## A Confederacy of Anomalies

A personal recollection of early years in lattice gauge theory <sup>1</sup>

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<sup>1</sup>not a balanced review! free after John Kennedy Toole, A Confederacy of Dunces, with one dunce (a + b) = 0

### Prologue

1969: September, PhD student of S.A. Wouthuysen in Amsterdam 1970: started working on my own, e.g. on:

Glashow, Weinberg, Salam, Veltman, 'Higgs', Wilson, massive Yang-Mills, partons, renormalization group, chiral anomaly,...

#### 1971:

my hero, Wilson, published article <sup>2</sup> which contained a non-perturbative approximation to a path integral (in the continuum)

- no longer a magic trick for deriving Feynman diagrams

### Prologue 1972

chiral anomalies should cancel <sup>3 4</sup>

judged this artificial – wanted non-perturbative UV-regulator

Spring: scalar field on a lattice

August: moved to LA, student of Robert J. Finkelstein at UCLA topic: massive YM in Schwinger's Source Theory

Christmas: non-Abelian gauge field on a lattice, 'Wilson action'

<sup>3</sup>C. Bouchiat, J. Iliopoulos, Ph. Meyer, An anomaly-free version of Weinberg's model, PLB38(1972)519

<sup>4</sup>D.J. Gross, R. Jackiw, *Effect of Anomalies on Quasi-Renormalizable* Theories, PRD6(1972)477

## 1973

testing perturbative continuum limit at one loop on  $U(1)_V \times U(1)_A$  gauge-Higgs model coupled to one Dirac field:  $\mathcal{L}_F = \bar{\psi}\gamma^{\mu}\partial_{\mu}\psi + \bar{\psi} \left[g \gamma^{\mu} V_{\mu} + g_5 i\gamma^{\mu}\gamma_5 A_{\mu} + G(\sigma + i\gamma_5 \pi)\right]\psi$ 

'anomalous chiral gauge theory'

spatial cubic lattice, continuous time (just QM)

develop lattice methods

vector and axial-vector Ward-Takahashi identities exactly valid at finite lattice spacing

velocity-of-light counterterm etc.

vector selfenergy ('vacuum polarization'): correct *Lorentz-invariant form* but factor 8 too large

VVA and VV $\pi$  triangle diagrams: zero

propagators and vertices periodic in spatial momentum space



fermion propagator

$$\left[i\gamma^0 k_0 + i\gamma^j P_j(k) + m\right]^{-1}, \quad P_j(k) = \frac{1}{a}\sin(ak_j)$$

 $8 = 2^3$  zeros of  $P_j(k)$  at  $ak_j = \{0, \pi\}$  contribute to continuum limit of vector selfenergy Tried to cure by:  $P_j(k) = \frac{2}{a} \sin\left(\frac{ak_j}{2}\right)$ discontinuous  $\Rightarrow$  non-local *non-local, non-Lorentz covariant, UV divergent* continuum limit

- multiplicity not-ignorable
- doublers real particles?

decided to leave it 'for a while'

November: article with Robert <sup>5</sup>

inspired by  $^{6}$  investigated possibility of massive solution of massless YM, using Schwinger-Dyson equations  $^{7}$ 

<sup>5</sup>Robert J. Finkelstein and JS, *Massive Gauge Field in Source Theory II*, Ann.Phys.88(1974)157

<sup>6</sup>J.M. Cornwall, R.E. Norton, *Spontaneous Symmetry Breaking without Scalar Mesons*, *Phys.Rev.D8*(1973)3338

<sup>7</sup>JS, Possibility that massless Yang-Mills fields generate massive vector particles, Phys.Rev.D10(1974)2473

Spring/Summer: seminar at UCLA by Wilson <sup>8</sup> bomb shell

lattice, confinement, strong coupling!

after seminar I did *not* tel him about my own lattice work, but mentioned problem with fermions. Wilson's answer:

add term  $\propto \partial_\mu \bar{\psi} \partial_\mu \psi$ 

chiral symmetry should come back in continuum limit

November: PhD, Thesis:

Massive Vector Particles with Yang-Mills Couplings

- Source Theory + the Schwinger-Dyson eqns. article (lattice regularization only briefly mentioned)

<sup>&</sup>lt;sup>8</sup>K.G. Wilson, Confinement of Quarks, Phys.Rev.D10(1974)2445 → E ∽ Q ~

January: back to Amsterdam

- depressed - could not write article about my lattice results

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turned to Regge pole theory description of experimental scattering results at CERN

together with new PhD student Luuk Karsten

## 1977

September: SLAC Summer Institute: lecture by Sidney Drell on Lattice Field Theory

Avoid fermion doubling: <sup>9</sup>  $P_j(k) = k_j$ ,  $ak_j \in \{-\pi, \pi\}$ 



I suggested problems (expected from  $P_j(k) = \frac{2}{a} \sin\left(\frac{ak_j}{2}\right)$ ):

non-local divergences in continuum limit of gauge theory

Drell: "You ought to write this up!"

 $\Rightarrow$  returned to LGT, with Luuk Karsten

<sup>9</sup>S.D. Drell, M. Weinstein, S. Yankielowicz, Strong Coupling Field Theories. 2. Fermions and Gauge Fields on a Lattice, Phys.Rev.D14(1976)1627 ( ) ⇒ ⇒ ⇒ ⊃ < ?

#### 1978

catch up with exponentially rising number of publications in particular  $^{10\ 11\ 12\ 13\ 14\ 15}$  (overlooked  $^{16\ 17})$ 

<sup>10</sup>K.G. Wilson, 13th International School of Subnuclear Physics: *New Phenomena in Subnuclear Physics*, CLNS-321 (1975)

<sup>11</sup>J.B. Kogut and L. Susskind, *Hamiltonian Formulation of Wilson's Lattice Gauge Theories*, Phys.Rev.D11(1975)395

<sup>12</sup>L. Susskind, Lattice fermions, Phys.Rev.D16(1977)3031

<sup>13</sup>B.E. Baaquie, *Gauge Fixing and Mass Renormalization in the Lattice Gauge Theory*, Phys.Rev.D16(1977)2612

<sup>14</sup>H.S. Sharatchandra, *The Continuum Limit of Lattice Gauge Theories in the Context of Renormalized Perturbation Theory*, Phys.Rev.D18(1978)2042

<sup>15</sup>J. Shigemitsu, *Spectrum calculations in lattice gauge theory using Wilson's fermion method*, Phys.Rev.D18(1978)1709

<sup>16</sup>M. Lüscher, Construction of a Selfadjoint, Strictly Positive Transfer Matrix for Euclidean Lattice Gauge Theories, Commun.Math.Phys.54(1977)283

<sup>17</sup>M. Creutz, Gauge Fixing, the Transfer Matrix, and Confinement on a Lattice, Phys.Rev.D15(1977)1128

calculated continuum limit of 'anomaly (triangle) diagrams' with Drell-Weinstein-Yankielowicz (DWY) fermion method & find indeed non-local, non-covariant, divergent contributions <sup>18</sup>

... also in gauge-field selfenergy diagrams <sup>19</sup>

- 'doubler fermions' replaced by non-local & non-covariant contributions

<sup>18</sup>L.H. Karsten, JS, *Axial Symmetry in Lattice Theories*, Nucl.Phys.B144(1978)536

<sup>19</sup>L.H. Karsten, JS, *The Vacuum Polarization with SLAC Lattice Fermions*, Phys.Lett.B85(1979)100 September: Cargèse Summer Institute

Wilson describes numerical RG computations <sup>20</sup>

& mentions MC computations by Creutz <sup>21</sup> (eye opener for me)

and strong coupling calculations, with Padé extrapolation to weak coupling, by Kogut, Pearson and Shigemitsu  $^{\rm 22}$ 

<sup>20</sup>K.G. Wilson, *Monte Carlo Calculations for the Lattice Gauge Theory*, NATO Sci.Ser.B 59 (1980) 363-402

<sup>21</sup>M. Creutz, Solving Quantized SU(2) Gauge Theory, Brookhaven,
 September 1979; Asymptotic Freedom Scales, Phys.Rev.Lett.45(1980)313
 <sup>22</sup>J.B. Kogut, R.B. Pearson, J. Shigemitsu, The QCD beta Function at
 Intermediate and Strong Coupling, Phys.Rev.Lett. 43(1979)484 (2000) (2000)

November: PhD Luuk Karsten. Thesis:

On Lattice Gauge Theories and On Backward Pion-Nucleon Scattering

- doubler fermions at  $ak_\mu=\pi$  are genuine particles
- chiral charges:  $g_5 
  ightarrow g_5 Q_5$ ,  $Q_5 = (-1)^n$ ,
- n = number of  $\pi$  s in four-vector ak giving sin $(ak_{\mu}) = 0$

$$\sum Q_5 = 0$$

- anomalies cancel in triangle diagrams !

in the  $U(1)_V \times U(1)_A$  model a Wilson-type mass term made gauge invariant with the Higgs field  $\sigma + i\pi$  led to the tree-graph level formula

$$m_{\mathrm{F}} = m + rac{r}{a} \sum_{\mu} [1 - \cos(ak_{\mu})]$$

with  $m=G\langle\sigma
angle$  and  $r/a=G_{
m W}\langle\sigma
angle$  (Wilson's r=1)

- in a continuum limit  $a \rightarrow 0$  with  $r = a \tilde{r}$  this would result in a mass spectrum:

$$m_{\rm F} = m + 2n\,\tilde{r}\,, \quad n = 1, 2, 3, 4$$

subsequently the doublers  $(n \ge 1)$  seem then be removable by sending  $\tilde{r} \to \infty$ , *except* for triangle diagrams where they become the chiral anomaly in the axial-vector WT identity (thesis and <sup>23</sup>)

<sup>23</sup>L.H. Karsten, *The Lattice Fermion Problem and Weak Coupling Perturbation Theory*, in Field Theoretical Methods in Elementary Physics, NATO Sci.Ser.B55(1980)235 However,  $\tilde{r}\to\infty$  would imply  $G_W\to\infty$  and one would be led to strong coupling dynamics . . .

Luuk moved to Stanford

we still had to write-up publish the results extend the calculations applying to QCD with Wilson fermions at fixed r

February: Banks and Caashers stimulating paper on chiral symmetry breaking  $^{\rm 24}$ 

May: strong coupling paper  $^{25}$  ( $\approx$  simultaneously with  $^{26}$ ) Hamiltonian formulation,  $N_{\rm f}$  flavors, N colors, two orders in 1/N

$$S_{\text{mass}} = \bar{\psi} M \psi - \bar{\psi} W \psi$$
  
$$\bar{\psi} W \psi = \frac{ar}{2} \sum_{x,j} \left[ \bar{\psi}(x) U_j(x) \psi(x+a_j) + \bar{\psi}(x+a_j) U_j^{\dagger}(x) \psi(x) \right]$$

Greensite and me discovered that our write-ups nearly had same title)

- anti-ferromagnetic ground state

Nambu-Goldstone bosons (including vector mesons)

 $m^2 \propto M \langle \bar{\psi} \psi \rangle \,, \quad \langle \bar{\psi} \psi \rangle \propto N$ 

<sup>24</sup>T. Banks, A. Casher, *Chiral Symmetry Breaking in Confining Theories*, NPB169(1980)103

<sup>25</sup>JS, Chiral Symmetry Breaking in QCD: Mesons as Spin Waves, NPB175(1980)307

increasing r: increase  $M > M_{\rm cr}$ , critical mass  $M_{\rm cr}(r,g) \propto r^2$ , (avoid broken parity phase when  $M < M_{\rm cr}$ )

for 0.7 <  $r \le 1$  vector mesons ( $\rho$ ,  $K^*$ ) have lost their NGB character, but *not* the charged pseudo scalars ( $\pi^{\pm}$ , K)

'current quark mass' of flavor a:  $m_a = M_a - M_{
m cr}$ 

$$m_{ab}^2 \propto (m_a + m_b), \quad a \neq b$$

neutral:  $m_{aa}^2 > m_{ab}^2$  in next order 1/N quantitatively: effective strength of anomaly ('U(1) problem')

'dynamical quark mass' of oldfashioned quark model  $m_{
m dyn}$ 

$$m_V \simeq 2 m_{\rm dyn} + m_a + m_b$$

#### 1980 cont.

for  $M = M_{\rm cr}$ , chiral *non-singlet* charge  $Q_{ab}^5$  commutes with effective Hamiltonian and creates zero energy GB state  $Q_{ab}^5|0\rangle$  out of vacuum

chiral *singlet* charge  $Q_{aa}^5$  is anomalous for r > 0

currents were derived from a lattice version of the Standard Model that generated the *r*-parameter in terms of a 'Wilson-Yukawa coupling'  $G_{\rm W}$  and the vev of the Higgs field  $\Phi$ :  $\tilde{r} = G_{\rm W} \langle \Phi \rangle$ 

decoupling doublers would require  $G_{\rm w} \rightarrow \infty$  ??

matrix elements of currents (meson decay constants  $f_{\pi}$ ,  $\gamma_{\rho}$ , etc) did not come out very well (neither very badly)

anomaly paper with Luuk <sup>27</sup> (errata in <sup>28</sup>)

- 
$$ar{\psi}(M_{
m cr}-W)\psi
ightarrow$$
 flavor-  $U(1)$  anomaly

first of No-Go proofs by Nielsen and Ninomiya <sup>29 30 31</sup>

<sup>28</sup>R. Groot, J. Hoek, JS, Normalization of Currents in Lattice QCD, NPB237(1984)111

<sup>29</sup>H.B. Nielsen, M. Ninomiya, *Absence of Neutrinos on a Lattice. 1. Proof by Homotopy Theory*, NPB185(1981)20; NPB195(1982)541 (erratum)

<sup>30</sup>Absence of Neutrinos on a Lattice. 2. Intuitive Topological Proof, NPB193(1981)173

<sup>31</sup>A No-Go Theorem for regularizing chiral fermions, PLB105(1981)219

<sup>&</sup>lt;sup>27</sup>Luuk H. Karsten, JS, *Lattice fermions: species doubling, chiral invariance and the triangle anomaly*, NPB183(1981)103

No-Go proof by Karsten <sup>32</sup> using the Poincaré-Hopf theorem which relates the index of a vector field on a manifold (in casu  $P_{\mu}(k)$ ) to the Euler number of the manifold (momentum space  $T^4$ )

- weak coupling calculations with staggered fermions by Sharatchandra, Thün and Weisz, including chiral anomaly and ratio of A-scales  $^{\rm 33}$ 

- papers by Lüscher, Münster and Weisz concerning the roughening transition (causing a  $-\pi/(12r)$  term in the potential), e.g. <sup>34</sup> <sup>35</sup> <sup>36</sup>

<sup>32</sup>L.H. Karsten, Lattice fermions in Euclidean space-time, PL104B(1981)315
 <sup>33</sup>H.S. Sharatchandra, H.J. Thun, P. Weisz, Susskind Fermions on a Euclidean Lattice. NPB192(1981)205

<sup>34</sup>G. Münster, P. Weisz, *On the Roughening Transition in Nonabelian Lattice Gauge Theories*, NPB180(1981)330

<sup>35</sup>M. Lüscher, G. Münster, P. Weisz, *How Thick Are Chromoelectric Flux Tubes* ?, NPB180(1981)1

<sup>36</sup>M. Lüscher, Symmetry Breaking Aspects of the Roughening Transition in Gauge Theories, NPB180(1981)317

with Kawamoto, Euclidean effective action at strong coupling (staggered- and Wilson-fermions):  $^{37\ 38}$ 

meson and baryon fields, large N approximation

- for r = 0 transformation to staggered-fermion form exhibits  $U(4N_{\rm f}) \times U(4N_{\rm f})$  symmetry, broken spontaneously to  $U(4N_{\rm f})$  with NG bosons and massive baryons

- for r = 1 effective action has chiral symmetry at  $M = M_{\rm cr}$  up to 2nd order in the NG fields, but at 4th order the symmetry is broken, as shown in the  $\pi$ - $\pi$  scattering amplitude

similar results (naive fermions): 39

<sup>37</sup>N. Kawamoto, JS, Effective Lagrangian and Dynamical Symmetry Breaking in Strongly Coupled Lattice QCD, NPB192(1981)100

<sup>38</sup>J. Hoek, N. Kawamoto, JS, *Baryons in the Effective Lagrangian of Strongly Coupled Lattice QCD*, NPB199(1982)495

<sup>39</sup>H. Kluberg-Stern, A. Morel, O. Napoly, B. Petersson, *Spontaneous Chiral* Symmetry Breaking for a U(N) Gauge Theory on a Lattice, NPB190(1981)504

#### 1983

spin, flavor and symmetry group of ('reduced') staggered fermions with method for loop calculation  $^{\rm 40}$ 

related work  $^{\rm 41}$   $^{\rm 42}$   $^{\rm 43}$ 

PhD Jaap Hoek 44 45 46

 $^{40}$ Cees van den Doel, JS, Dynamical Symmetry Breaking in Two Flavor SU(N) and SO(N) Lattice Gauge Theories, NPB228(1983) 122

<sup>41</sup>F. Gliozzi, *Spinor Algebra of the One Component Lattice Fermions*, NPB204(1982)419

<sup>42</sup>H. Kluberg-Stern, A. Morel, O. Napoly and B. Petersson, *Flavors of Lagrangian Susskind Fermions*, NPB220 (1983)447

<sup>43</sup>P. Becher & H. Joos, *Geometric fermions* 

<sup>44</sup>Thesis: Effective Action Calculation in Lattice QCD, Amsterdam, 1983
 <sup>45</sup>Strong Coupling Expansion of the Generating Functional for Gauge
 Systems on a Lattice with Arbitrary Sources, J.Comp.Phys. 49(1983)265
 <sup>46</sup>Strong Coupling Expansion of the SU(3) and U(3) Effective Actions,
 J.Comp.Phys. 54(1984)245

September: Cargèse Summer Institute

Wilson lectured 47

- I mentioned to him my work on LGT at UCLA in 1972-1973

- Wilson's generous answer: ask your promotor to send me a description of what you did and I shall mention it in future reviews

and so it was done

<sup>47</sup>K.G. Wilson, *Monte Carlo Renormalization Group and the Three-Dimensional Ising Model*, Progress in Gauge Field Theory, NATO Sci.Ser.B 115 (1984) 589-604 Swift in  $^{48}$  used essentially the same method as I did (cf.  $^{49})$  for deriving currents for the 'spin wave paper'

a U(1) version was analyzed a few years later by Hands and Carpenter  $^{\rm 50}$ 

<sup>48</sup>P.V.D. Swift, The Electroweak Theory on the Lattice, PLB145(1984)256
 <sup>49</sup>JS, Fermions on a lattice, Act.Phys.Pol.B17(1986)531
 <sup>50</sup>S.J. Hands, D.B. Carpenter, Lattice Sigma Model and Fermion Doubling, NPB266(1986)285

## $\gtrsim 1984$

lots of inspiring papers were being published, on Monte-Carlo-RG scaling, string tension and potential, hadron spectrum, critical temperature

- desire for similar computational work in Amsterdam and the arrival of the Cyber 205 supercomputer led to taking part in larger collaborations  $^{51}$   $^{52}$ 

Peter Hasenfratz sent me gauge-field configurations:

"they are beautiful !"

 $^{51}A.$  König, K.H. Mütter, K. Schilling, JS, Large distance propagators for hadrons on a 56  $\times$  16  $^3$  lattice, PLB157(1985)421

<sup>52</sup>K.C. Bowler, F. Gutbrod, P. Hasenfratz, U. Heller, F. Karsch, R.D. Kenway, I. Montvay, G.S. Pawley, JS, D.J. Wallace, *The* β-function and potential at  $\beta = 6.0$  and 6.3

#### 1984 - 1986

more on staggered fermions: 53 54 55 56 57

related work: 58 59

<sup>53</sup>Maarten F.L. Golterman, JS, *Relation Between QCD Parameters on the Lattice and in the Continuum*, PLB140(1984)392

<sup>54</sup>M.F.L. Golterman, JS, *Self-Energy and Flavor Interpretation of Staggered Fermions*, NPB245(1984)61

<sup>55</sup>M.F.L. Golterman, JS, *Lattice Baryons With Staggered Fermions*, NPB255(1985)328

<sup>56</sup>M.F.L. Golterman, *Staggered Mesons*, NPB273(1986)663

<sup>57</sup>M.F.L. Golterman, *Irreducible Representations of the Staggered Fermion Symmetry Group*, NPB278(1986)417

<sup>58</sup>G.W. Kilcup, S.R. Sharpe, *A Tool Kit for Staggered Fermions*, NPB283(1987)493

<sup>59</sup>A. Coste, C. Korthals Altes, O. Napoly, *Calculation of the Nonabelian Chiral Anomaly on the Lattice*, NPB289(1987)645  $\gtrsim$  1986 U(1) problem and topological charge can instantons explain the large mass of the  $\eta'$  meson?

anomalous divergence of the axial flavor-singlet current

$$egin{array}{rcl} \partial_\mu \left( ar{\psi} i \gamma_\mu \gamma_5 \psi 
ight) &=& 2m \, ar{\psi} i \gamma_5 \psi + 2i N_{
m f} \, q & U(1) \; {
m current} \ q &=& rac{1}{32 \pi^2} \, \epsilon_{\kappa \lambda \mu 
u} \; {
m Tr} \; {
m G}_{\kappa \lambda} \, {
m G}_{\mu 
u} & {
m topological} \; {
m charge \; density} \end{array}$$

derived lattice version <sup>60</sup> of the Witten-Veneziano relation

$$\begin{split} m_{\eta'}^2 - \frac{1}{2} m_{\eta}^2 - \frac{1}{2} m_{\pi^0}^2 &= \frac{6}{f_{\pi}^2} \chi \qquad \qquad N_{\rm f} = 3 \\ \chi &= \int d^4 x \langle q(x) \, q(0) \rangle_{|_{\rm quenched}} \\ &\simeq (180 \, {\rm MeV})^4 \quad {\rm topological \ susceptibility} \end{split}$$

U(1) problem and topological charge (cont.)

which led to a 'fermionic determination' of topological charge  $^{61\ 62}$ 

$$\overline{Q} = \operatorname{Tr} \kappa_P m \gamma_5 (\not \! D + m + M_{\rm cr} - W)^{-1} \quad \text{Wilson fermions}$$
$$= \frac{1}{4} \operatorname{Tr} \kappa_P m \Gamma_5 (\not \! D + m)^{-1} \quad \text{staggered fermions}$$

<sup>61</sup>J.C. Vink, Staggered Fermions, Topological Charge and Topological Susceptibility in Lattice QCD, PLB212(1988)483; and references there-in <sup>62</sup>M.L. Laursen, JS, J.C. Vink, Small-Scale Instantons, Staggered Fermions and the Topological Susceptibility, NPB343(1990)522; and references there-in U(1) problem and topological charge (cont.)

other definitions: geometrical <sup>63</sup> <sup>64</sup> 'cooling' e.g. <sup>65</sup> <sup>66</sup> <sup>67</sup>

<sup>63</sup>M. Lüscher, Topology of Lattice Gauge Fields, CMP85(1982)39
 <sup>64</sup>A. Philips, D. Stone, Lattice Gauge Fields, Principal Bundles and the Calculation of Topological Charge, CMP103(1986)399

<sup>65</sup>Y. Iwasaki, T. Yosie, *Instantons and Topological Charge in Lattice Gauge Theory*, PLB131(1983)159

<sup>66</sup>M. Teper, Instantons in the Quantized SU(2) Vacuum: A Lattice Monte Carlo Investigation, PLB162(1985)357; The Topological Susceptibility in SU(2) Lattice Gauge Theory: An Exploratory Study, PLB171(1986)81

<sup>67</sup>E-M. Ilgenfritz, M.L. Laursen, G. Schierholz, M. Muller-Preussker, H. Schiller, *First Evidence for the Existence of Instantons in the Quantized SU(2) Lattice Vacuum*, NPB268(1986)693 in a model with only quarks or leptons with (possibly!) decoupled doublers, anomalies may be carried by the Higgs field  $^{68}$ 

decoupling of doublers was studied by focusing on Higgs-fermion interaction without gauge fields

<sup>68</sup>E. D'Hoker and E. Fahri, *Decoupling a fermion in the standard* electro-weak theory, NPB248(1984)77

### $\gtrsim$ 1986 Standard Model: Wilson-Yukawa (cont.)

SU(2) invariant combinations  $\Phi^{\dagger}P_{\rm L}\Psi$  and  $\bar{\Psi}\Phi P_{\rm R}$  put in the Wilson-Yukawa 'mass term'

$$S_{\text{Fmass}} = (y + 4w)\bar{\Psi}_{x}(\Phi_{x}P_{\text{R}} + \Phi_{x}^{\dagger}P_{\text{L}})\Psi_{x}$$
$$-\frac{w}{2}\sum_{\mu} \left[\bar{\Psi}_{x}\Phi_{x}P_{\text{R}}\Psi_{x+\hat{\mu}} + \bar{\Psi}_{x}\Phi_{x+\hat{\mu}}^{\dagger}P_{\text{L}}\Psi_{x+\hat{\mu}} + \text{h.c.}\right]$$

- continuum limit at the weak coupling FM(W)-PMW phase boundary: as expected, doublers are heavier than target fermions but not really removable because of triviality <sup>69</sup>

<sup>&</sup>lt;sup>69</sup>W. Bock, A.K. De, C. Frick, K. Jansen, T. Trappenburg, *Search for an upperbound of the renormalized Yukawa coupling in a lattice fermion-Higgs model*, NPB371(1992)683

## $\gtrsim$ 1986 Standard Model: Wilson-Yukawa (cont.)

- at the disjoint strong coupling FM(S)-PMS phase boundary: lots of papers by groups of aficionados (& a nice Amsterdam-Jülich collaboration) with the conclusion:

- removing doublers possible but on the way out they become fermion-Higgs bound-state SU(2) singlets !  $^{70\ 71}$ 

- only at the FM(W)-PMW phase boundary it is possible to *approximate* the physics of the Standard Model (doublers still present)

mirror fermion models appear to fare better in this respect <sup>72</sup>

<sup>70</sup>M.F.L. Golterman, D. Petcher, JS, *Fermion-interactions in models with Strong Wilson-Yukawa Couplings*, NPB370(1992)51 and references therein <sup>71</sup>W. Bock, A.K. De, JS, *Fermion masses at strong Wilson-Yukawa Couplings in the Symmetric Phase*, NPB388(1992)243 and references therein <sup>72</sup>C. Frick, L. Lin, I. Montvay, G. Münster, M. Plagge, T. Trappenberg, H. Wittig, *Numerical simulation of heavy fermions in an*  $SU(2)_{\rm L} \otimes SU(2)_{\rm R}$ *symmetric Yukawa model*, NPB397(1993)431  $\gtrsim$  1992 Standard Model: staggered fermion

... it did not work out well



## $\gtrsim 1992$ changing the game

D.B. Kaplan: domain walls (Y. Shamir use for QCD)R. Narayanan & H. Neuberger: overlapGinsparg-Wilson: Lüscher

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#### meanwhile

we enjoyed work, e.g. by

... too many names, gave up