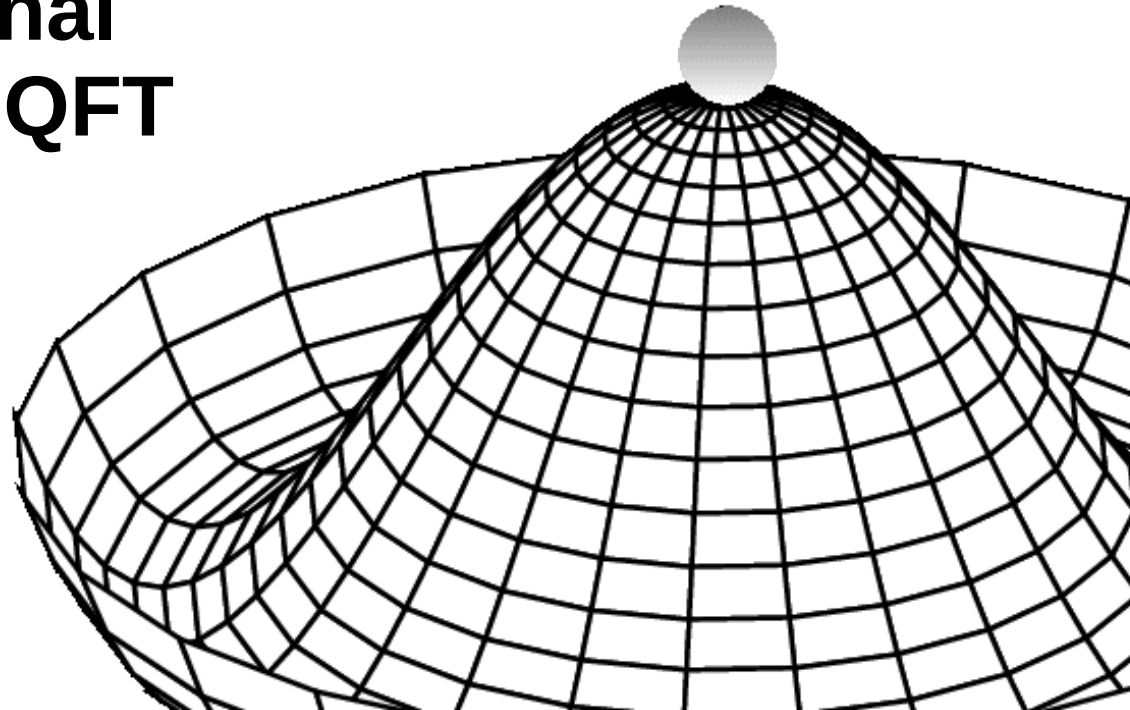


Aspects of dimensional reduction in thermal QFT

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Work in collaboration with:
M. Chala, J. Santiago, J. C. Criado and J. López



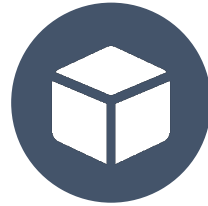
Outline



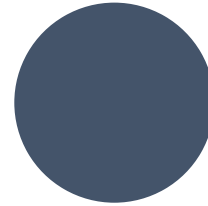
**Need for
thermal QFT**



**Formal
aspects**



**3dEFT
approach**



**Future
applications**



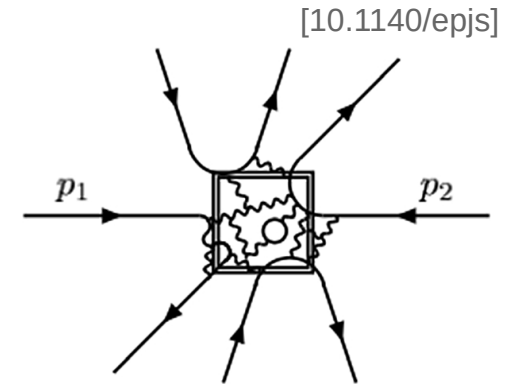
Need for thermal QFT

- The Standard Model is usually studied at **zero temperature** in a **non-zero temperature Universe**
- At finite temperature, collective statistical effects appear

Standard Model of Elementary Particles

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	=2.2 MeV/c ²	=1.28 GeV/c ²	=173.1 GeV/c ²	0	=124.97 GeV/c ²
charge	2/3	2/3	2/3	0	0
spin	1/2	1/2	1/2	1	0
	u up	c charm	t top	g gluon	H higgs
QUARKS	d down	s strange	b bottom	γ photon	
	=4.7 MeV/c ²	=96 MeV/c ²	=4.18 GeV/c ²	0	
	-1/3	-1/3	-1/3	0	
	1/2	1/2	1/2	1	
	e electron	μ muon	τ tau	Z Z boson	
LEPTONS	=0.511 MeV/c ²	=105.66 MeV/c ²	=1.7768 GeV/c ²	=91.19 GeV/c ²	
	-1	-1	-1	0	
	1/2	1/2	1/2	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
	<1.0 eV/c ²	<0.17 MeV/c ²	<18.2 MeV/c ²	=80.360 GeV/c ²	
	0	0	0	±1	
	1/2	1/2	1/2	1	
				W W boson	

SCALAR BOSONS (Higgs)
GAUGE BOSONS VECTOR BOSONS (photon, Z, W)

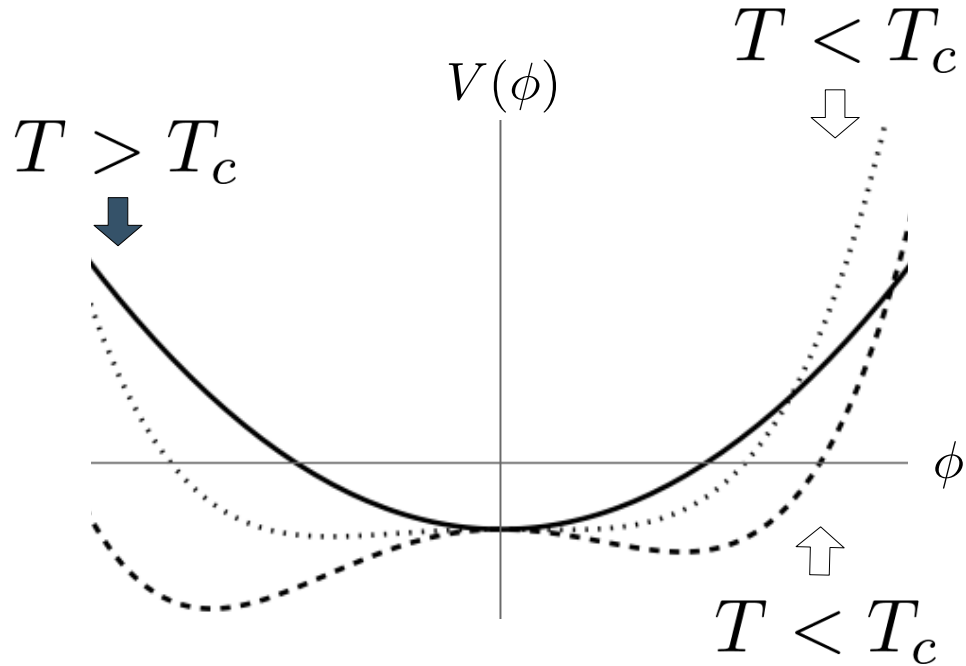


$$Z^{-1} \text{Tr} (e^{-\beta \mathcal{H}} \langle p_1 \dots | p_2 \dots \rangle)$$

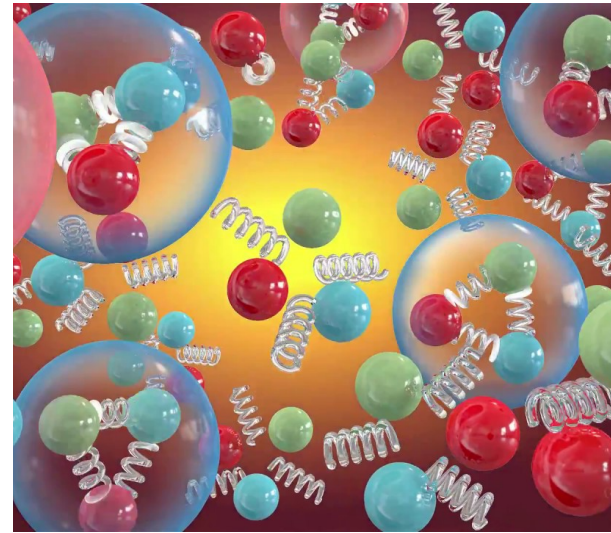
Thermal average



Need for thermal QFT



Phase transition



Quark-gluon
plasma



Formal aspects

- **Generating functional** ($J=0$) in QFT:

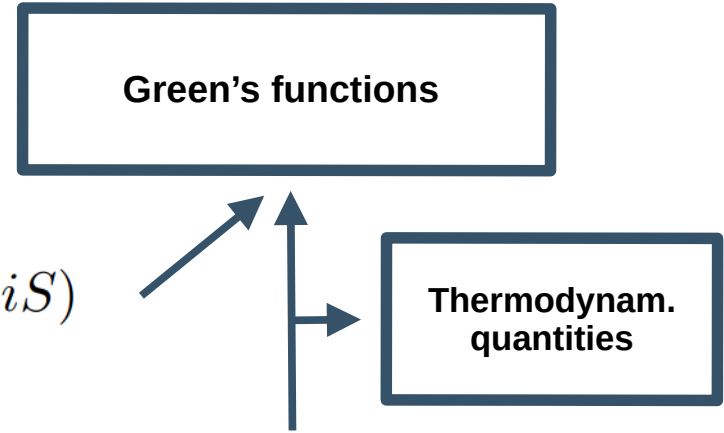
$$\mathcal{Z}[0] = \langle q' t | q 0 \rangle = \langle q' 0 | e^{-i\mathcal{H}t} | q 0 \rangle = \mathcal{N} \int \mathcal{D}q \exp(iS)$$

- **Partition function** in quantum statistical mechanics:

$$\mathcal{Z}_{\text{th}} = \text{Tr} (e^{-\beta\mathcal{H}}) = \sum_q \langle q 0 | e^{-\beta\mathcal{H}} | q 0 \rangle = \mathcal{N} \int_{q(0)=q(-i\beta)} \mathcal{D}q \exp(-S_E)$$

where

$$S = \int d^4x \mathcal{L}(q, \dot{q}) \quad \text{and} \quad S_E = iS = \int_0^\beta d\tau \int d^3x \mathcal{L}_E(q_E, \dot{q}_E)$$





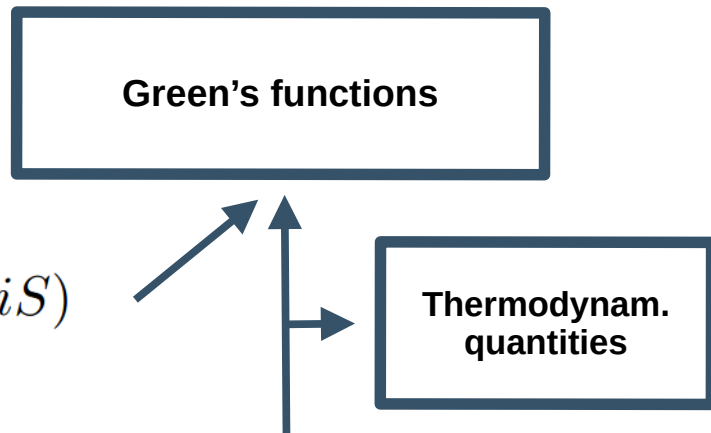
Formal aspects

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QFT at finite temperature = Euclidean QFT at zero temperature with periodic time

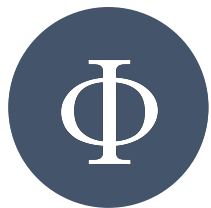


Formal aspects

Check out this review by M. Laine and A. Vuorinen [1701:01554]

Summary

	$T = 0$	$T \neq 0$
Space-time	Minkowski $\mathbb{R}^3 \times \mathbb{R}^1$ ($-\infty < t < \infty$)	Euclidean $\mathbb{R}^3 \times \mathcal{S}^1$ ($0 < (\tau = it) \leq \beta$)
Bosonic fields	$\phi(x)$	$\sum_{n=-\infty}^{\infty} \phi_n(\mathbf{x}) \exp(i2\pi n T \tau)$
Bosonic momentum	p^μ	$P = (2\pi n T, \mathbf{p})$
Fermionic fields	$\psi(x)$	$\sum_{n=-\infty}^{\infty} \psi_n(\mathbf{x}) \exp(i\pi(2n+1)T\tau)$
Fermionic momentum	p^μ	$P = (\pi(2n+1)T, \mathbf{p})$



Formal aspects

Feynman rules

	$T = 0$	$T \neq 0$
Vertex	Same	Same
Product of momenta	$p^\mu q_\mu$	$-P \cdot Q$
Propagator	$\frac{i}{p^2 - m^2}$	$\frac{1}{P^2 + m^2}$
Loop integral	$\int \frac{d^4 q}{(2\pi)^4}$	$\oint_Q \equiv T \sum_{n=-\infty}^{\infty} \int \frac{d^3 q}{(2\pi)^3}$





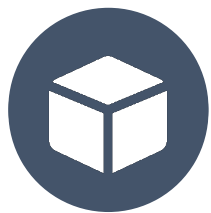
Formal aspects: $\lambda\phi^4$ model

$$\mathcal{L}_M = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{m^2}{2} \phi^2 - \lambda \phi^4$$

↓ Thermal bath

$$\mathcal{L}_E = \sum_{n=-\infty}^{\infty} \left[\frac{1}{2} (\partial_i \phi_n)^2 + \frac{1}{2} (m^2 + (2\pi nT)^2) \phi_n^2 \right] + \text{coupling between modes}$$

- At energies below the temperature scale, heavy thermal modes can be **integrated out**
- We are left with an **effective 3D field theory (3dEFT)** for the **zero mode** only



3dEFT approach

Analogous calculation without effective operators
[2101:05528]

Full theory:

$$\mathcal{L}_E = \frac{1}{2} (\partial_\tau \phi)^2 + \frac{1}{2} (\partial_i \phi)^2 + \frac{m^2}{2} \phi^2 + \lambda \phi^4$$

Off-shell
matching

$$\mathcal{O}(\lambda^2, P^4)$$

$$m^2 \sim \lambda T^2$$

$$P_0 = 0$$



$$\mathcal{L}_{(4)} = \frac{K}{2} (\partial_i \varphi)^2 + \frac{m_3^2}{2} \varphi^2 + \lambda_3 \varphi^4,$$

$$\mathcal{L}_{(6)} = \beta_{61} \dot{\varphi}^2 \varphi \partial^2 \varphi + \beta_{62} \varphi^3 \partial^2 \varphi,$$

$$\mathcal{L}_{(8)} = \alpha_{82} \varphi^2 \partial_i \partial_j \varphi \partial^i \partial^j \varphi + \beta_{82} \varphi^3 \partial^4 \varphi + \beta_{83} \varphi^2 \partial^2 \varphi \partial^2 \varphi.$$

$$\mathcal{L}_3 = \mathcal{L}_{(4)} + \mathcal{L}_{(6)} + \mathcal{L}_{(8)}$$

See talk by J. López!

3d EFT:

(up to 4 deriv. and 4 field
insertions)



3dEFT approach

- We first compute the **2-point 1PI diagrams** in both theories:

$$\boxed{\varphi \rightarrow \varphi} \quad (\text{tree-level})$$

$$\text{---}\blacktriangleright\text{---} = -K\mathbf{p}^2 - m_3^2 + 2\beta_{61}\mathbf{p}^4$$

$$\boxed{\phi \rightarrow \phi} \quad (\text{tree-level})$$

$$\text{---}\blacktriangleright\text{---} = -(P^2 + m^2)$$



3dEFT approach

$\phi \rightarrow \phi$ (1-loop)

$$\begin{aligned}
 \text{Diagram: a horizontal line with a circle loop attached to a central vertex} &= -\frac{4!\lambda}{2} \int' \frac{1}{Q^2 + m^2} = -\frac{4!\lambda}{2} \int' \left[\frac{1}{Q^2} \left(1 - \frac{m^2}{Q^2} + \mathcal{O}\left(\frac{m^4}{Q^4}\right) \right) \right] \\
 &= -\frac{4!\lambda}{2} \left(\int' \frac{1}{Q^2} - m^2 \int' \frac{1}{Q^4} \right) + \mathcal{O}(\lambda^3)
 \end{aligned}$$

1) Hard region expansion in $\frac{m^2}{(\pi T)^2}$. Double hard region expansion at 2-loop level [2311.13630]

2) Write dashed sum-integrals as un-dashed $\int'_{Q'} f(\mathbf{q}, Q_0) = \int'_{Q} f(\mathbf{q}, Q_0) - \int_{\vec{q}} f(\mathbf{q}, 0)$ scaleless vanish in dim. reg.



3dEFT approach

3) Use (scheme-dependent) master bosonic 1-loop sum-integral formula in $d=3$ limit:

$$I_{\alpha}^{\beta\gamma}(d) \equiv \int_Q \frac{(Q_0^2)^{\beta} (\mathbf{q}^2)^{\gamma}}{(Q^2)^{\alpha}} = \left(\frac{e^{\gamma_E} \mu^2}{4\pi} \right)^{\epsilon} \frac{2T(2\pi T)^{d-2\alpha+2\beta+2\gamma}}{(4\pi)^{d/2}} \quad [2101.05528]$$
$$\frac{\Gamma(d/2 + \gamma)\Gamma(-d/2 + \alpha - \gamma)}{\Gamma(d/2)\Gamma(\alpha)} \zeta(-d + 2\alpha - 2\beta - 2\gamma)$$

We need:

$$I_1^{00}(3) = \frac{T^2}{12}$$

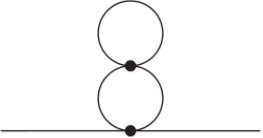

$$I_2^{00}(3) = \frac{1}{16\pi^2} \left(\frac{1}{\epsilon} + \log\left(\frac{\mu^2}{T^2}\right) + \psi^{(0)}\left(\frac{1}{2}\right) + \text{constants} \right)$$

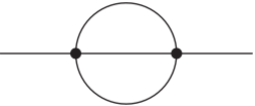

Requires renormalization



3dEFT approach

$\phi \rightarrow \phi$ (2-loop)


$$= \frac{(-4!\lambda)^2}{4} \not\int_{Q K} \frac{1}{(Q^2 + m^2)^2(K^2 + m^2)} = \frac{(-4!\lambda)^2}{4} I_2^{00}(d) I_1^{00}(d) + \mathcal{O}(\lambda^3)$$



$$= \frac{(-4!\lambda)^2}{2} \not\int_{Q K} \frac{1}{(Q^2 + m^2)(K^2 + m^2)((P - Q - K)^2 + m^2)}$$




3dEFT approach

Sum-integrals like these appear in the 2-loop sunset diagram:

$$L_{s_1 s_2 s_3}^{s_4 s_5}(d) \equiv \int_{QK} \frac{(Q_0)^{s_4} (K_0)^{s_5}}{(Q^2)^{s_1} (K^2)^{s_2} ((Q+K)^2)^{s_3}}$$

We do **not** have an **easy and systematic** way to compute these. We find some specific examples in the literature using IBP relations and tensor reduction formulae. [1207.4042, 1208.0284]

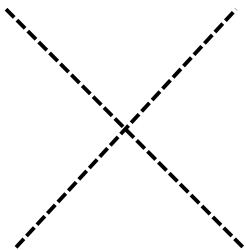
Would be useful to have codes to solve some families of loop sum-integrals for generic indices.



3dEFT approach

- Now, we compute the **4-point 1PI diagrams** in both theories:

$\varphi\varphi \rightarrow \varphi\varphi$ (tree-level)



$$\begin{aligned} &= -4!\lambda_3 \\ &+ 4\alpha_{82} \left[(p_1 \cdot p_2)^2 + (p_1 \cdot p_3)^2 + (p_1 \cdot p_4)^2 + (p_2 \cdot p_3)^2 + (p_2 \cdot p_4)^2 + (p_3 \cdot p_4)^2 \right] \\ &+ 6\beta_{62} [p_1^2 + p_2^2 + p_3^2 + p_4^2] \\ &+ 6\beta_{82} [p_1^4 + p_2^4 + p_3^4 + p_4^4] \\ &+ 4\beta_{83} [p_2^2 p_1^2 + p_3^2 p_1^2 + p_4^2 p_1^2 + p_2^2 p_3^2 + p_2^2 p_4^2 + p_3^2 p_4^2] \Big|_{p_1+p_2+p_3+p_4=0} \end{aligned}$$

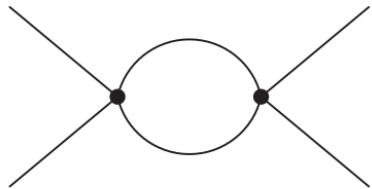


3dEFT approach

$\phi\phi \rightarrow \phi\phi$ (tree-level)

$$\text{X} = -4!\lambda$$

$\phi\phi \rightarrow \phi\phi$ (1-loop)

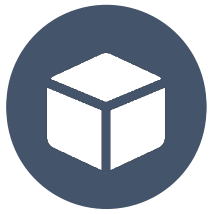


$$= \frac{(-4!\lambda)^2}{2} \sum_{j=\{s,t,u\}} \not\int' \frac{1}{(Q^2 + m^2)((Q + P_j)^2 + m^2)}$$

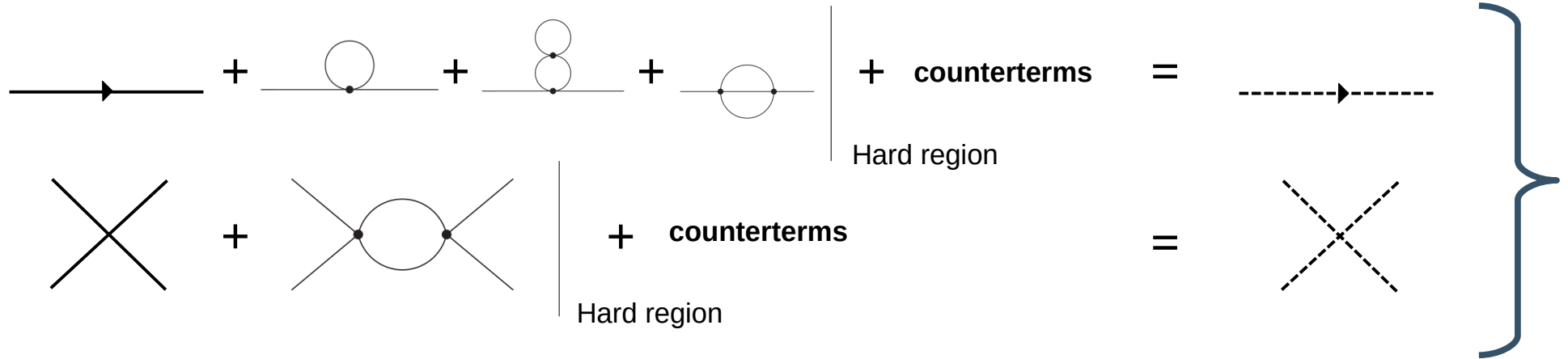
Can be written in terms of

$I_\alpha^{\beta\gamma}(d)$

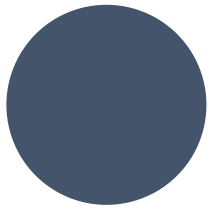




3dEFT approach



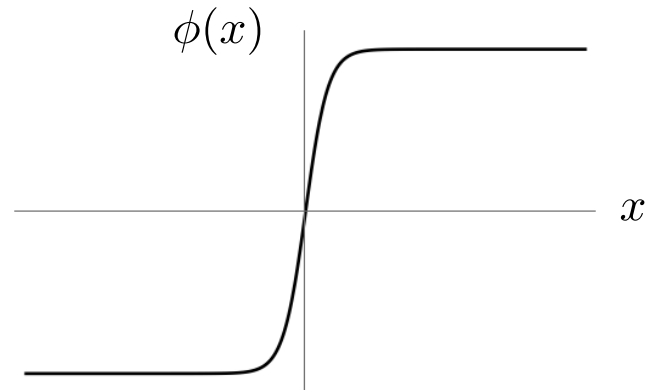
- Wilson coefficients are matched at scale $\mu \sim T$
- 3dEFT coefficients are then run (RGE) to a lower energy scale $\mu \sim m \sim \sqrt{\lambda T}$

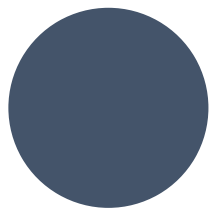


Future applications

Why such a simple model?

- We want to test our EFT-building procedures and **gain some intuition** about 3D Lagrangians
- What are the consequences of adding effective operators with **higher order derivatives**?
 - In 4D theories with asymmetric potentials we find **non-constant classical field configurations** called *bounces*.
 - What kind of classical solutions can we find in a 3dEFT?





Future applications

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- What are the consequences of adding effective operators with **higher order derivatives**?

What are our future goals?

- Cosmological phase transitions and gravitational waves production pre-CMB (?) [2305.02357]
- Apply our knowledge to the Standard Model to study the EW phase transition

Still A LOT for me to learn :)

Thank you for your attention!
¡Gracias por vuestra atención!