



Anisotropic excited bottomonia from a basis of smeared operators

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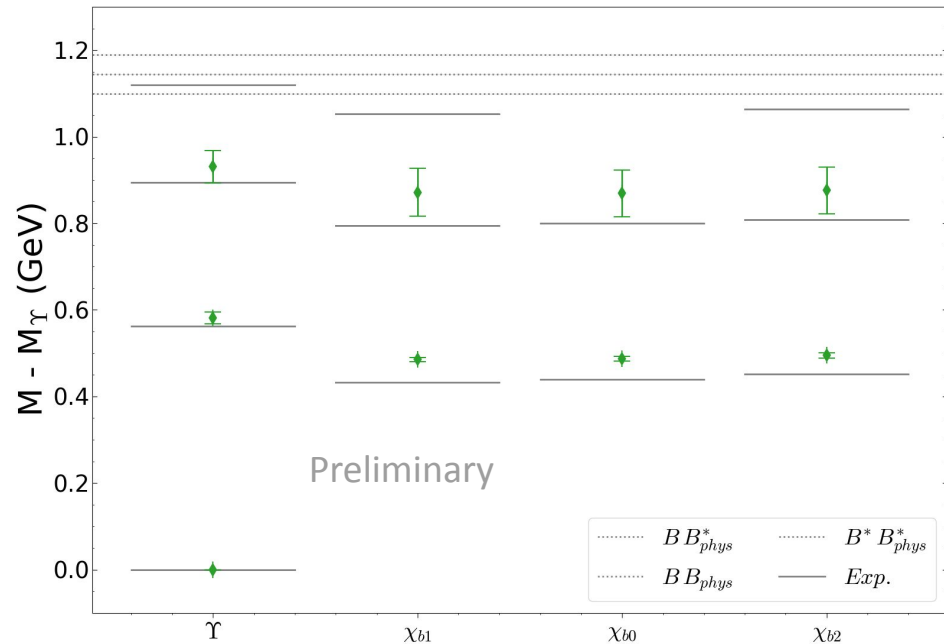
QCD at non-zero temperature

Lattice 2024, the University of Liverpool

31st June 2024

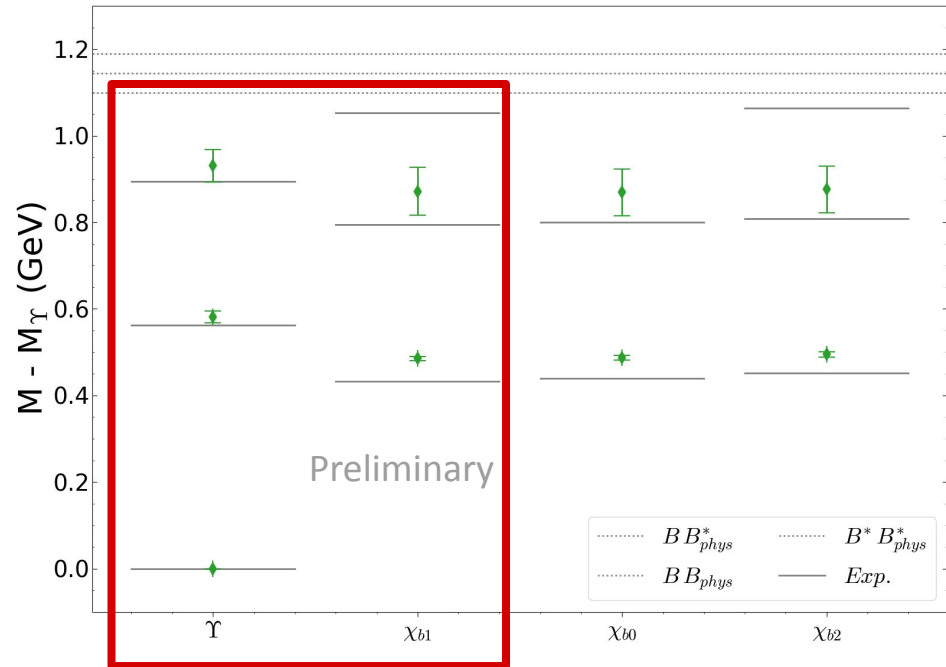
Bottomonium spectrum

- Experimentally accessible
- Phenomenologically relevant to
 - heavy-ion experiments
 - Quark gluon plasma
 - Early universe
- Easily computable
 - via lattice NRQCD
 - statistically well-behaved due to scale separation



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Ensemble Details

Generation 2L FASTSUM

N_T	128	64	56	48	40	36	32	28	24	20	16
Temperature (MeV)	47	95	109	127	152	169	190	217	253	304	380
# Wall Sources	16	16	16	20	24	24	32	28	24	20	16

Action details:

- **Gauge: Symanzik-improved, tree-level tadpole**
- **Fermion: Wilson-clover, tree-level tadpole, stout-links**
- **Same parameters as HadSpec Collaboration**
- **Approx. 1000 configurations at each temperature**
- **NRQCD action for bottom quarks**
 - Incorporating $O(v^4)$ corrections
 - Tree-level matching coefficients

$$m_\pi \sim 236 \text{ MeV}, \xi \sim 3.5, T_c \sim 167 \text{ MeV}$$

Excited State spectroscopy

Generalised EigenValue Problem - GEVP

- Build correlation matrix of two point functions

$$G_{ij}(\tau) = \langle \Omega | \mathcal{O}_i \mathcal{O}_j^\dagger | \Omega \rangle = \sum_{\alpha} \frac{Z_i^{\alpha} Z_j^{\alpha \dagger}}{2 E_{\alpha}} e^{-E_{\alpha} \tau}$$

- Solve generalised eigenvalue problems

$$\begin{aligned} G_{ij}(\tau_0 + \delta_{\tau}) u_j^{\alpha} &= e^{-E_{\alpha} \delta_{\tau}} G_{ij}(\tau_0) u_j^{\alpha} \\ v_i^{\alpha} G_{ij}(\tau_0 + \delta_{\tau}) &= e^{-E_{\alpha} \delta_{\tau}} v_i^{\alpha} G_{ij}(\tau_0) \end{aligned}$$

- Construct Projected Correlator

$$G_{\alpha}(\tau) = v_i^{\alpha} G_{ij}(\tau) u_j^{\alpha}$$

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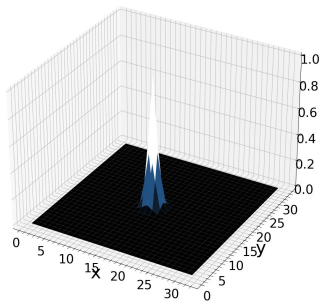
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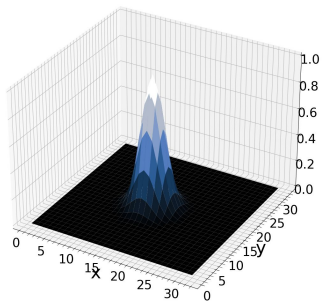
GEVP - Operator Basis

Four widths of Gaussian and `excited` operator

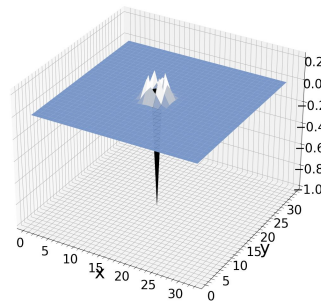
X_1_000



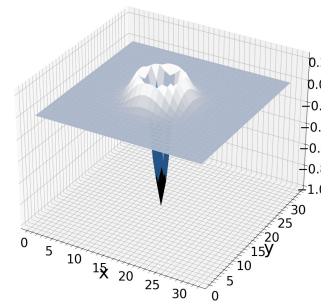
X_2_50



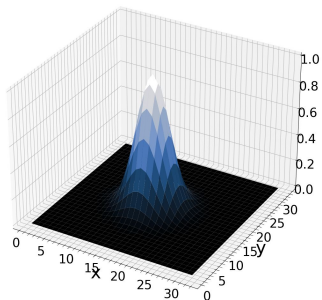
E_1_000



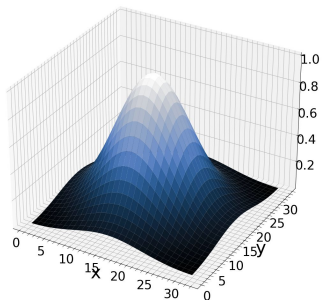
E_2_50



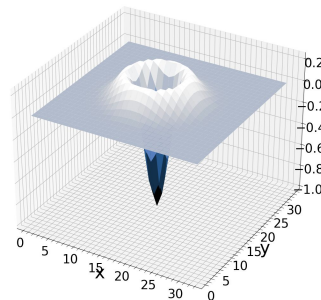
X_3_50



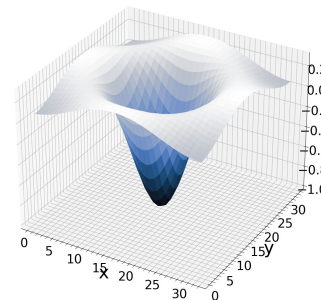
X_8_00



E_3_50

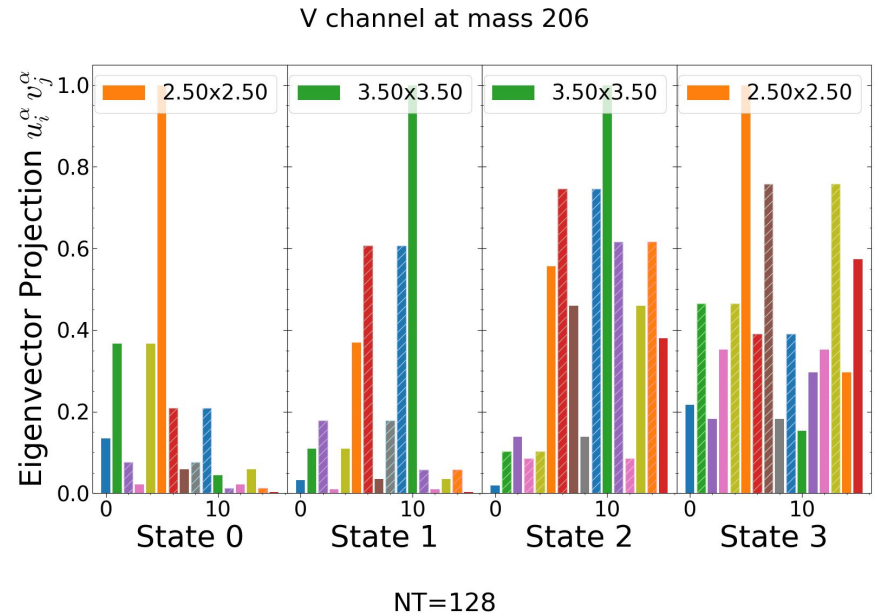


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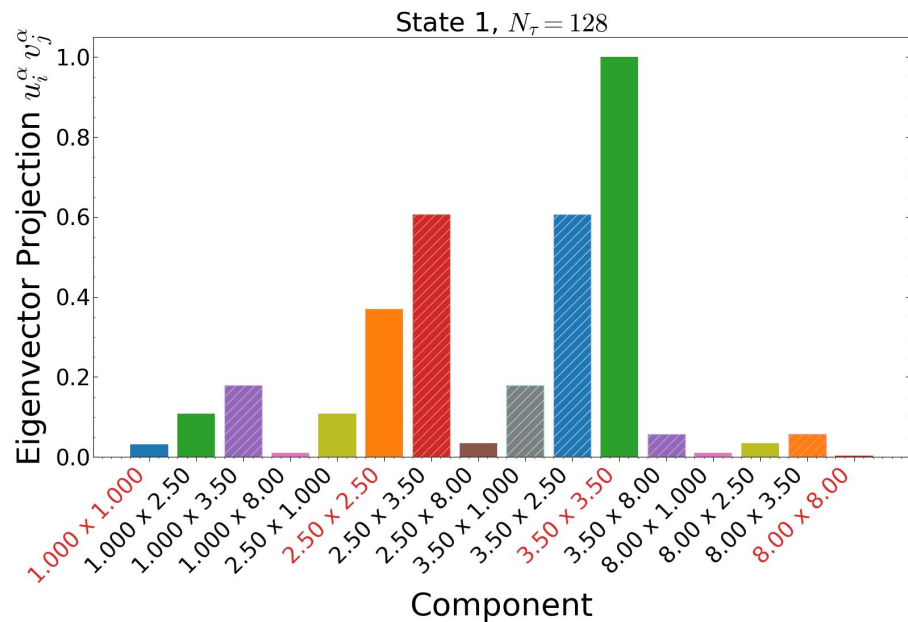
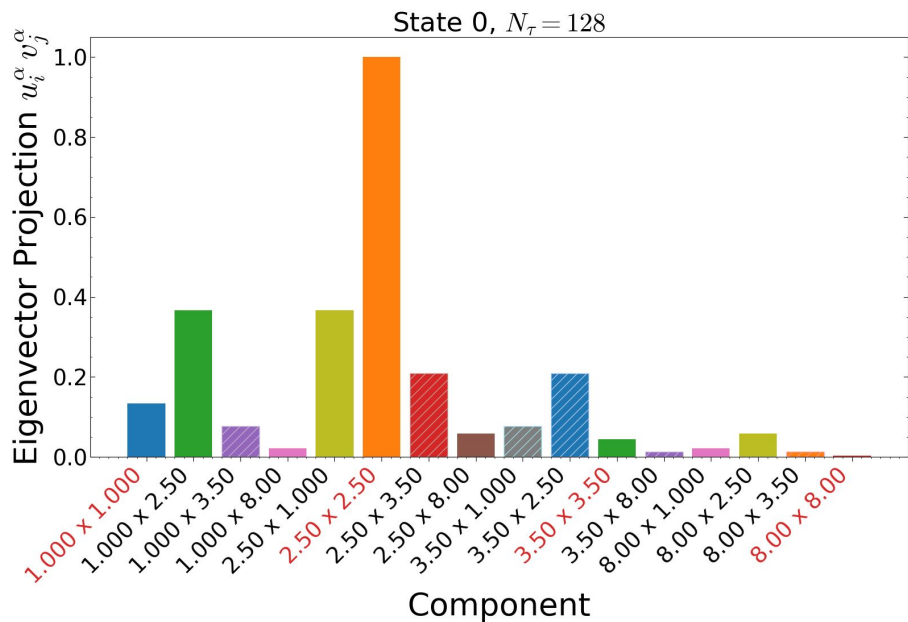


Eigenvectors

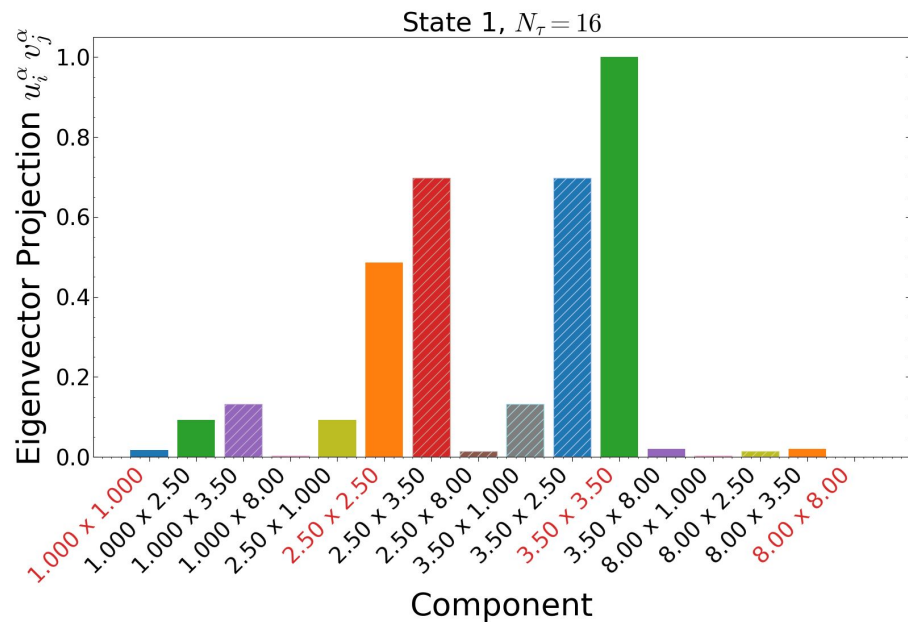
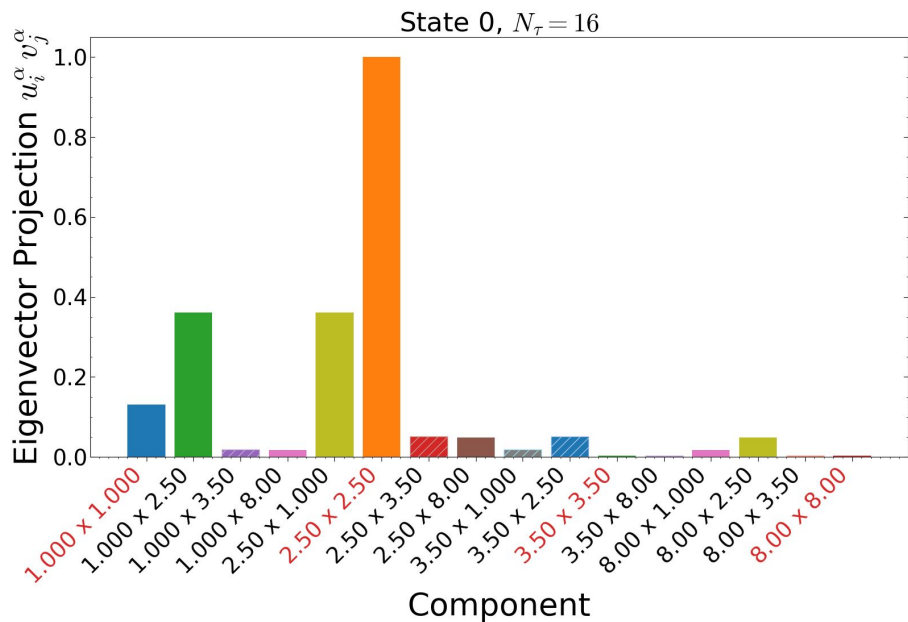
- Related to overlap of each ‘operator’ with each state
- Examine eigenvectors to see how they change as temperature increases
- Plots have the largest contribution is normalised to one, and negative contributions are ‘hashed’



Eigenvectors

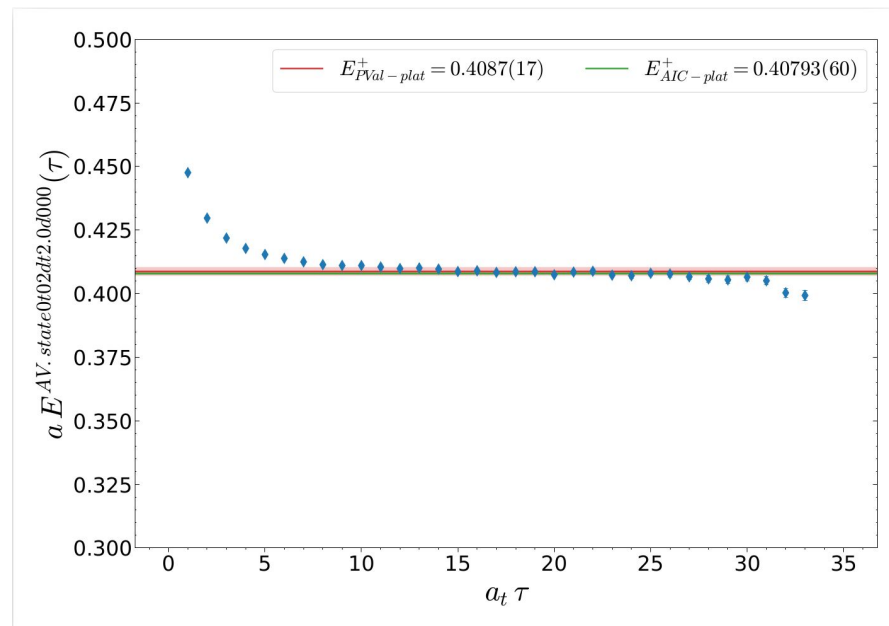
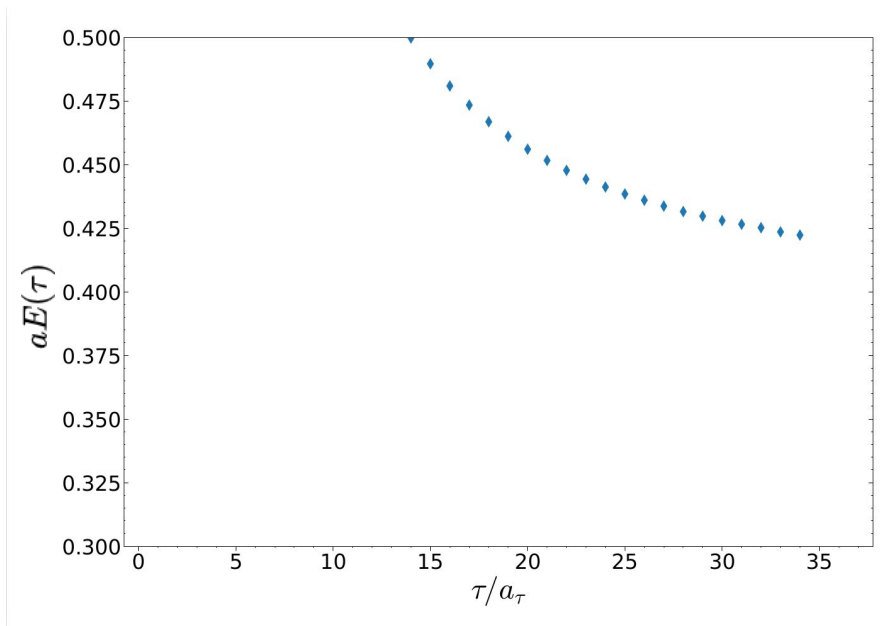


Eigenvectors



Improved state isolation

$N_t = 36$, $\chi_{b1}(1P)$



Spectral Representation

of NRQCD correlator

$$G(\tau) = \int_0^\infty \frac{d\omega}{2\pi} e^{-\omega\tau} \rho(\omega)$$

Model spectral function $\rho(\omega)$ using a delta-function of the ground state.

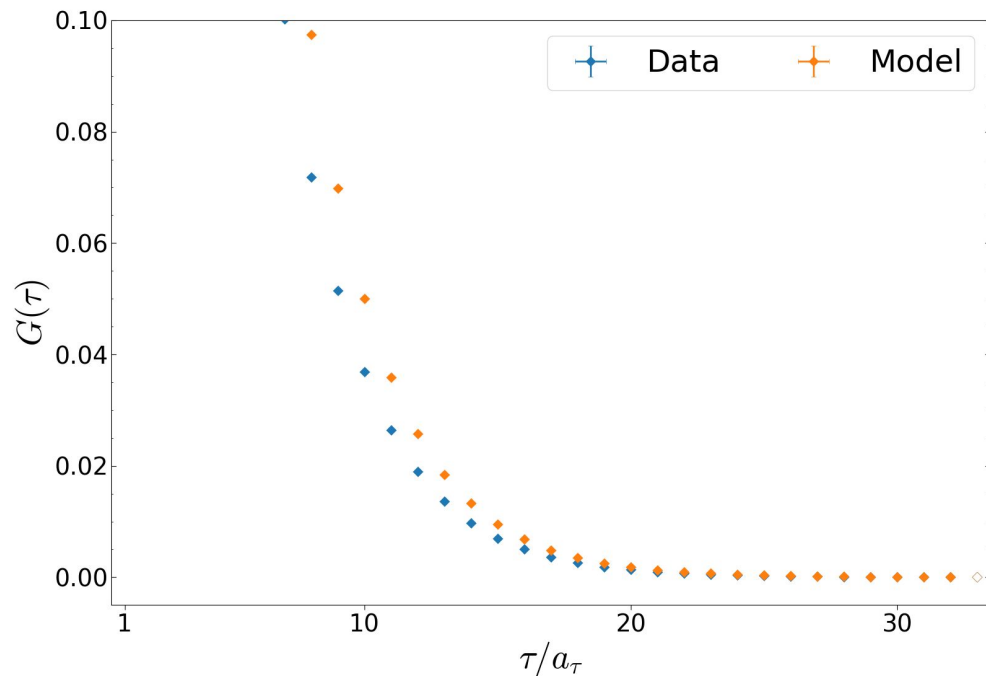
Construct single ratio

$$r(\tau; T, T_0) = \frac{G(\tau; T)}{G_{\text{model}}(\tau; T, T_0)}$$

And hence double ratio

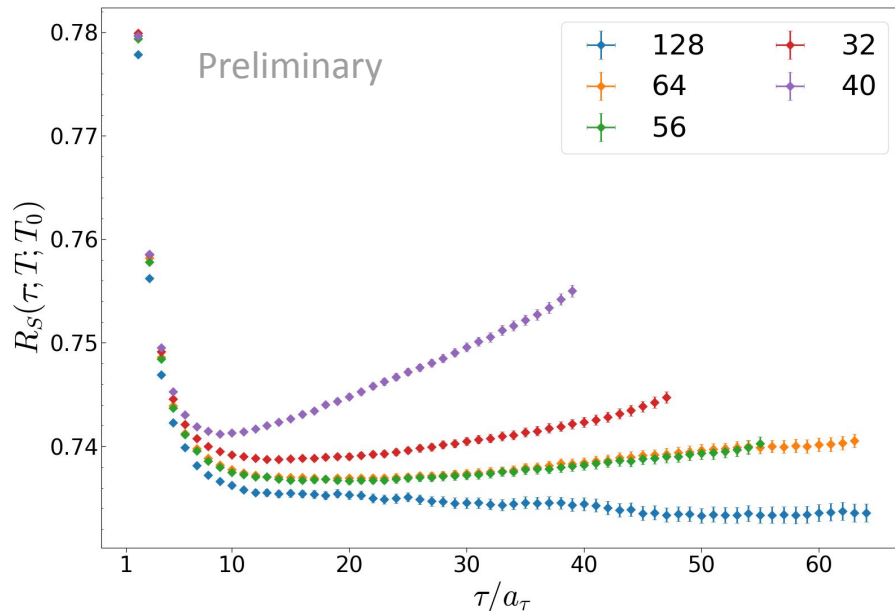
$$R(\tau; T, T_0) = \frac{r(\tau; T, T_0)}{r(\tau; T_0, T_0)}$$

Describes the `change` in spectral function $\rho(\omega)$



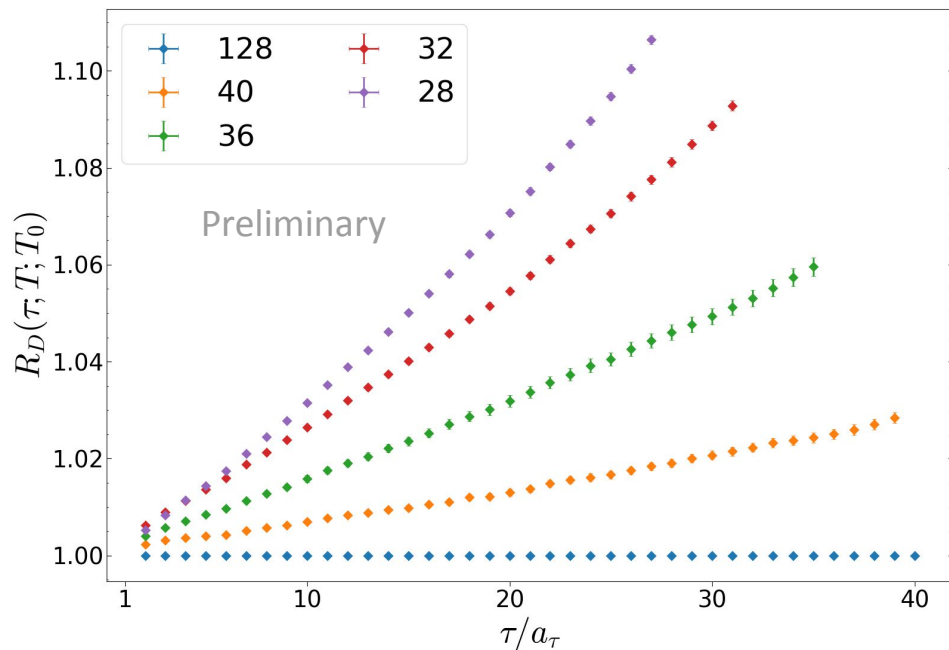
Single & Double Ratio

- Single Ratio shows how similar to zero-temperature
 - Excited states still present
 - Constant if $\rho(\omega)$ is a delta-function
- Double Ratio
 - Removes excited state effect
 - Differences from **one** show difference in correlator

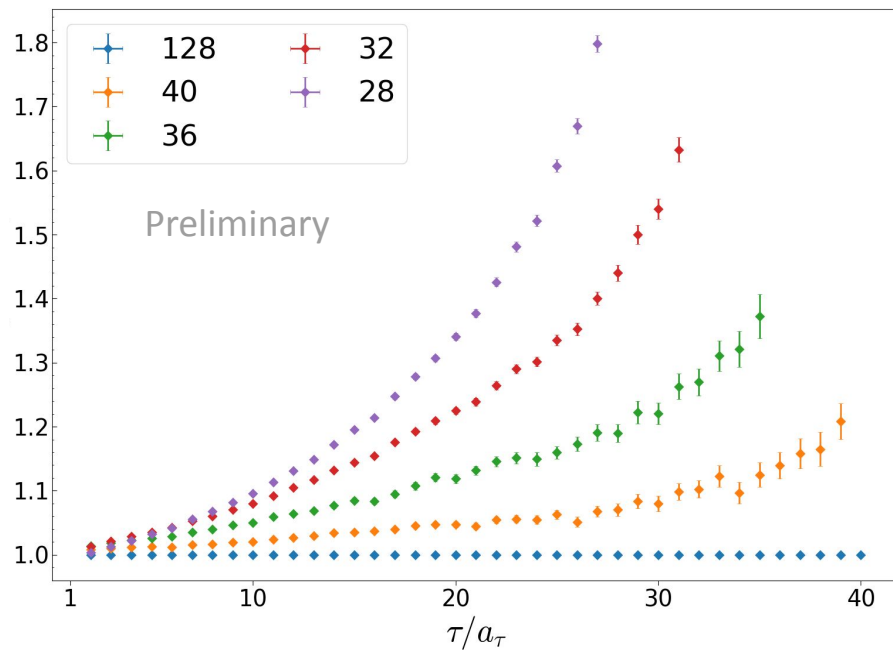


Double Ratio

Differences from **one** show difference in correlator



$\Upsilon(1S)$

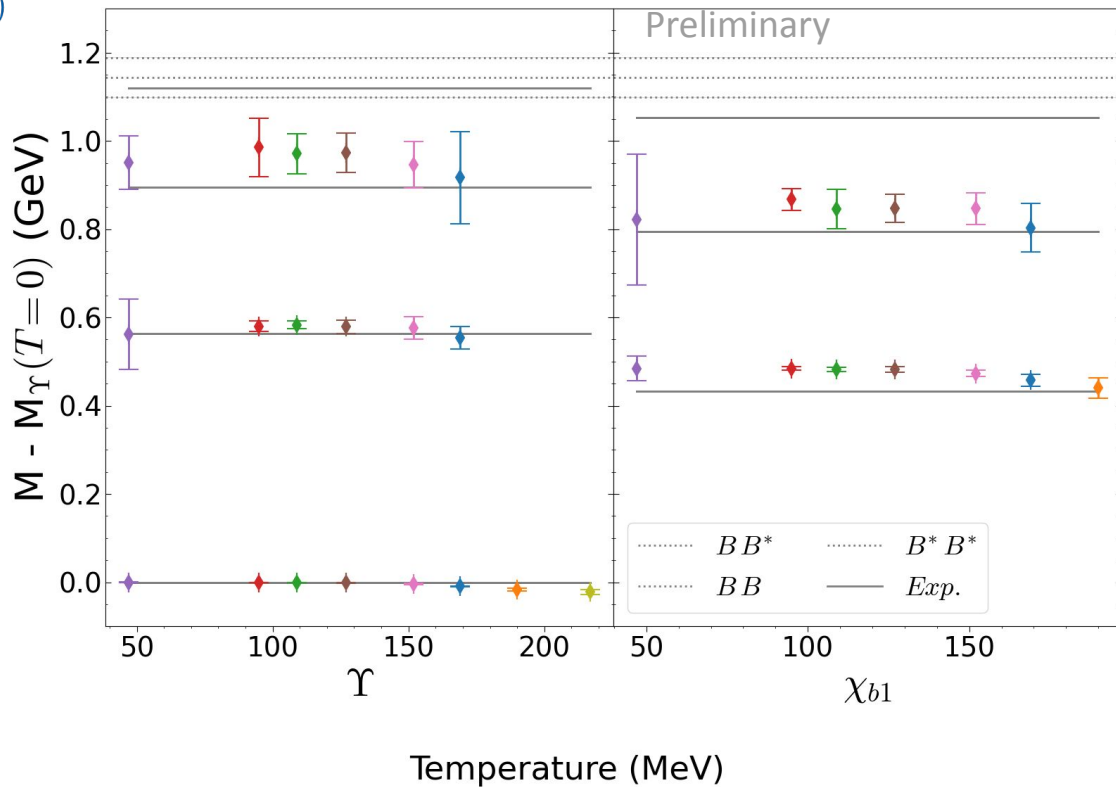


$\Upsilon(2S)$

Mass Spectrum Results

Subtract zero-temperature $\Upsilon(1S)$

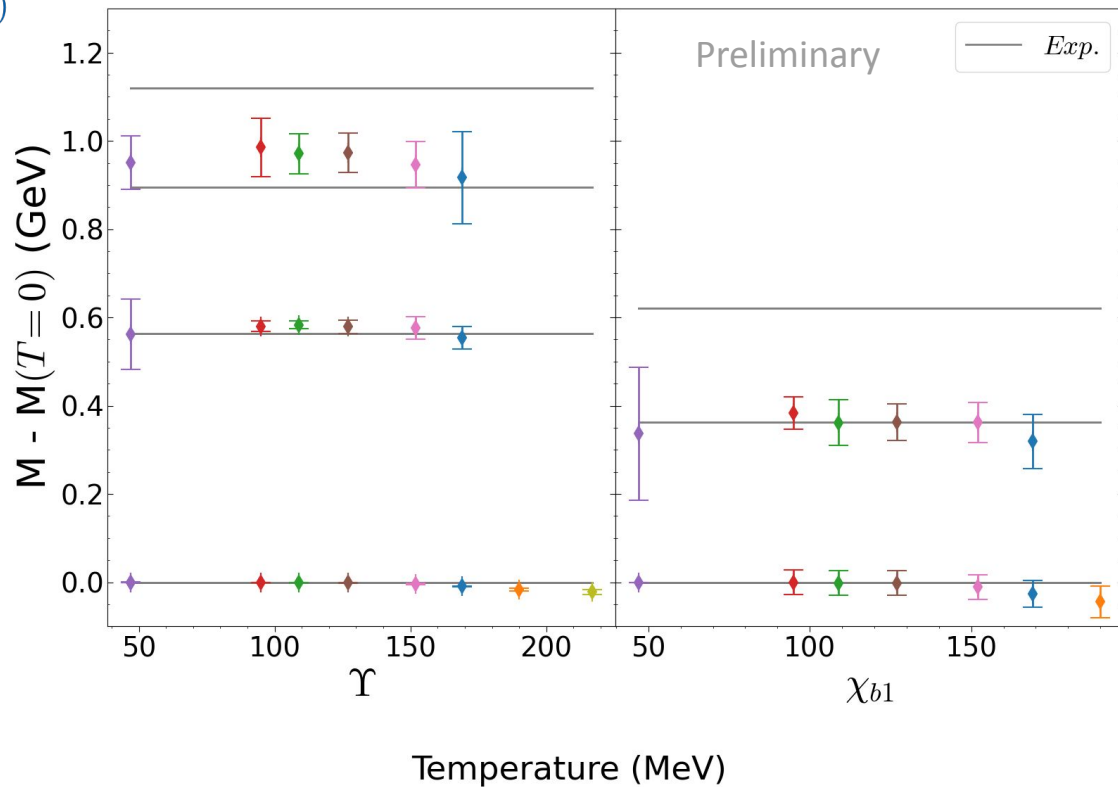
- Double Ratio informs trust in standard (multi-) exponential fits
- $$\sum_i A_i e^{-E_i \tau}$$
- Model averaging techniques used to give robust determination of energy.



Mass Spectrum Results

Subtract zero-temperature $\Upsilon(1S)$
or $\chi_{b1}(1P)$

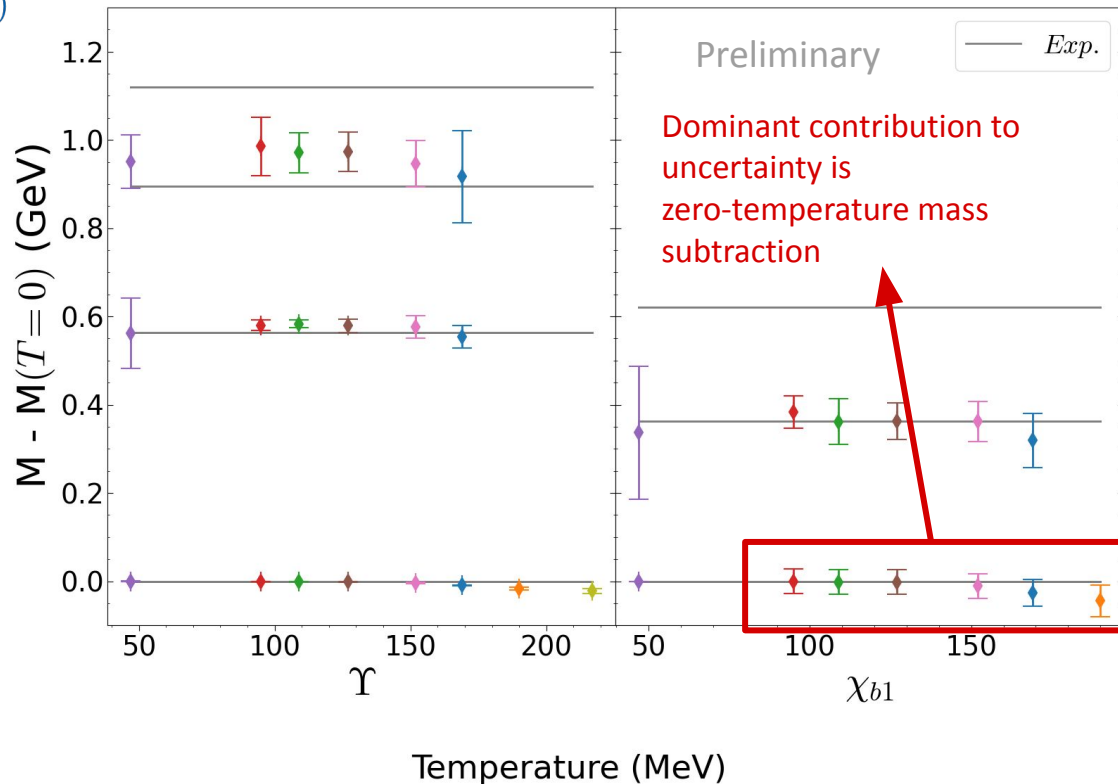
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`Moments`

`Time-Derivative Moments`

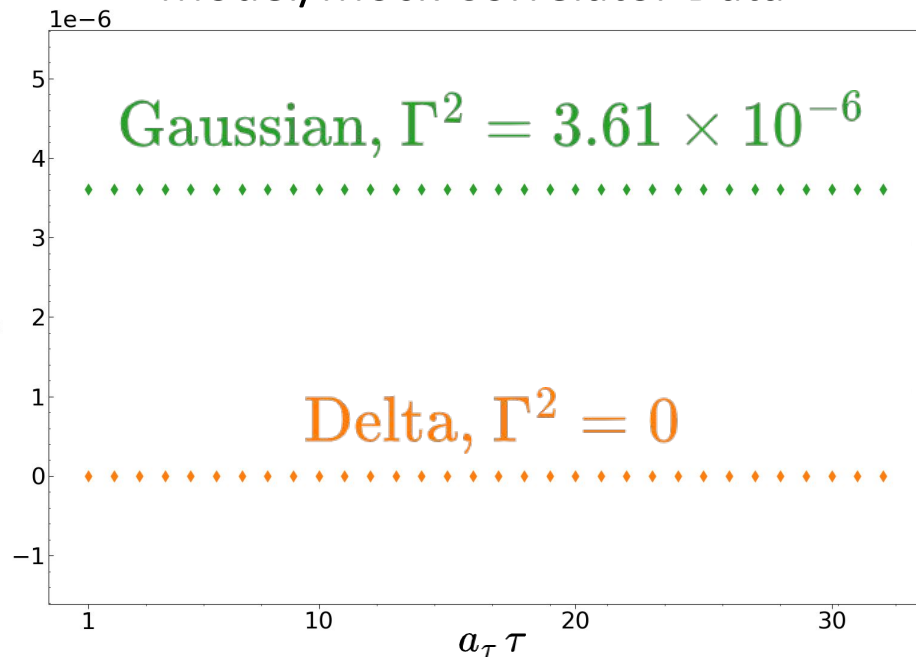
$$G(\tau) = \int_0^\infty \frac{d\omega}{2\pi} e^{-\omega\tau} \rho(\omega)$$

If $\rho(\omega)$ is Gaussian with width Γ and mean E , second log-derivative is

$$\begin{aligned} \frac{d^2 \log(G(\tau))}{d\tau^2} &= \frac{G''(\tau)}{G(\tau)} - \left(\frac{G'(\tau)}{G(\tau)} \right)^2 \\ &= \cancel{E^2} + \Gamma^2 - \cancel{(E)^2} \\ &= \Gamma^2 \end{aligned}$$

This is the difference between 2nd and 1st non-central moments of a Gaussian

Model/Mock Correlator Data



`Moments`

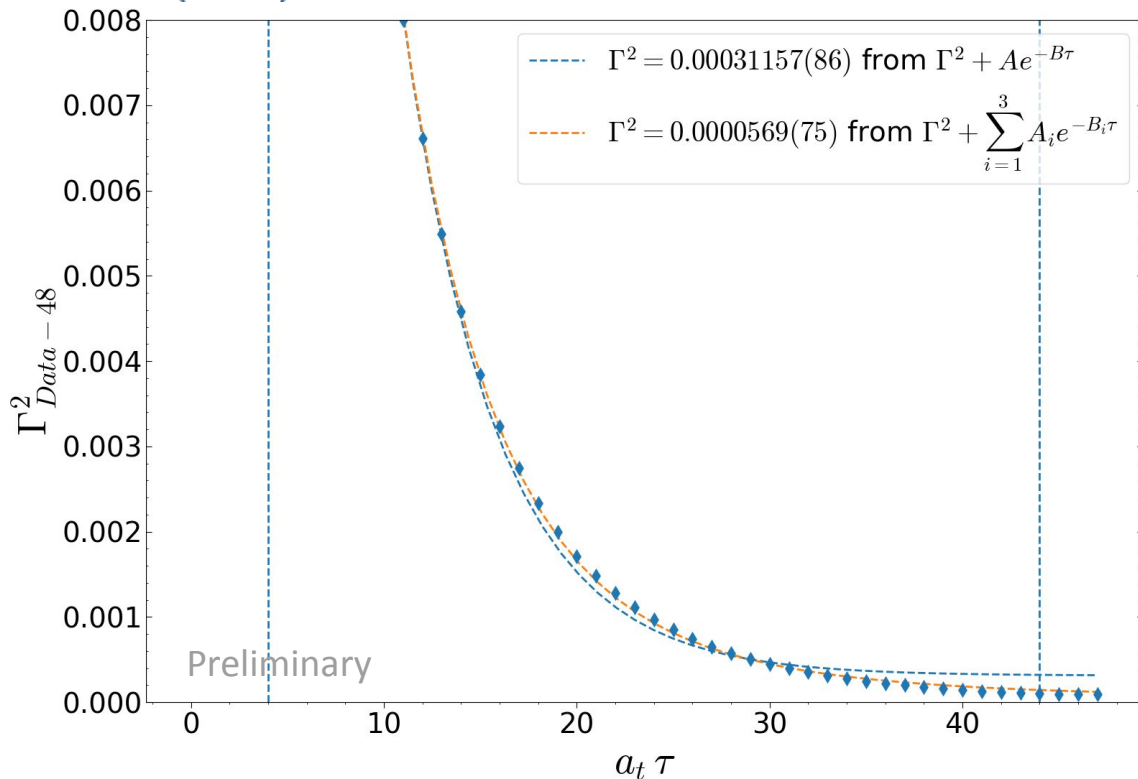
Point-Point

- Excited states shift form
- Fit with function

$$\Gamma^2 + \sum_{i=1}^N A_i e^{-B_i \tau}$$

- Easier at higher temperatures as Γ^2 becomes larger
- This is an upper bound only

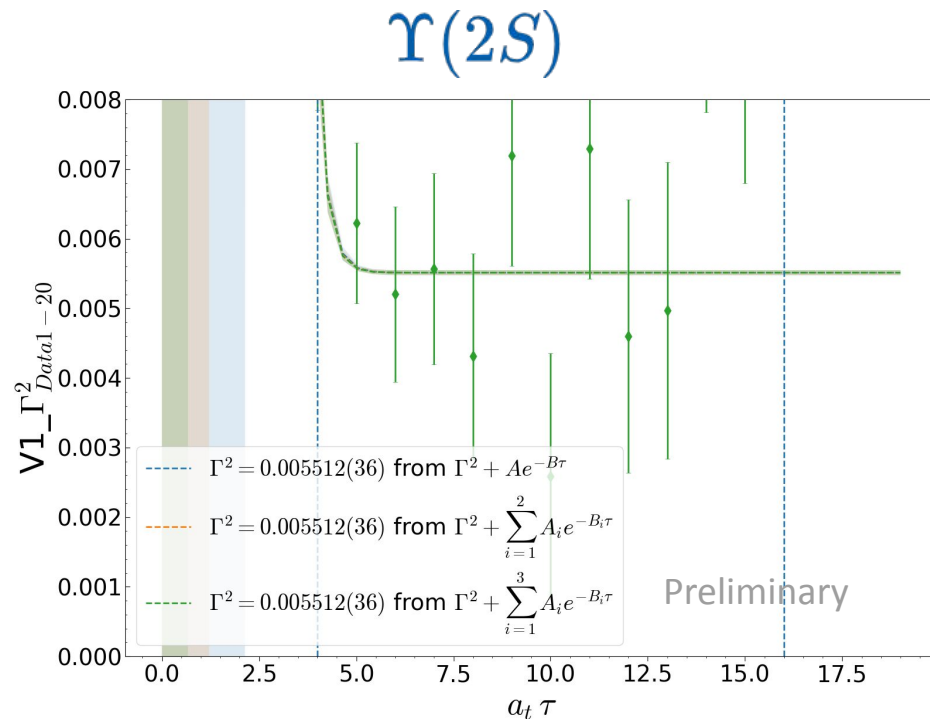
$\Upsilon(1S)$ Point-Point Correlator



`Moments`

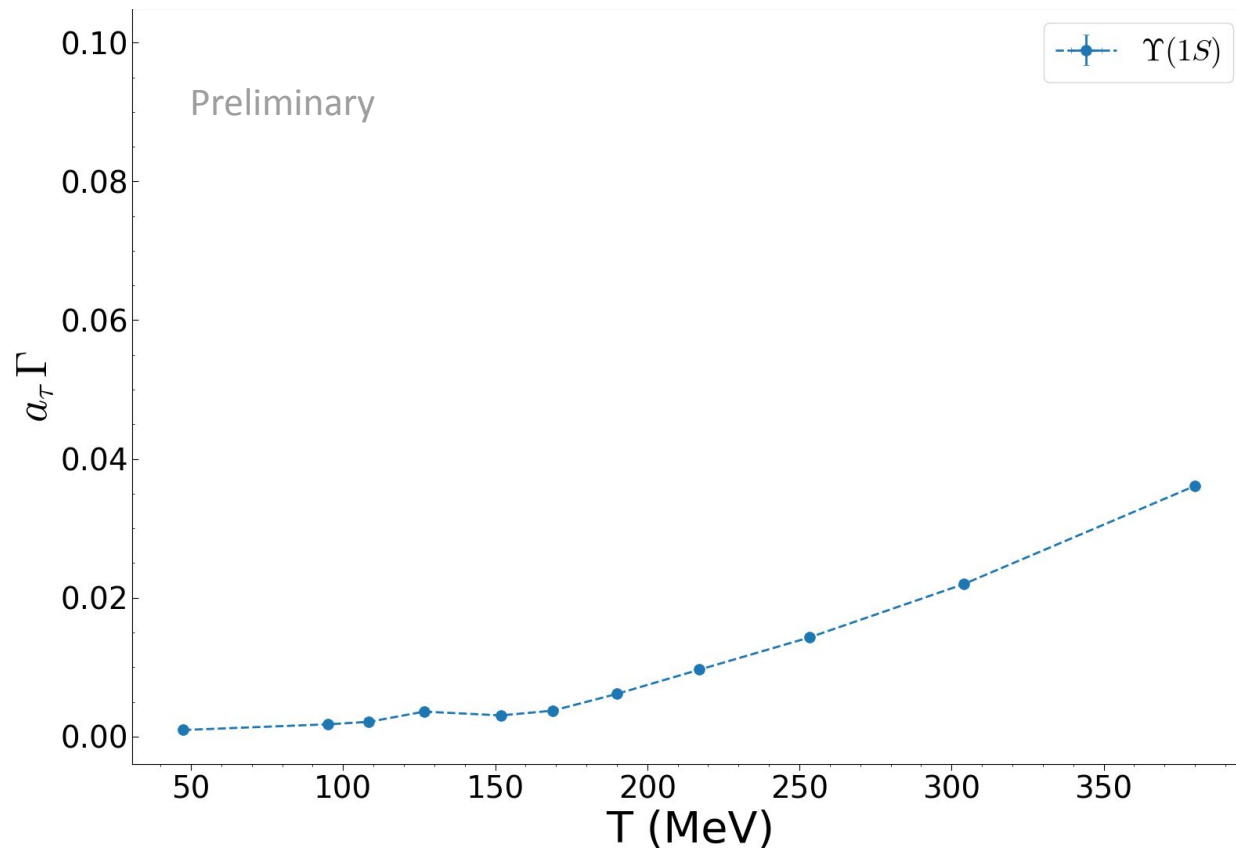
GEVP

- Apply `moments` method to GEVP projected correlators
- GEVP essential for access to excited states for moments
- Method is fairly robust against noise
 - Constant Γ^2 term helps
 - Exponential terms not well constrained
 - More statistics ongoing



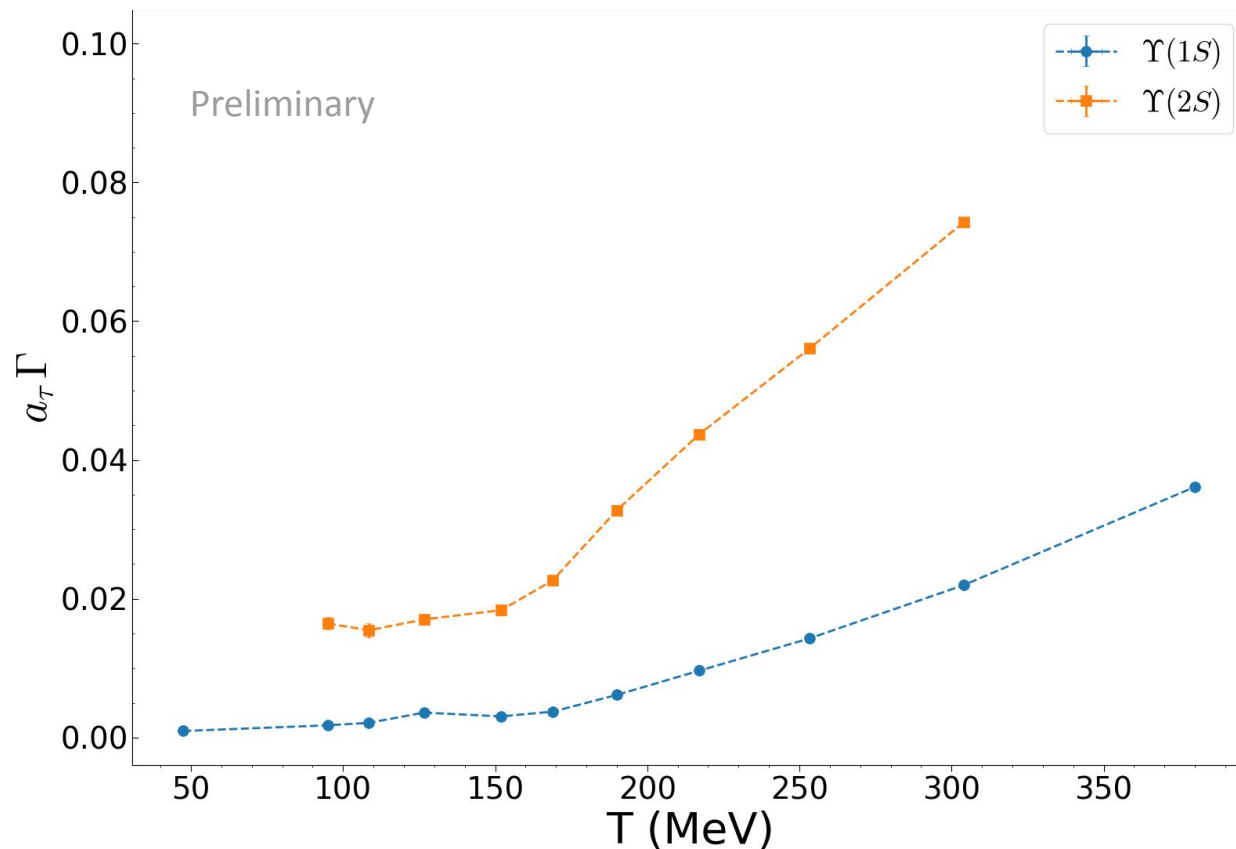
`Moments`

GEVP



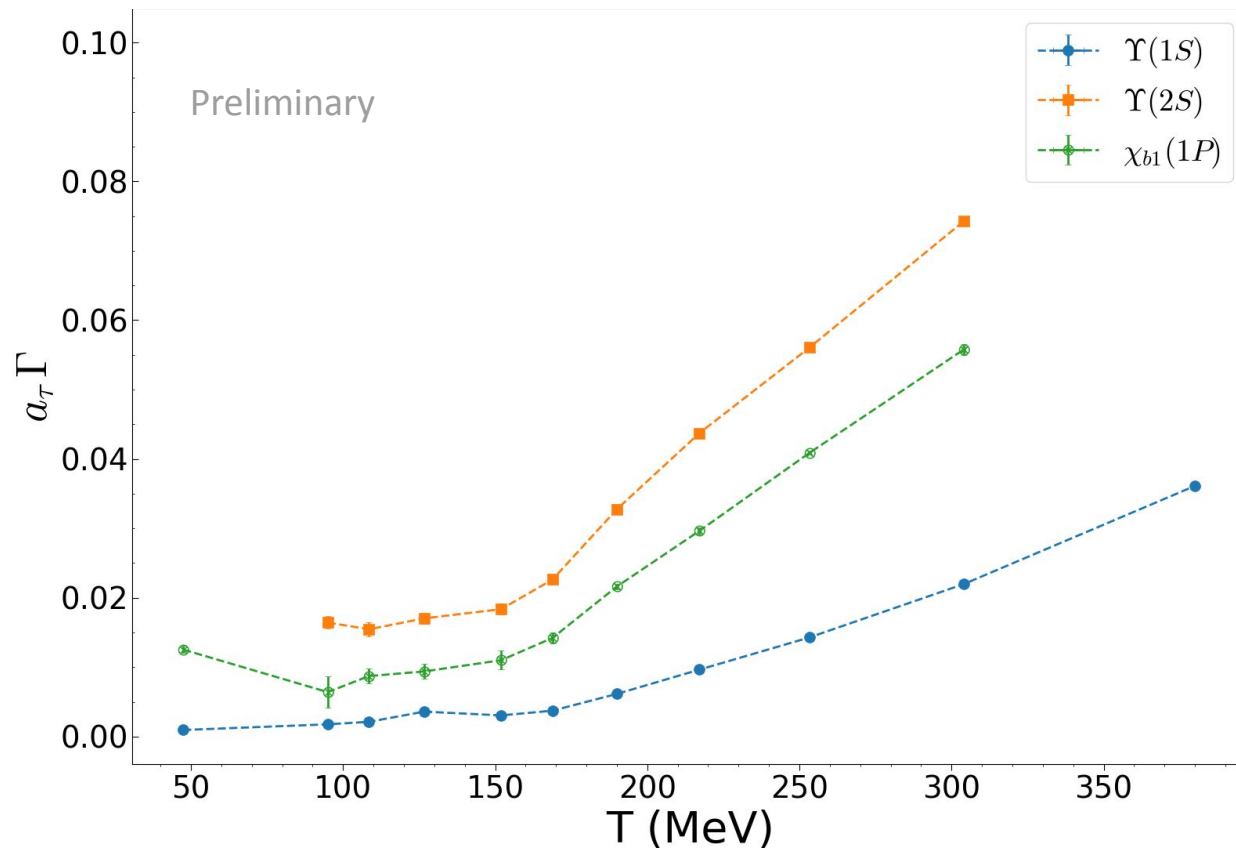
`Moments`

GEVP



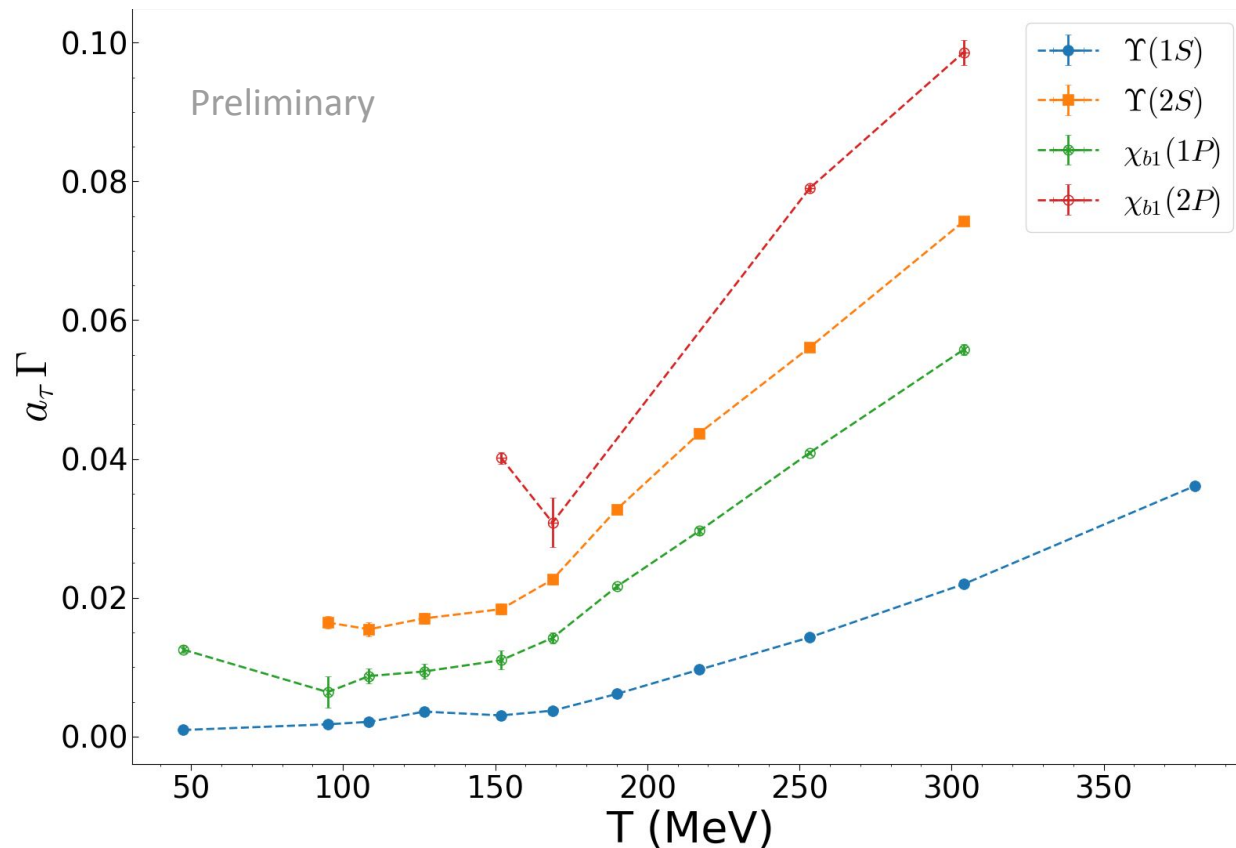
`Moments`

GEVP



'Moments'

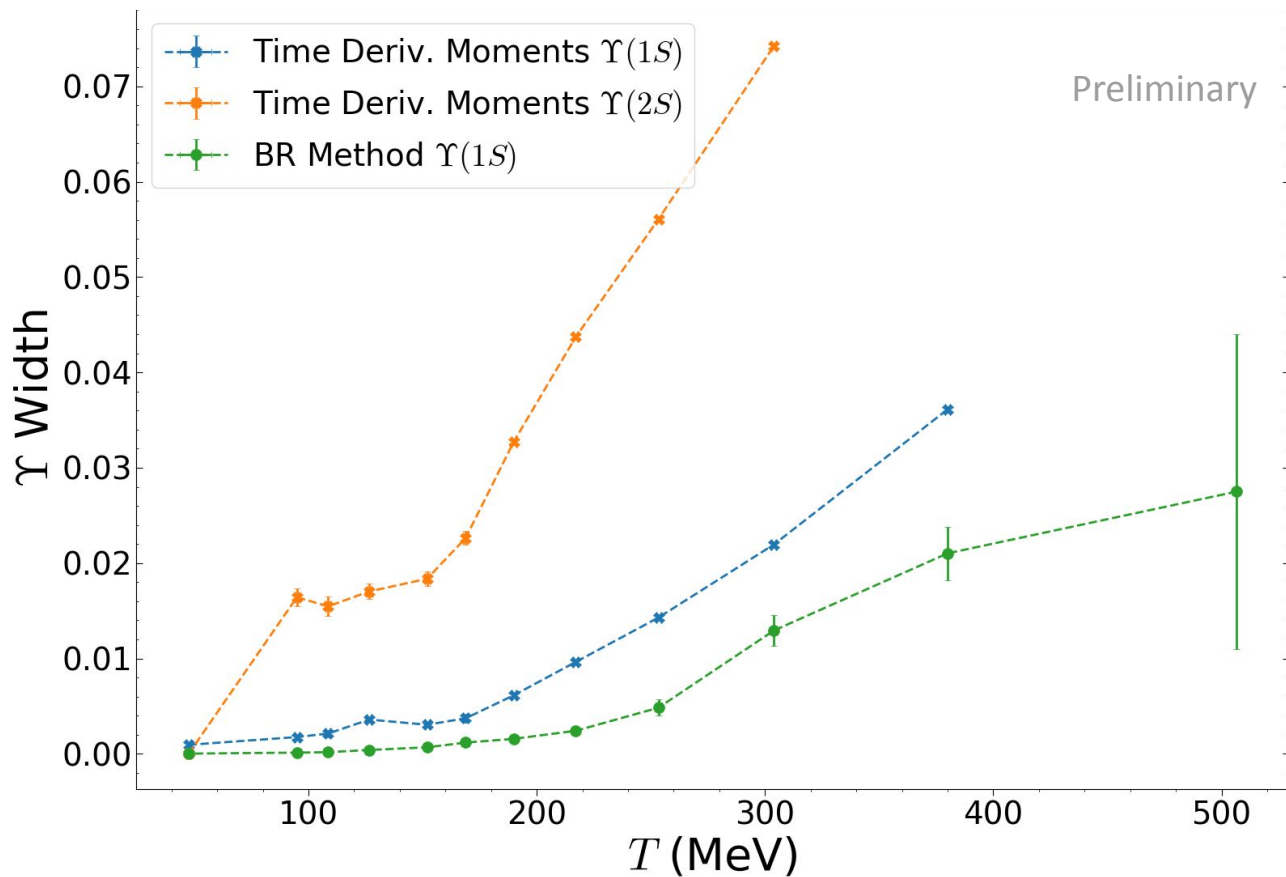
GEVP



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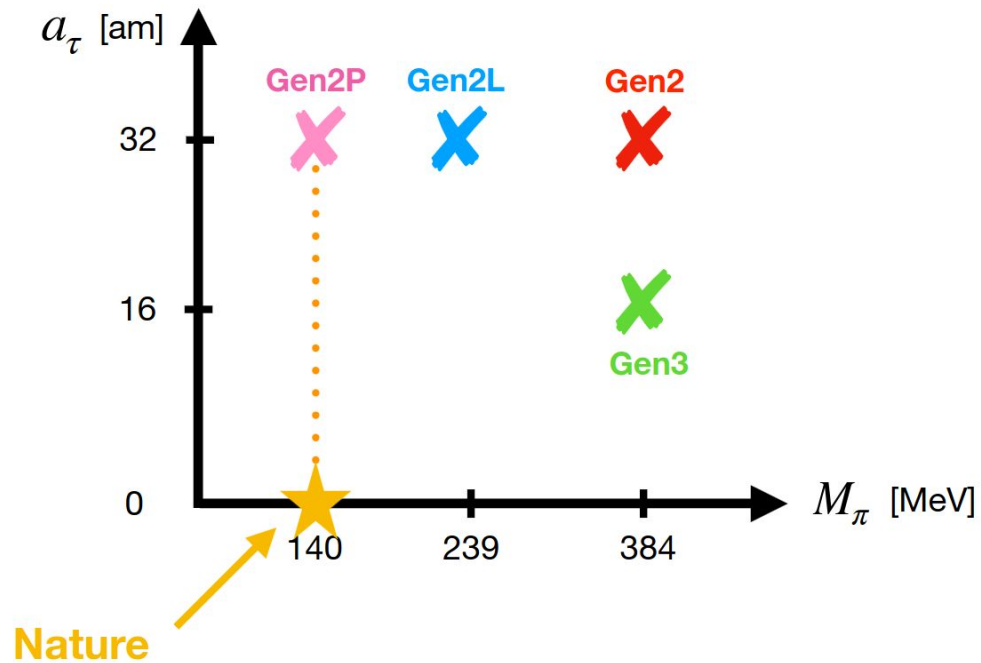
Comparison

- Bayesian Reconstruction method
- Moments method for ground & excited states
- Encouraging similarity between methods
- Excited state is broader than ground state



Summary

- Presented results for the mass of Υ and χ_{b1} excited states using a basis of ‘smeared’ operators
 - At zero and finite temperature
- (Re-)introduced ‘moments’ method to examine ‘widths’ of ground state (Gaussian) spectral functions
- Applied ‘moments’ to GEVP projected correlators
- GEVP of smeared operators was successful in allowing use of the ‘moments’ method for excited states
- Systematics of method not fully explored for this study (GEVP correlators)





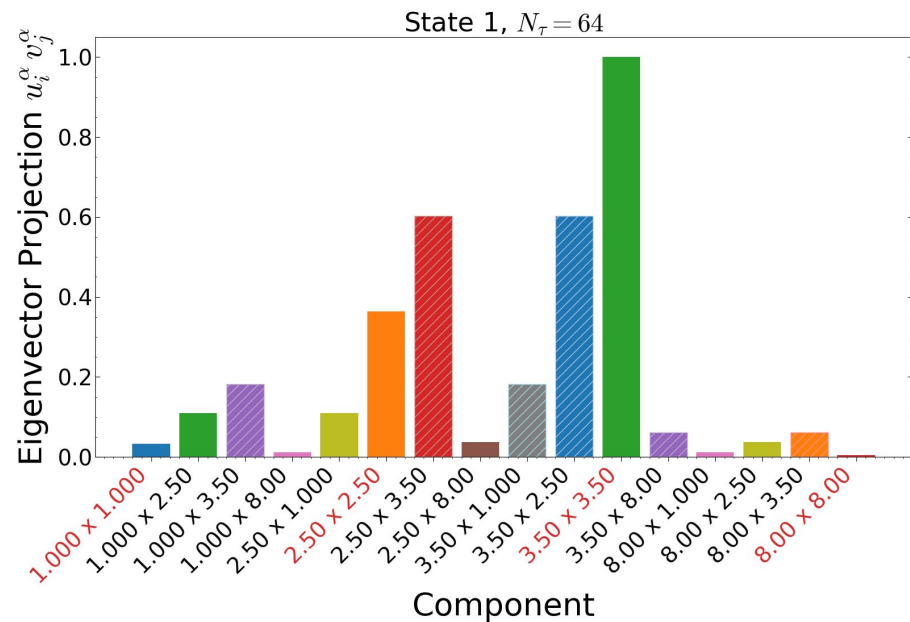
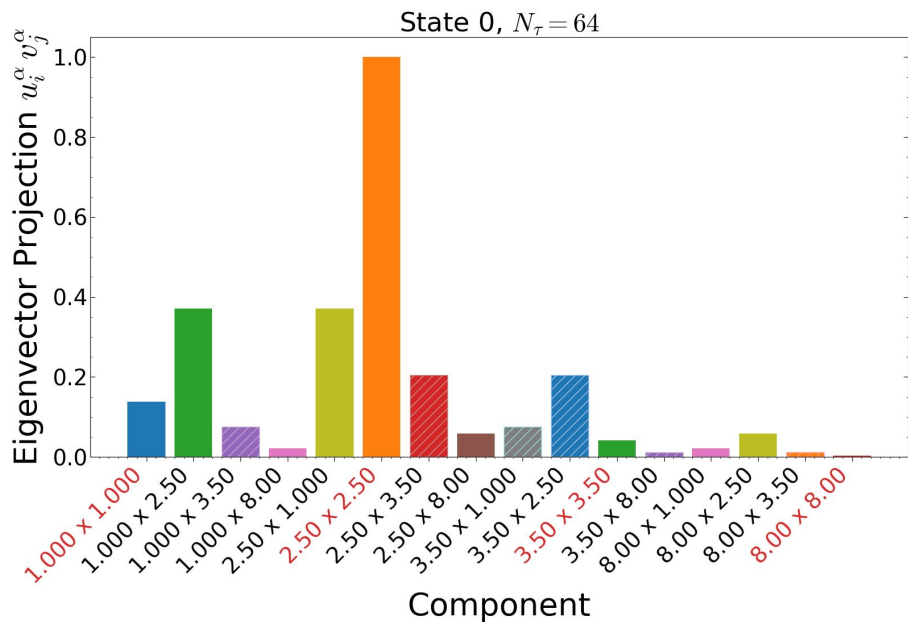
Trinity College Dublin

Coláiste na Tríonóide, Baile Átha Cliath

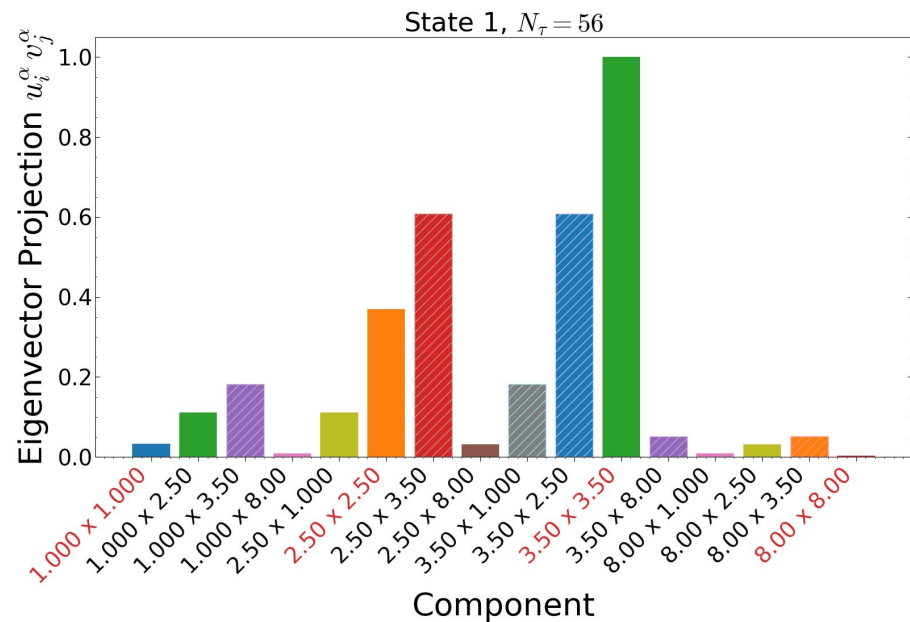
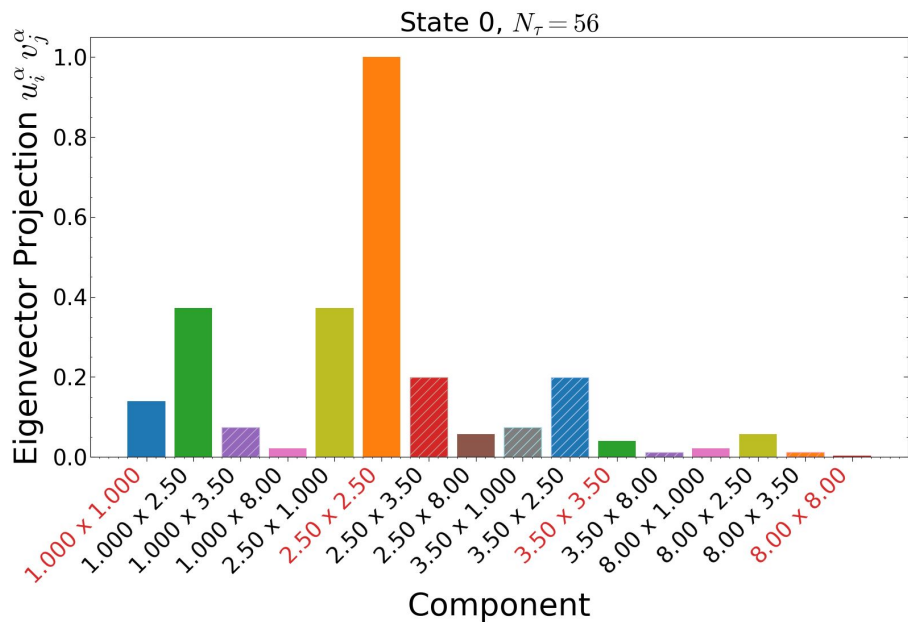
The University of Dublin

Additional Slides

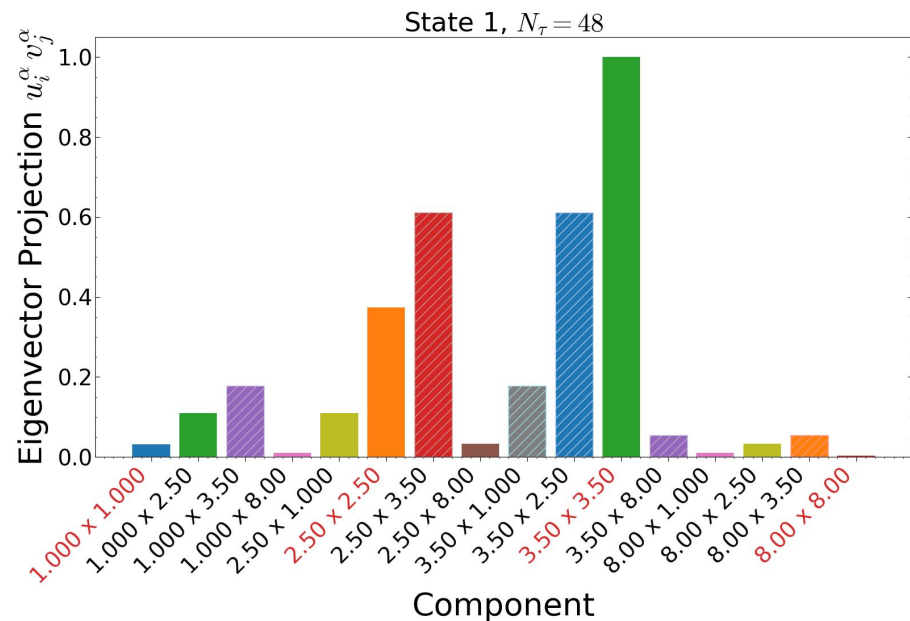
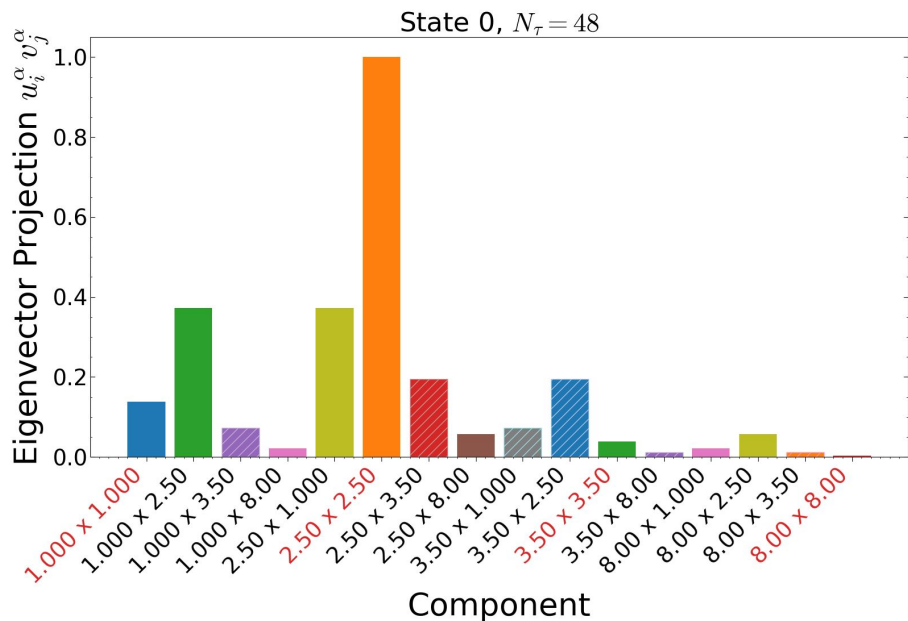
Eigenvectors



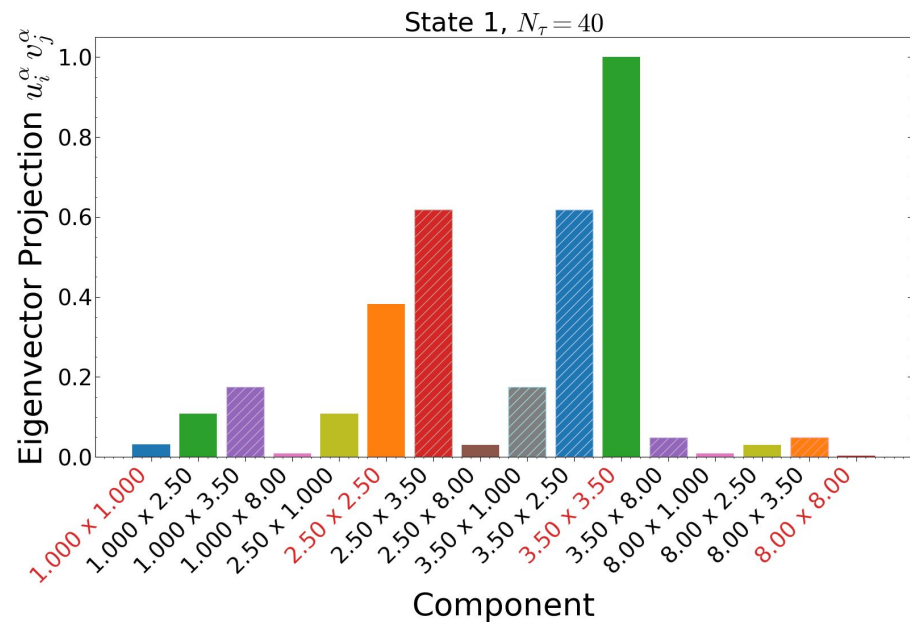
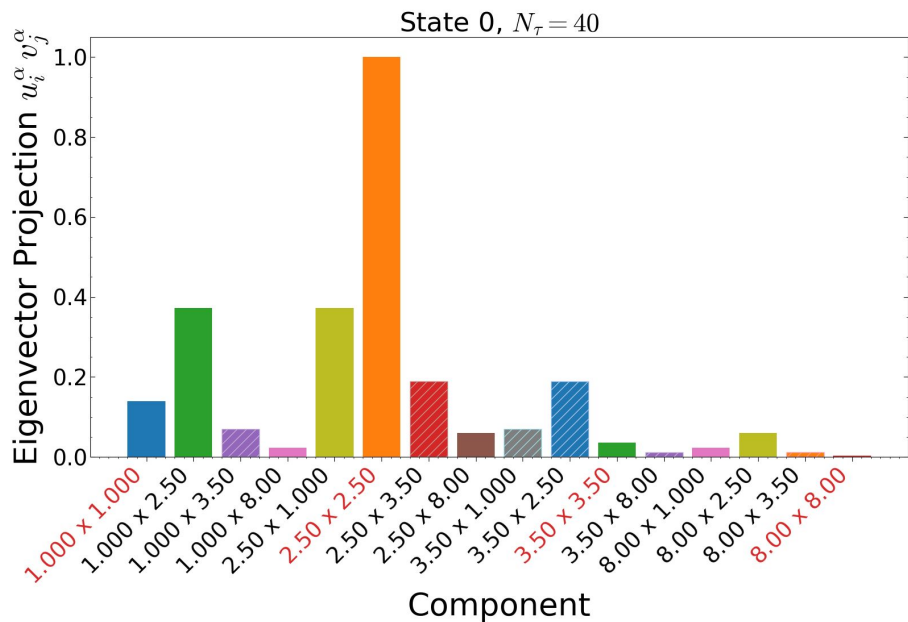
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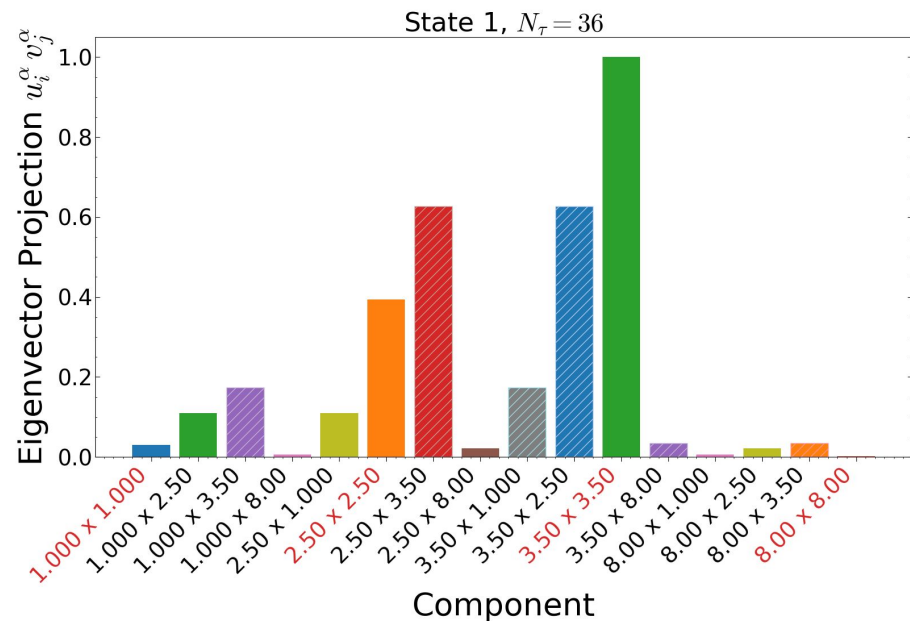
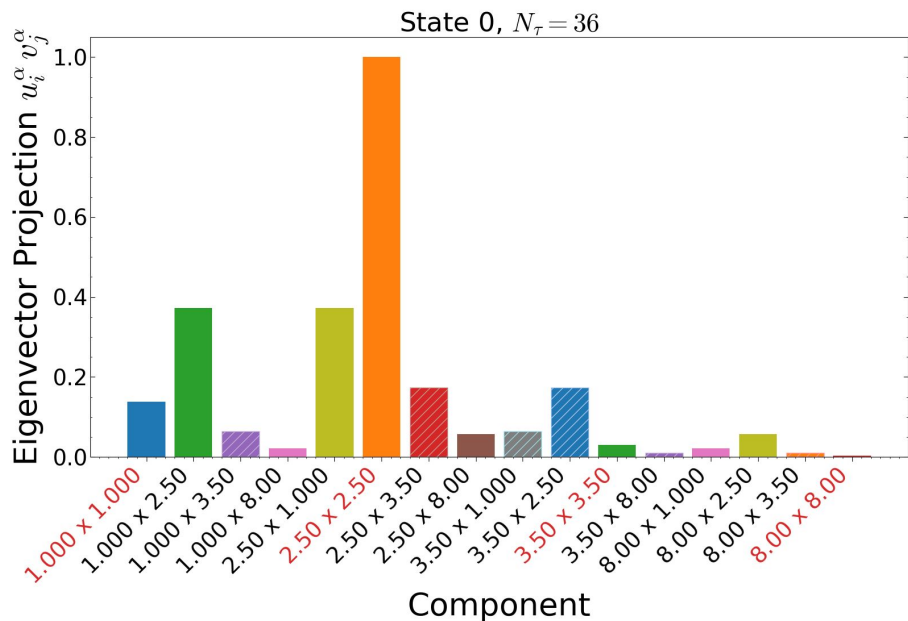
Eigenvectors



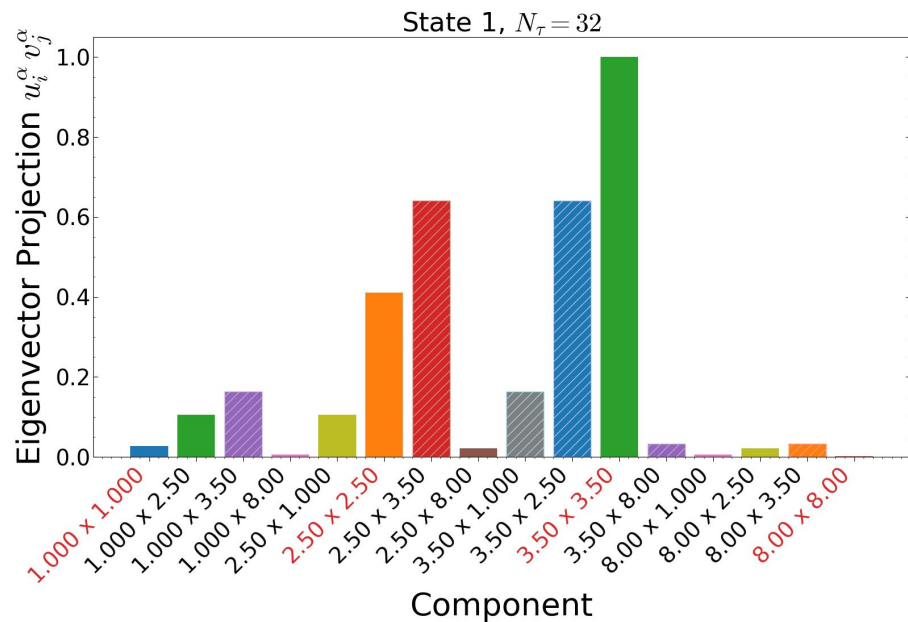
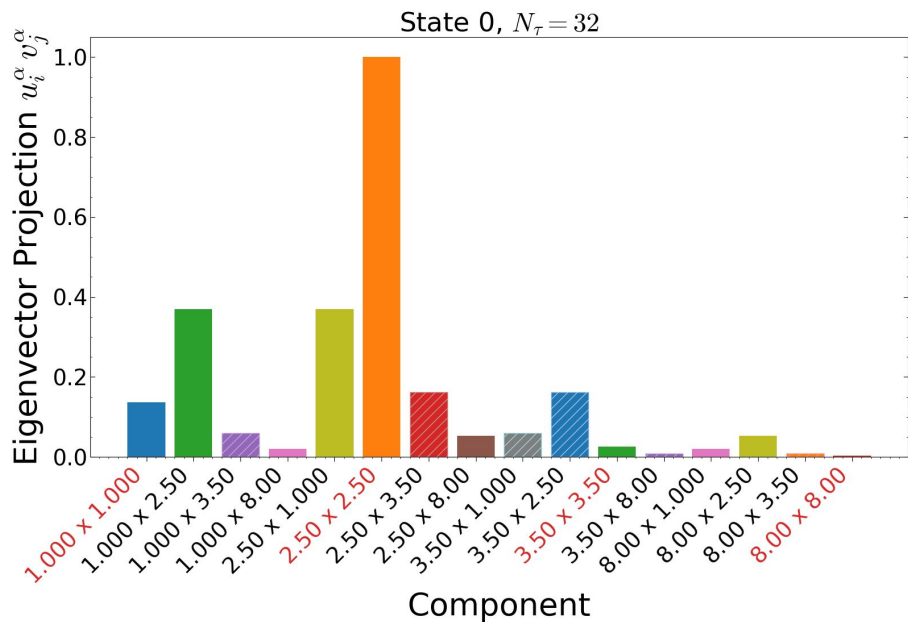
Eigenvectors



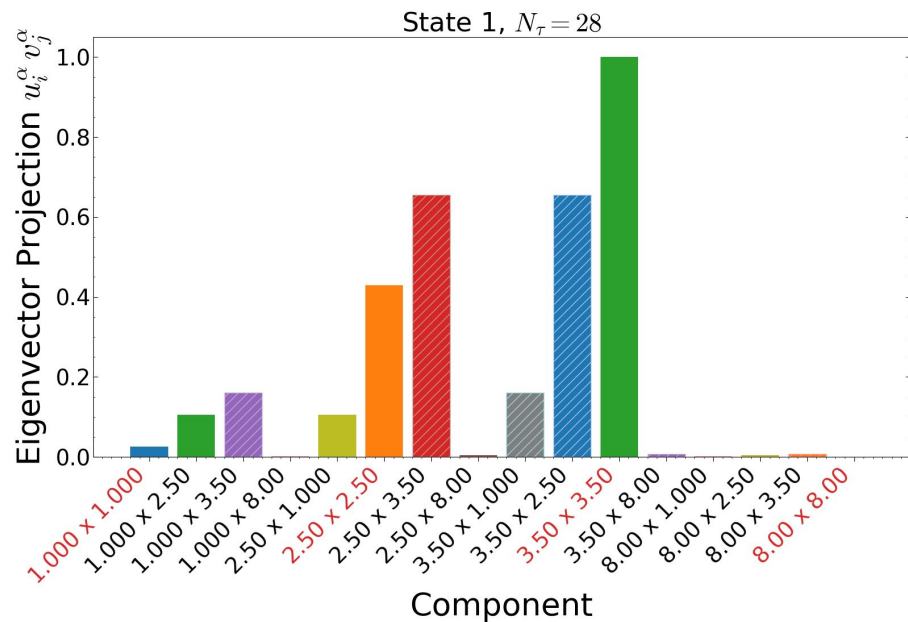
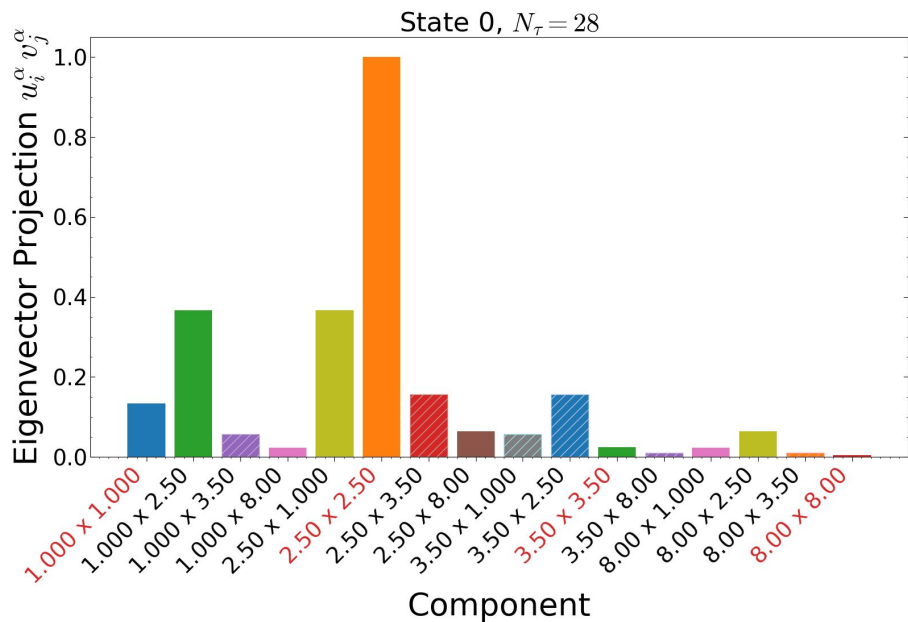
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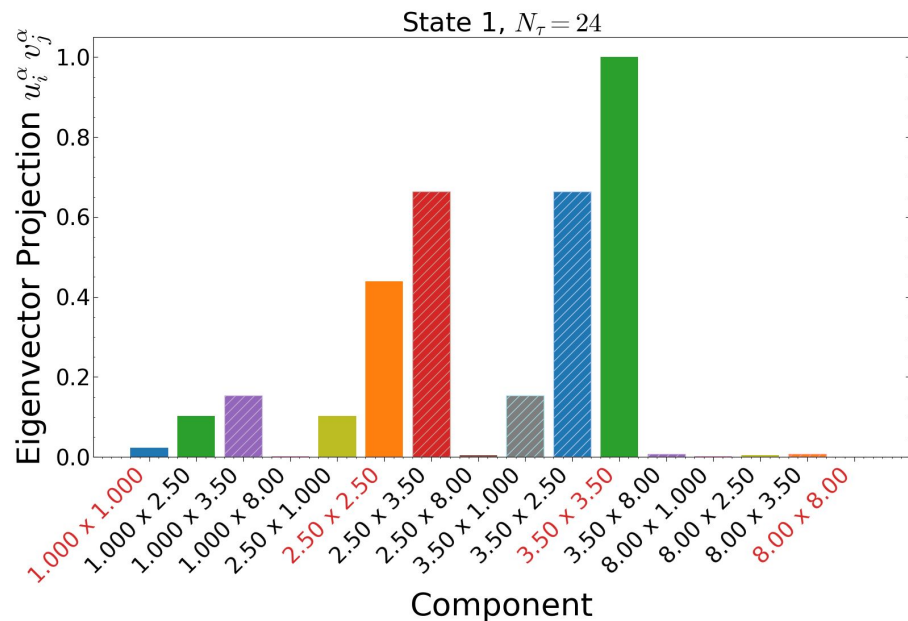
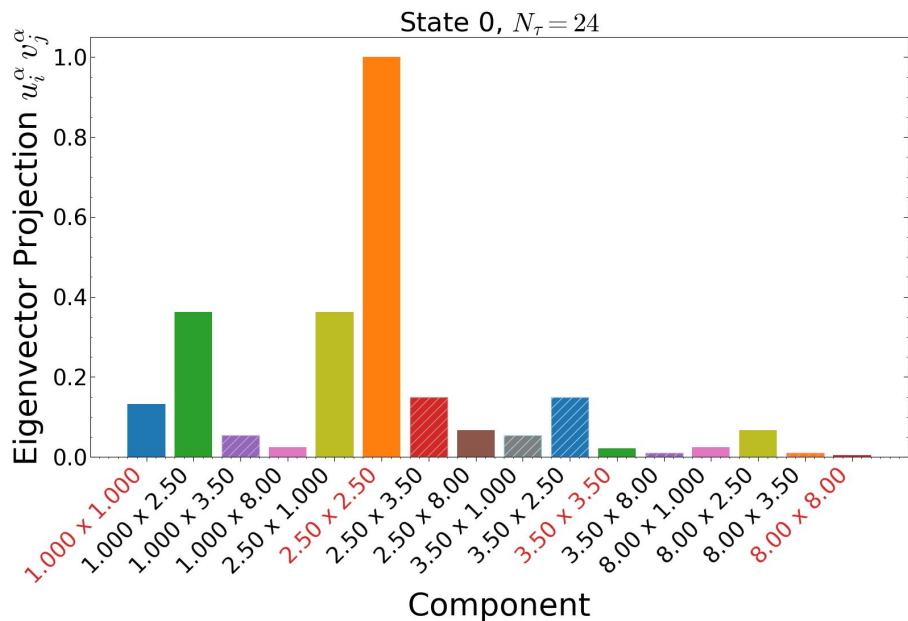
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