

# Low-Lying Spectrum of Two-Dimensional Adjoint QCD from the Lattice

William Jay, Manki Kim, Patrick Oare, Phiala Shanahan, Neill Warrington

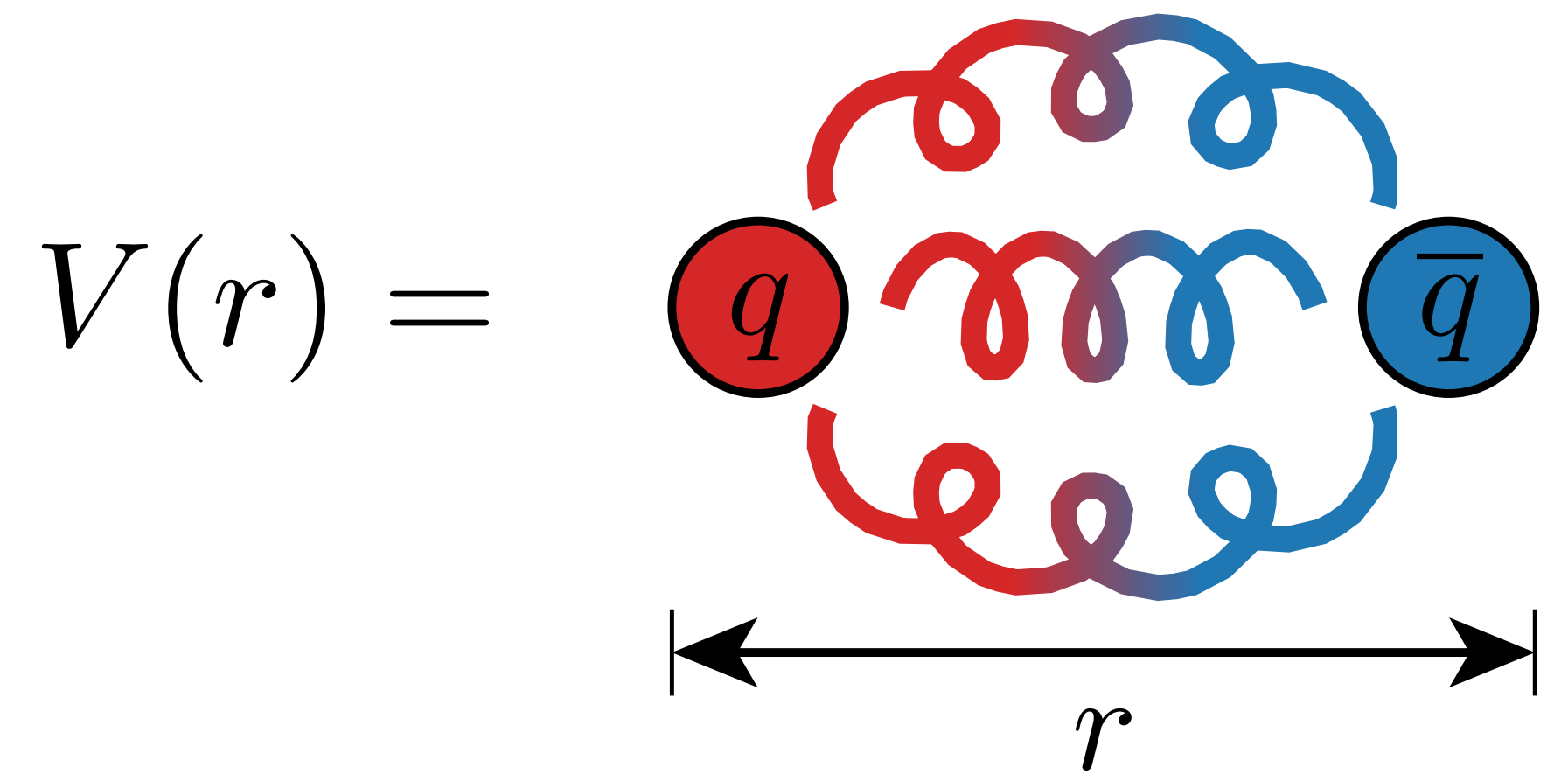
July 30th, 2024

Lattice 2024, University of Liverpool



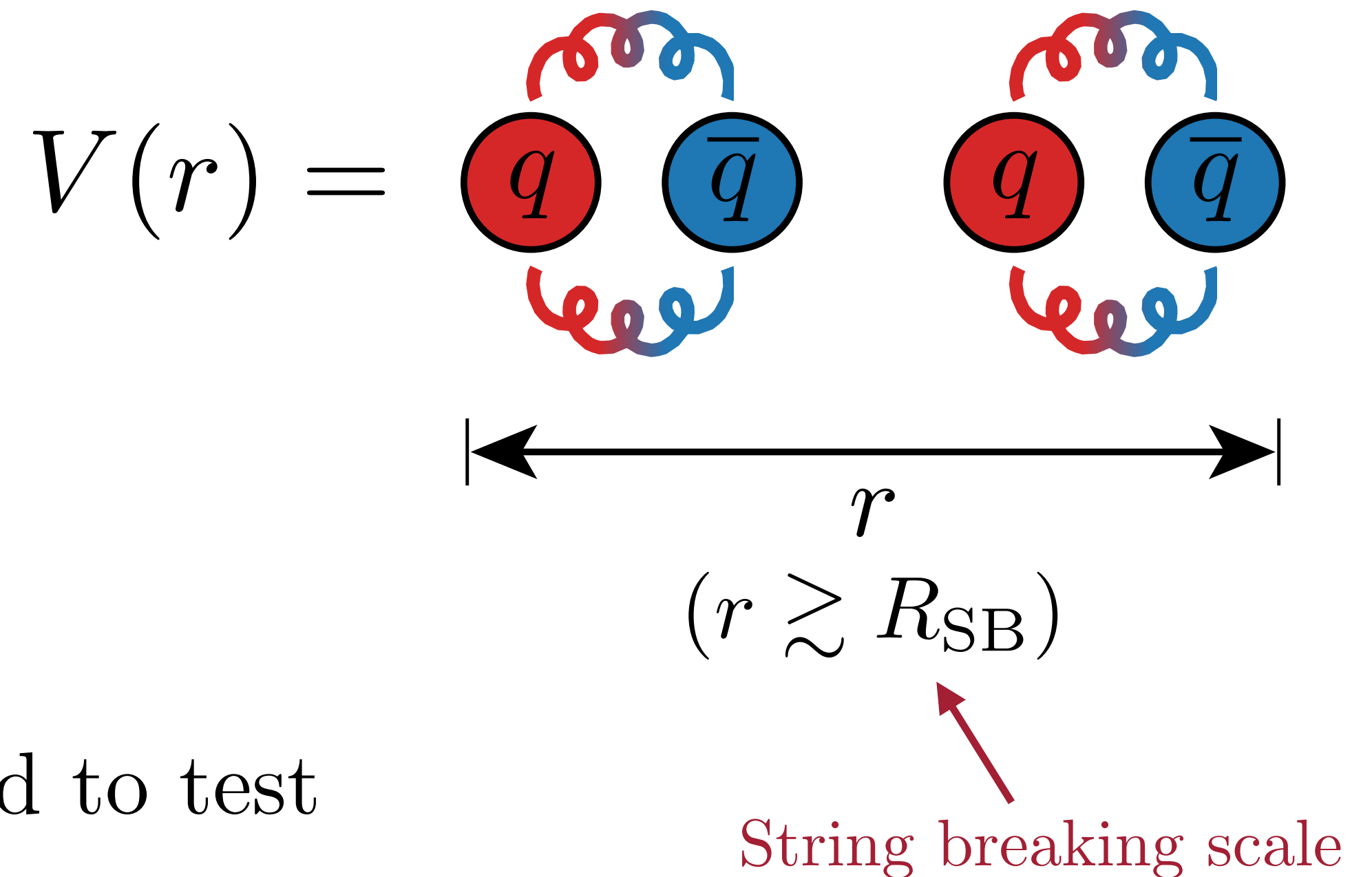
# Confinement in Gauge Theories

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- With this definition, pure  $SU(N)$  gauge theory is confining.



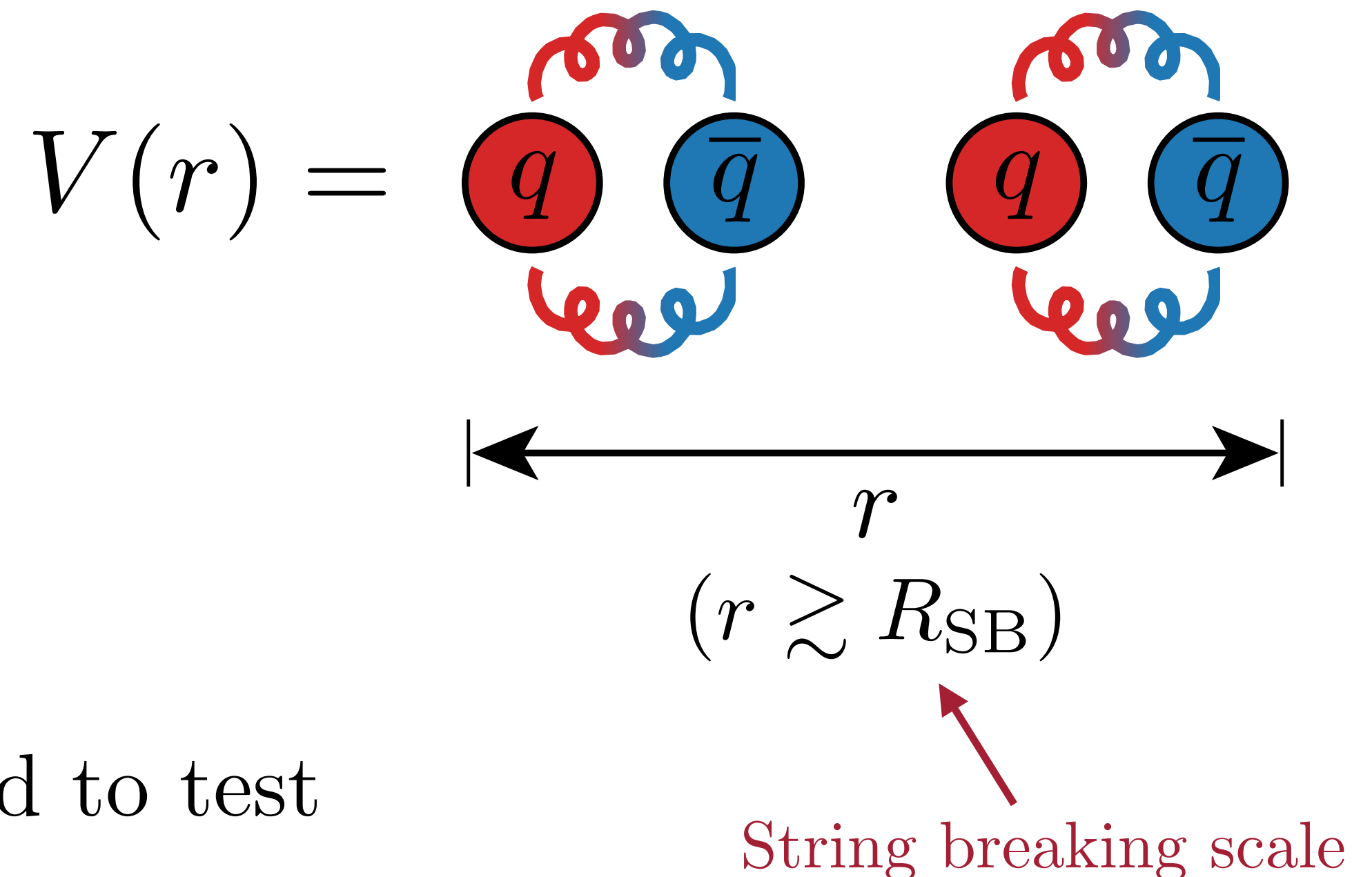
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- QCD undergoes **string breaking** at large  $r$ .
  - ▶ Energetically favorable for sea quarks to bind to test charges and create a pair of mesons.
- How do adjoint fields influence the confining nature of a theory? Pure  $SU(N)$  gauge theory confines; what about theories with adjoint matter?



# Two-Dimensional Adjoint QCD (QCD<sub>2</sub>)

- In  $d = 2$  spacetime dimensions, pure SU(N) gauge theory is solvable and does not have propagating local degrees of freedom.

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- The simplest two-dimensional theory with dynamical adjoint degrees of freedom is **two-dimensional adjoint QCD** (QCD<sub>2</sub>).
  - ▶ Theory of a Majorana fermion coupled to an SU(N) gauge field in the adjoint representation with Minkowski action:

$$S_{\text{QCD}_2}[\psi, G] = \int d^2x \text{Tr} \left[ \frac{1}{2g^2} G_{\mu\nu}(x) G^{\mu\nu}(x) + \bar{\psi}(x) (i\gamma^\mu D_\mu - m) \psi(x) \right]$$

Majorana fermion field  
 $\psi(x) = \psi^a(x) t^a$



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Z. Komargodski *et. al.*  
Symmetries and strings of  
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- Toy model for confinement: confining for  $m > 0$ , deconfining for  $m = 0$ .

Majorana fermion field  
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J. Donahue, *et. al.* Confining  
Strings, Infinite Statistics and  
Integrability. *Phys. Rev. D*  
101 (2020) 8, 081901.

R. Dempsey *et. al.* Adjoint  
Majorana QCD<sub>2</sub> at Finite N.  
JHEP 04 (2023) 107.

(Computation of the  
spectrum with DLCQ  
at  $N = 2,3,4$ ).

# QCD<sub>2</sub> on the lattice

G. Bergner, S. Piemonte, M. Ünsal  
JHEP 07 (2024) 048.

- First study of QCD<sub>2</sub> with LGT was published this year.
  - ▶ Computed Polyakov loop  $\langle P \rangle$ , static quark potential, chiral condensate, and string tension for a number of  $(N, g, m)$ .
  - ▶ Used two discretizations: Wilson and reweighted overlap.



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  - ▶ Used two discretizations: Wilson and reweighted overlap.
- Our calculation uses a Wilson discretization on a Euclidean  $L \times T$  lattice, and the Wilson action for the gauge field  $U_\mu(x)$ .

▶ Dirac operator:

$$D_W(x, y) = 1 \delta_{x,y} - \kappa \sum_{\mu=1}^2 \left[ V_\mu(x)(1 - \gamma_\mu)\delta_{x+\hat{\mu},y} + V_\mu^T(y)(1 + \gamma_\mu)\delta_{x-\hat{\mu},y} \right]$$

Hopping parameter  $\kappa = \frac{1}{2m + 4}$


adjoint links  $V_\mu^{ab}(x) \equiv 2 \text{Tr}[U_\mu^\dagger(x)t^a U_\mu(x)t^b]$

# LGT Setup ( $N = 2$ colors)

$L$	$T$	$\beta$	$\kappa$	$N_{\text{cfgs}}$
10	10	3.125	0.138889	90
12	12	4.5	0.15	90
14	14	6.125	0.159091	90
16	16	8.0	0.166667	90
18	18	10.125	0.173077	90
★ 20	20	12.5	0.178571	90
22	22	15.125	0.183333	90
24	24	18.0	0.1875	76
26	26	21.125	0.191176	58
28	28	24.5	0.194444	45

# Majorana Fermions on the Lattice

$$\begin{aligned}\langle \mathcal{O}_1 \dots \mathcal{O}_k \rangle &= \frac{1}{Z} \int DU D\psi D\bar{\psi} e^{-S_g[U] - \frac{1}{2} \int \bar{\psi} D_W \psi} \mathcal{O}_1 \dots \mathcal{O}_k \\ &= \frac{1}{Z} \int DU e^{-S_g[U]} \text{pf } D_W \langle \mathcal{O}_1 \dots \mathcal{O}_k \rangle_F\end{aligned}$$

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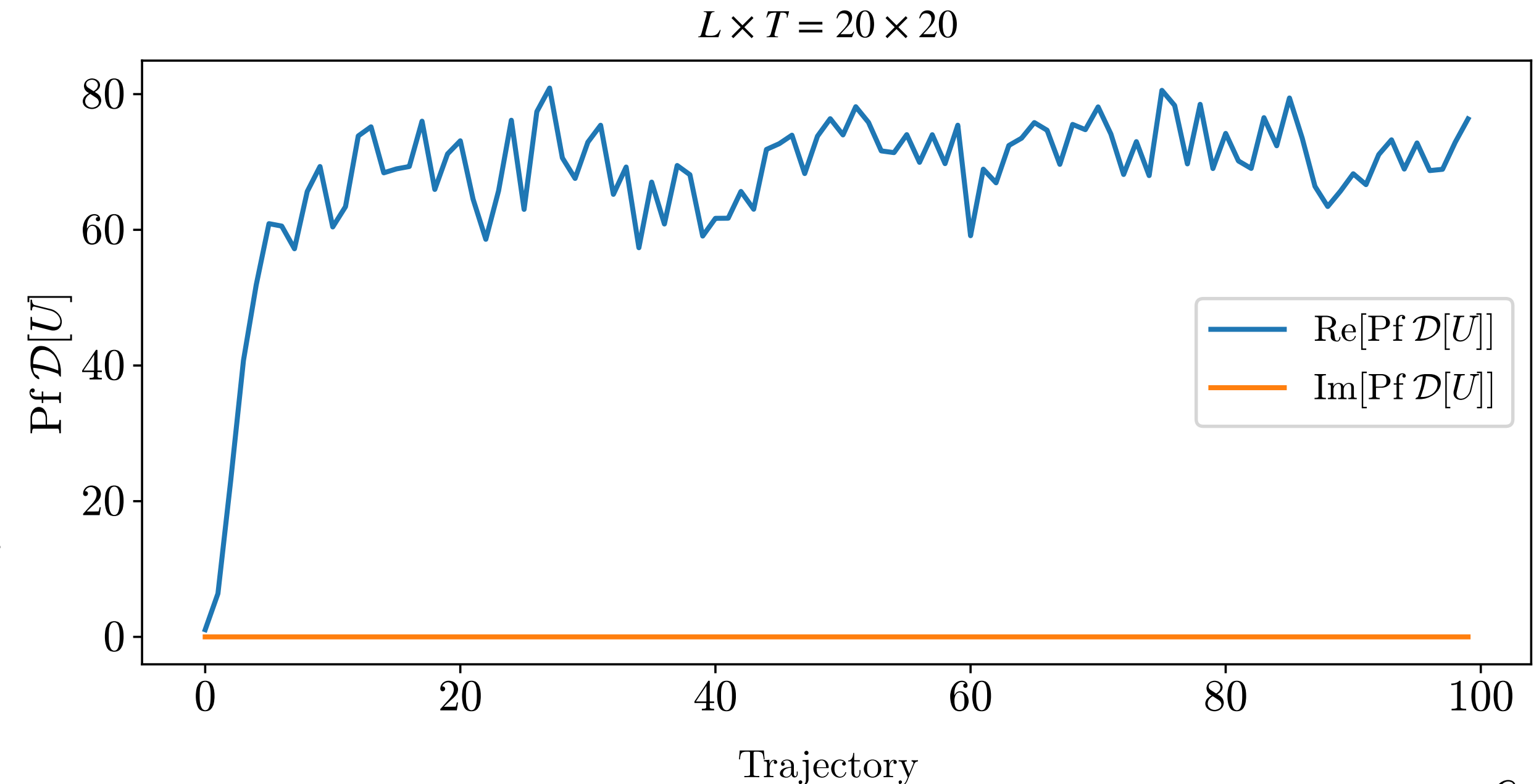
- The Pfaffian  $\text{pf } D_W$  is *not necessarily positive*: how to interpret as a probability density for sampling?
  1. Check or prove that  $\text{pf } D_W > 0$ .
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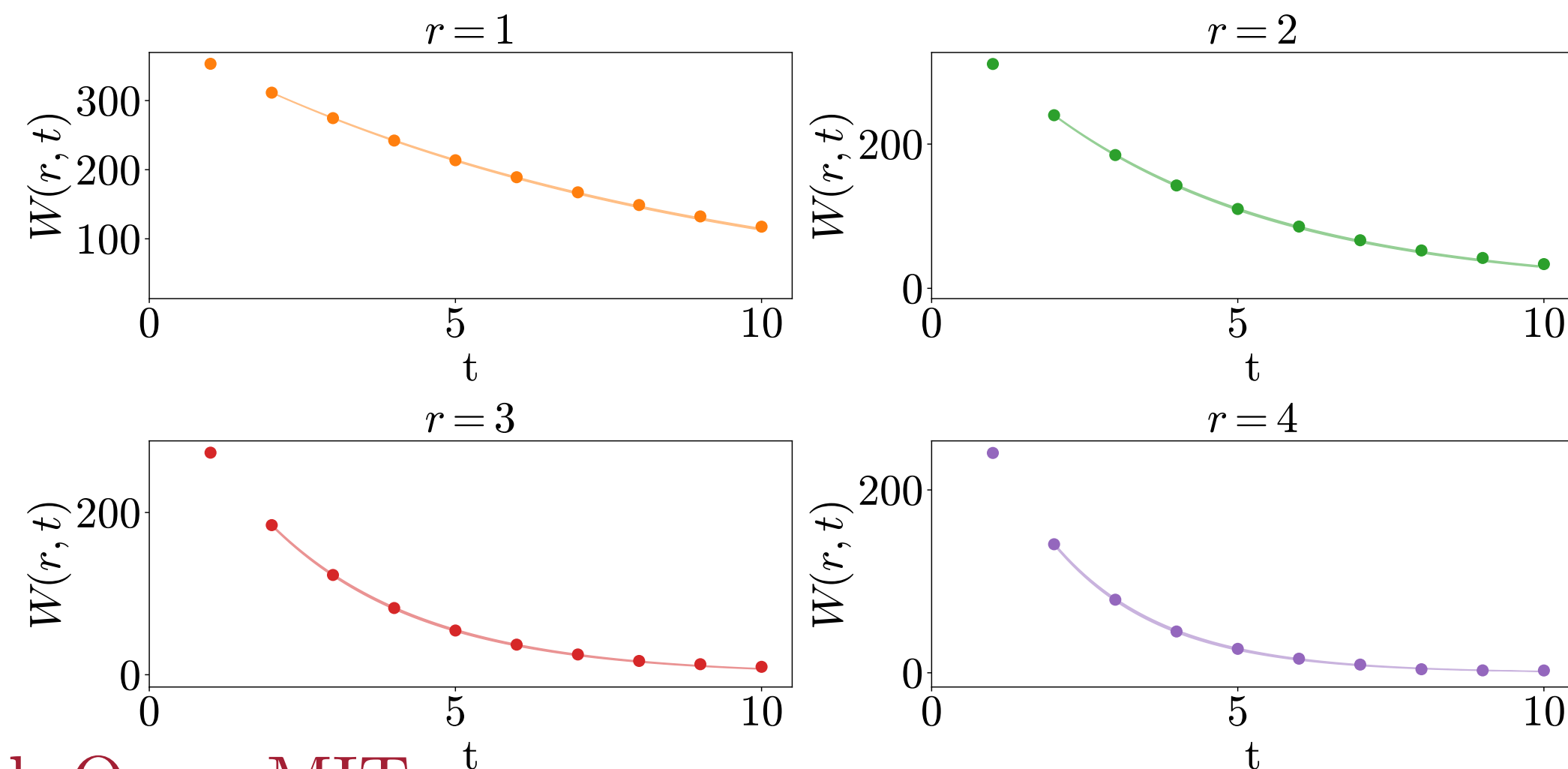
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# Static Quark Potential

- The static quark potential  $V(r)$  is extracted by fitting the  $r \times t$  Wilson loop  $W(r, t)$  to the functional form,

$$W(r, t) = Ce^{-V(r)t}$$



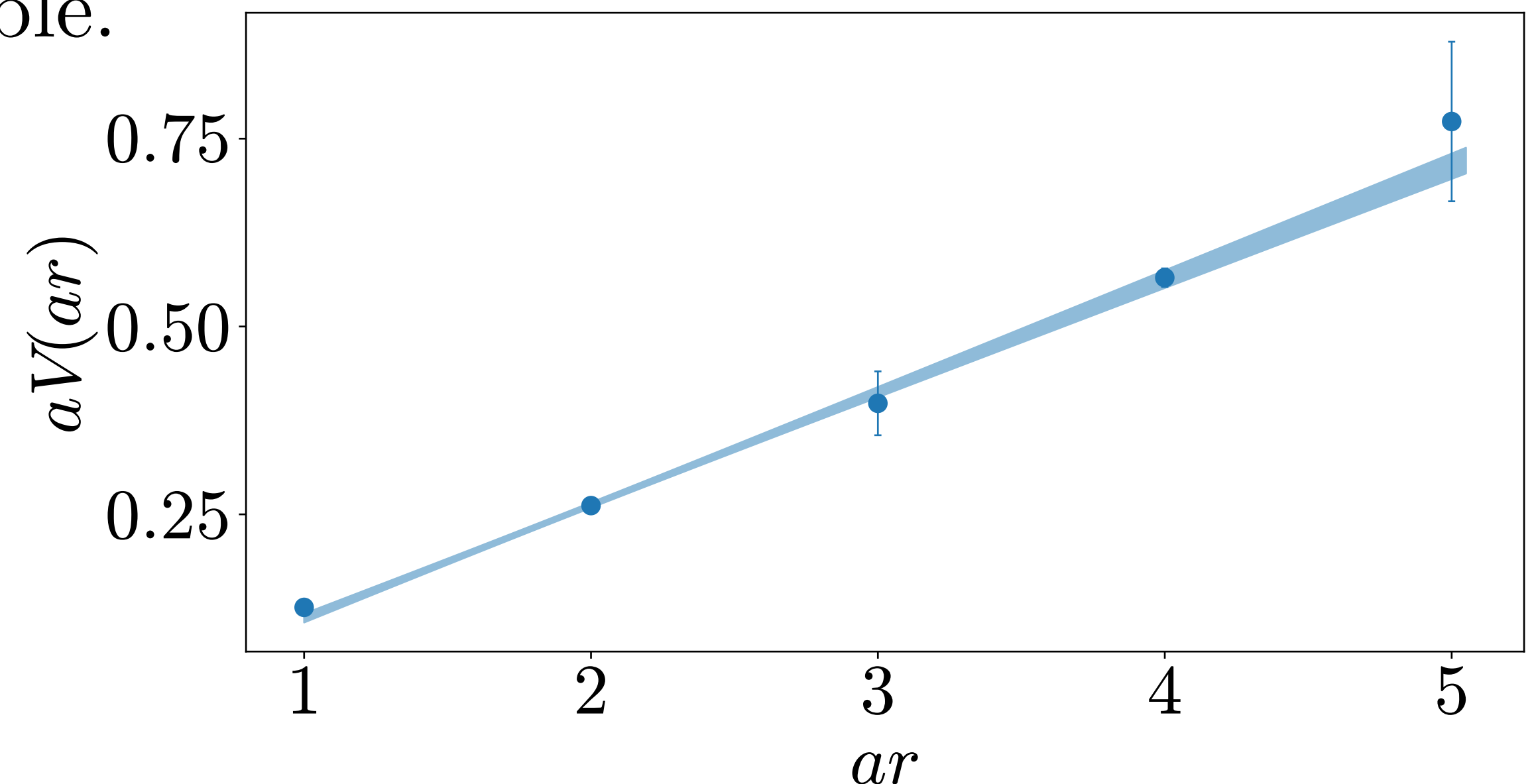
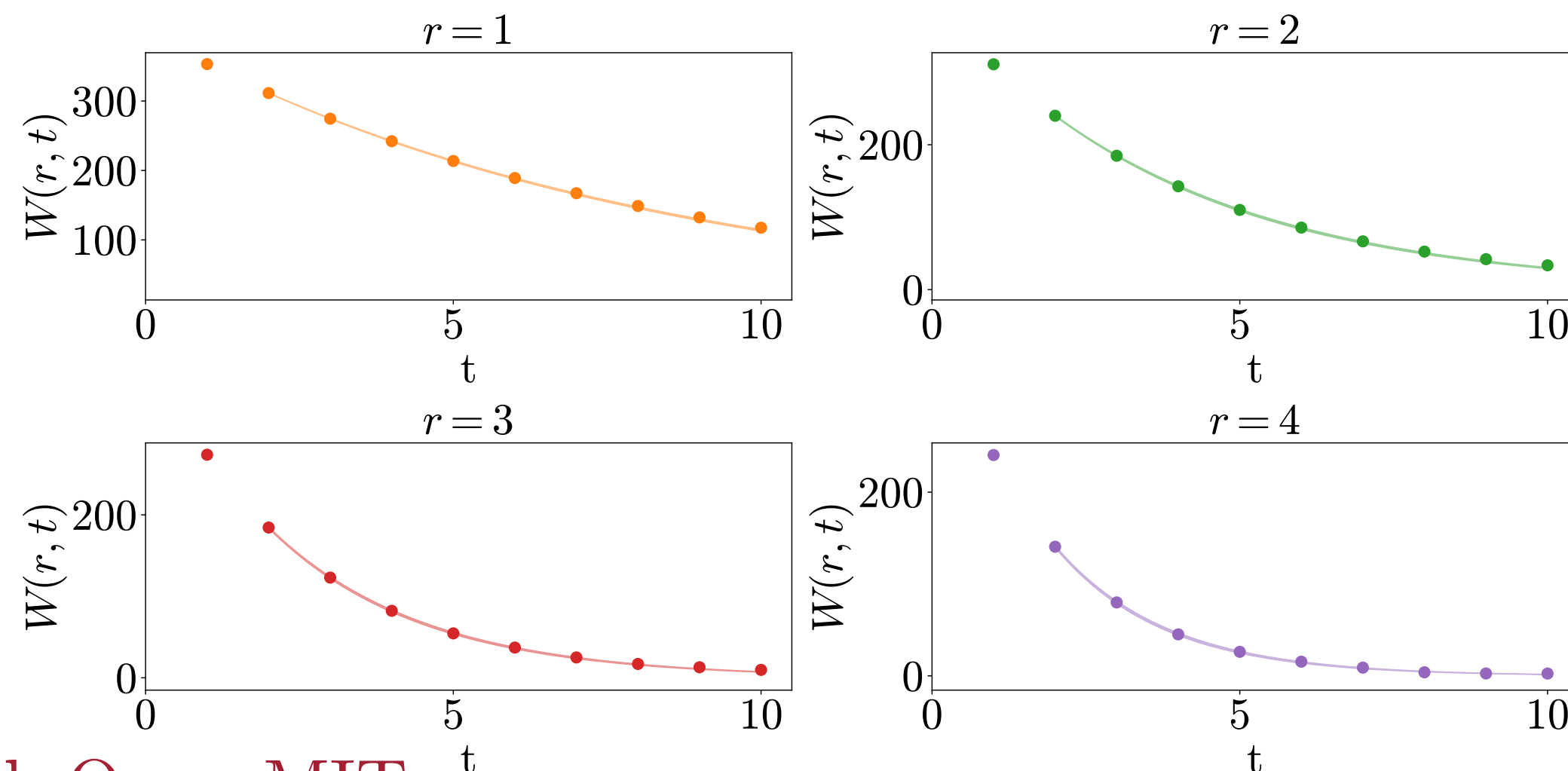


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- $V(r)$  is fit to the linear ansatz  $V(r) = A + \sigma r$  to extract the string tension  $\sigma$ .
  - $\sigma$  is used to set the scale of each ensemble.



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  - ▶ In  $d = 2$ ,  $\{1, \gamma^1\}$  are parity even, and  $\{\gamma_5, \gamma^0\}$  are parity odd.

$$\chi_\Gamma(t; p) = \frac{1}{\sqrt{L}} \sum_x e^{-ipx} \bar{\psi}(x, t) \Gamma \psi(x, t)$$

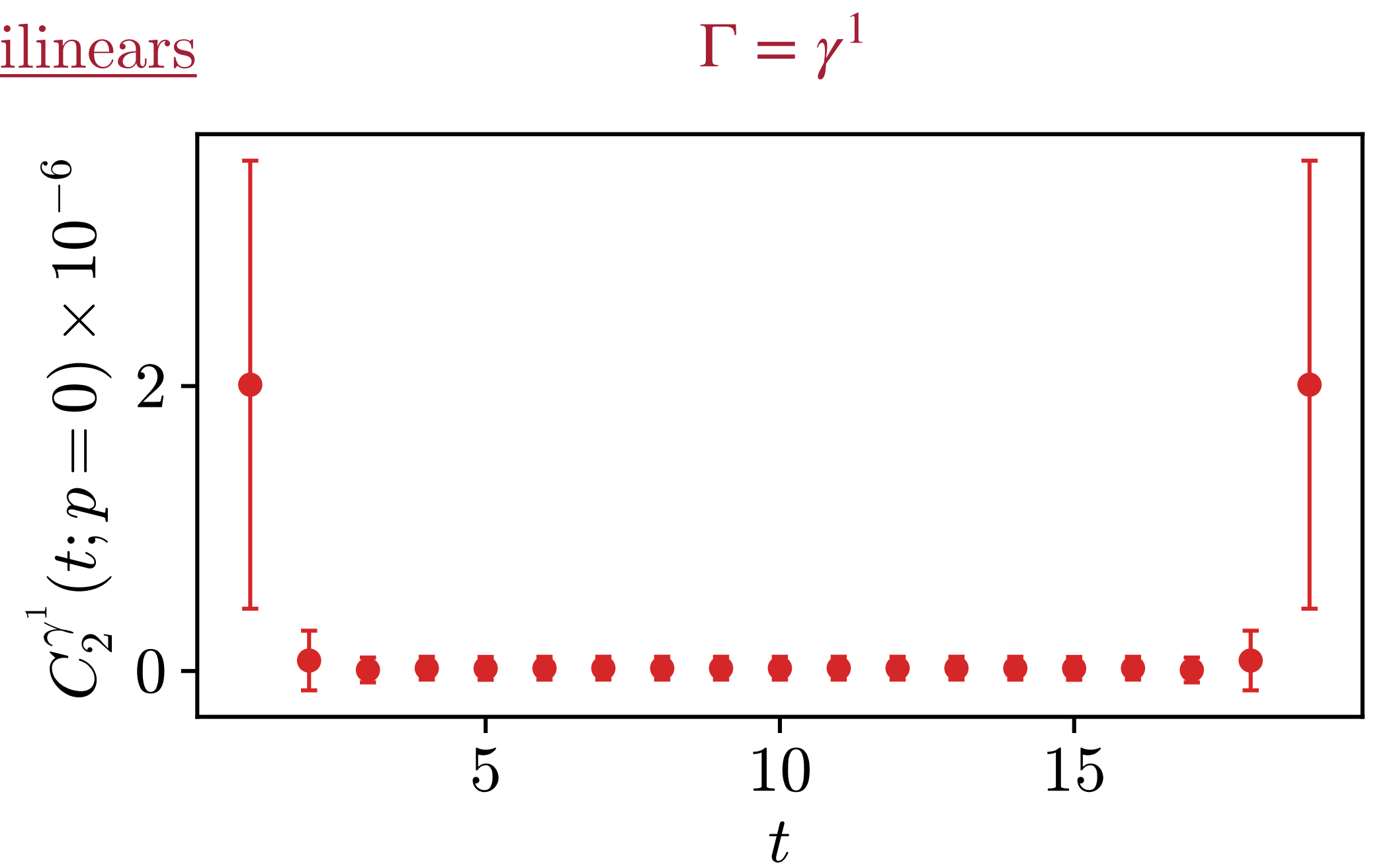
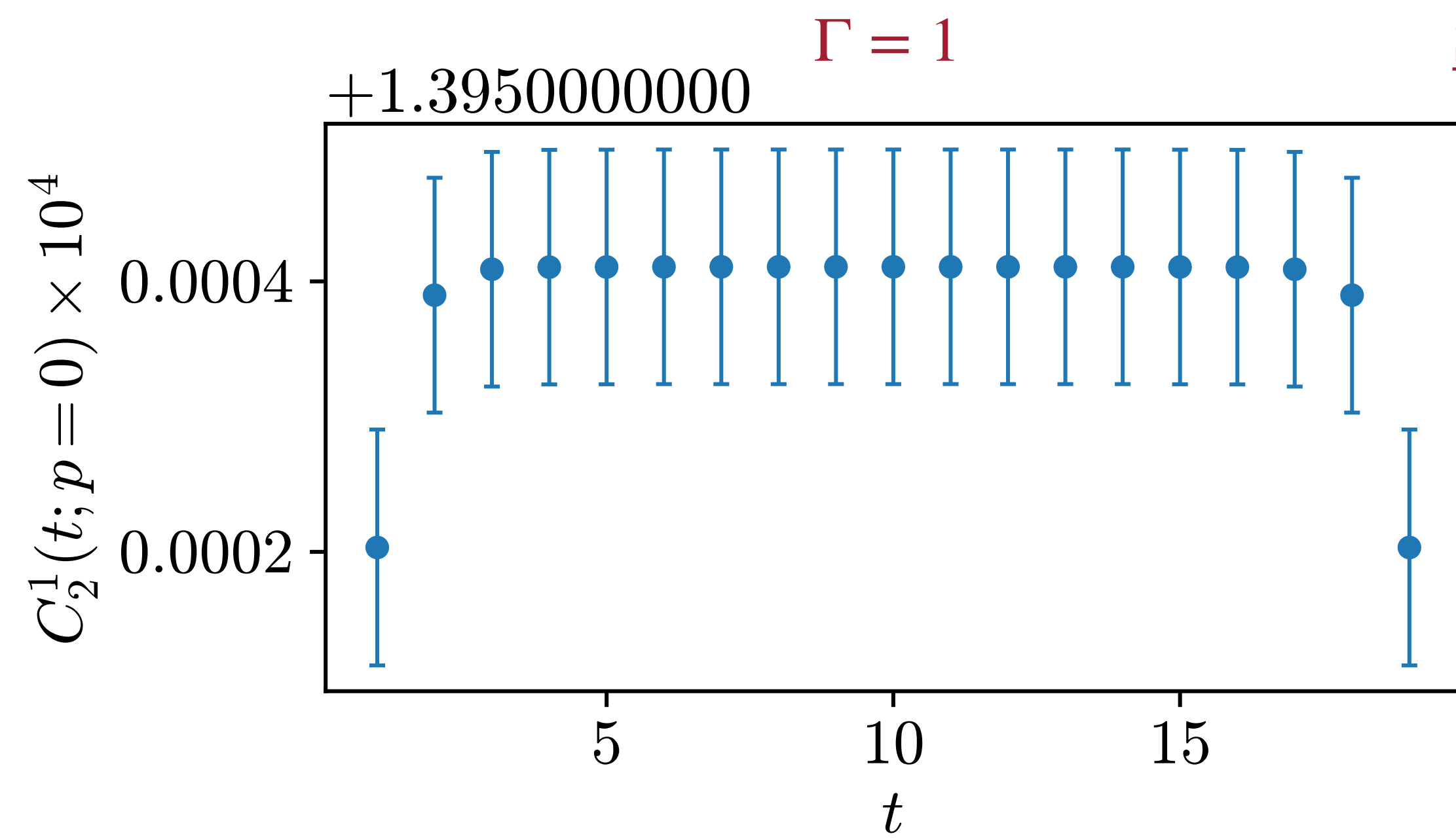
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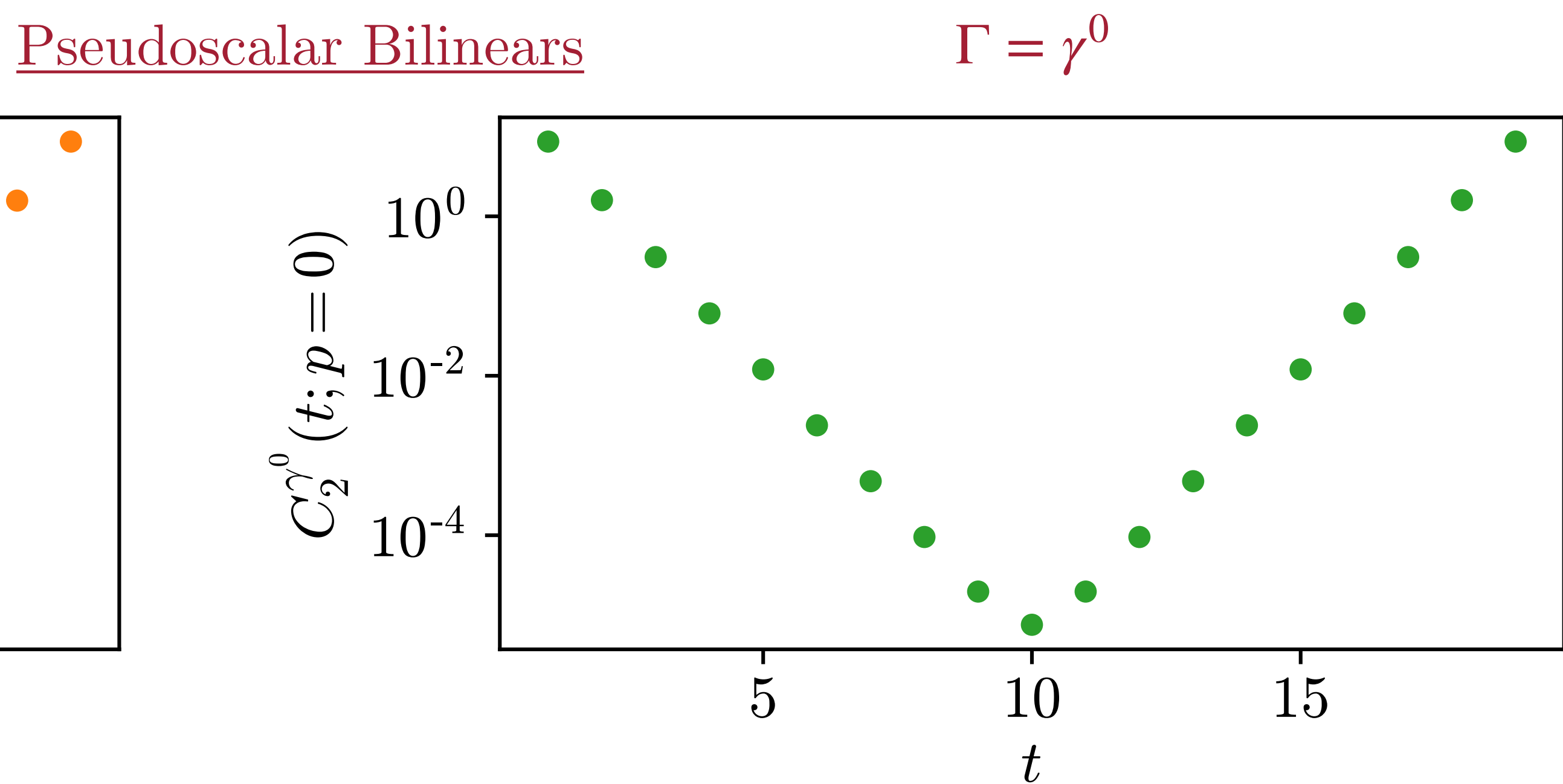
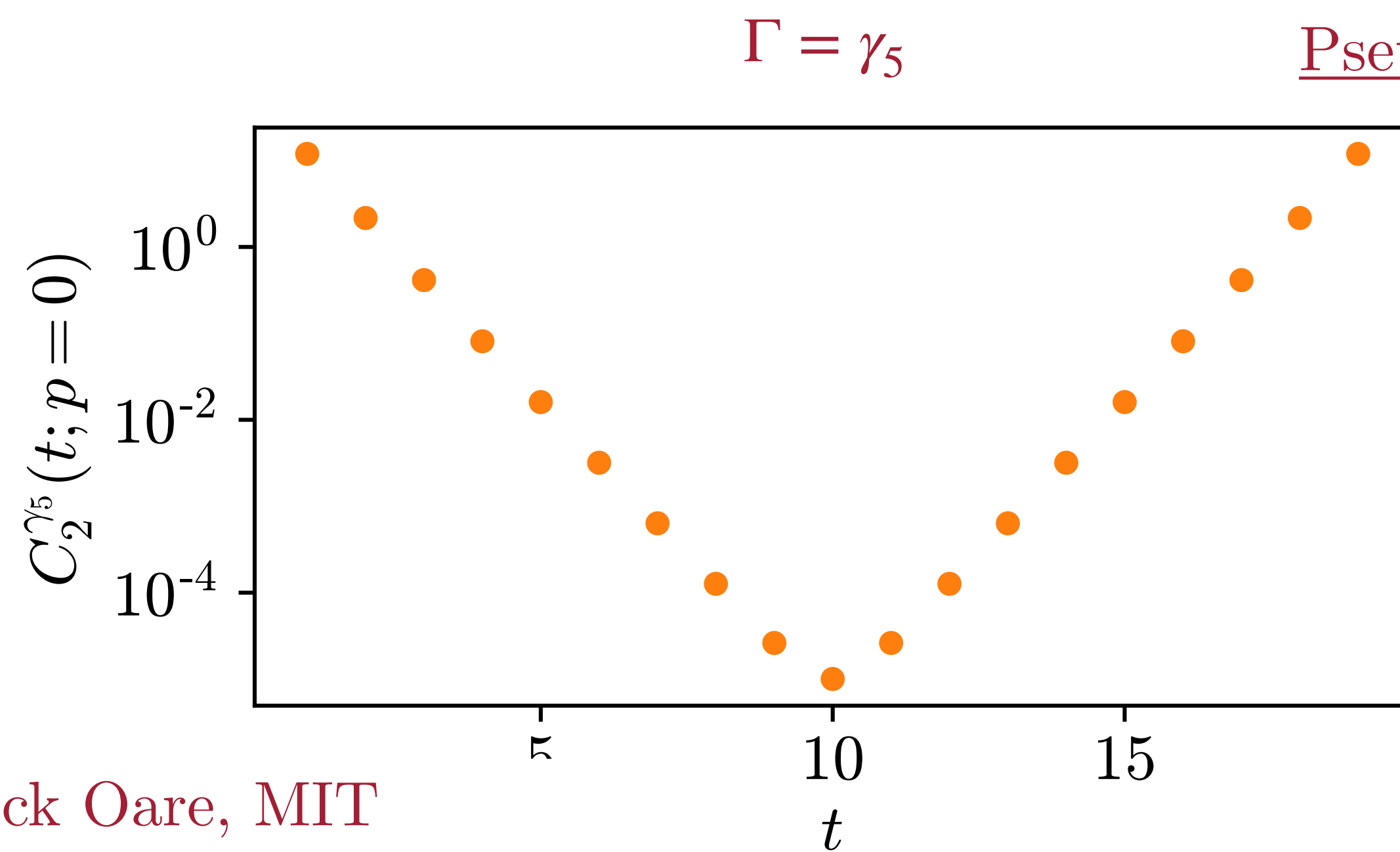
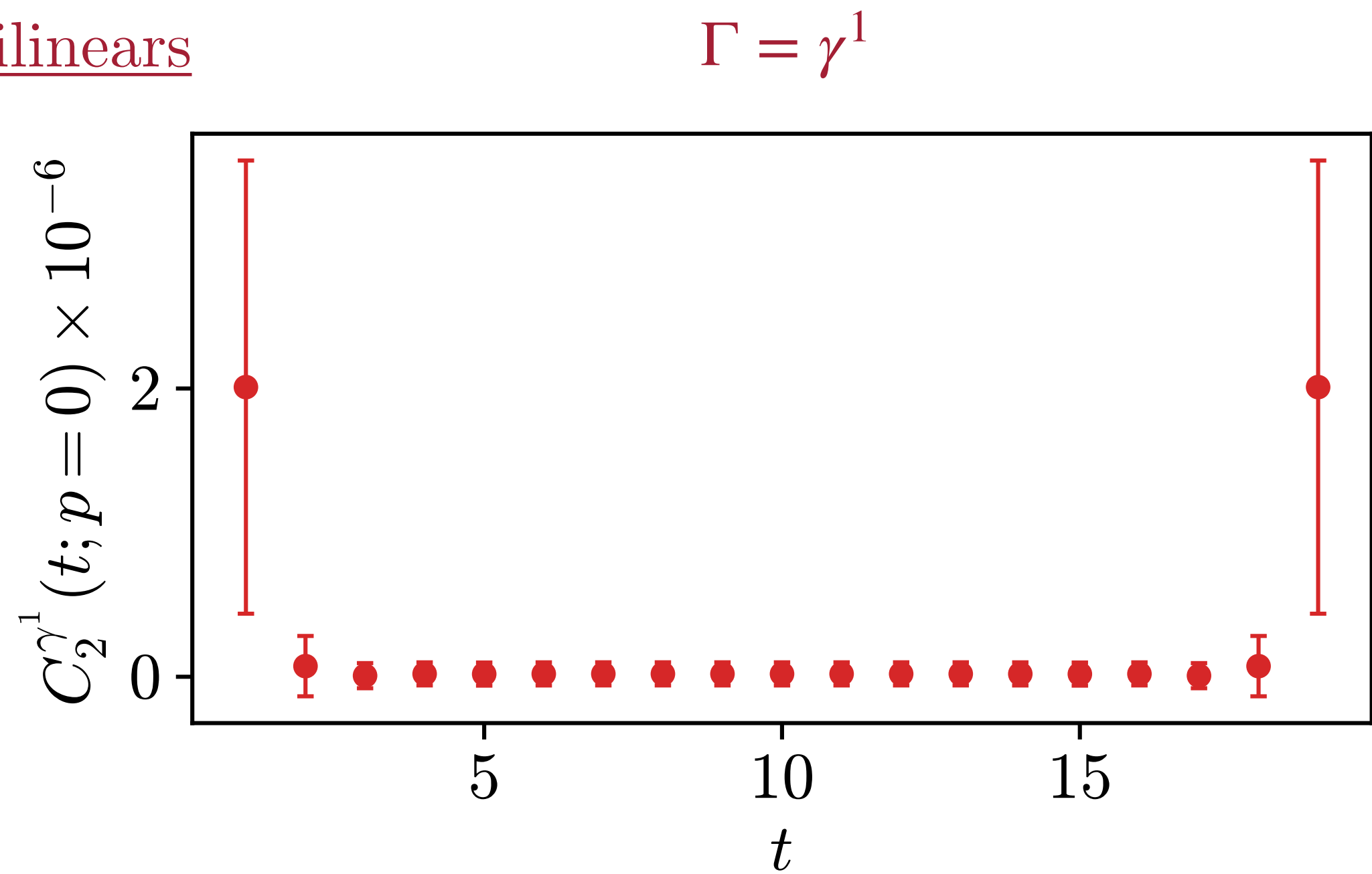
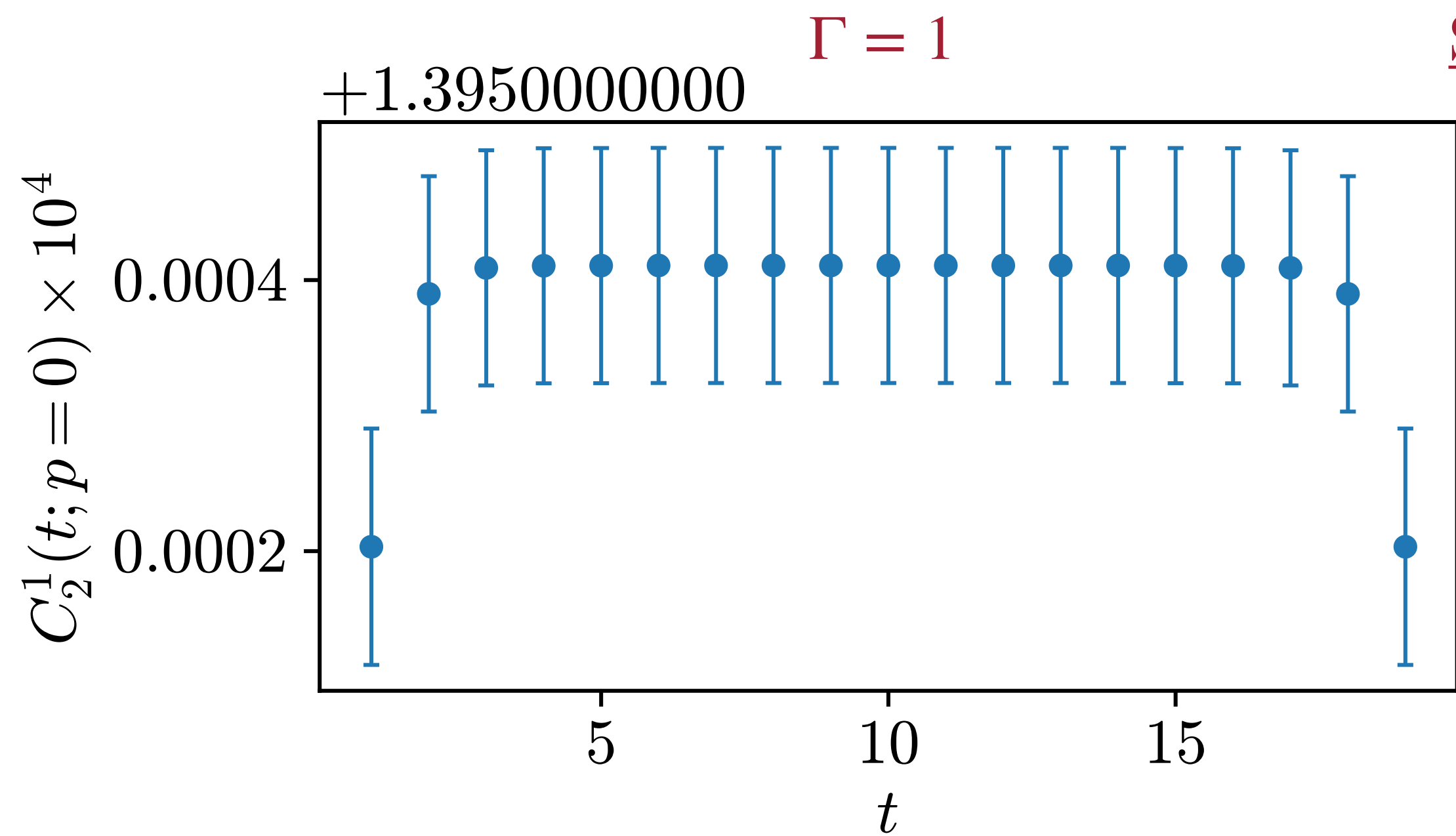
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- Compute two-point correlators at fixed momentum  $p$  to extract the ground state in each sector.

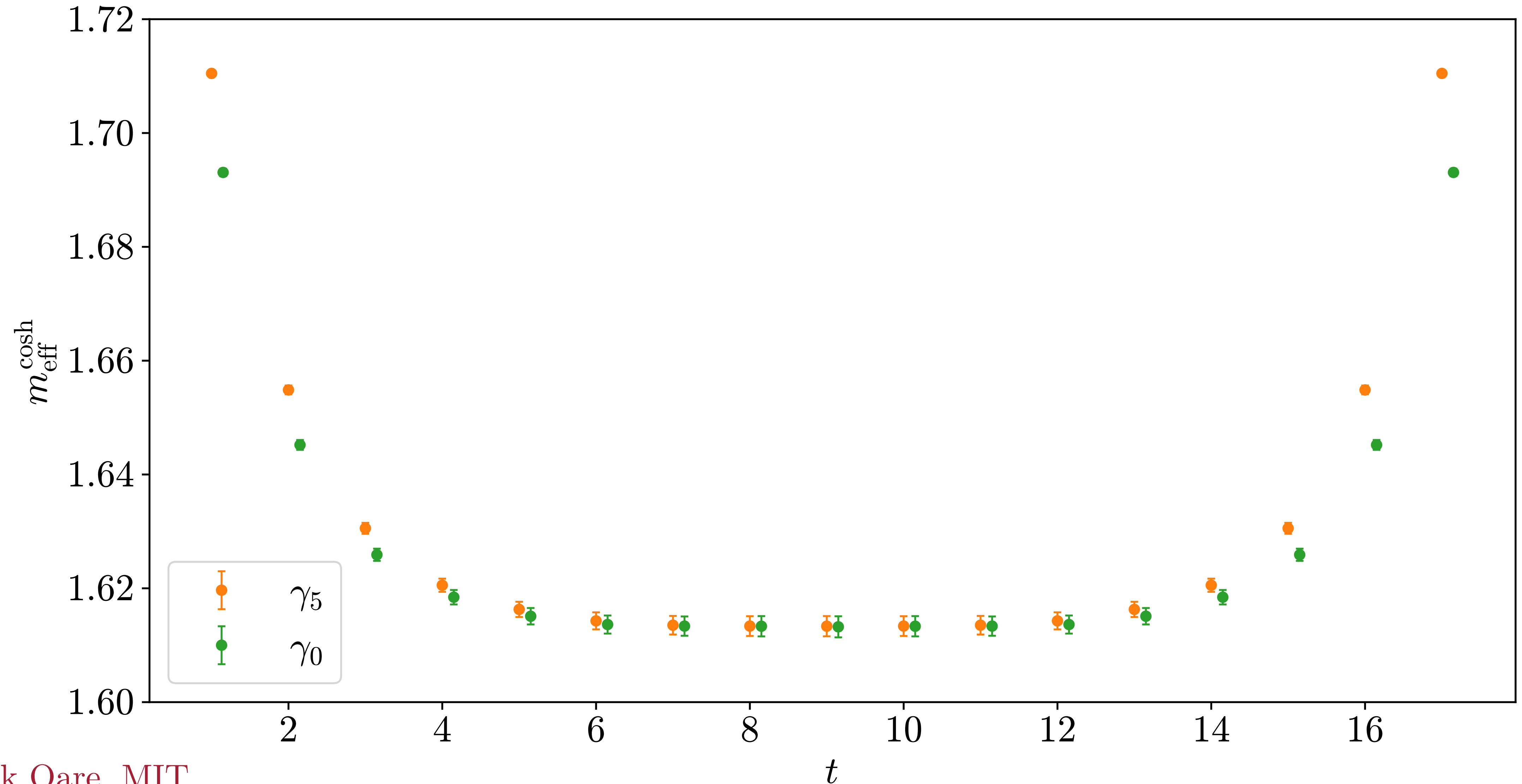
$$C_2^\Gamma(t; p) = \frac{1}{T} \sum_s \langle \chi_\Gamma(t + s; p) \bar{\chi}_\Gamma(s; p) \rangle.$$



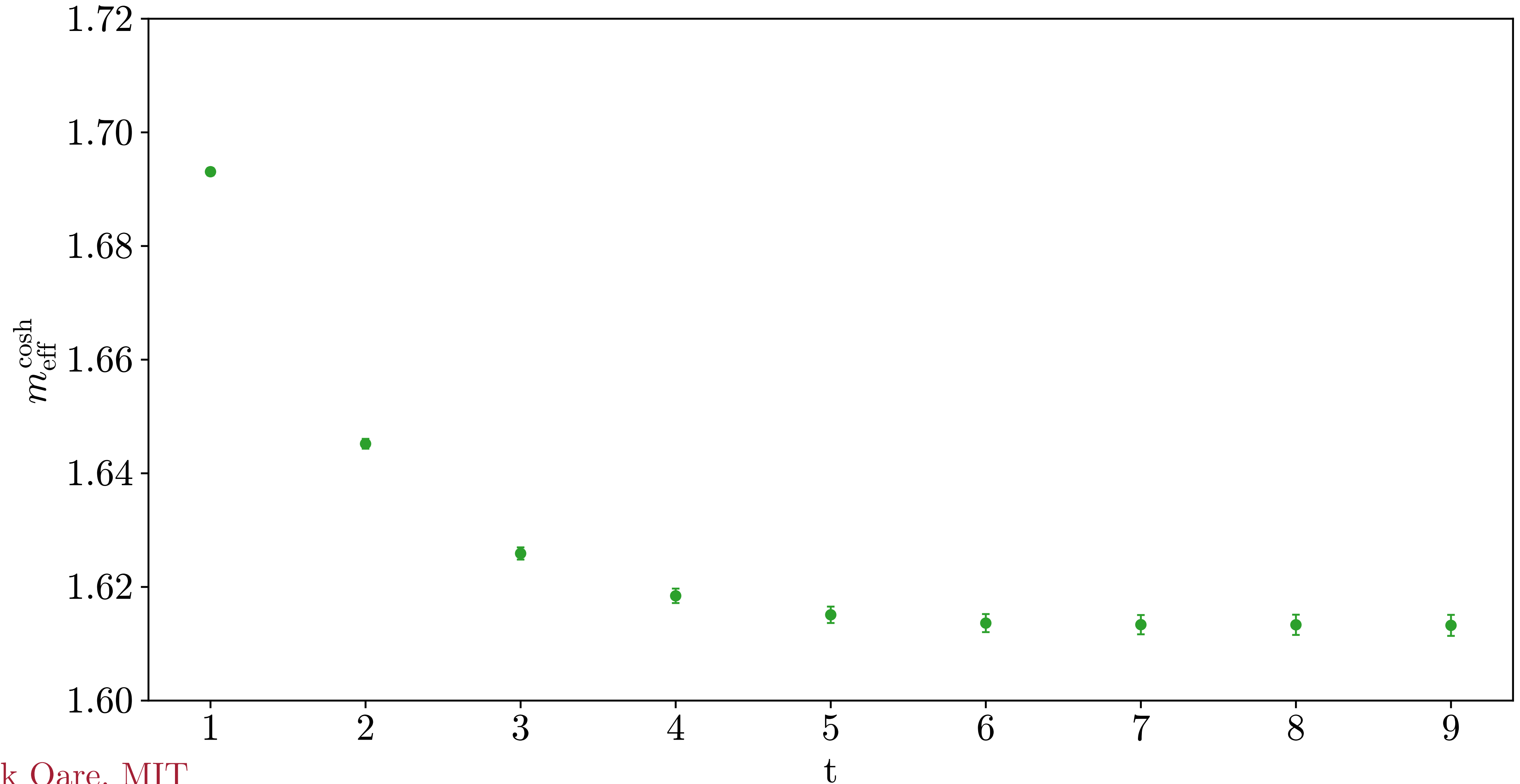




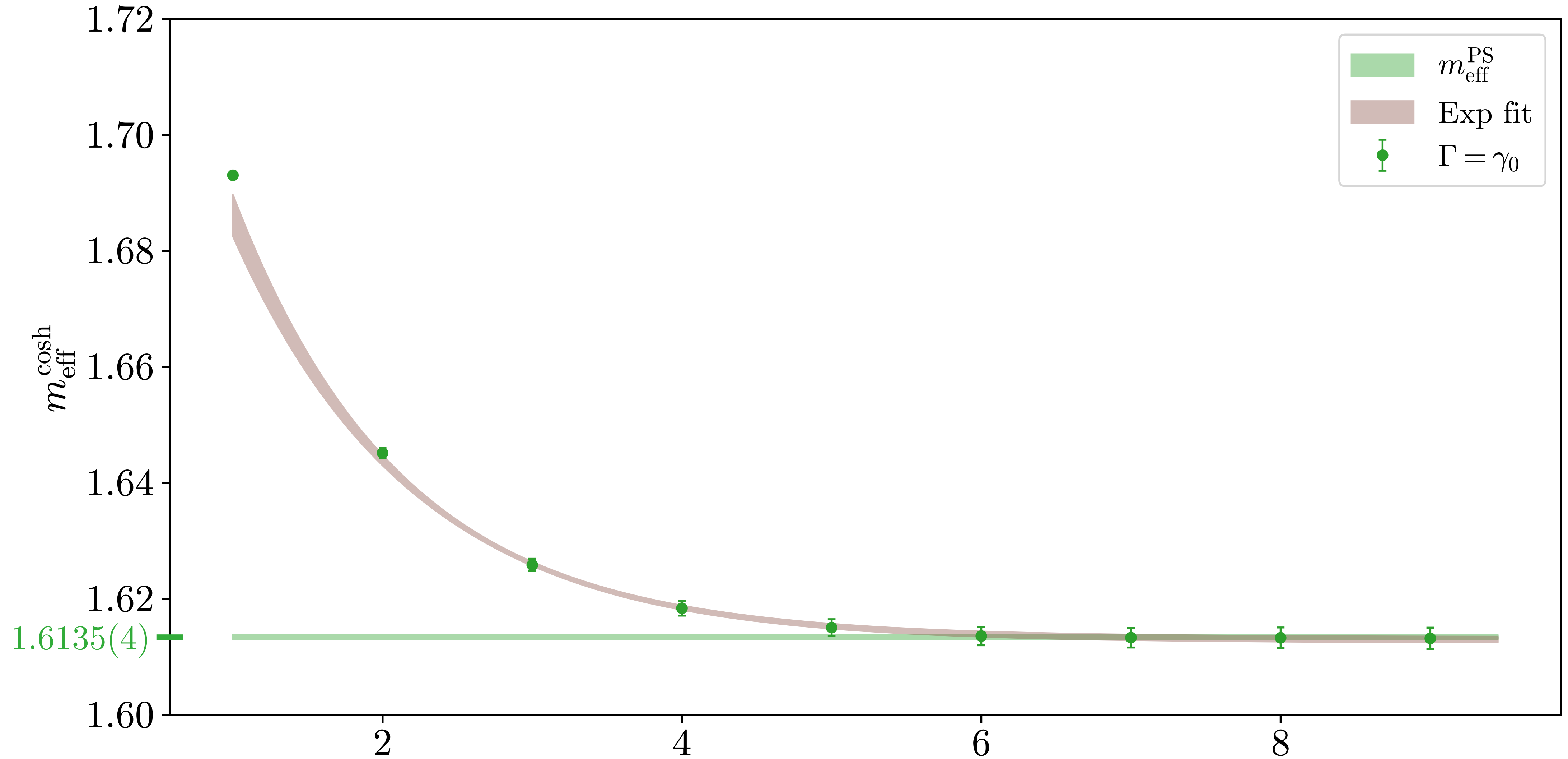
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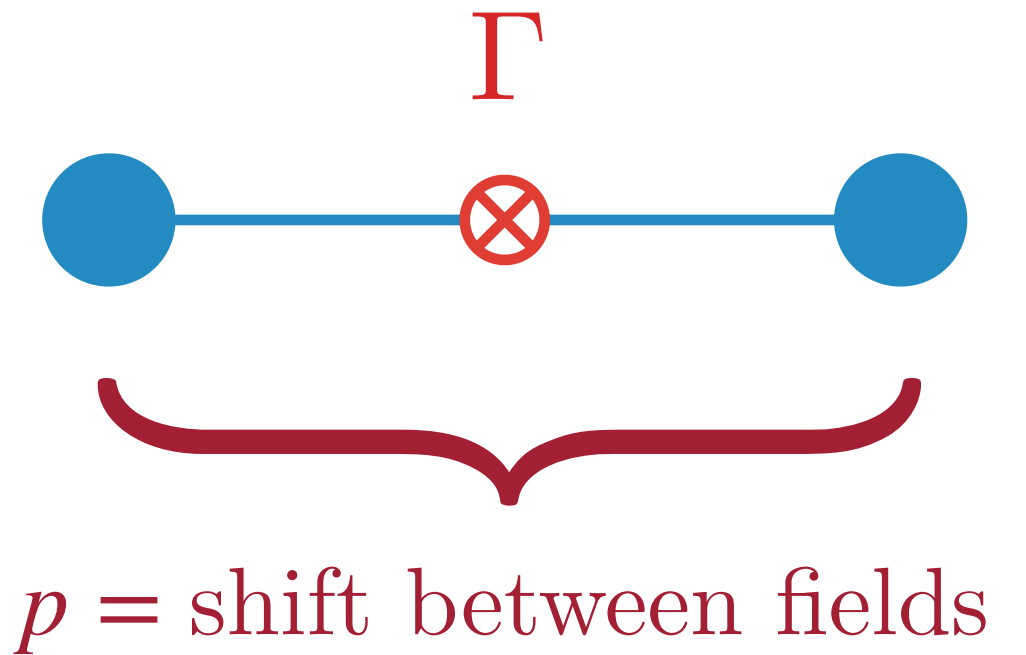


# Low-Lying Spectrum and the GEVP

- Compute the low-lying spectrum with the operator basis,

$$B_{\Gamma}^p(t) = \sum_x \bar{\psi}(x) W_{\text{Adj}}(x, x + p\hat{0}) \Gamma \psi(x + p\hat{0})$$

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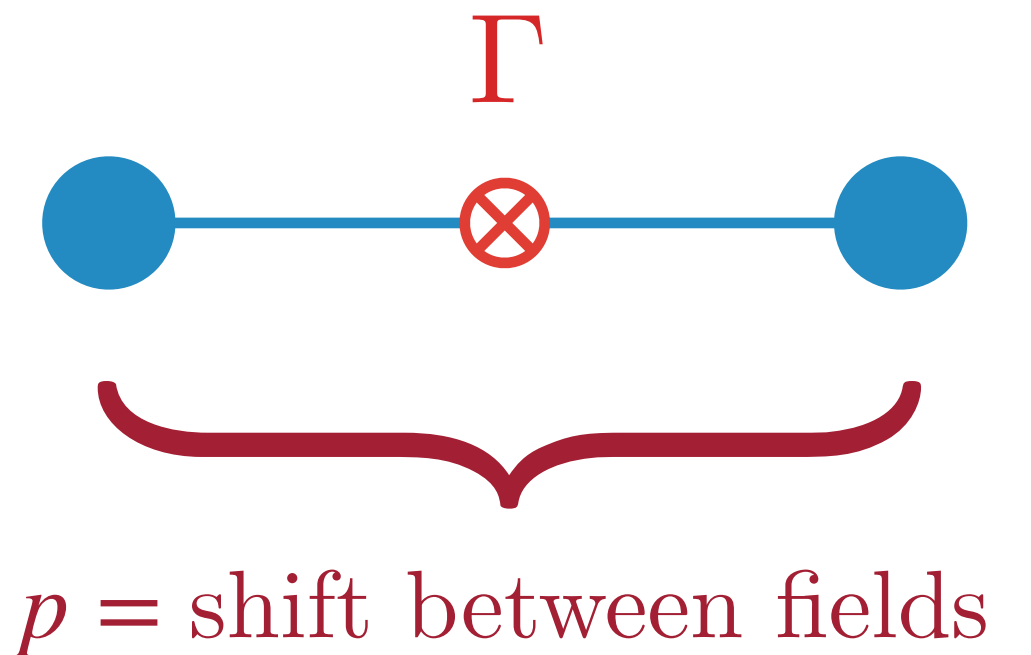
- Construct the correlator matrix,

$$C_{pq}^{\Gamma}(t) = \frac{1}{T} \sum_s \langle \mathcal{B}_p^{\Gamma}(t+s) (\mathcal{B}_q^{\Lambda, \Gamma})^{\dagger}(s) \rangle$$

and solve the Generalized Eigenvalue Problem (GEVP) to evaluate variational bounds  $E_n^{(\text{eff})}$  on the low-lying spectrum,

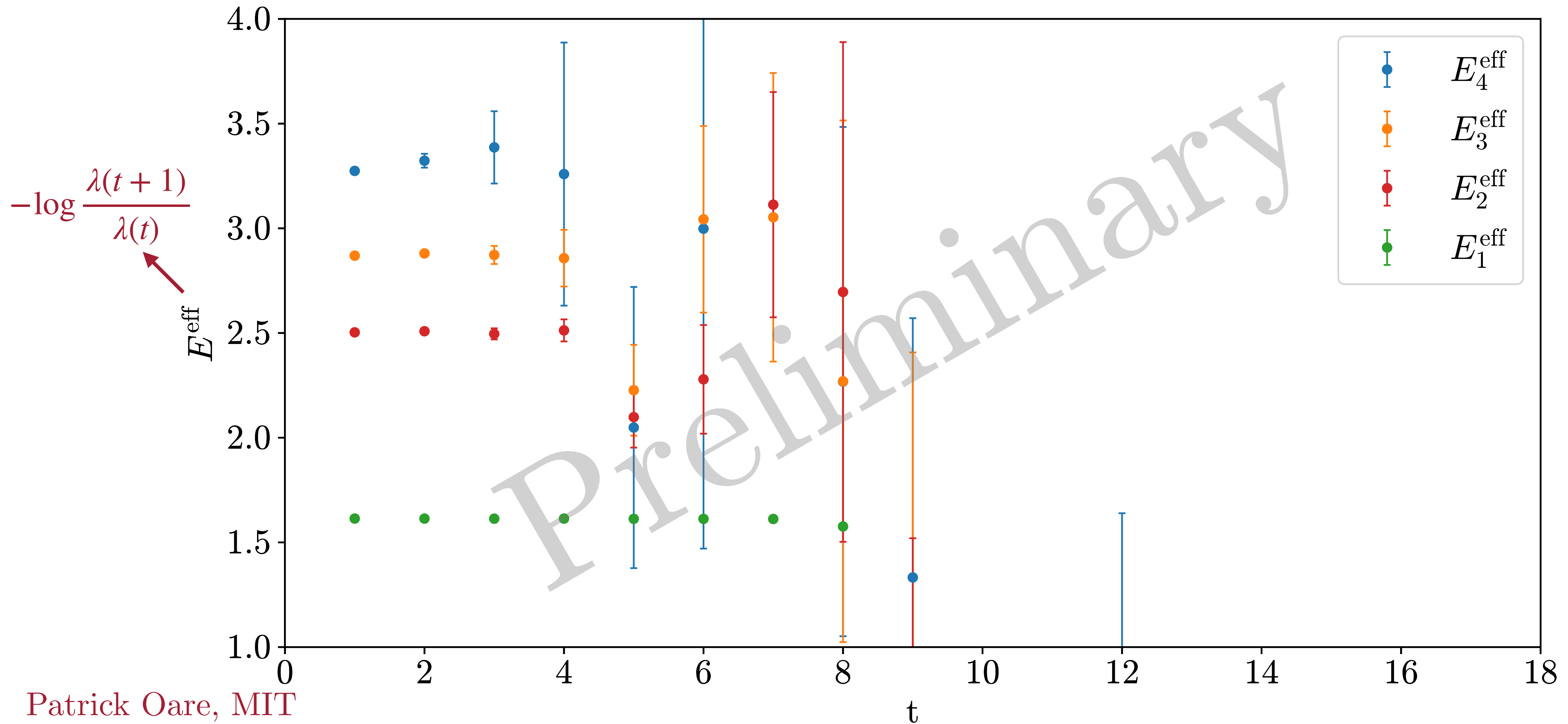
$$C(t) \vec{v}^{(k)}(t, t_0) = \lambda^{(k)}(t, t_0) C(t_0) \vec{v}^{(k)}(t, t_0)$$

$\swarrow \lambda^{(k)}(t, t_0) \sim e^{-E_k(t-t_0)}$



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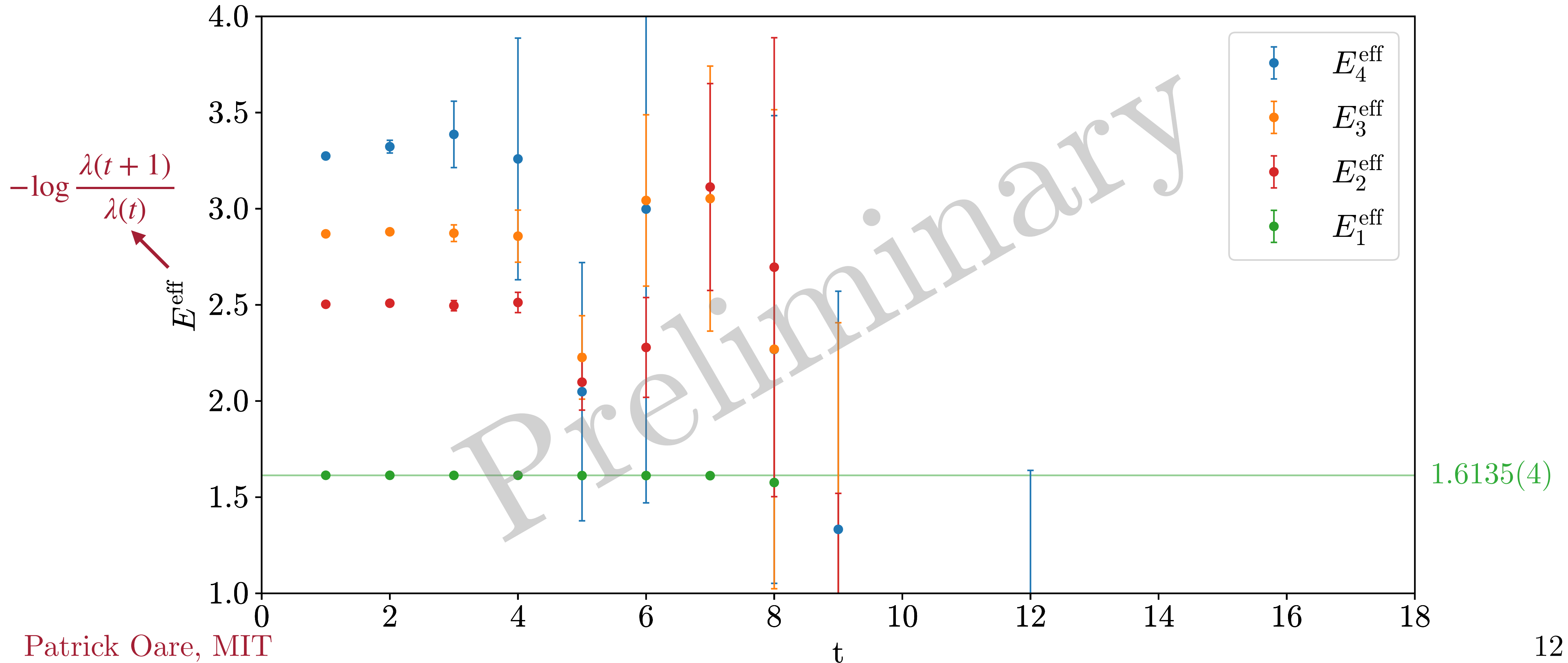
$\Gamma = \gamma_5$ ,  $t_0 = 0$ , maximum shift 3.





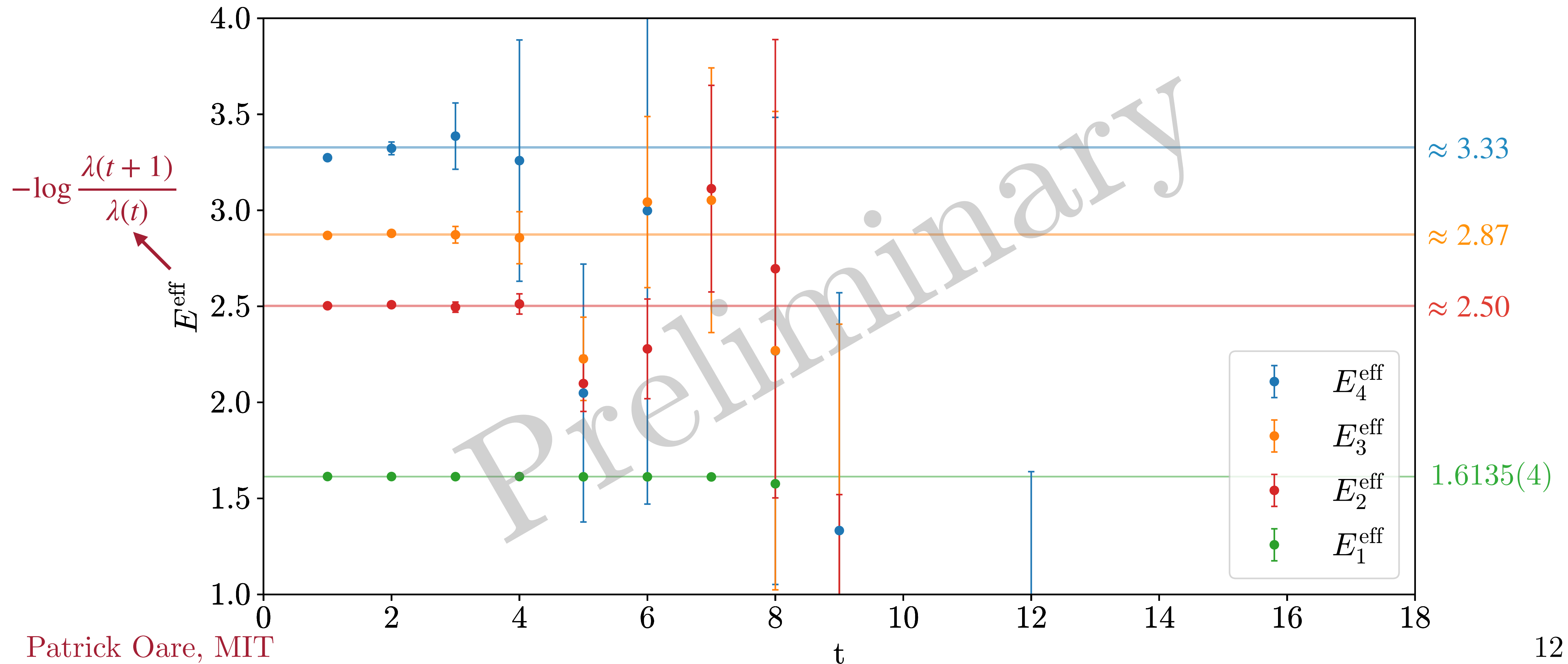
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# Four-Fermion Operators

- Two additional operators may be added to the  $\text{QCD}_2$  action consistent with its symmetries ( $\mathcal{O}_1 = \mathcal{O}_2$  for  $N = 2$  colors):

$$\mathcal{O}_1 = (\text{Tr } \bar{\psi}\psi)^2 \qquad \mathcal{O}_2 = \text{Tr } (\bar{\psi}\gamma^\mu\psi\bar{\psi}\gamma_\mu\psi)$$

A. Cherman *et al.*,  
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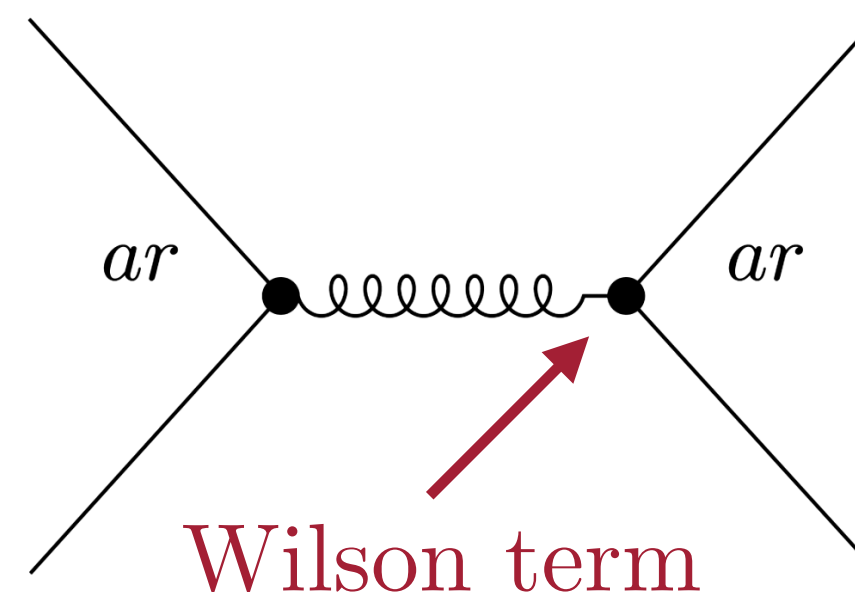
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- Certain discretizations of  $\text{QCD}_2$  action may radiatively generate these operators.
  - ▶ Non-zero  $\langle \mathcal{O}_1 \rangle$  indicates that a coupling for  $\mathcal{O}_1$  is generated at finite  $a$ .

A. Cherman, M. Neuzil.  
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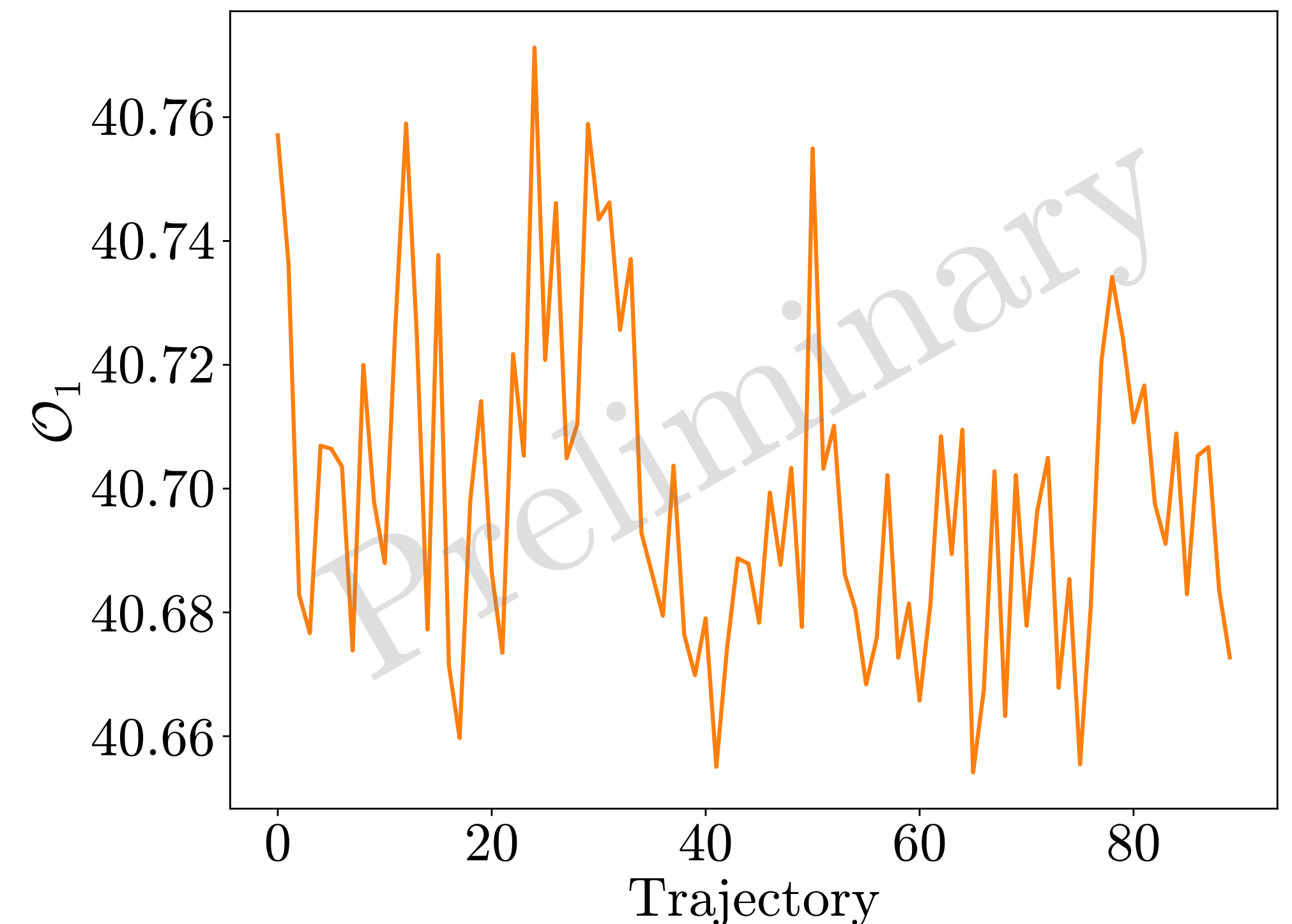
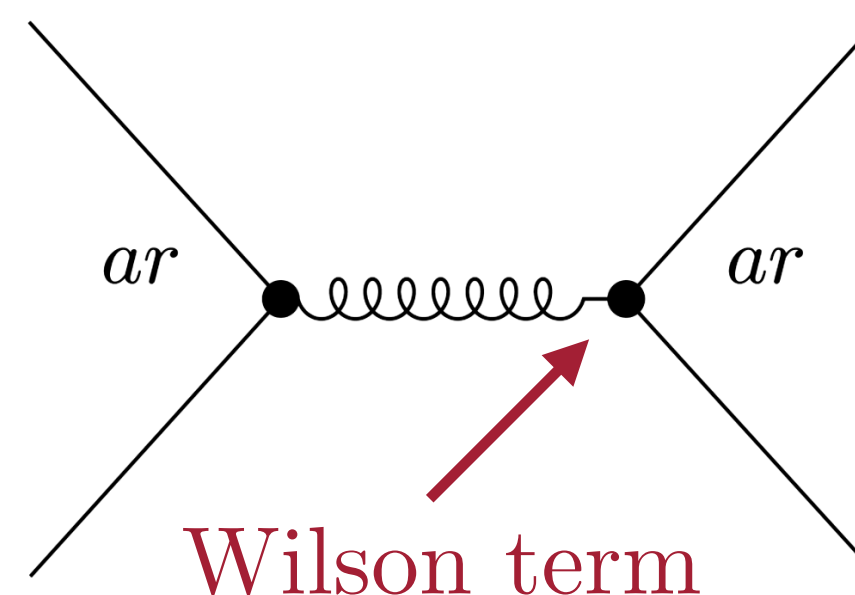
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# Conclusion

- We have presented an ongoing computation of the low-lying spectrum of  $\text{QCD}_2$  using Lattice Gauge Theory.
  - Extrapolation to the continuum limit is still required.
- We are currently in the process of scaling up the calculation.
  - Generating additional ensembles with larger numbers of lattice sites and numbers of colors.
  - Generating additional configurations on each ensemble used in this work.
- Further investigation of the four-fermion operators  $\mathcal{O}_1$  and  $\mathcal{O}_2$  are required: how does  $\langle \mathcal{O}_i \rangle$  scale as  $a \rightarrow 0$ ?

# Backup Slides

# Finite-Volume Dispersion

