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# Pseudoscalar Screening Mass at Finite Temperature and Magnetic Field

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## Motivation



[Source: P. Costa, et al., arXiv:1712.08387v1]

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### Effects of magnetic field on strong interacting matter

Magnetic catalysis (MC) Inverse Magnetic catalysis (IMC) Observed at zero temperature Observed near  $T_{pc}$ Chiral condensate suppressed Chiral condensate enhanced lattice results show decreased T<sub>nc</sub> led to theoretical prediction of increased  $T_{pc}$ Valence guarks dominate Sea quarks dominate [Buividovich et al., PhysLetB.2009.11.017, [Bali et al., Phys.Rev.D.86.071502, D'Elia et al., PhysRevD 82.051501, Bruckman et al., JHEP04(2013)112, H.-T. Ding et al., PhysRevD.104.014505] H.-T. Ding et al., Phys.Rev.D105.034514]



Pseudoscalar Screening Mass at Finite Temperature and Magnetic Field

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Ward Identity: relates chiral condensate with chiral susceptibility [H.-T. Ding et al.

PhysRevD.104.014505]

$$(m_{u}+m_{d})\chi_{\pi^{0}}=\langle\bar{\psi}\psi\rangle_{u}+\langle\bar{\psi}\psi\rangle_{d}$$

Chiral susceptibility  $\chi_H$ : integrated screening correlator

$$\chi_H = \int dz \ C_H(z)$$

Screening correlators  $C_H(z)$ : two point correlation function of pseudoscalar meson H propagating in spatial direction z

$$C_H(z,T) = \int_0^{1/T} d\tau \int dx \int dy \langle J_H^{\dagger}(x,y,z,\tau) J_H(0,0,0,0) \rangle$$

At large distances, the screening correlator decays exponentially giving us the screening mass  ${\cal M}$ 

$$z \to \infty$$
,  $C(z) \to e^{-zM(T)}$ 

The inverse of a screening mass M is the screening length, i.e., the spatial distance beyond which the effects of a test hadron are effectively screened.

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- How does the long distance behavior of correlation function depend on the magnetic field?
- Does it have a similar behavior like chiral condensate?







H.-T. Ding et al., Phys.Rev.D105,034514

- A long history of screening mass analysis at zero magnetic field DeTar, Kogut, Phys.Rev.Lett.59.399, Gottlieb et al., Phys.Rev.Lett.59.1881, Forcrand et al., Phys.Rev.D63.054501, Gavai et al., Phys.Rev.D67.034501, [HotQCD,EPJC,71(2):113,2011, PhysRevD.100.094510]
- Screening mass at finite magnetic field: H.-T. Ding et al., Phys.Rev.D105,034514 described the behavior of lattice screening mass on a wide temperature range at  $M_\pi=220~MeV$
- Our goal:
  - 1 zoom in near critical point to observe dominant effect
  - 2 continuum extrapolated pseudoscalar screening mass near T<sub>pc</sub> at physical quark mass.









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Lattice setup			

- A tree-level improved Symanzik gauge action lattices with  $N_f = 2 + 1$  with Highly improved staggered quarks (HISQ) and uniform magnetic fields along the *z* direction used.
- $N_{\tau} = 8, 12, 16$  with aspect ratio  $N_{\sigma}/N\tau = 4$  fixed.
- Scale set by fixing the kaon decay constant  $f_K = 156.1/\sqrt{2}$  MeV. HotQCD, Phys.Rev.D90,094503, PhysRevD.100.094510
- The strange quark mass  $m_s$  tuned to its physical value by fixing  $M_{\eta_{s\bar{s}}} = 686$  MeV and light quark mass  $m_l = m_s/27$ .
- The magnetic flux on the lattice is quantized using the formulation

$$eB = \frac{6\pi N_b}{N_\sigma^2 a^2} = 6\pi N_b T^2 \frac{N_\tau^2}{N_\sigma^2}$$
(1)

- Temperatures range: T = 145 to 166 MeV
- Magnetic flux: Nb = 1, 2, 3, 4, 6, 12, 16, 24
- Magnetic field strengths range: eB = 0 to 0.8 GeV<sup>2</sup>

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# Screening mass on lattice

The periodic boundary on lattice results in a symmetric behavior of the correlator given by

$$C(n_z) = \sum_i A_i \cosh\left(M_i(n_z - N_\sigma/2)\right)$$
(2)



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Continuum ext	randation		

- Use the lattice observable O for each  $N_{\tau}$  and perform a 2D interpolation in T-B plane.
- Given our action has a correction of  $O(a^2)$ , we perform the continuum extrapolation in  $N_{\tau}^2$  using a linear ansatz on two finest lattices and quadratic ansatz on all three lattices

$$O_{N_{\tau}} = O_{cont}^{linear} + \frac{b}{N_{\tau}^2}$$
(3)

$$O_{N_{\tau}} = O_{cont}^{quadratic} + \frac{b}{N_{\tau}^2} + \frac{c}{N_{\tau}^4}$$
(4)

• Considering the values obtained from linear and quadratic methods as independent values, we obtain the continuum screening mass

$$O_{cont} = \frac{1}{2} (O_{cont}^{linear} + O_{cont}^{quadratic})$$
(5)

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### Neutral meson: Screening Mass vs Magnetic Field





- H. Ding et al., 2022 showed dominatingly decrease screening mass at very high and very low temperatures.
- Near T<sub>pc</sub>, a minima beyond which the screening mass increases
- The magnetic field at which minima occurs decreases with increased temperature
- The minima shifts at larger magnetic field for mesons with larger mass

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# Neutral meson: Screening Mass vs Temperature





- Near T<sub>pc</sub>, screening mass at fixed magnetic field cross each other
- Crossing temperature increase with increasing meson mass
- Screening mass crossing not observed in our range temperature for ss meson

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# Charged meson: Screening Mass vs Magnetic Field



- The charged meson screening mass steadily increase with increasing magnetic field strength
- The increase seem to approach a constant with increasing magnetic field strength

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# Charged meson: Screening Mass vs Temperature



• The charged meson screening mass at constant magnetic field strength increase with temperature

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# Summary

- Numerical estimates of continuum extrapolated neutral and charged meson screening mass
- Continnuum extrapolated neutral pseudoscalar screening mass near the  $T_{pc}$  indicate large distance behavior have similar behavior like chiral condensate, i.e.,
  - 1 early deviation from a T=0 mass at larger magnetic field
  - 2 suppression of the magnetic effect with increased meson mass
- Charged pseudoscalar screening mass increase with increasing magnetic field.

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