

# Update on Glueballs

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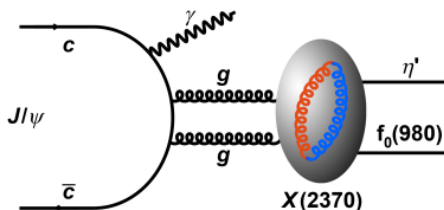
Lattice 2024: Liverpool, England

August 3, 2024



# Why an update on glueballs?

- BESIII first-time determination of  $0^{-+}$  quantum numbers of  $X(2370)$  in PRL **132**, 181901 (2024)



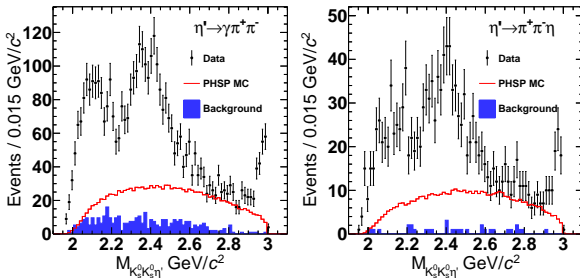
- BESIII earlier observation of  $X(2370)$  in  $J/\psi \rightarrow \gamma\pi^+\pi^-\eta'$  [PRL **106**, 072002 (2011)]
- mass consistent with lightest  $0^{-+}$  glueball from lattice QCD
- gluon rich environment

# $X(2370)$ from BESIII

- partial wave analysis of  $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta'$  gives

$$m = 2395 \pm 11(\text{stat})_{-94}^{+26}(\text{syst}) \text{ MeV}/c^2$$

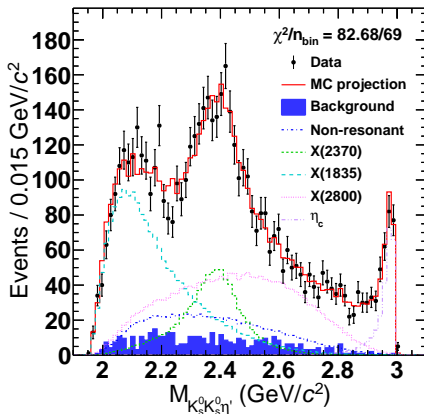
$$\Gamma = 188_{-17}^{+18}(\text{stat})_{-33}^{+124}(\text{syst}) \text{ MeV}$$



- unfortunate name  $X(2370)$  from earlier paper

# $X(2370)$ from BESIII

- optimal PWA fit:  $X(1835)$ ,  $X(2370)$ ,  $\eta_c$  and broad  $0^{-+}$   $X(2800)$   
Breit-Wigner decays through  $f_0(980)\eta'$  to  $(K_S^0 K_S^0)_S \eta'$  and  $(K_S^0 K_S^0)_D \eta'$  with nonresonant components
- statistical significance of  $X(2370)$  is  $> 11.7\sigma$



# Identifying glueballs in experiments

- difficult to identify glueballs
- mass ratios only really known in pure gluon theory
- expect flavor symmetric decays, but differing quark masses leads to differing phase spaces
- no rigorous predictions on decay patterns and their branching ratios
- glueball decays could be similar to that of charmonium
- states could be admixtures with quark-antiquark states

## Some past glueball candidates

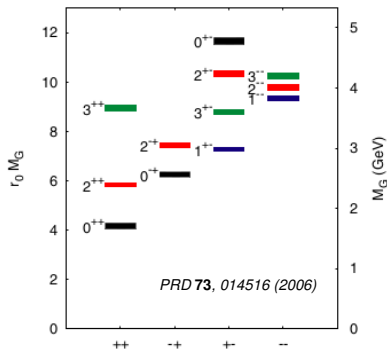
- light scalar candidates  $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1710)$  (MarkII in 1980s, Crystal Barrel in 1990s)
- narrow  $\xi(2230)$  tensor glueball candidate due to good flavor symmetric decay property from MarkIII in 1980s/BESI in 1990s  
→ not confirmed by BESII nor BESIII with much higher statistics
- odderon (odd  $C$ -parity) from D0 and TOTEM [PRL **127**, 062003 (2021)]

# Pure-gauge glueballs in lattice QCD

- pure-gauge glueball spectrum
  - long history in lattice QCD
- calculations date back to early days of lattice QCD (1970s)
- Monte Carlo computations very noisy, rapid correlator falloffs
  - need for large amount of statistics
- progress from extended smeared operators in 1980s/1990s
  - M. Teper, C. Michael, D. Weingarten, among others
- anisotropic lattice: better temporal resolution of correlators
  - CM + M. Peardon late 1990s

# Pure-gauge glueball spectrum

- pure-gauge glueball masses (no quarks)
- mass **ratios** well determined
- setting scale ambiguous since no quarks: string tension from Cornell potential
- lightest scalar  
~ 1600 – 1700 MeV
- states are not fixed numbers of gluons
- scalar mass gap has \$1 million bounty (Clay Mathematics Institute)





# Excited states from correlation matrices

- energies from temporal correlations  $C_{ij}(t) = \langle 0 | \bar{O}_i(t) O_j(0) | 0 \rangle$
- in finite volume, energies are discrete (neglect wrap-around)

$$C_{ij}(t) = \sum_n Z_i^{(n)} Z_j^{(n)*} e^{-E_n t}, \quad Z_j^{(n)} = \langle 0 | O_j | n \rangle$$

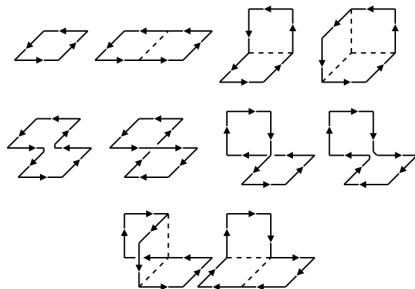
- not practical to do fits using above form
- define new correlation matrix  $\tilde{C}(t)$  using a single rotation

$$\tilde{C}(t) = U^\dagger C(\tau_0)^{-1/2} C(t) C(\tau_0)^{-1/2} U$$

- columns of  $U$  are eigenvectors of  $C(\tau_0)^{-1/2} C(\tau_D) C(\tau_0)^{-1/2}$
- choose  $\tau_0$  and  $\tau_D$  large enough so  $\tilde{C}(t)$  diagonal for  $t > \tau_D$
- 2-exponential fits to  $\tilde{C}_{\alpha\alpha}(t)$  yield energies  $E_\alpha$  and overlaps  $Z_j^{(n)}$

# Glueball operators

- important: use **good** operators for signal before noise growth
- glueball operators: gauge-invariant loops of link variables  $U_\mu(x)$
- Teper fuzzing: large links from links+staples
- multiple sizes for radial structure
- different shapes for orbital structure



# Qualitative features of glueball spectrum

- spectrum qualitatively understood in terms of interpolating operators of minimal dimension (Jaffe, Johnson, Ryzak, Ann. Phys. 168, 344 (1986))

- dimension 4:

$$\text{Tr} F_{\mu\nu} F_{\alpha\beta} \Rightarrow 0^{++}, 0^{-+}, 2^{++}, 2^{-+}$$

- dimension 5:

$$\text{Tr} F_{\mu\nu} D_{\rho} F_{\alpha\beta} \Rightarrow 1^{++}, 3^{++}$$

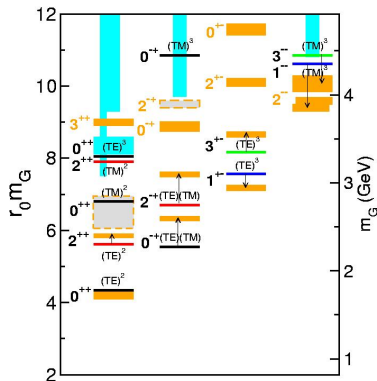
- dimension 6:

$$\begin{aligned} \text{Tr} F_{\mu\nu} F_{\rho\omega} F_{\alpha\beta} &\Rightarrow 0^{\pm\pm}, 1^{\pm\pm}, 2^{\pm\pm}, 3^{\pm-} \\ \text{Tr} F_{\mu\nu} \{D_{\rho}, D_{\omega}\} F_{\alpha\beta} &\Rightarrow 1^{-+}, 3^{-+}, 4^{\pm+} \end{aligned}$$

- of lightest 6 states, 4 have the  $J^{PC}$  of the dimension 4 operators
- absence of low-lying  $0^{\pm-}, 1^{-+}$  glueballs explained

# Glueballs from MIT bag model

- qualitative agreement with bag model
- constituent gluons are TE or TM modes in spherical cavity
- Hartree modes with residual perturbative interactions
- center-of-mass correction
- parameter modifications 1983  $\rightarrow$  1993  
 $\alpha_s : 1.0 \rightarrow 0.5$   
bag parameter  $B^{1/4} :$   
230 MeV  $\rightarrow$  280 MeV

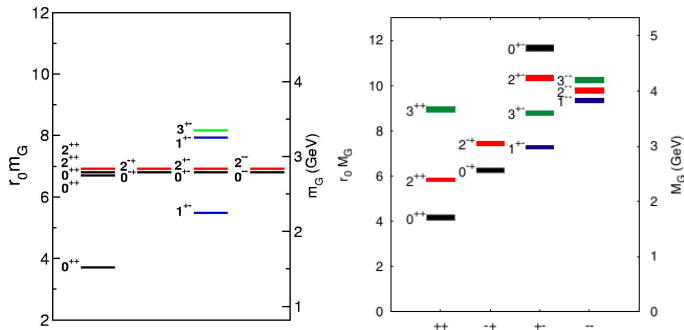


Carlson, Hansson, Peterson, PRD27, 1556 (1983);

J. Kuti (private communication)

# Glueballs from Isgur-Paton flux tube model

- disagreement with one particular string model
- Isgur, Paton, PRD31, 2910 (1985)

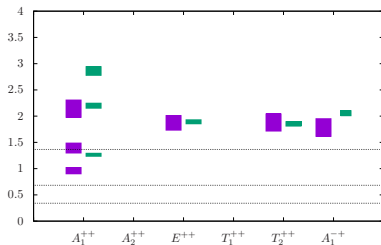


# Why are glueballs with quarks so hard in lattice QCD?

- must extract all levels lying below glueballs of interest
- many 2-meson, 3-meson, 4-meson levels expected below
- 2-meson correlators require timeslice-to-timeslice propagators
- glueballs expected to be resonances
- glueballs require high statistics: difficult with quarks
- scalar sector requires large VEV subtraction

# Glueballs with $N_f = 4$ light quarks

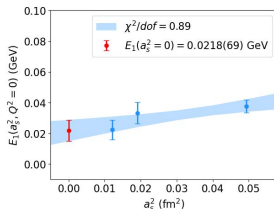
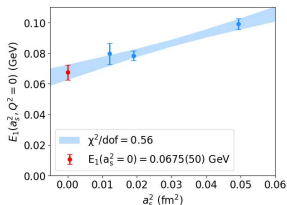
- Athenodorou et al. arXiv:2308.10054 [hep-lat]
- examine effect of quark loops on glueball spectrum
- several ensembles used with  $m_\pi \approx 250$  MeV
- GEVP used on correlation matrices with **only** glueball operators
- no meson-meson operators!! (really need these)



- conclusions: scalar channel lowers toward  $2\pi$ , tensor and pseudoscalar spectrum only slightly affected

# Radiative decay of the scalar glueball

- Zou et al. arXiv:2404.01564 [hep-lat]
- **quenched** approximation, 3 gauge ensembles  
 $a_s \sim 0.11, 0.14, 0.22$  fm, continuum extrapolation
- evaluate EM transition matrix element  $\langle S | J_{\text{em}}^\mu | V \rangle$
- multipole expansion: two form factors  $E_1(Q^2)$  and  $C_1(Q^2)$
- get widths from  $E_1(0)$ :  $Q^2 \rightarrow 0$ ,  $a \rightarrow 0$  limits taken

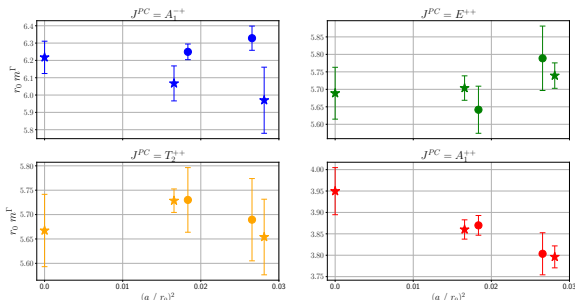


- find  $\Gamma(J/\psi \rightarrow \gamma G) = 0.578(86)$  keV with  
 $\text{Br}(J/\psi \rightarrow \gamma G) = 6.2(9) \times 10^{-3}$  and  $\Gamma(G \rightarrow \gamma \phi) = 0.074(47)$  keV
- conclude  $J\psi \rightarrow \gamma G \rightarrow \gamma \gamma \phi$  not detectable by BESIII



# Error reduction algorithm

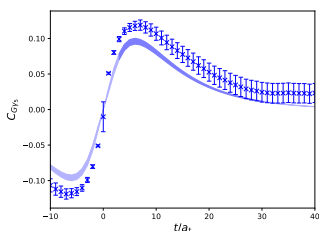
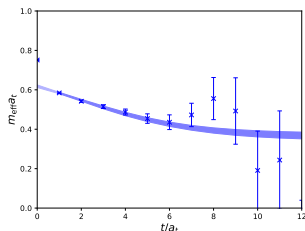
- Barca et al. arXiv:2406.12656 [hep-lat]
- new multi-level sampling procedure proposed for error reduction of glueball correlators in pure gauge theory



- state-of-the-art stars, new method circles
- no reduction in glueball mass errors, but significant error reduction in large- $t$  correlators to improve confidence in plateau estimates

# Glueball- $\eta$ mixing

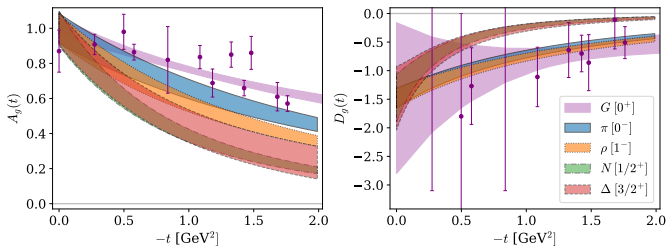
- X. Jiang et al., PRD **107**, 094510 (2023)
- studied mixing of  $0^{-+}$  glueball and pseudoscalar  $\bar{q}q$  meson
- $16^3 \times 128$  anisotropic  $N_f = 2$  lattice with  $m_\pi \approx 350$  MeV
- used distillation, no GEVP, diagonal and cross correlators
- obtained very small  $3.5^\circ$  mixing angle glueball- $\eta$  from cross correlator



- (left) Effective mass  $0^{-+}$  glueball (right) Cross correlator (shifted horizontally) with fit from temporal derivative

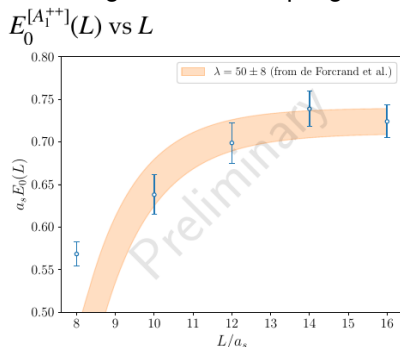
# Gravitational form factors of glueballs

- D. Pefkou (poster), Abbott, Hackett, Romero-Lopez, Shanahan
- use gravitational form factors to probe structure of glueballs
- obtained from energy-momentum tensor matrix elements
- express matrix element of  $T_{\mu\nu}$  in scalar glueball state in terms of form factors  $A(Q^2)$  and  $D(Q^2)$
- results also for  $\pi, \rho, N, \Delta$
- preliminary results in pure gauge theory  $24^3 \times 48$  lattice  
 $a = 0.1$  fm



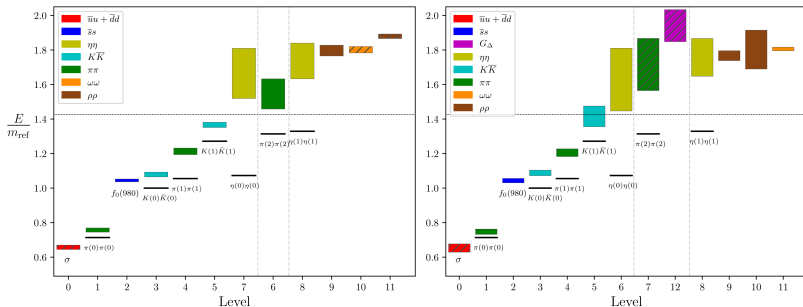
# Scalar glueball scattering

- M. Hansen parallel talk (with M. Bruno and A. Rago)
- calculate finite-volume energies in Yang-Mills to extract  $GG \rightarrow GG$  amplitudes
- anisotropic lattice, use of multi-level algorithm, scale set using  $t_0$
- obtained volume dependence of  $A_1^{++}$  energy (single and 2 glueball operators)
- use Lüscher relation to get trilinear coupling  $\lambda$  from energies



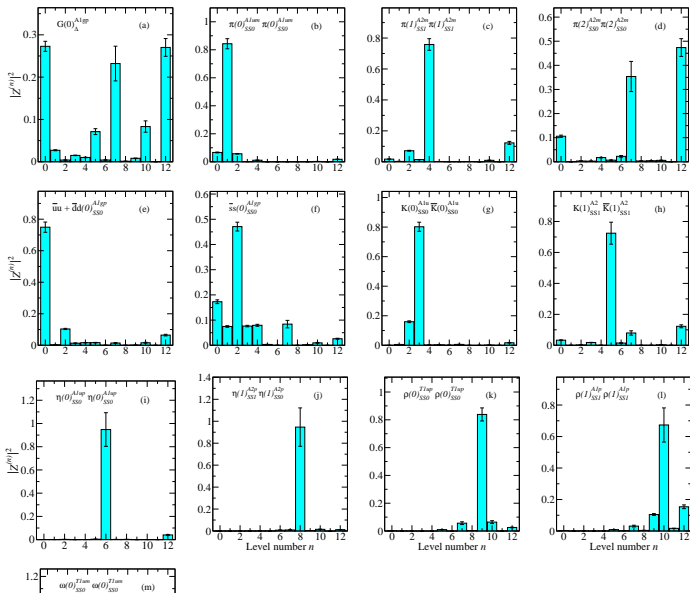
# Scalar glueball in $N_f = 2 + 1$ QCD

- R. Brett et al. arXiv:1909.07306 [hep-lat]
- $24^3 \times 128$  anisotropic lattice,  $m_\pi \sim 390$  MeV
- $A_{1g}^+$  spectrum show below:
  - (left) 4  $\bar{q}q$  operators, 10 two-meson operators
  - (right) added glueball operator



- bad news for the scalar glueball?
- need for Lüscher analysis (ongoing)

# $A_{1g}^+$ overlaps



# Conclusion

- glueballs very challenging in full QCD
- recent study suggests no scalar state below 2 GeV is pure or dominant glueball
- pseudoscalar glueball is new focus of attention
- identifying glueballs in experiments is challenging too
- long history in lattice QCD and still very active