

Qubit Regularization: Asymptotic Freedom via New Renormalization Group Flows

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Lattice Gauge Theory Work in Progress



***Supported by:
US Department of Energy***



Quantum Field Theory with Quantum Computers

Quantum Field Theory with Quantum Computers

Growing field of research within our community!

Plenary Talk 2018, Preskill

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Two types of
opportunities:



How to reformulate the QFT



How to use the quantum
hardware

Quantum Field Theory with Quantum Computers

Growing field of research within our community!

Plenary Talk 2018, Preskill



This talk is about showing you a glimpse of how quantum computing is already helping us learn to formulate “old” problems with “new” variables and understand the field deeper.

Lagrangian \longrightarrow Hamiltonian

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“Local” lattice Hilbert spaces play an important role

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Traditional: \mathcal{H}_{Trad} $\dim(\mathcal{H}_{Trad}) = \infty$

“Digital” quantum computer: \mathcal{H}_Q $\dim(\mathcal{H}_Q) = \text{finite}$

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qubit regularization

Hersh Singh, *SC Phys.Rev.D* 100 (2019) 5, 054505
Hanqing Liu, *SC Symmetry* 14 (2022) 2, 305

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Continuum QFT emerges through the usual limiting process

continuum limit

$$a \rightarrow 0$$

Thermodynamic limit

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Many papers seem to assume this is necessary!

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This is not required at least in some cases!

An interesting alternate approach is the D-theory

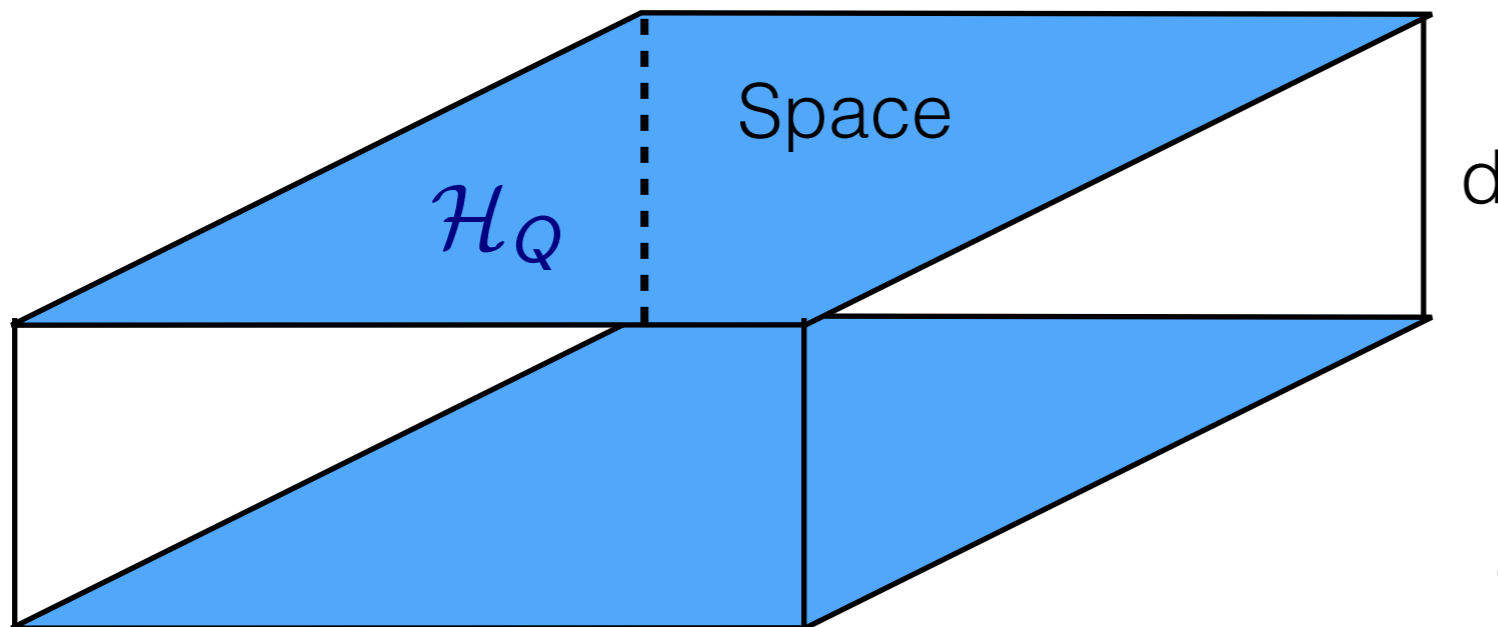
Lattice 1998, Plenary talk by Wiese.

Brower, SC, Riederer, Wiese, NPB 693 (2004), 149

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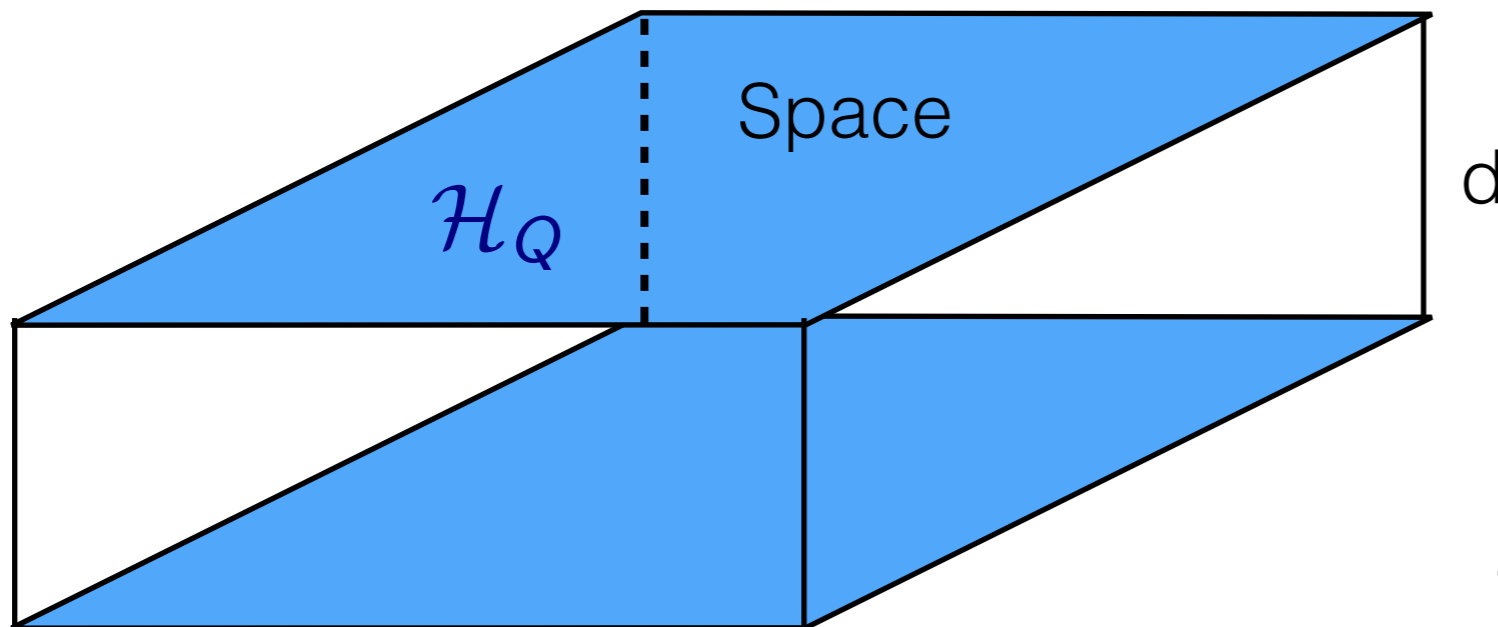


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RG plays an important role!

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12 Parallel Talks on Monday

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Various types of qubit regularizations.

How to prepare initial states.

How to evolve systems in real time on a real hardware!

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Apologies: I do not plan on reviewing these results!

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(Two examples, Euclidean space)

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(Two examples, Euclidean space)

We will discover a new a type of RG flow!

Collaborators



Siew



Maiti



Liu



Singh



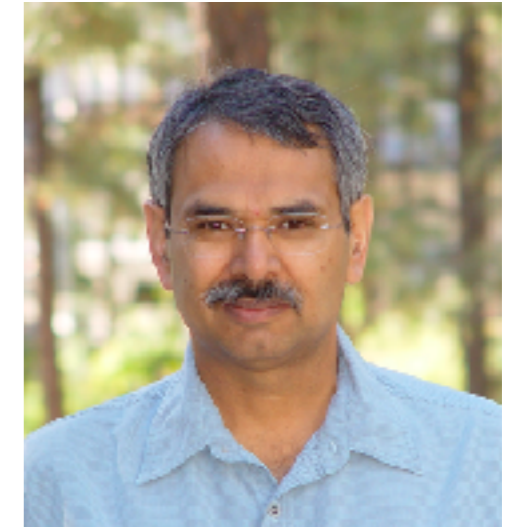
Marinkovic



Banerjee



Bhattacharya



Gupta

Outline

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Asymptotic Freedom as an RG flow

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Recovering Asymptotic Freedom with finite Hilbert space

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Example 1: $O(3)$ Non-linear sigma model

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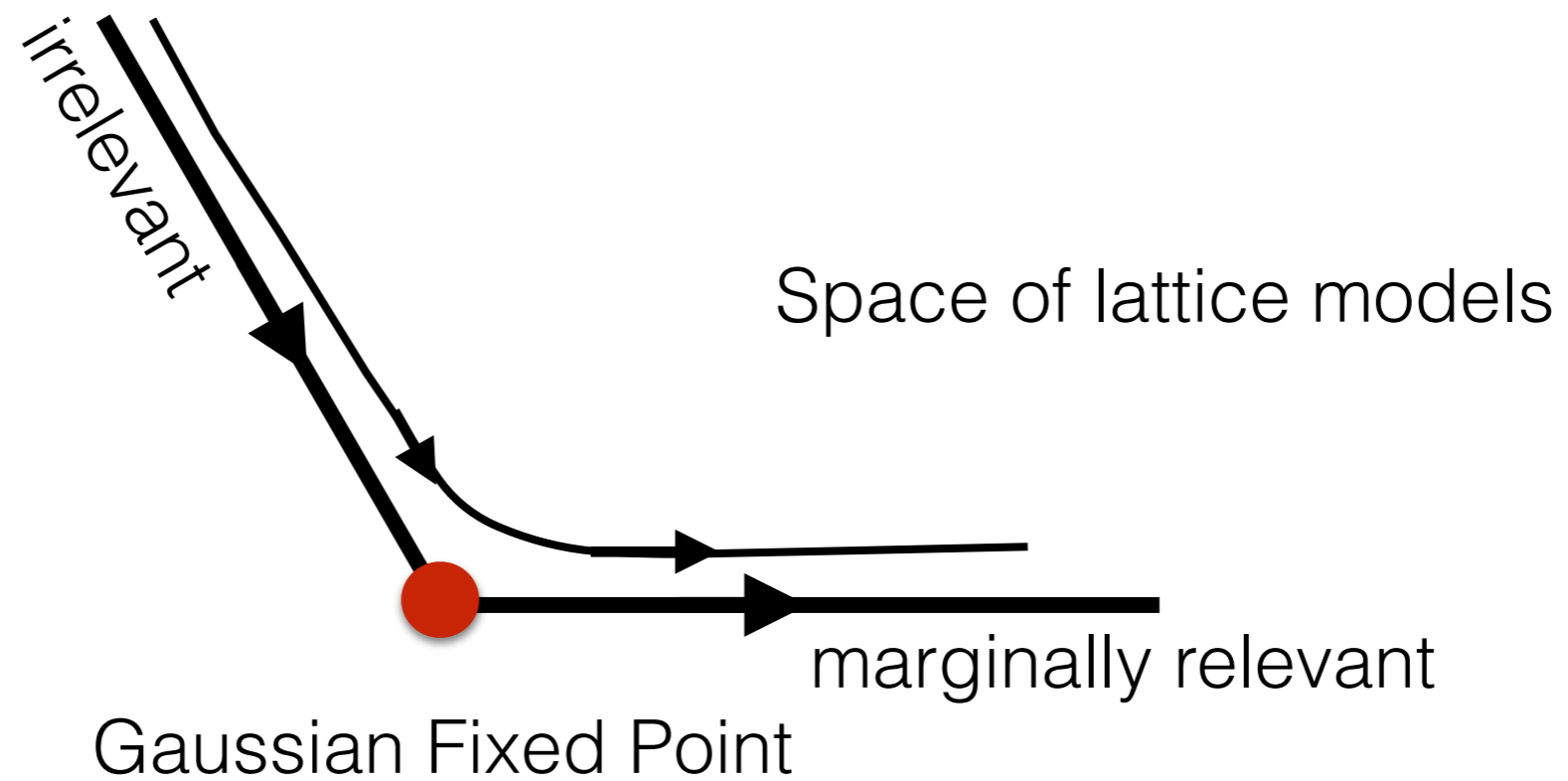
Qubit Regularization of Asymptotic Freedom in Gauge Theories

Conclusions

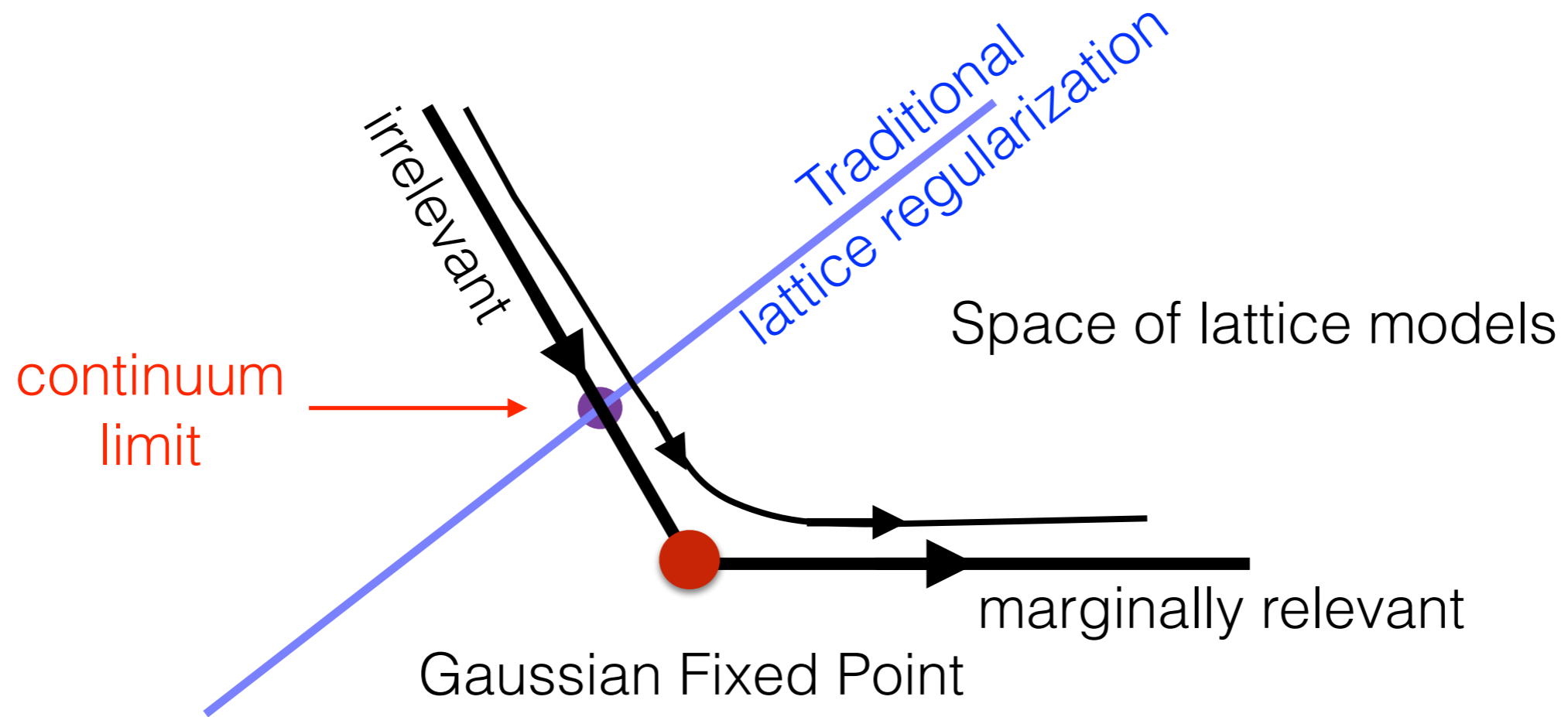
Asymptotic Freedom as an RG Flow

Traditional Lattice Regularization

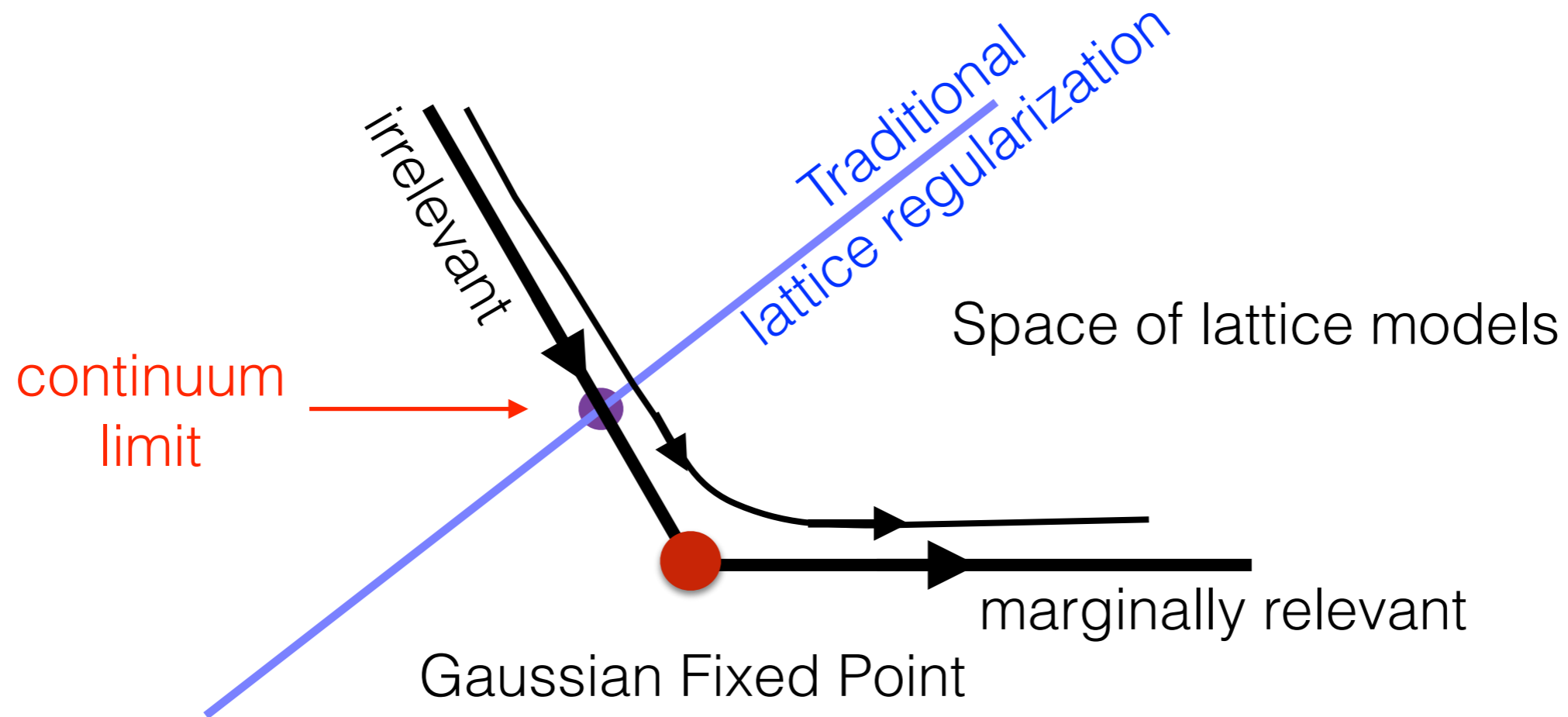
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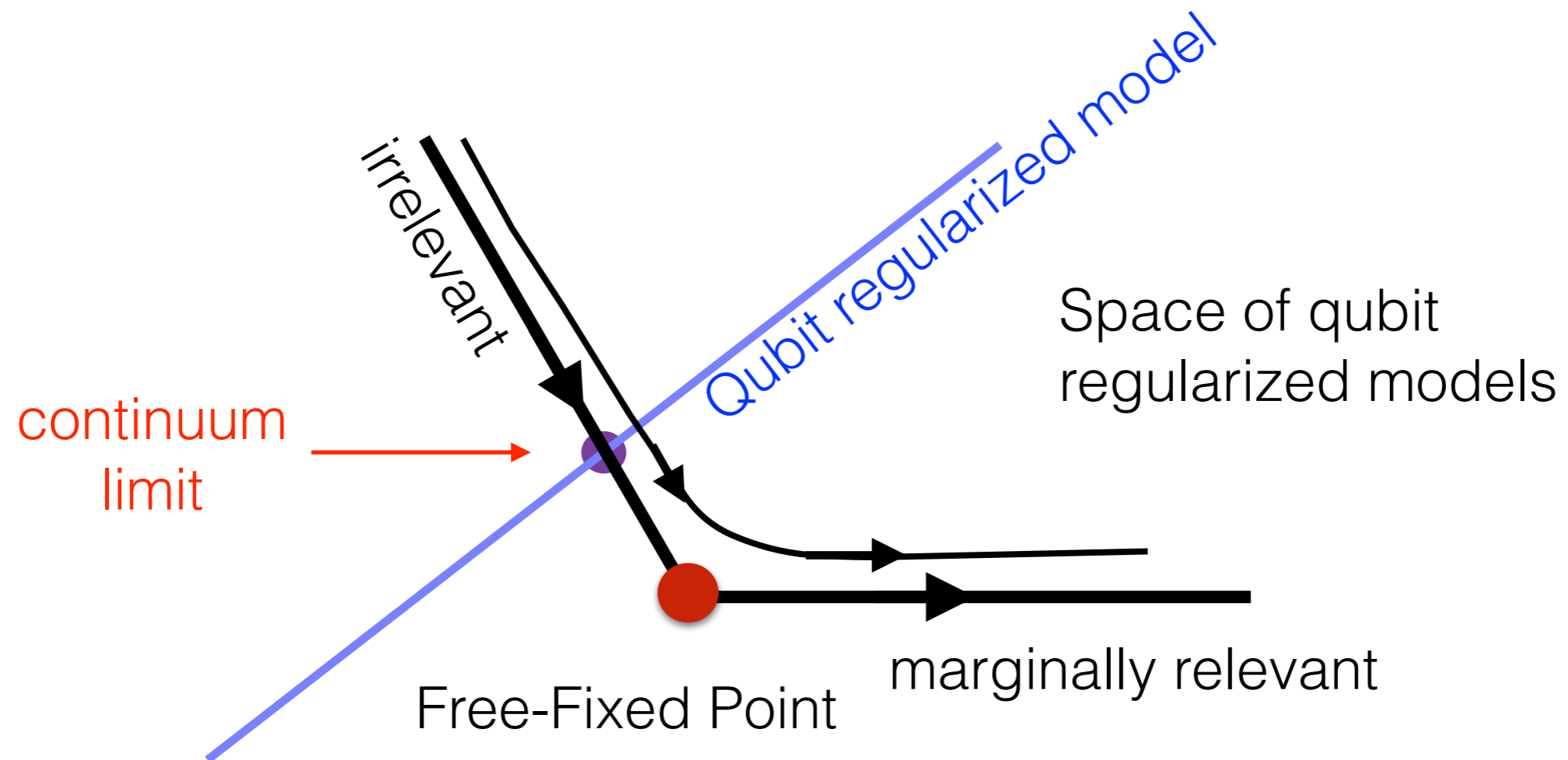
From an RG perspective, the limit

$$\mathcal{H}_Q \longrightarrow \mathcal{H}_{\text{Trad}}$$

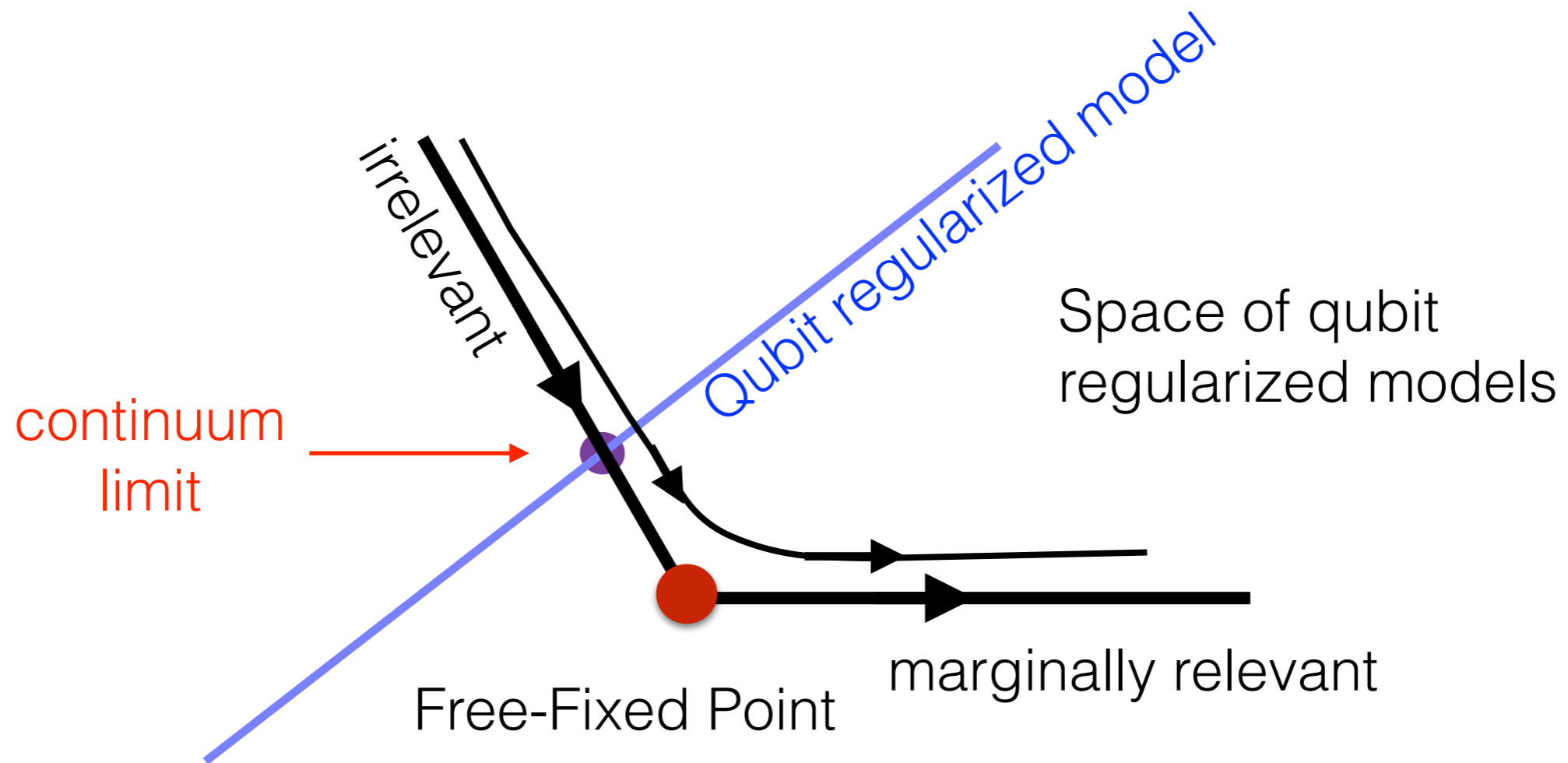
may not be necessary!

We may be able to fine tune to
the free critical surface within \mathcal{H}_Q

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Lesson from RG:

Explore the space of qubit models,
don't be stuck with the traditional Hamiltonian.

Recovering Asymptotic Freedom with a finite Hilbert space

Example: $O(3)$ Non-linear sigma model

Traditional Lattice Action:

$$S_L = -\beta \sum_{\langle (x,\tau), (y,\tau') \rangle} \vec{\phi}_{x,\tau} \cdot \vec{\phi}_{y,\tau'}$$

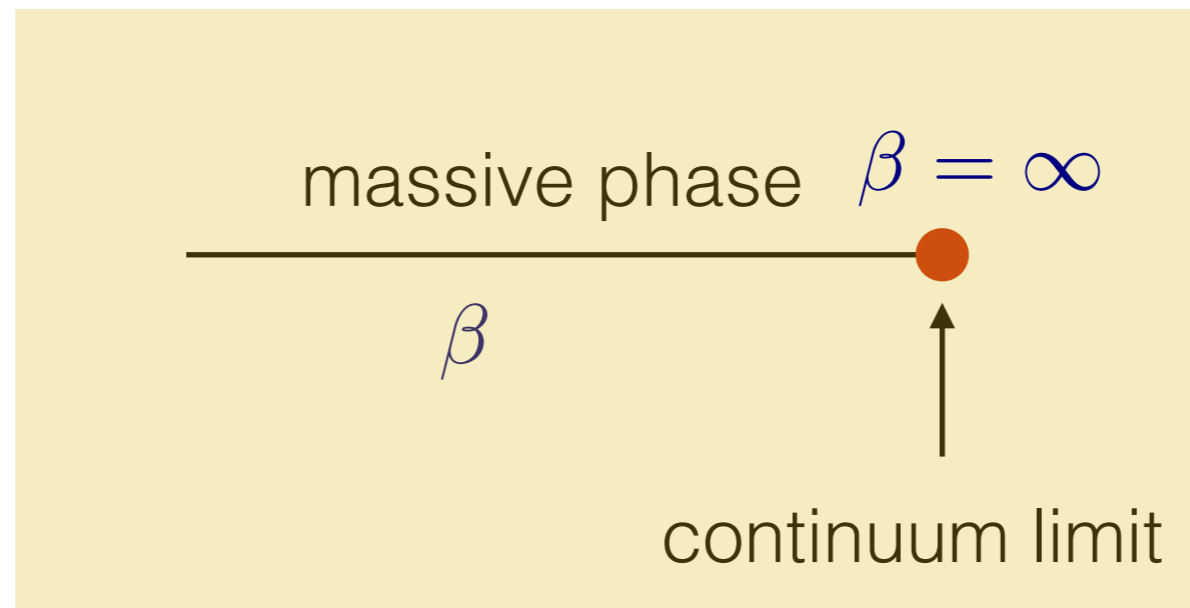
$\vec{\phi}_{x,\tau}$ = three component unit vector

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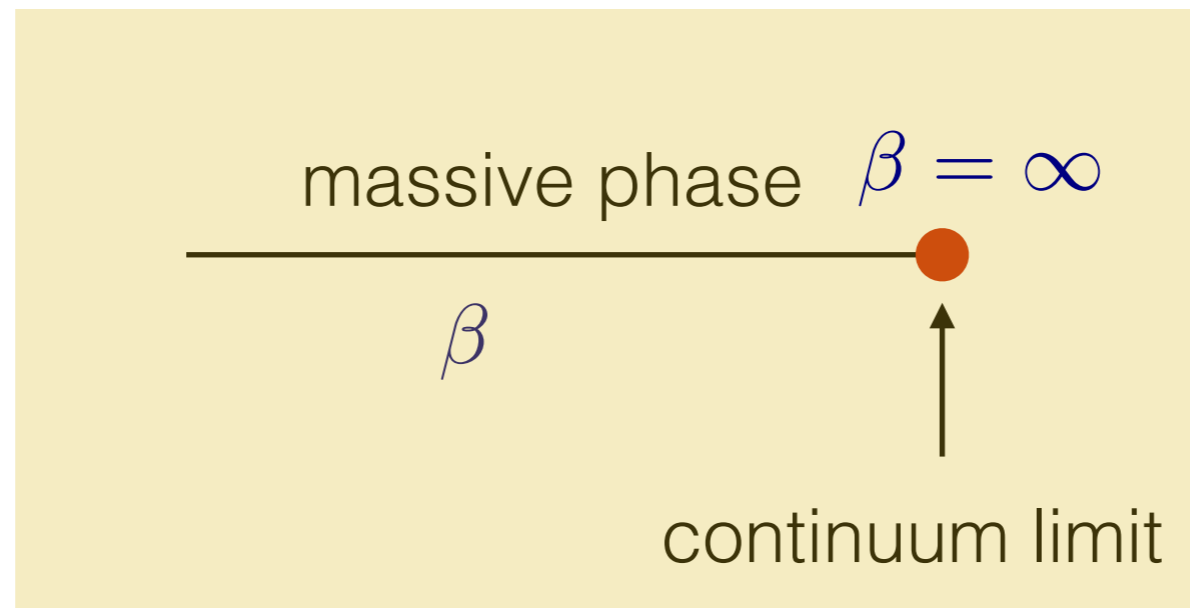
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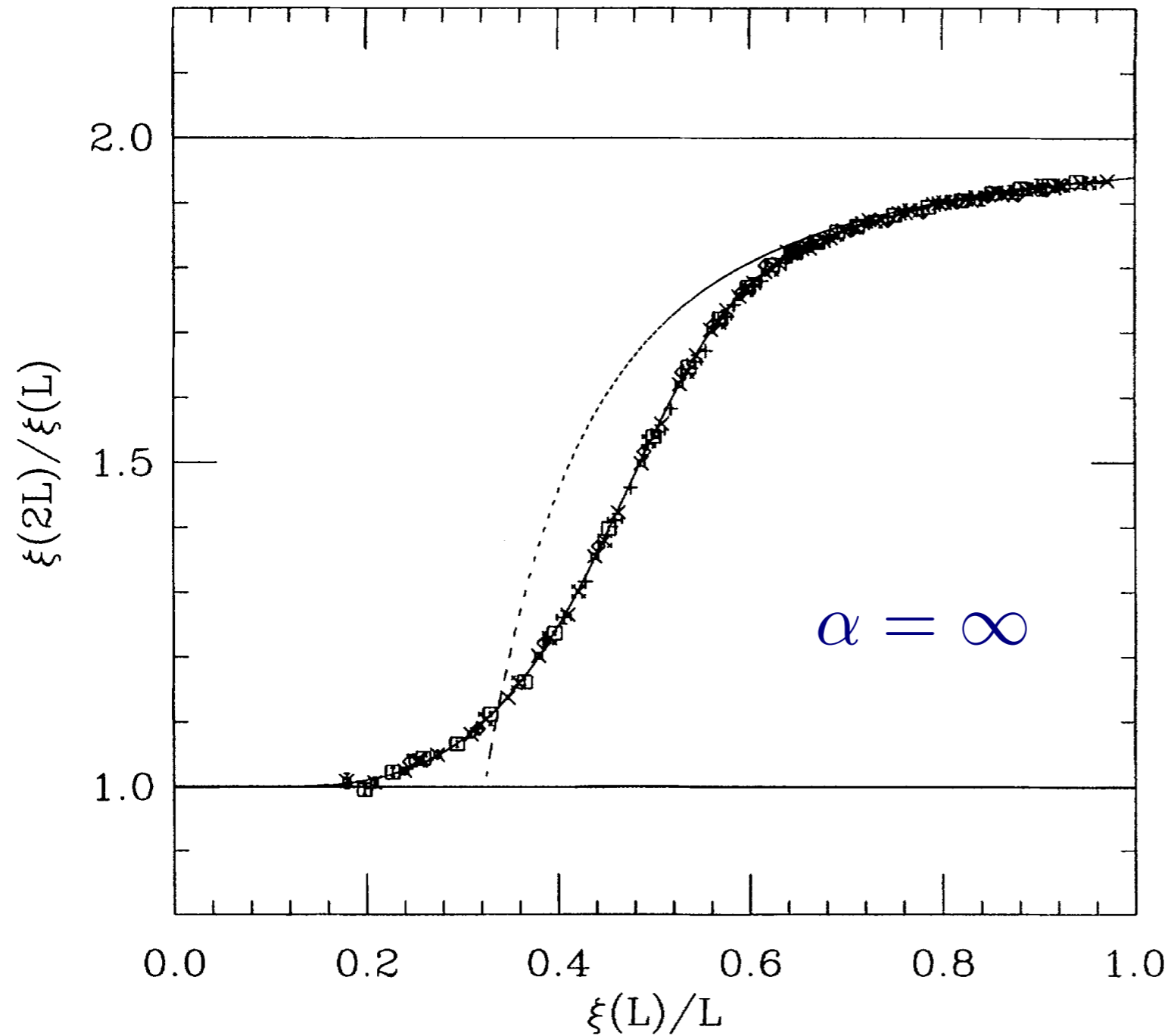
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At the critical point we get the desired free theory!

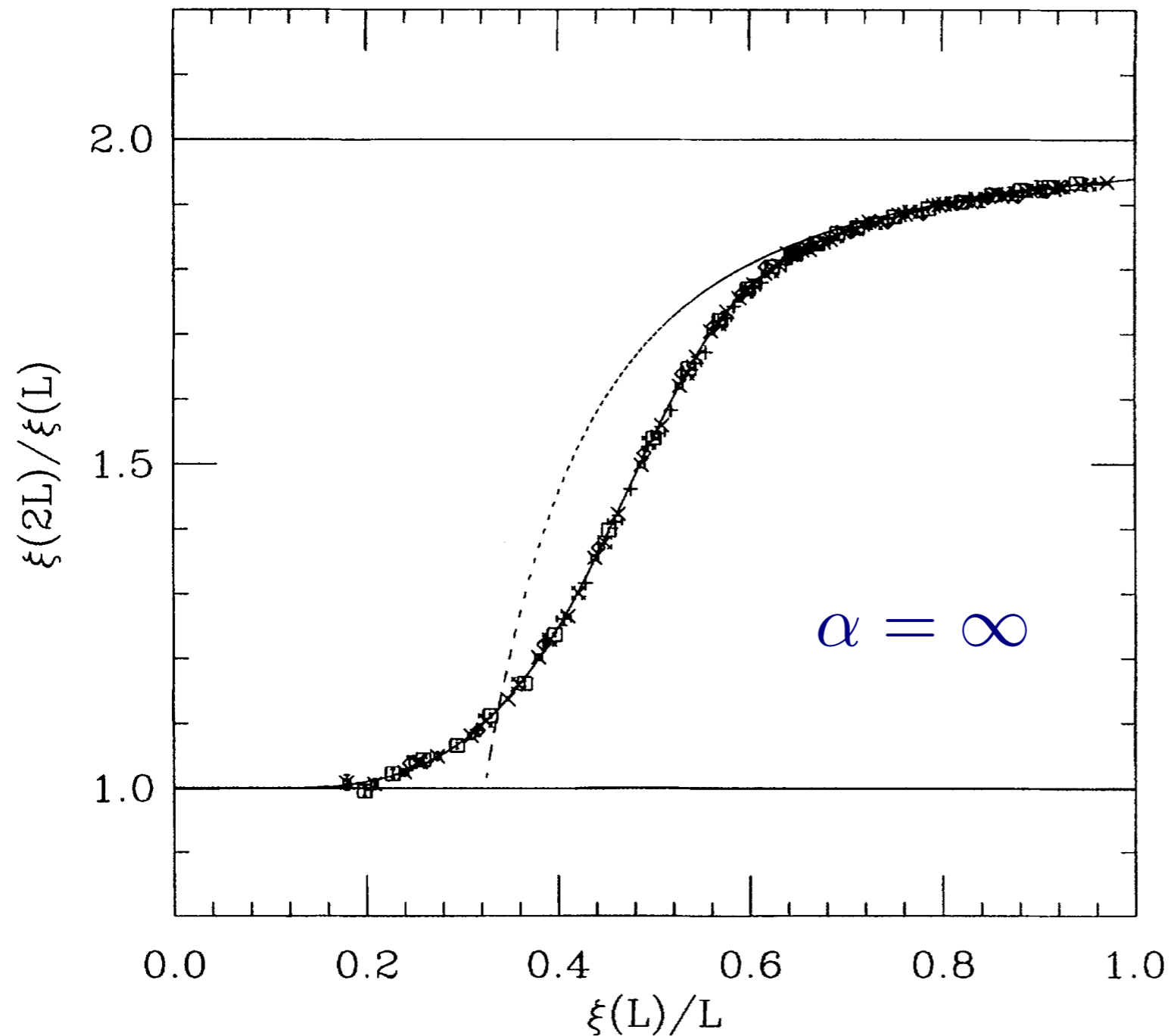
Continuum Physics: Universal step-scaling function

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Caracciolo, et.al., PRL 75, 1891 (1995)

Continuum Physics: Universal step-scaling function



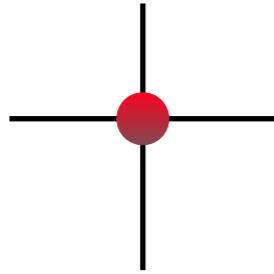
Can we reproduce this continuum physics in a lattice model with a finite Hilbert space?

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Qubit Regularization of $SO(3)$ fields

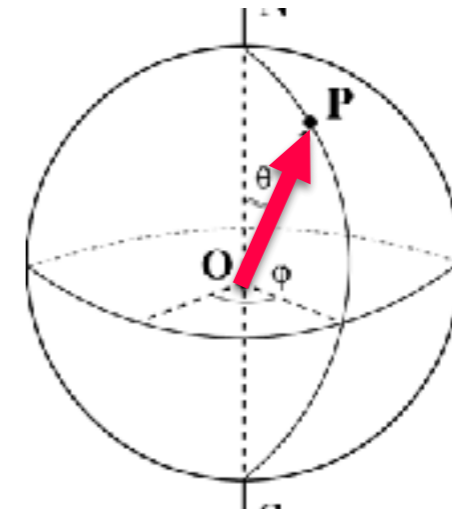
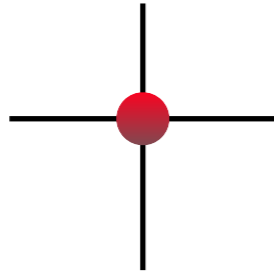
Qubit Regularization of $SO(3)$ fields

Local site Hilbert space
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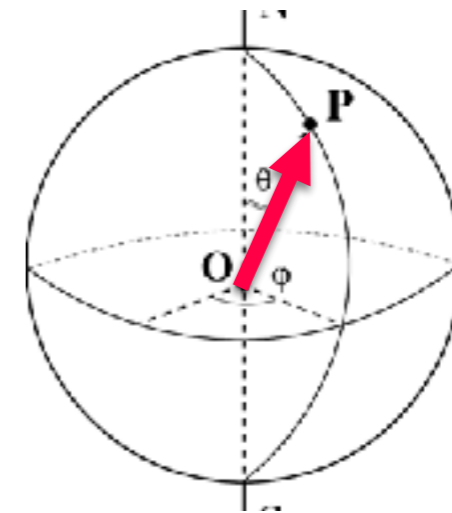
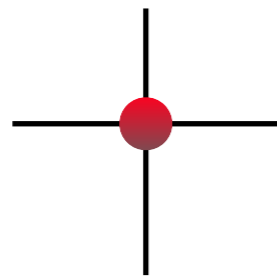


$$\vec{\phi}_x = (\phi_x^1, \phi_x^2, \phi_x^3)$$

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Basis of the traditional Hilbert space $\mathcal{H}_{\text{Trad}}$:

$$\int d\Omega |\theta, \varphi\rangle \langle \theta, \varphi| = 1$$

“position basis”

$$\sum_{l=0}^{\infty} \sum_{m=-l}^l |l, m\rangle \langle l, m| = 1$$

“momentum basis”

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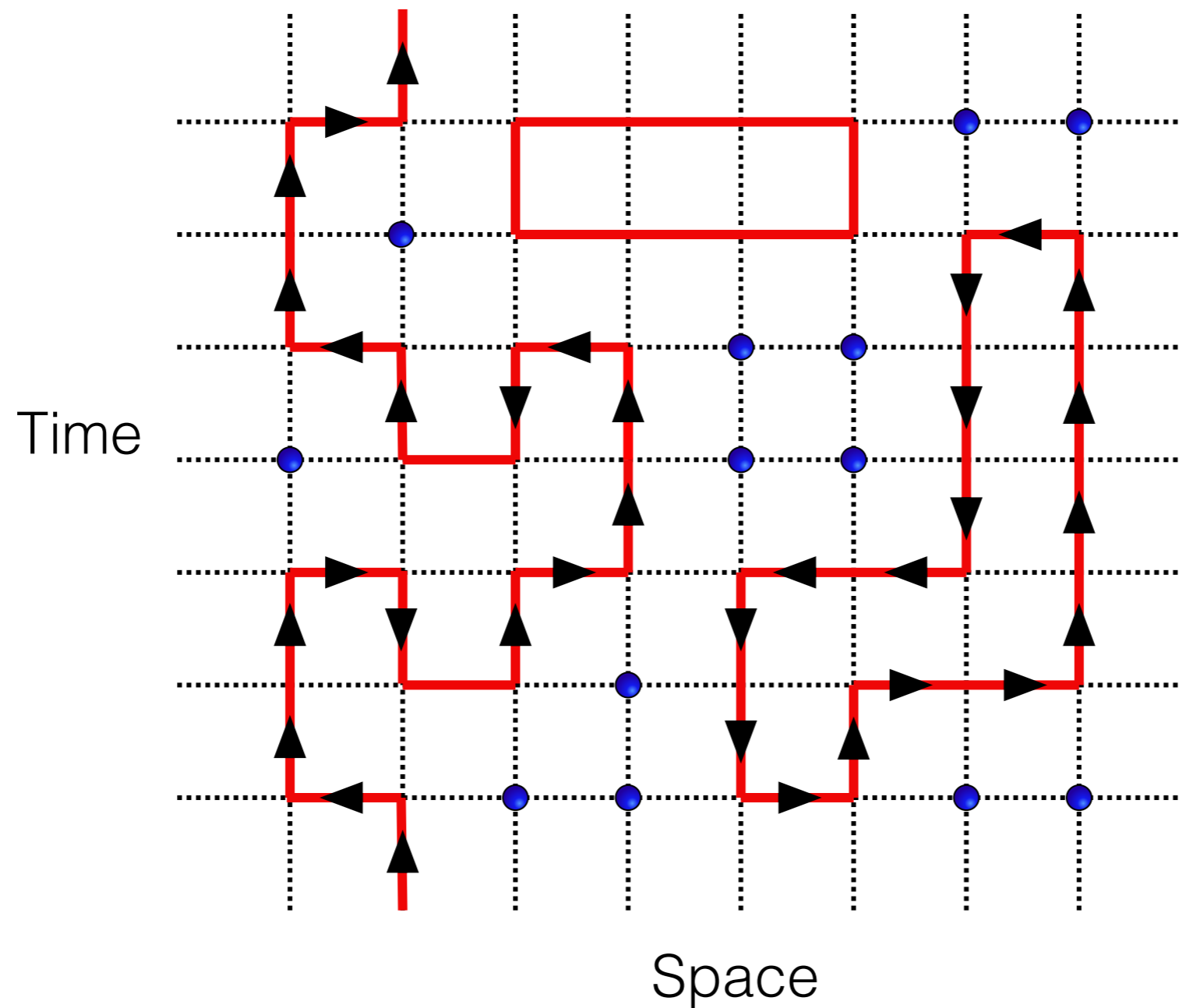
($\dim(\mathcal{H}_Q) = 4$) Two qubits per lattice site 

Space-time Euclidean configurations in the
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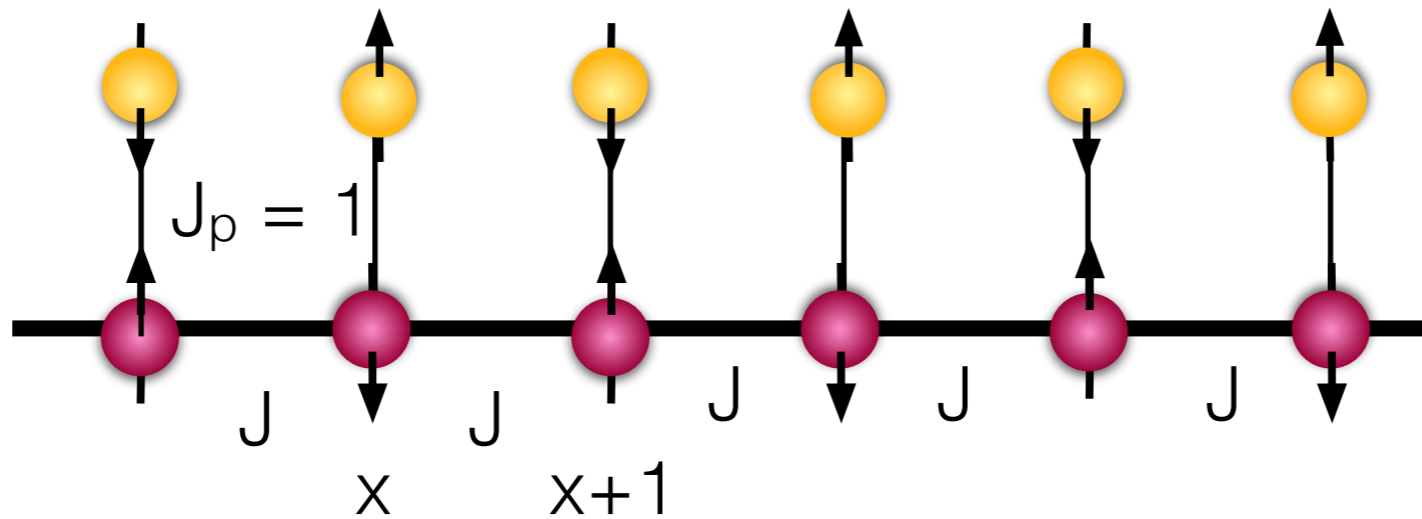
Qubit regularized model

Bhattacharya, Buser, SC, Gupta, Singh PRL 126 (2021), 172001

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Heisenberg-Comb

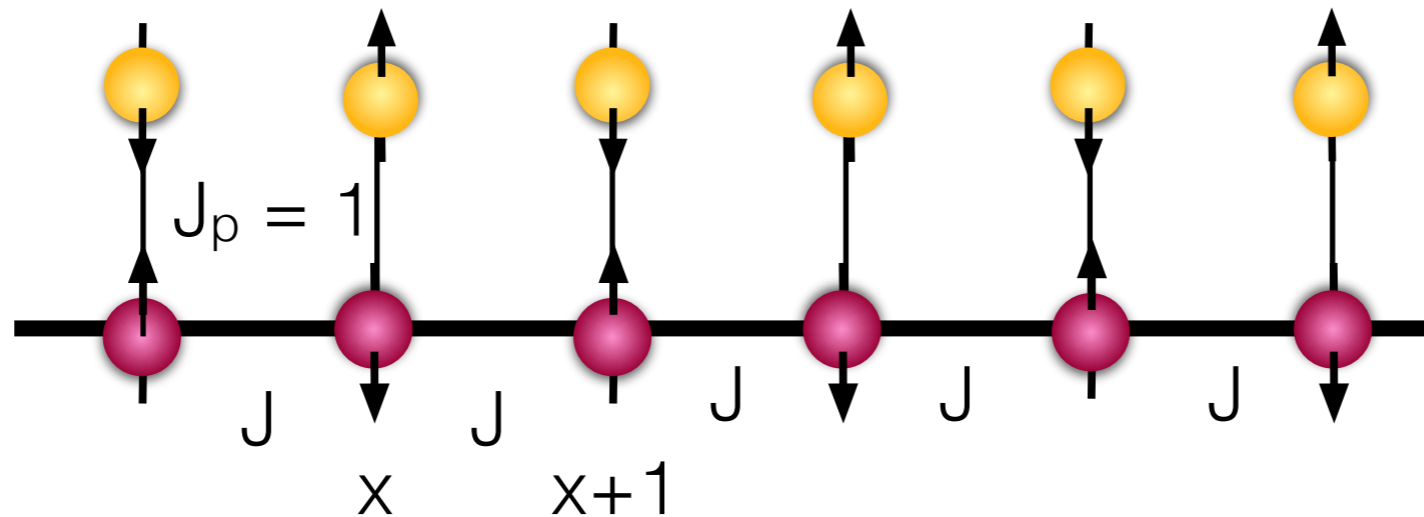


$$H = \sum_x J \mathbf{S}_{x,1} \cdot \mathbf{S}_{x+1,1} + \mathbf{S}_{x,1} \cdot \mathbf{S}_{x,2}$$

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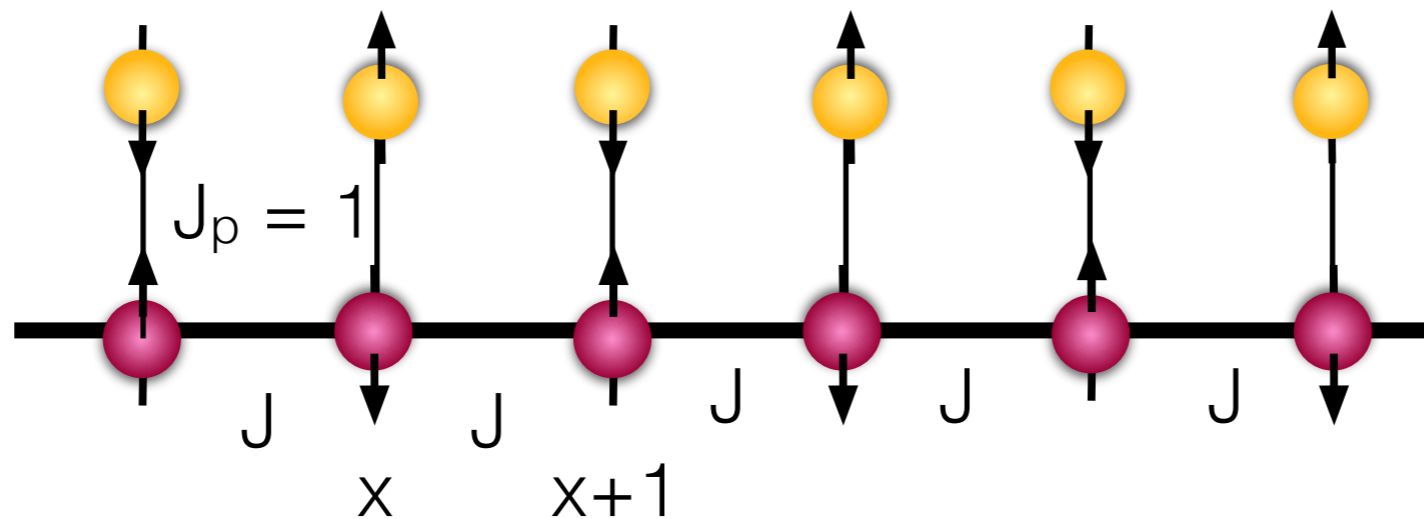
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Quantum Critical Point: $J \rightarrow \infty$ (Spin-1/2 Chain)

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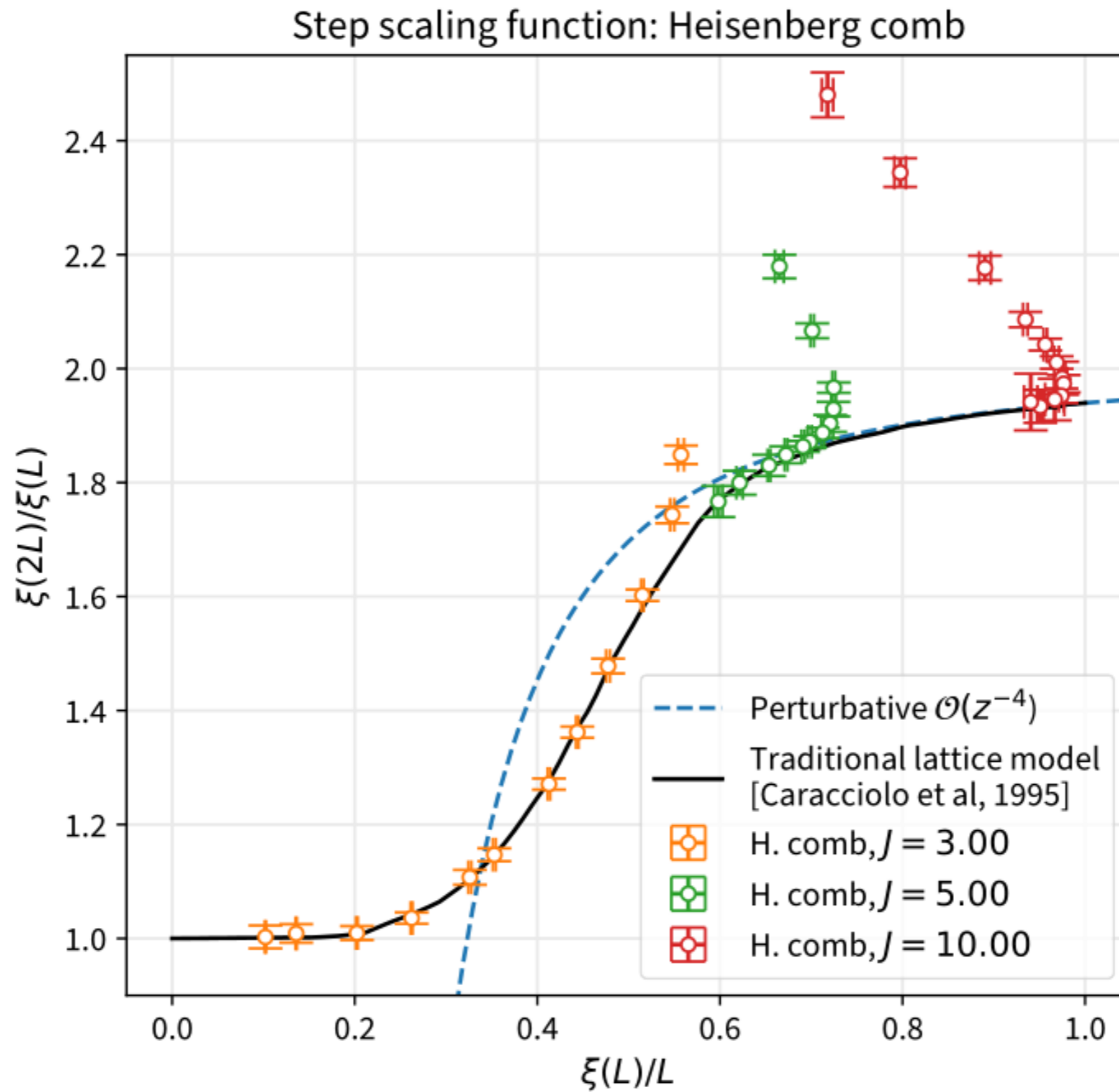
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At the critical point we get a decoupled critical theory that is not the desired free theory!

Universal Step Scaling Function

Universal Step Scaling Function



Traditional Model at

$$\beta \rightarrow \infty$$

$$\mathcal{H}_{\text{Trad}} = \textit{infty}$$

\equiv

Heisenberg Comb at

$$J \rightarrow \infty$$

$$\mathcal{H}_Q = 4$$

Example: BKT Transition

Lattice Action:

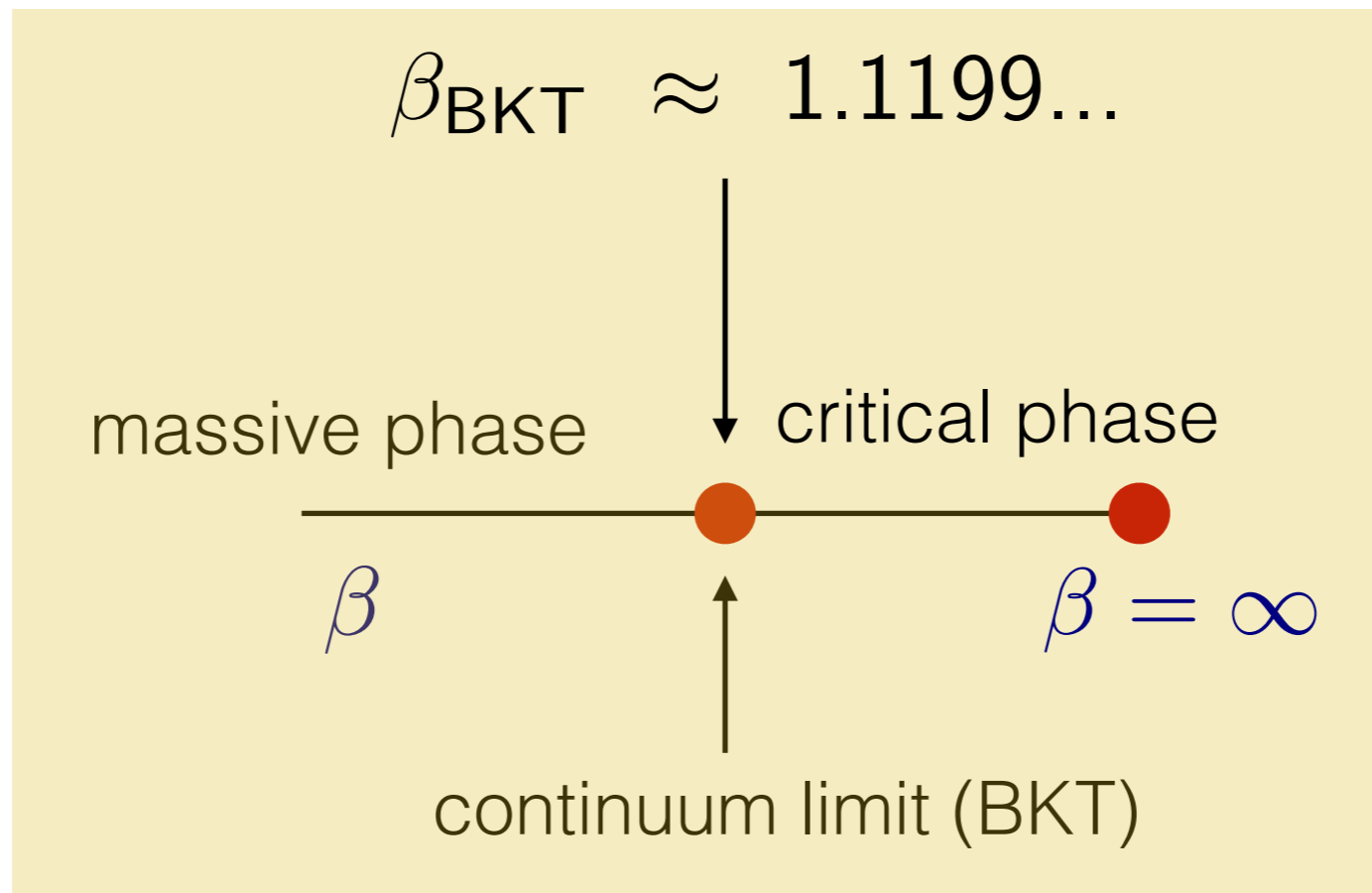
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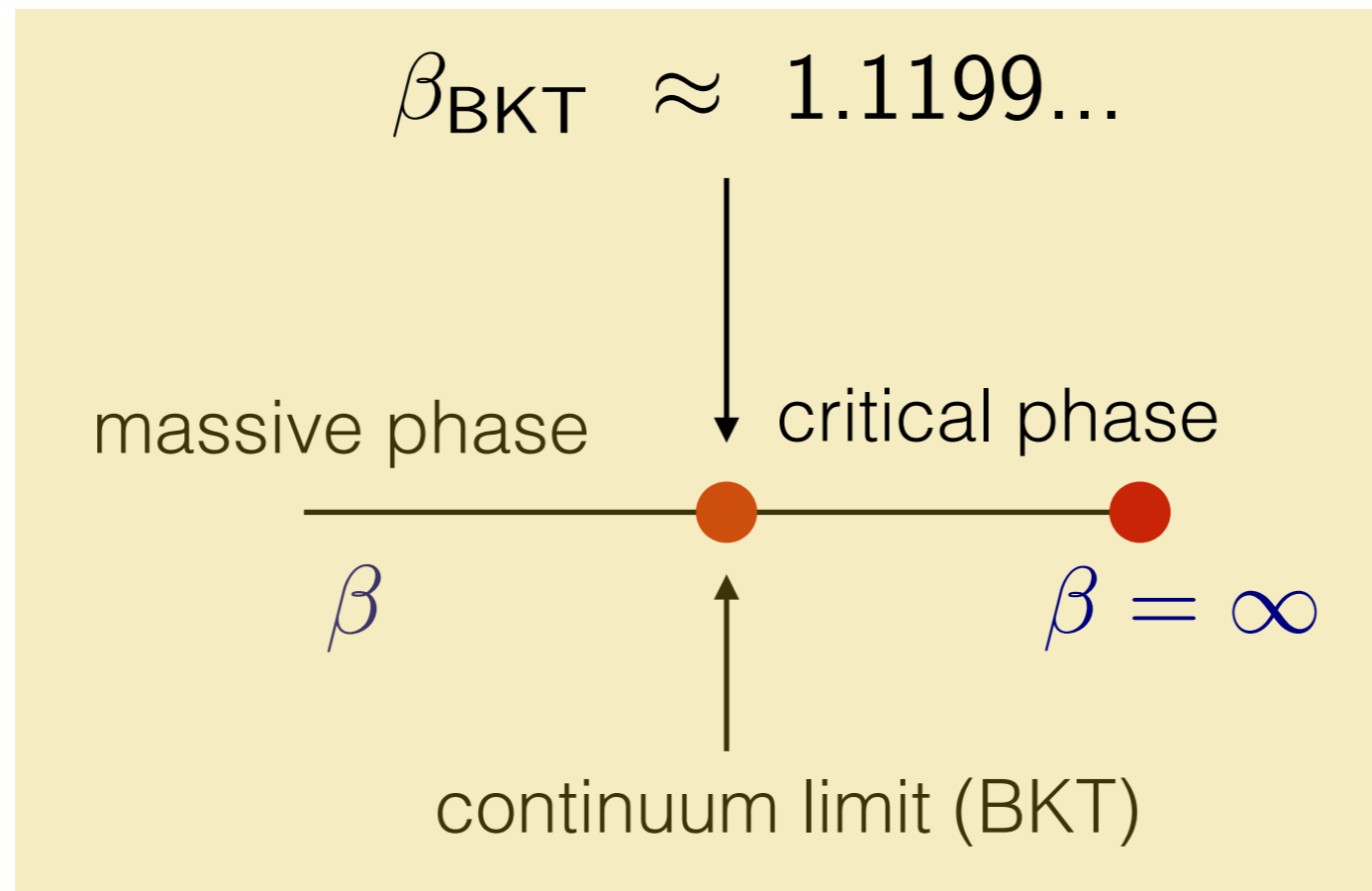
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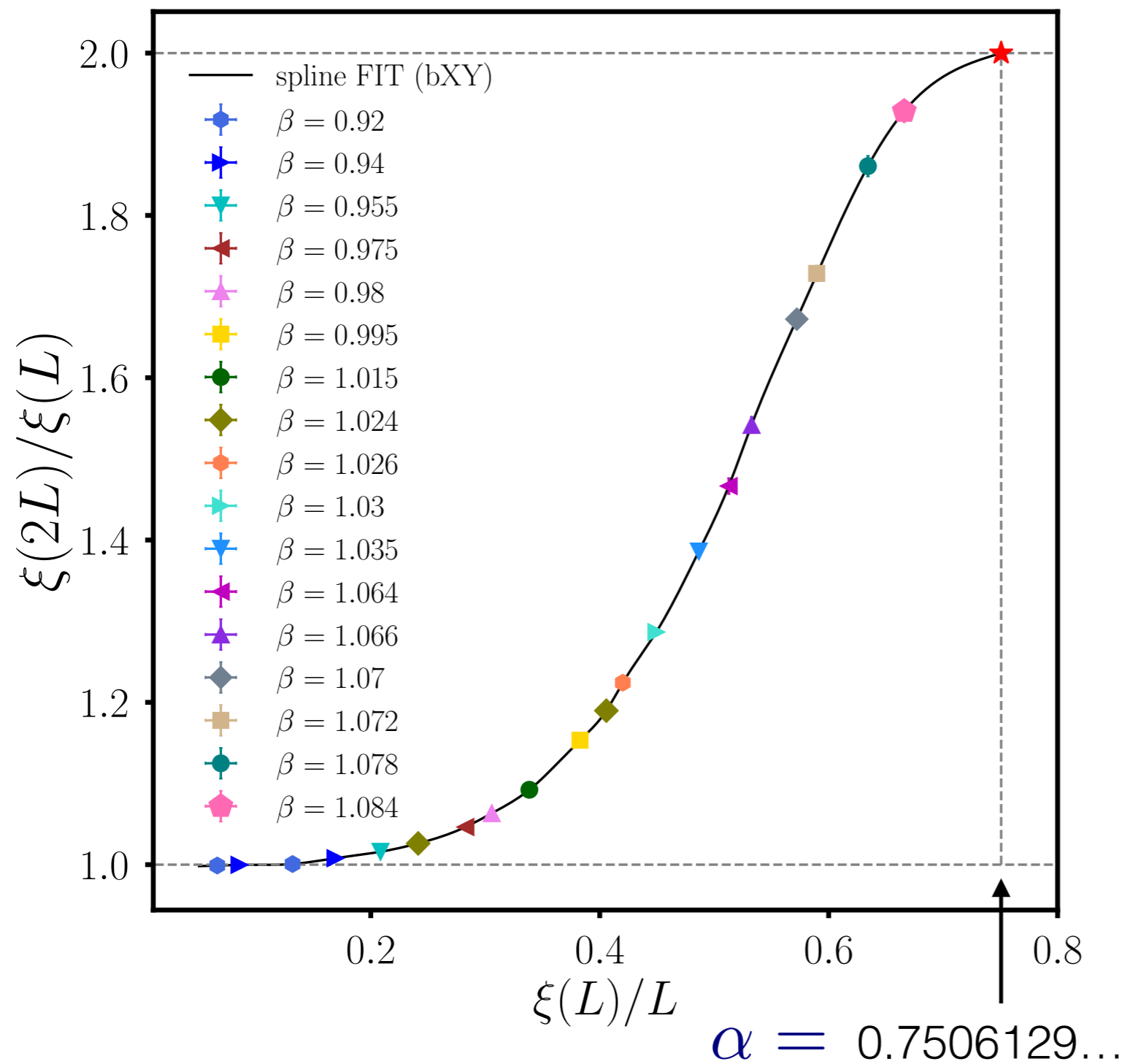
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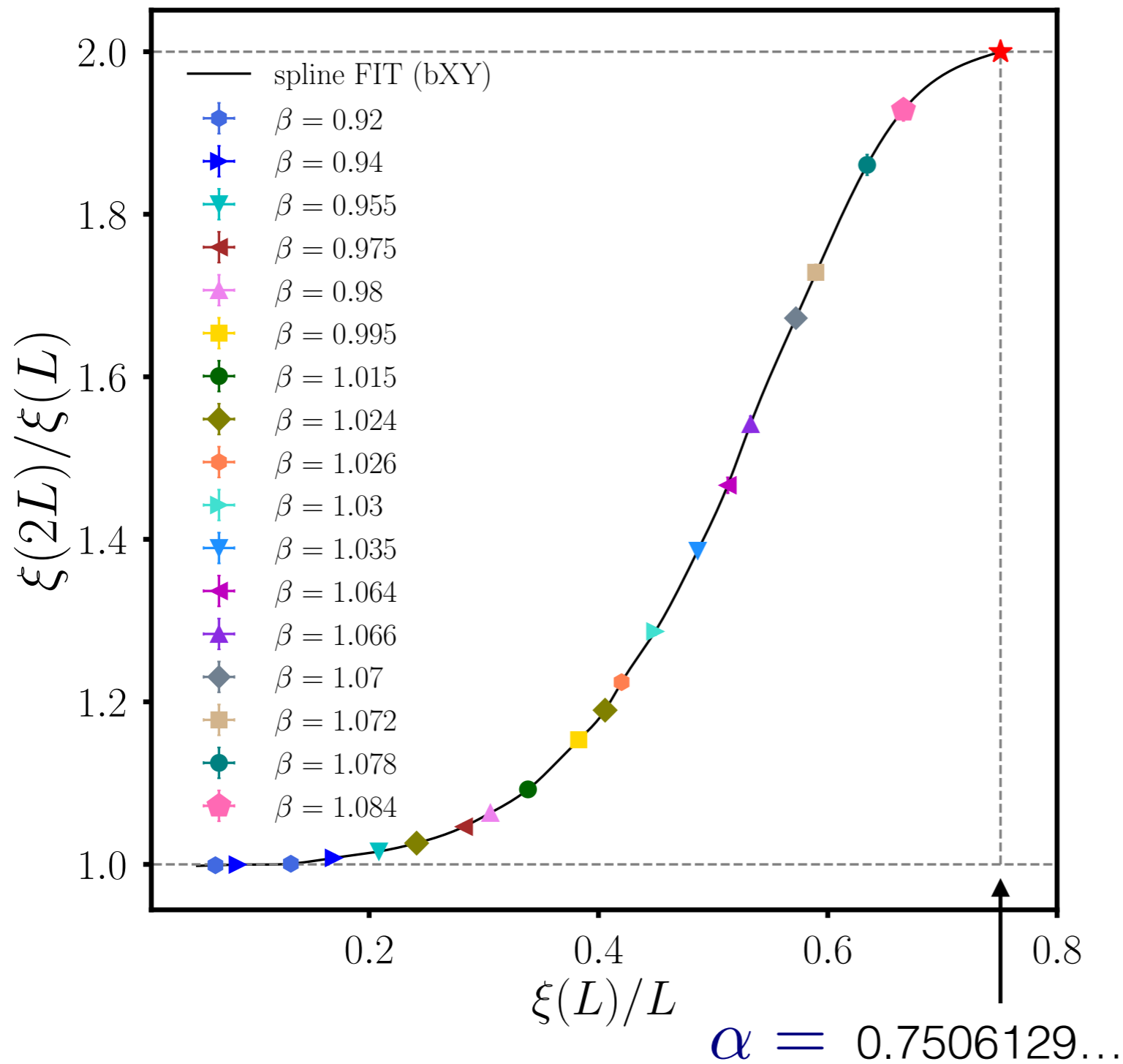
At the critical point we get a Gaussian theory!

The Step-Scaling function at the BKT transition using the traditional lattice model



Hasenbusch, cond-mat/0506552v2 (2008)

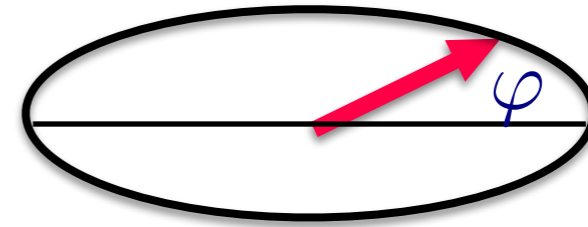
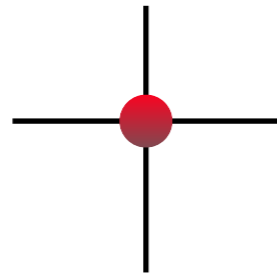
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Qubit Regularization of SO(2) fields

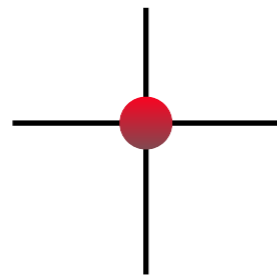
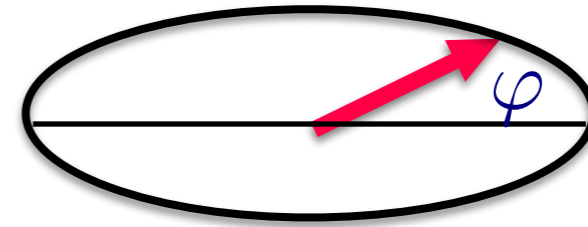
Local site Hilbert space describes a quantum particle on a circle



$$\vec{\phi}_x = (\cos(\varphi_x), \sin(\varphi_x))$$

Qubit Regularization of SO(2) fields

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A simple qubit regularization scheme is $Q = \{0, 0, 1, -1\}$

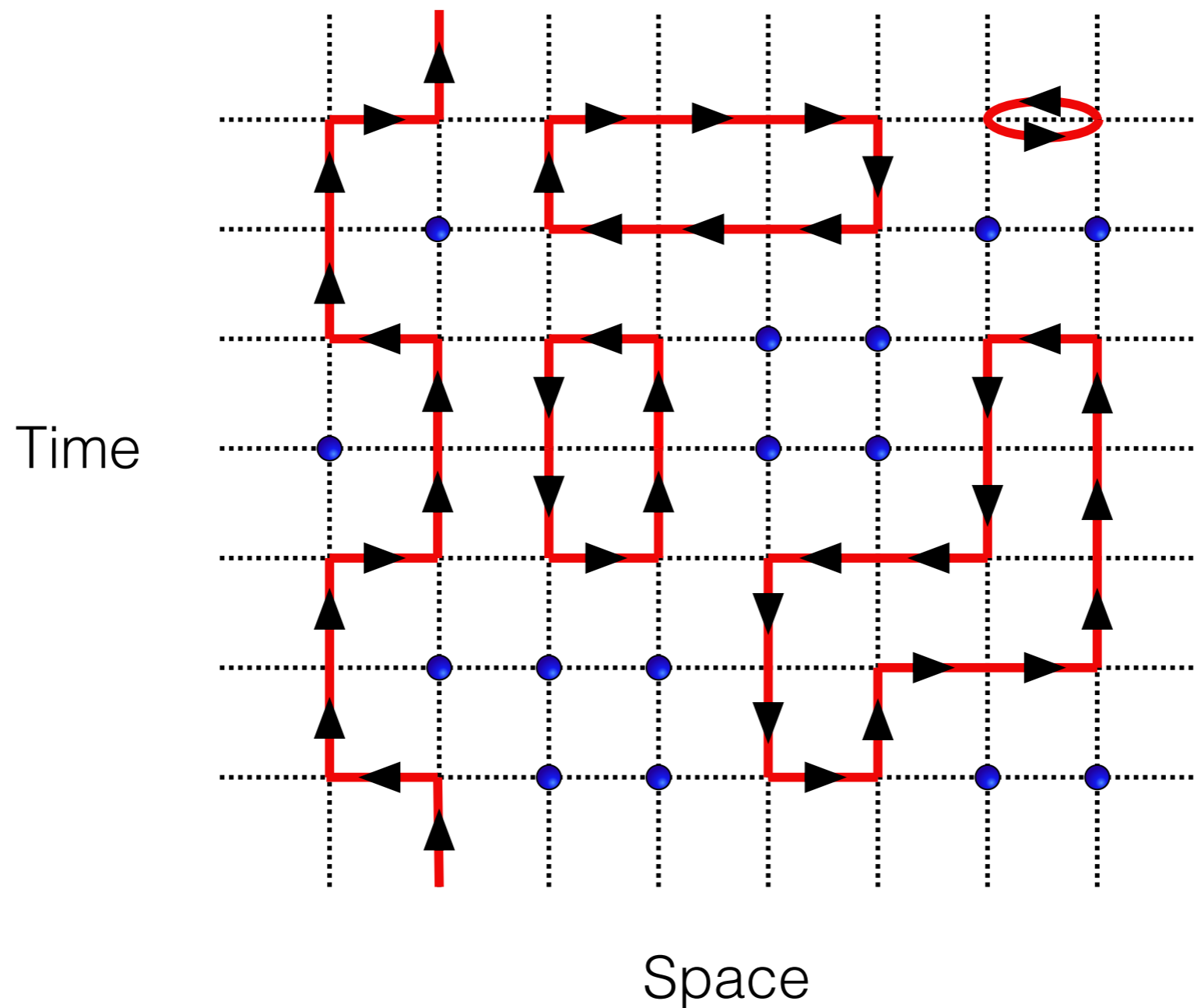
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Space-time Euclidean configurations in the
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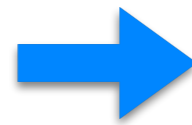
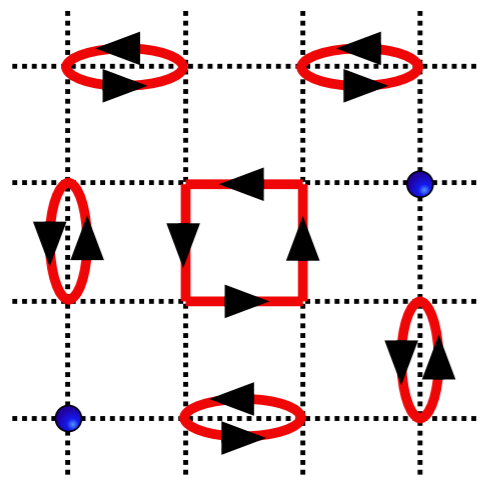
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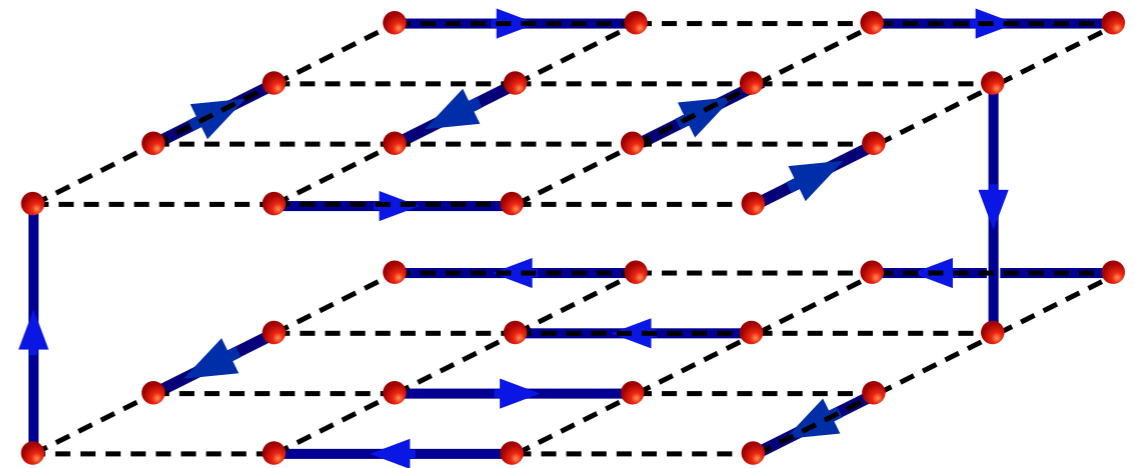
Mapping to two layers of closed pack dimer model

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worldline configuration

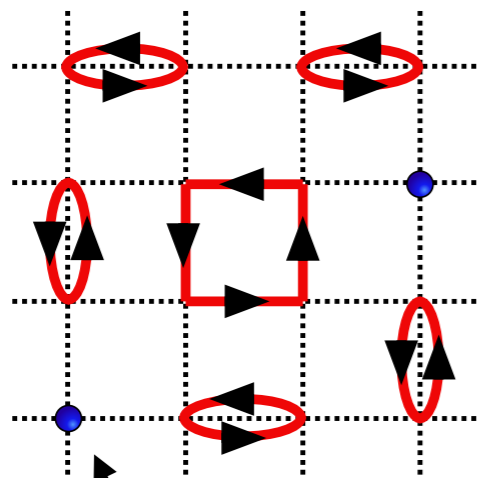


closed packed dimer configurations

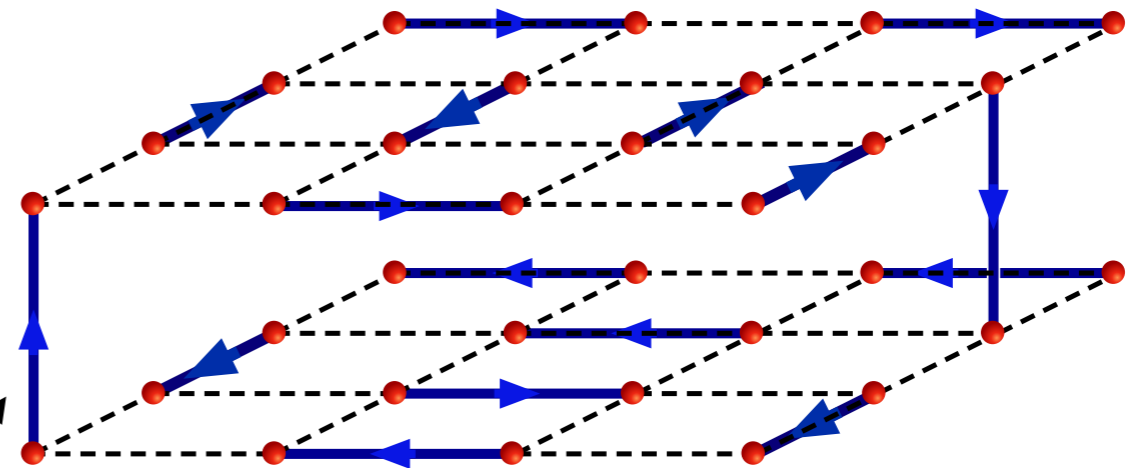


Mapping to two layers of closed pack dimer model

worldline configuration



closed packed dimer configurations



Singlets are mapped to inter-layer dimers

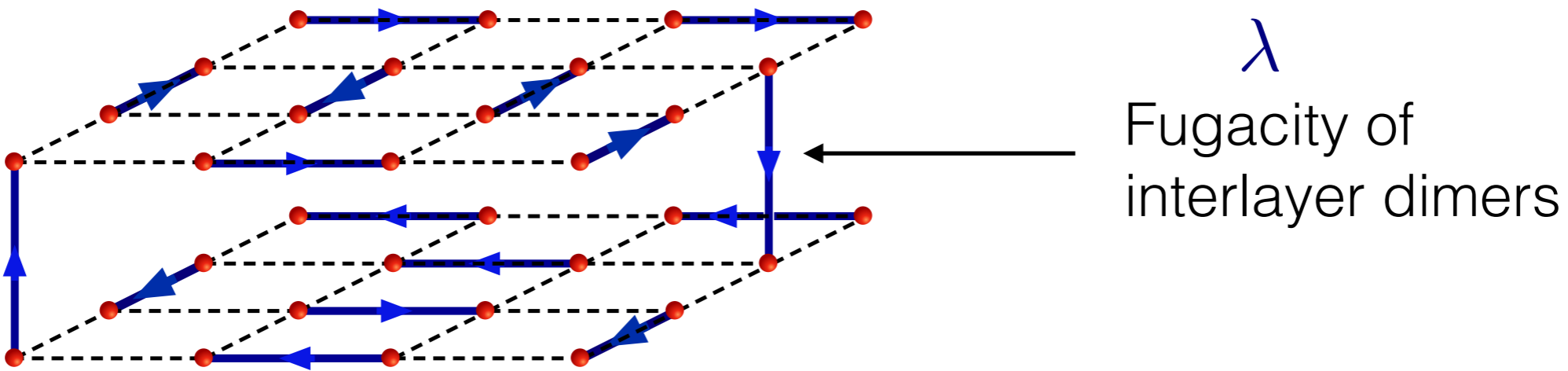
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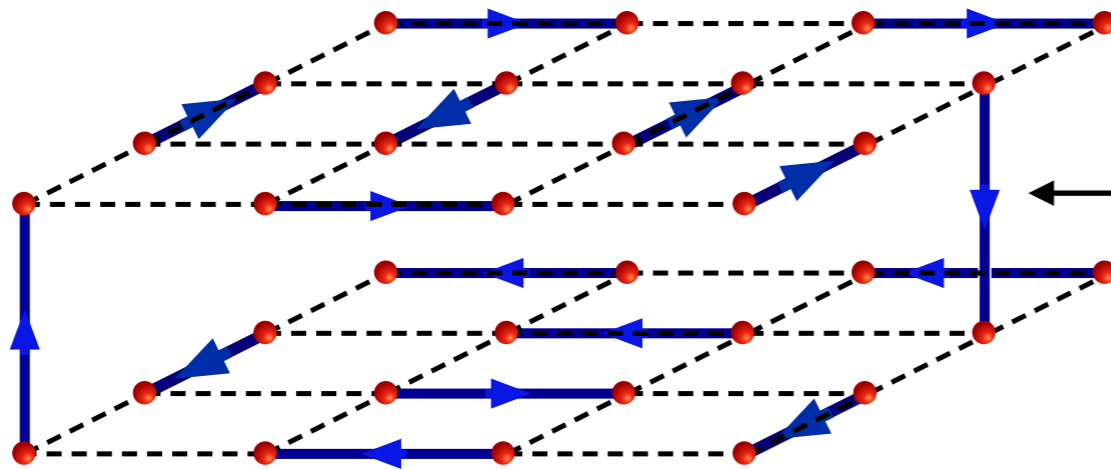
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closed packed dimer configurations



λ
Fugacity of
interlayer dimers

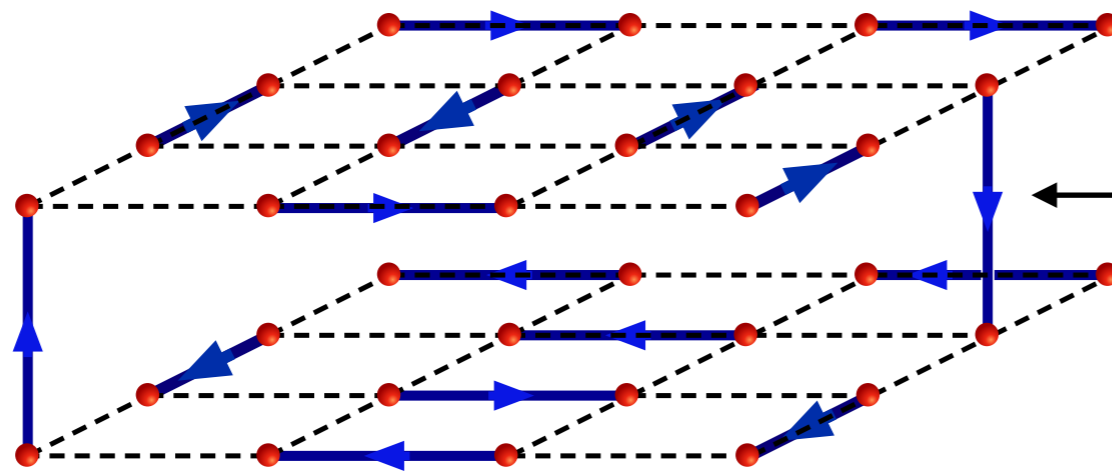
Partition function

$$Z = \sum_C \lambda^{N_I}$$

Qubit regularized model

Maiti, Banerjee, SC, Marinkovic PRL 132 (2024), 041601

closed packed dimer configurations



λ
Fugacity of
interlayer dimers

Partition function

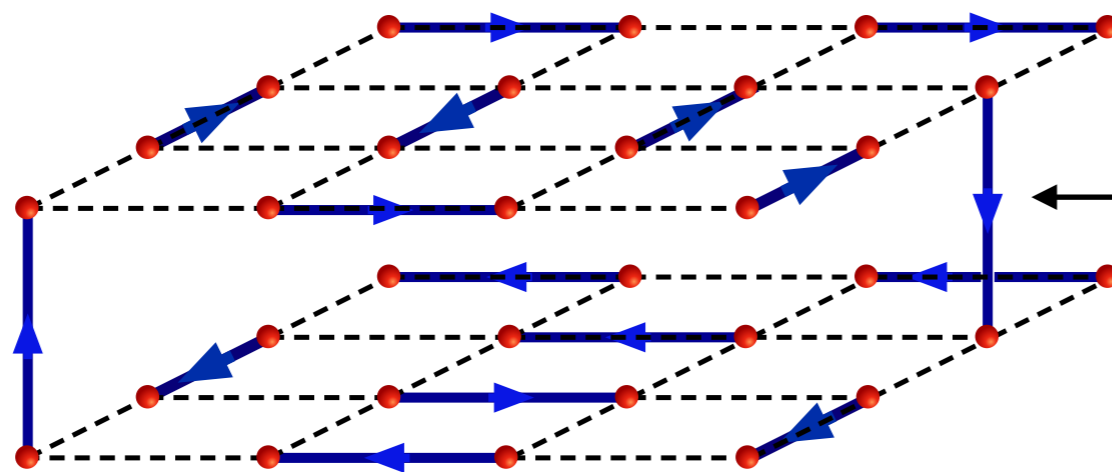
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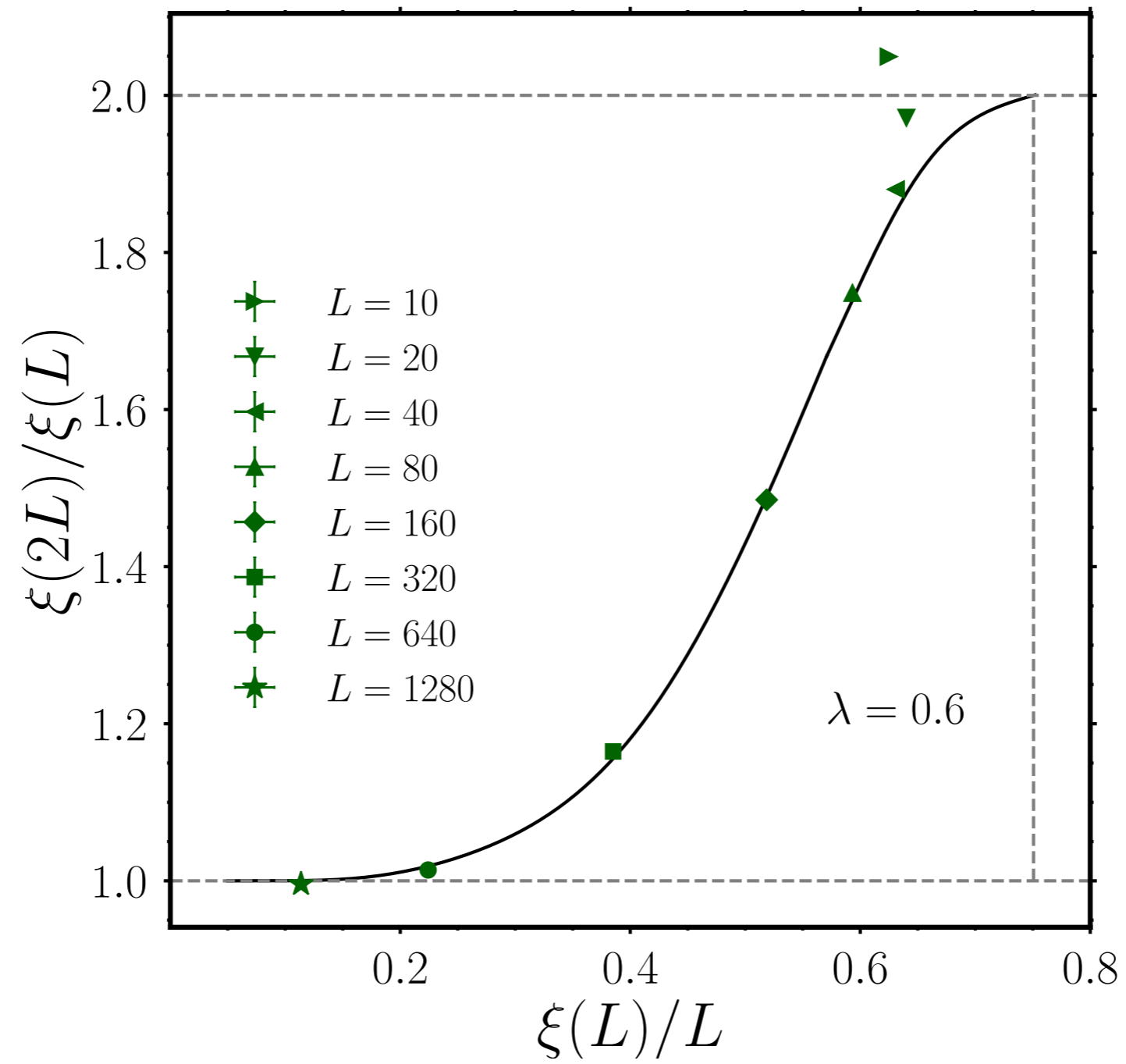
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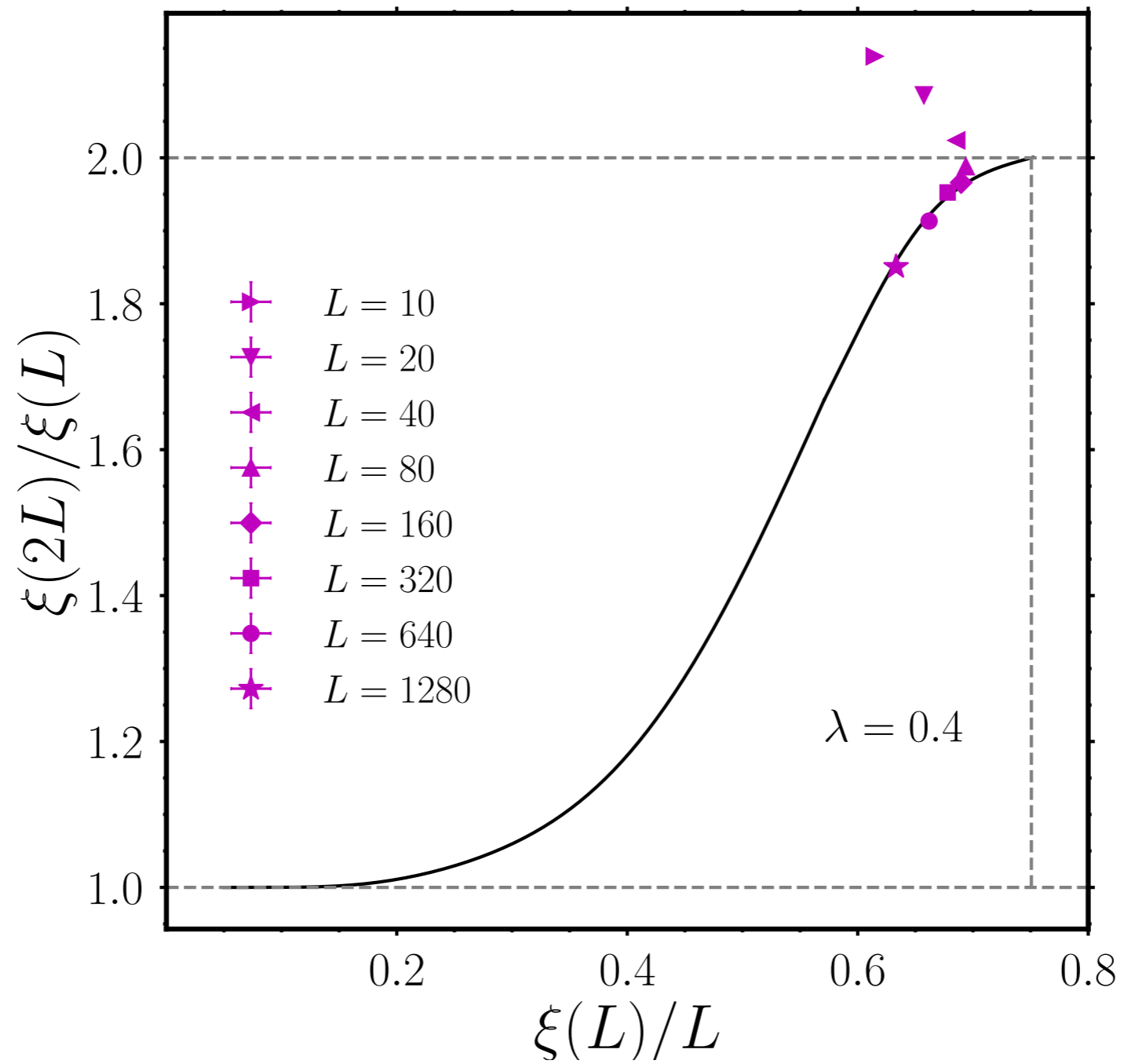
Continuum Limit: $\lambda \rightarrow 0$

At the critical point we get a decoupled critical theory that is again not the desired free theory!

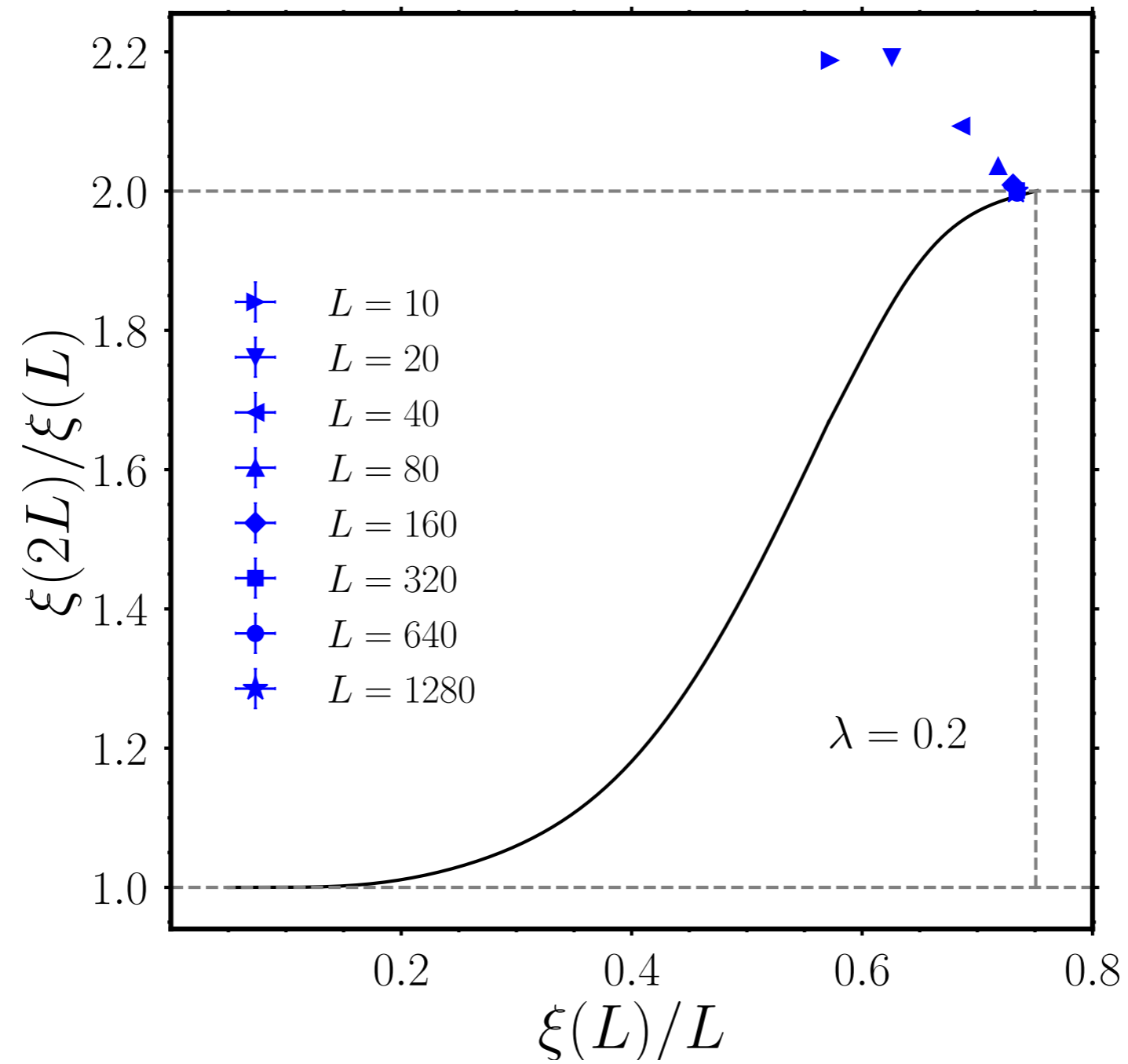
Step-scaling function of the qubit model



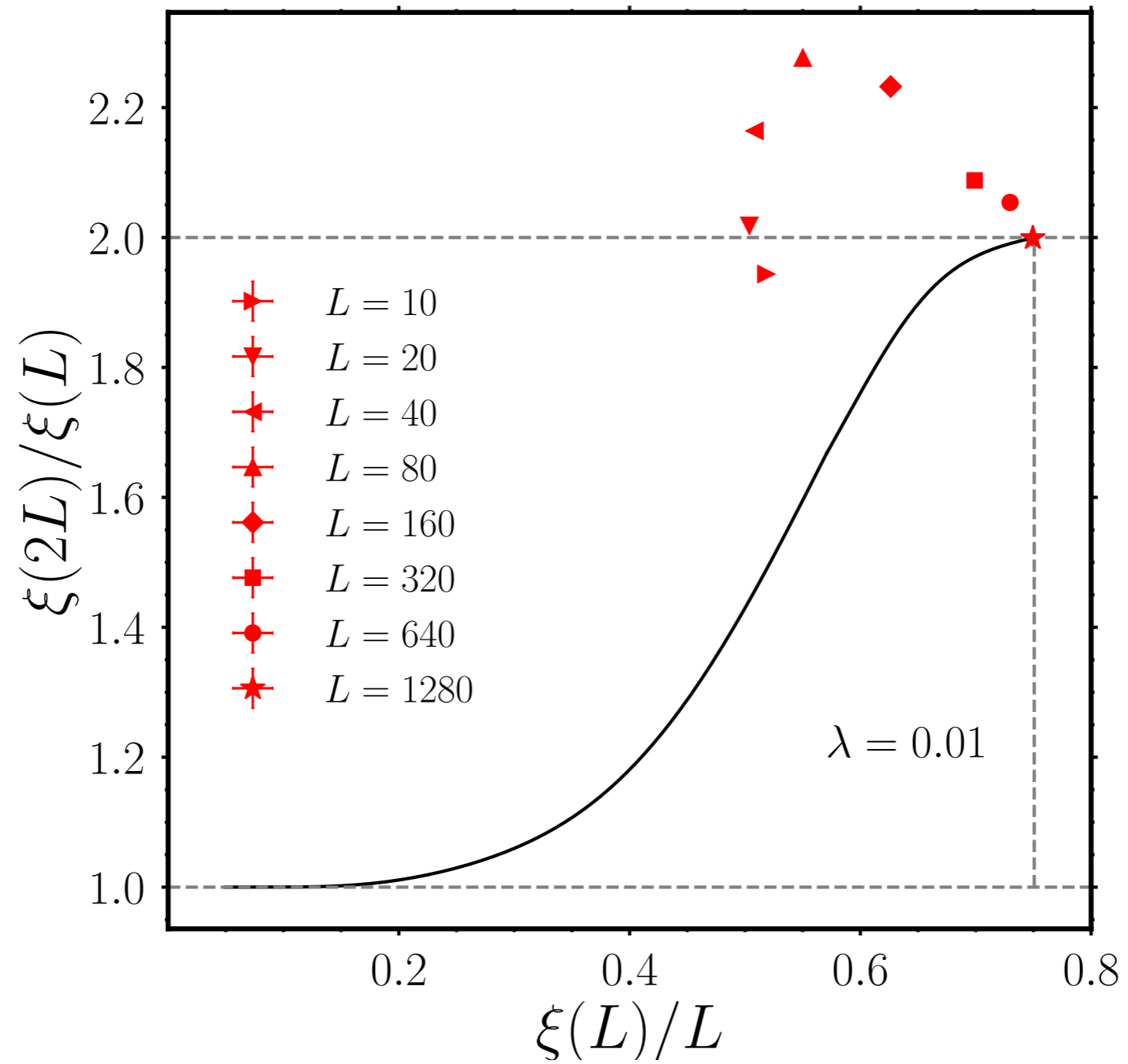
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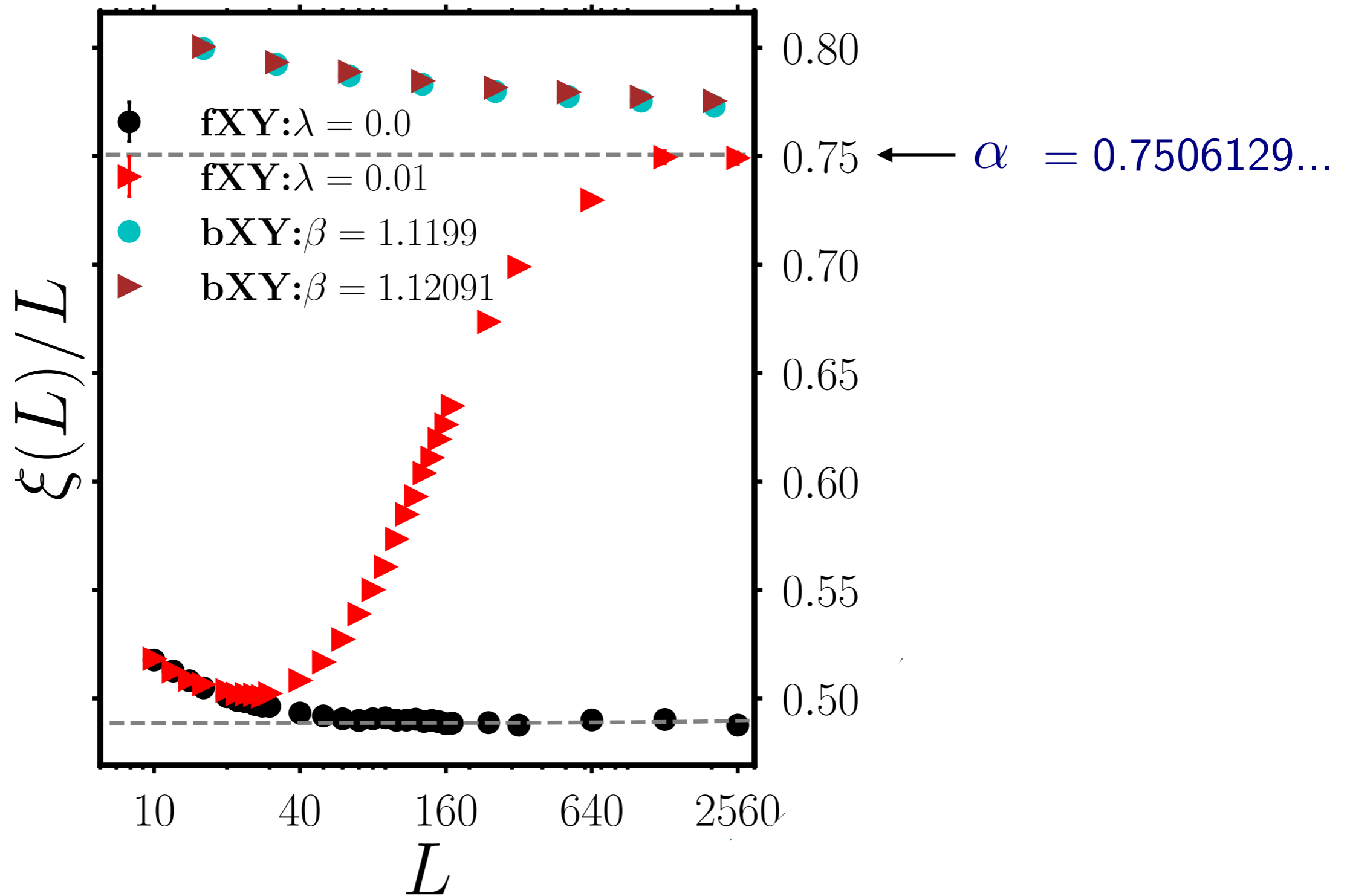
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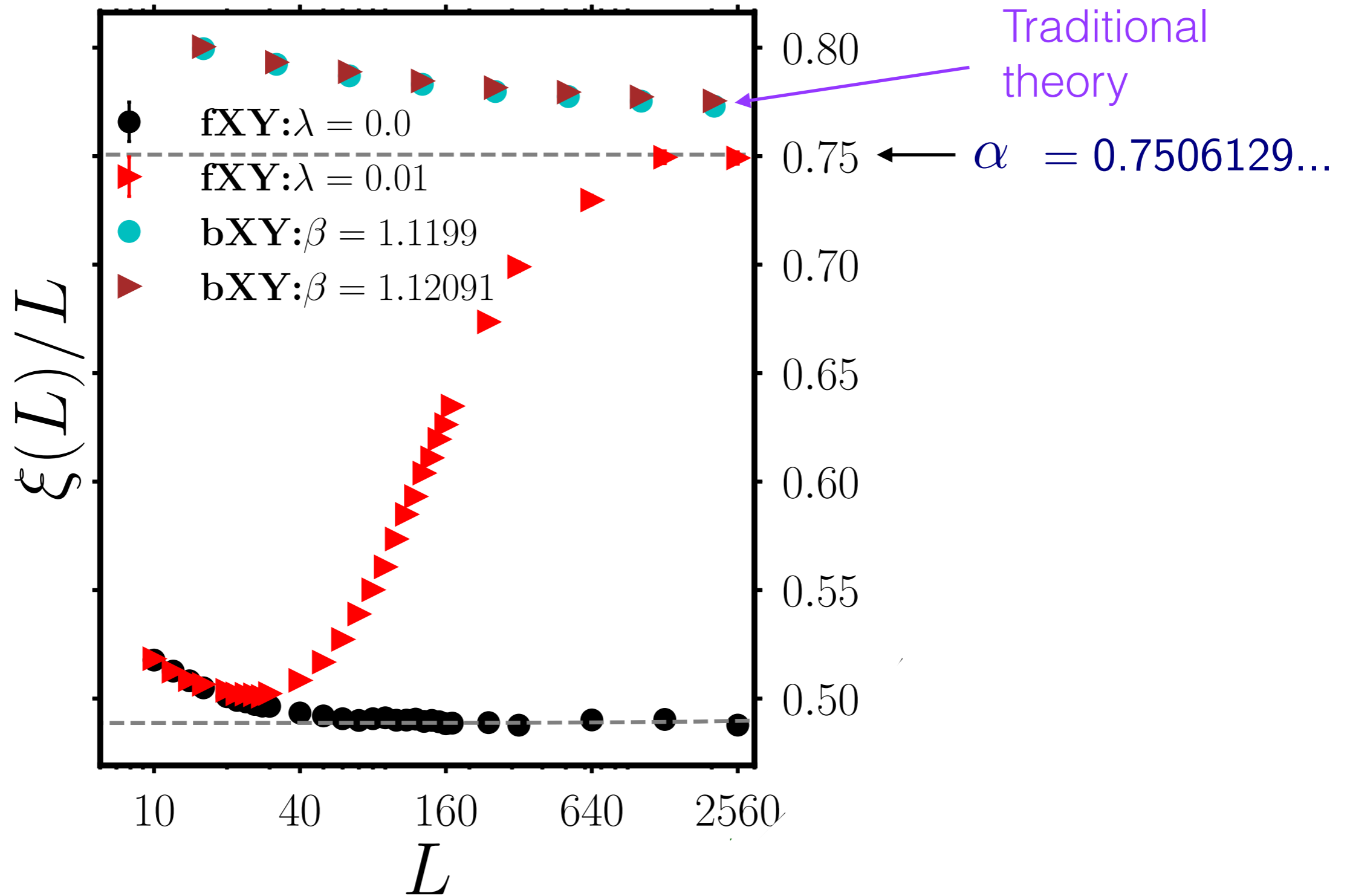
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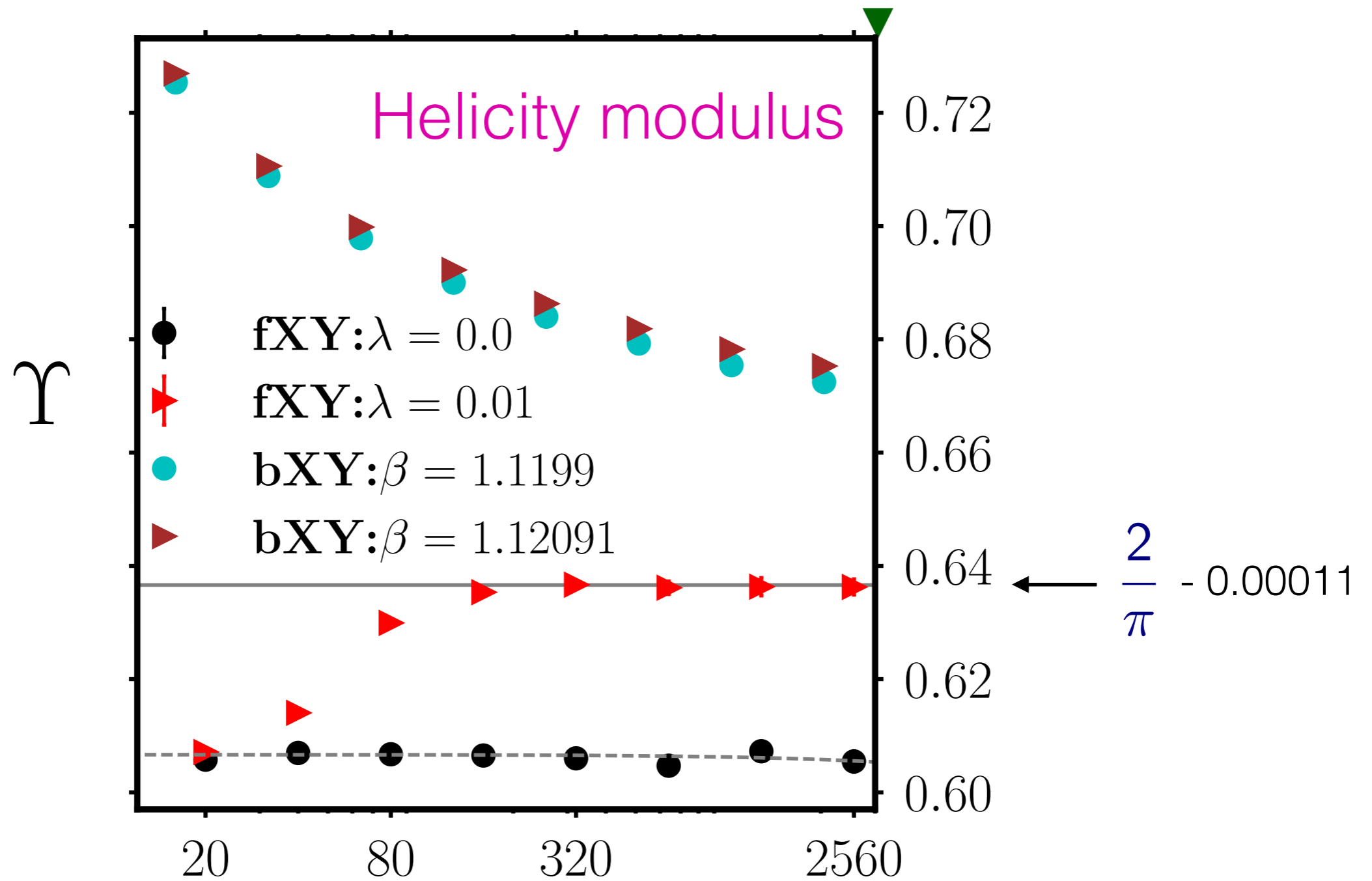
Decoupled vs. Coupled Theory



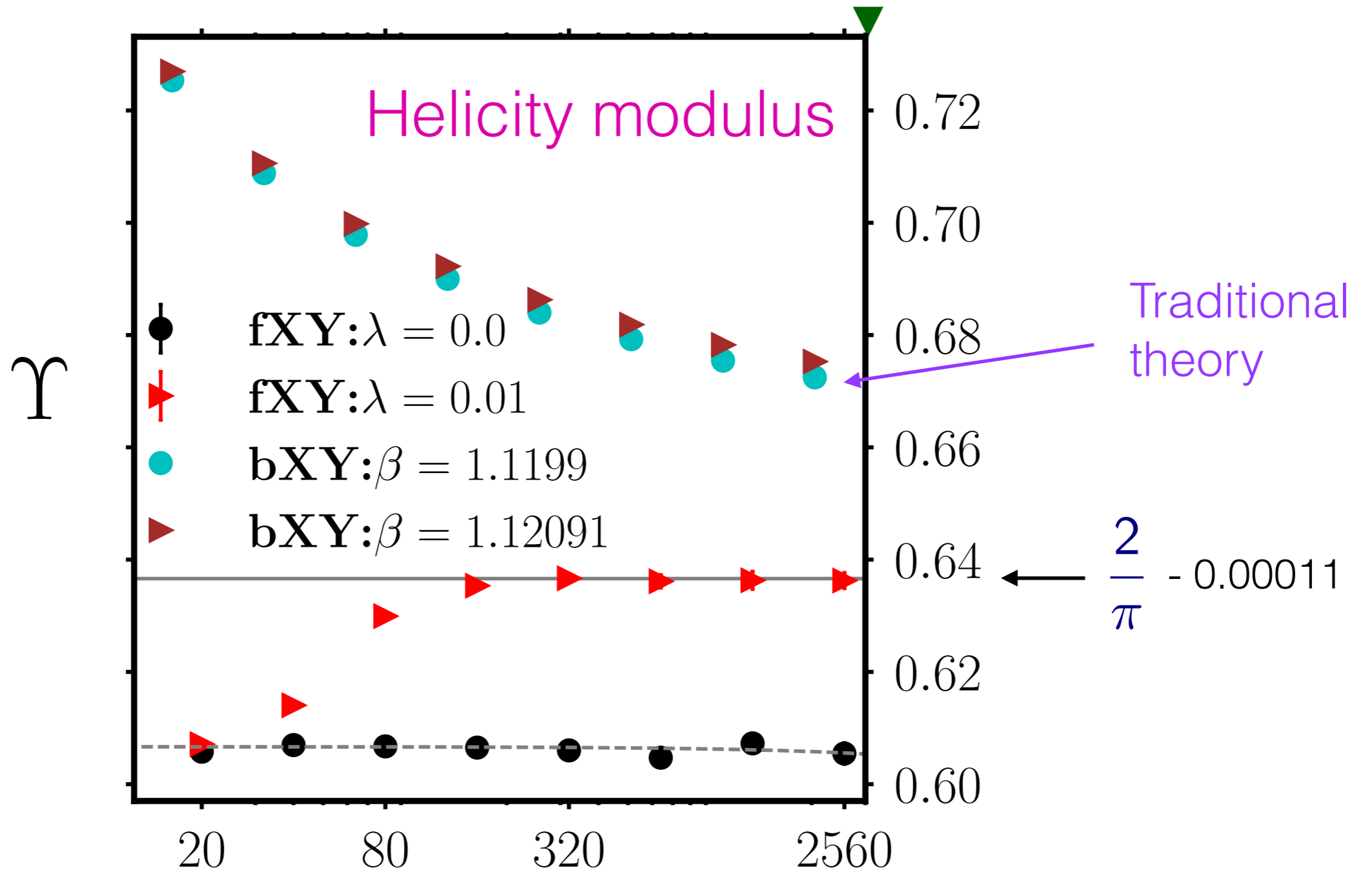
Decoupled vs. Coupled Theory



Decoupled vs. Coupled Theory



Decoupled vs. Coupled Theory



Massive Phase of the XY model
at the BKT transition

$$\mathcal{H}_{\text{Trad}} = \text{infty}$$

≡

Two layers of coupled
closed pack dimers

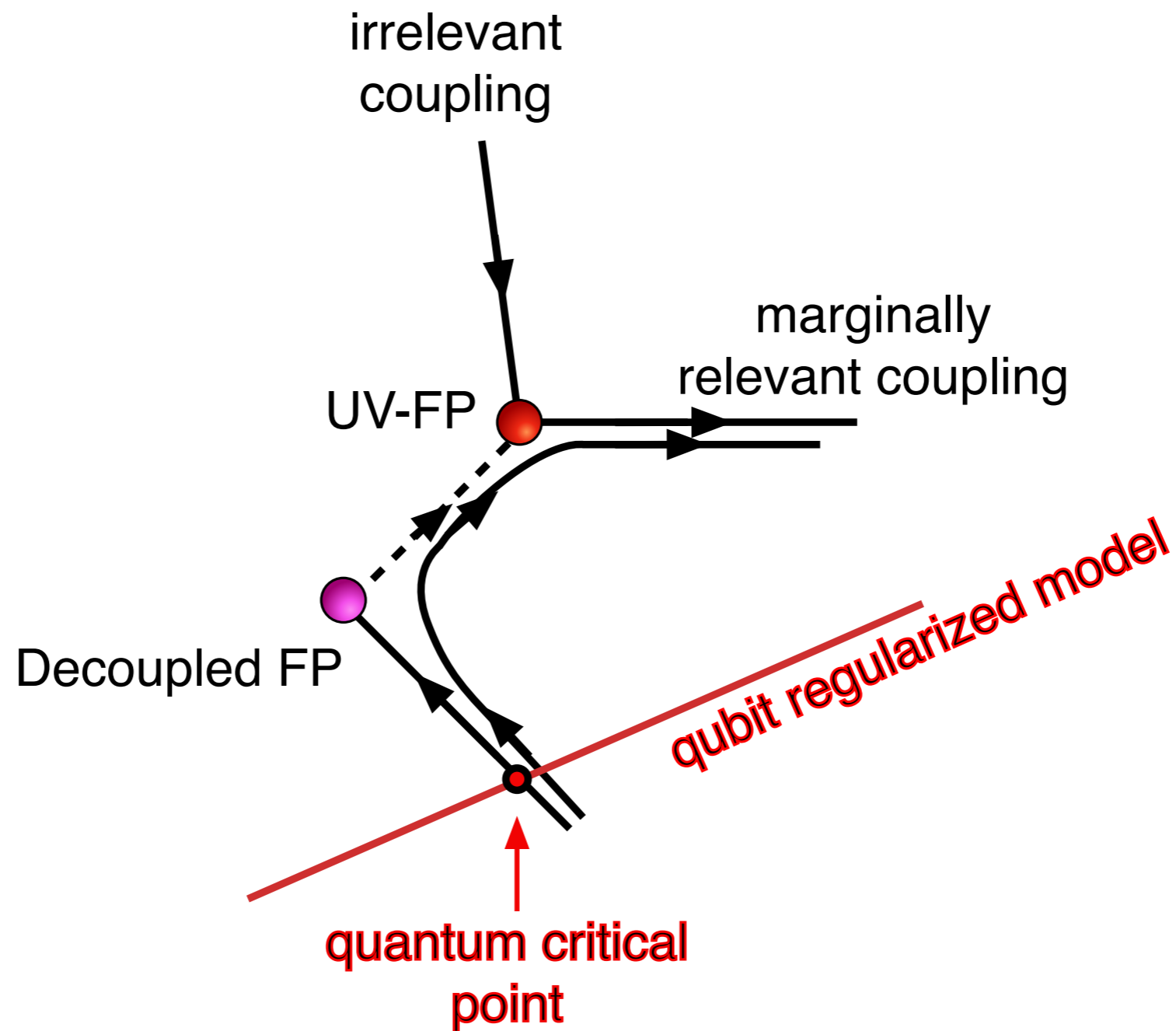
$$\lambda \rightarrow 0$$

$$\mathcal{H}_{\text{Q}} = 4$$

A Novel-RG flow for Asymptotic Freedom

In both these examples, asymptotic freedom
is recovered via new type of RG flow

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Qubit Regularization of Asymptotic Freedom in Gauge Theories

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Asymptotic Freedom of Yang Mills theory



Deconfined
phase at high
temperatures

Confined massive
phase at zero
temperatures

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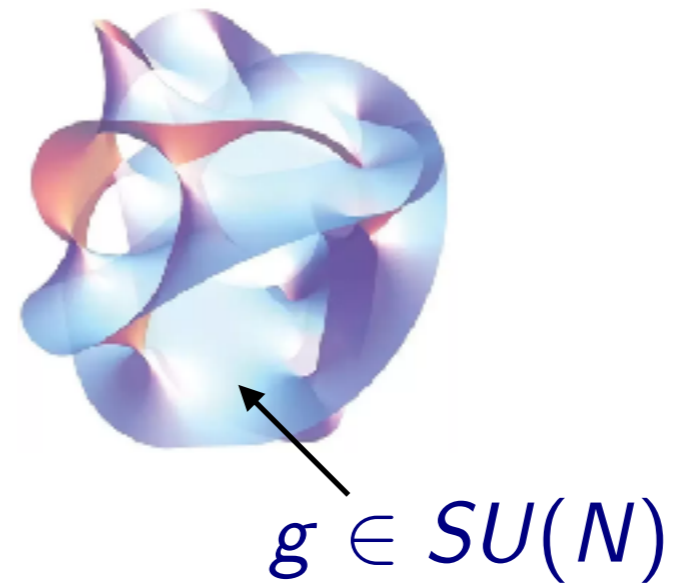
Confined massive
phase at zero
temperatures

Classical qubit models must already show this!

Qubit Regularization of $SU(N)$ gauge fields

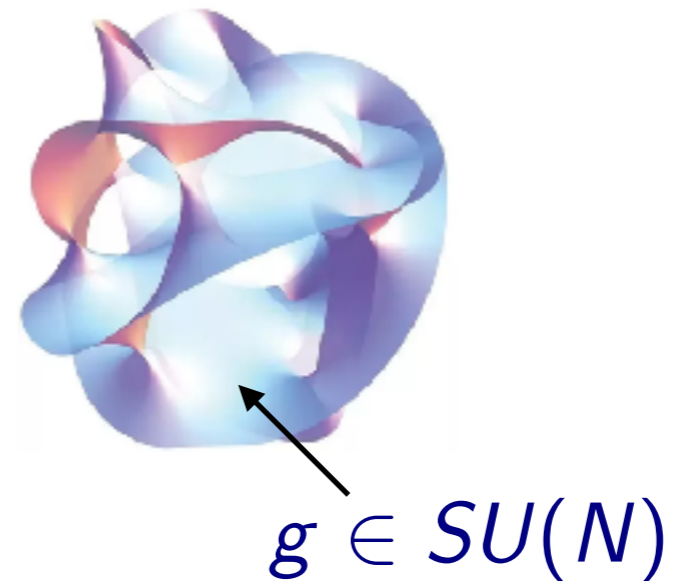
Qubit Regularization of $SU(N)$ gauge fields

Local link Hilbert space describes a quantum particle on the surface of the $SU(N)$ group manifold



Qubit Regularization of SU(N) gauge fields

Local link Hilbert space describes a quantum particle on the surface of the SU(N) group manifold



Basis of the full Hilbert space $\mathcal{H}_{\text{Trad}}$:

$$\int [dg] |g\rangle\langle g| = I$$

“position basis”

$$\sum_{\lambda} \sum_{i,j} |D_{ij}^{\lambda}\rangle\langle D_{ij}^{\lambda}| = \mathbb{I}$$

“momentum basis”

λ labels distinct irreps of SU(N)

This means the traditional link Hilbert space is given by

$$\mathcal{H}_{\text{Trad}} = \bigoplus_{\lambda} \mathcal{H}_{\lambda} \quad \leftarrow \text{Peter-Weyl Theorem}$$

where $\mathcal{H}_{\lambda} = V_{\lambda} \otimes V_{\lambda}^*$ is spanned by $\{|D_{ij}^{\lambda}\rangle\}$, $i, j = 1, 2, \dots, d_{\lambda}$

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Qubit Regularized Hilbert Space $\mathcal{H}_Q = \bigoplus_{\lambda \in Q} V_{\lambda} \otimes V_{\lambda}^*$

$$\dim(\mathcal{H}_Q) = \sum_{\lambda \in Q} (d_{\lambda})^2$$

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A simple qubit regularization involves $Q = \{1, \underbrace{\square, \begin{array}{|c|} \hline \square \\ \hline \end{array}}, \dots\}$

Hanqing Liu, SC Symmetry 14 (2022) 2 305,



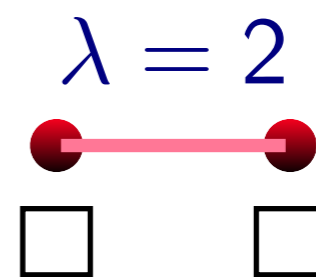
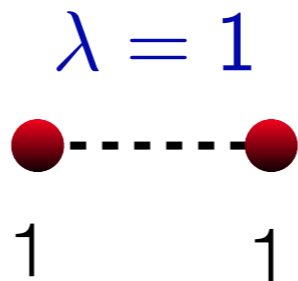
All anti-symmetric irreps

SU(2):

$$Q = \{1, 2\}$$

\mathcal{H}_Q

=



SU(2):
 $Q = \{1, 2\}$

$$\mathcal{H}_Q = \begin{array}{c} \lambda = 1 \\ \bullet \text{---} \bullet \\ 1 \quad 1 \end{array} \oplus \begin{array}{c} \lambda = 2 \\ \bullet \text{---} \bullet \\ \square \quad \square \end{array}$$

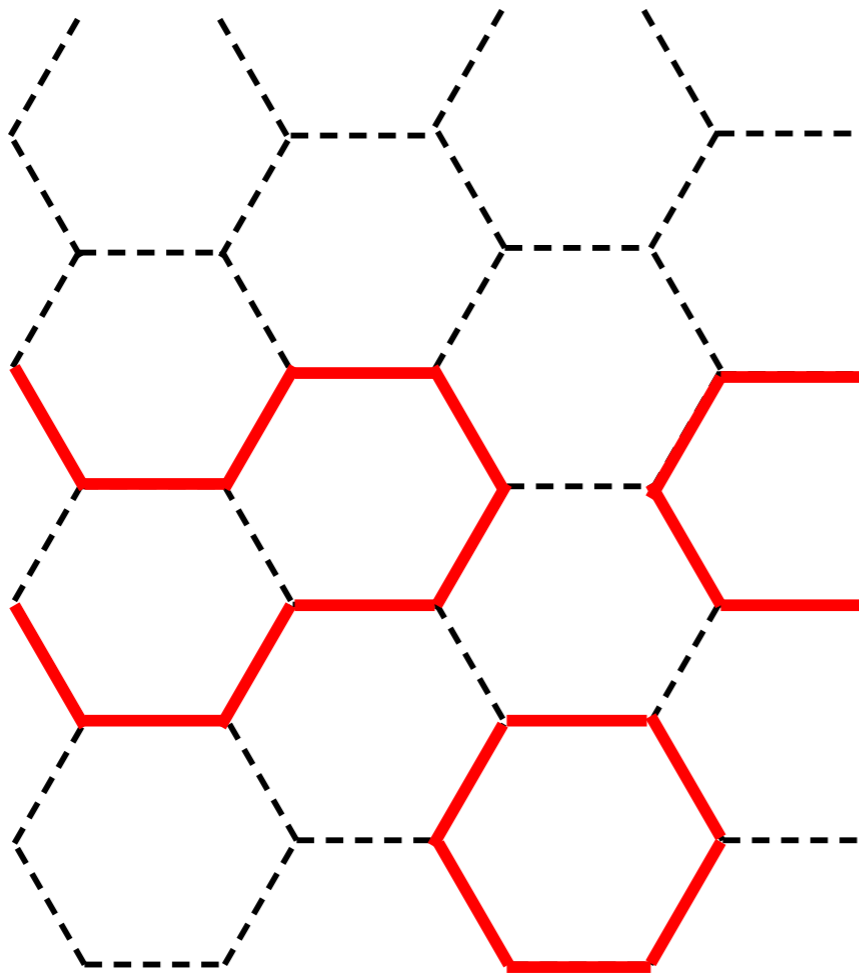
The physical Hilbert Space $\mathcal{H}_{\text{phys}}$ only involves gauge invariant states

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1 1
 \oplus
 \square \square

The physical Hilbert Space $\mathcal{H}_{\text{phys}}$ only involves gauge invariant states

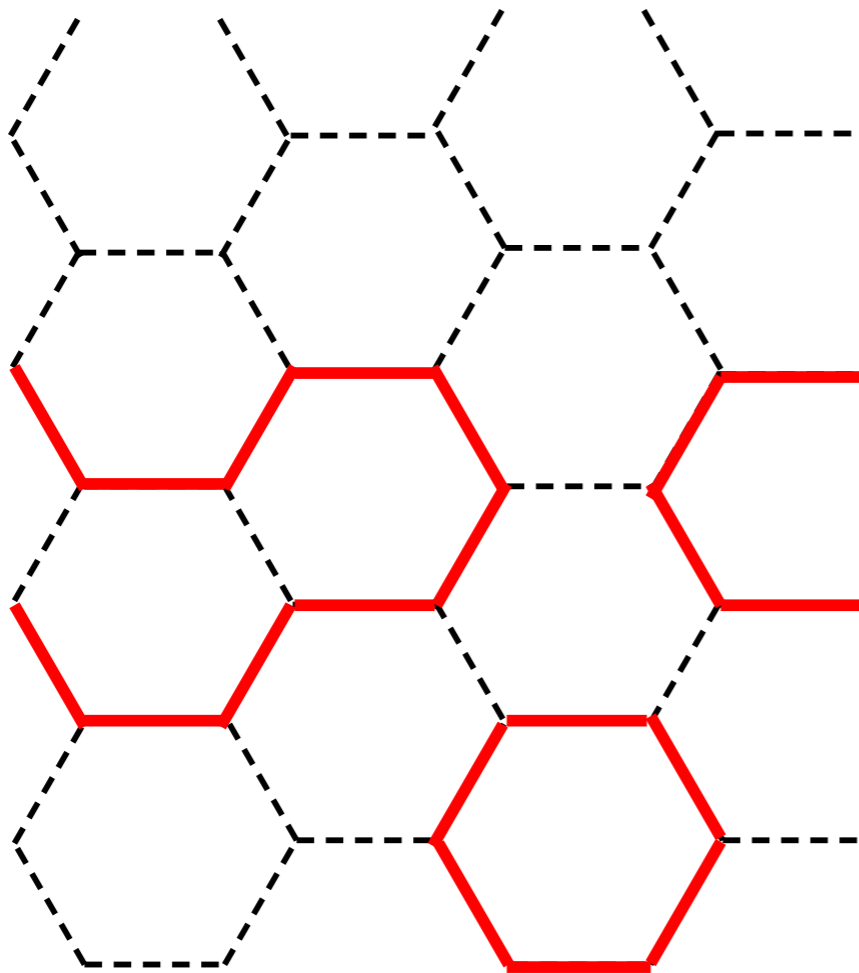


Gauge invariant
 pure gauge state

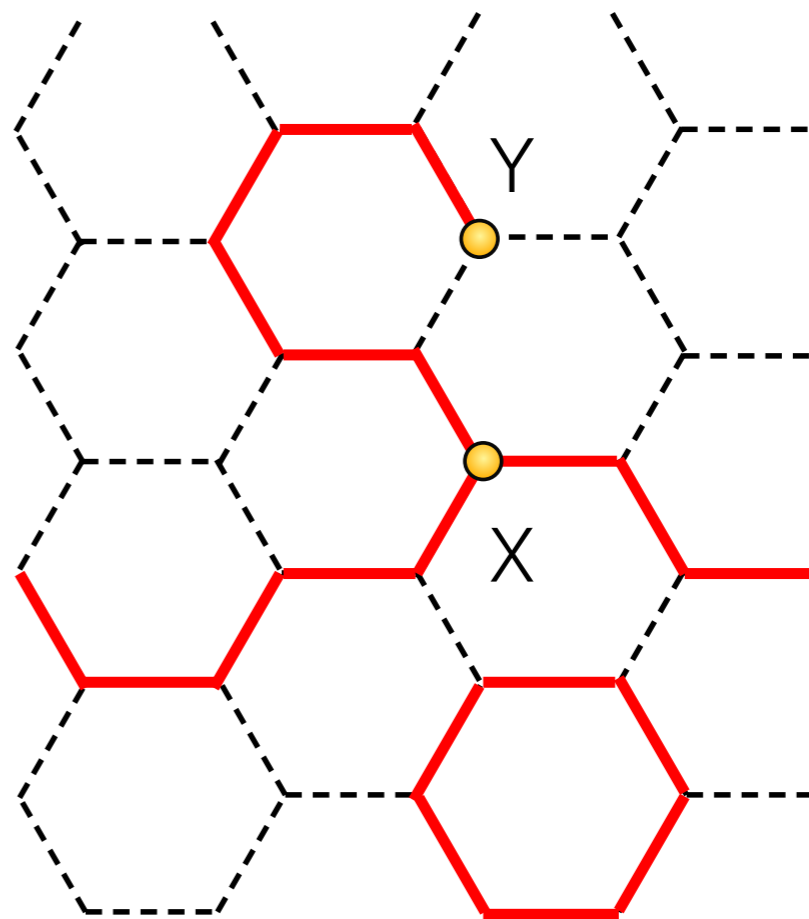
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Gauge invariant
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Gauge invariant state
with two quarks

SU(3):

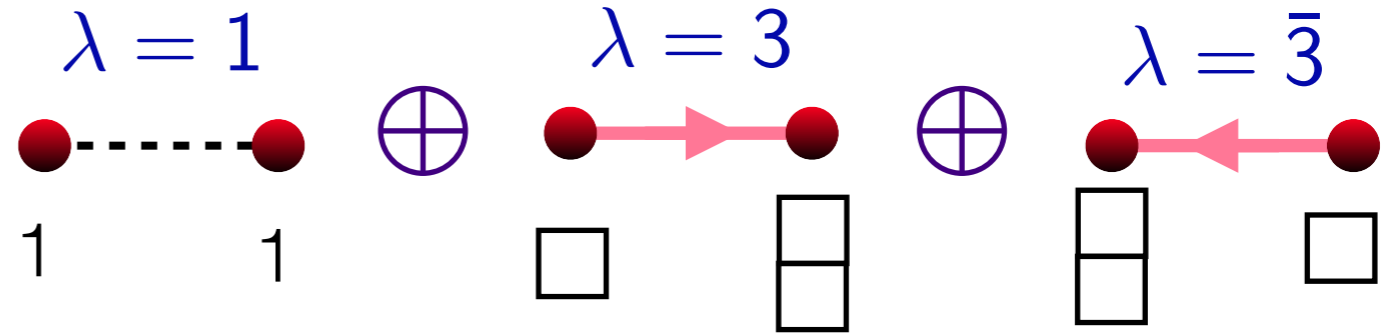
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$$\mathcal{H}_Q = \begin{array}{c} \lambda = 1 \\ \bullet \cdots \bullet \\ 1 \quad 1 \end{array} \oplus \begin{array}{c} \lambda = 3 \\ \bullet \xrightarrow{\quad} \bullet \\ \square \quad \begin{array}{c} \square \\ \square \end{array} \end{array} \oplus \begin{array}{c} \lambda = \bar{3} \\ \bullet \xleftarrow{\quad} \bullet \\ \begin{array}{c} \square \\ \square \end{array} \quad \square \end{array}$$

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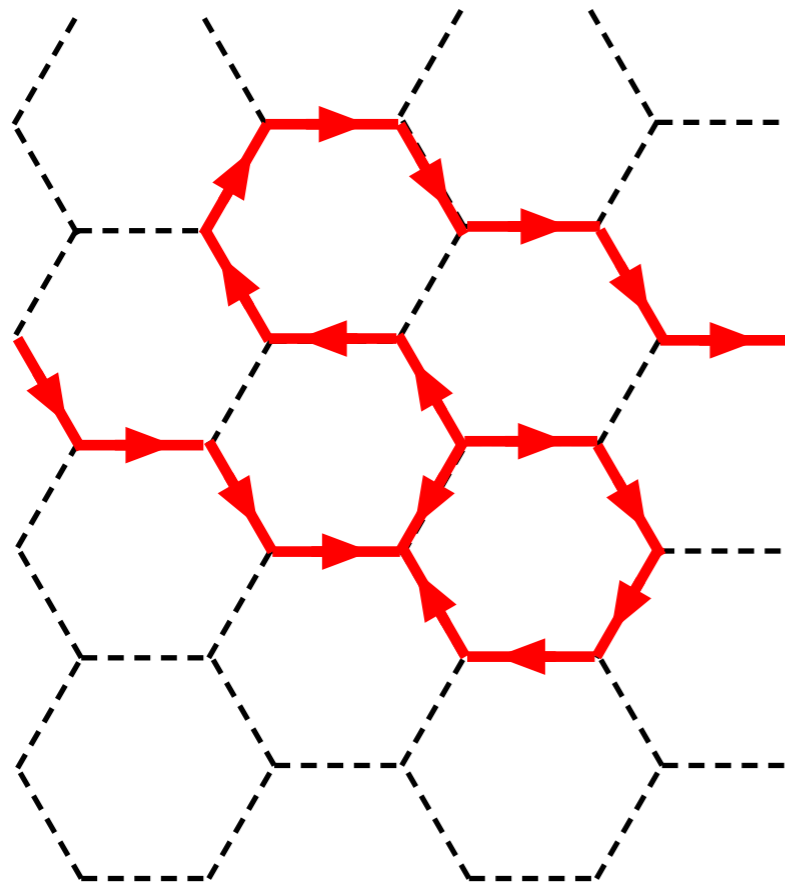
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The physical Hilbert Space $\mathcal{H}_{\text{phys}}$ only involves gauge invariant states

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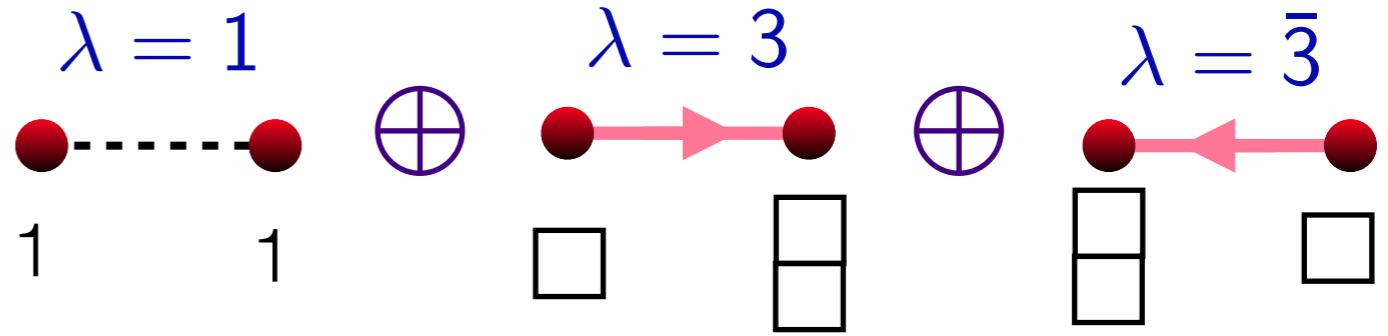


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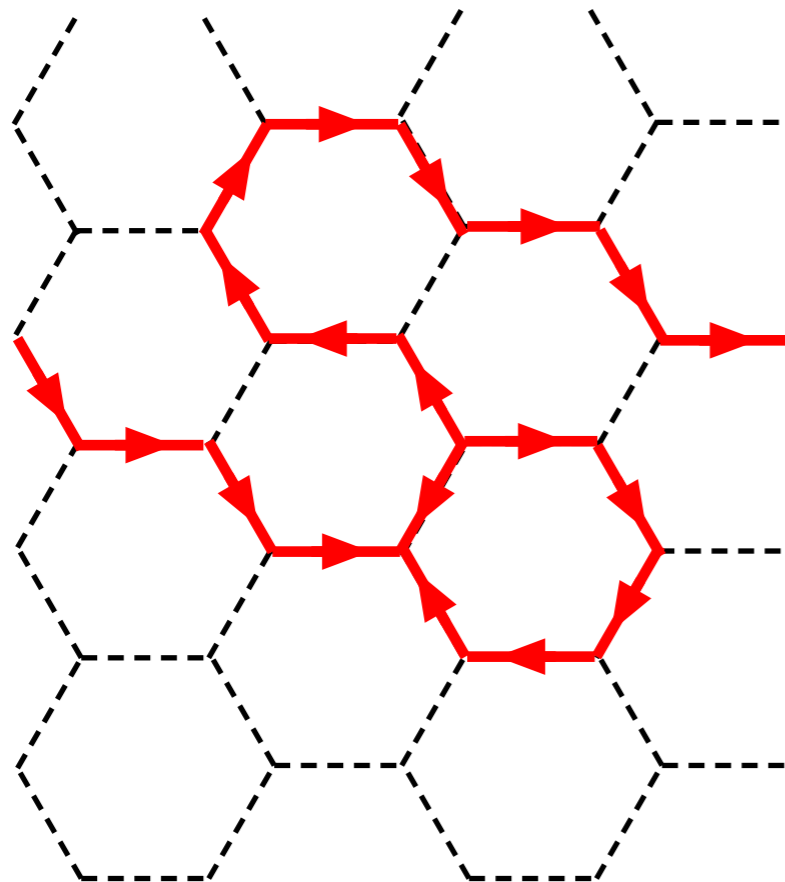
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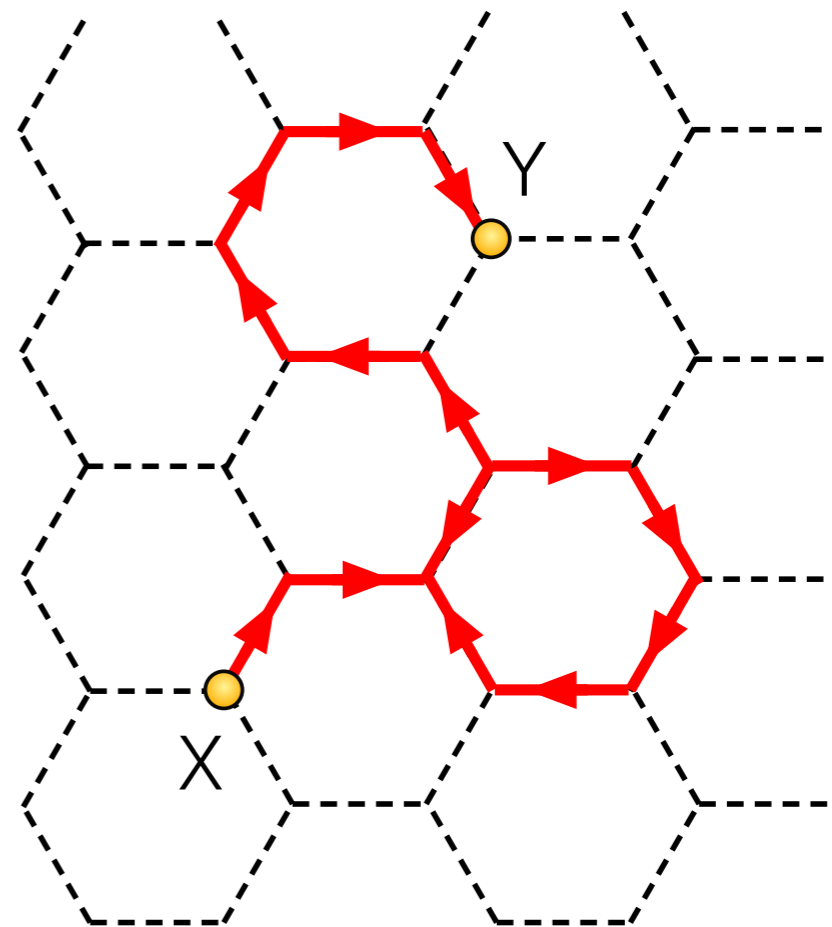


The physical Hilbert Space

$\mathcal{H}_{\text{phys}}$ only involves gauge invariant states



Gauge invariant pure gauge state



Gauge invariant state with a quarks and an anti-quark.

Non-Abelian
Gauge Theories

“momentum basis”


Generalized
Quantum Dimer Models

Non-Abelian
Gauge Theories

“momentum basis”


Generalized
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Easy to study classical dimer models
with Monte Carlo methods!

Non-Abelian Gauge Theories $\xrightarrow{\text{“momentum basis”}}$ Generalized Quantum Dimer Models

Easy to study classical dimer models with Monte Carlo methods!

A “confinement” observable accessible in dimer models!

$$\chi = \frac{1}{V} \sum_{x,y} \frac{Z_{xy}}{Z}$$

$$Z = \text{Tr} \left(e^{-\beta H_{cl}} \right) \quad Z_{xy} = \text{Tr}_{xy} \left(e^{-\beta H_{cl}} \right)$$

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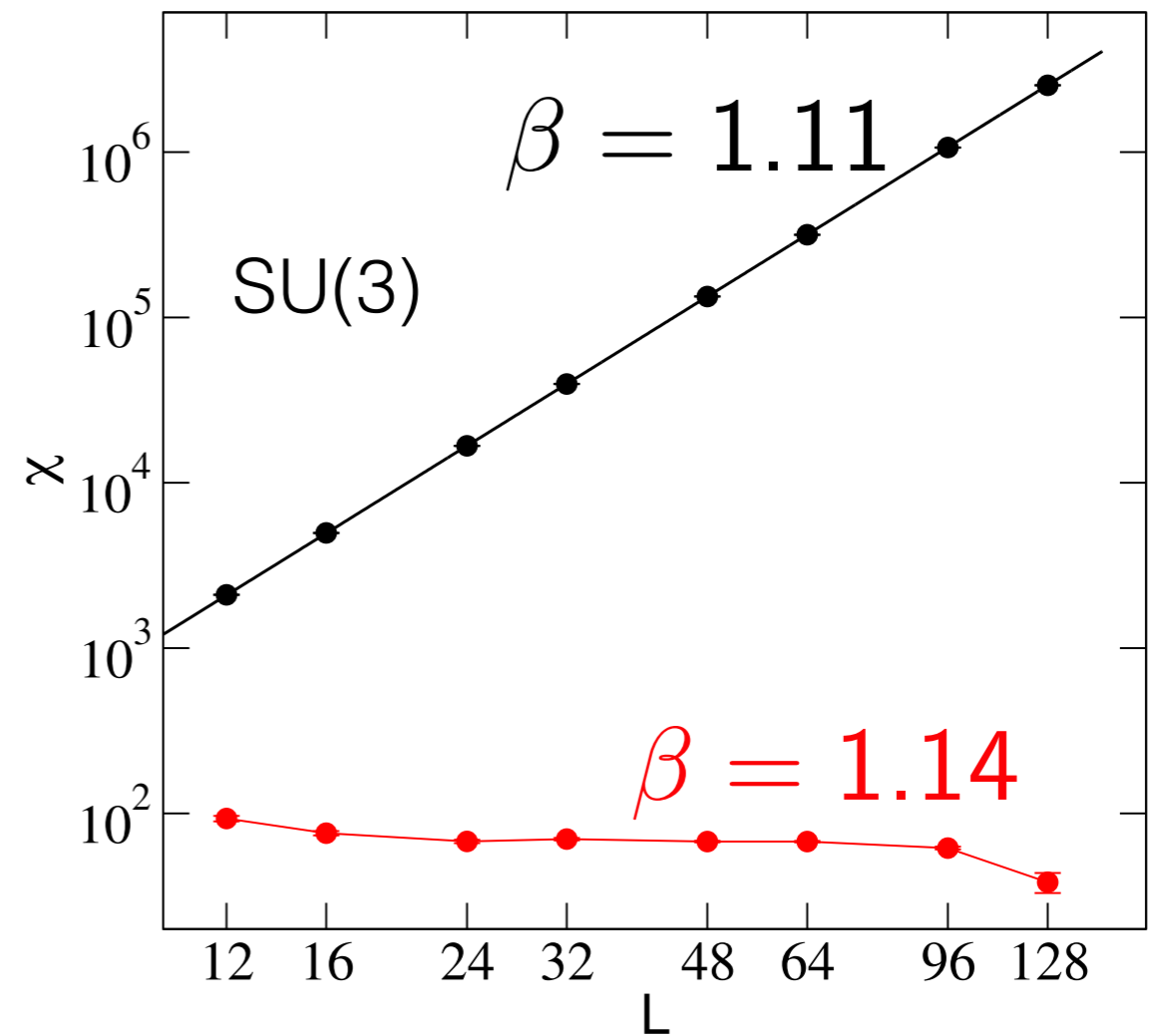
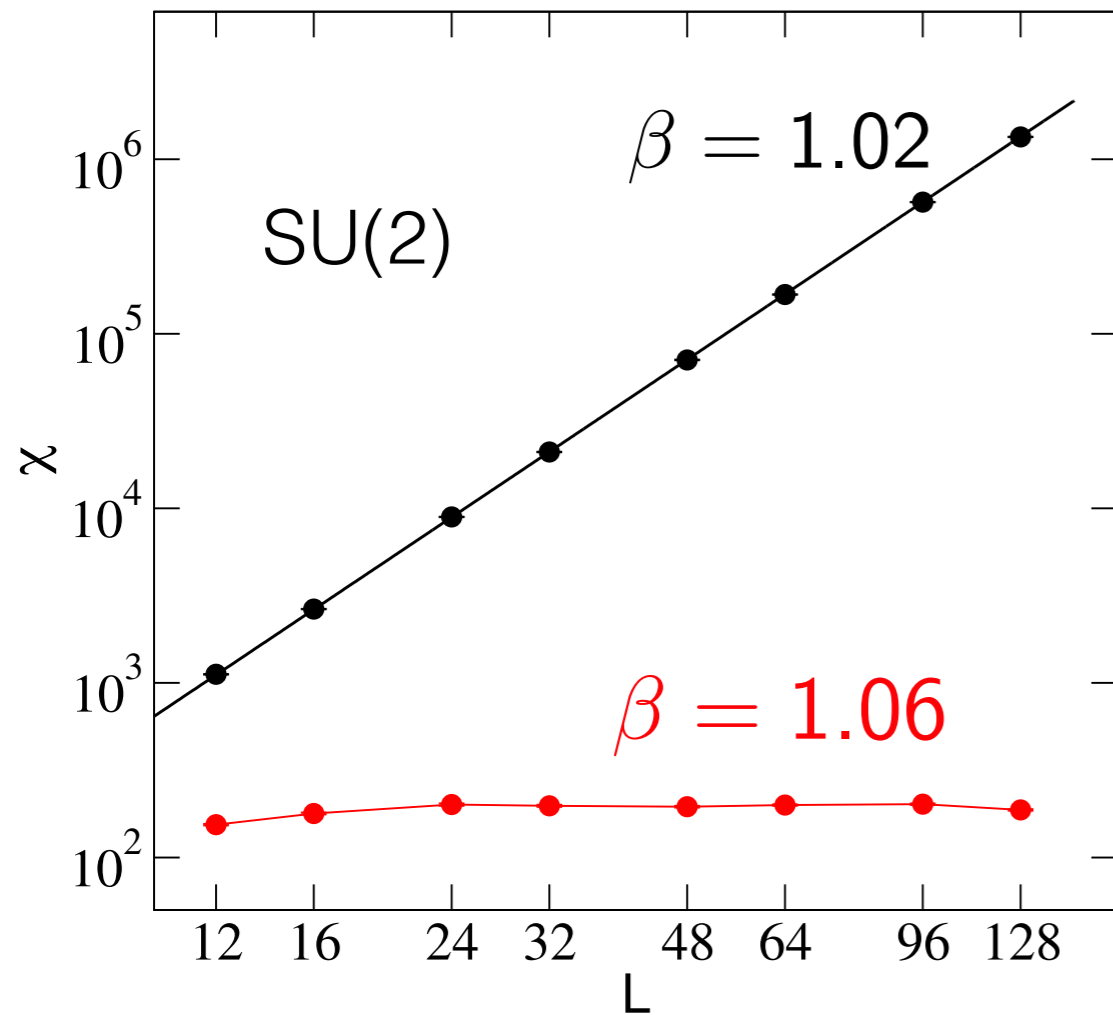
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Deconfined phase: $\chi \sim L^3$ Confined phase: $\chi \sim \text{Const}$

Confinement-Deconfinement Transition in $(3+1)d$

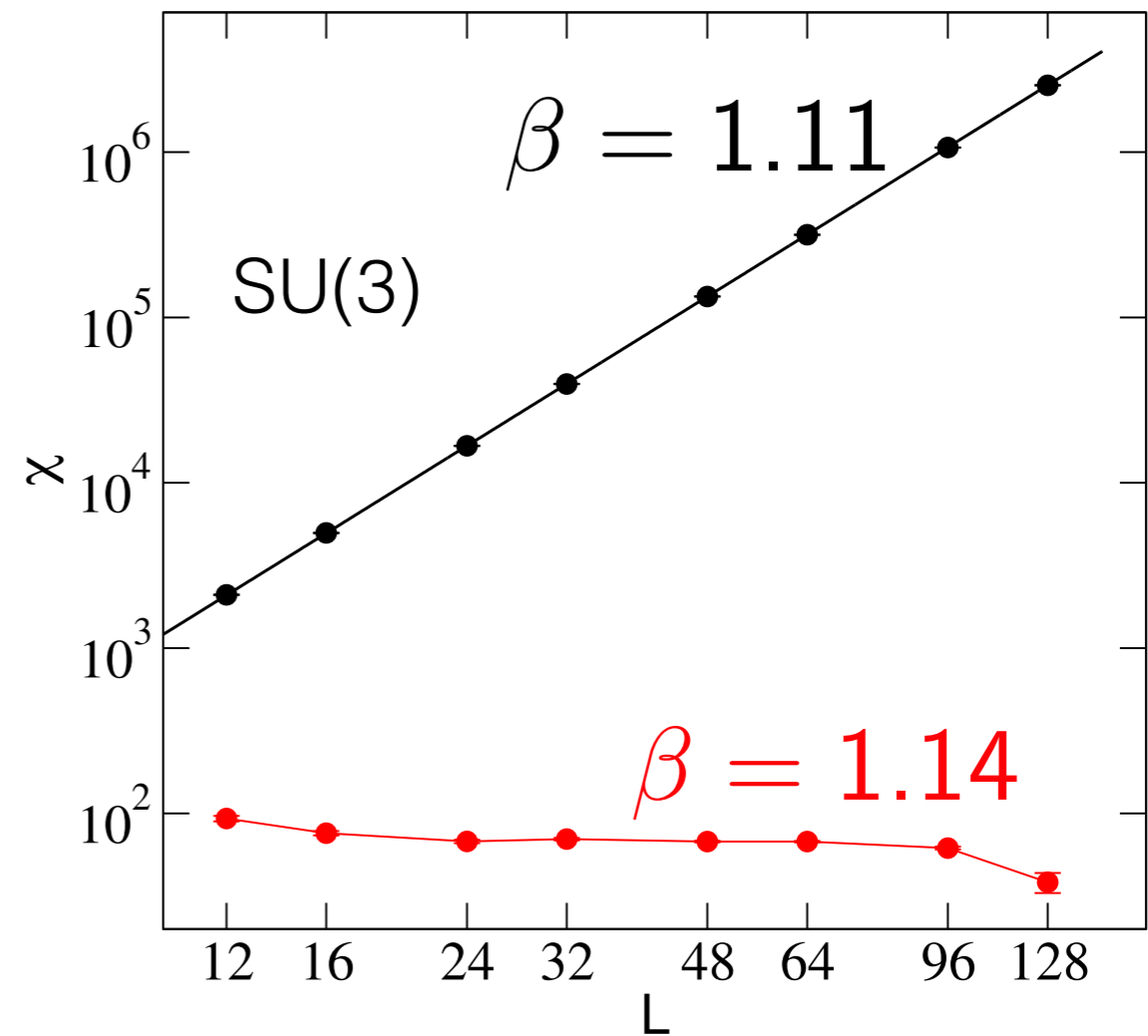
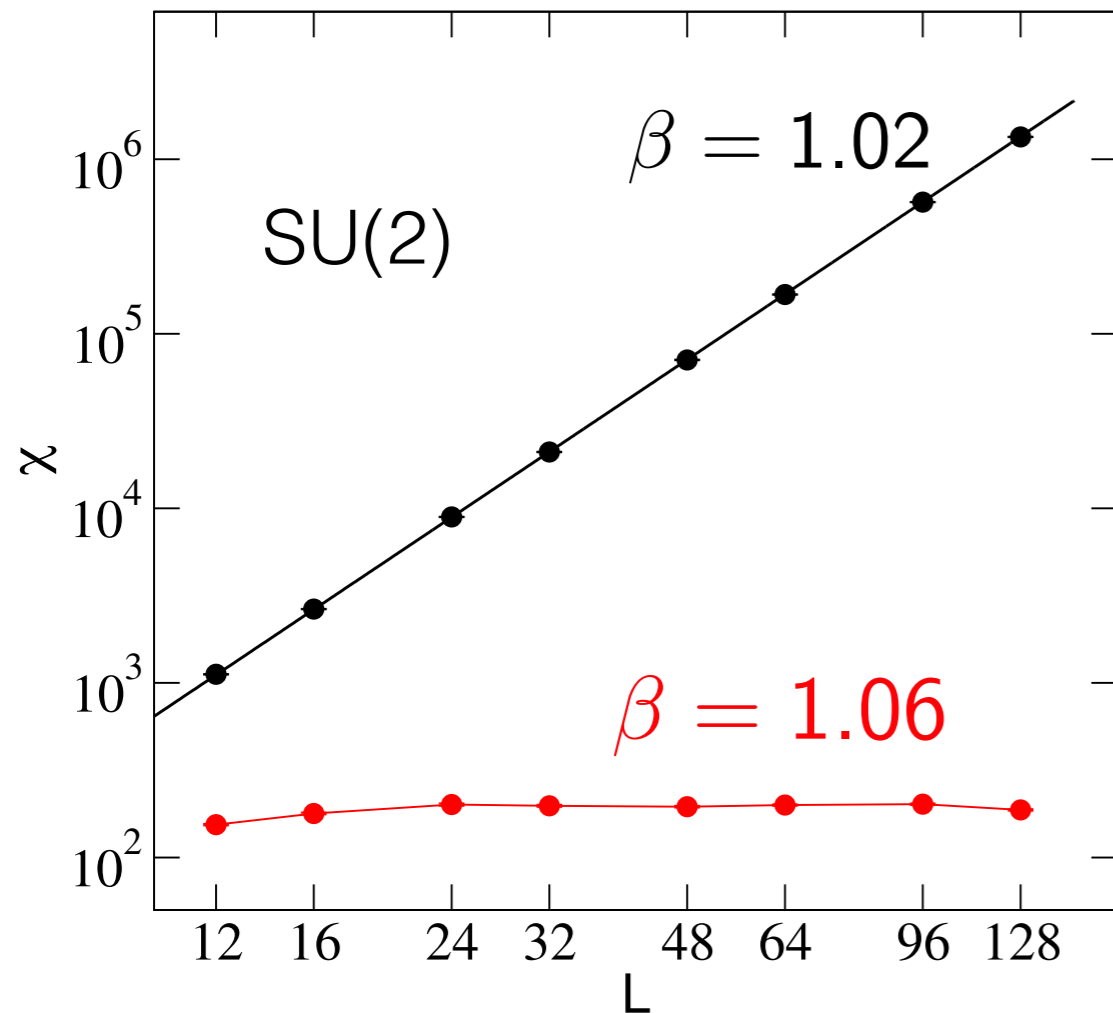
Confinement-Deconfinement Transition in (3+1)d

Preliminary Results



Confinement-Deconfinement Transition in (3+1)d

Preliminary Results



Both seem like first order transitions.

Easy to construct quantum dimer models
free of sign problems!

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Are there quantum critical points separating
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Does asymptotic freedom of Yang-Mills theory
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$SU(N)$ gauge theories with massless staggered matter in $1+1$ dimensions is also very interesting.

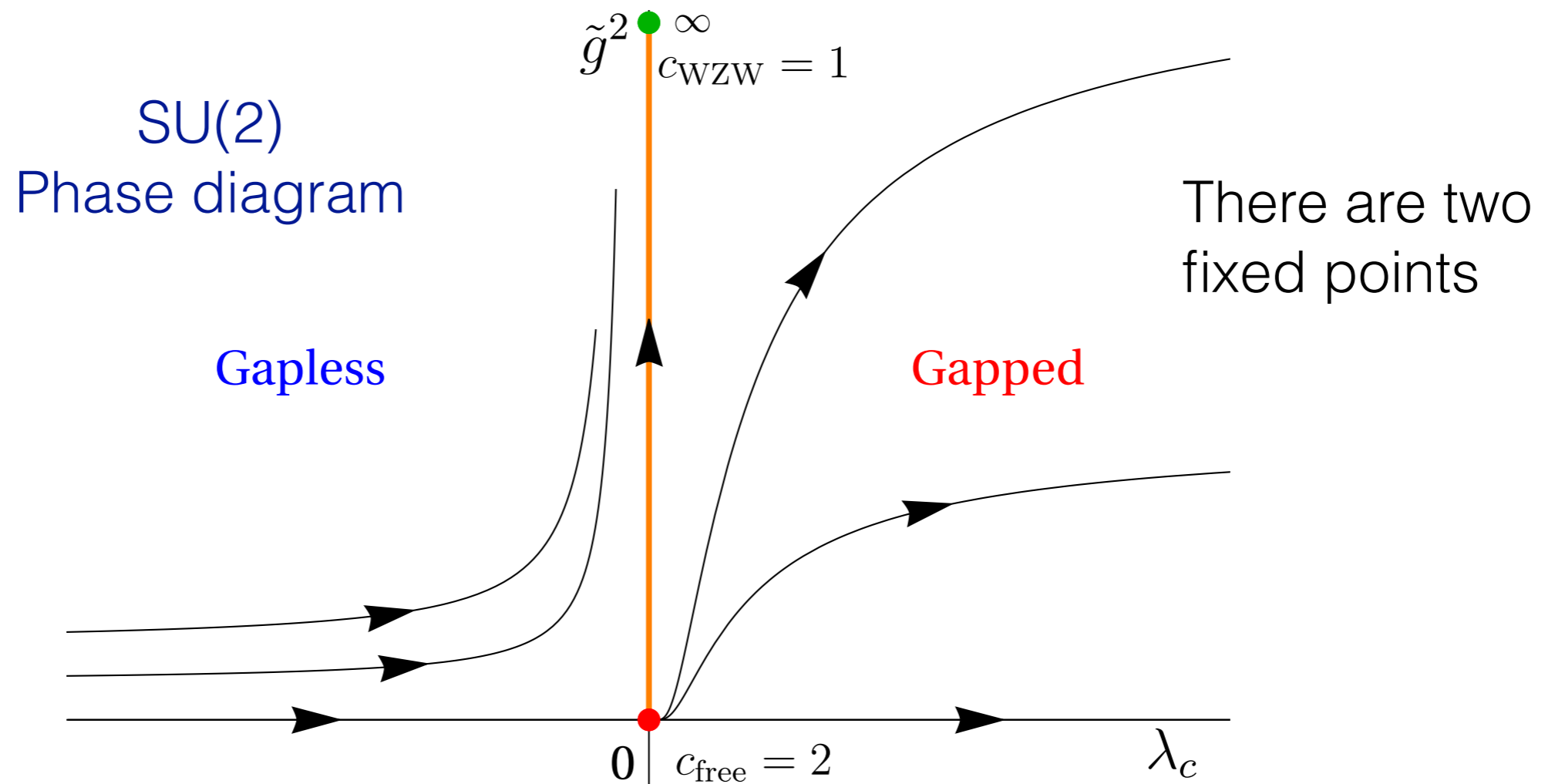
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There are two
fixed points

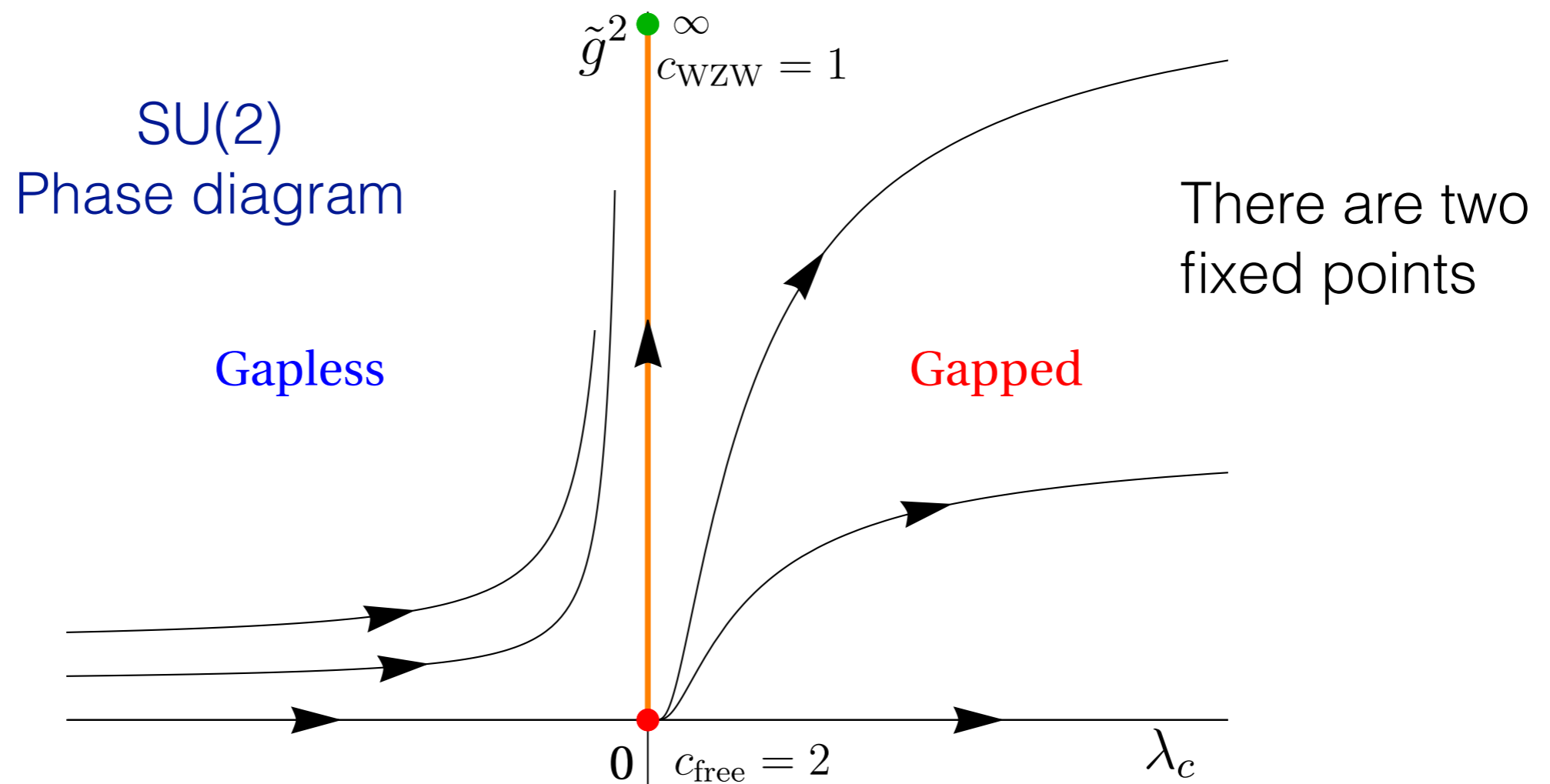
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How can we recover this physics with qubit regularization?

Liu, Bhattacharya, SC, Gupta, arXiv:2312.17734

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Can Yang-Mills theory arise at a quantum critical point of quantum dimer models?