

Pseudoscalar Screening Mass at Finite Temperature and Magnetic Field

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Motivation

[Source: P. Costa, et al., arXiv:1712.08387v1]

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Effects of magnetic field on strong interacting matter

Rleated reviews: Endrodi arXiv:2406.19780 Andersen et al.,
Rev.Mod.Phvs. 88. Rev.Mod.Phys. 025001 Kharzeev et al. Springer Vol. 871, 201

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Ward Identity: relates chiral condensate with chiral susceptibility [H.-T. Ding et al.

PhysRevD.104.014505]

$$
(m_u+m_d)\,\chi_{\pi^0}=\langle\bar{\psi}\psi\rangle_u+\langle\bar{\psi}\psi\rangle_d
$$

Chiral susceptibility χ _H: integrated screening correlator

$$
\chi_H = \int dz \ C_H(z)
$$

Screening correlators $C_H(z)$: two point correlation function of pseudoscalar meson H propagating in spatial direction z

$$
C_H(z, T) = \int_0^{1/T} d\tau \int dx \int dy \langle J_H^{\dagger}(x, y, z, \tau) J_H(0, 0, 0, 0) \rangle
$$

At large distances, the screening correlator decays exponentially giving us the screening mass M

$$
z \to \infty
$$
, $C(z) \to e^{-zM(T)}$

The inverse of a screening mass M is the screening length, i.e., the spatial distance beyond which the effects of a test hadron are effectively screened.

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- *•* How does the long distance behavior of correlation function depend on the magnetic field?
- *•* Does it have a similar behavior like chiral condensate?

H.-T. Ding et al., Phys.Rev.D105,034514

- *•* A long history of screening mass analysis at zero magnetic field DeTar, Kogut, Phys.Rev.Lett.59.399, Gottlieb et al.,Phys.Rev.Lett.59.1881, Forcrand et al.,Phys.Rev.D63.054501, Gavai et al., Phys.Rev.D67.034501, [HotQCD,EPJC,71(2):113,2011, PhysRevD.100.094510]
- Screening mass at finite magnetic field: H.-T. Ding et al., Phys.Rev.D105,034514 described the behavior of lattice screening mass on a wide temperature range at M*π* = 220 MeV
- *•* Our goal:
	- 1 zoom in near critical point to observe dominant effect
	- 2 continuum extrapolated pseudoscalar screening mass near T_{pc} at physical quark mass. $\mathbf{A} \equiv \mathbf{A} + \mathbf{A} + \mathbf{B} + \mathbf{A} + \math$

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- A tree-level improved Symanzik gauge action lattices with $N_f = 2 + 1$ with Highly improved staggered quarks (HISQ) and uniform magnetic fields along the z direction used.
- $N_{\tau} = 8, 12, 16$ with aspect ratio $N_{\sigma}/N_{\tau} = 4$ fixed.
- Scale set by fixing the kaon decay constant $f_K = 156.1/\sqrt{2}$ MeV. HotQCD, Phys.Rev.D90,094503, PhysRevD.100.094510
- The strange quark mass m_s tuned to its physical value by fixing $M_{n_{\rm s\bar{s}}} = 686$ MeV and light quark mass $m_l = m_s/27$.
- *•* The magnetic flux on the lattice is quantized using the formulation

$$
eB = \frac{6\pi N_b}{N_\sigma^2 a^2} = 6\pi N_b T^2 \frac{N_\tau^2}{N_\sigma^2}
$$
 (1)

- Temperatures range: $T = 145$ to 166 MeV
- *•* Magnetic flux: Nb = 1, 2, 3, 4, 6, 12, 16, 24
- Magnetic field strengths range: $eB = 0$ to 0.8 GeV^2

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The periodic boundary on lattice results in a symmetric behavior of the correlator given by

$$
C(n_z) = \sum_i A_i \cosh\left(M_i(n_z - N_{\sigma}/2)\right) \tag{2}
$$

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- *•* Use the lattice observable O for each N*^τ* and perform a 2D interpolation in T-B plane.
- Given our action has a correction of $O(a^2)$, we perform the continuum extrapolation in \mathcal{N}_{τ}^2 using a linear ansatz on two finest lattices and quadratic ansatz on all three lattices

$$
O_{N_{\tau}} = O_{cont}^{linear} + \frac{b}{N_{\tau}^2} \tag{3}
$$

$$
O_{N_{\tau}} = O_{cont}^{quadratic} + \frac{b}{N_{\tau}^2} + \frac{c}{N_{\tau}^4}
$$
 (4)

• Considering the values obtained from linear and quadratic methods as independent values, we obtain the continuum screening mass

$$
O_{cont} = \frac{1}{2} (O_{cont}^{linear} + O_{cont}^{quadratic})
$$
\n(5)

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 $E = \Omega Q$

Neutral meson: Screening Mass vs Magnetic Field

- *•* H. Ding et al., 2022 showed dominatingly decrease screening mass at very high and very low temperatures.
- Near T_{pc}, a minima beyond which the screening mass increases
- *•* The magnetic field at which minima occurs decreases with increased temperature
- *•* The minima shifts at larger magnetic field for mesons with larger mass

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- Near T_{pc}, screening mass at fixed magnetic field cross each other
- *•* Crossing temperature increase with increasing meson mass
- *•* Screening mass crossing not observed in our range temperature for ss meson

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• The charged meson screening mass steadily increase with increasing magnetic field strength

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• The increase seem to approach a constant with increasing magnetic field strength

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• The charged meson screening mass at constant magnetic field strength increase with temperature

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- *•* Numerical estimates of continuum extrapolated neutral and charged meson screening mass
- *•* Continnuum extrapolated neutral pseudoscalar screening mass near the T_{pc} indicate large distance behavior have similar behavior like chiral condensate, i.e.,

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 \Rightarrow

- \bullet early deviation from a T=0 mass at larger magnetic field
- 2 suppression of the magnetic effect with increased meson mass
- Charged pseudoscalar screening mass increase with increasing magnetic field.

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