

First-order phase transition in dynamical QCD at imaginary isospin

Gergely Endrődi, Guy D. Moore, Alessandro Sciarra



UNIVERSITÄT
BIELEFELD



TECHNISCHE
UNIVERSITÄT
DARMSTADT



GOETHE
UNIVERSITÄT
FRANKFURT AM MAIN

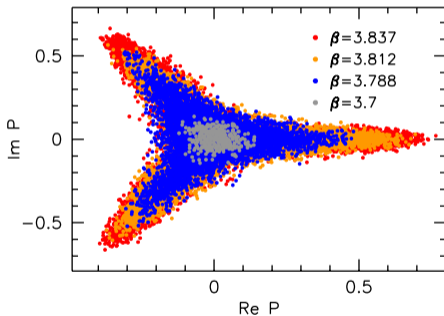


CRC-TR 211
Strong-interaction matter
under extreme conditions

Lattice'24 @ Liverpool, July 30, 2024

Appetizer

- ▶ Polyakov loop histogram in a QCD simulation with dynamical fermions



Outline

- ▶ theory: center-symmetric three-flavor QCD at imaginary isospin
- ▶ finite size scaling analysis of observables
- ▶ mass dependence and phase diagram
- ▶ summary

Center symmetry

Center symmetry

- ▶ pure gauge theory

$$Z = \int \mathcal{D}U e^{-\beta S_g}$$

- ▶ center transformation

$$\forall \mathbf{n} : \quad U_4(\mathbf{n}, n_4) \rightarrow z U_4(\mathbf{n}, n_4) \quad z \in Z(3) = \{1, \mathbb{1}e^{i2\pi/3}, \mathbb{1}e^{-i2\pi/3}\} \subset \text{SU}(3)$$

- ▶ invariance

$$S_g \rightarrow S_g$$

- ▶ with Polyakov loop as order parameter

$$P = \frac{1}{V} \sum_{\mathbf{n}} \prod_{n_4=1}^{N_t} U_4(\mathbf{n}, n_4) \rightarrow z P$$

Center symmetry

- ▶ pure gauge theory

$$Z = \int \mathcal{D}U e^{-\beta S_g}$$

- ▶ center transformation

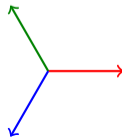
$$\forall \mathbf{n} : \quad U_4(\mathbf{n}, n_4) \rightarrow z U_4(\mathbf{n}, n_4) \quad z \in Z(3) = \{1, \mathbb{1}e^{i2\pi/3}, \mathbb{1}e^{-i2\pi/3}\} \subset \text{SU}(3)$$

- ▶ invariance

$$S_g \rightarrow S_g$$

- ▶ with Polyakov loop as order parameter, revealing SSB at high T

$$P = \frac{1}{V} \sum_{\mathbf{n}} \prod_{n_4=1}^{N_t} U_4(\mathbf{n}, n_4) \rightarrow z P$$



Center symmetry

- ▶ dynamical QCD

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(0) + m]$$

Center symmetry

- ▶ dynamical QCD

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(0) + m]$$

- ▶ fermion determinant breaks center symmetry and prefers: real sector at zero chemical potential

$$\log \det(\not{D}(0) + m) \stackrel{\text{large } m \text{ expansion}}{\approx} P\text{-independent} + \frac{1}{2m^{N_t}} \text{Re} \left(P \right) + \mathcal{O}(m^{-2N_t})$$



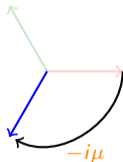
Center symmetry

- ▶ dynamical QCD

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(\mu) + m]$$

- ▶ fermion determinant breaks center symmetry and prefers:
real sector at zero chemical potential
complex center sector at imaginary chemical potential $i\mu$

$$\log \det(\not{D}(\mu) + m) \stackrel{\text{large } m \text{ expansion}}{\approx} P\text{-independent} + \frac{1}{2m^{N_t}} \operatorname{Re} \left(P e^{i\mu/T} \right) + \mathcal{O}(m^{-2N_t})$$



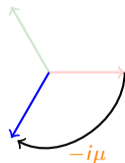
Center symmetry

- ▶ dynamical QCD

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(\mu) + m]$$

- ▶ fermion determinant breaks center symmetry and prefers:
real sector at zero chemical potential
complex center sector at imaginary chemical potential $i\mu$

$$\log \det(\not{D}(\mu) + m) \stackrel{\text{large } m \text{ expansion}}{\approx} P\text{-independent} + \frac{1}{2m^{N_t}} \operatorname{Re} \left(P e^{i\mu/T} \right) + \mathcal{O}(m^{-2N_t})$$



- ▶ including $\mu = i2\pi T/3$ equivalent to center trafo $z = \mathbb{1} e^{-i2\pi/3}$

Center-symmetric QCD

Center-symmetric QCD

- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

Center-symmetric QCD

- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

- ▶ now center transformation merely permutes u, d, s flavors

Center-symmetric QCD

- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

- ▶ now center transformation merely permutes u, d, s flavors
- ▶ exact center symmetry, just like pure gauge theory

Center-symmetric QCD

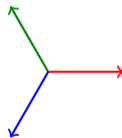
- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

- ▶ now center transformation merely permutes u, d, s flavors
- ▶ exact center symmetry, just like pure gauge theory



Center-symmetric QCD

- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

- ▶ now center transformation merely permutes u, d, s flavors

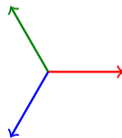
- ▶ exact center symmetry, just like pure gauge theory

- ▶ first observed in PNJL model [Kouno et al. '12](#)

general discussion on symmetries [Cherman et al. '17](#)

hints for first-order transition on the lattice [Iritani et al. '15](#)

scenarios for deconfinement, chiral symm. restoration [Yonekura '19](#) [Tanizaki et al. '17](#)



Center-symmetric QCD

- ▶ three degenerate quarks with imaginary isospin

$$\mathcal{Z} = \int \mathcal{D}U e^{-\beta S_g} \det[\not{D}(i2\pi T/3) + m] \det[\not{D}(-i2\pi T/3) + m] \det[\not{D}(0) + m]$$

- ▶ amounts to imaginary isospin chemical potential $i\mu_I = 4\pi T/3$

$$\mu_u = \mu_I/2, \quad \mu_d = -\mu_I/2, \quad \mu_s = 0$$

- ▶ now center transformation merely permutes u, d, s flavors

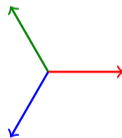
- ▶ exact center symmetry, just like pure gauge theory

- ▶ first observed in PNJL model [Kouno et al. '12](#)

general discussion on symmetries [Cherman et al. '17](#)

hints for first-order transition on the lattice [Iritani et al. '15](#)

scenarios for deconfinement, chiral symm. restoration [Yonekura '19](#) [Tanizaki et al. '17](#)



- ▶ goal: demonstrate first-order deconfinement and chiral transition

Results

Lattice setup and rotated Polyakov loop

- ▶ rooted stout-smearred staggered quarks

$$m_{u,d,s} = m_{ud}^{\text{phys}}, \quad m_s^{\text{phys}}, \quad 3 \cdot m_s^{\text{phys}}$$

$$N_s^3 \times N_t = 16^3 \times 6 \dots 32^3 \times 6$$

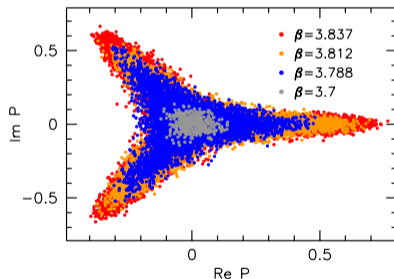
Lattice setup and rotated Polyakov loop

- ▶ rooted stout-smearred staggered quarks

$$m_{u,d,s} = m_{ud}^{\text{phys}}, \quad m_s^{\text{phys}}, \quad 3 \cdot m_s^{\text{phys}}$$

$$N_s^3 \times N_t = 16^3 \times 6 \dots 32^3 \times 6$$

- ▶ histogram of P in the complex plane



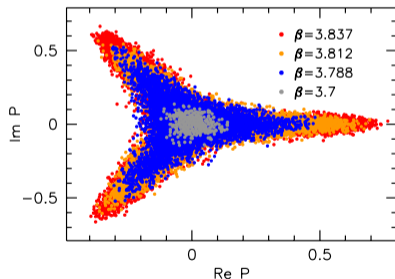
Lattice setup and rotated Polyakov loop

- ▶ rooted stout-smearred staggered quarks

$$m_{u,d,s} = m_{ud}^{\text{phys}}, \quad m_s^{\text{phys}}, \quad 3 \cdot m_s^{\text{phys}}$$

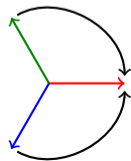
$$N_s^3 \times N_t = 16^3 \times 6 \dots 32^3 \times 6$$

- ▶ histogram of P in the complex plane



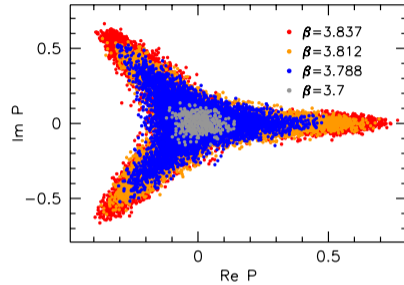
- ▶ rotated Polyakov loop $z_U \in Z(3)$

$$\bar{P} = \text{Re}(z_U P), \quad -\pi/3 \leq \arg(z_U P) \leq \pi/3$$



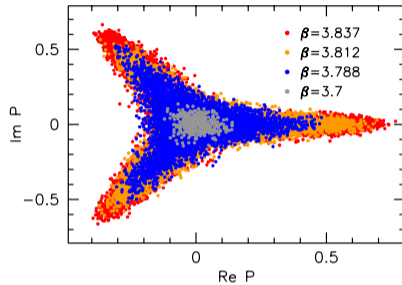
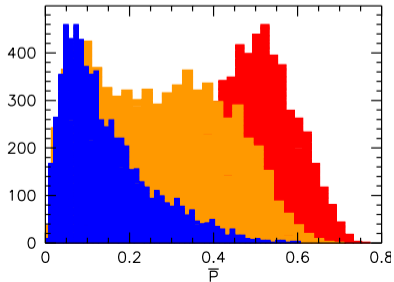
Histograms and MC history

► for $m_{u,d,s} = m_s^{\text{phys}}$, $16^3 \times 6$



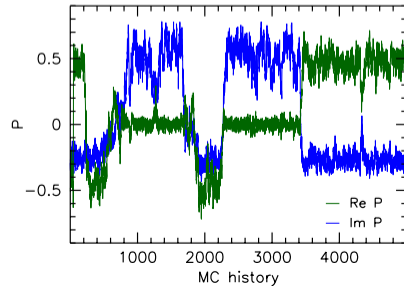
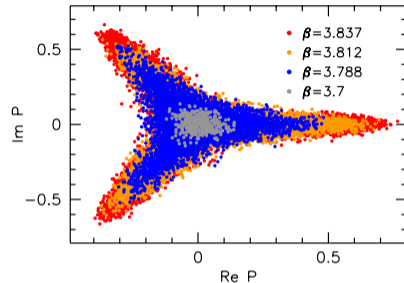
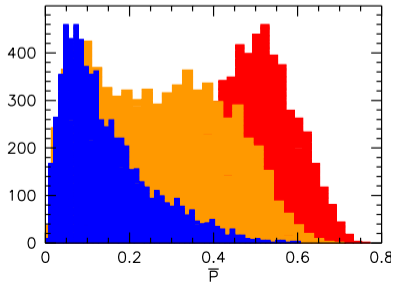
Histograms and MC history

- ▶ for $m_{u,d,s} = m_s^{\text{phys}}$, $16^3 \times 6$
- ▶ histogram of \bar{P}



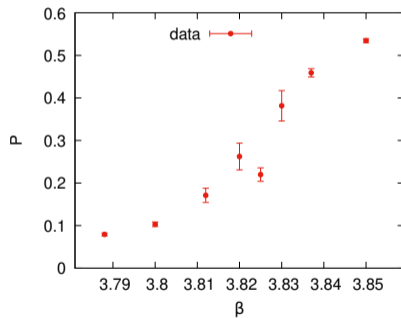
Histograms and MC history

- ▶ for $m_{u,d,s} = m_s^{\text{phys}}$, $16^3 \times 6$
- ▶ histogram of \bar{P}
- ▶ Monte-Carlo history of P around β_c



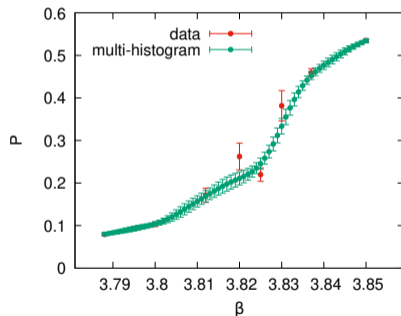
Multi-histogram method

- ▶ raw data for \bar{P} on $16^3 \times 6$, $m_{u,d,s} = m_s^{\text{phys}}$



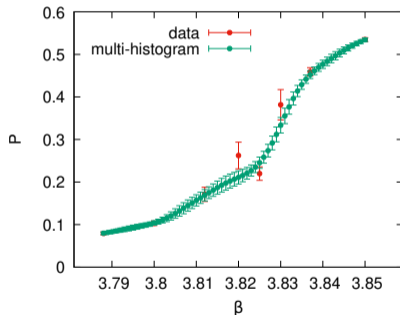
Multi-histogram method

- ▶ raw data for \bar{P} on $16^3 \times 6$, $m_{u,d,s} = m_s^{\text{phys}}$
- ▶ interpolation via multi-histogram method [Newman, Barkema](#)



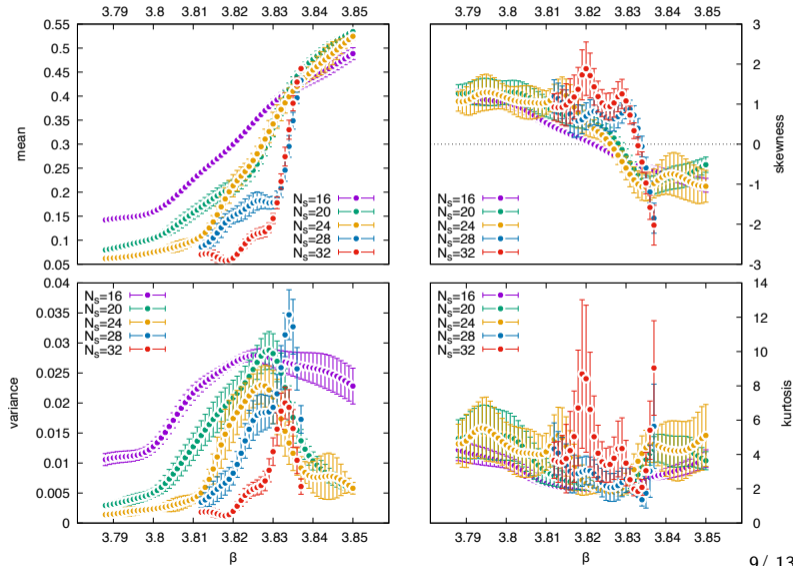
Multi-histogram method

- ▶ raw data for \bar{P} on $16^3 \times 6$, $m_{u,d,s} = m_s^{\text{phys}}$
- ▶ interpolation via multi-histogram method [Newman, Barkema](#)



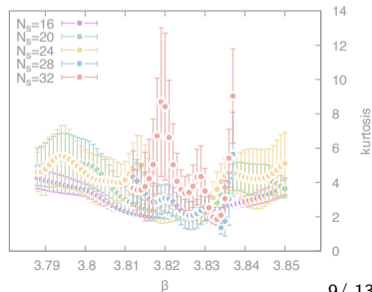
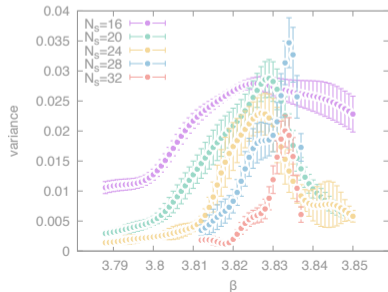
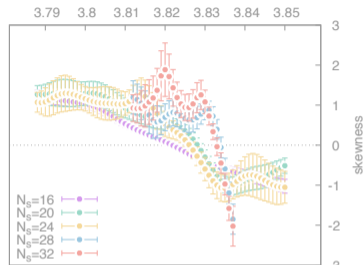
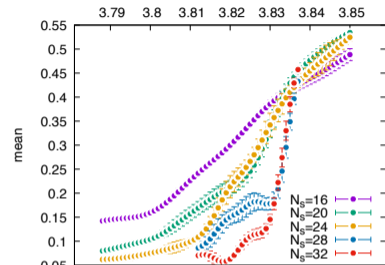
- ▶ allows interpolation of higher moments

Finite size scaling



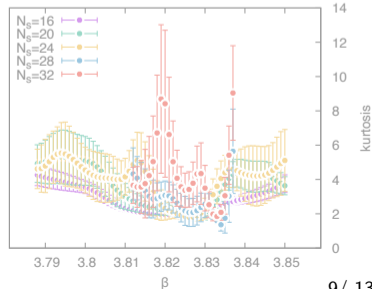
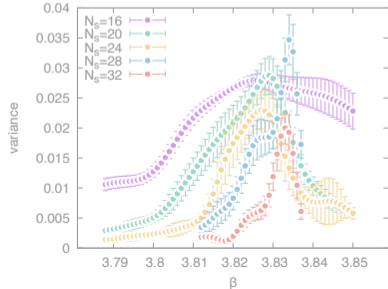
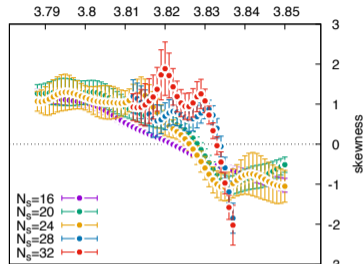
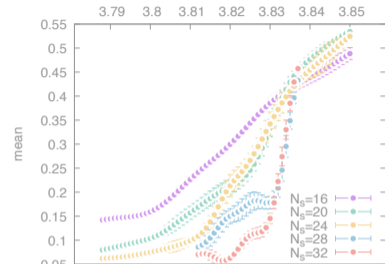
Finite size scaling

- ▶ \bar{P} gets steeper as $V \rightarrow \infty$



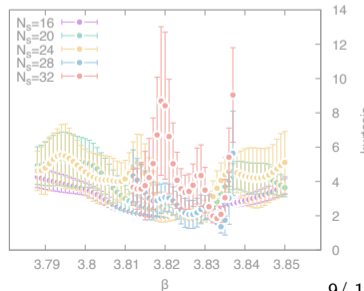
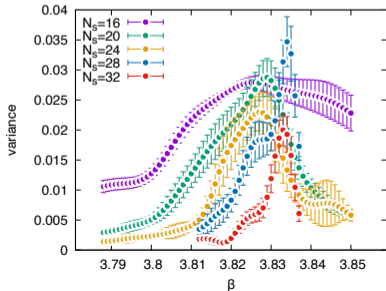
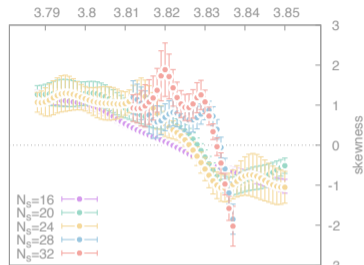
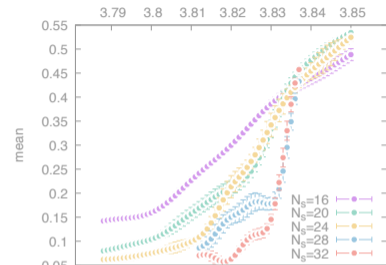
Finite size scaling

- ▶ \bar{P} gets steeper as $V \rightarrow \infty$
- ▶ skewness = 0 to define T_c



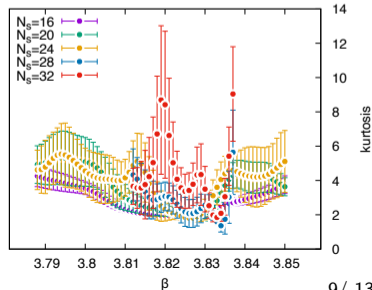
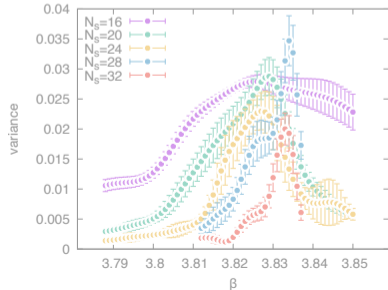
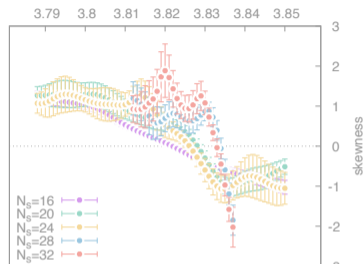
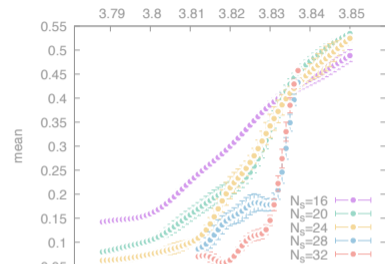
Finite size scaling

- ▶ \bar{P} gets steeper as $V \rightarrow \infty$
- ▶ skewness = 0 to define T_c
- ▶ $\chi_{\bar{P}} \propto V$ at $T = T_c$



Finite size scaling

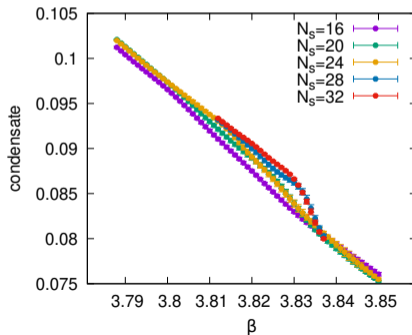
- ▶ \bar{P} gets steeper as $V \rightarrow \infty$
- ▶ skewness = 0 to define T_c
- ▶ $\chi_{\bar{P}} \propto V$ at $T = T_c$
- ▶ Binder ~ 1 at $T = T_c$



Chiral symmetry restoration

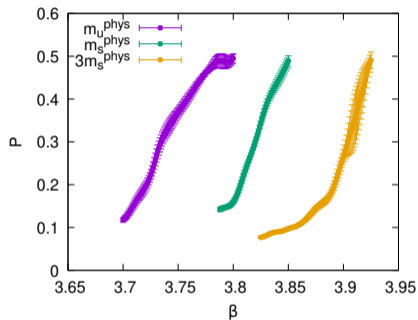
- ▶ how is chiral symmetry restoration affected by first-order deconfining phase transition?
- ▶ center-symmetric quark condensate [Cherman et al. '17](#)

$$\bar{\psi}\psi = \bar{\psi}_u\psi_u + \bar{\psi}_d\psi_d + \bar{\psi}_s\psi_s$$



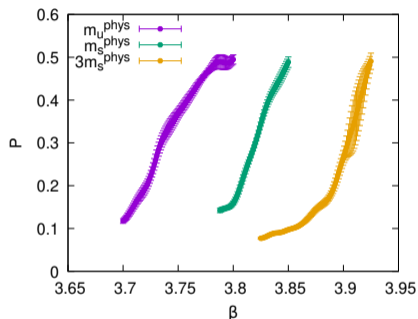
Impact of quark mass and scale

- ▶ first-order nature of transition not expected to be affected by $m_{u,d,s}$
- ▶ scan for $16^3 \times 6$



Impact of quark mass and scale

- ▶ first-order nature of transition not expected to be affected by $m_{u,d,s}$
- ▶ scan for $16^3 \times 6$

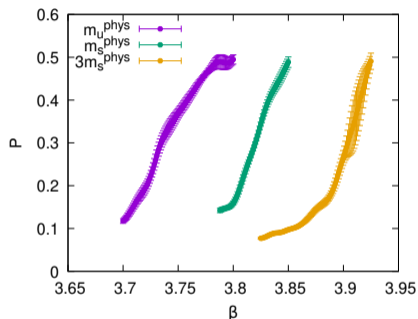


- ▶ calculate scale at $T = \mu = 0$ via w_0

$m_{u,d,s}$	m_{ud}^{phys}	m_s^{phys}	$3 \cdot m_s^{\text{phys}}$
$T_c w_0$	0.239(1)	0.248(1)	0.2505(5)

Impact of quark mass and scale

- ▶ first-order nature of transition not expected to be affected by $m_{u,d,s}$
- ▶ scan for $16^3 \times 6$



- ▶ calculate scale at $T = \mu = 0$ via w_0 comparison

$m_{u,d,s}$	m_{ud}^{phys}	m_s^{phys}	$3 \cdot m_s^{\text{phys}}$	pure gauge	$m_{ud}^{\text{phys}}, m_s^{\text{phys}}$
$T_c w_0$	0.239(1)	0.248(1)	0.2505(5)	0.2507(2) <small>⌘ Borsányi et al. '22</small>	≈ 0.14 <small>⌘ Borsányi et al. '12</small>

Phase diagram

