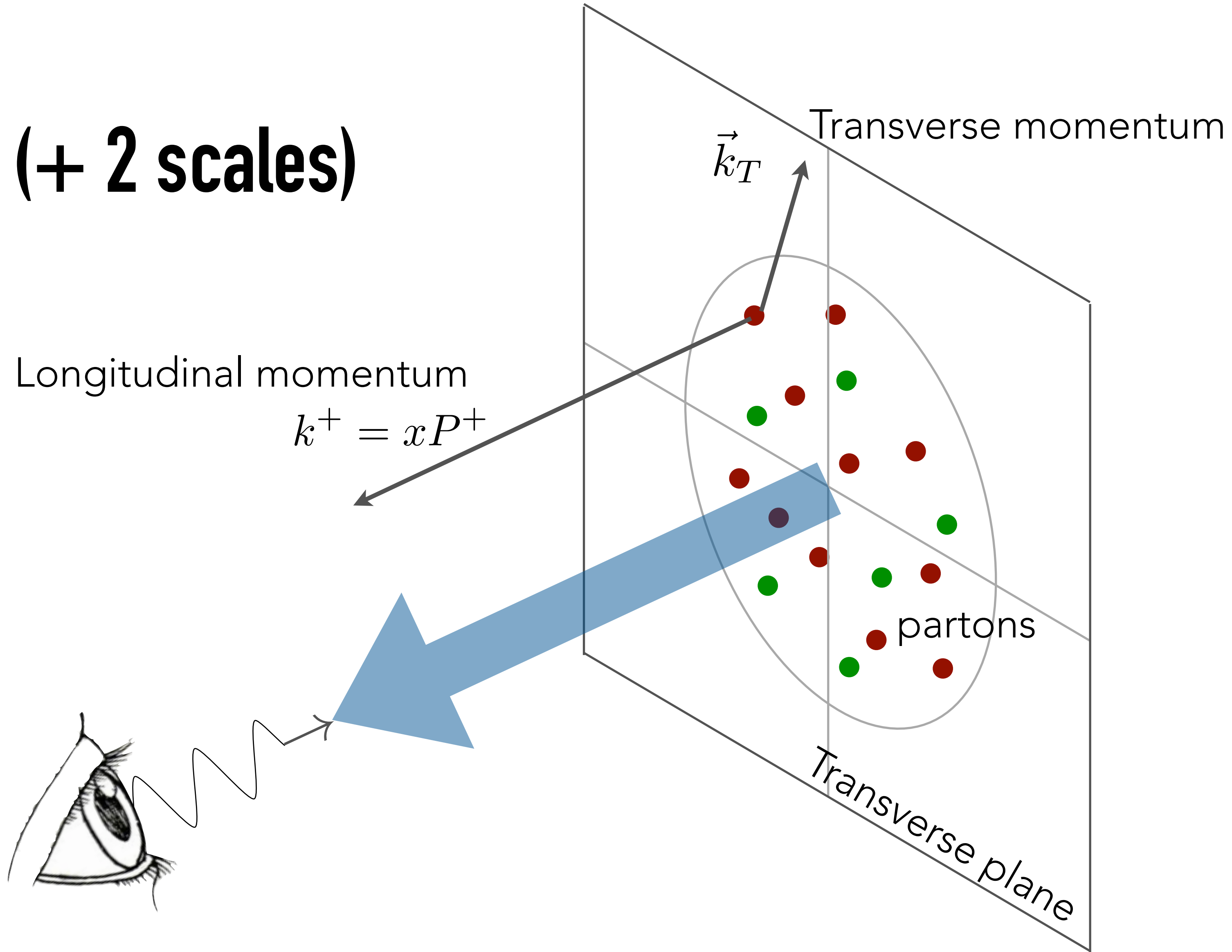
Fit with the inclusion of a strong parity violating parton distribution function

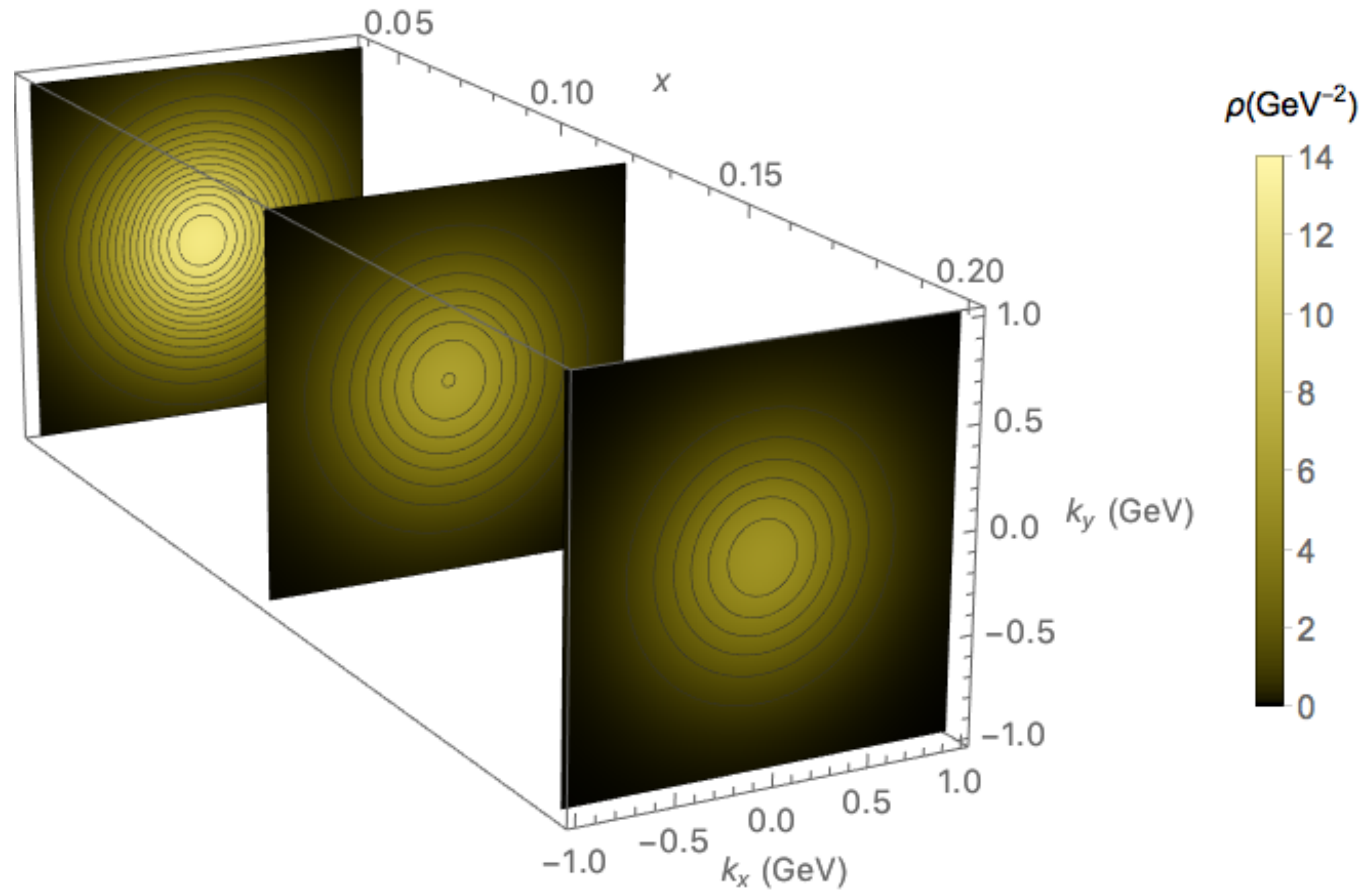
3-DIMENSIONAL MAPS

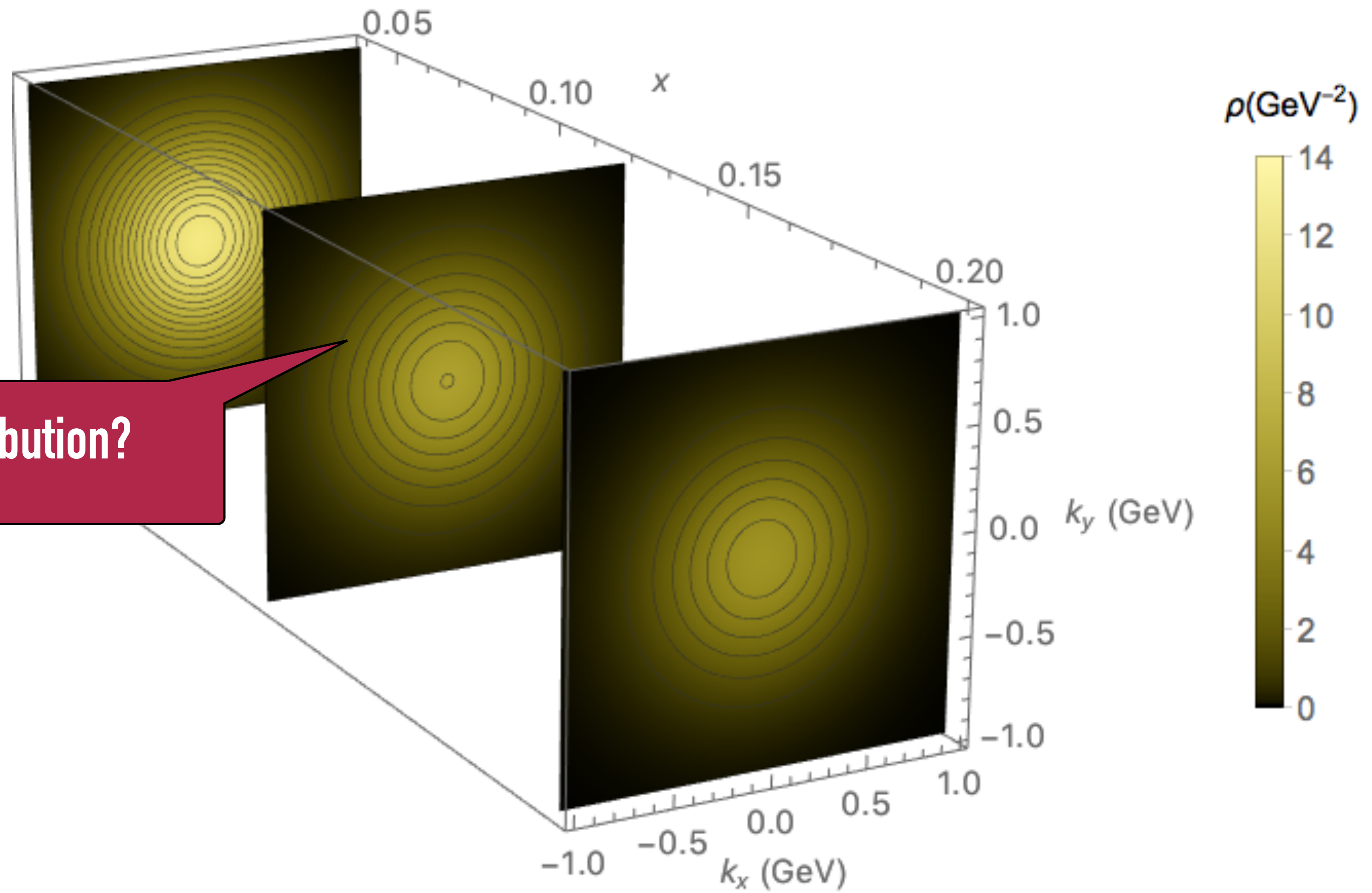
Transverse-Momentum Distributions

$$f(x, \vec{k}_T)$$

3 dimensional (+ 2 scales)



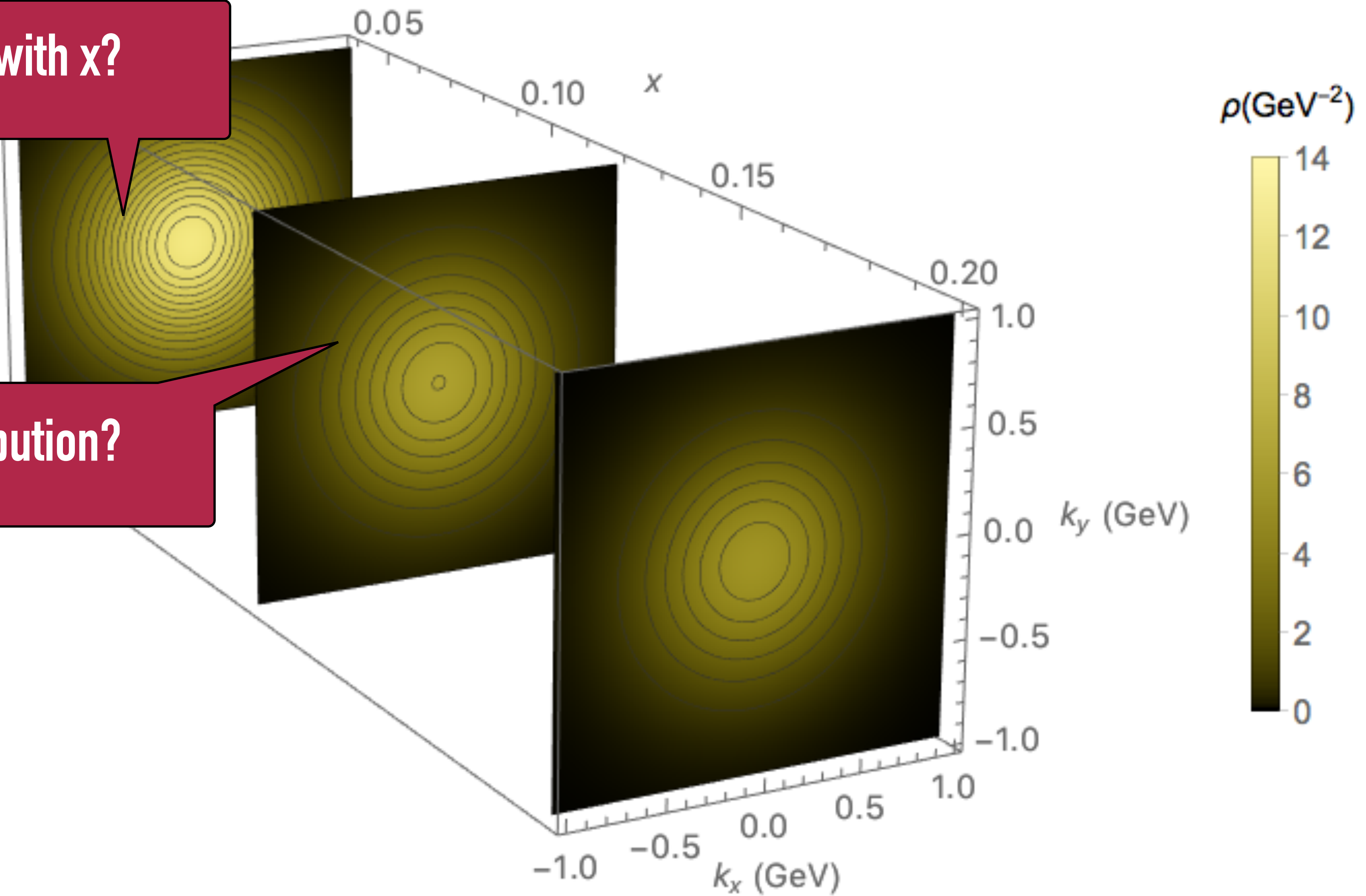




How "wide" is the distribution?

How does it change with x ?

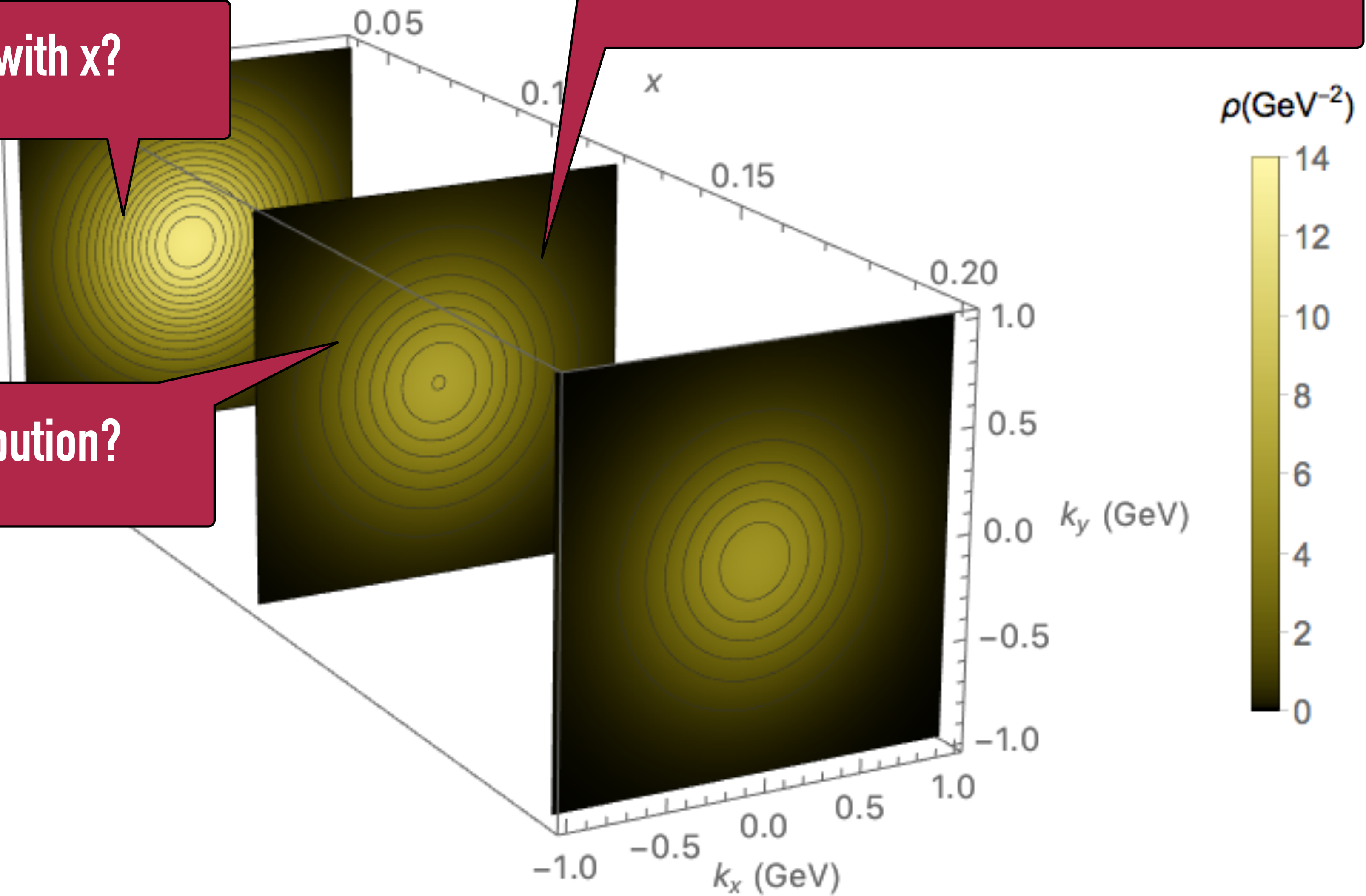
How "wide" is the distribution?



How does it change with x ?

Is there a difference between flavors?

How "wide" is the distribution?

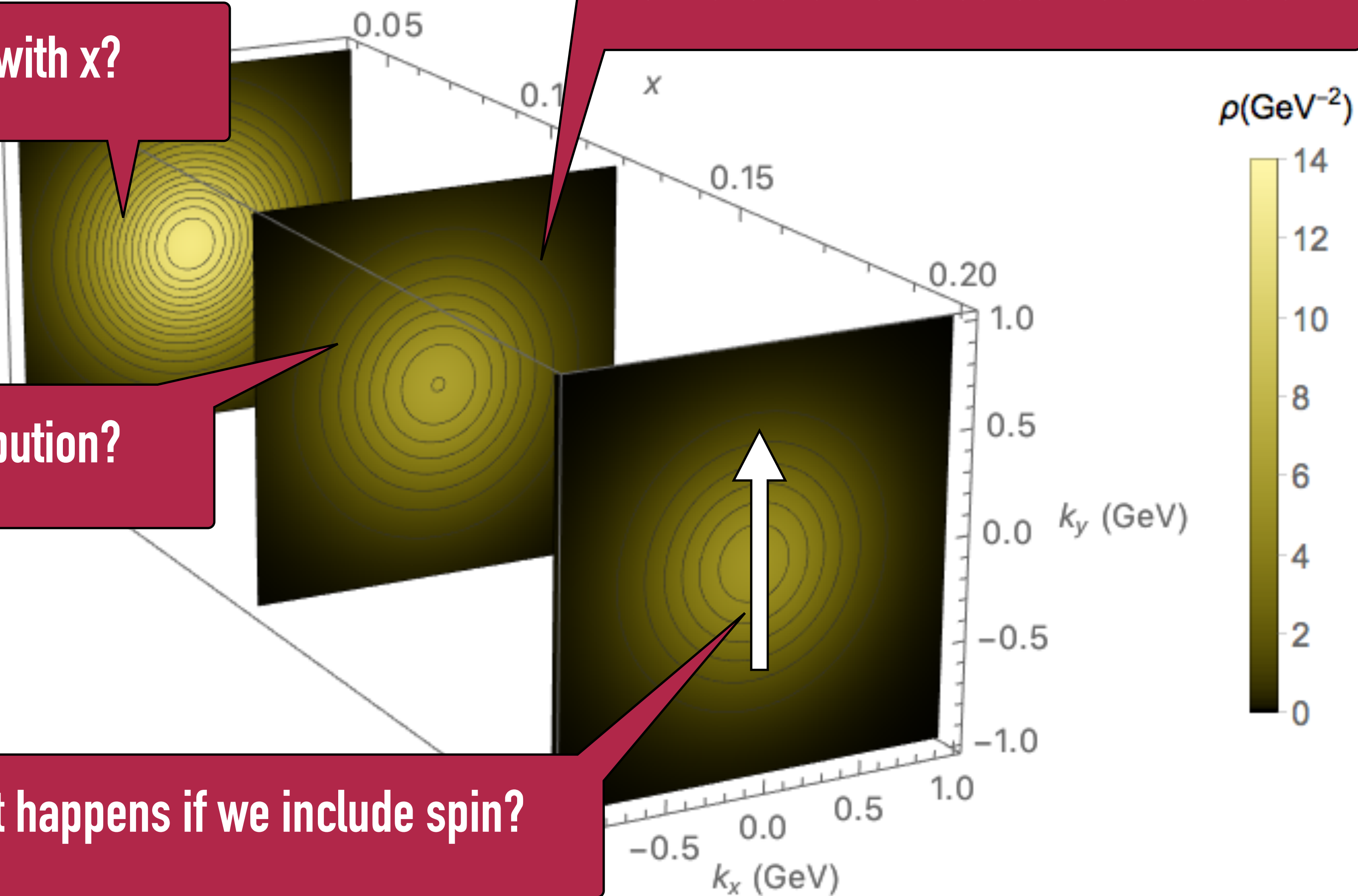


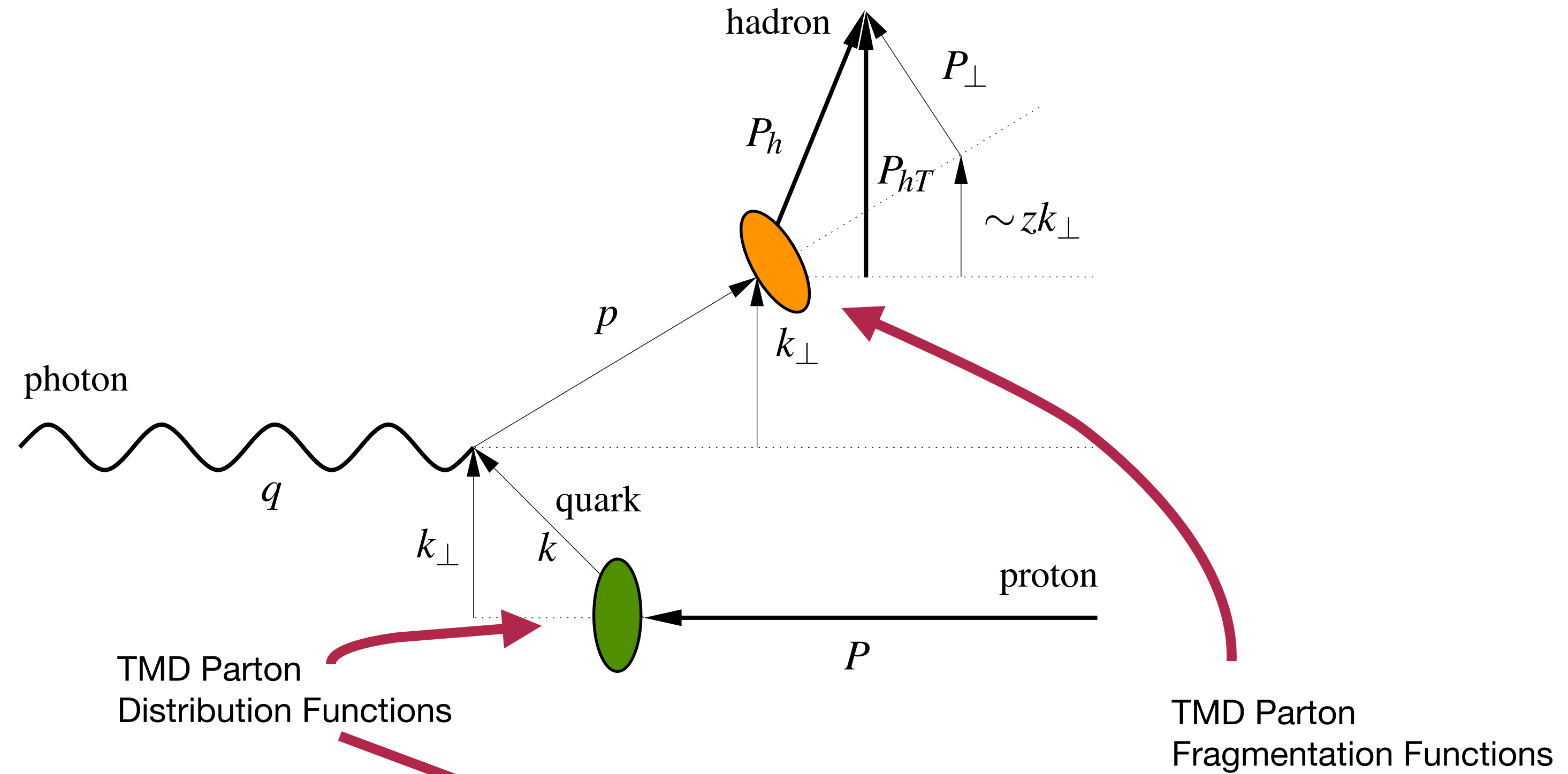
How does it change with x ?

Is there a difference between flavors?

How "wide" is the distribution?

What happens if we include spin?





$$\begin{aligned}
 & F_{UU,T}(x, z, \mathbf{P}_{hT}^2, Q^2) \\
 &= x \sum_q \mathcal{H}_{UU,T}^q(Q^2, \mu^2) \int d^2 \mathbf{k}_{\perp} d^2 \mathbf{P}_{\perp} f_1^a(x, \mathbf{k}_{\perp}^2; \mu^2) D_1^{a \rightarrow h}(z, \mathbf{P}_{\perp}^2; \mu^2) \delta(z \mathbf{k}_{\perp} - \mathbf{P}_{hT} + \mathbf{P}_{\perp}) \\
 &= x \sum_a \mathcal{H}_{UU,T}^a(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{P}_{h\perp}|) \hat{f}_1^a(x, z^2 b_{\perp}^2; \mu^2) \hat{D}_1^{a \rightarrow h}(z, b_{\perp}^2; \mu^2)
 \end{aligned}$$

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu b_*) e^{\int_{\mu b_*}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu b_*} \right)^{K_{\text{resum}} + g_K}$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K}$$

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

perturbative Sudakov
form factor

$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K}$$

collinear PDF

Collins-Soper kernel
(perturbative and
nonperturbative)

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$

matching coefficients
(perturbative)

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

perturbative Sudakov form factor

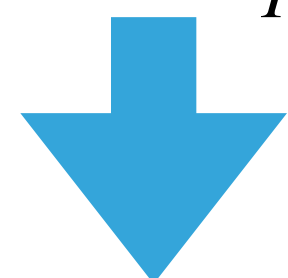
$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K}$$

collinear PDF

Collins-Soper kernel (perturbative and nonperturbative)

matching coefficients (perturbative)

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$



$$\mu_{b_*} = \frac{2e^{-\gamma_E}}{\bar{b}_*}$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

$$\hat{f}_1^a(x, |\mathbf{b}_T|; \mu, \zeta) = \int d^2\mathbf{k}_\perp e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^a(x, \mathbf{k}_\perp^2; \mu, \zeta)$$

perturbative Sudakov form factor

$$\hat{f}_1^a(x, b_T^2; \mu_f, \zeta_f) = [C \otimes f_1](x, \mu_{b_*}) e^{\int_{\mu_{b_*}}^{\mu_f} \frac{d\mu}{\mu} (\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta_f}}{\mu})} \left(\frac{\sqrt{\zeta_f}}{\mu_{b_*}} \right)^{K_{\text{resum}} + g_K} f_{1NP}(x, b_T^2; \zeta_f, Q_0)$$

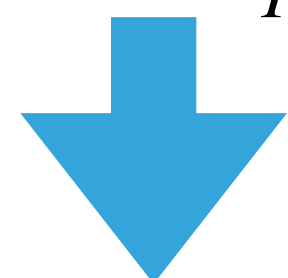
collinear PDF

Collins-Soper kernel (perturbative and nonperturbative)

nonperturbative part of TMD

matching coefficients (perturbative)

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T}$$



$$\mu_{b_*} = \frac{2e^{-\gamma_E}}{\bar{b}_*}$$

[see, e.g., Collins, "Foundations of Perturbative QCD" \(11\)](#)
[TMD collaboration, "TMD Handbook," arXiv:2304.03302](#)

	Accuracy	SIDIS HERMES	SIDIS COMPASS	DY fixed target	DY collider	N of points	χ^2/N_{points}
Pavia 2017 arXiv:1703.10157	NLL	✓	✓	✓	✓	8059	1.55
SV 2019 arXiv:1912.06532	N ³ LL ⁻	✓	✓	✓	✓	1039	1.06
MAP22 arXiv:2206.07598	N ³ LL ⁻	✓	✓	✓	✓	2031	1.06

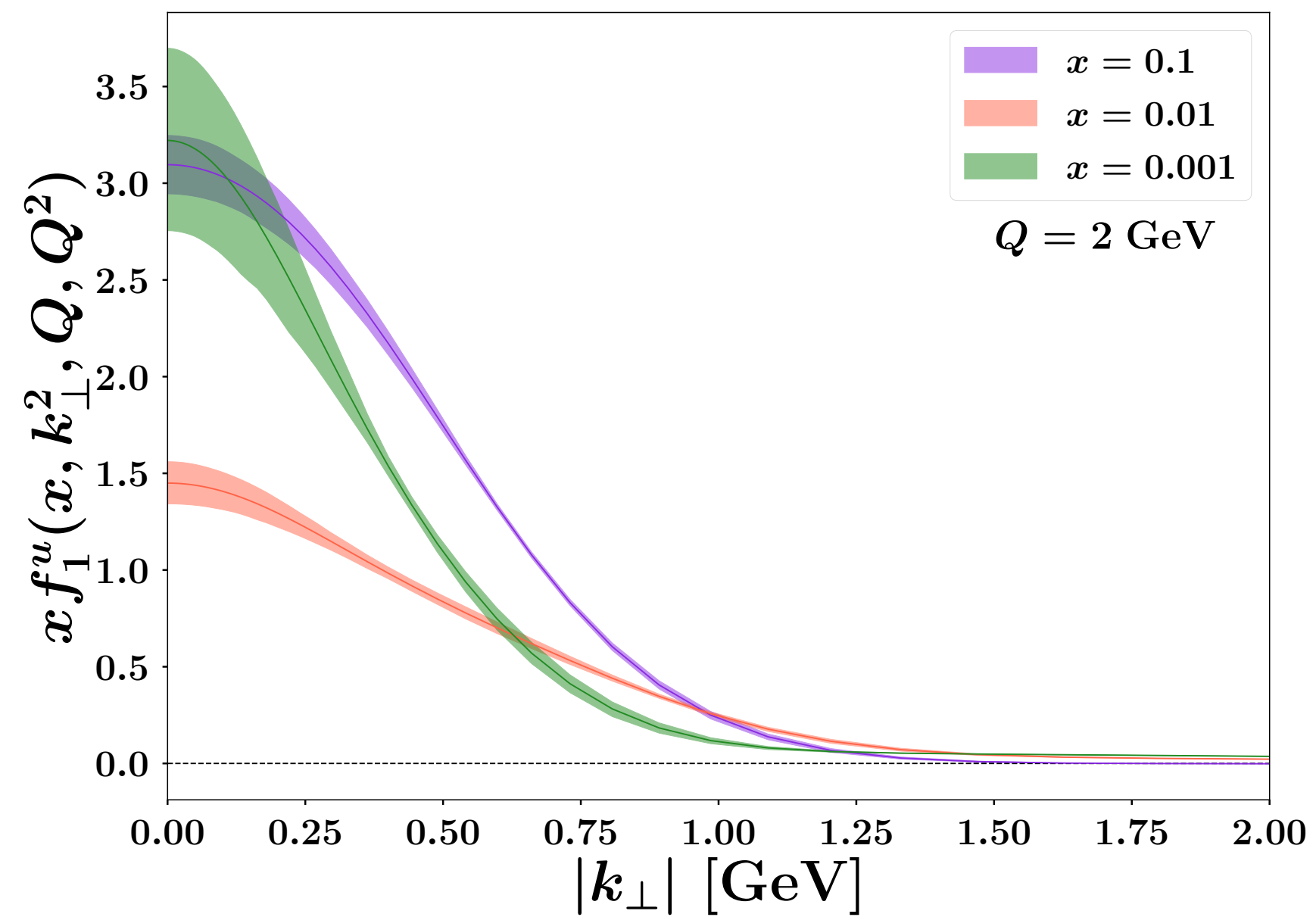


FIG. 13: The TMD PDF of the up quark in a proton at $\mu = \sqrt{\zeta} = Q = 2$ GeV (left panel) and 10 GeV (right panel) as a function of the partonic transverse momentum $|k_\perp|$ for $x = 0.001, 0.01$ and 0.1 . The uncertainty bands represent the 68% CL.

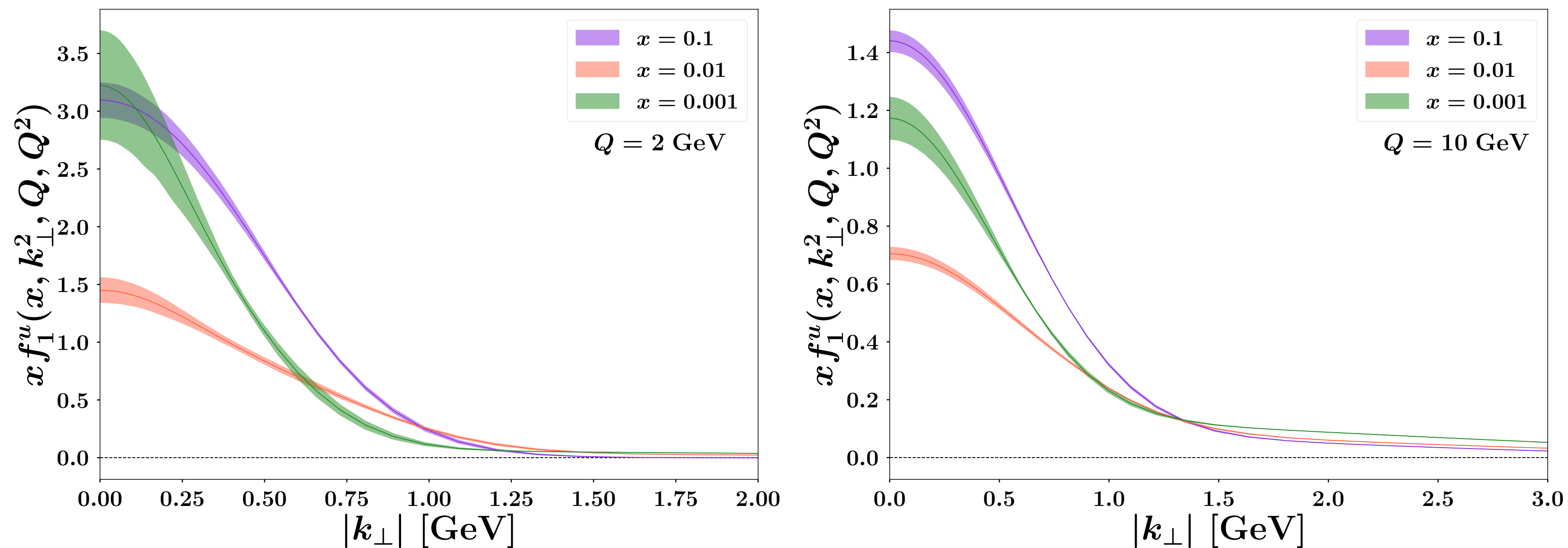
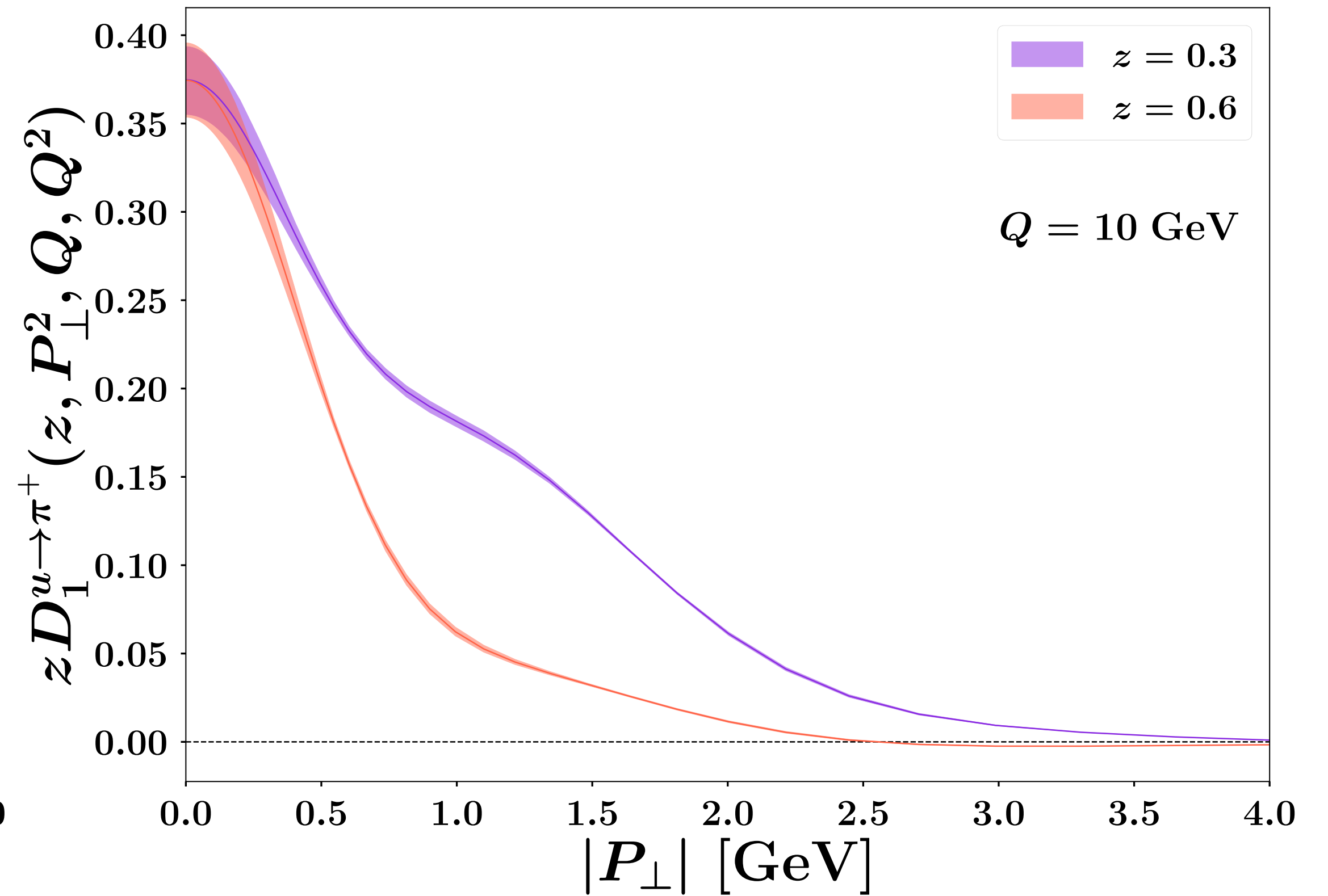
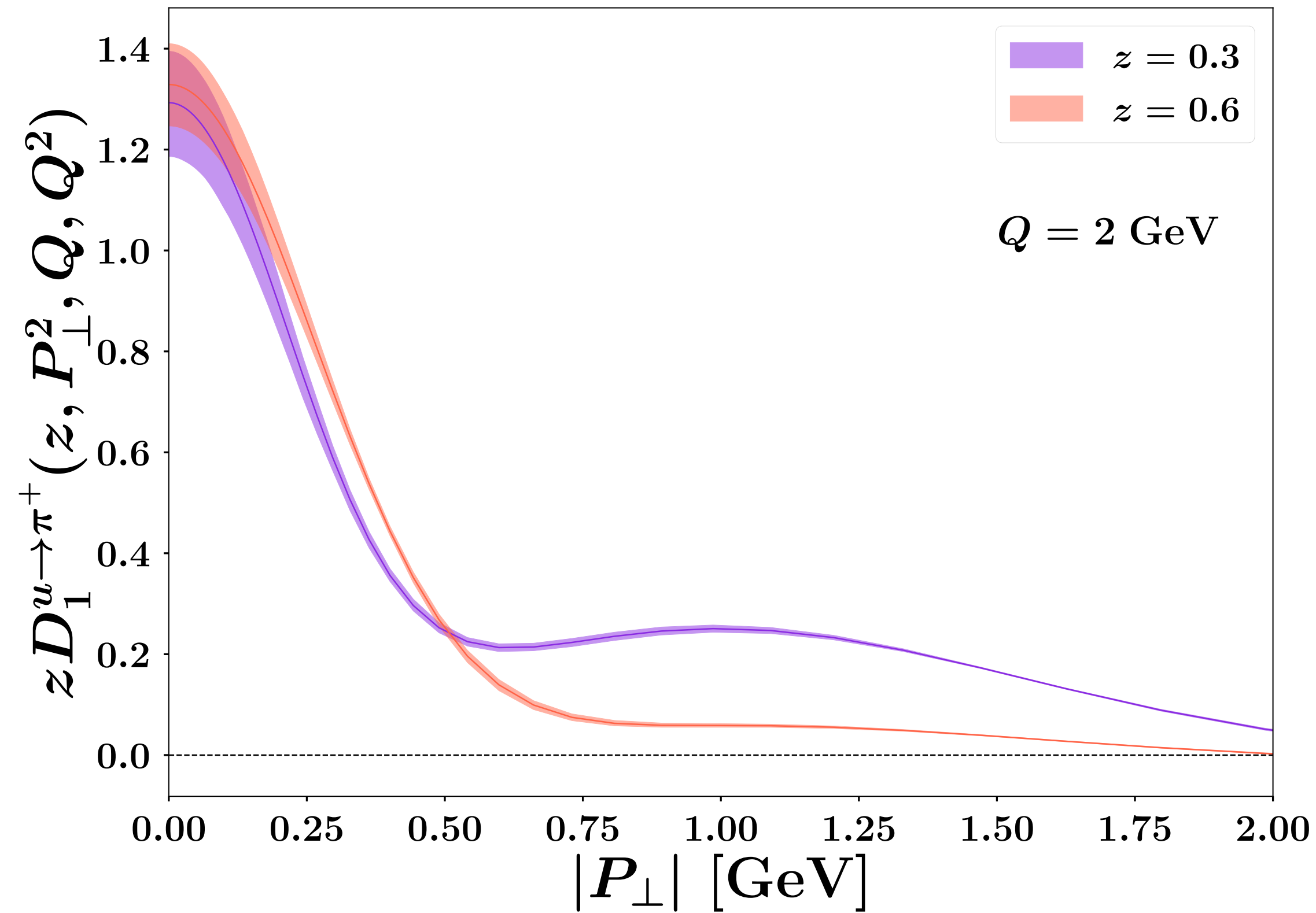
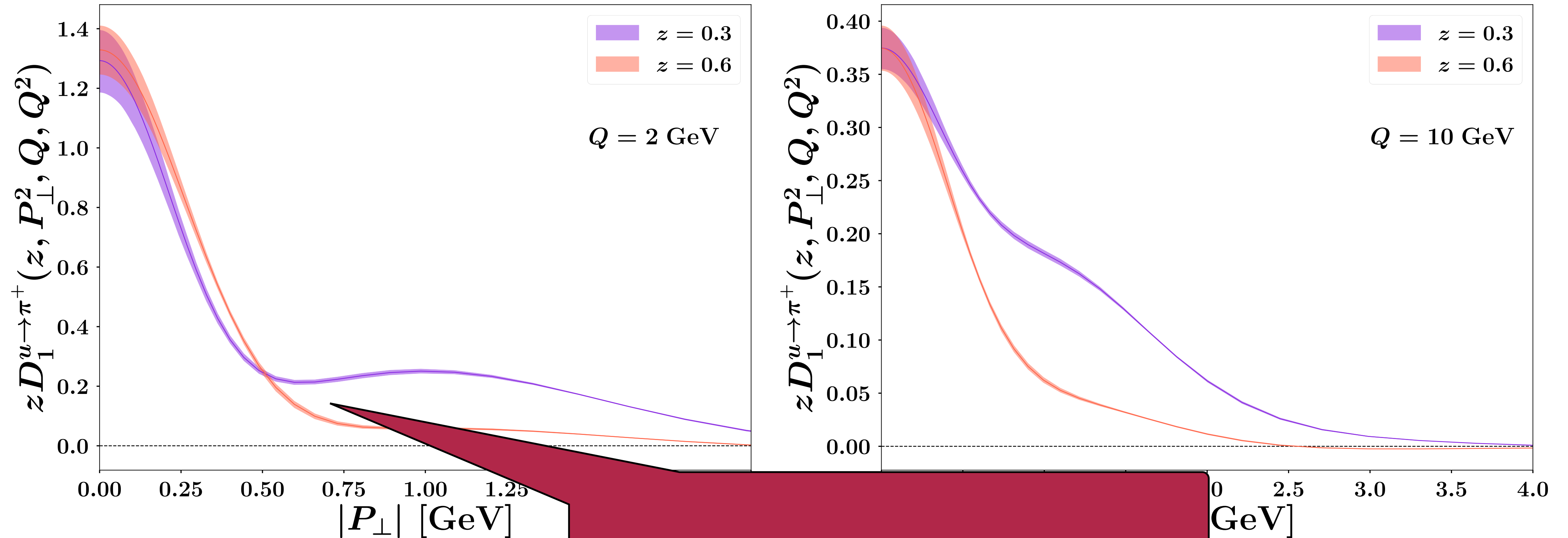
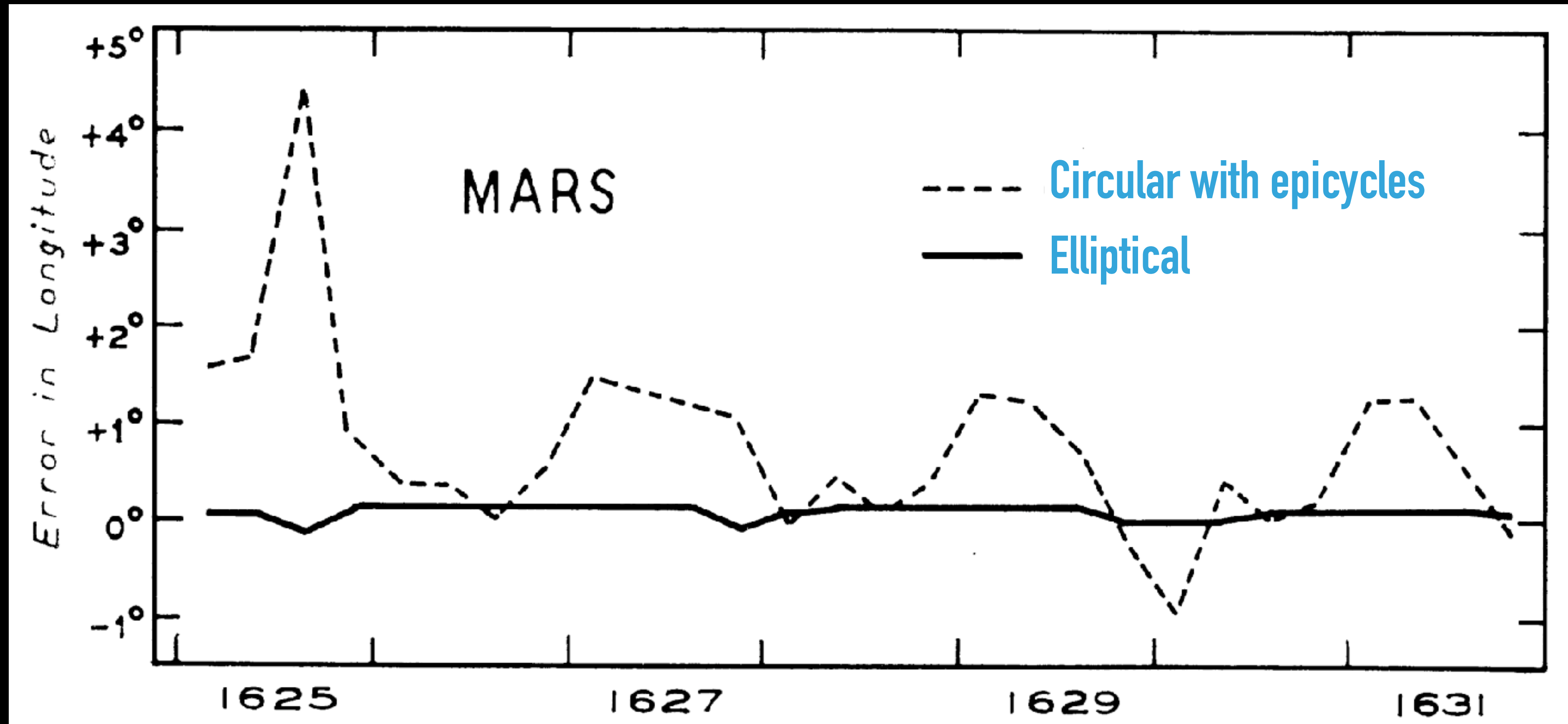


FIG. 13: The TMD PDF of the up quark in a proton at $\mu = \sqrt{\zeta} = Q = 2 \text{ GeV}$ (left panel) and 10 GeV (right panel) as a function of the partonic transverse momentum $|k_\perp|$ for $x = 0.001, 0.01$ and 0.1 . The uncertainty bands represent the 68% CL.

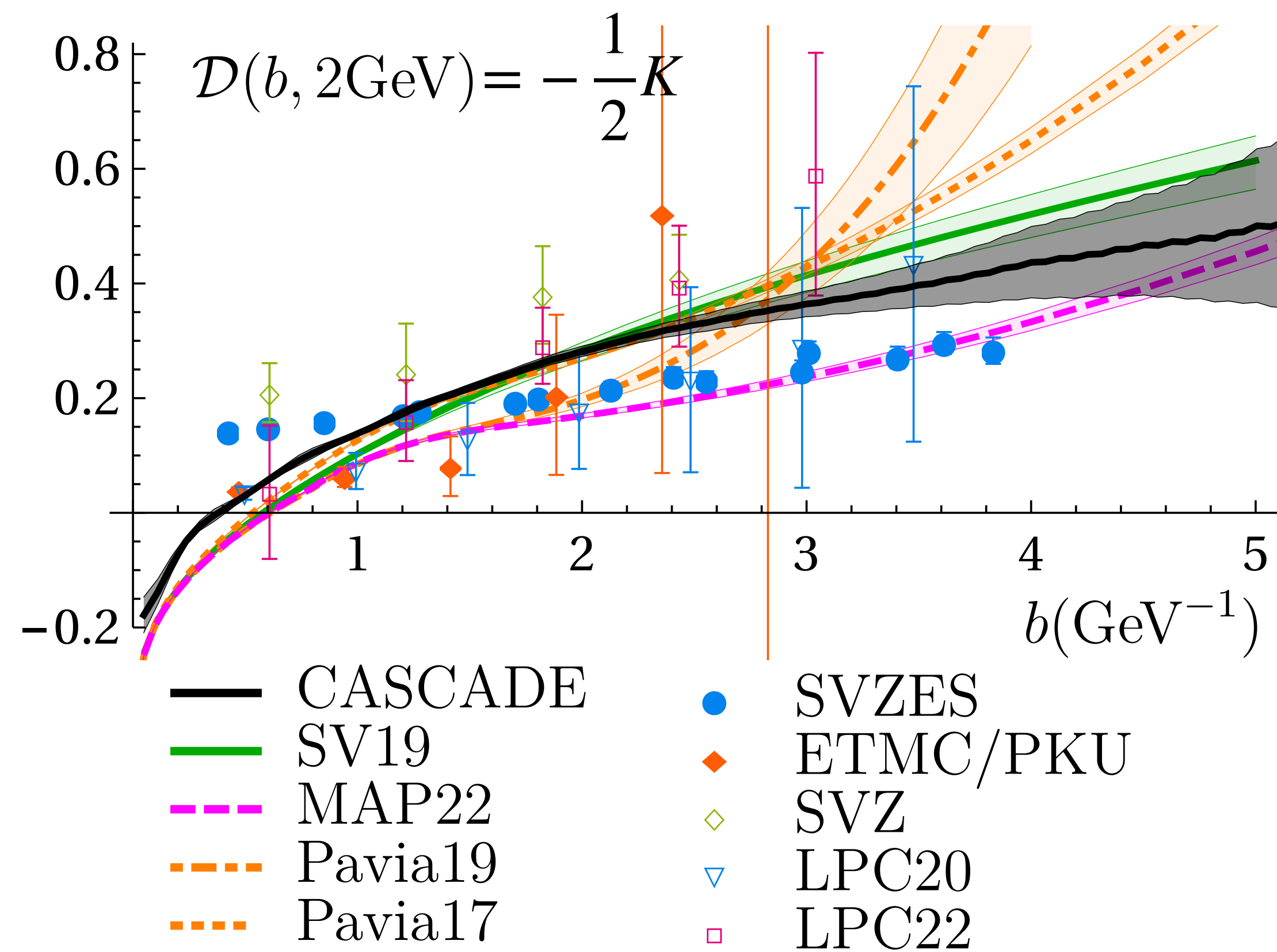




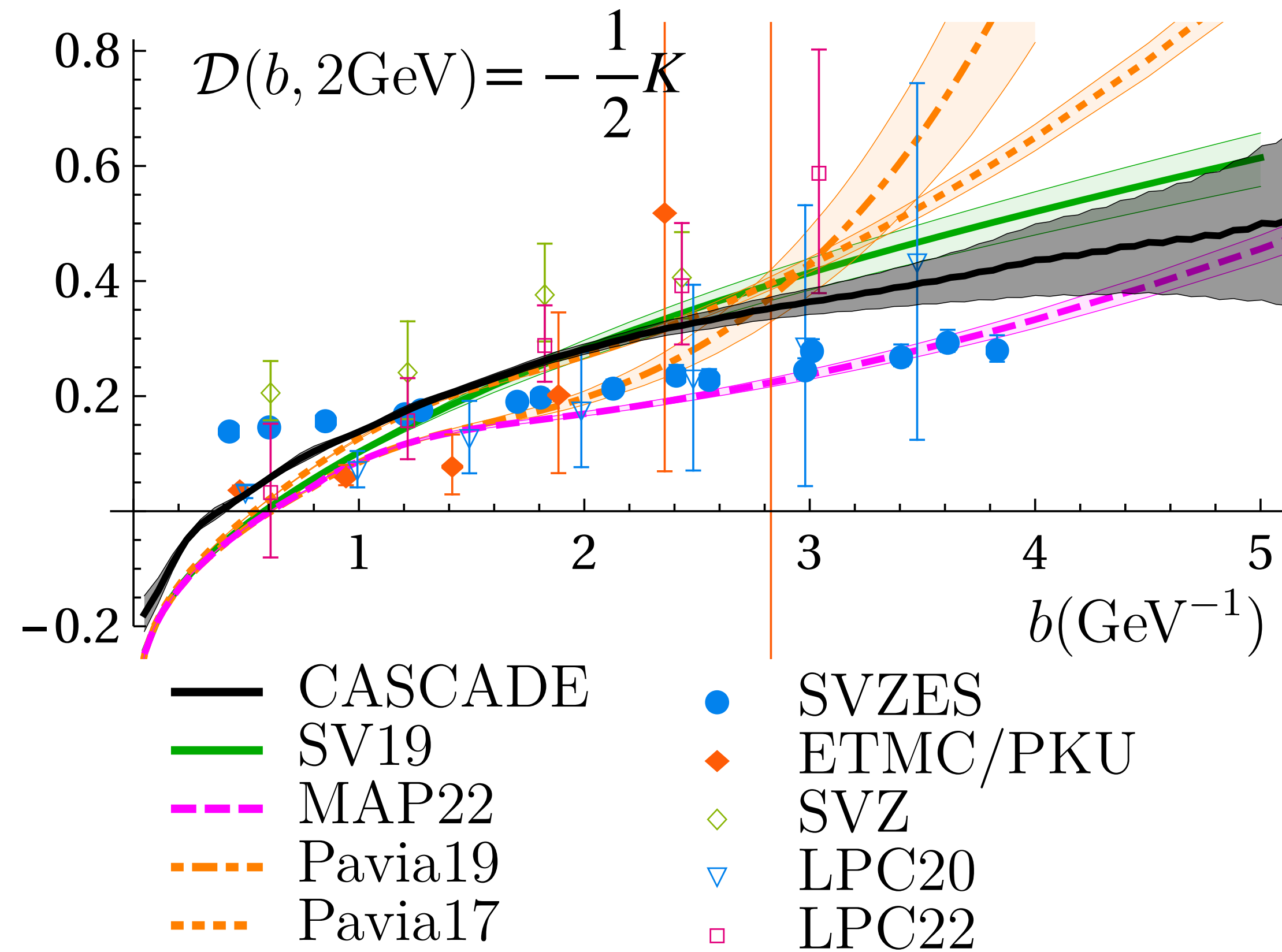
Different from a simple Gaussian
(relation with talk of P. Skands)



[Bermudez Martinez, Vladimirov, arXiv:2206.01105](#)

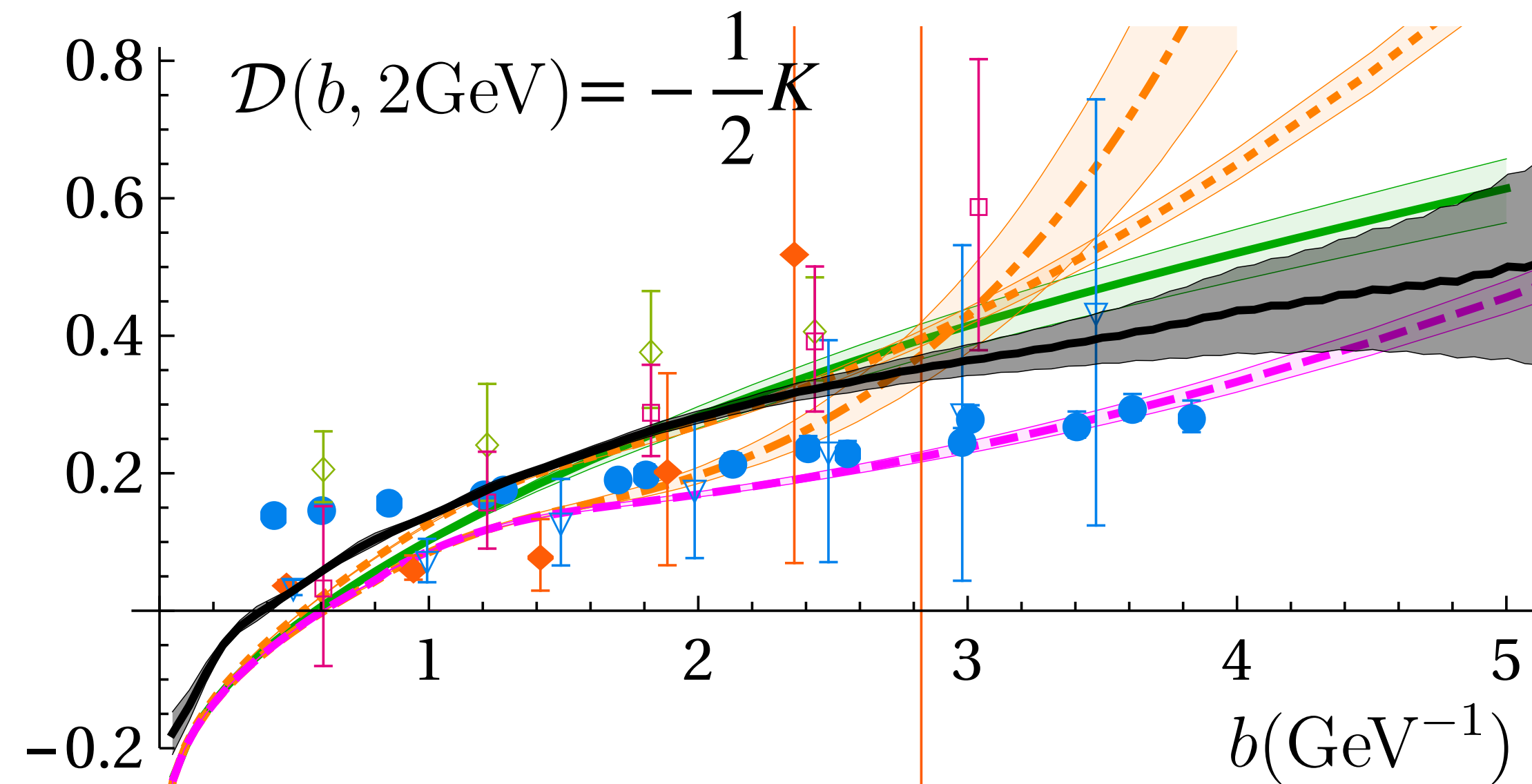


Bermudez Martinez, Vladimirov, arXiv:2206.01105



TMD phenomenology

Bermudez Martinez, Vladimirov, arXiv:2206.01105

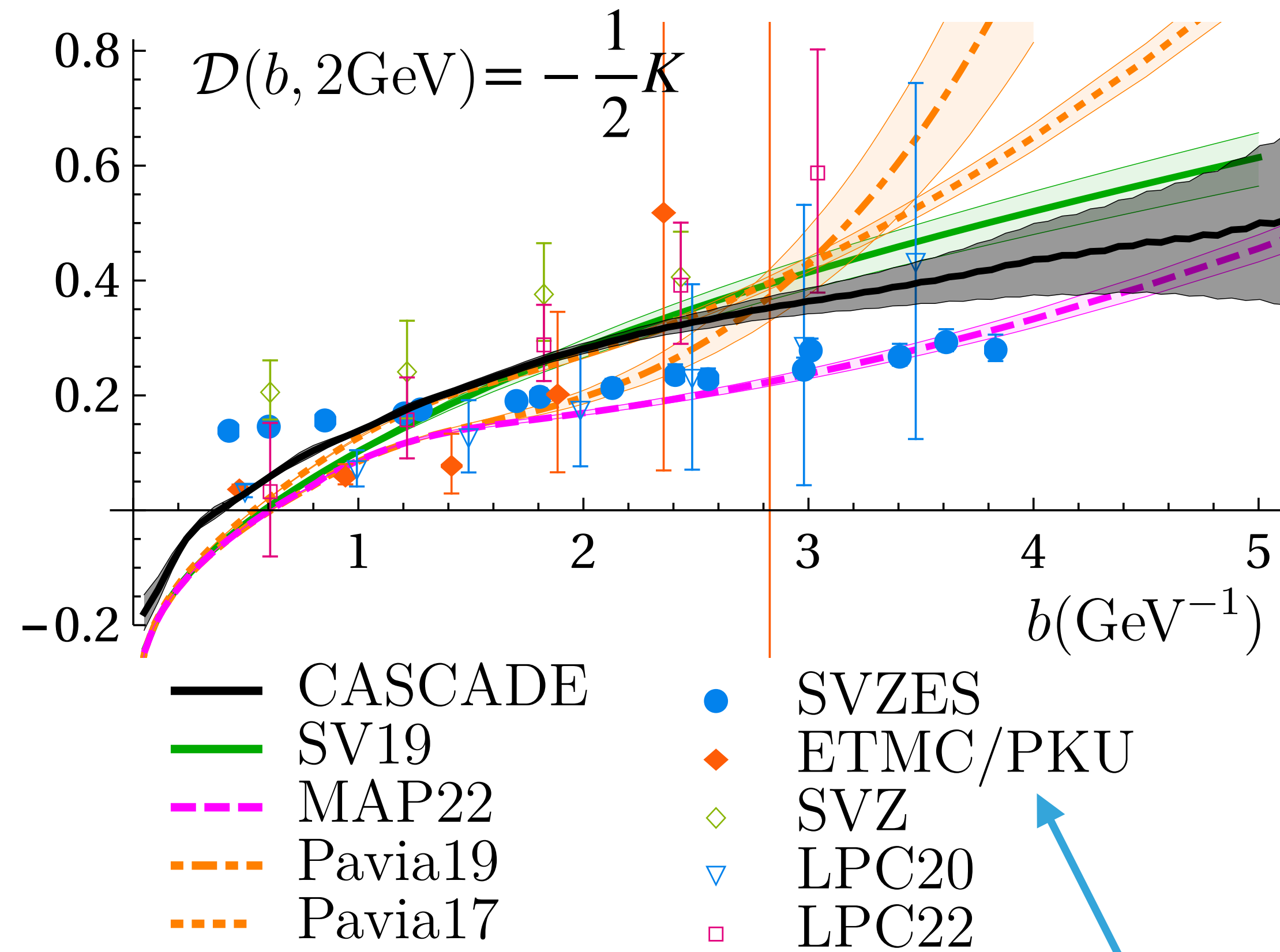


- CASCADE
- SV19
- - - MAP22
- - - Pavia19
- - - Pavia17
- SVZES
- ◆ ETMC/PKU
- ◇ SVZ
- ▽ LPC20
- LPC22

TMD phenomenology

Lattice QCD

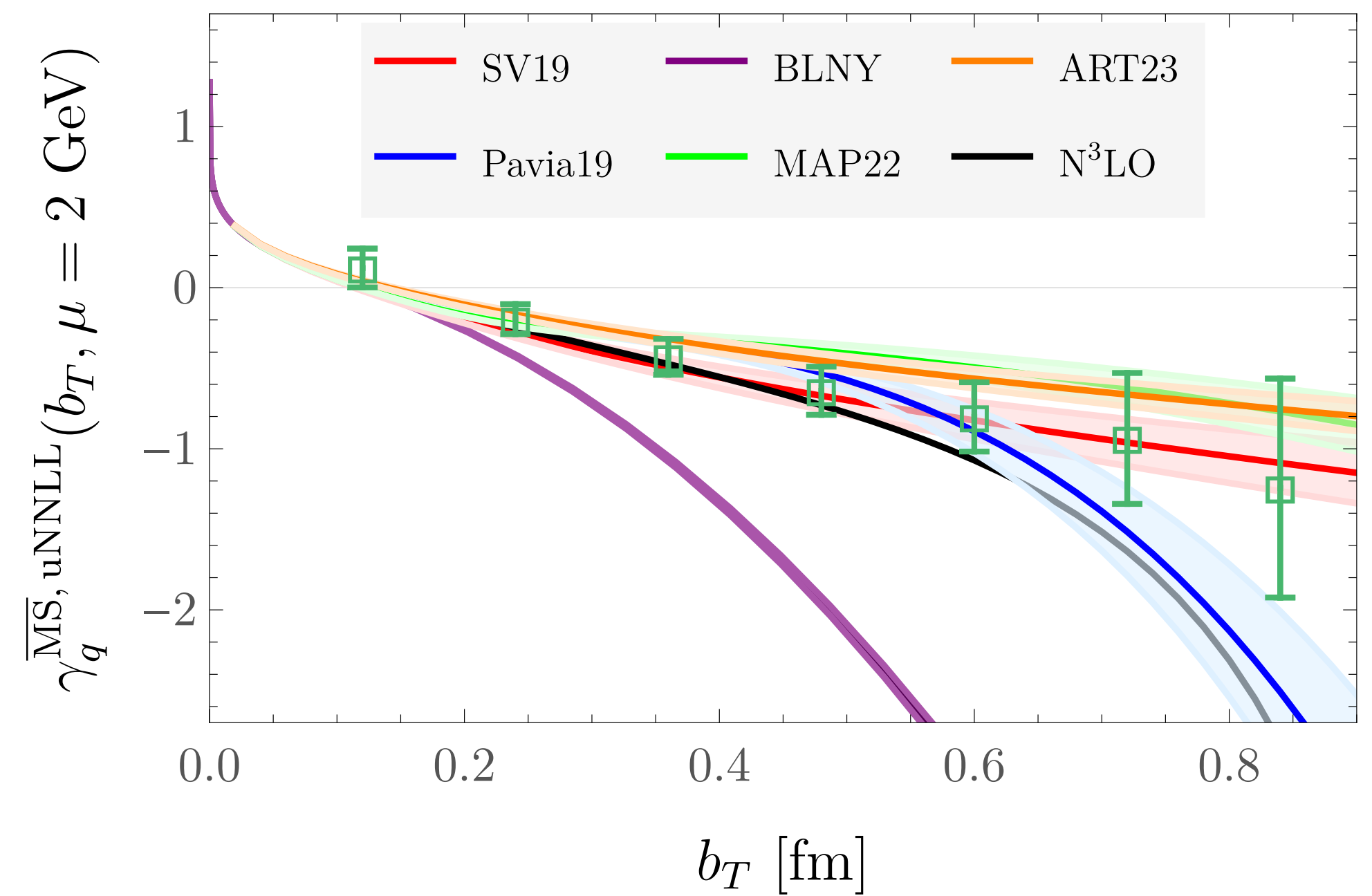
Bermudez Martinez, Vladimirov, arXiv:2206.01105



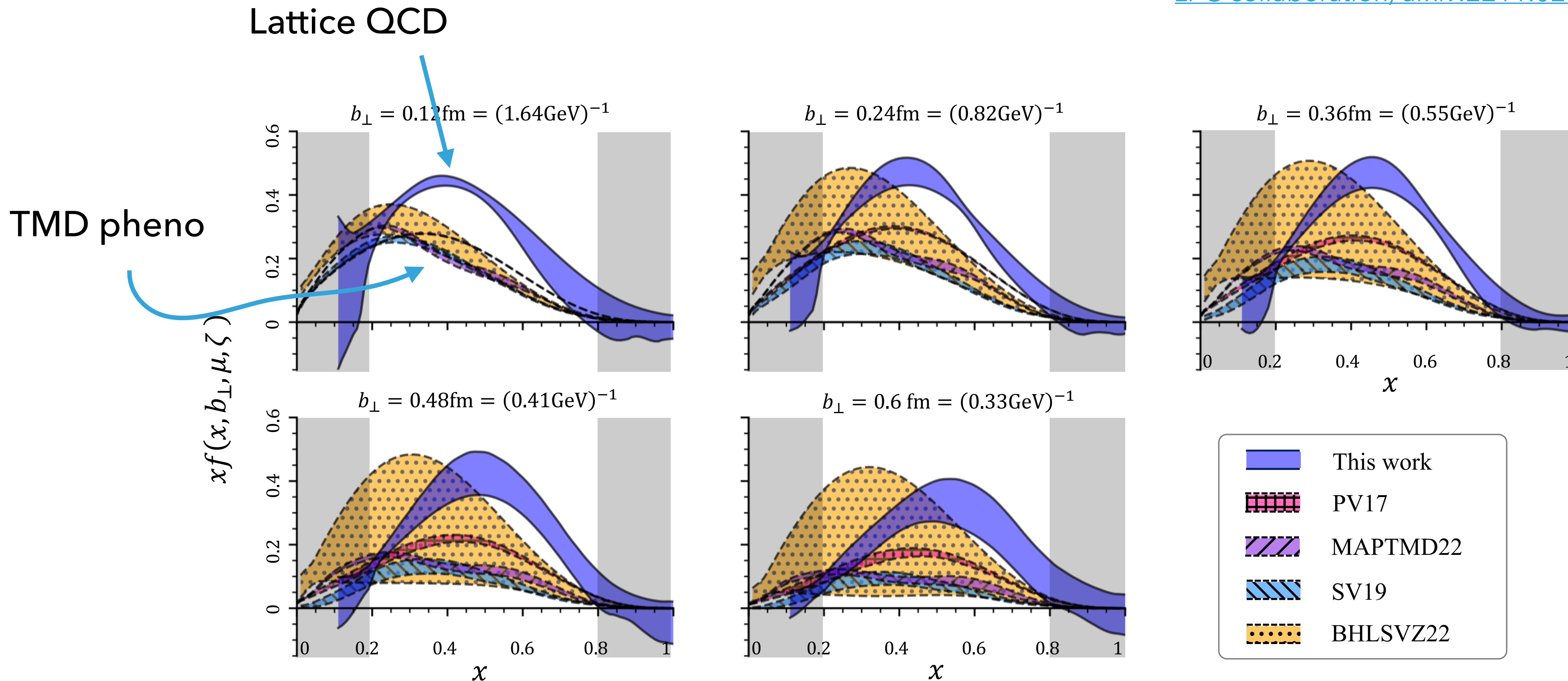
TMD phenomenology

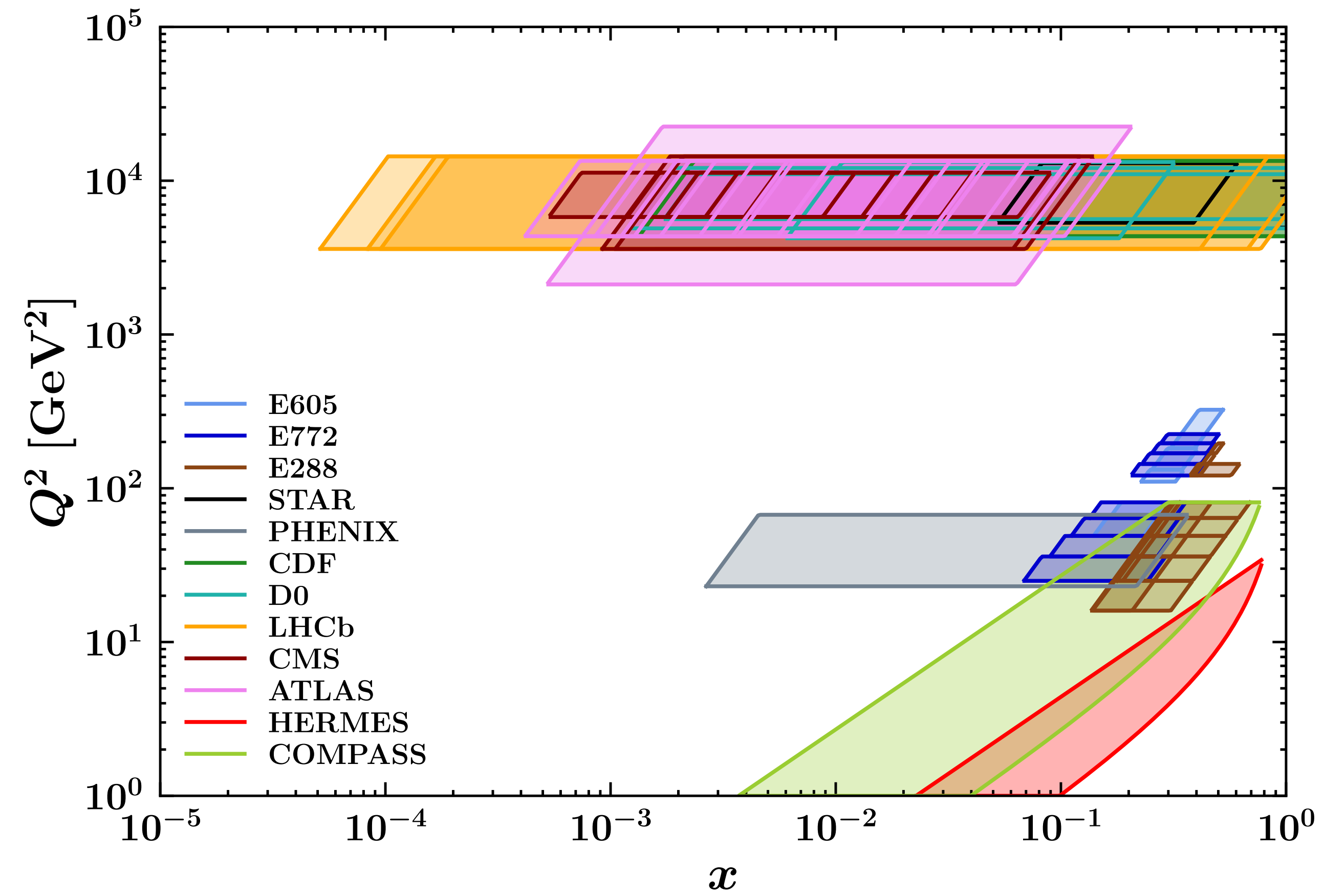
Lattice QCD

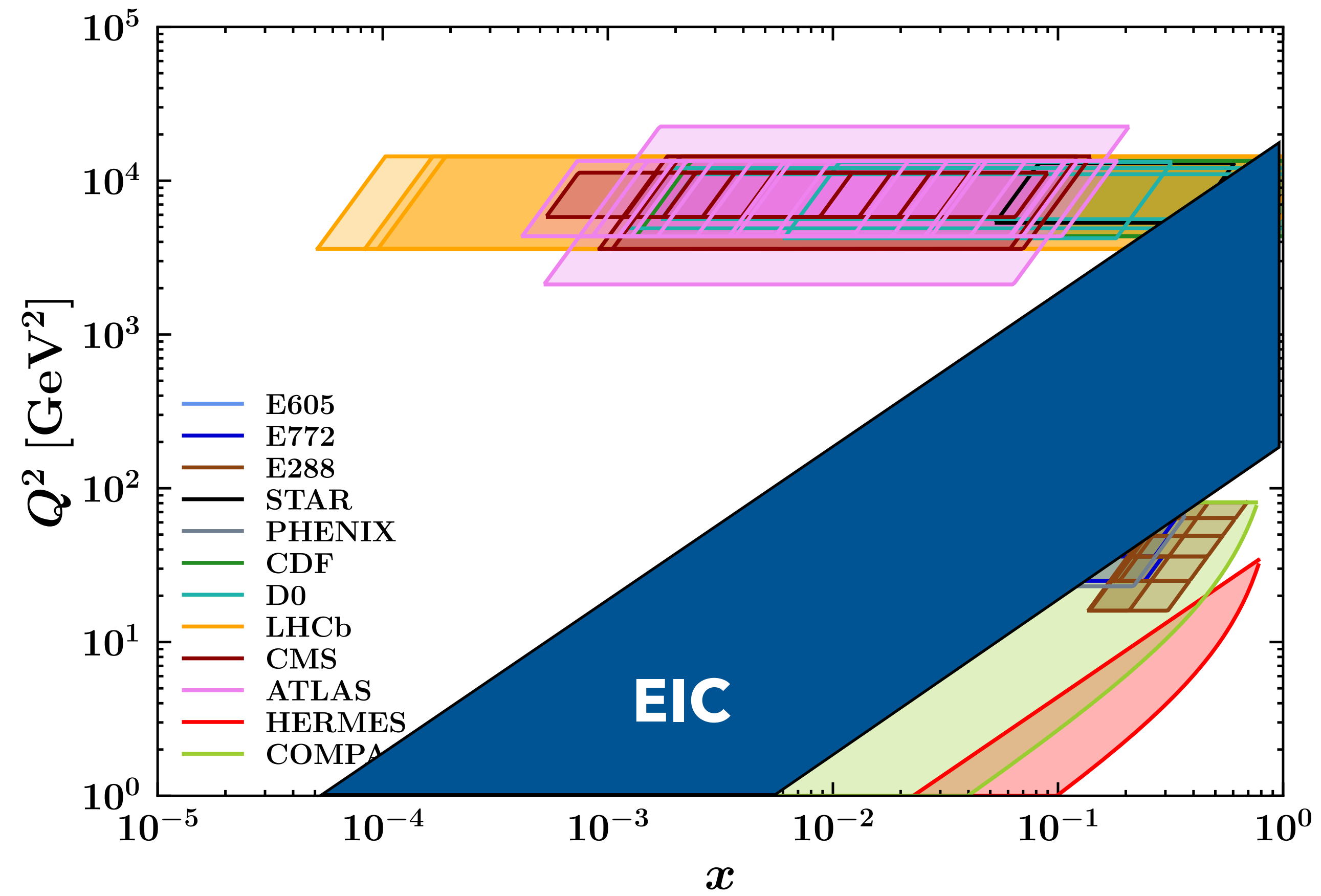
Avkhadiev, Shanahan, Wagman, Zhao, arXiv:2307.12359

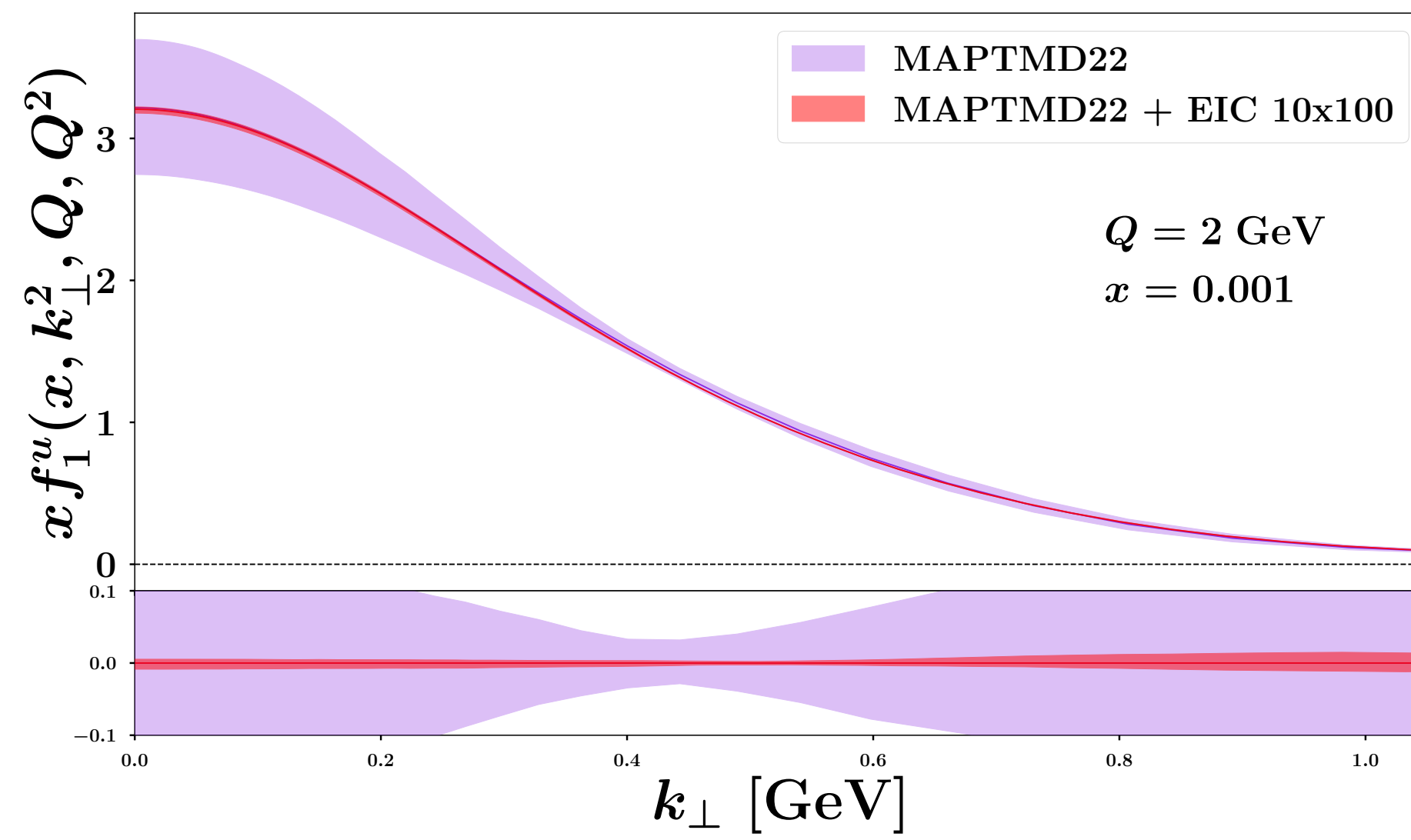
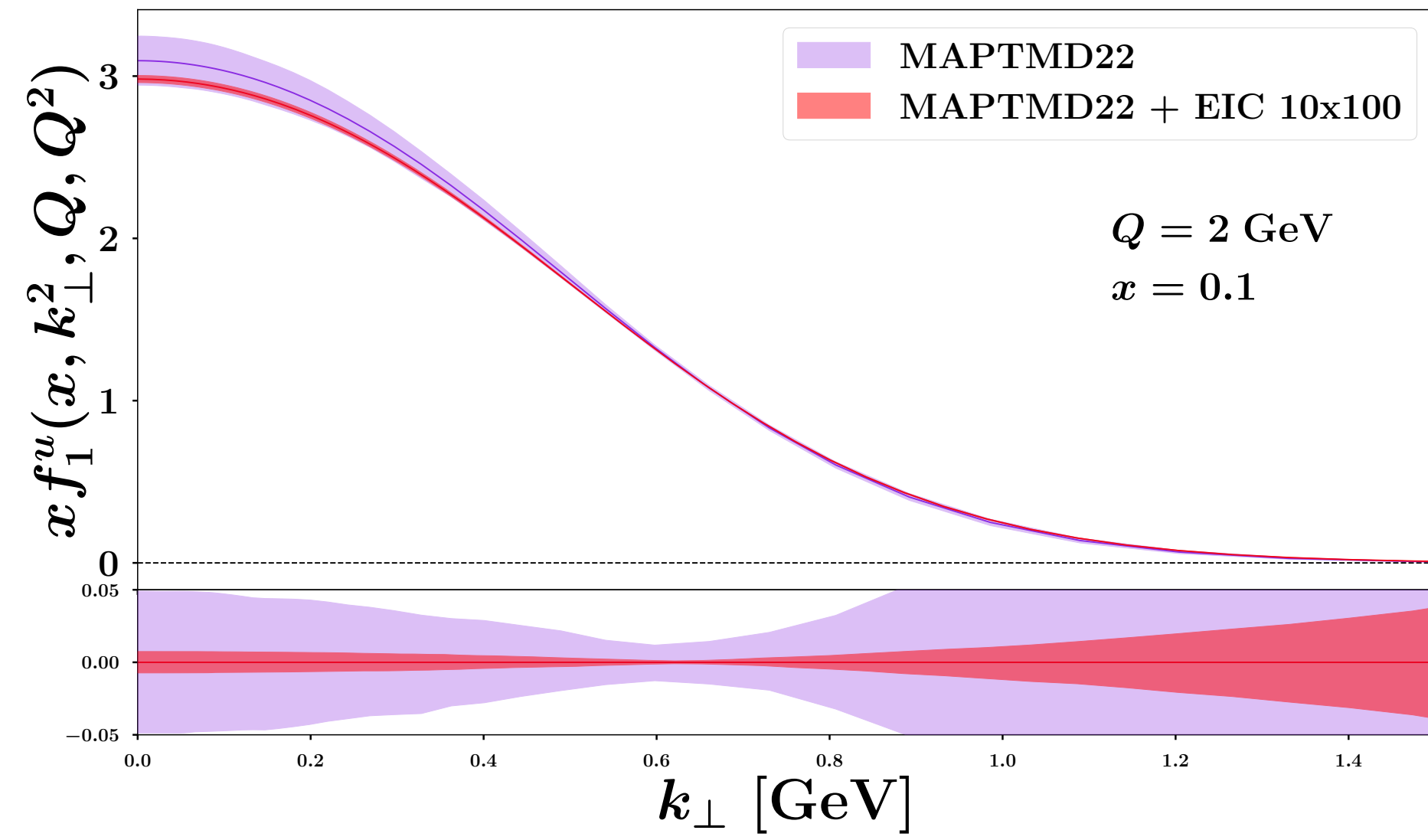


[LPC collaboration, arxiv:2211.02340](#)



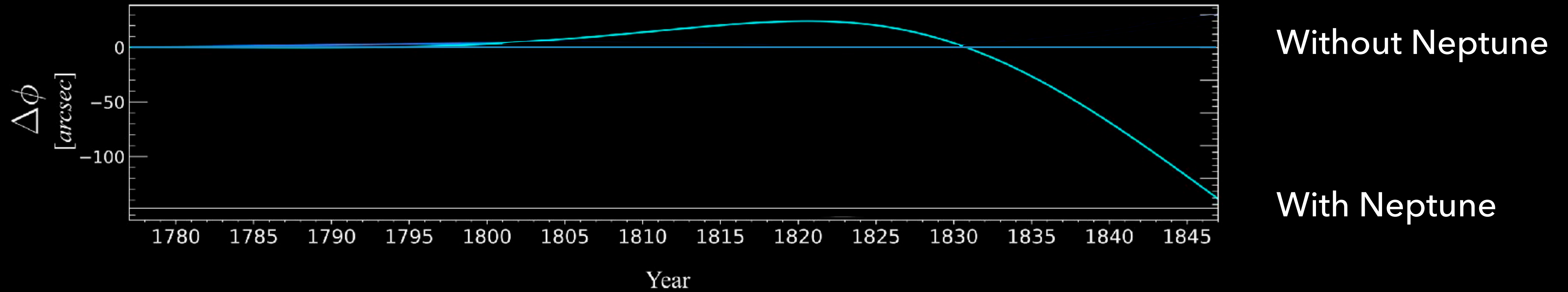




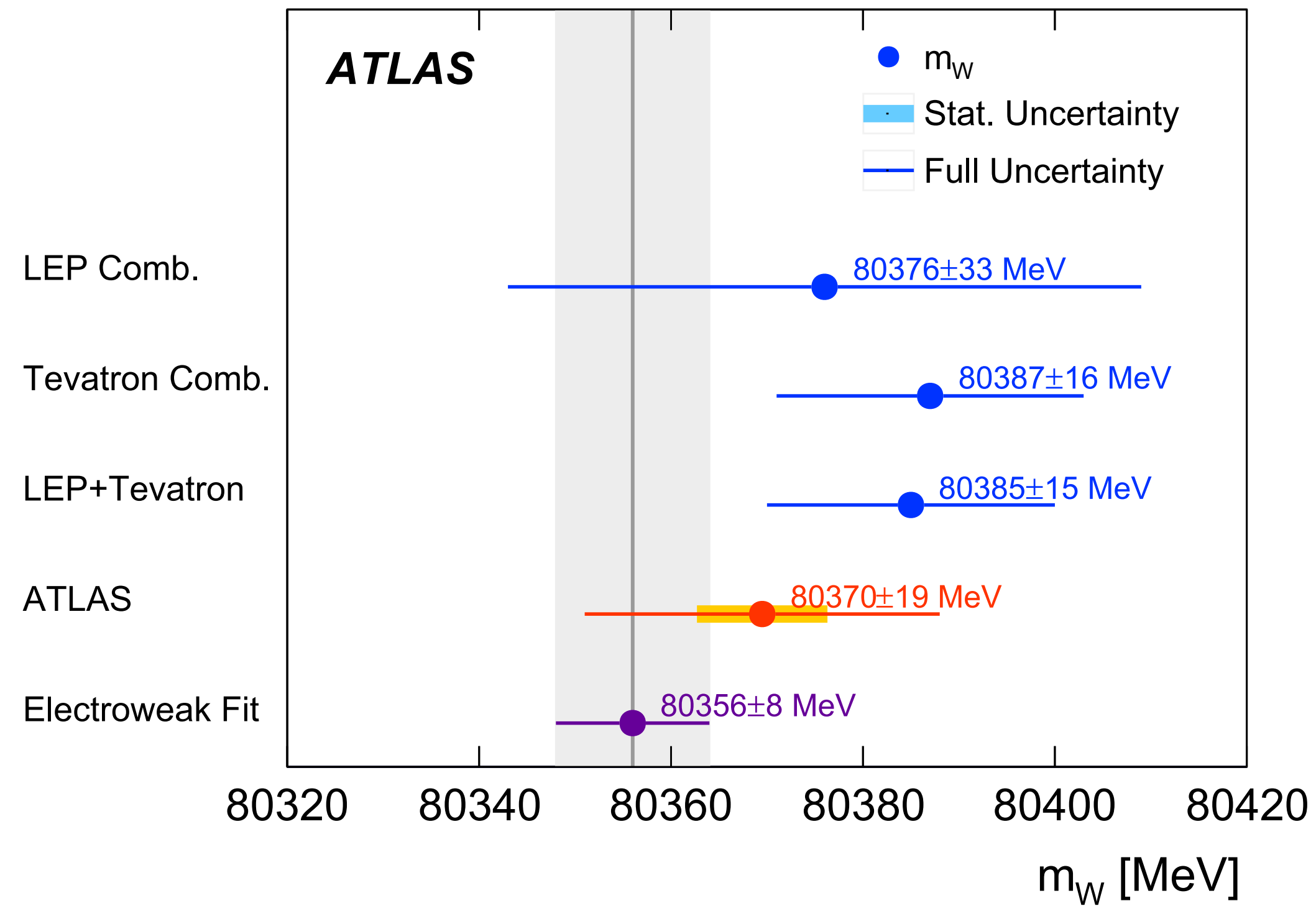


EIC

Uranus's longitude predictions



[ATLAS Collab. arXiv:1701.07240](https://arxiv.org/abs/1701.07240)

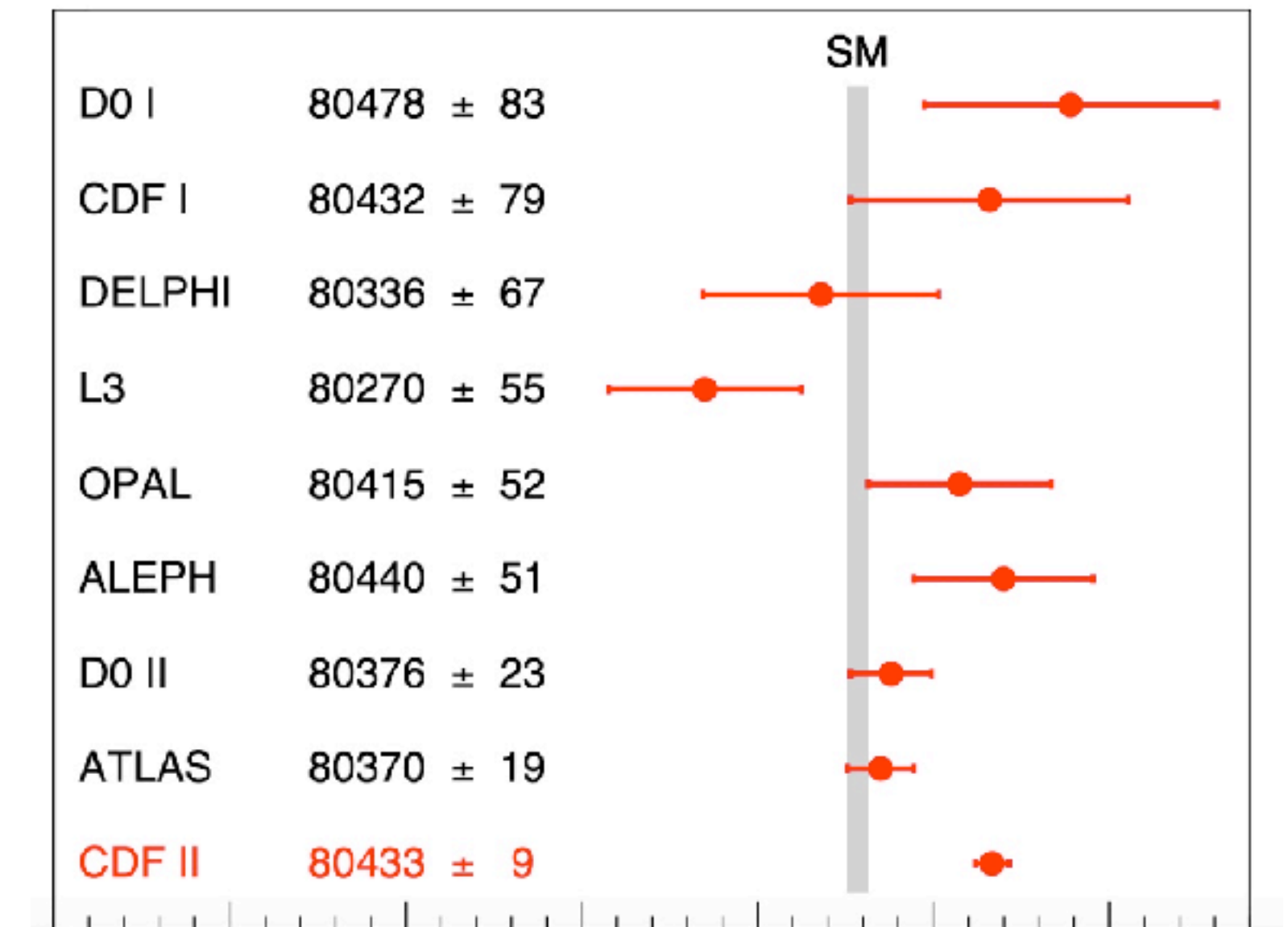
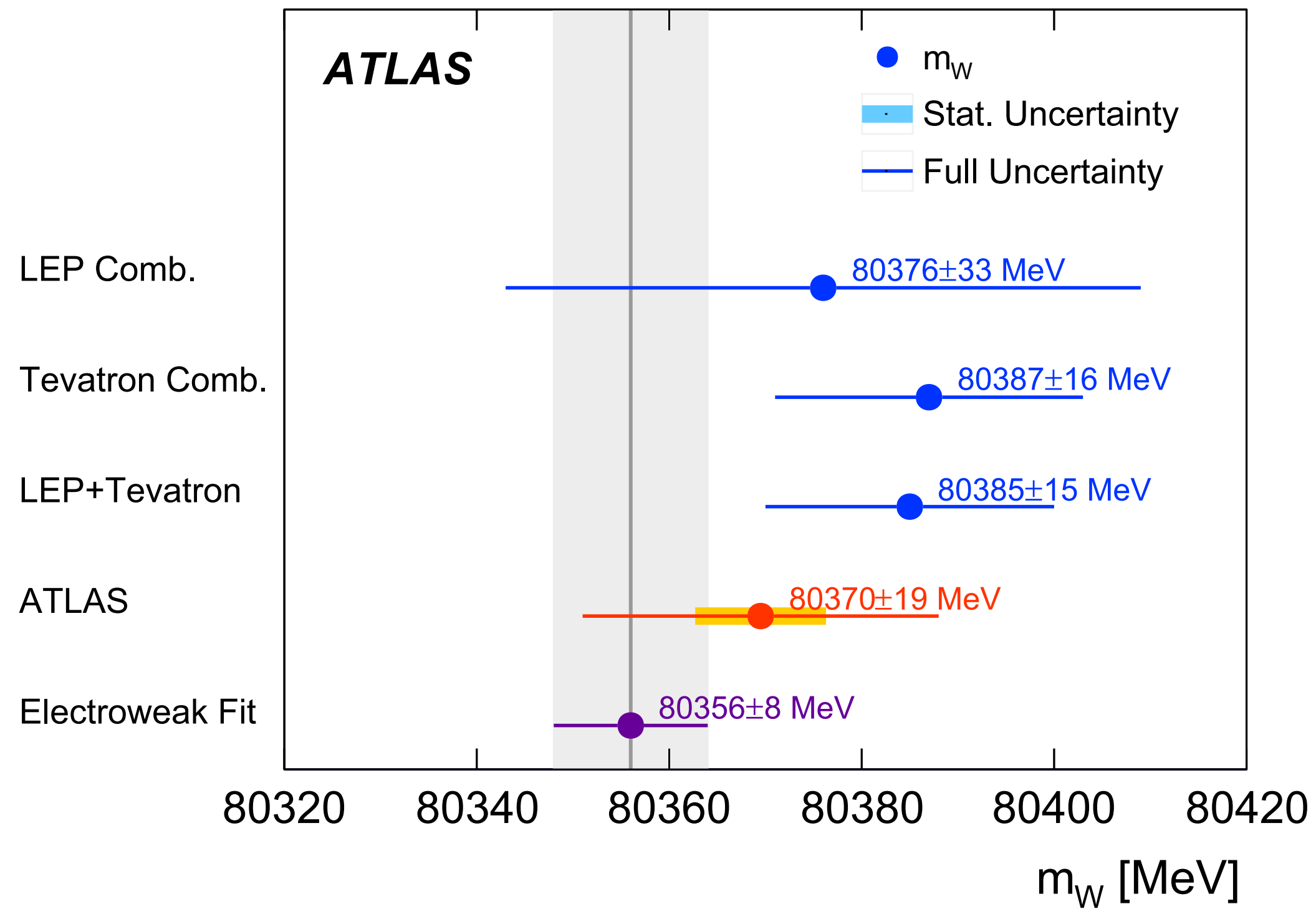


$$\begin{aligned}
 m_W &= 80370 \pm 7 \text{ (stat.)} \pm 11 \text{ (exp. syst.)} \pm 14 \text{ (mod. syst.) MeV} \\
 &= 80370 \pm 19 \text{ MeV,}
 \end{aligned}$$

$$m_{W^+} - m_{W^-} = -29 \pm 28 \text{ MeV.}$$

[ATLAS Collab. arXiv:1701.07240](https://arxiv.org/abs/1701.07240)

[CDF Collab.. Science 2022](https://arxiv.org/abs/2106.11917)



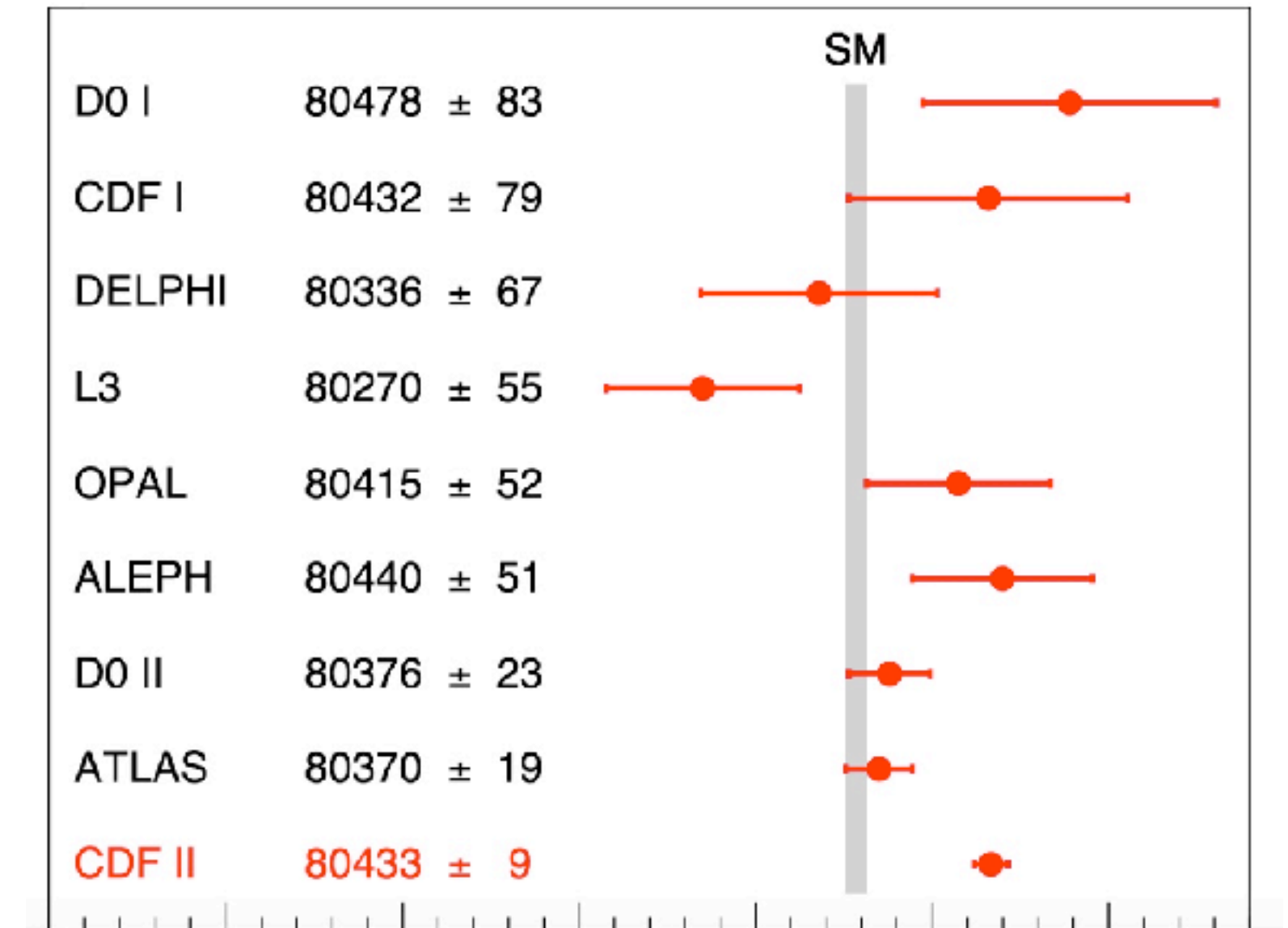
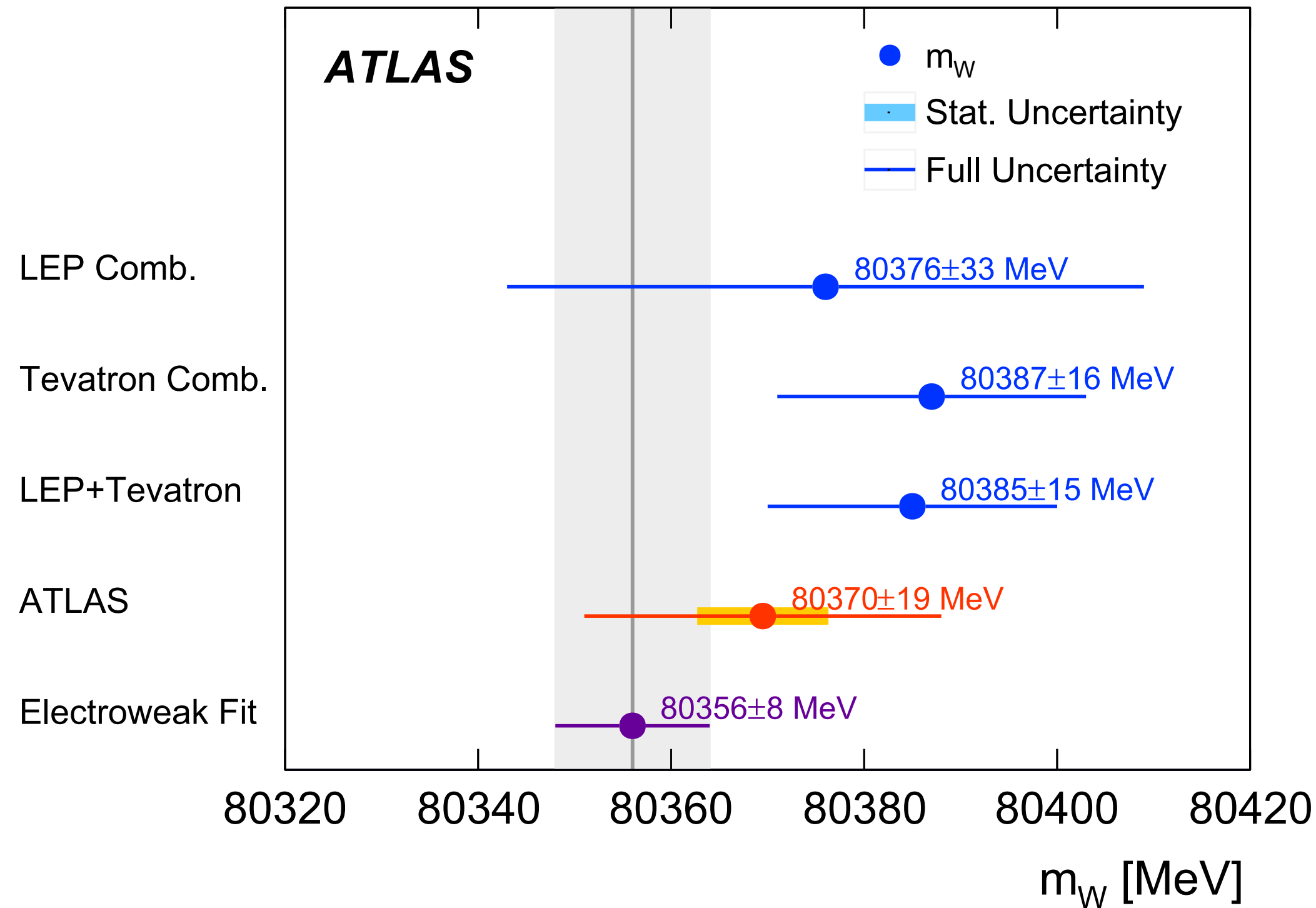
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[ATLAS Collab. arXiv:1701.07240](#)

[CDF Collab.. Science 2022](#)



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$$= 80370 \pm 19 \text{ MeV,}$$

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All analyses assume that TMDs are not flavor dependent. What happens if they are?

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

Try some judicious choices of flavour dependent widths and check

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Set	u_v	d_v	u_s	d_s	s
1	0.34	0.26	0.46	0.59	0.32
2	0.34	0.46	0.56	0.32	0.51
3	0.55	0.34	0.33	0.55	0.30
4	0.53	0.49	0.37	0.22	0.52
5	0.42	0.38	0.29	0.57	0.27

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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narrow, medium, large

narrow, large, narrow

large, narrow, large

large, medium, narrow

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[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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They all describe the
Z spectrum very well

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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They all describe the Z spectrum very well



	ΔM_{W+}		ΔM_{W-}	
Set	m_T	$p_{T\ell}$	m_T	$p_{T\ell}$
1	0	-1	-2	3
2	0	-6	-2	0
3	-1	9	-2	-4
4	0	0	-2	-4
5	0	4	-1	-3

[Bacchetta, Bozzi, Radici, Ritzmann, Signori, arXiv:1807.02101](#)

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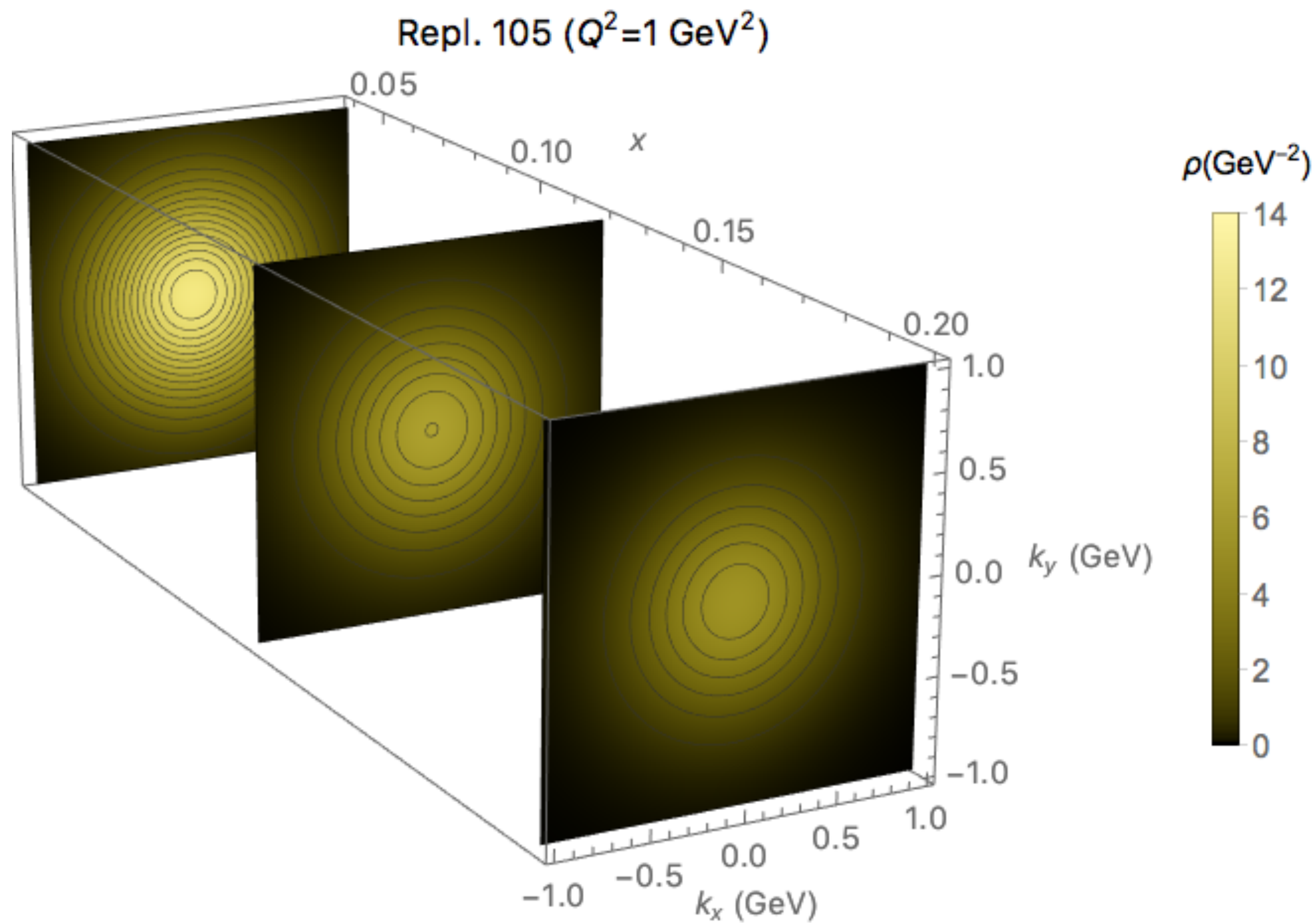
narrow, medium, large
 narrow, large, narrow
 large, narrow, large
 large, medium, narrow
 medium, narrow, large

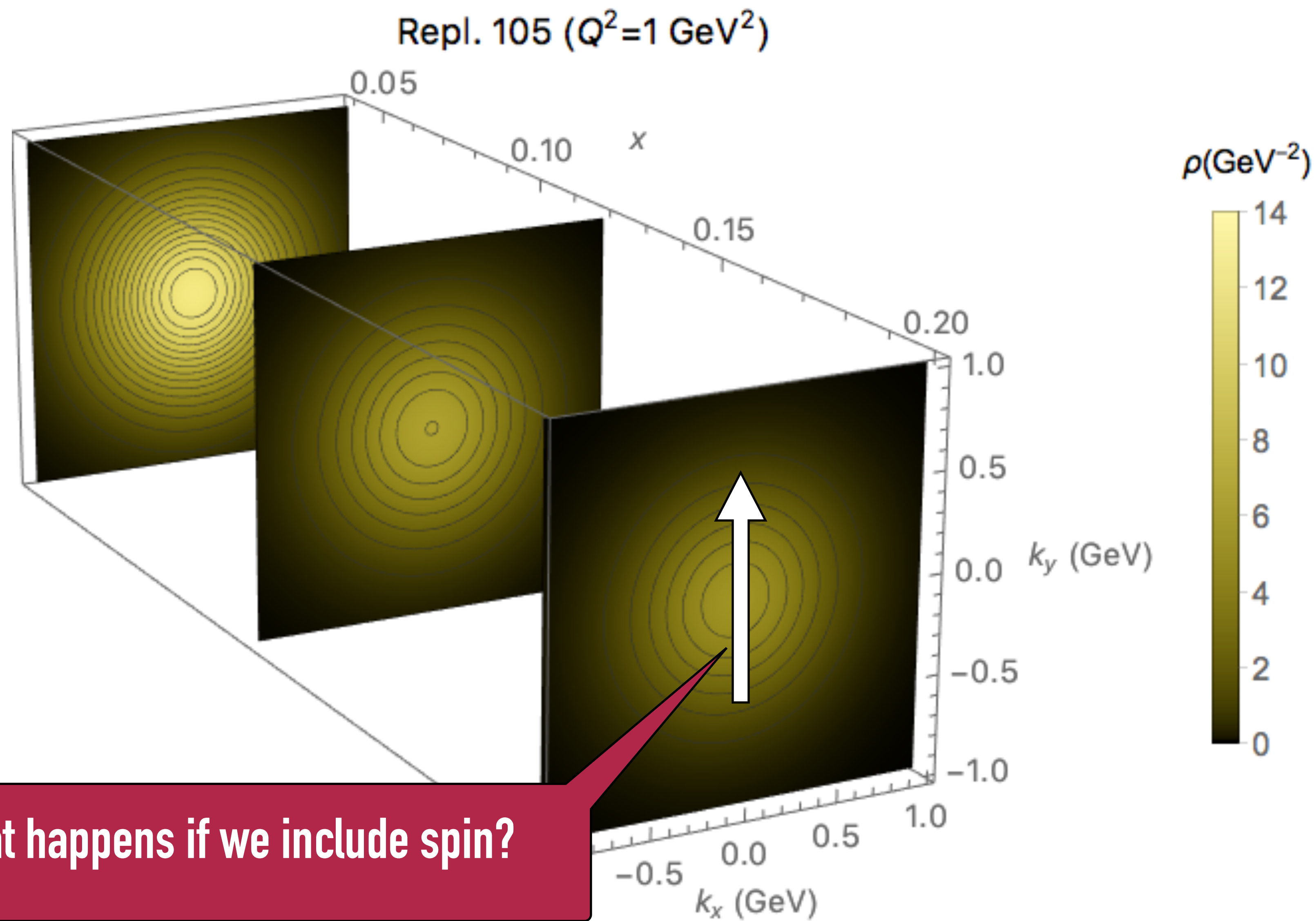
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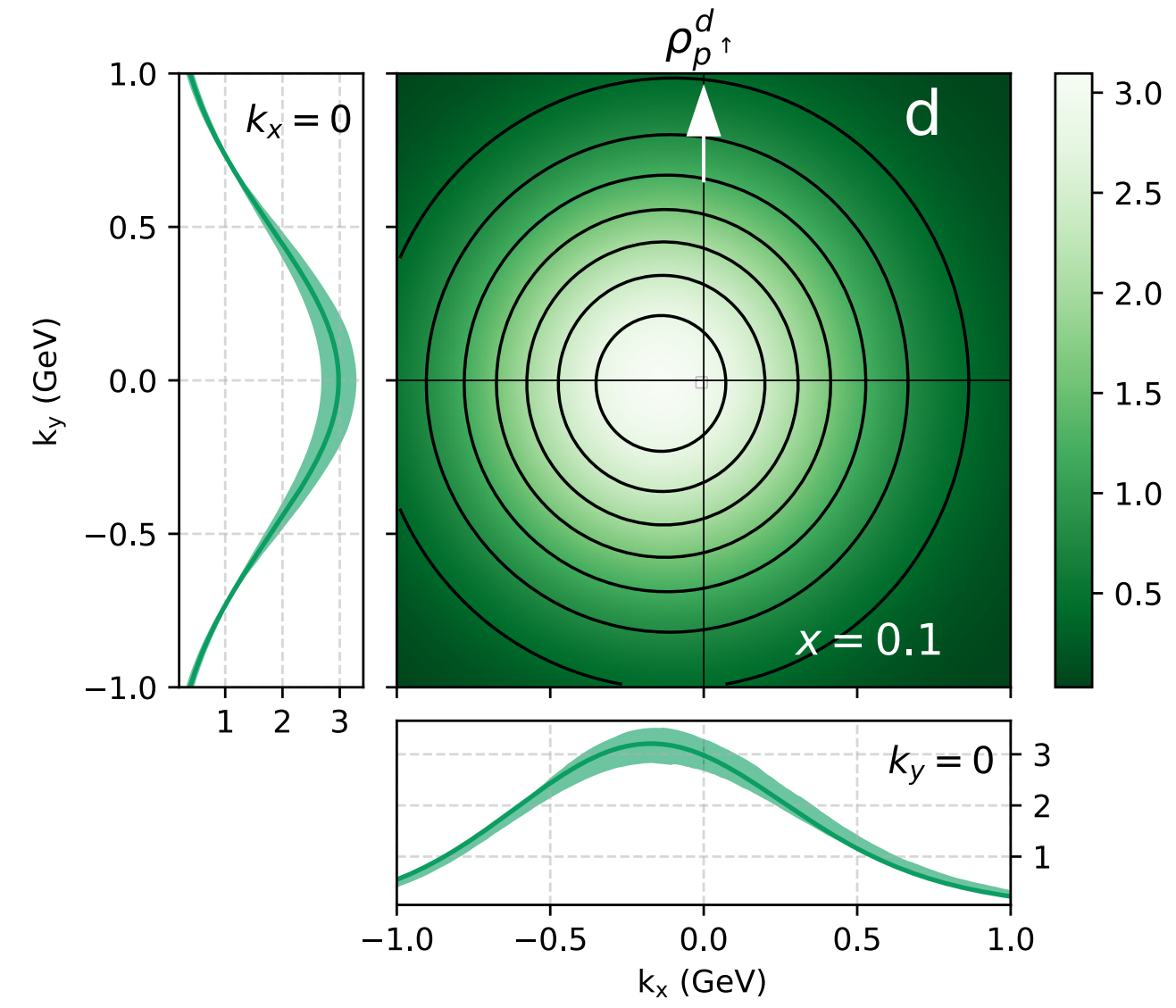
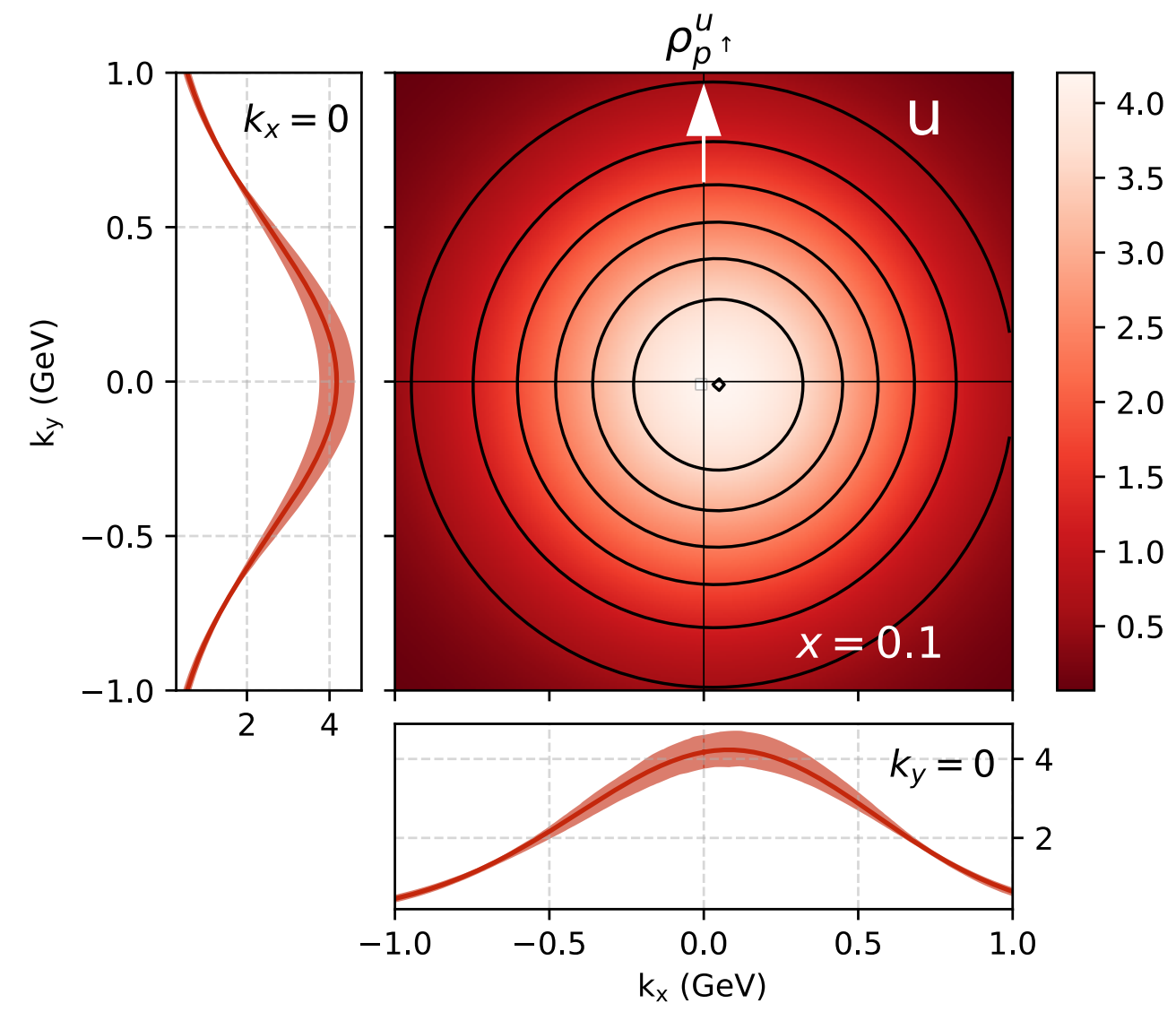
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Not taking into account the flavor dependence of TMDs can lead to errors in the determination of the W mass, of the order of a few MeVs



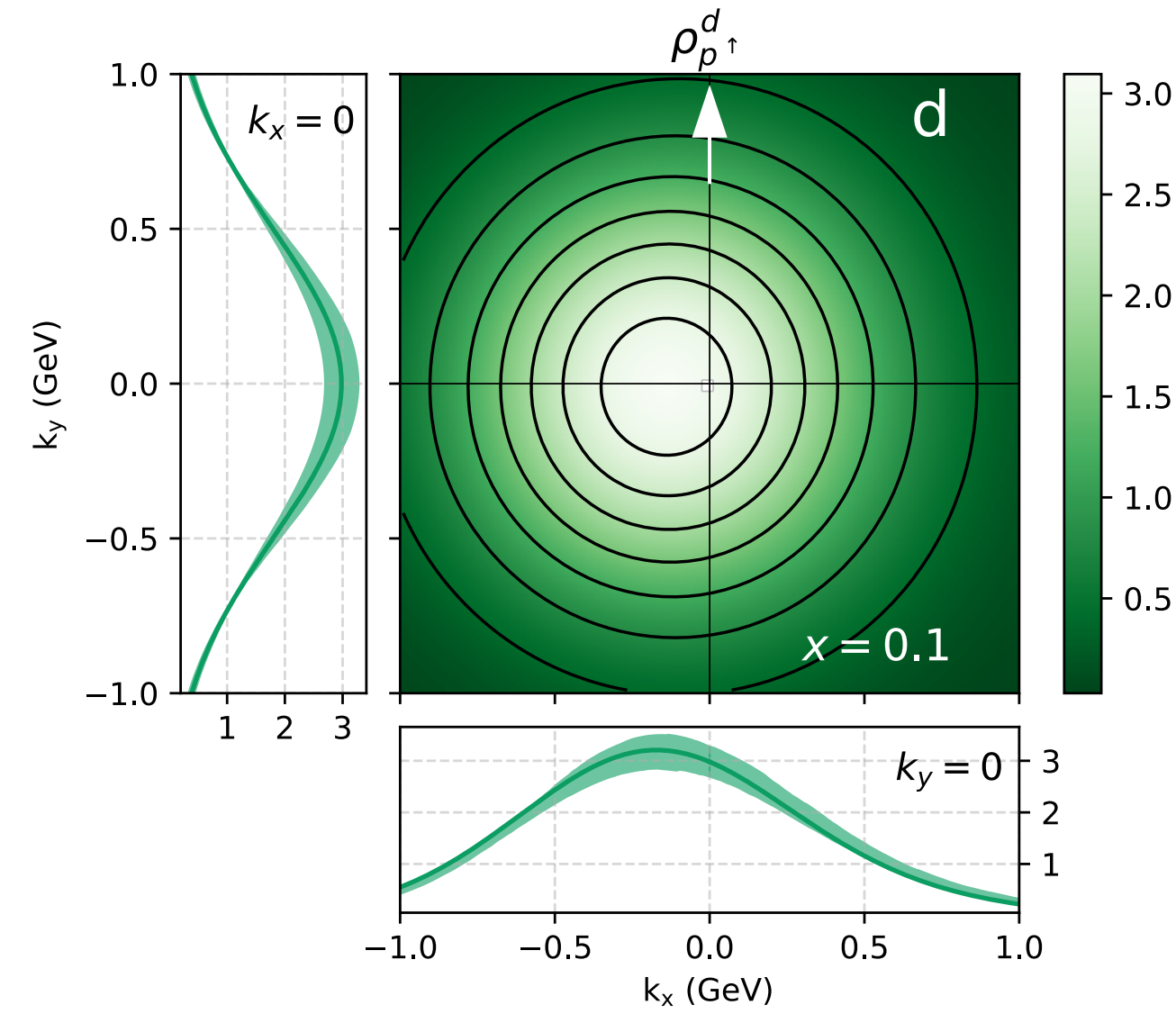
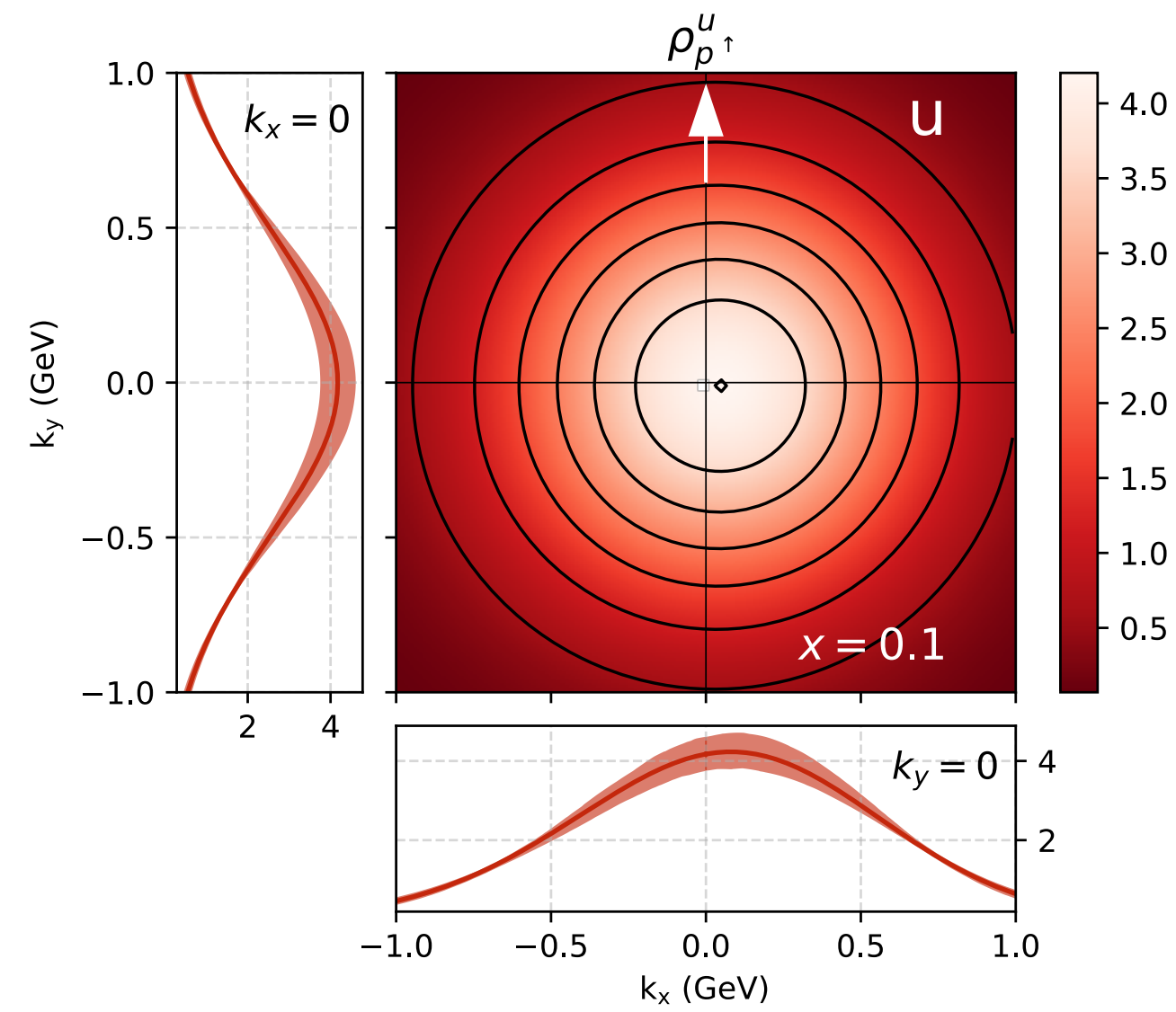


What happens if we include spin?



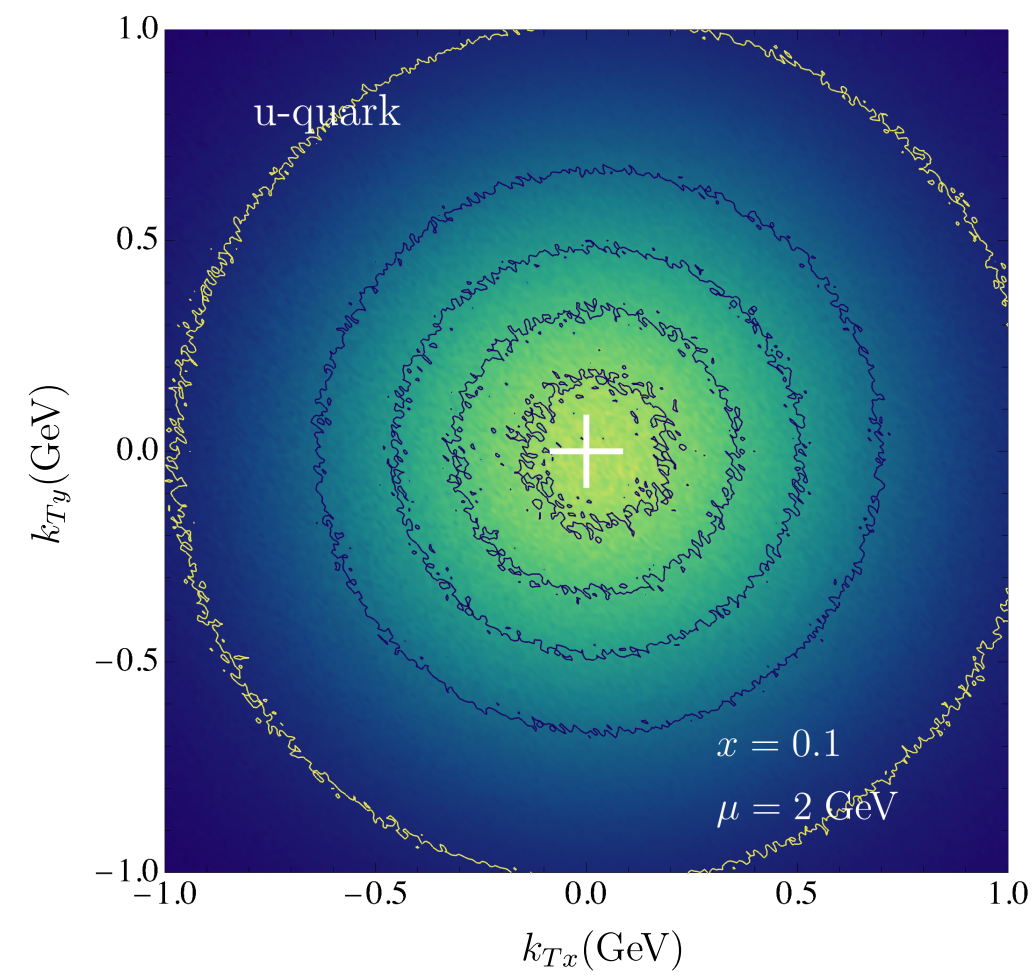
$Q = 2\text{GeV}$

[Bacchetta, Delcarro,
Pisano, Radici,
arXiv:2004.14278](#)

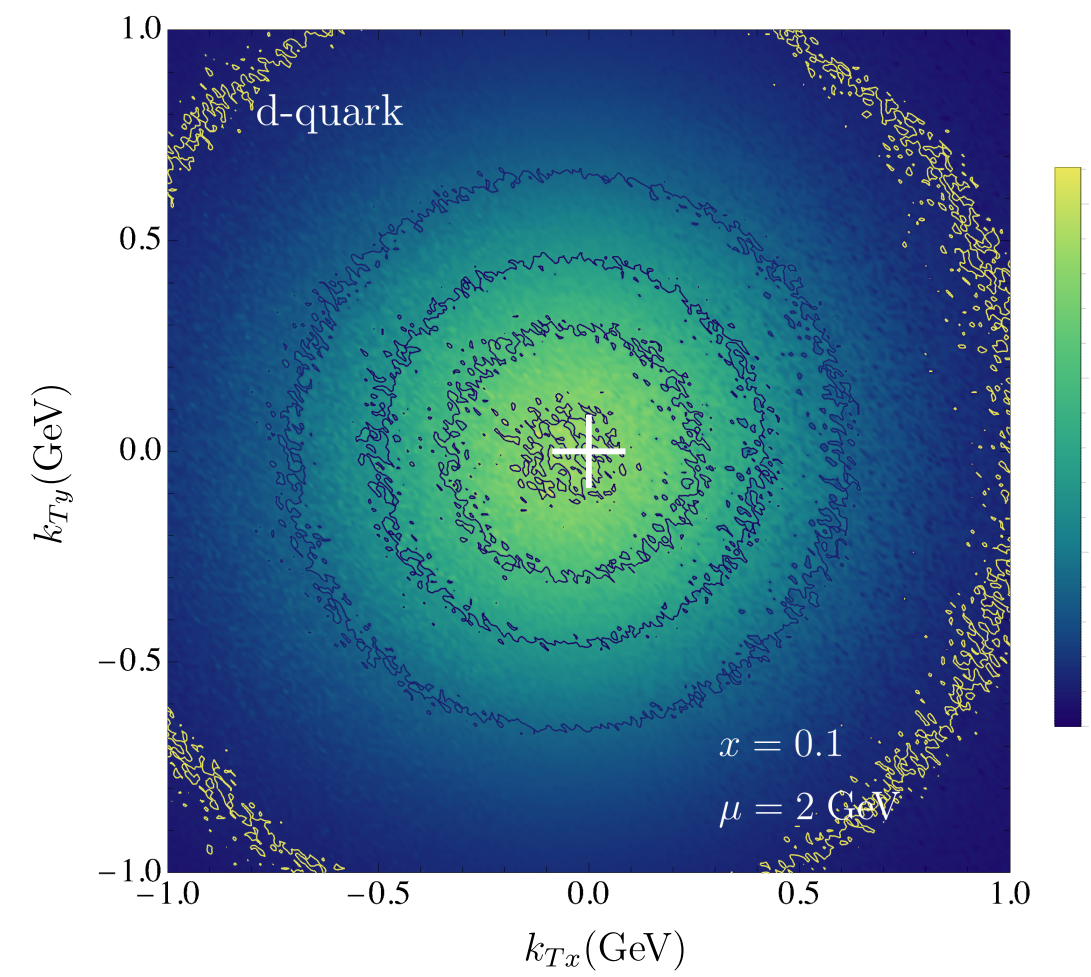


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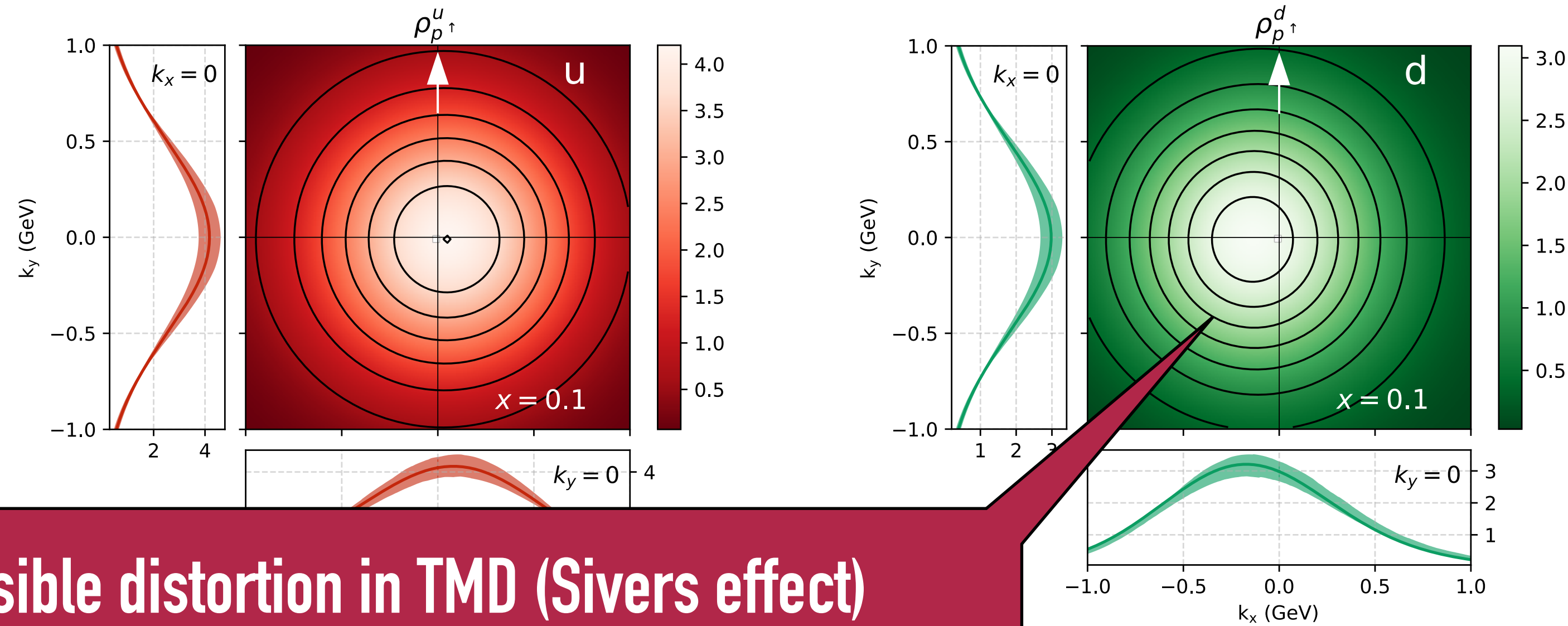


(a)



(b)

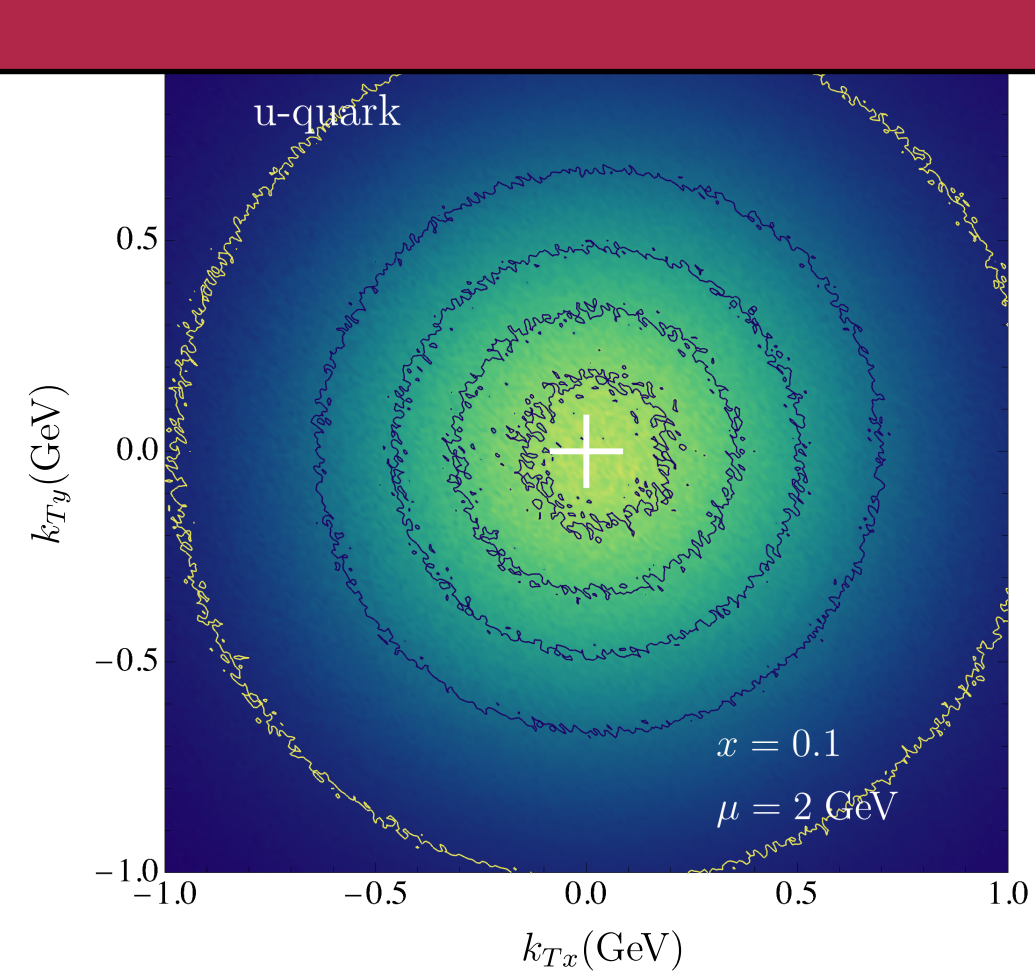
[Bury, Prokudin, Vladimirov, arXiv:2103.03270](#)



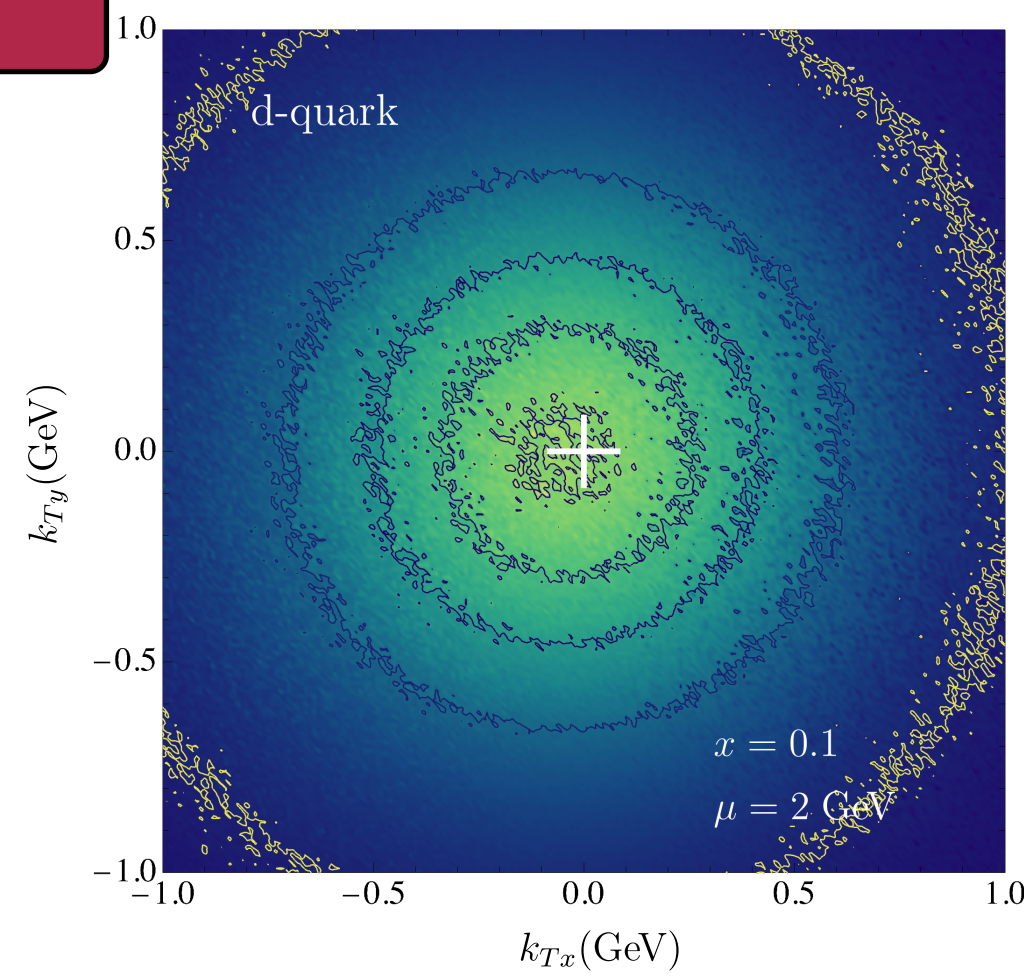
$Q = 2 \text{ GeV}$

[Bacchetta, Delcarro, Pisano, Radici, arXiv:2004.14278](#)

Possible distortion in TMD (Sivers effect)



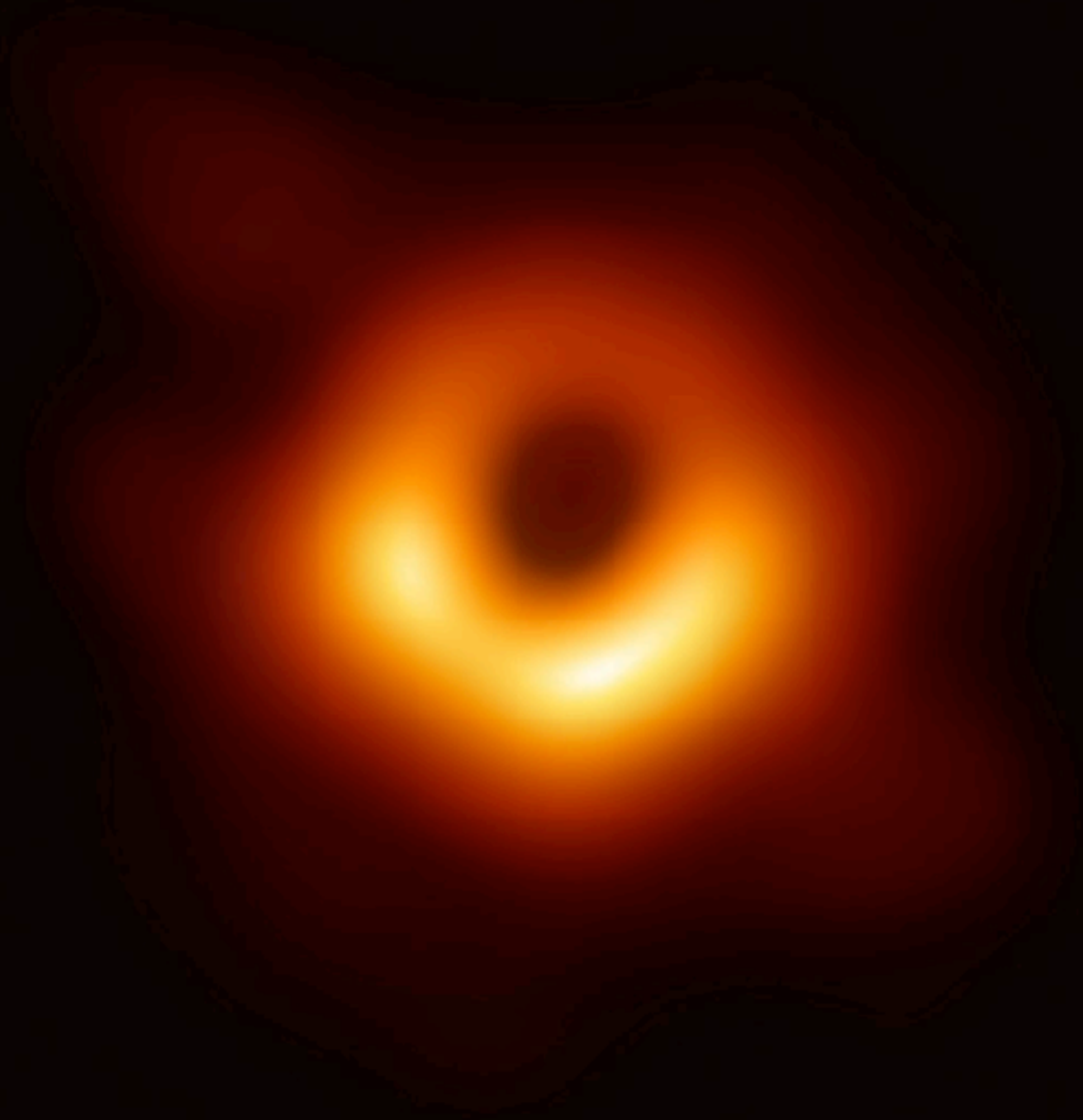
(a)



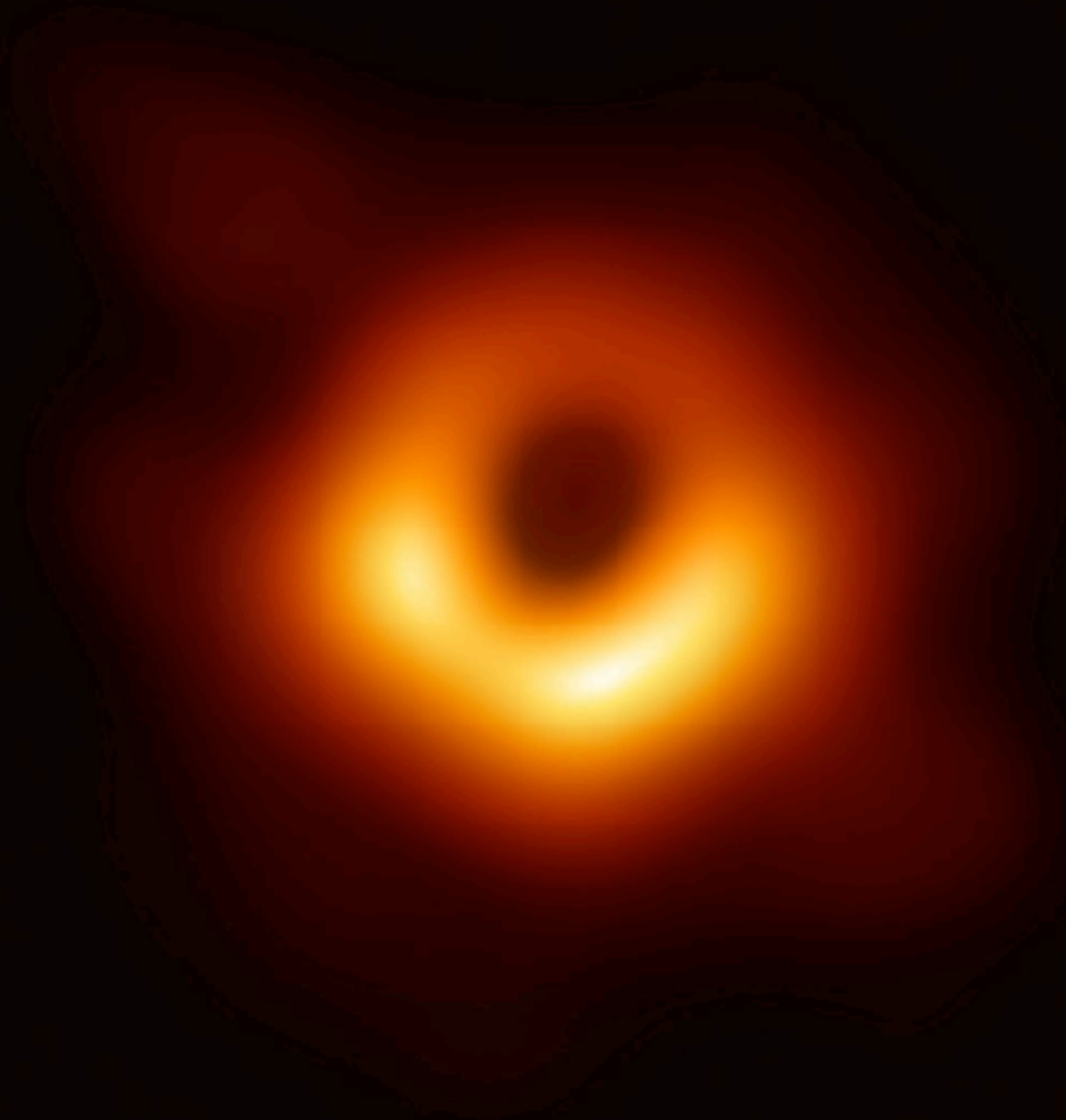
(b)

[Bury, Prokudin, Vladimirov, arXiv:2103.03270](#)

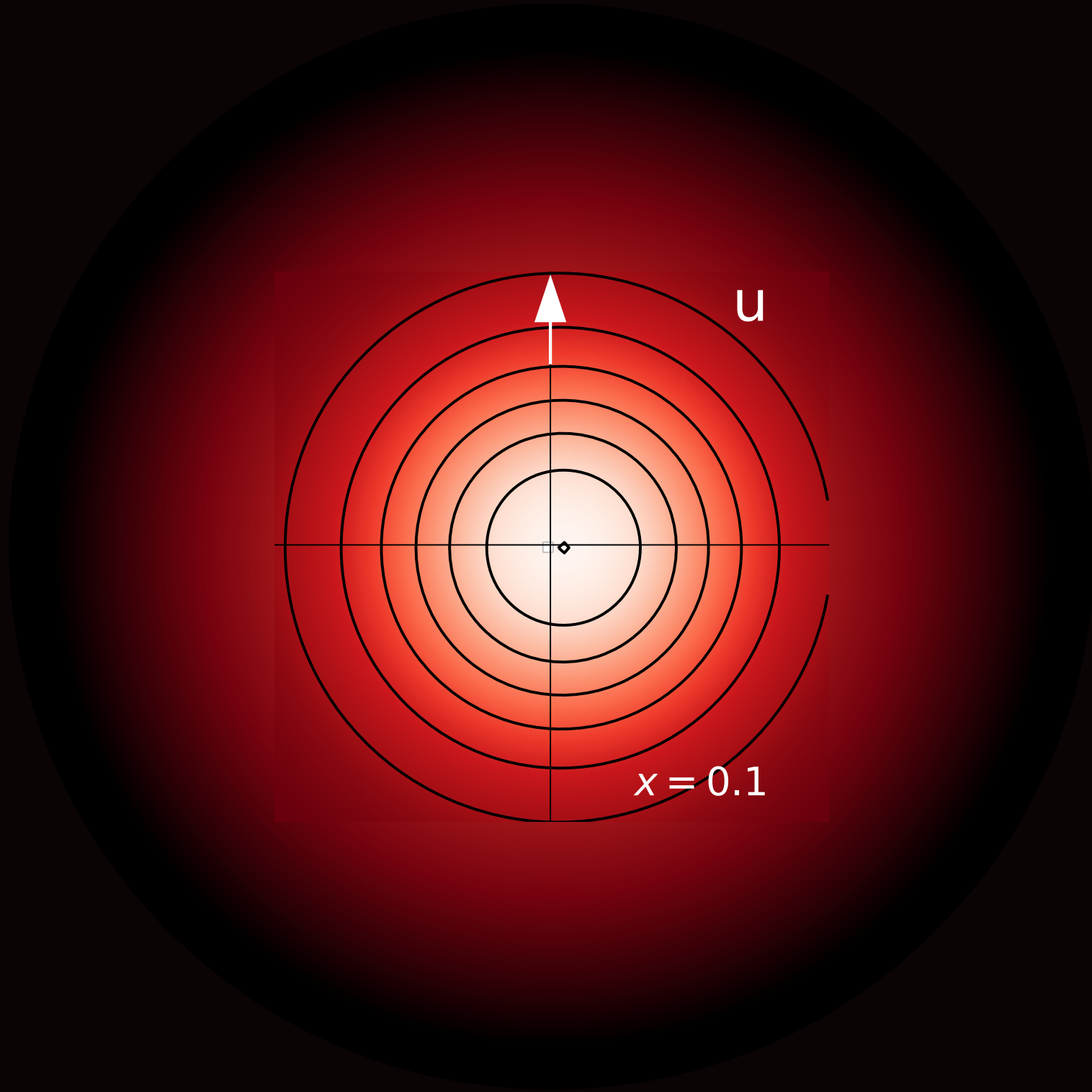
A picture of a black hole (2019)



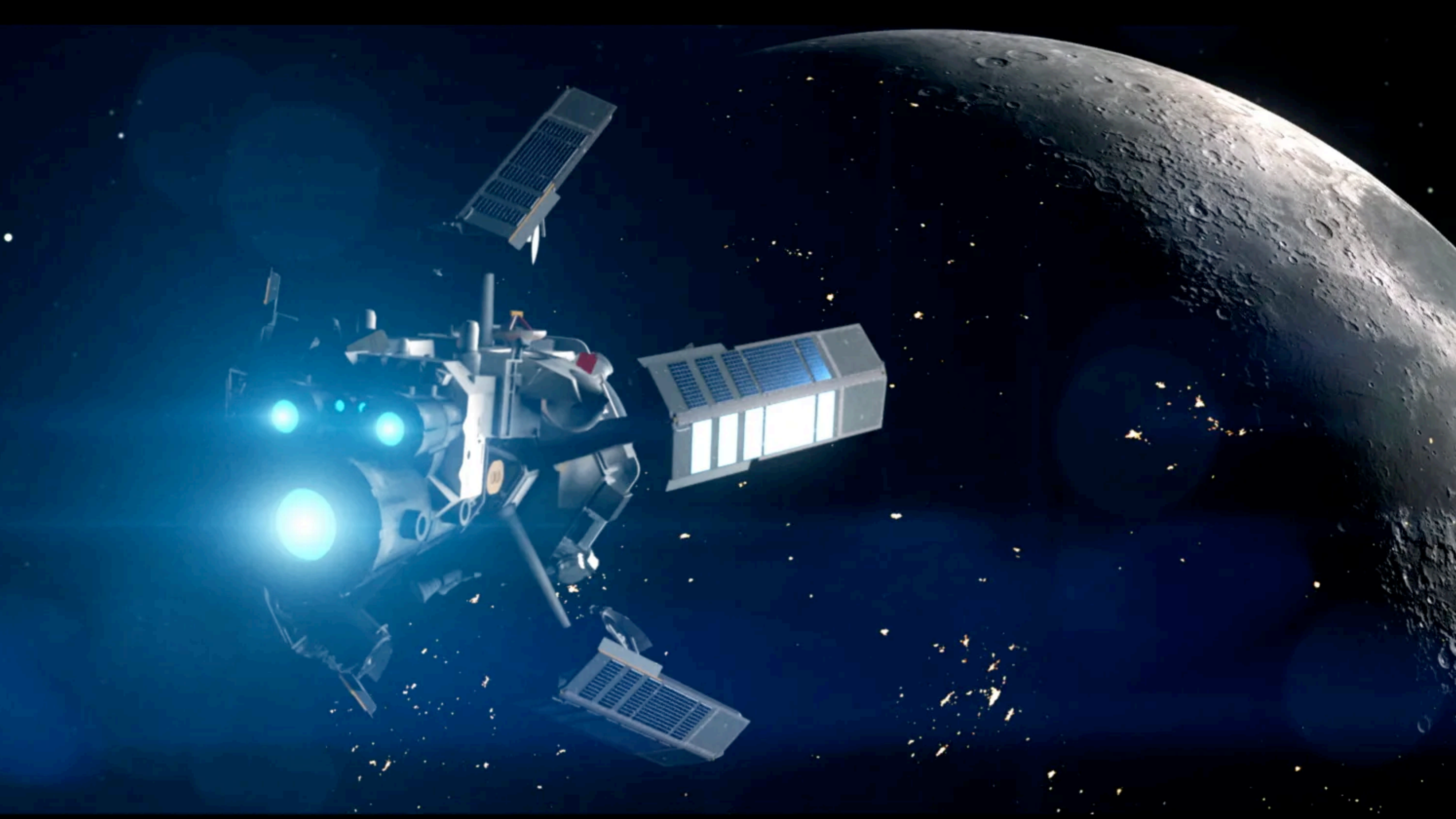
A picture of a black hole (2019)



A picture of a proton (2020)



Our world is made of electrons, photons, quarks, and gluons:
I believe we will find ways to use them before we use the Higgs bosons or black holes.



- ▶ The EIC will be a groundbreaking machine for QCD studies

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- ▶ I discussed some opportunities to study the multidimensional structure of nucleons, but there are many more

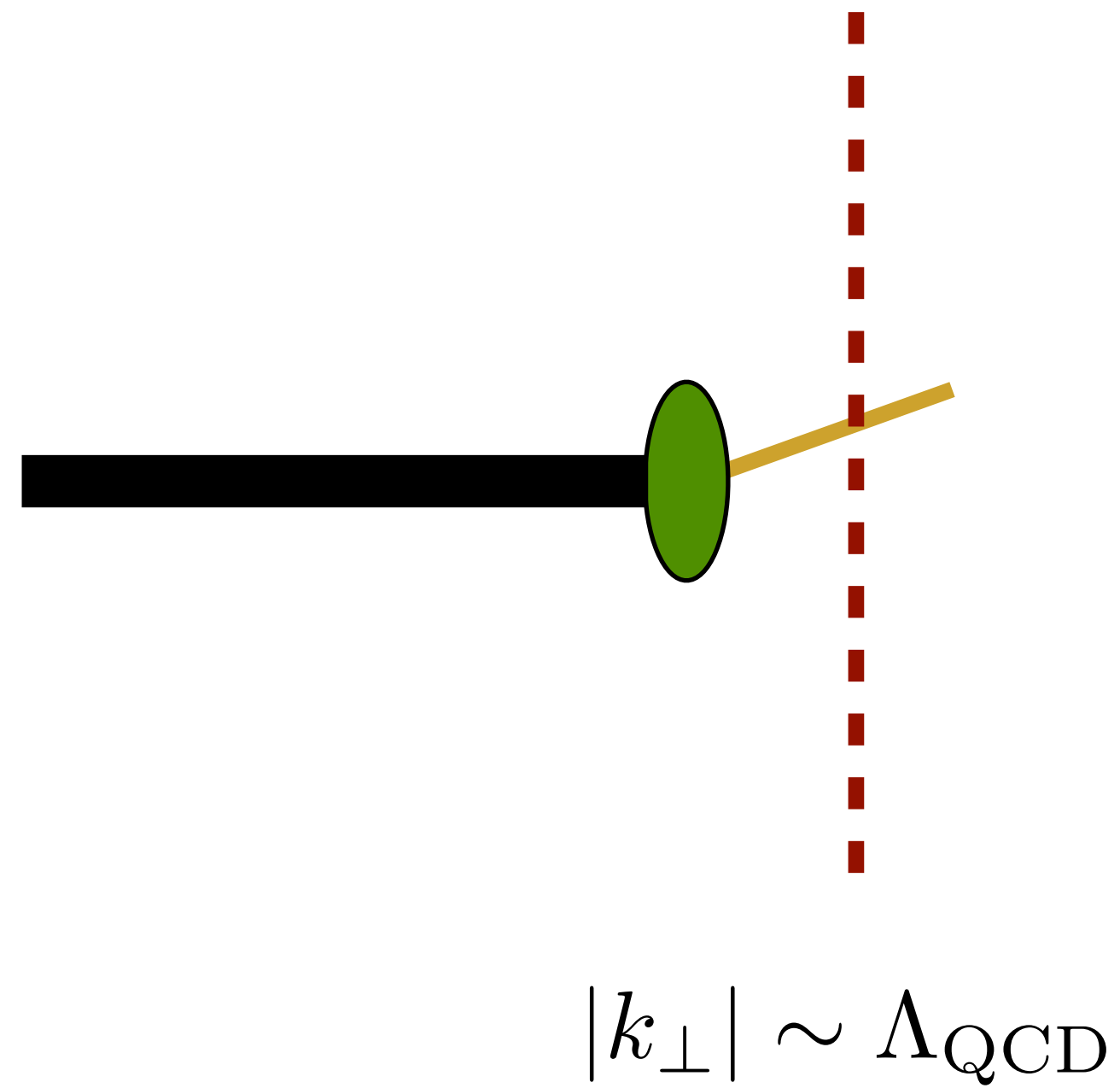
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- ▶ Results can be used to check lattice QCD predictions and look for new physics
- ▶ The long-term goal is the capability of computing the multidimensional structure of the nucleon, and eventually of the nucleus, and the hadronization process, all based on QCD

BACKUP

DIFFERENT CONTRIBUTIONS TO TRANSVERSE MOMENTUM

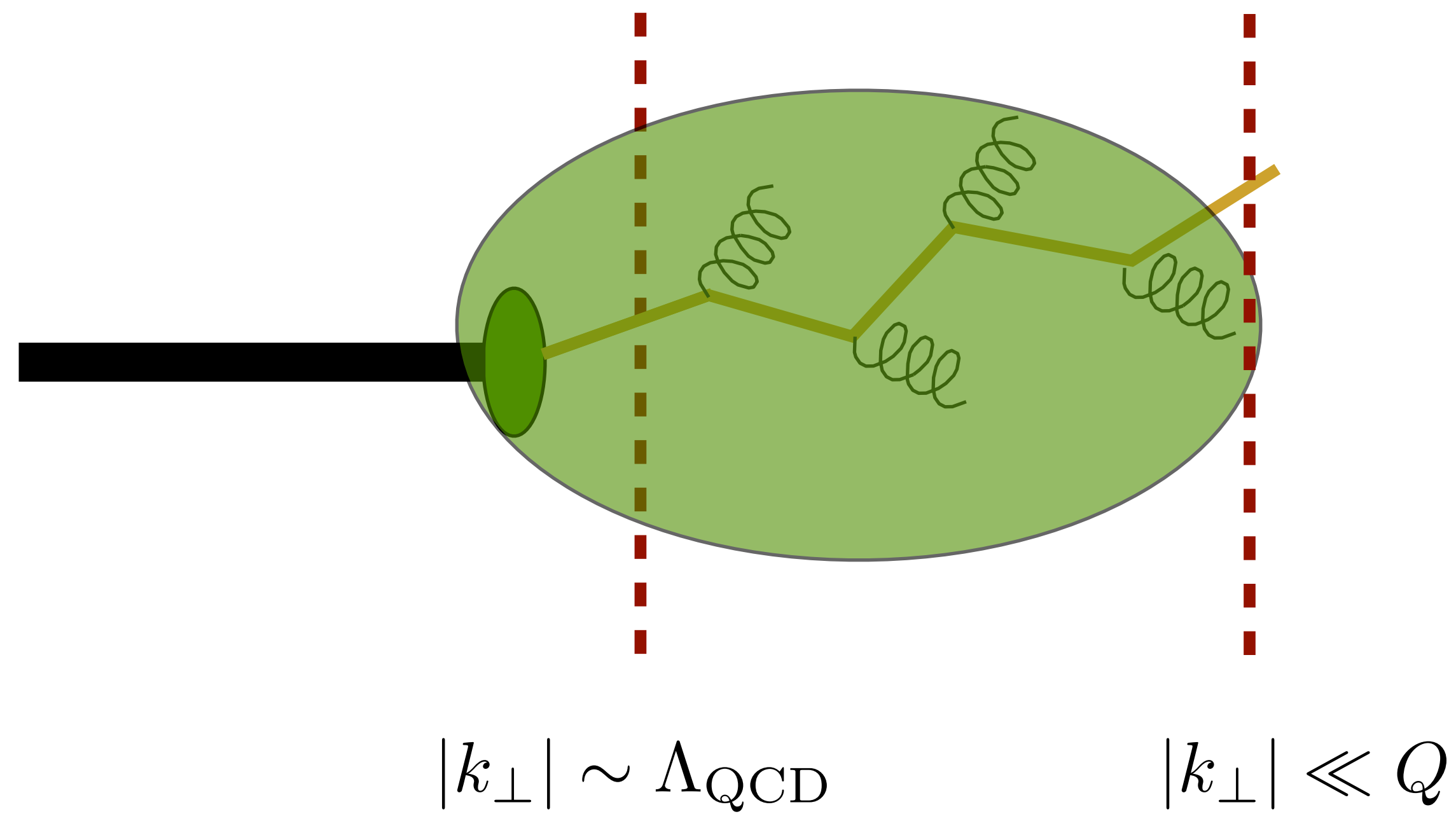
“intrinsic”
transverse
momentum



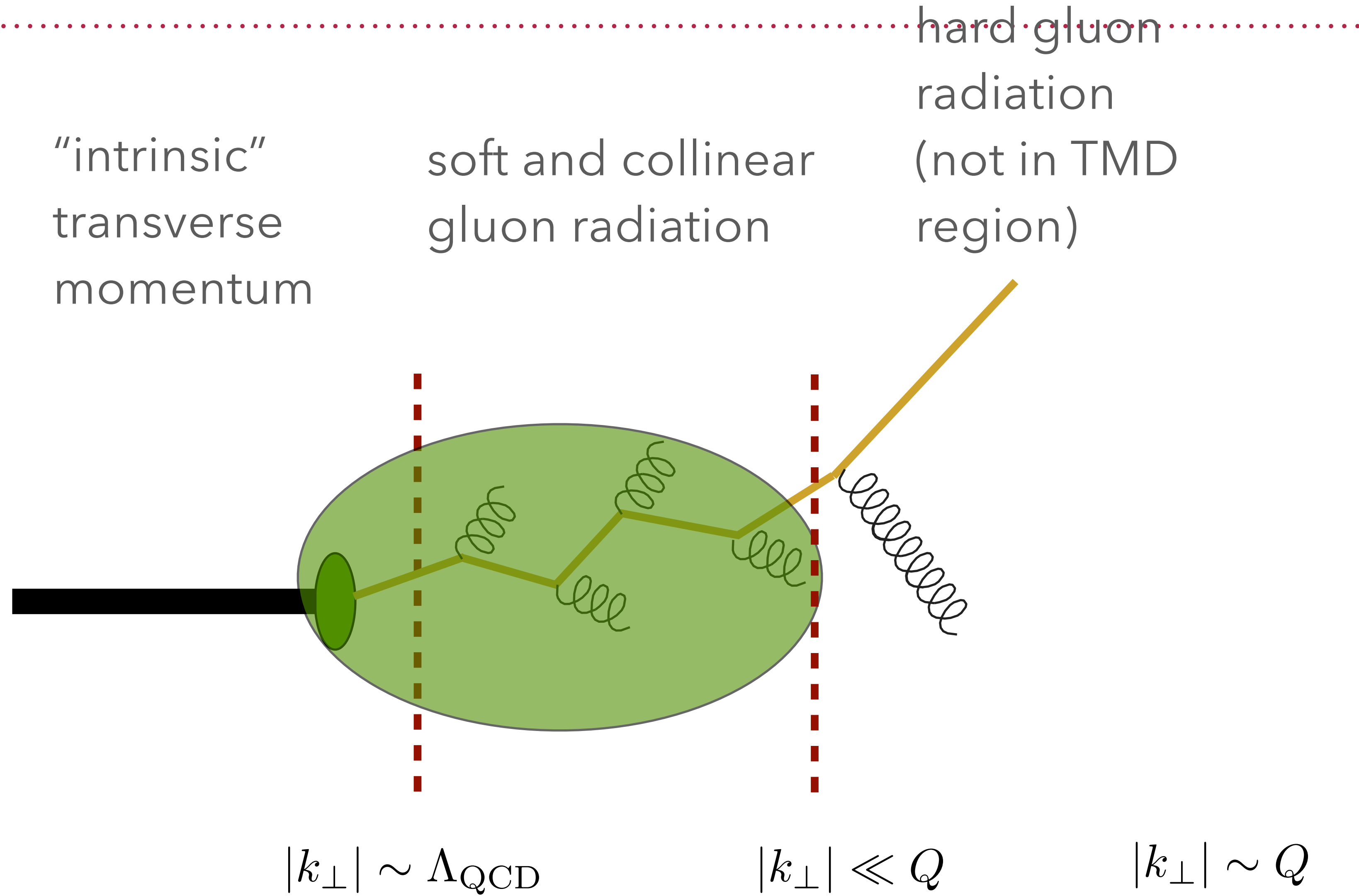
DIFFERENT CONTRIBUTIONS TO TRANSVERSE MOMENTUM

“intrinsic”
transverse
momentum

soft and collinear
gluon radiation



DIFFERENT CONTRIBUTIONS TO TRANSVERSE MOMENTUM



Sudakov form factor

$$\text{LL} \quad \alpha_S^n \ln^{2n} \left(\frac{Q^2}{\mu_b^2} \right)$$

Sudakov form factor

LL $\alpha_S^n \ln^{2n} \left(\frac{Q^2}{\mu_b^2} \right)$

NLL $\alpha_S^n \ln^{2n} \left(\frac{Q^2}{\mu_b^2} \right), \quad \alpha_S^n \ln^{2n-1} \left(\frac{Q^2}{\mu_b^2} \right)$

Sudakov form factor

matching coeff.

LL $\alpha_S^n \ln^{2n} \left(\frac{Q^2}{\mu_b^2} \right)$

C^0

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$$C^0$$

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$$\left(C^0 + \alpha_S C^1 \right)$$

Sudakov form factor

matching coeff.

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$(C^0 + \alpha_S C^1)$

the difference between the two is formally NNLL $\alpha_S^n \ln^{2n-2} \left(\frac{Q^2}{\mu_b^2} \right)$

$$\log(Q^2 b_T^2) \rightarrow \log(Q^2 b_T^2 + 1)$$

*[see, e.g., Bozzi, Catani, De Florian, Grazzini
hep-ph/0302104](#)*

$$\log(Q^2 b_T^2) \rightarrow \log(Q^2 b_T^2 + 1)$$

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[hep-ph/0302104](https://arxiv.org/abs/hep-ph/0302104)

$$b_*(b_c(b_T)) = \sqrt{\frac{b_T^2 + b_0^2/(C_5^2 Q^2)}{1 + b_T^2/b_{\max}^2 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}} \quad b_{\min} \equiv b_*(b_c(0)) = \frac{b_0}{C_5 Q} \sqrt{\frac{1}{1 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}}$$

Collins et al.
[arXiv:1605.00671](https://arxiv.org/abs/1605.00671)

$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\max}^2}}$$

Collins, Soper, Sterman, NPB250 (85)

$$\mu_0 = 1 \text{ GeV}$$

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Collins, Soper, Sterman, NPB250 (85)

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$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\text{max}}^2}}$$

Collins, Soper, Sterman, NPB250 (85)

$$\mu_b = 2e^{-\gamma_E}/b_*$$

$$\bar{b}_* \equiv b_{\text{max}} \left(\frac{1 - e^{-b_T^4/b_{\text{max}}^4}}{1 - e^{-b_T^4/b_{\text{min}}^4}} \right)^{1/4}$$

$$b_{\text{max}} = 2e^{-\gamma_E}$$

$$b_{\text{min}} = \frac{2e^{-\gamma_E}}{Q}$$

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$$b_{\text{min}} = \frac{2e^{-\gamma_E}}{Q}$$

These are all choices that should be at some point checked/challenged

$$\hat{f}_1^q(x, b_T; \mu^2) = \sum_i (C_{qi} \otimes f_1^i)(x, b_*; \mu_b) e^{\tilde{S}(b_*; \mu_b, \mu)} e^{g_K(b_T) \ln \frac{\mu}{\mu_0}} \hat{f}_{\text{NP}}^q(x, b_T)$$

$$\mu_0 = 1 \text{ GeV}$$

$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\text{max}}^2}} \quad \text{Collins, Soper, Sterman, NPB250 (85)}$$

$$\mu_b = 2e^{-\gamma_E}/b_* \quad \bar{b}_* \equiv b_{\text{max}} \left(\frac{1 - e^{-b_T^4/b_{\text{max}}^4}}{1 - e^{-b_T^4/b_{\text{min}}^4}} \right)^{1/4} \quad b_{\text{max}} = 2e^{-\gamma_E}$$

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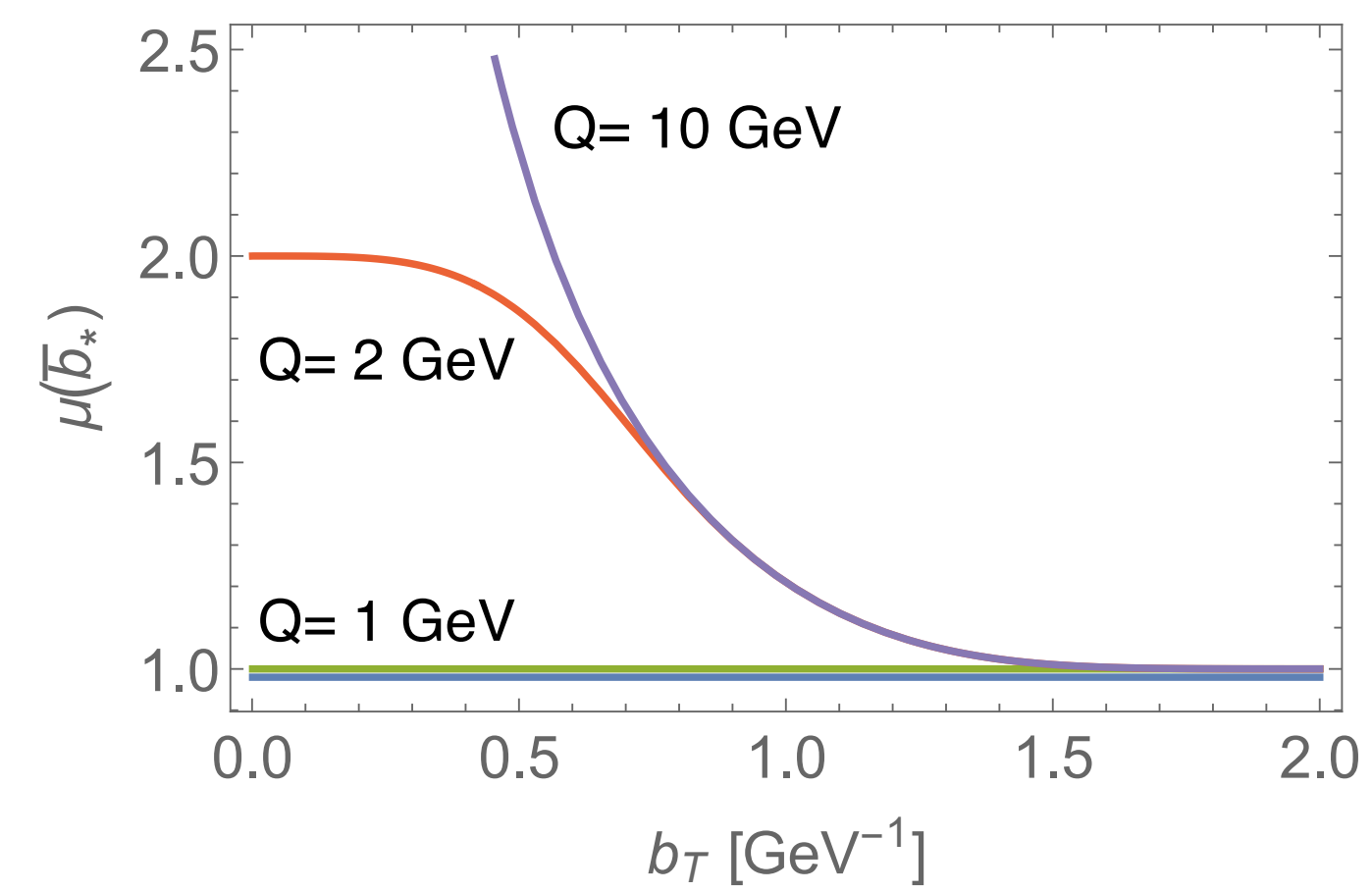
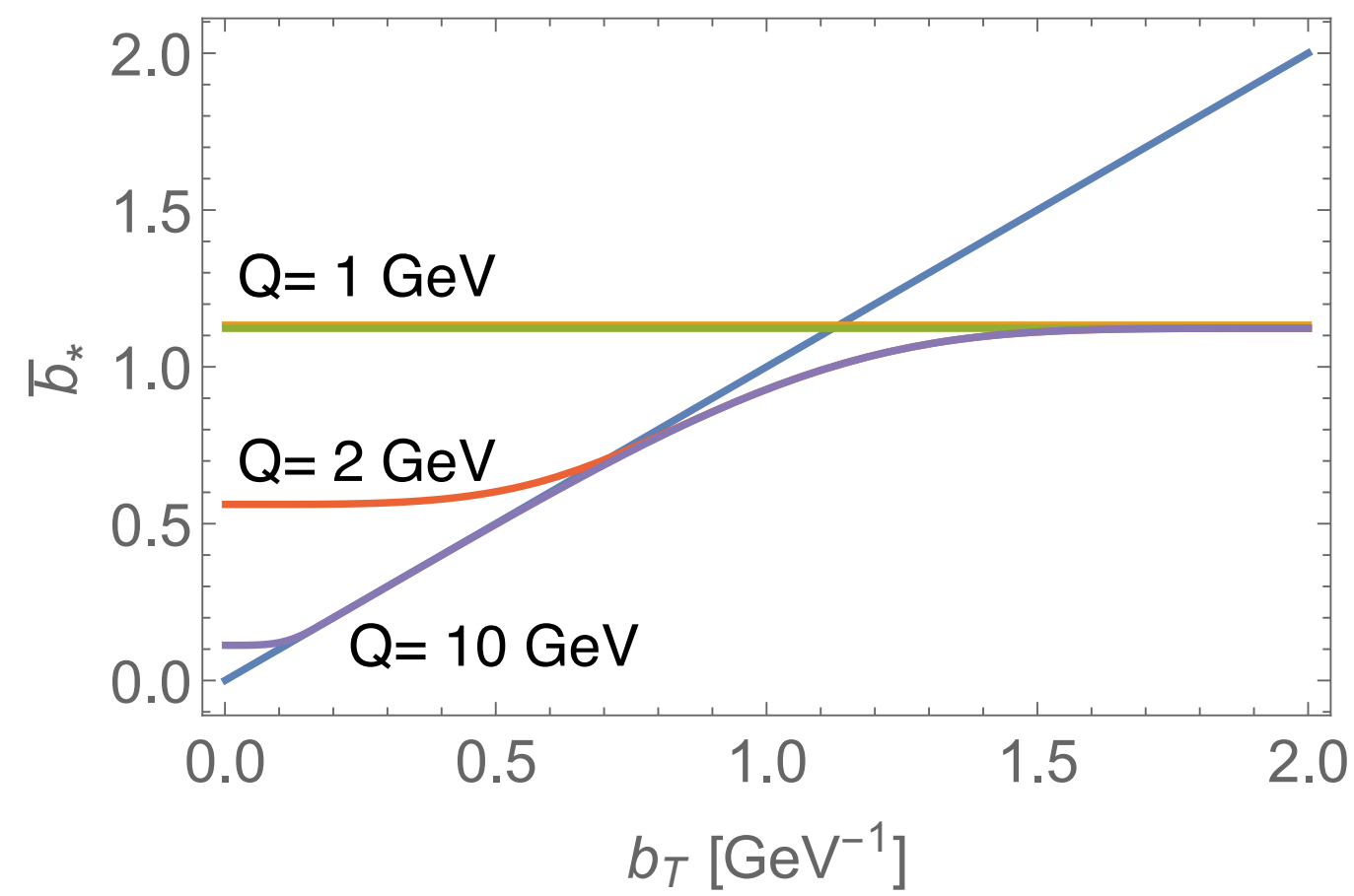
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$$\mu_b = 2e^{-\gamma_E} / b_*$$

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$$b_{\max} = 2e^{-\gamma_E}$$

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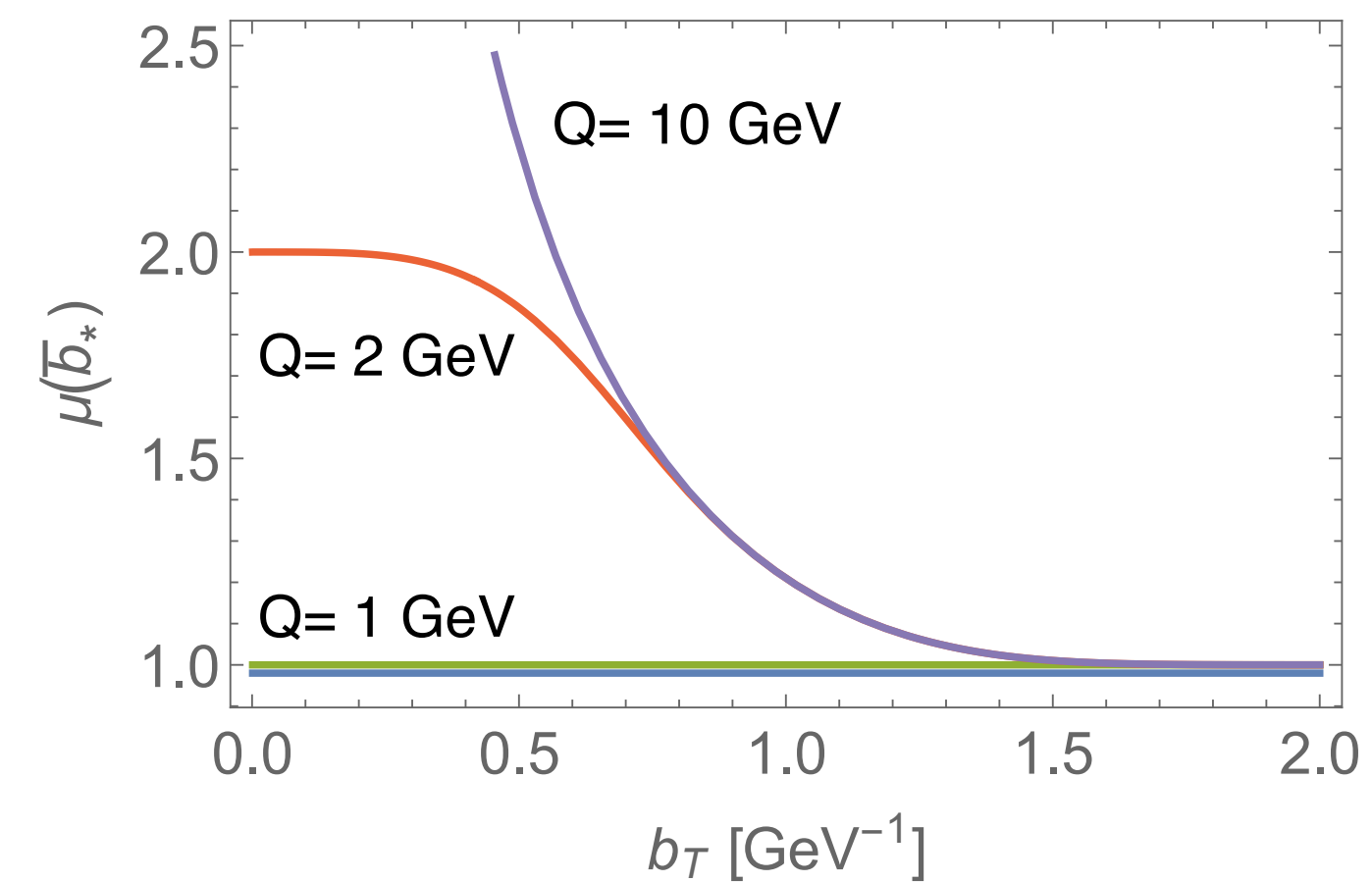
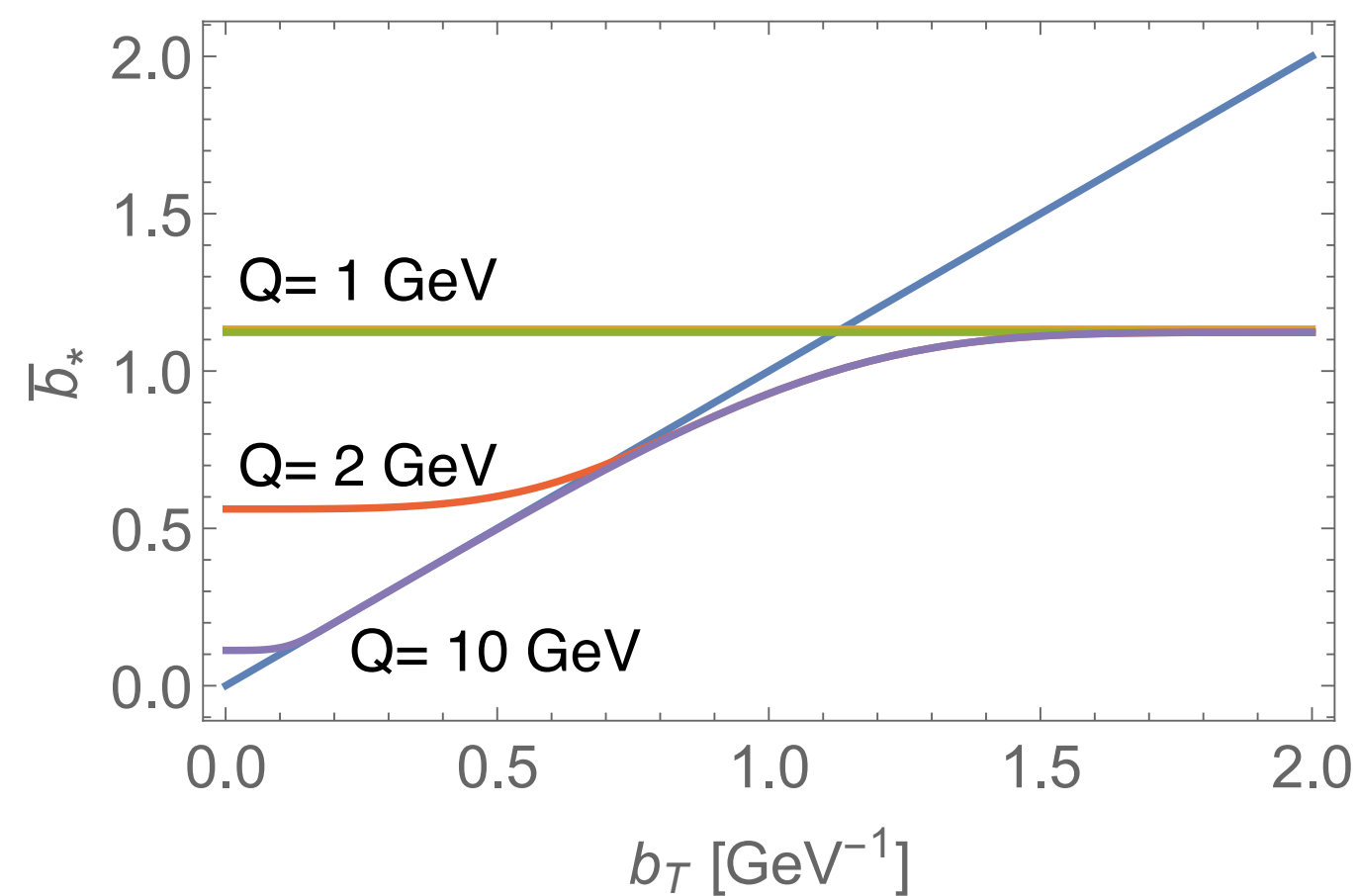


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$$b_{\max} = 2e^{-\gamma_E}$$

$$b_{\min} = \frac{2e^{-\gamma_E}}{Q}$$



No significant effect at high Q , but large effect at low Q
(inhibits perturbative contribution)



**ENHANCEMENT
PREFACTOR** = $\frac{\left. \frac{d\sigma}{dx dz dQ^2} \right|_{\text{nonmix.}}}{\int \text{TMD } d^2 P_{hT}}$



**ENHANCEMENT
PREFACTOR** = $\frac{d\sigma}{dx dz dQ^2} \Big|_{\text{nonmix.}}}{\int \text{TMD } d^2P_{hT}}$

The prefactor is independent of the fitting parameters



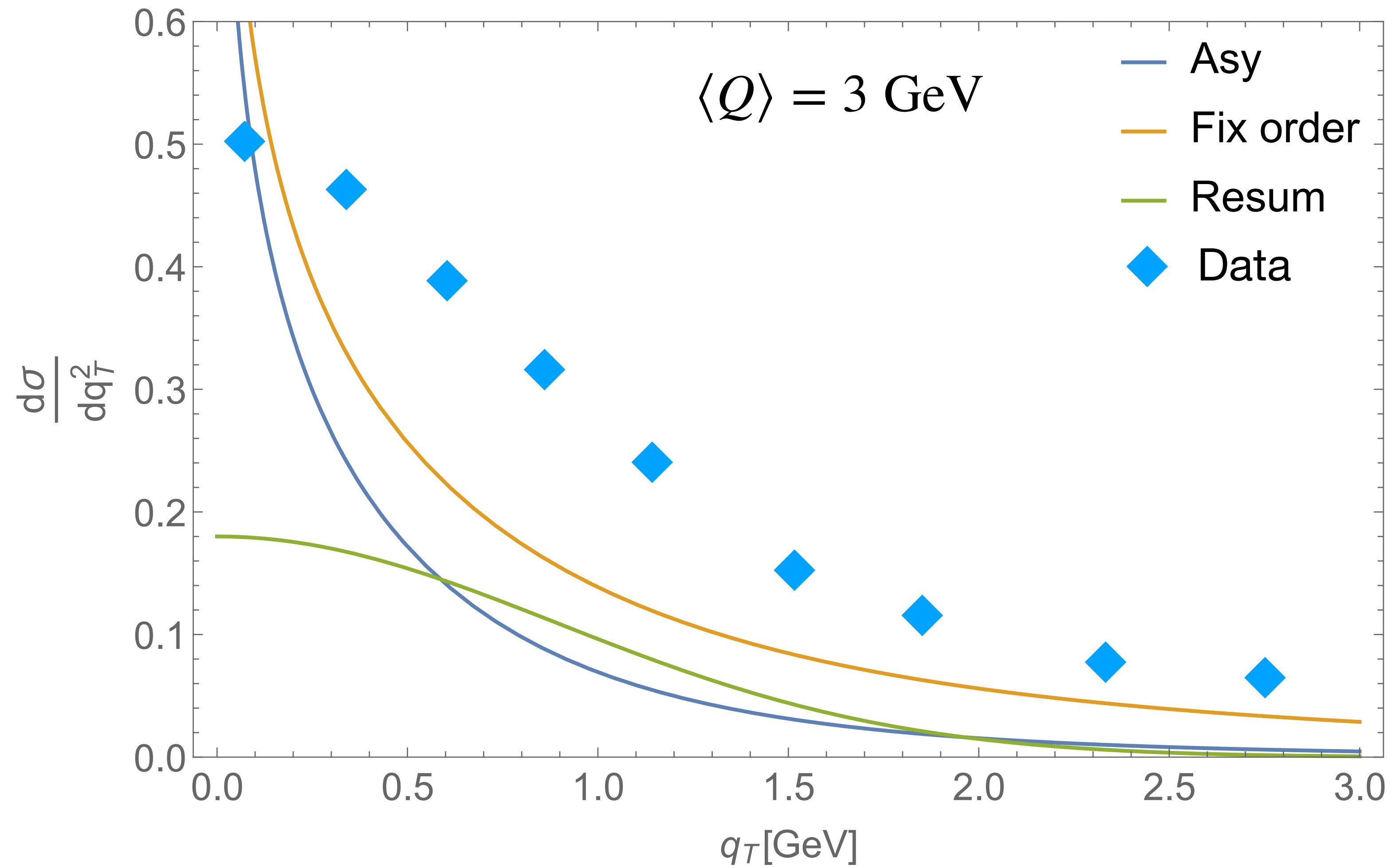
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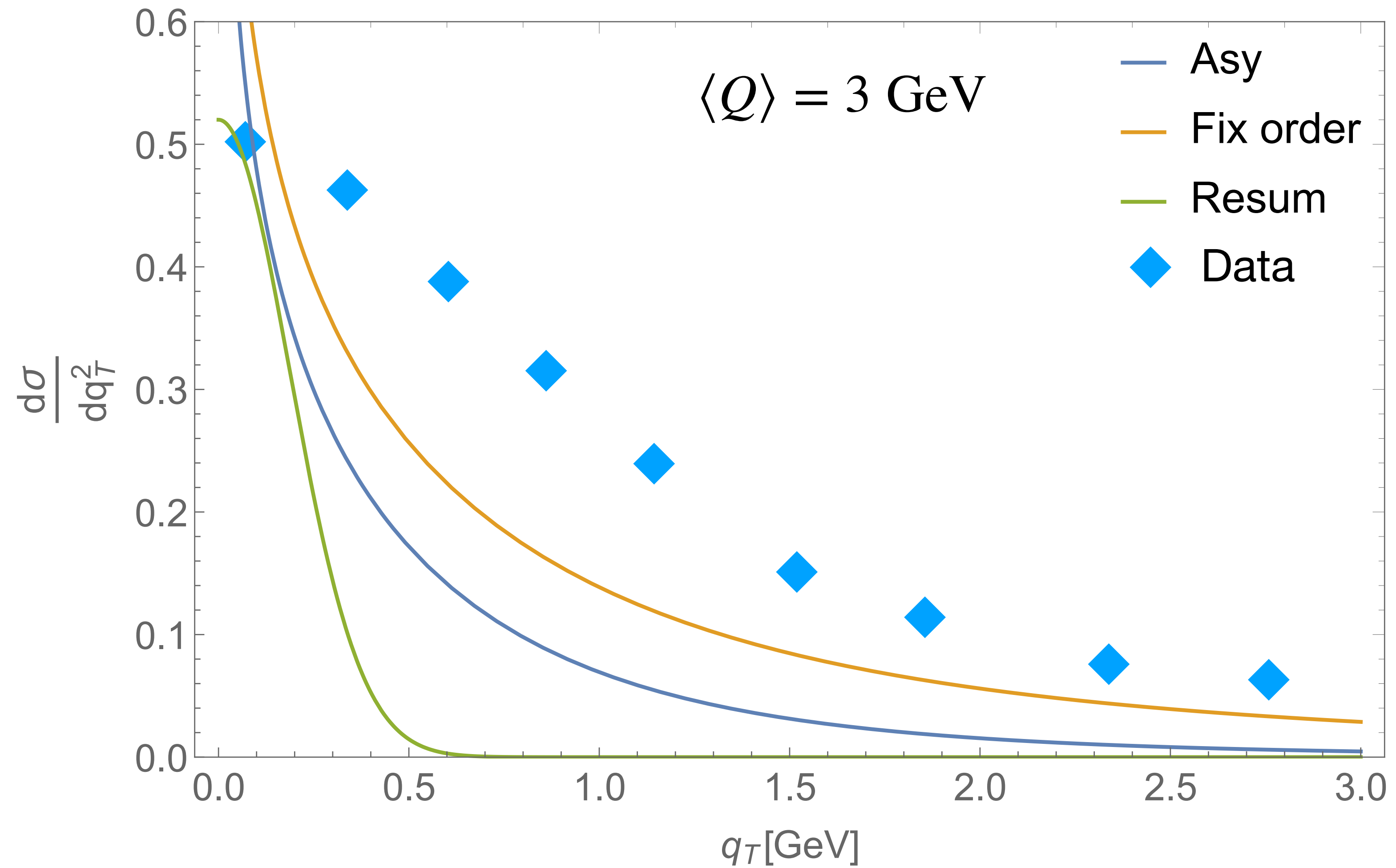
**Higher-order corrections decrease the
role of the TMD region.
We need to enhance it with a prefactor.**

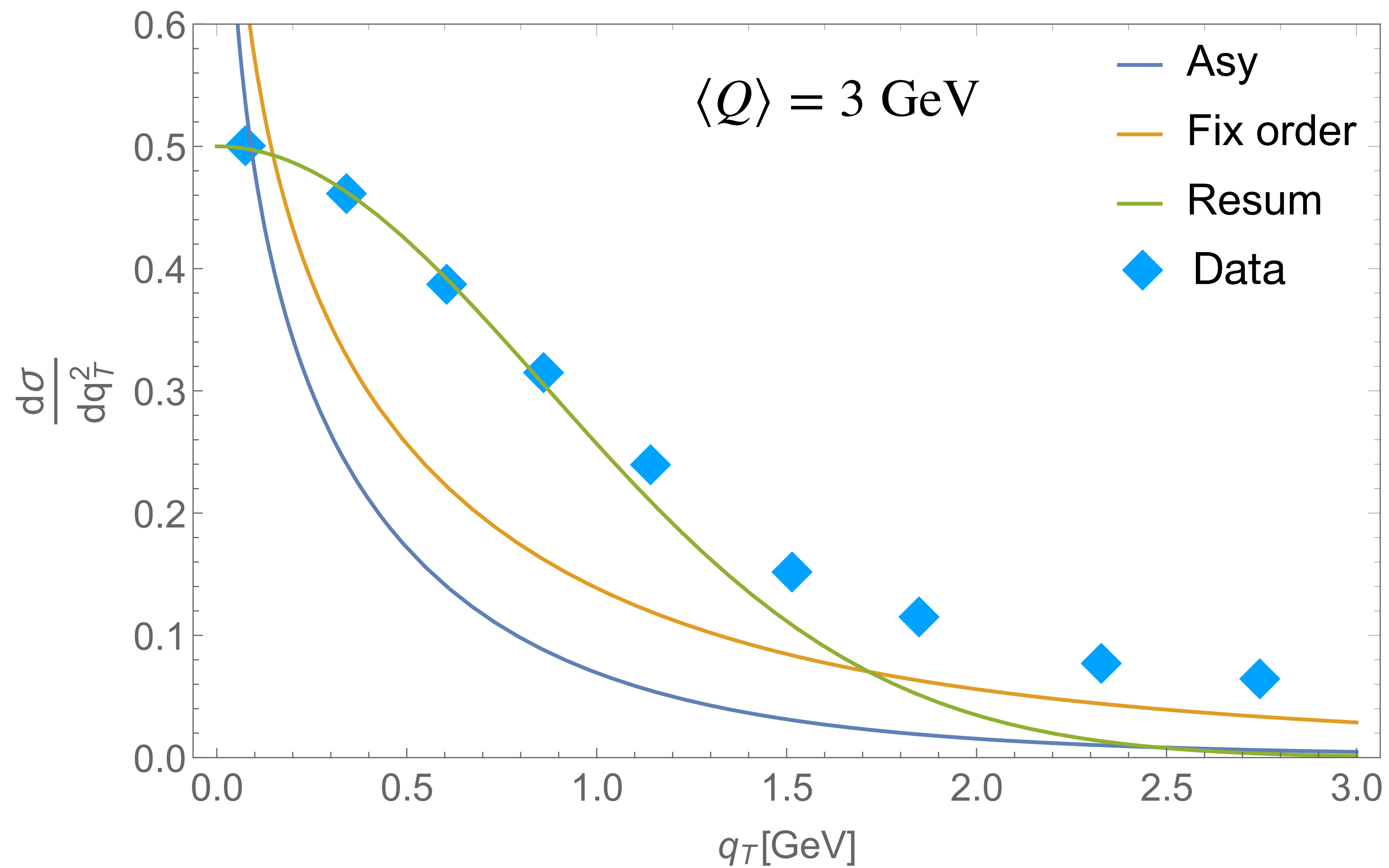
$$\frac{d\sigma^h}{dx dQ^2 dz} \Big|_{O(\alpha_s^1)} = \sigma_0 \sum_{ff'} \frac{e_f^2}{z^2} (\delta_{f'f} + \delta_{f'g}) \frac{\alpha_s}{\pi} \left\{ \left[D_1^{h/f'} \otimes C_1^{f'f} \otimes f_1^{f/N} \right] (x, z, Q) \right. \\ \left. + \frac{1-y}{1+(1-y)^2} \left[D_1^{h/f'} \otimes C_L^{f'f} \otimes f_1^{f/N} \right] (x, z, Q) \right\},$$

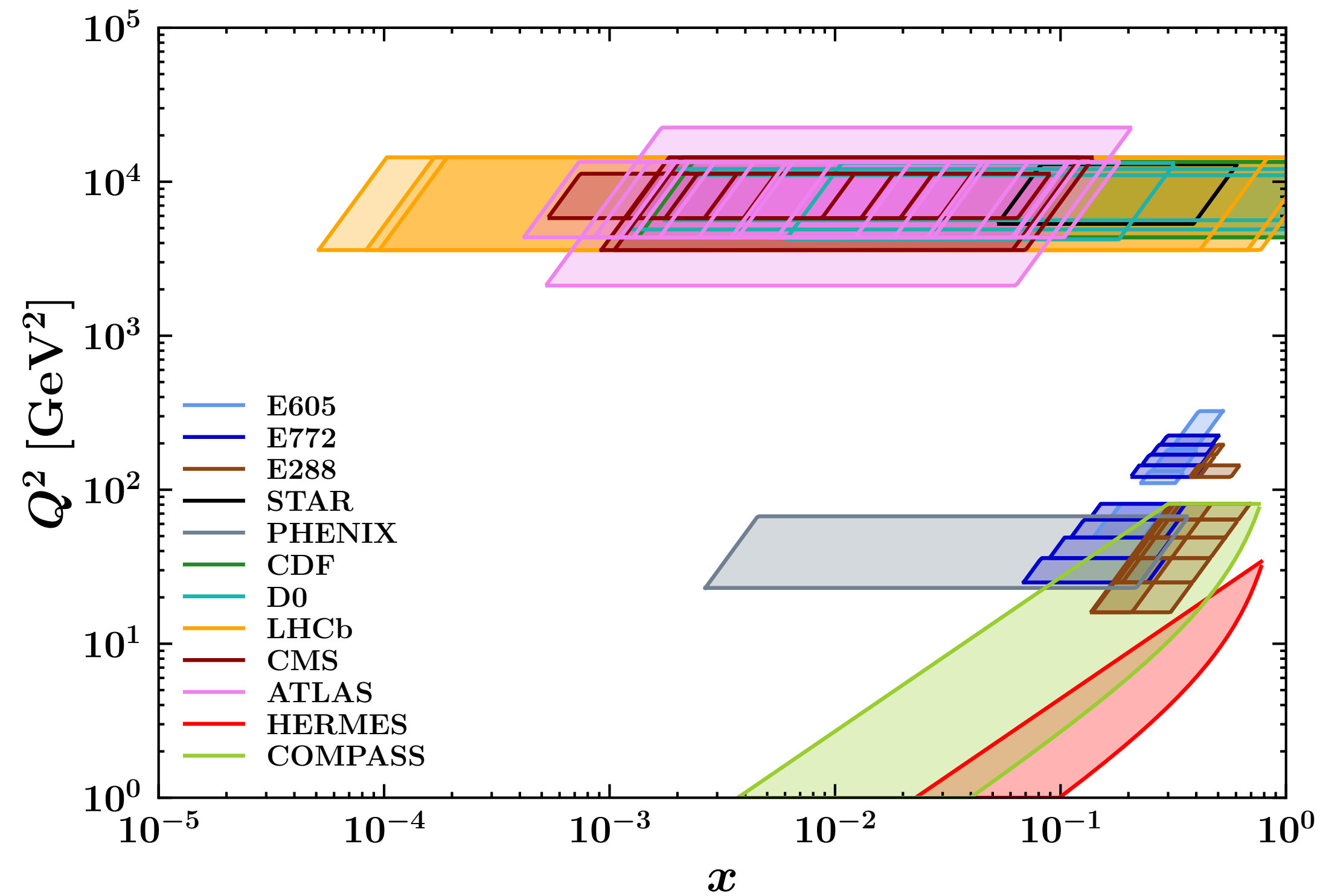
$$C_1^{qq} = \frac{C_F}{2} \left\{ -8\delta(1-x)\delta(1-z) \right. \\ \left. + \delta(1-x) \left[P_{qq}(z) \ln \frac{Q^2}{\mu_F^2} + L_1(z) + L_2(z) + (1-z) \right] \right. \\ \left. + \delta(1-z) \left[P_{qq}(x) \ln \frac{Q^2}{\mu^2} + L_1(x) - L_2(x) + (1-x) \right] \right. \\ \left. + 2 \frac{1}{(1-x)_+} \frac{1}{(1-z)_+} - \frac{1+z}{(1-x)_+} \frac{1+x}{(1-z)_+} + 2(1+xz) \right\},$$



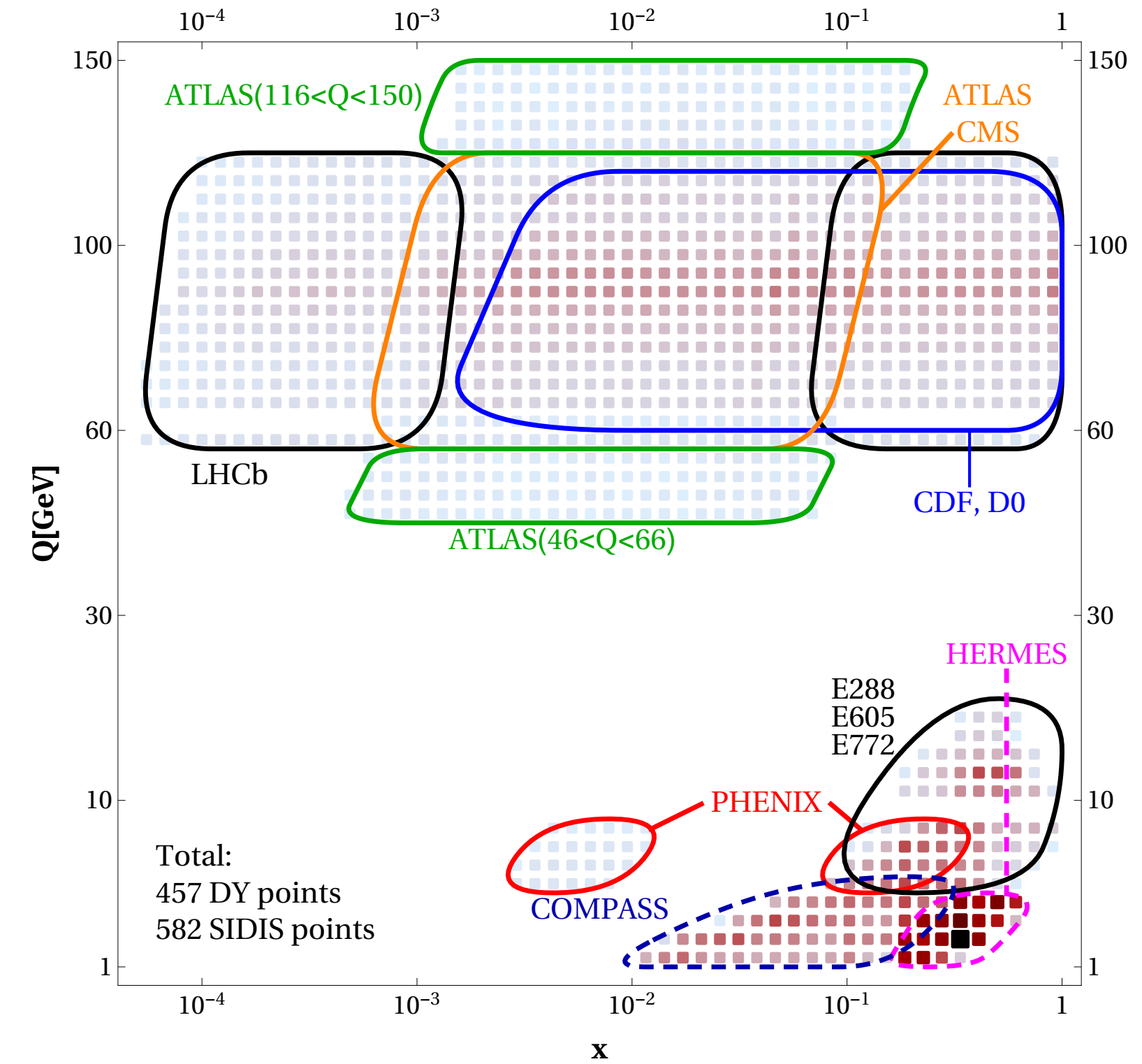
SOLUTION 1: RESTRICT TMD REGION



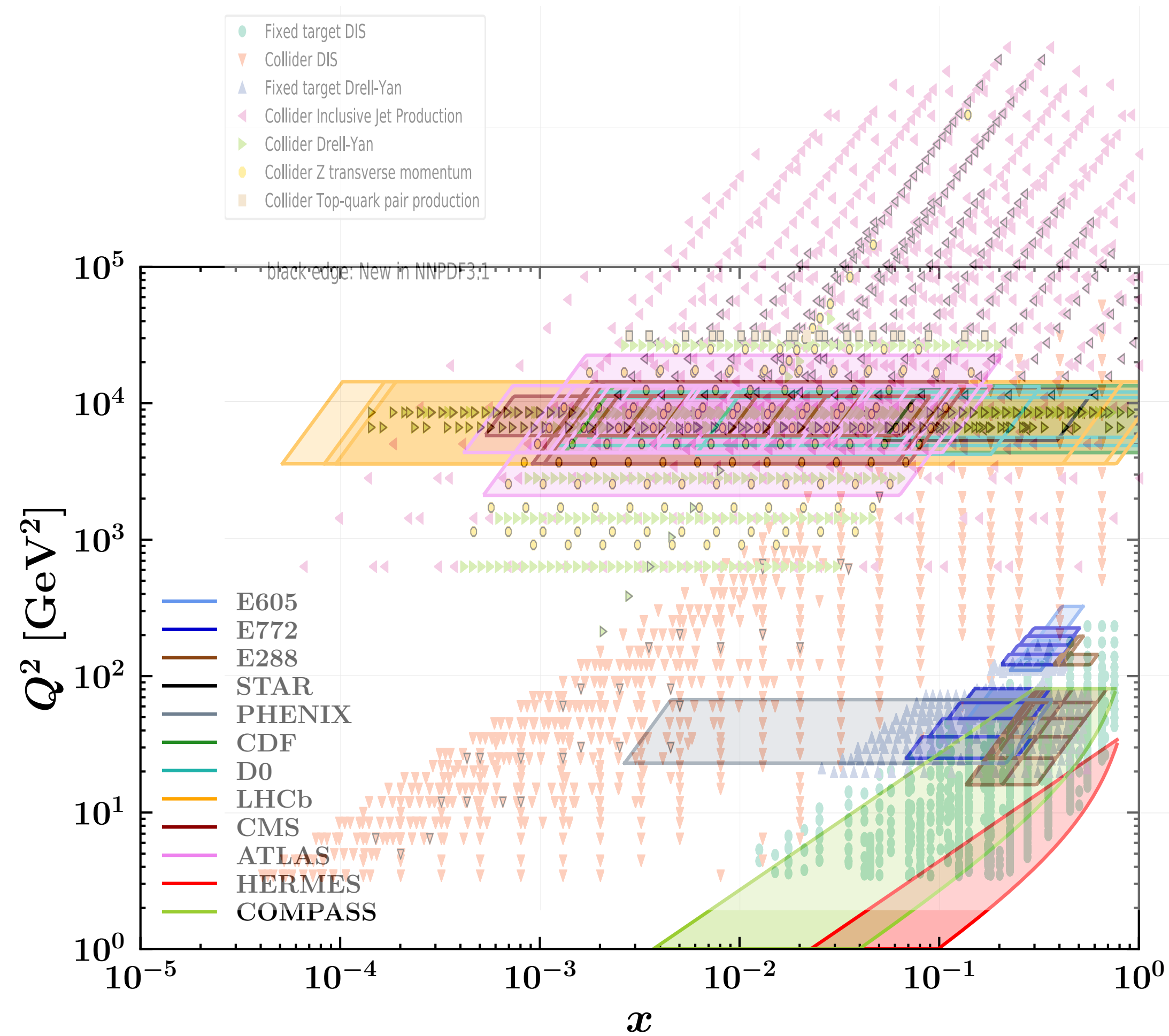




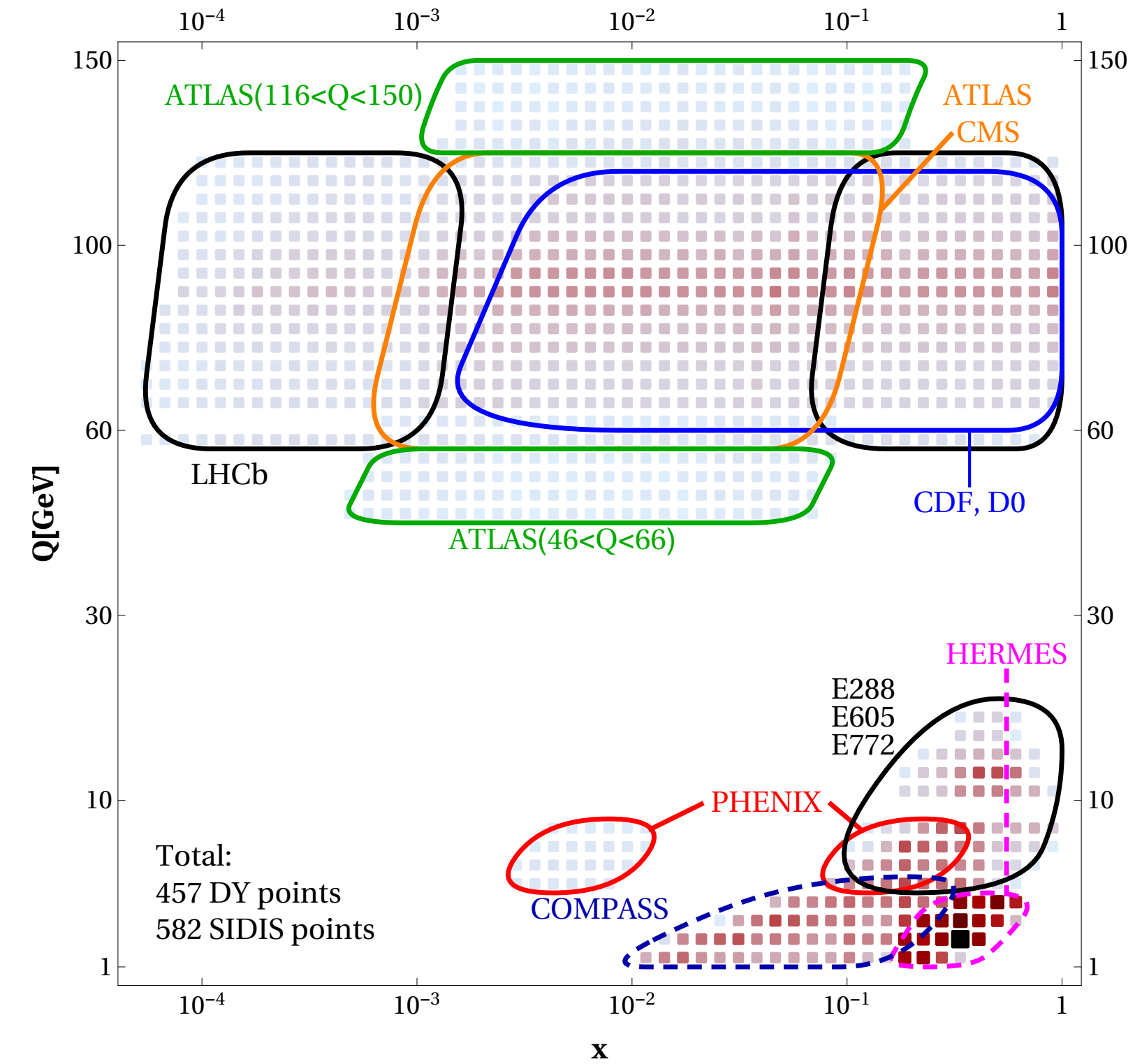
[MAP Collaboration](#)
[Bacchetta, Bertone, Bissolotti, Bozzi, Cerutti,](#)
[Piacenza, Radici, Signori, arXiv:2206.07598](#)



[Scimemi, Vladimirov,](#)
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[arXiv:1912.06532](#)

Data set	N_{dat}	χ_D^2/N_{dat}	$\chi_\lambda^2/N_{\text{dat}}$	χ_0^2/N_{dat}
Tevatron total	71	0.87	0.06	0.93
LHCb total	21	1.15	0.3	1.45
ATLAS total	72	4.56	0.48	5.05
CMS total	78	0.53	0.02	0.55
PHENIX 200	2	2.21	0.88	3.08
STAR 510	7	1.05	0.10	1.15
DY collider total	251	1.86	0.2	2.06
DY fixed-target total	233	0.85	0.4	1.24
HERMES total	344	0.48	0.23	0.71
COMPASS total	1203	0.62	0.3	0.92
SIDIS total	1547	0.59	0.28	0.87
Total	2031	0.77	0.29	1.06

<https://github.com/MapCollaboration/NangaParbat>



☰ README.md



Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

Download

You can obtain NangaParbat directly from the github repository:

<https://github.com/MapCollaboration/NangaParbat>

For the last development branch you can clone the master code:

```
git clone git@github.com:MapCollaboration/NangaParbat.git
```

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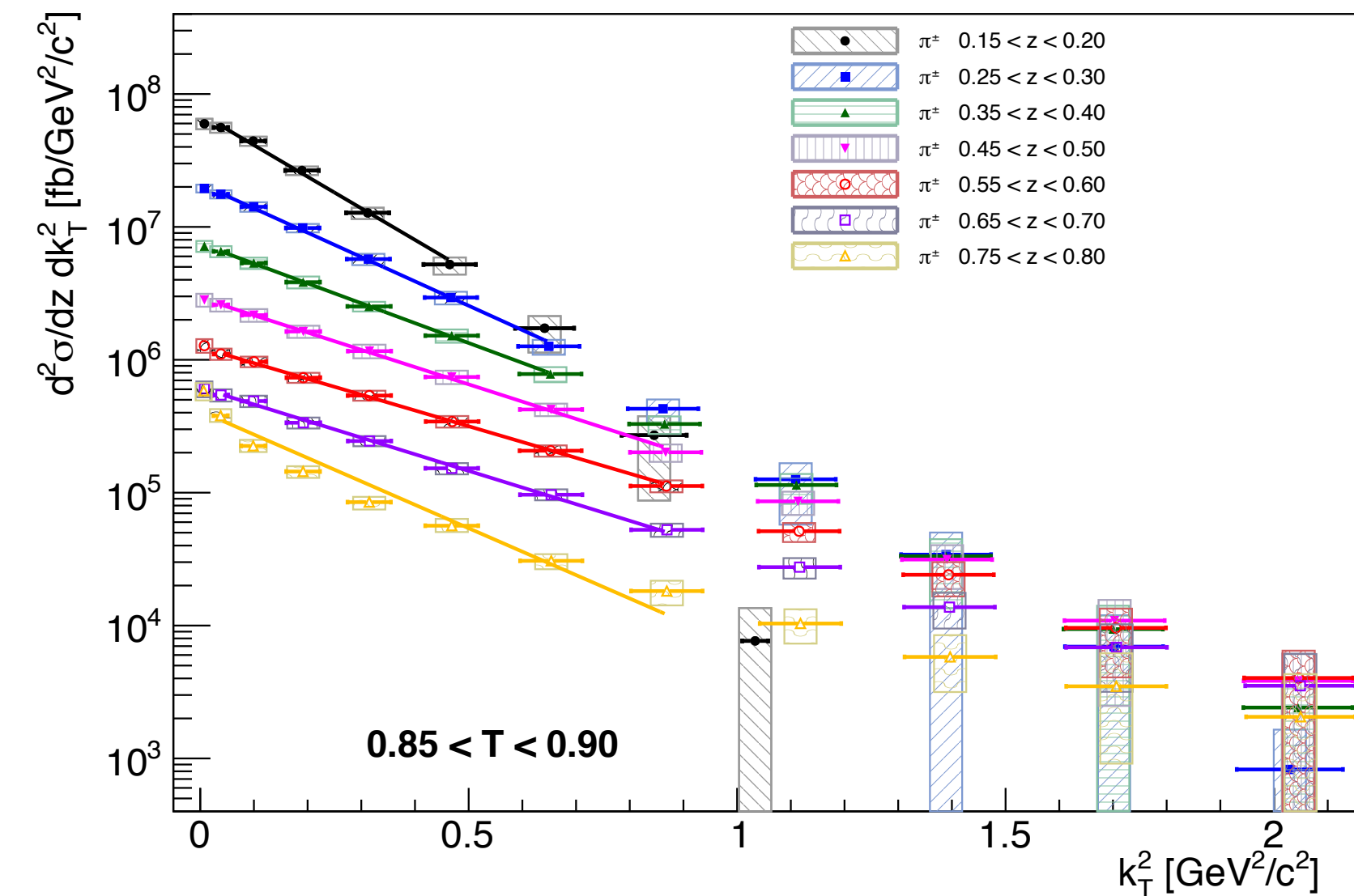
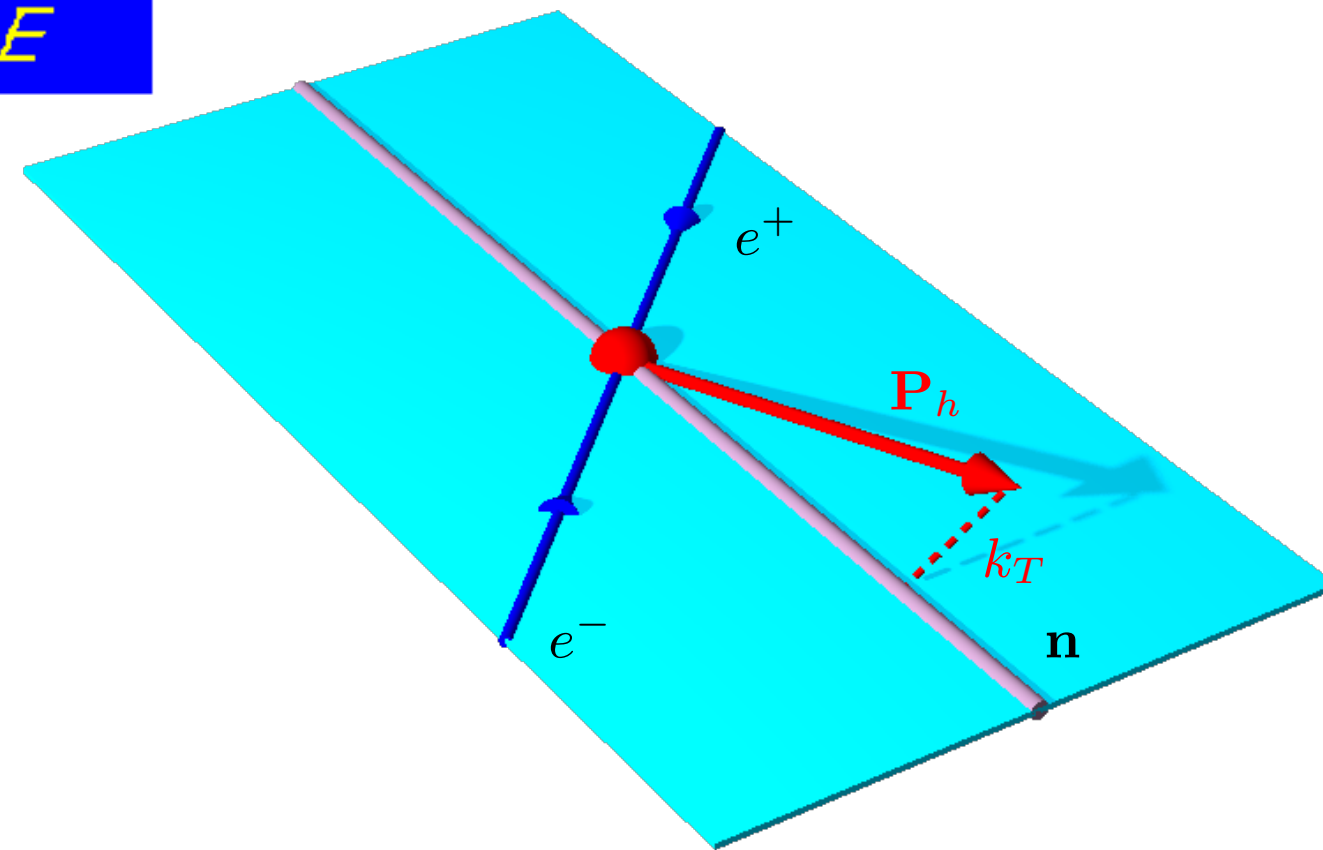
Also:

ARTEMIDE

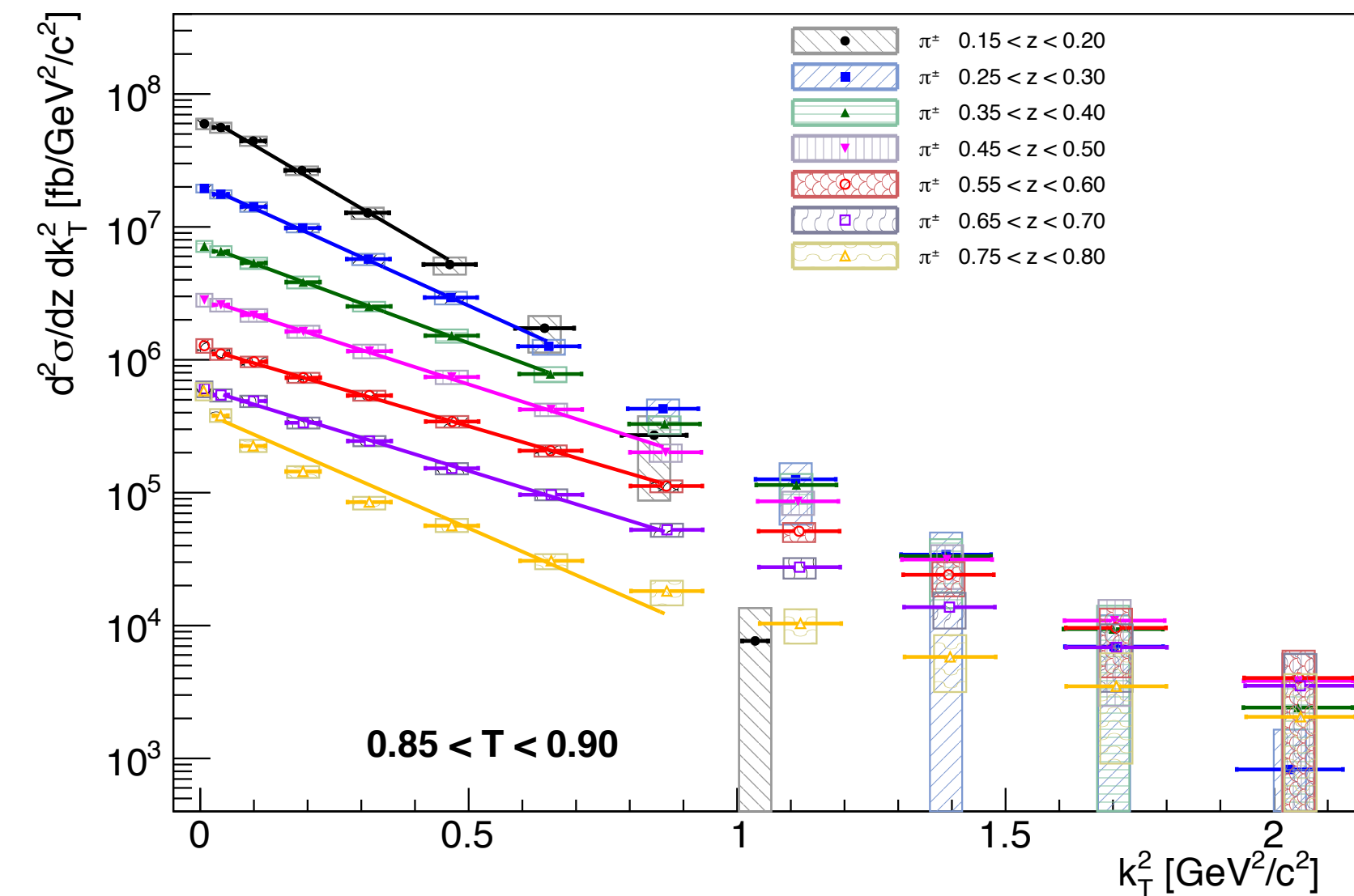
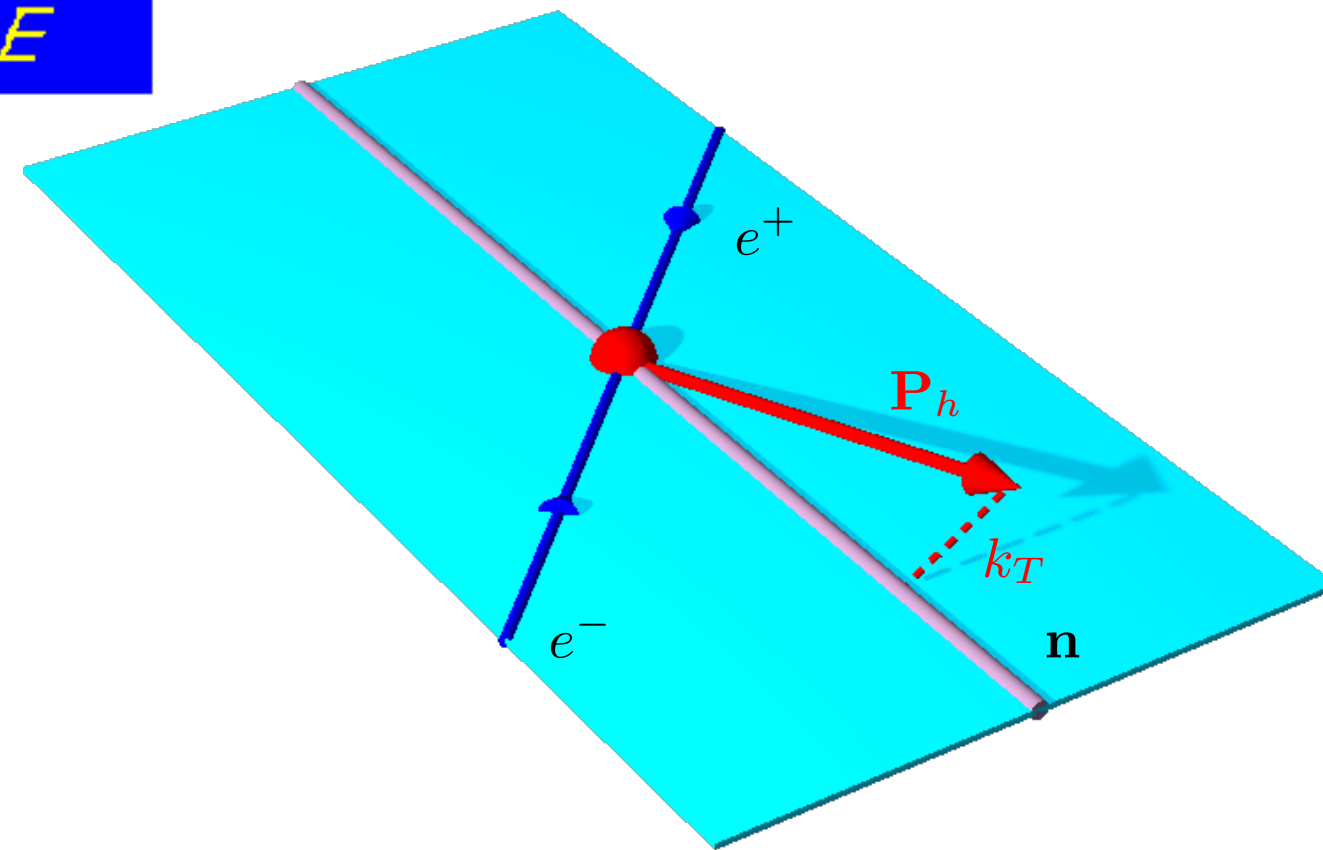
<https://teorica.fis.ucm.es/artemide/>

TMDLIB

<https://tmdlib.hepforge.org/>



First direct measurement of TMD effects in fragmentation functions
Makes use of thrust axis: the formalism should take it into account

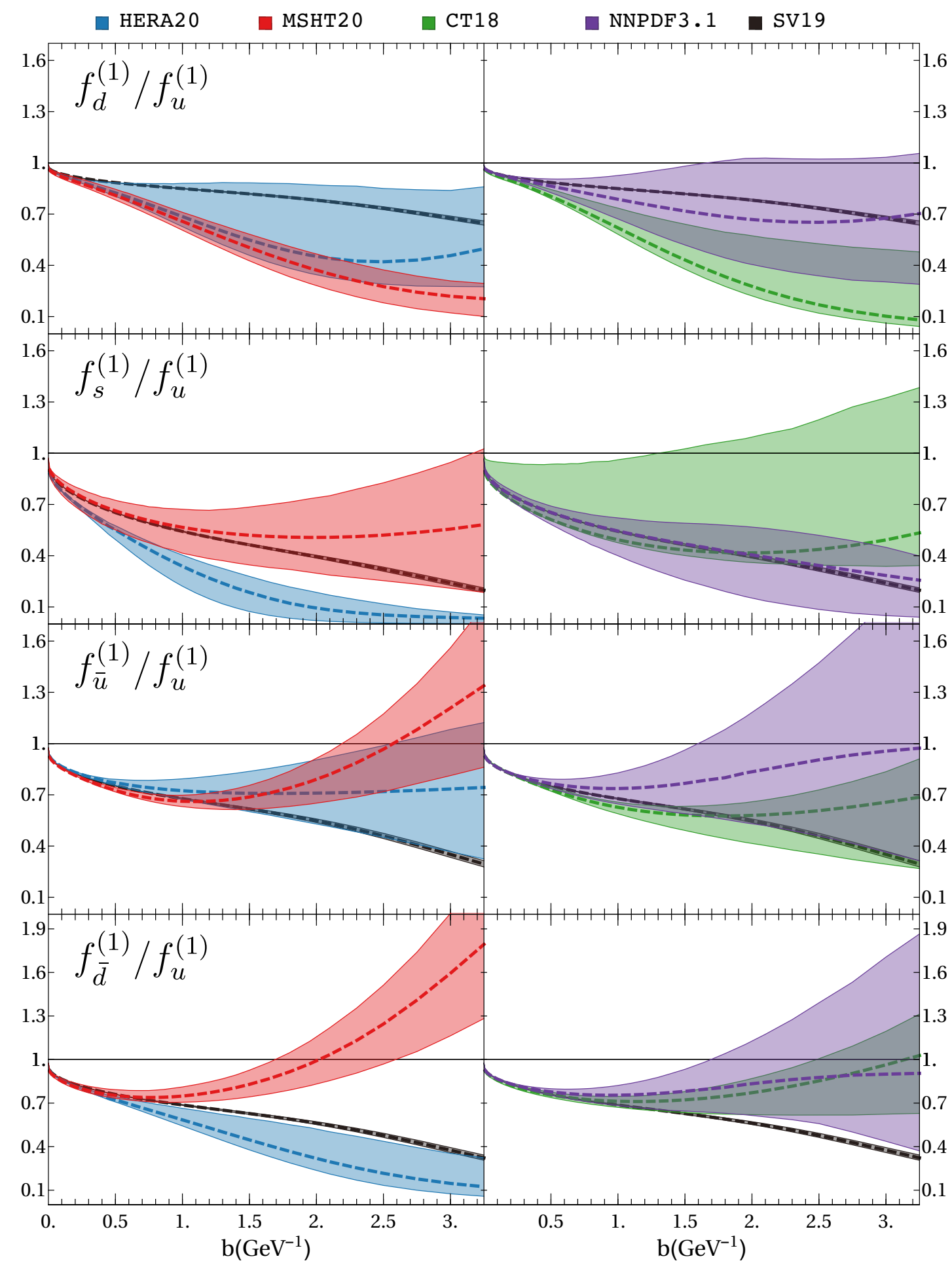


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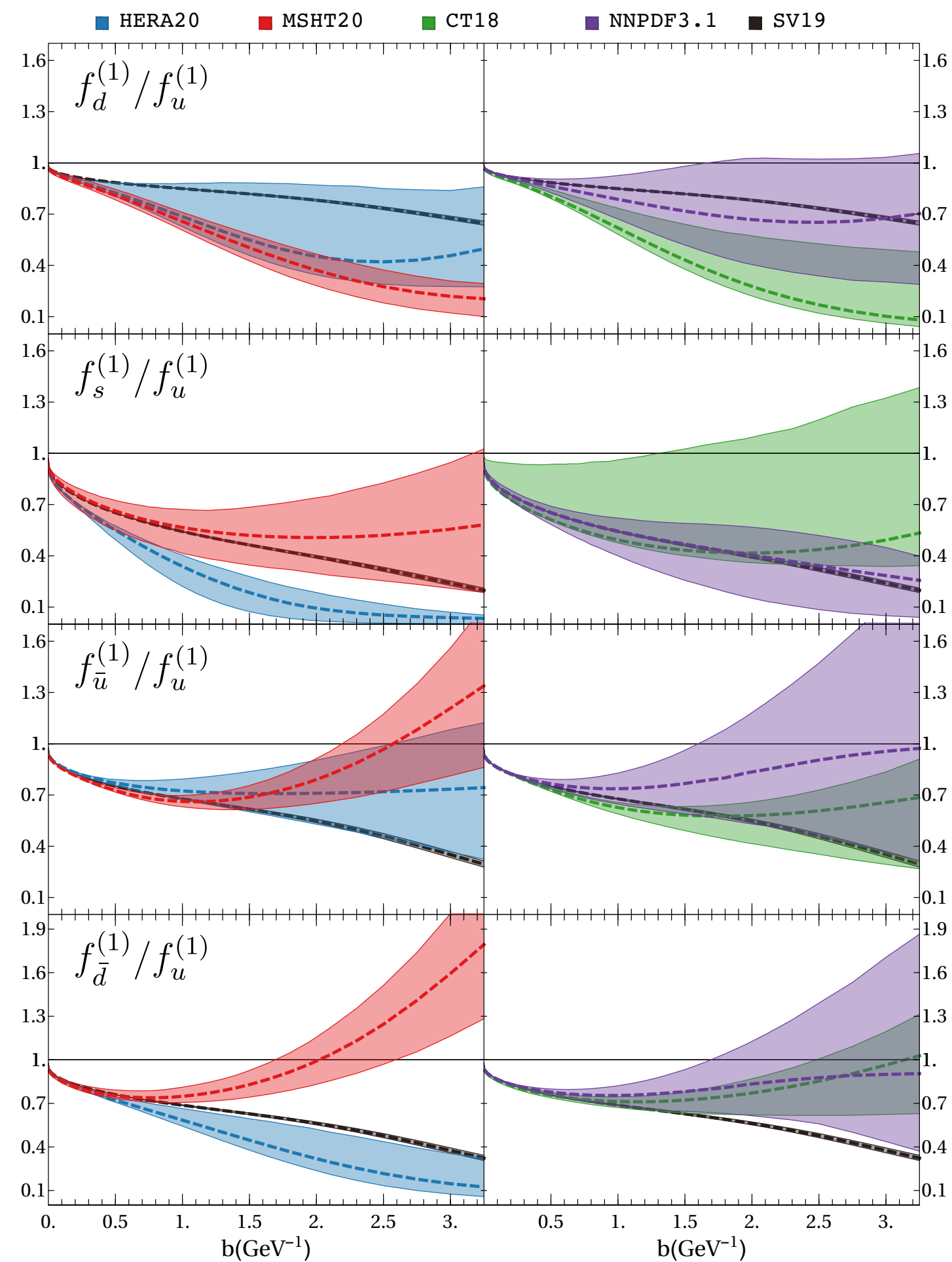
See <https://arxiv.org/abs/2206.08876>

RECENT STUDY WITH FLAVOR DEPENDENCE

Bury, Hautmann, Leal-Gomez, Scimemi, Vladimirov, Zurita, arxiv:2201.07114



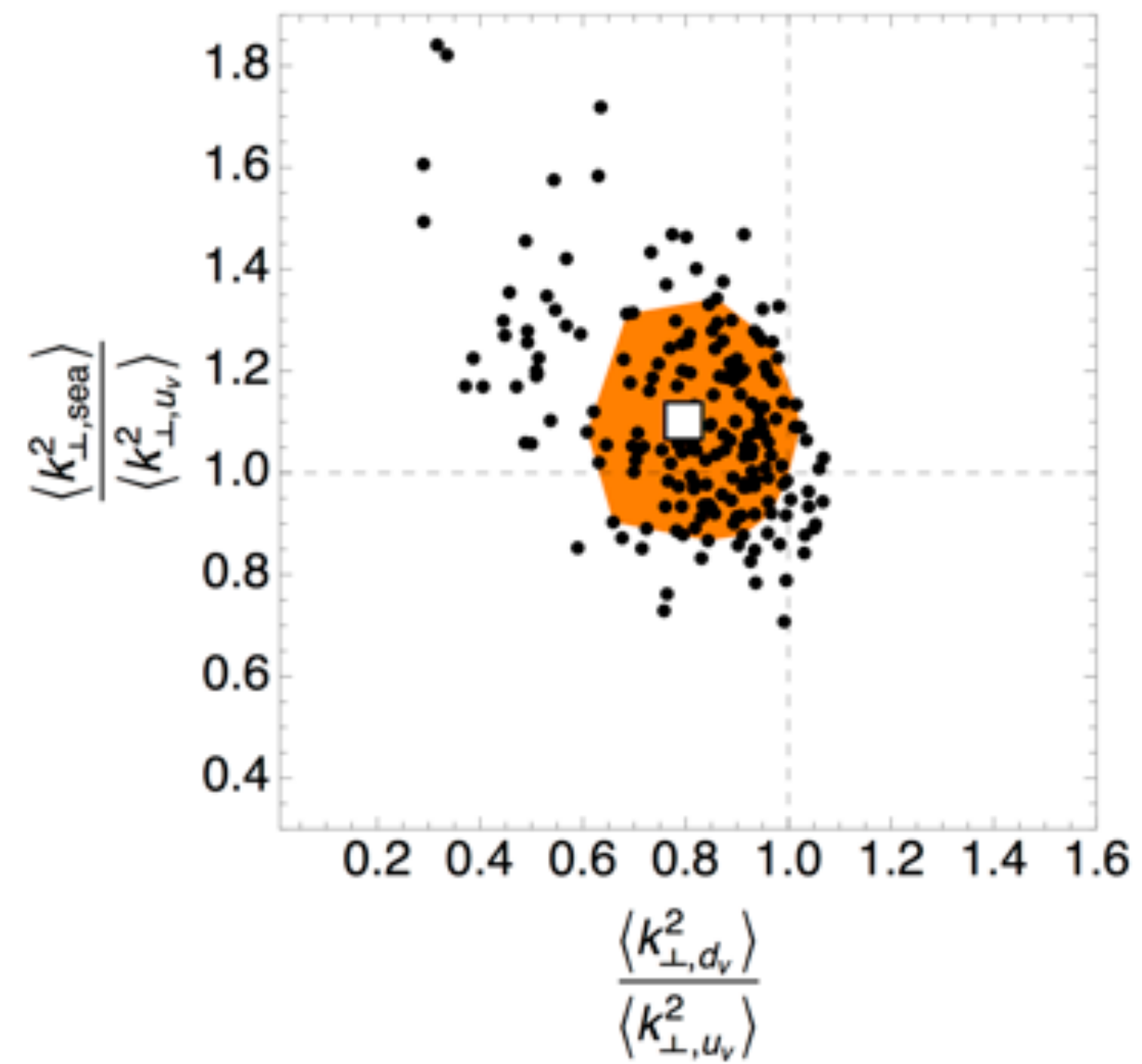
Bury, Hautmann, Leal-Gomez, Scimemi, Vladimirov, Zurita, arxiv:2201.07114



There seems to be a lot of room for flavor dependence. Different collinear PDFs lead to different results...

Signori, Bacchetta, Radici, Schnell JHEP 1311 (13)

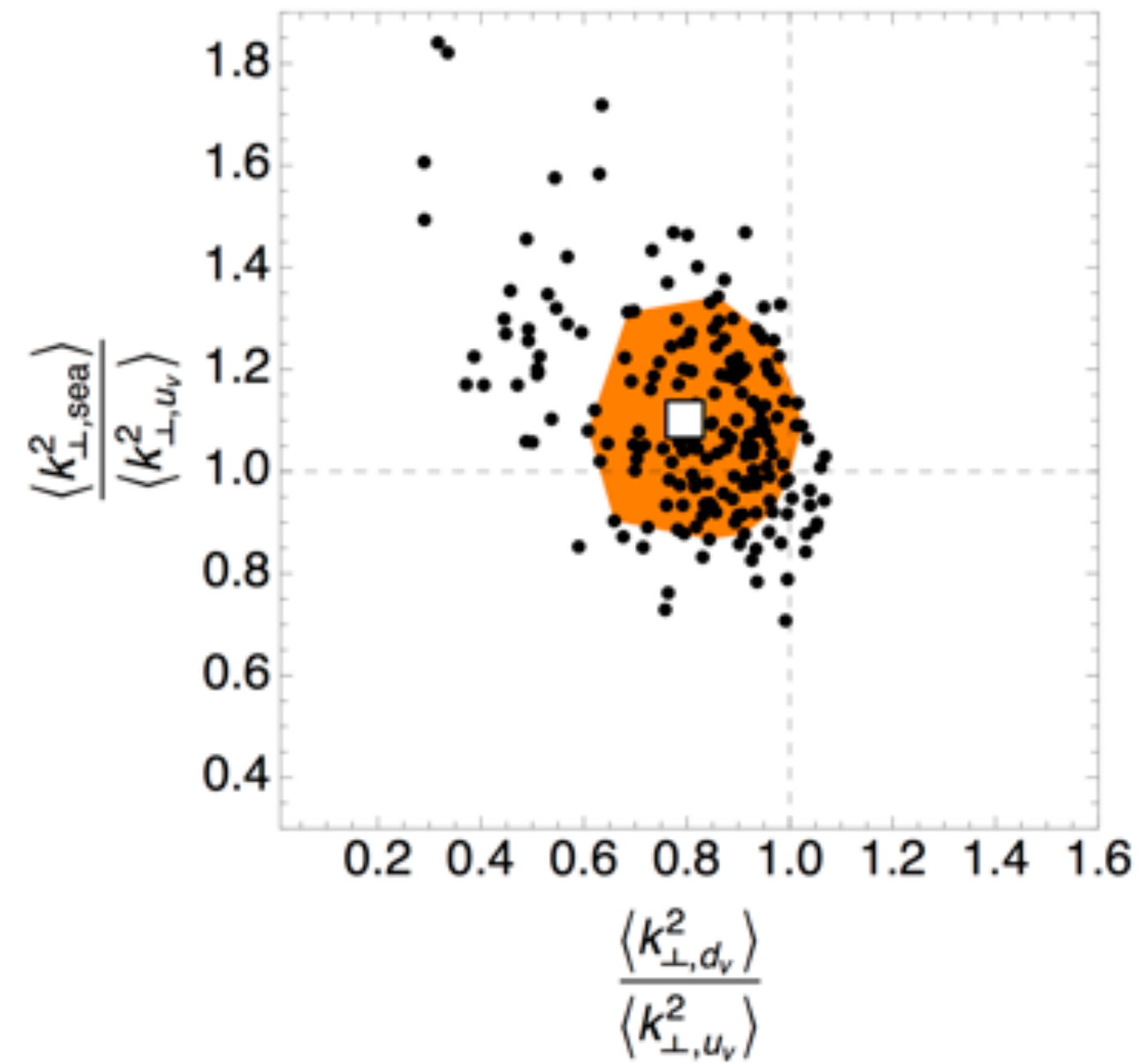
Ratio of width of sea /
width of up valence



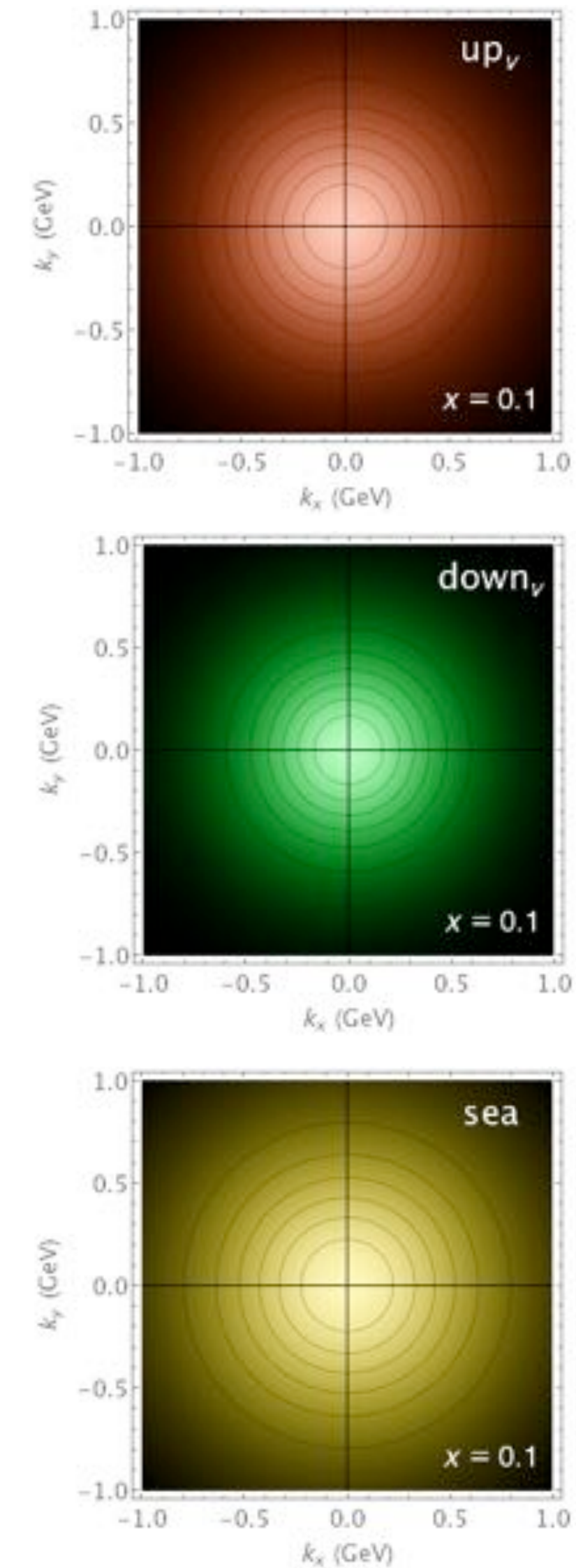
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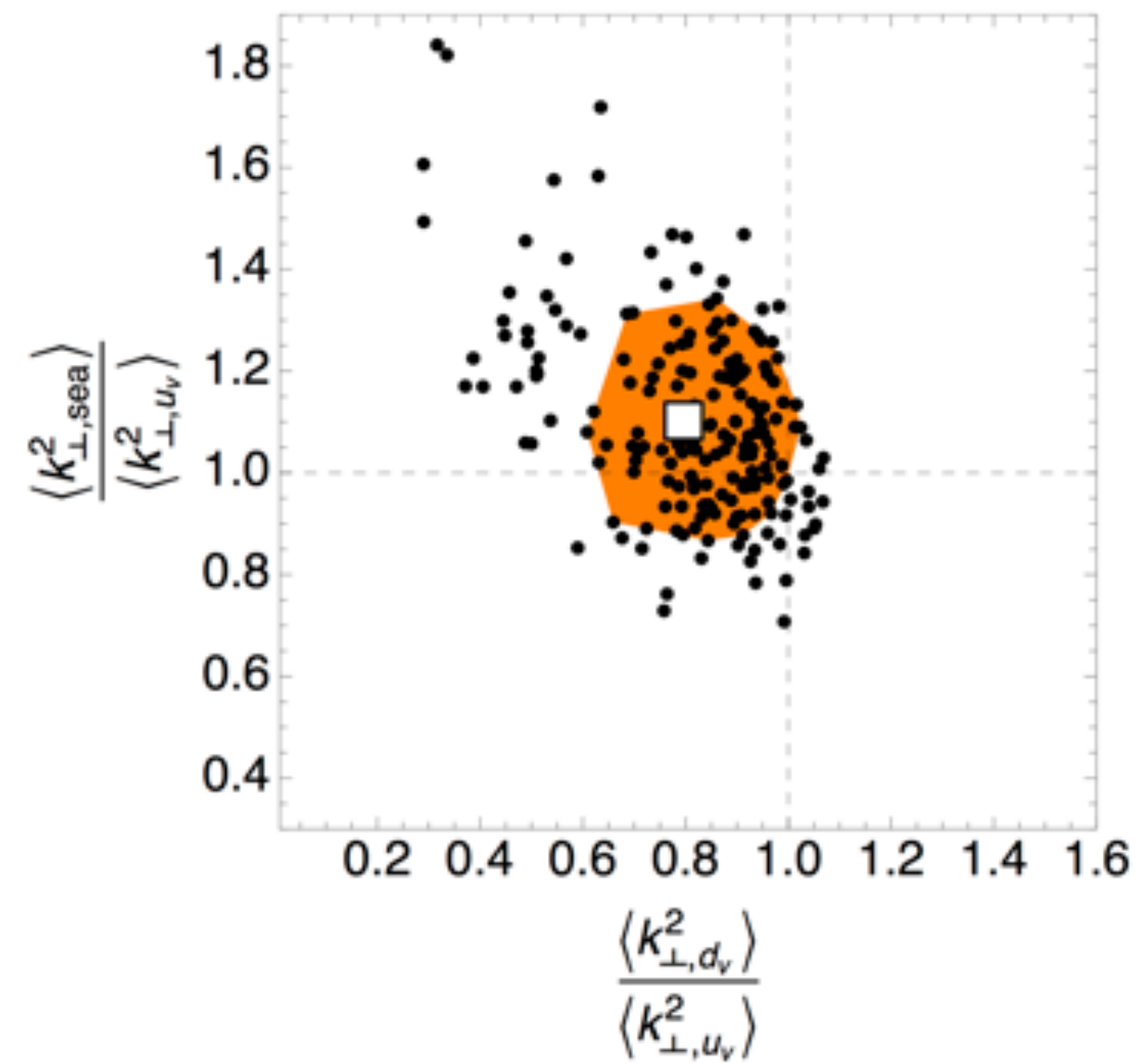


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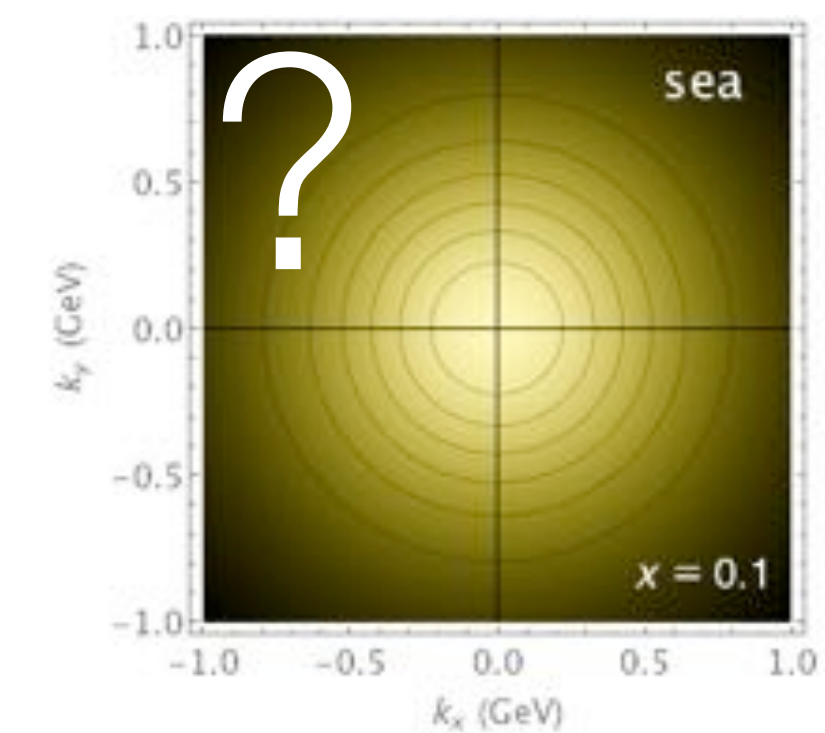
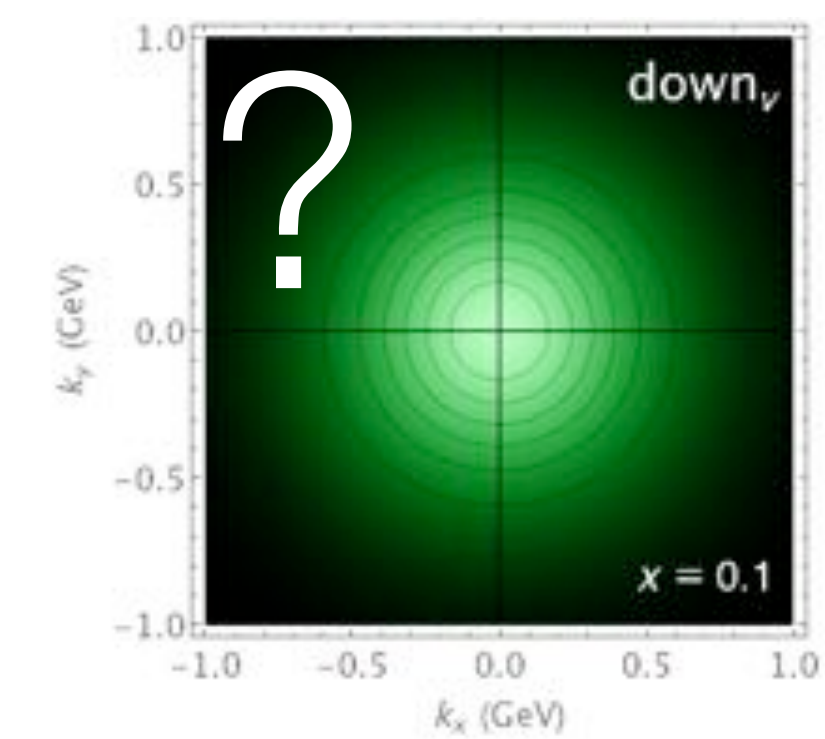
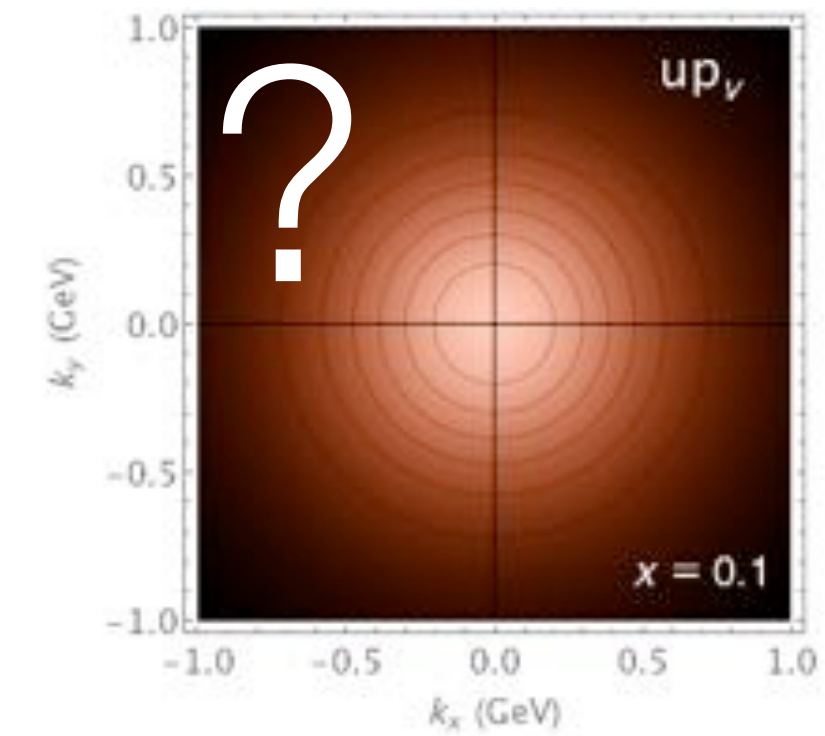


Signori, Bacchetta, Radici, Schnell JHEP 1311 (13)

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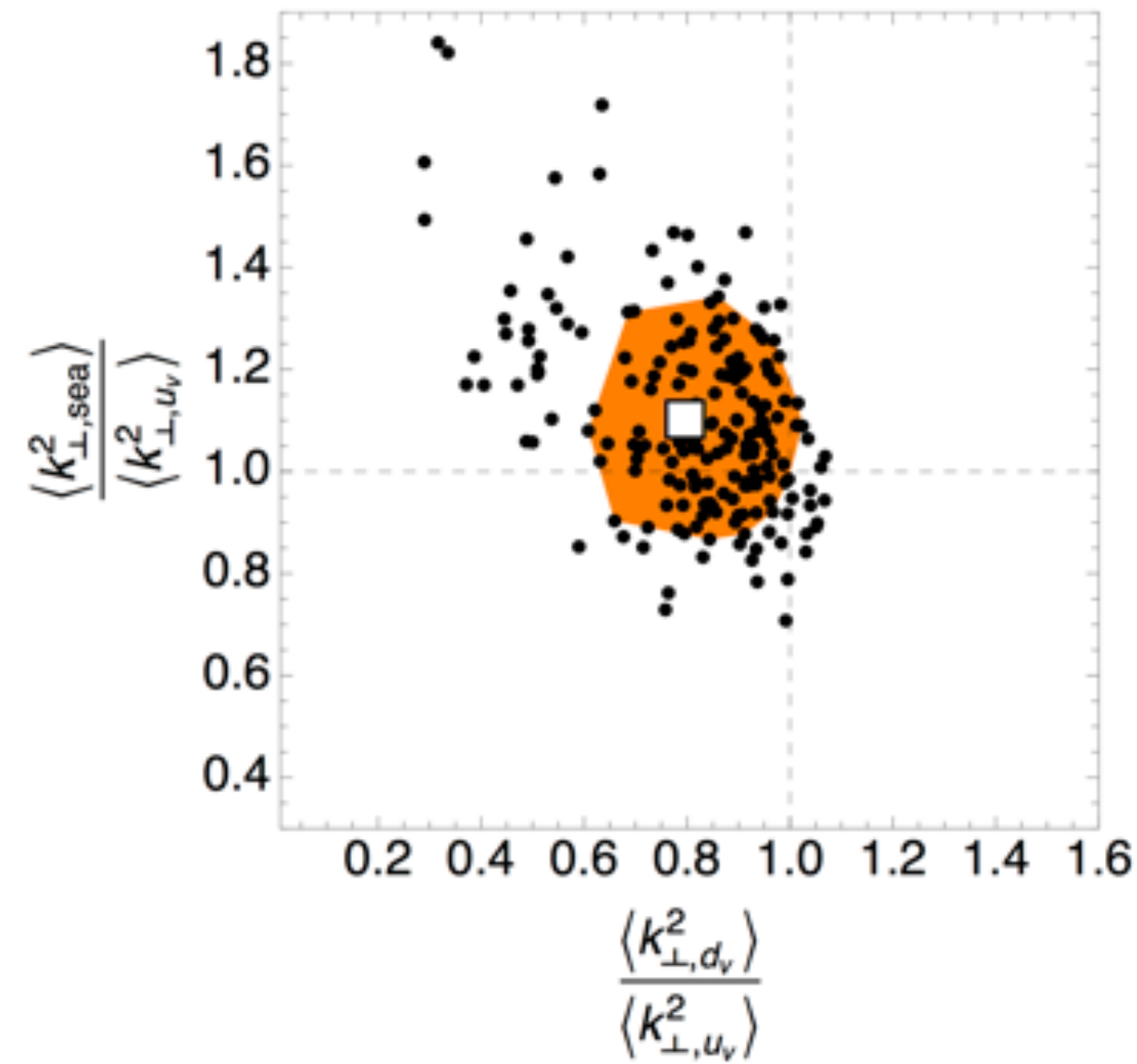


Ratio width of down valence/
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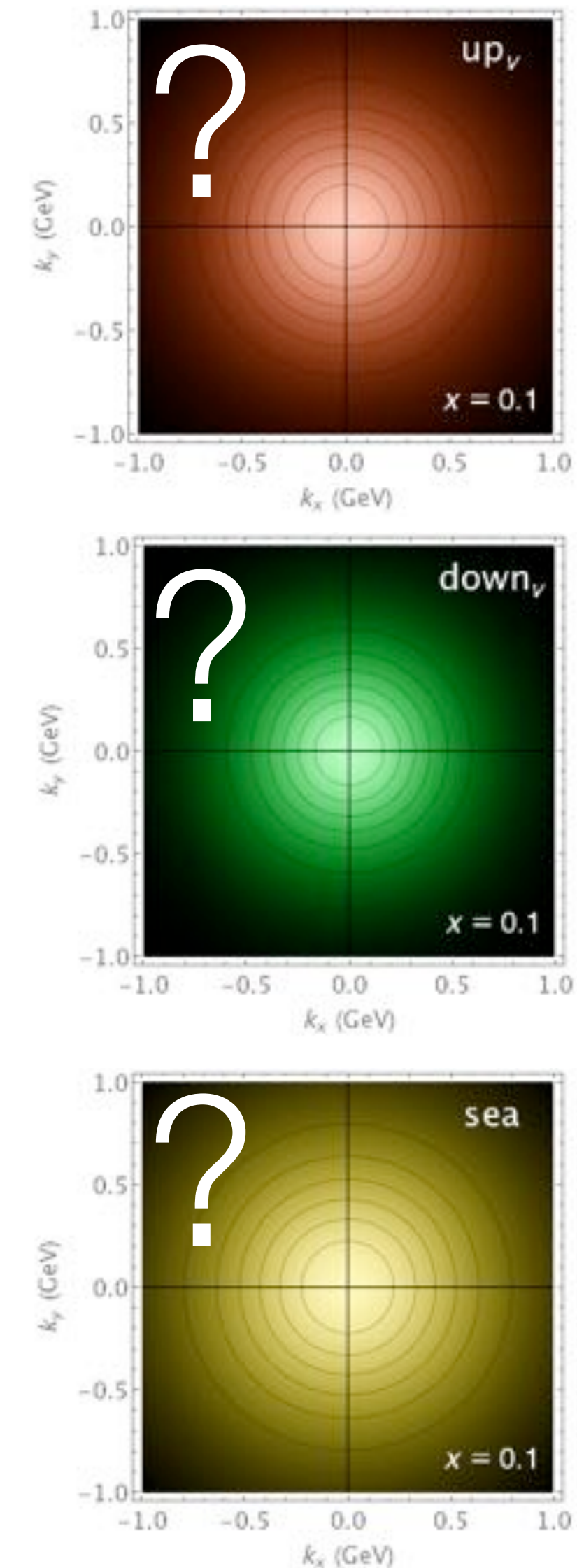
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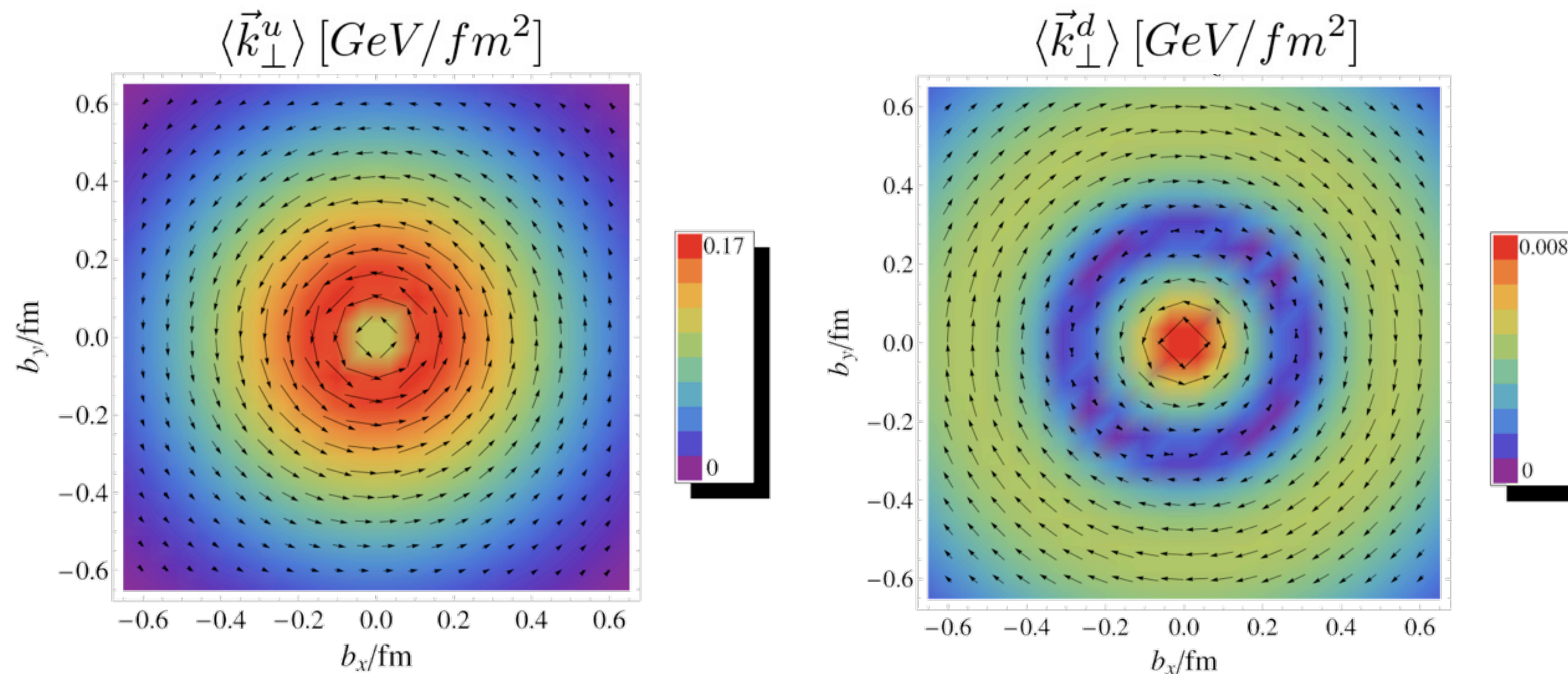
Ratio width of down valence/
width of up valence

There is room for flavour dependence,
but we don't control it well

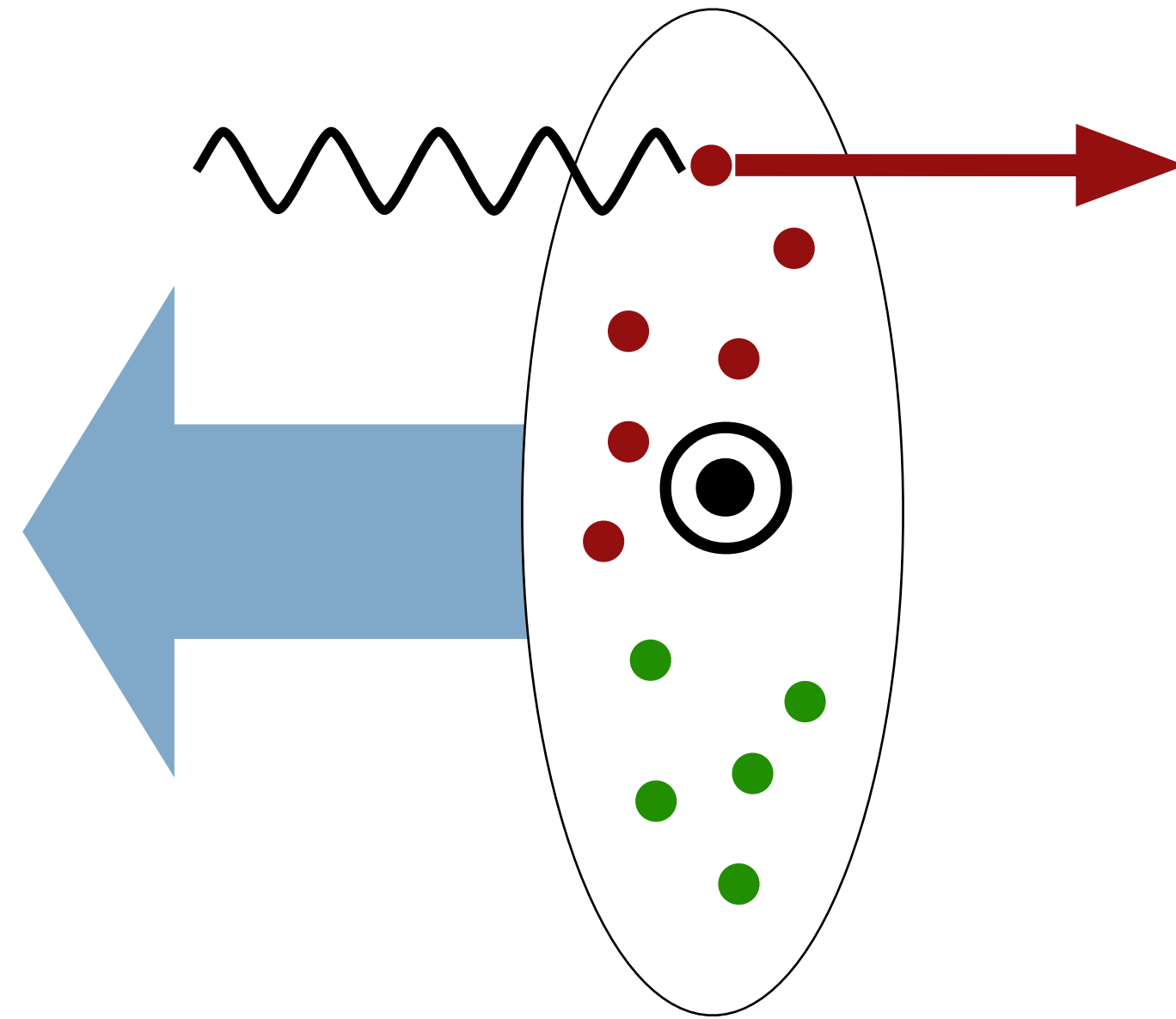


Only way to provide direct access to partonic orbital angular momentum

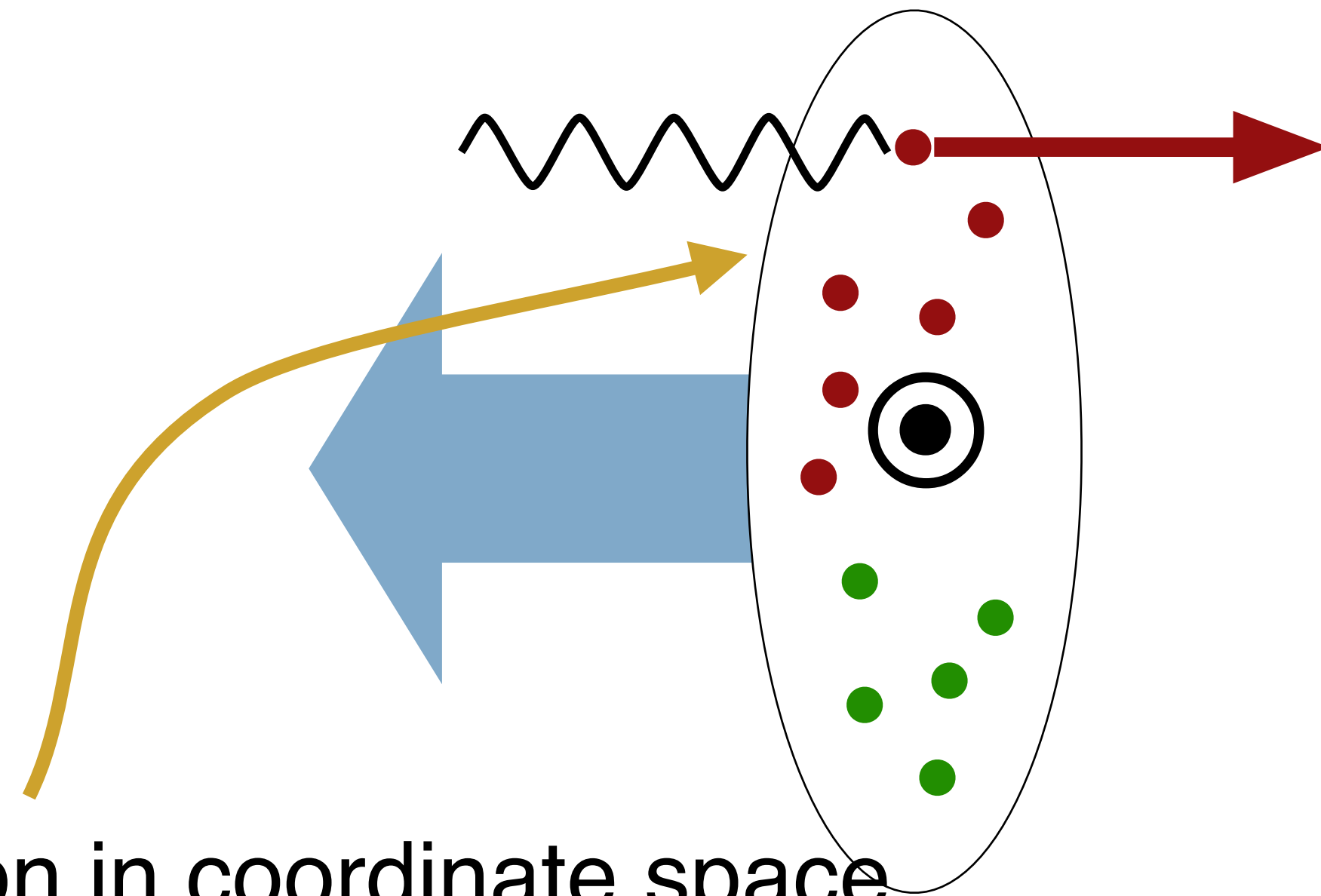
$$\mathcal{L}_z^q = \int dx d^2\vec{k}_\perp d^2\vec{b}_\perp (\vec{b}_\perp \times \vec{k}_\perp) \rho_{LU}^q(\vec{b}_\perp, \vec{k}_\perp, x)$$



based on Pasquini, Lorcé, Xiong, Yuan, PRD 85 (12)

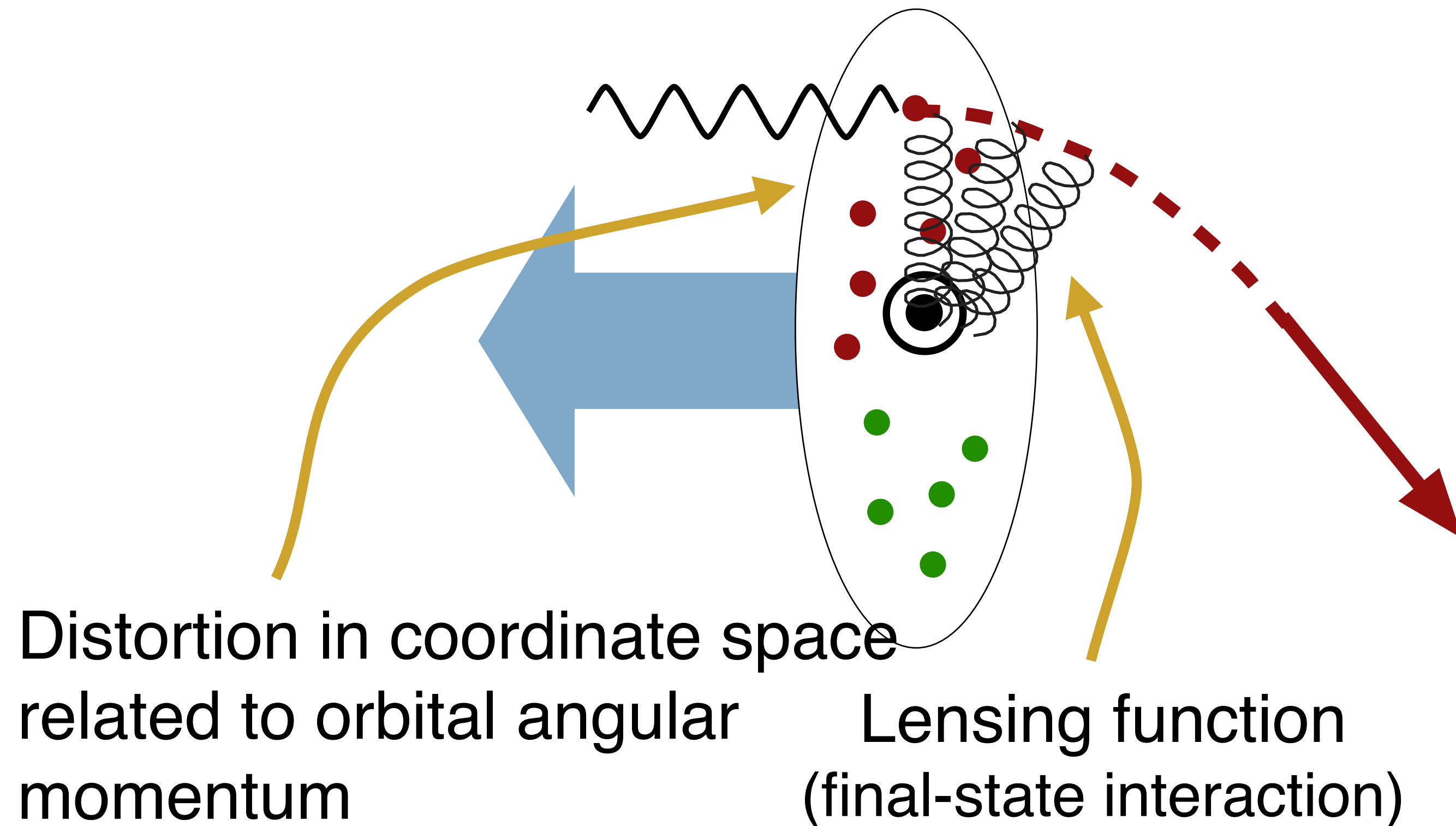


based on Burkardt, PRD66 (02)

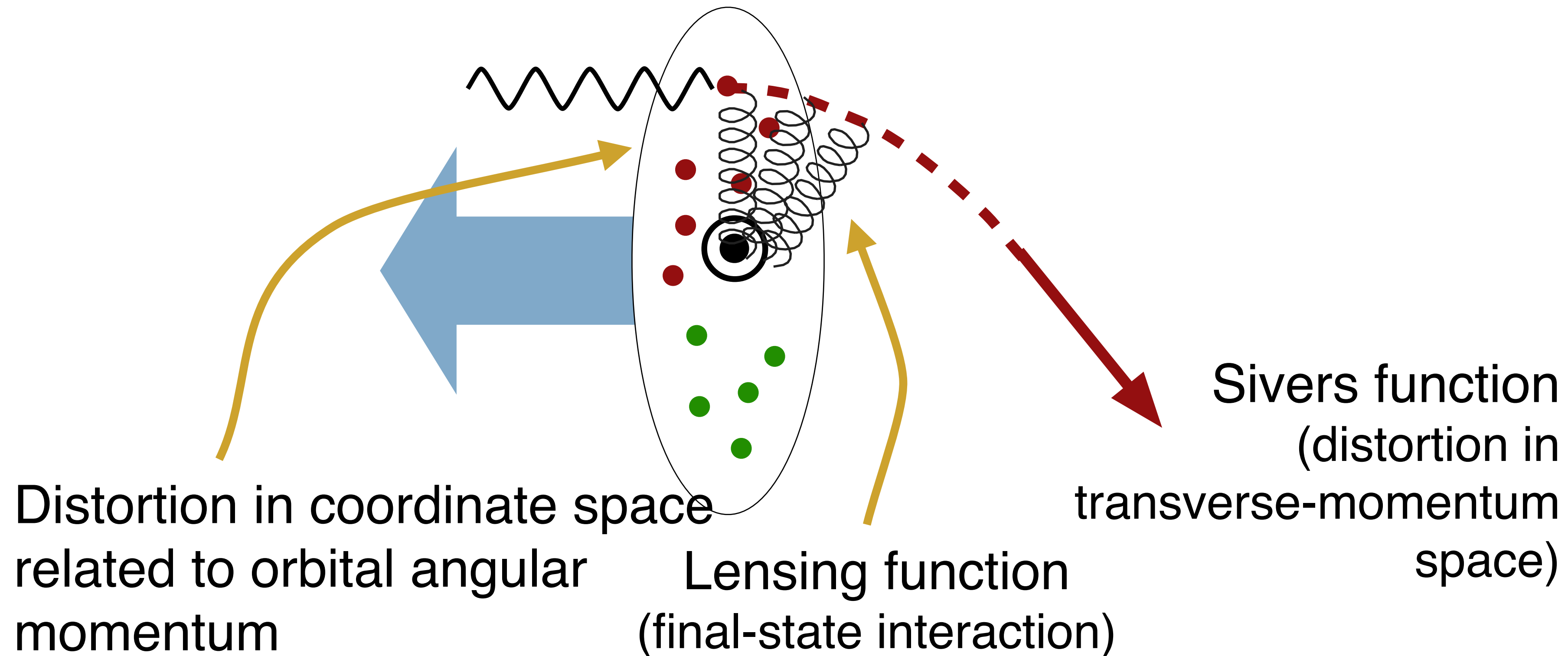


Distortion in coordinate space
related to orbital angular
momentum

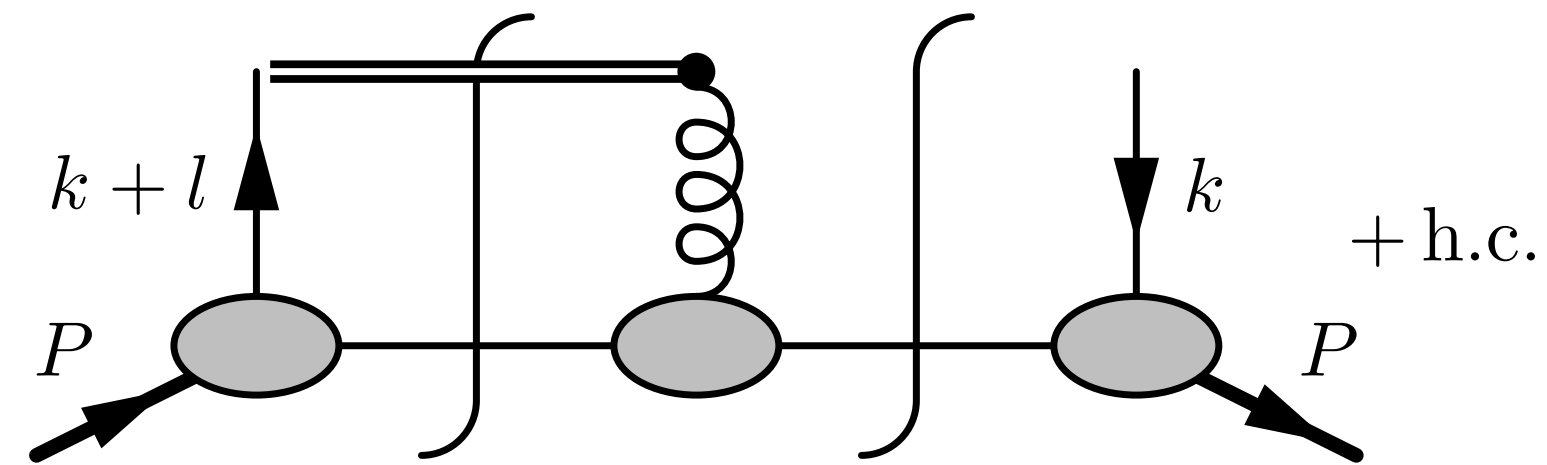
[based on Burkardt, PRD66 \(02\)](#)



[based on Burkardt, PRD66 \(02\)](#)



[based on Burkardt, PRD66 \(02\)](#)



$$f_{1T}^{\perp(0)a}(x; Q_L^2) = -\frac{3MC_F\alpha_S}{2(1-x)} E^a(x, 0, 0; Q_L^2)$$

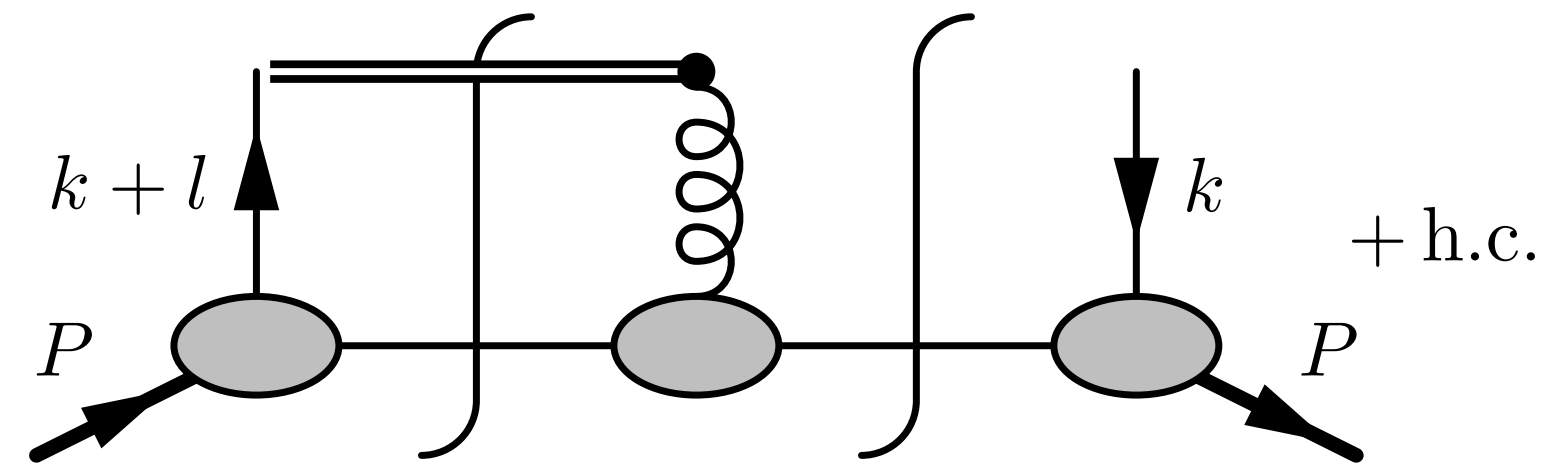
[Burkardt, Hwang, PRD69 \(04\)](#)

[Lu, Schmidt, PRD75 \(07\)](#)

[Bacchetta, Conti, Radici, PRD 78 \(08\)](#)

This relation holds only in simple models

[Bacchetta, Pasquini, Rodini, https://arxiv.org/abs/1907.06960](https://arxiv.org/abs/1907.06960)



$$f_{1T}^{\perp(0)a}(x; Q_L^2) = -\frac{3MC_F\alpha_S}{2(1-x)} E^a(x, 0, 0; Q_L^2)$$

Lensing function (flavor independent)

[Burkardt, Hwang, PRD69 \(04\)](#)

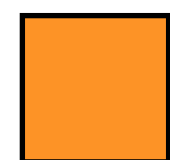
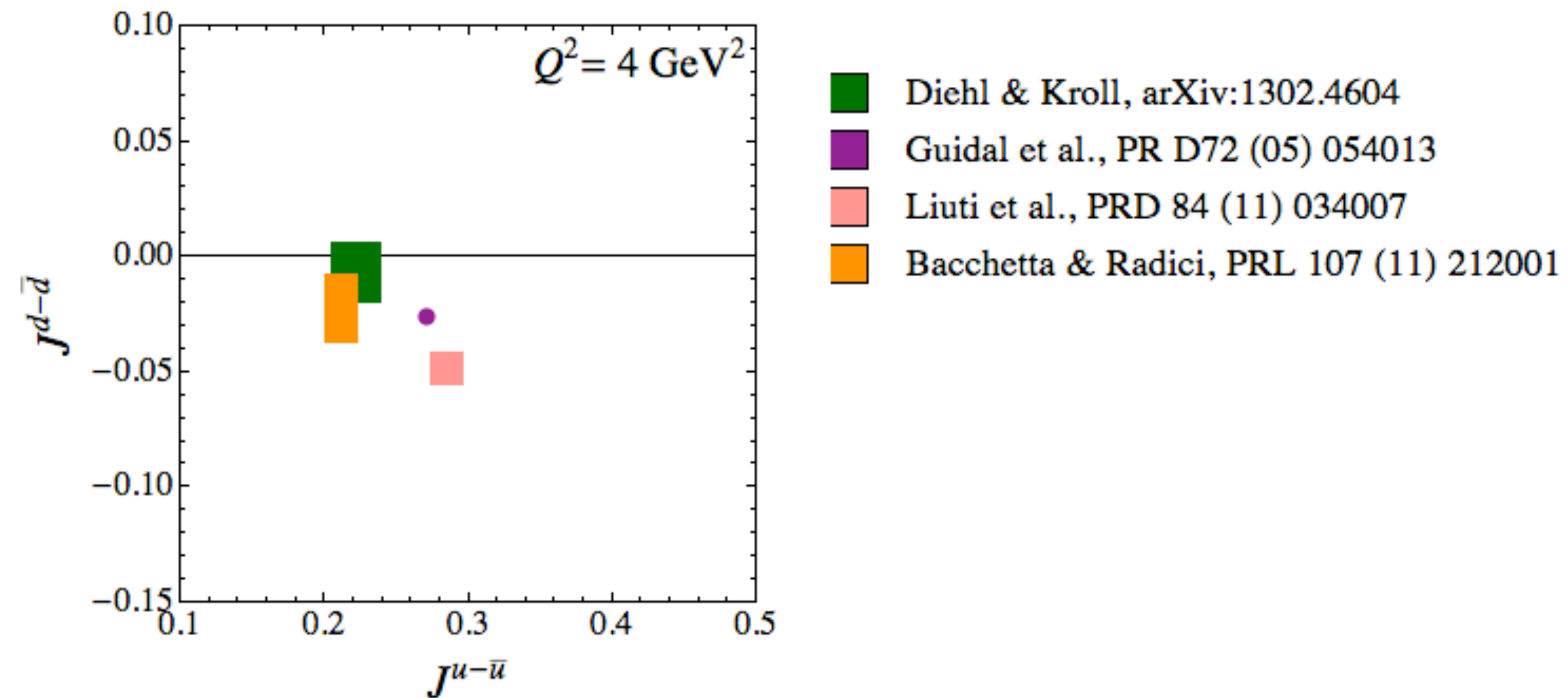
[Lu, Schmidt, PRD75 \(07\)](#)

[Bacchetta, Conti, Radici, PRD 78 \(08\)](#)

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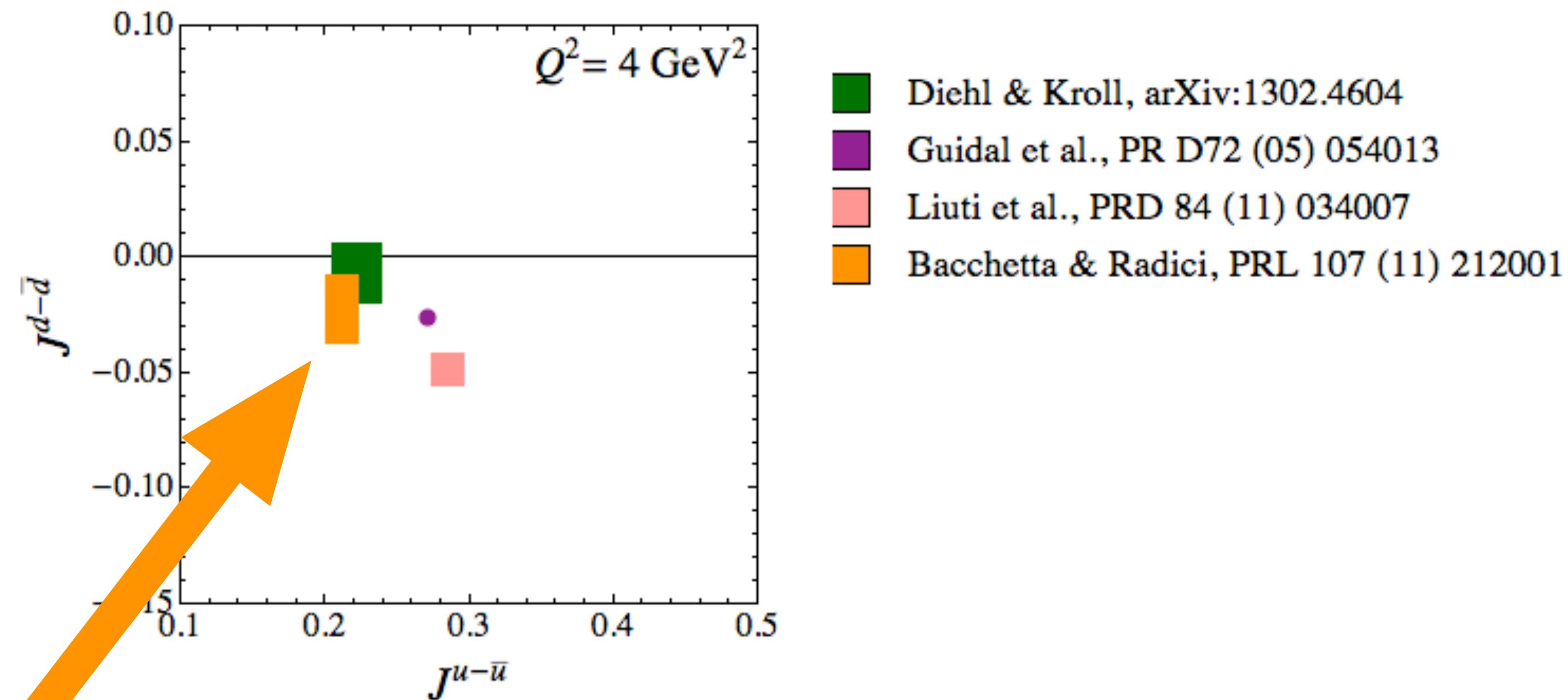
[Bacchetta, Pasquini, Rodini, https://arxiv.org/abs/1907.06960](https://arxiv.org/abs/1907.06960)

Other results obtained through form factors + assumptions

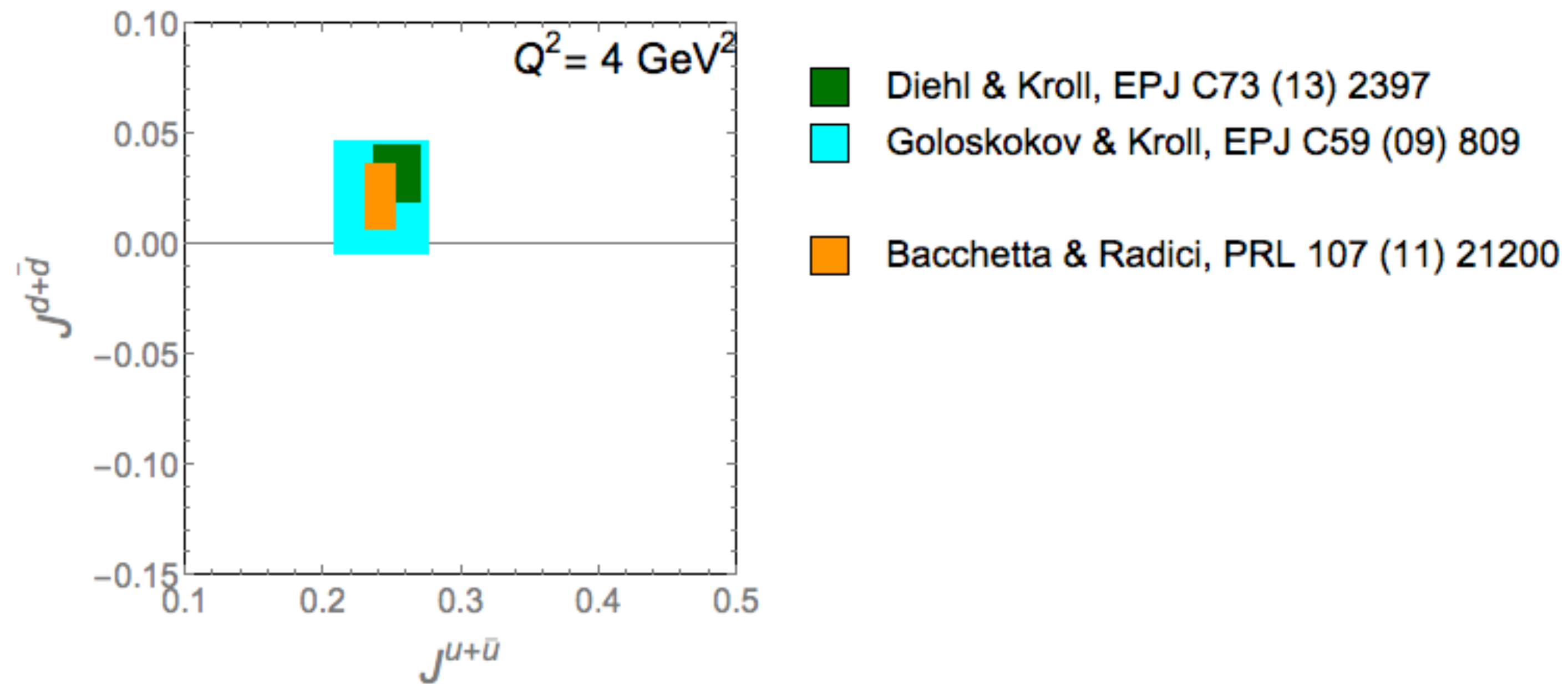


Estimate of angular momentum based on model assumptions + Sivers fit

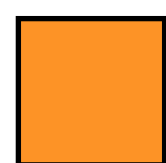
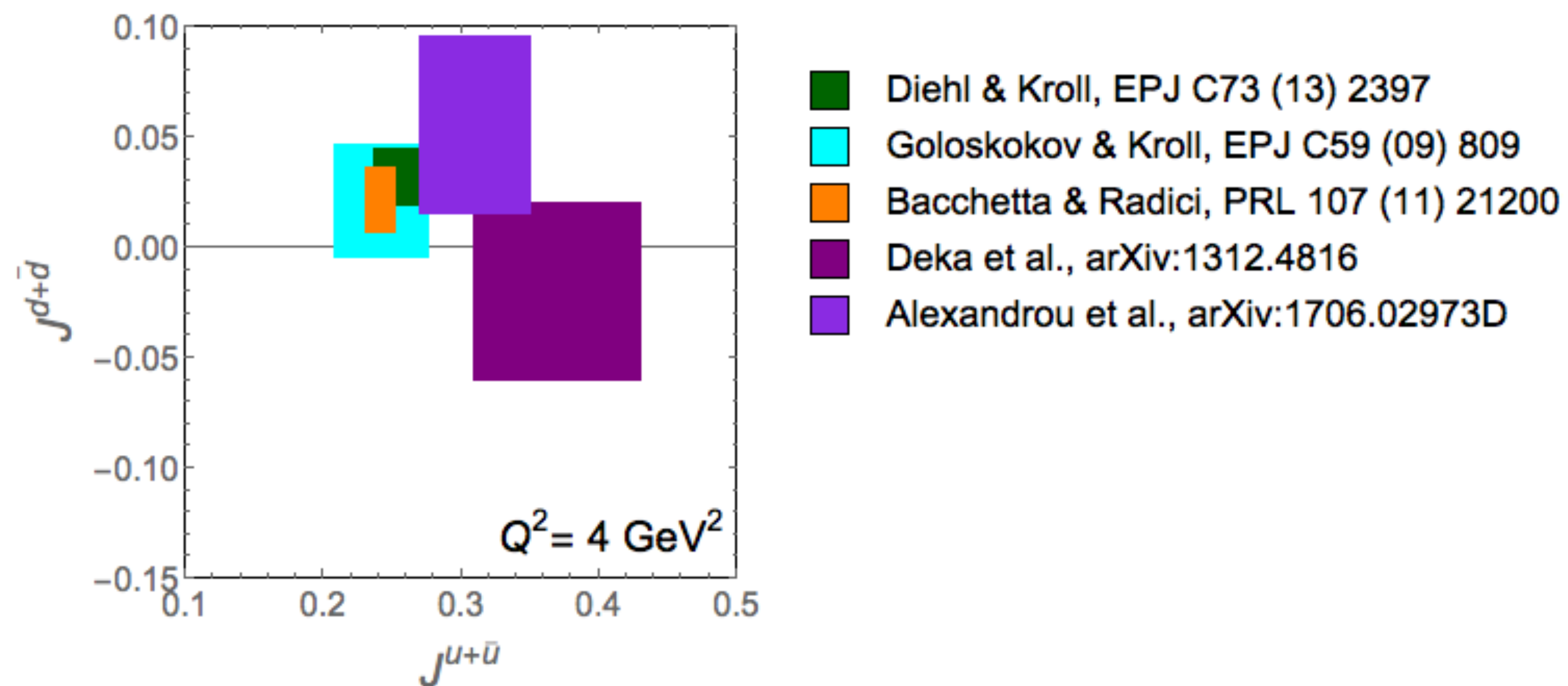
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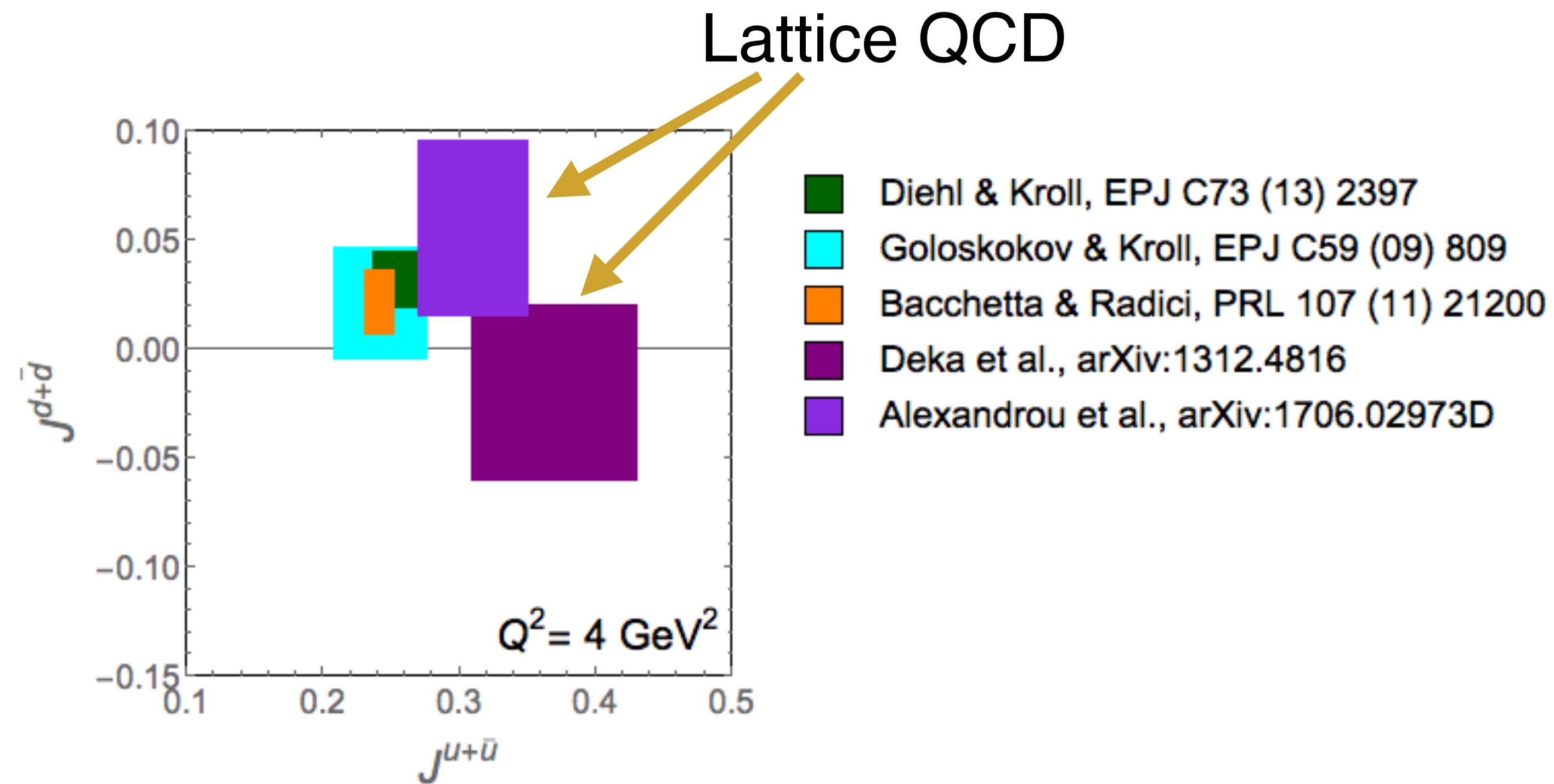
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