Finite temperature QCD phase transition with 3 flavors of Möbius domain wall fermions

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In collaboration with

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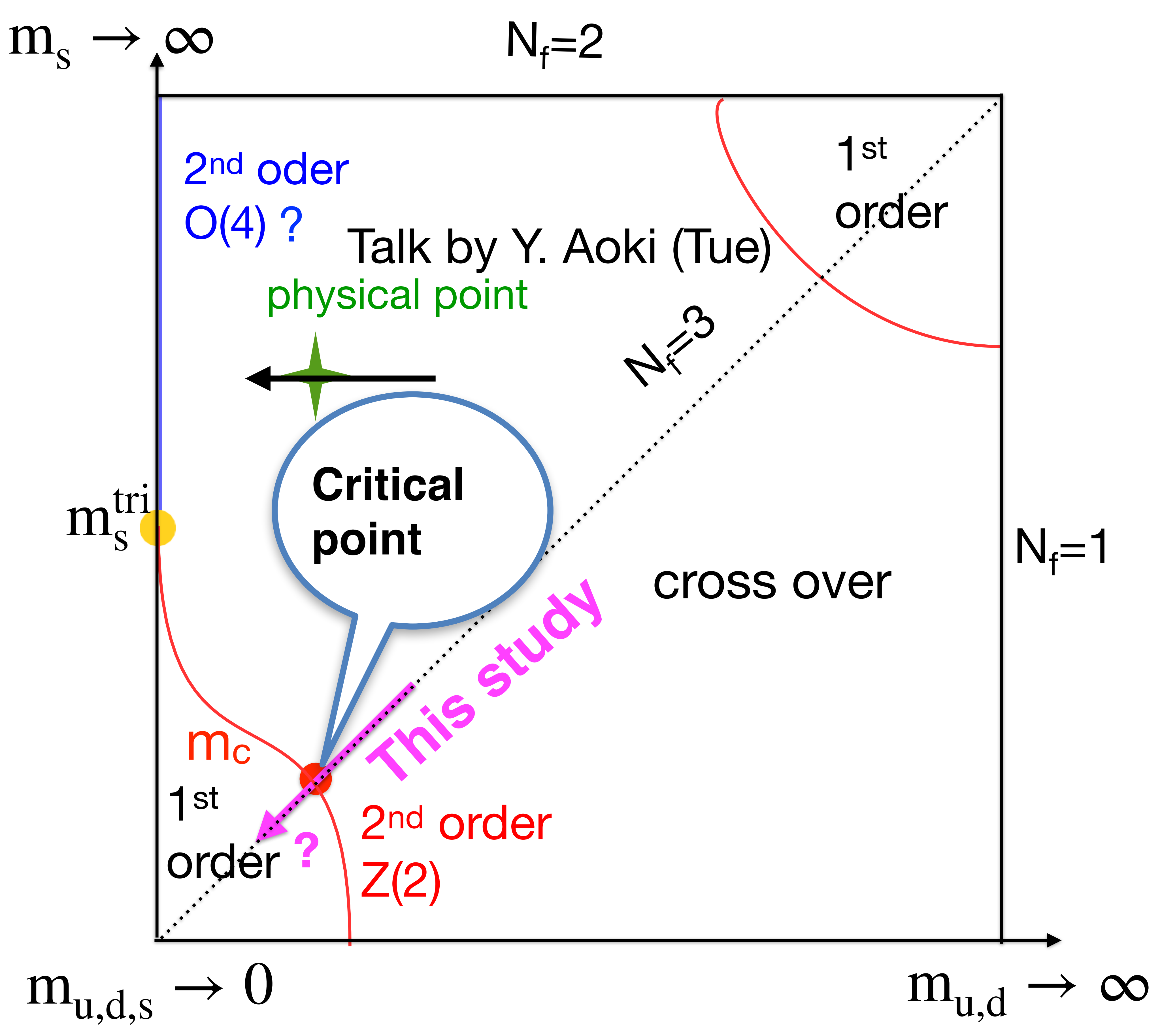
Lattice 2024

The nature of QCD phase transition at $\mu_B = 0$

Columbia plot



- ϵ expansion: 1st order phase transition in the chiral limit for $N_f = 3$

  [R. D. Pisarski, F. Wilczek PRD 84]

- Possible 2nd order phase transition in the $N_f = 3$ chiral limit:

  [G. Fejos, PRD 22]
  [S. R. Kousvos, A. Stergiou SciPost 23]
  [J. Bernhardt, C. S. Fischer PRD 23]
  [R. D. Pisarski, F. Rennecke PRL 24]
  [G. Fejos, T. Hatsuda arXiv:2404.00554]

Need to be checked by lattice QCD

Critical point on the $N_f = 3$ chiral region:

- Location? (existence?)
- Universality class?

Previous Nf=3 lattice QCD studies

Action	N_t	$m_\pi^{Z_2}$ [MeV]	Ref.
Staggered,standard	4	290	Karsch et al. (2001)
Staggered,standard	6	150	de Forcrand et al. (2007)
Staggered, HISQ	6	$\lesssim 50$	Bazavov et al. (2017)
Staggered, stout	4-6	$0?$	Varnhost (2014)
Wilson, standard	4	$\lesssim 670$	Iwasaki et al. (1996)
Wilson-Clover	6-10	$\lesssim 170$	Jin et al. (2017)
Wilson-Clover	6-12	$\lesssim 110$	Kuramashi et al. (2020)

➡ Strong cutoff and discretization effects

Evidence for chiral limit to feature 2nd order PT with staggered and HISQ fermion [F. Cuteri et al. JHEP 2021; S. Sharma et al. PRD 2022]

We propose to use chiral fermion (Mobius domain wall fermion)

- Exact chiral symmetry at finite a for infinite Ls
- Reduced χ_{SB} parameterized by residual mass when Ls is finite

Lattice Setup

- N_f=3 Mobius Domain Wall Fermion
- Tree-level Symanzik improved gauge action and stout smearing

★T=0:

$$\beta = 4.0, 24^3 \times 48 \times 16: \quad 0.02 \le am_q \le 0.045,\ am_{\mathrm{res}}(\mathrm{estimated}) \approx 0.006$$

$$\beta = 4.1, 24^3 \times 48 \times 16: \quad 0.015 \le am_q \le 0.040$$

$$\beta = 4.17, 32^3 \times 64 \times 16: \quad 0.012 \le am_q \le 0.026:$$

★T=121(2) MeV ($\beta = 4.0$, ($a = 0.1361(20)$ fm, determined from Wilson flow t_0))

$$N_s^3 \times 12 \times 16: \quad N_s = 48, -0.004 \le am_q \le -0.003$$

$$N_s = 36, -0.005 \le am_q \le 0.001$$

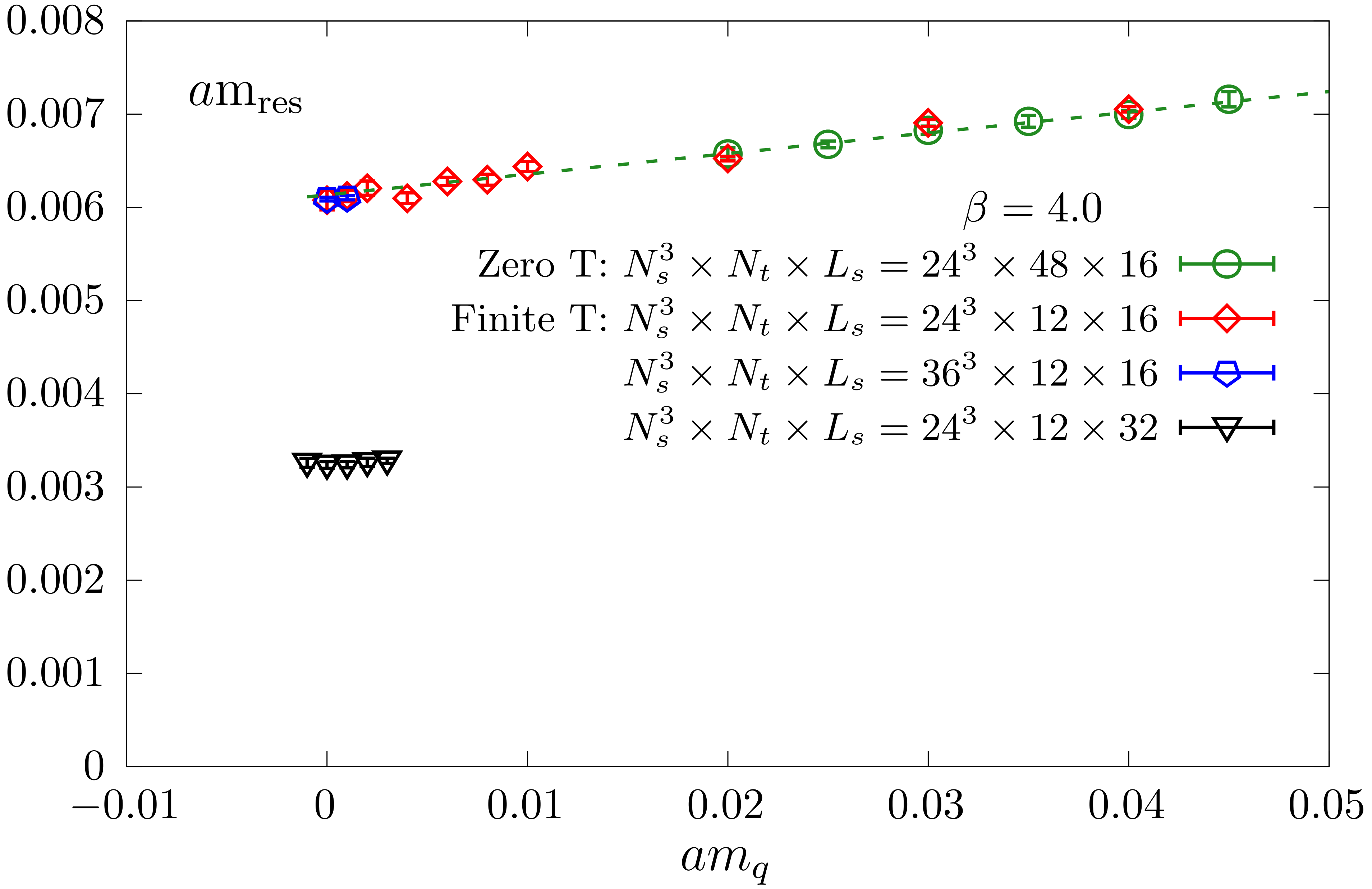
$$N_s = 24, -0.006 \le am_q \le 0.1$$

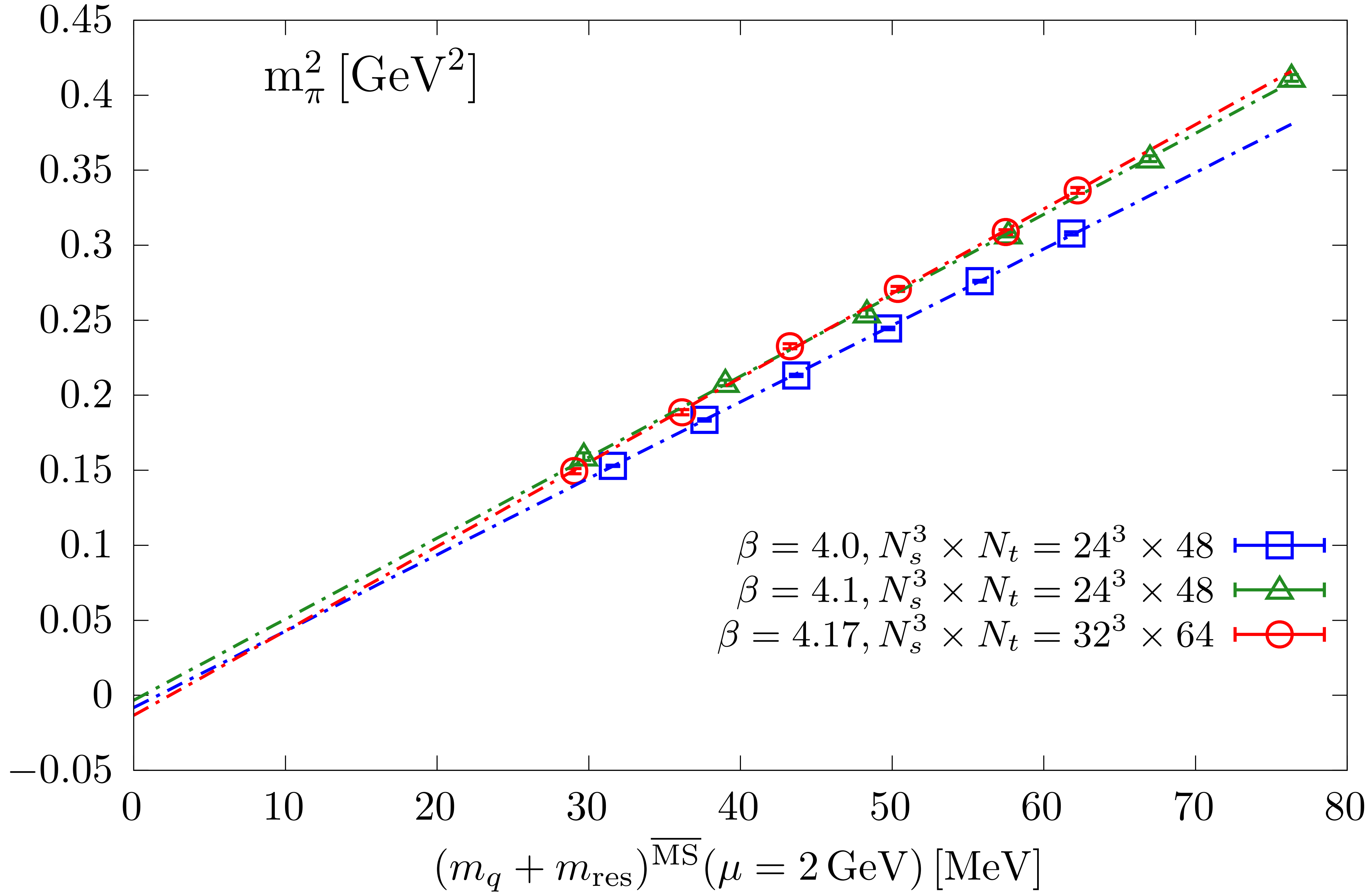
$$N_s^3 \times 12 \times 32: \quad N_s = 24, -0.001 \le am_q \le 0.003$$

Residual chiral symmetry breaking

$$m_{\rm res} = \frac{\left\langle \sum_{\vec{x}} J^a_{5q}(\vec{x},t) P^a(\vec{0},0) \right\rangle}{\left\langle \sum_{\vec{x}} P^a(\vec{x},t) P^a(\vec{0},0) \right\rangle} \Bigg|_{t \geq t_{min}}$$

- **For finite Ls chiral symmetry is broken, leading to an additive renormalization of the mass:** $m_q \to m_q + m_{\rm res}$





- m_{res} **has a linear dependence on** m_q **(lattice artifacts)**
  - **At** $am_q = 0$**:** $am_{\rm res} = 0.00613(9)$ **for** $L_s = 16$, $am_{\rm res} = 0.00324(3)$ **for** $L_s = 32$
- **At strong coupling,** m_{res} **dominated by gauge field dislocations, suppressed by** $1/L_s$
- m_π **almost vanishes at chiral limit**

Chiral condensate

$$\langle\bar{\psi}\psi\rangle|_{\rm DWF} \sim \langle\bar{\psi}\psi\rangle|_{\rm cont.} + C^D\frac{m_q + xm_{res}}{a^2} + \ldots$$

[S. Sharpe, arXiv: 0706.0218]

- $x = \mathcal{O}(1)$ but $x \neq 1$
- Additive divergence remains by $m = m_q + m_{res} \to 0$:

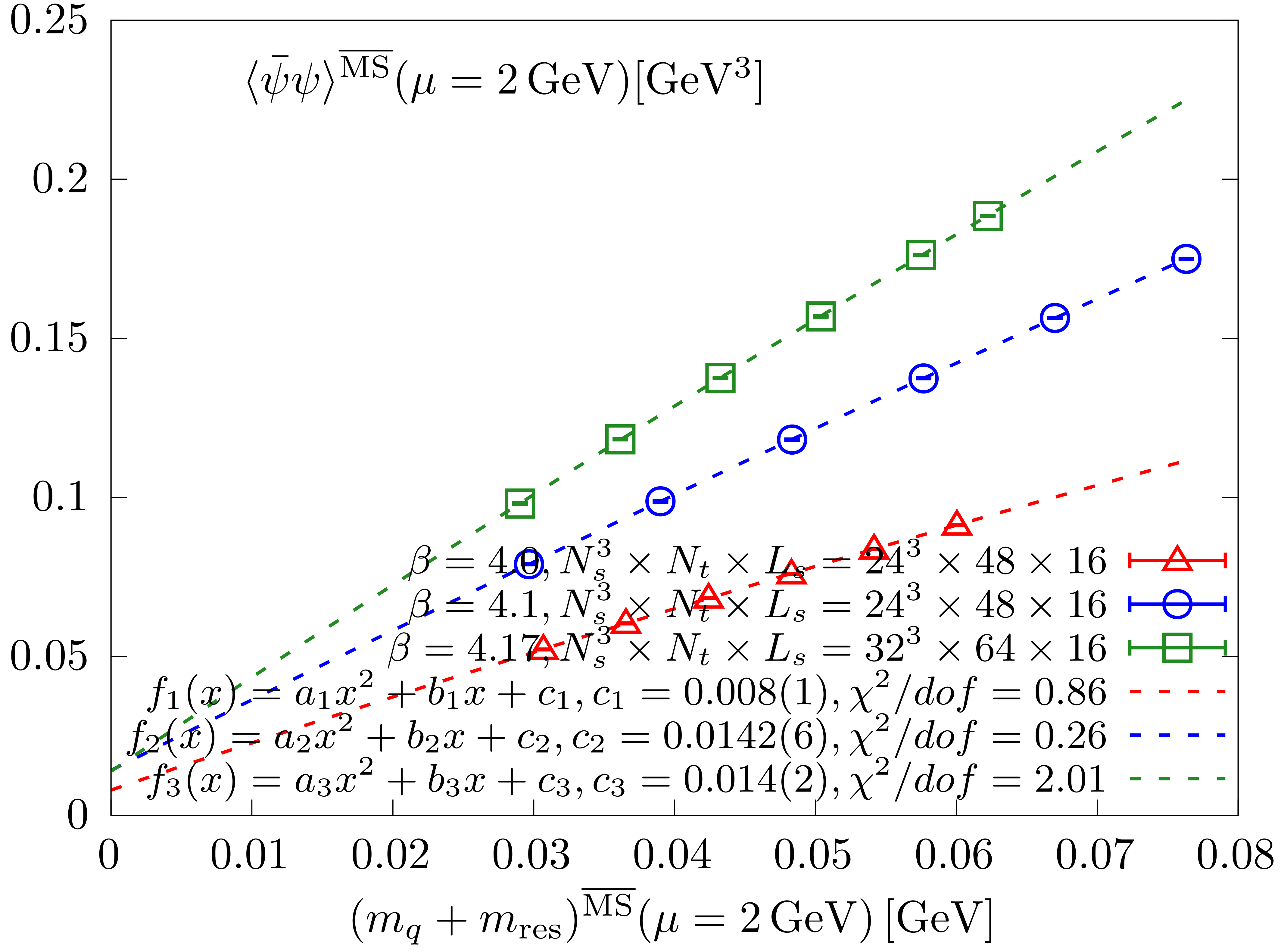
$$\lim_{m\to 0}\lim_{L\to 0}\langle\bar{\psi}\psi\rangle|_{DWF} \sim \langle\bar{\psi}\psi\rangle|_{\rm cont.} + C^D\frac{(x-1)m_{res}}{a^2}$$

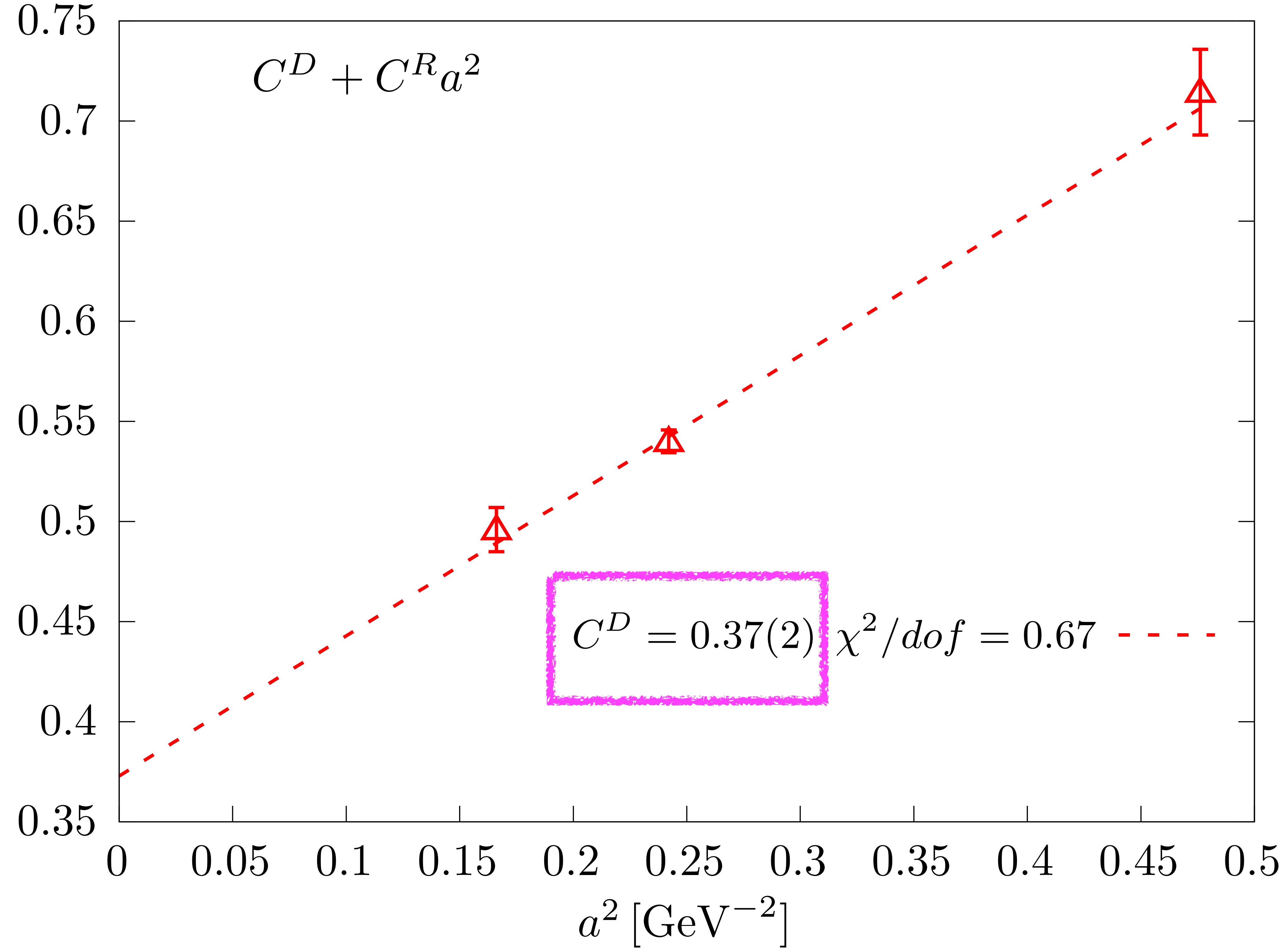
Two ways of subtracting divergences:

- $\langle\bar{\psi}\psi\rangle^{\rm ren} = Z_m^{-1}\left[\langle\bar{\psi}\psi\rangle^{\rm T>0} - \langle\bar{\psi}\psi\rangle^{\rm T=0}\right]$
- $\langle\bar{\psi}\psi\rangle^{\rm ren} = Z_m^{-1}\left[\langle\bar{\psi}\psi\rangle - C^D\frac{m_q + xm_{res}}{a^2}\right]$, **If we know C^D and x**

Zero T results

Calculate C^D for subtracting UV divergence term $C^D \frac{m_q + x m_{res}}{a^2}$



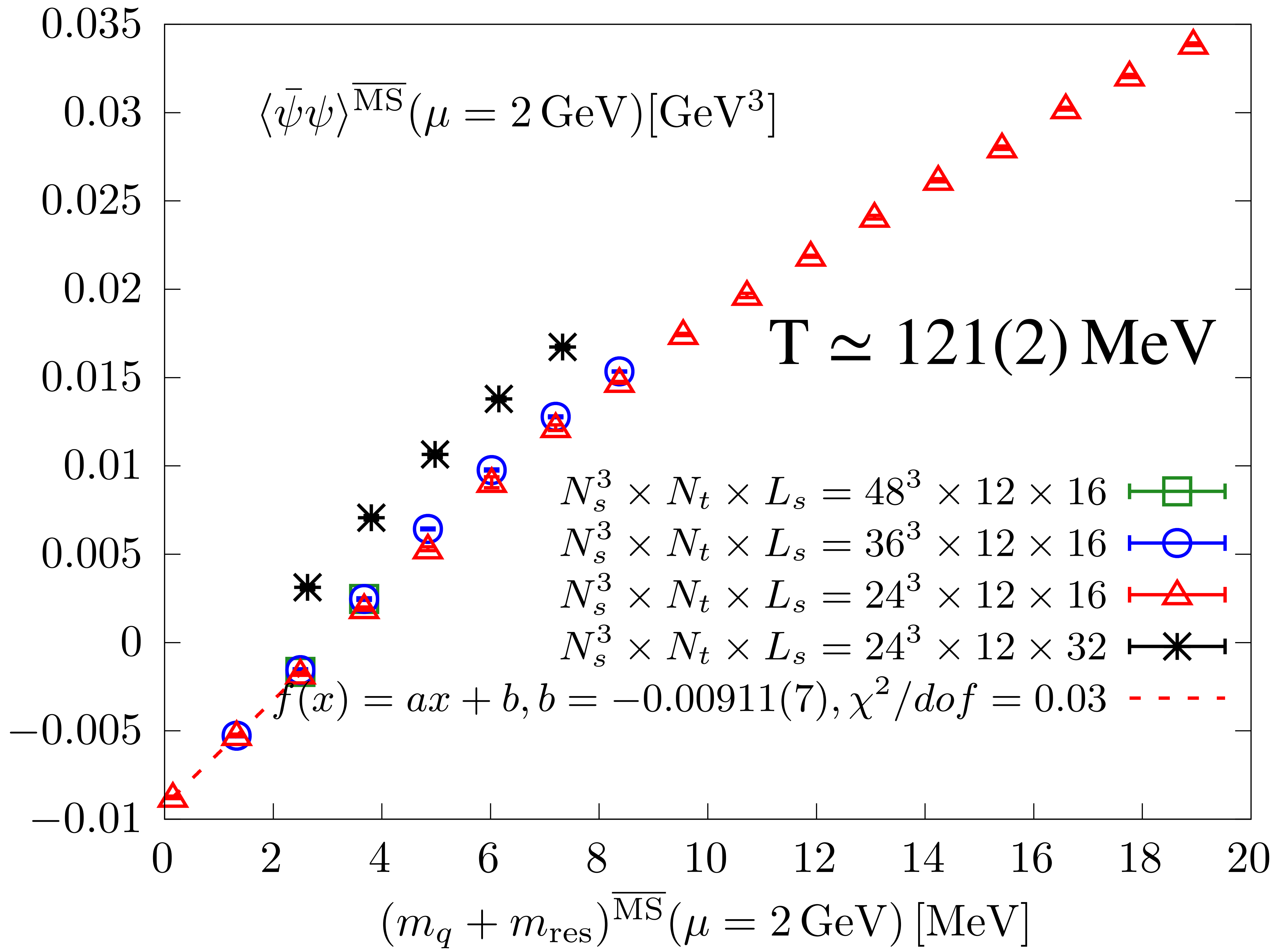


$$\langle\bar{\psi}\psi\rangle(m_q + m_{res}) = \langle\bar{\psi}\psi\rangle(0) + C^D \frac{m_q + x m_{res}}{a^2} + C^R(m_q + m_{res}) + A(m_q + m_{res})^2$$

$$= \langle\bar{\psi}\psi\rangle(0) + (C^D + C^R a^2)\frac{m_q + m_{res}}{a^2} + C^D \frac{(x-1)m_{res}}{a^2} + A(m_q + m_{res})^2$$

Finite T results

Calculate x for subtracting UV divergence term $C^D \dfrac{m_q + x m_{res}}{a^2}$



$$\lim_{(m_q+m_{res})\to 0} \langle\bar{\psi}\psi\rangle|_{DWF} \sim \langle\bar{\psi}\psi\rangle|_{cont.} + C^D \frac{(x-1)m_{res}}{a^2}$$

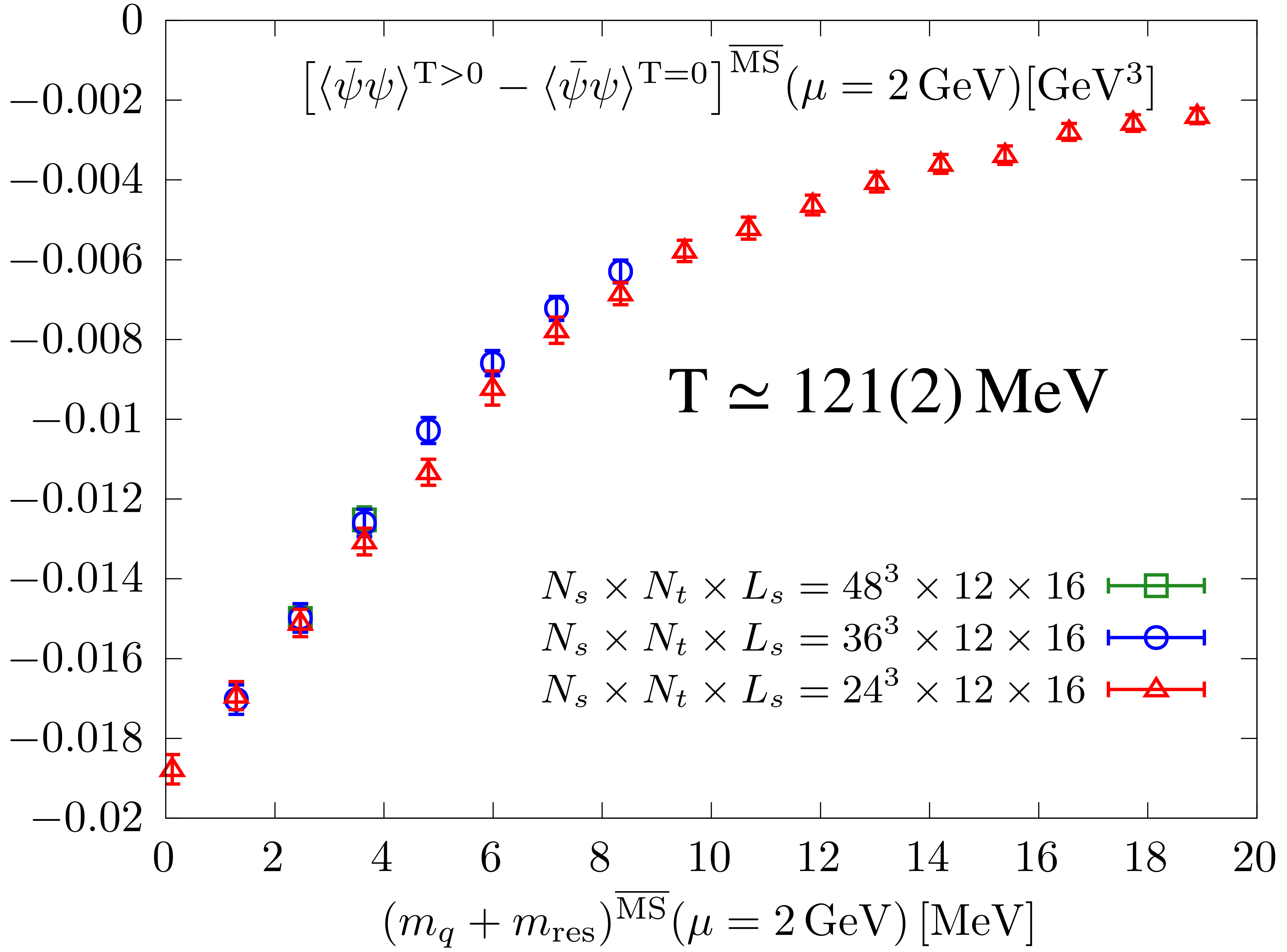
$x = -0.6(1)$ **for T>Tc**

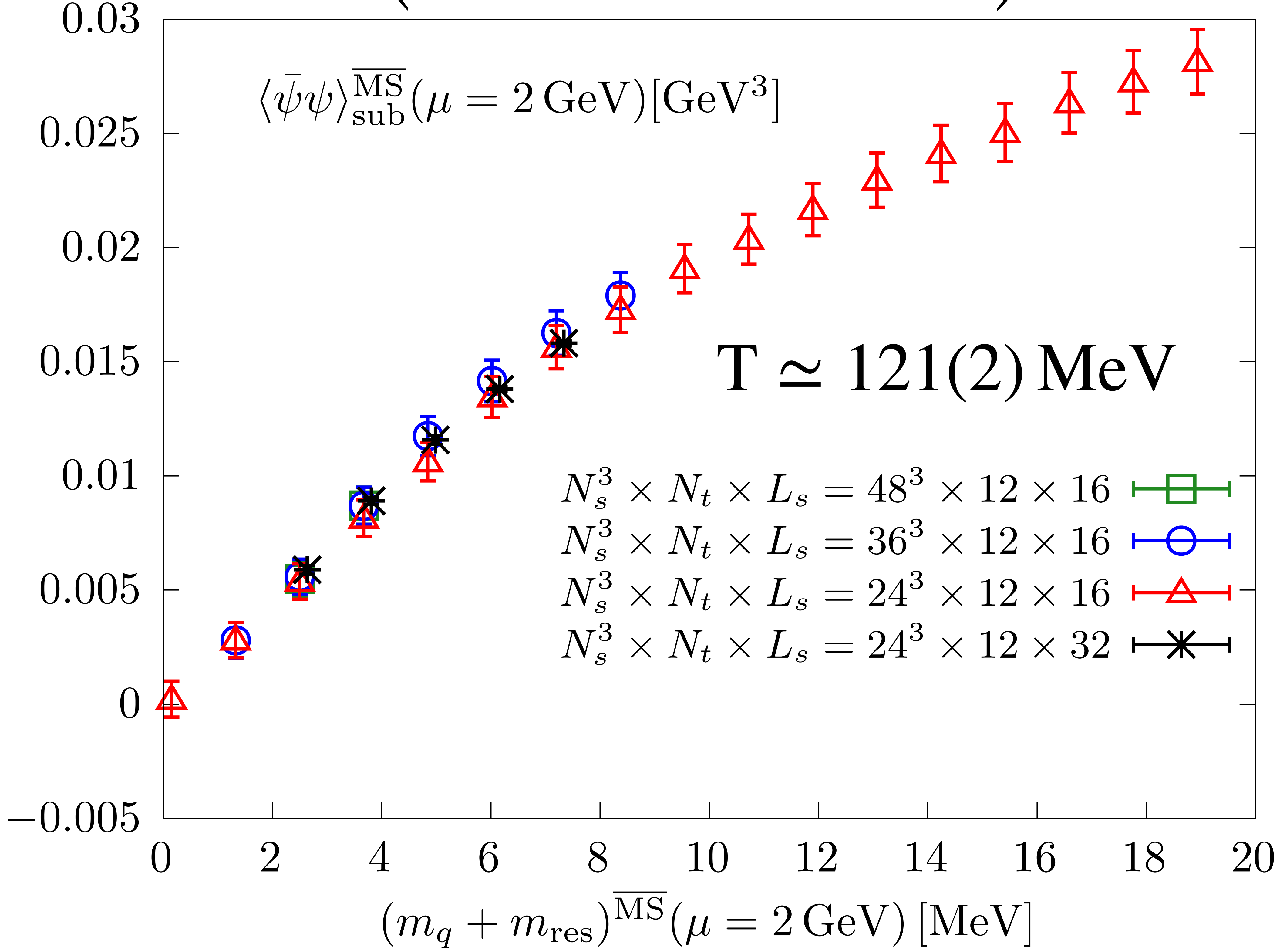
$T_c = 98^{+3}_{-6}\,\mathrm{MeV}$

[S. Sharma et al. PRD 2022]

Renormalized chiral condensate

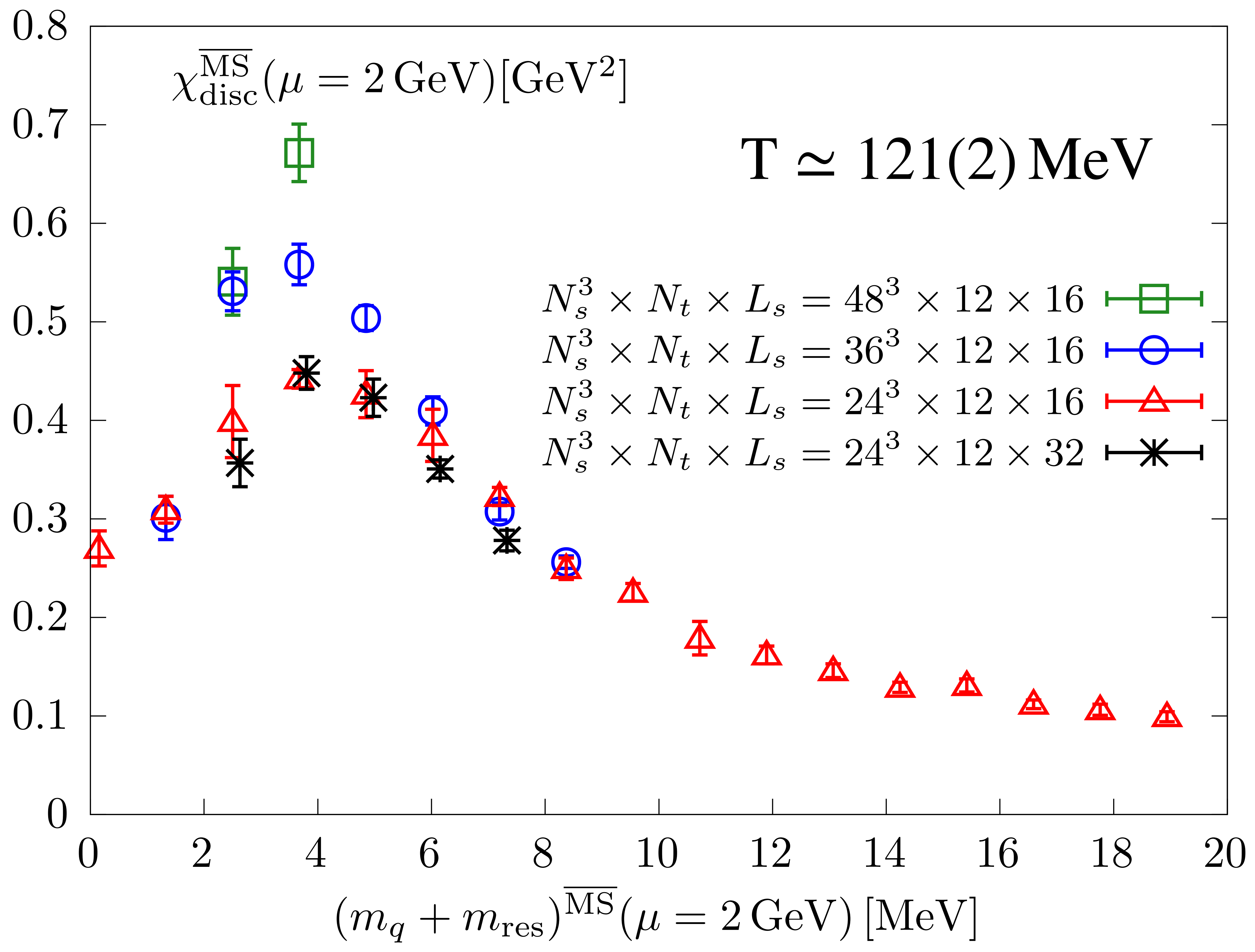
$$\langle\bar{\psi}\psi\rangle_{\rm sub}^{\overline{\rm MS}}(2\,{\rm GeV}) = \left(\langle\bar{\psi}\psi\rangle - C^D\frac{m_q + x m_{res}}{a^2}\right)/Z_m^{\overline{\rm MS}}(2\,{\rm GeV})$$





- **Subtracted chiral condensate vanishes in the chiral limit**
- **Finite volume effect is visible at low T**
- **Subtracted chiral condensate remains the same with same total quark mass for varying Ls**

Disconnected chiral susceptibility



- **Large finite volume effect near the transition point, but the change in peak height is not as large as anticipated from a real phase transition**

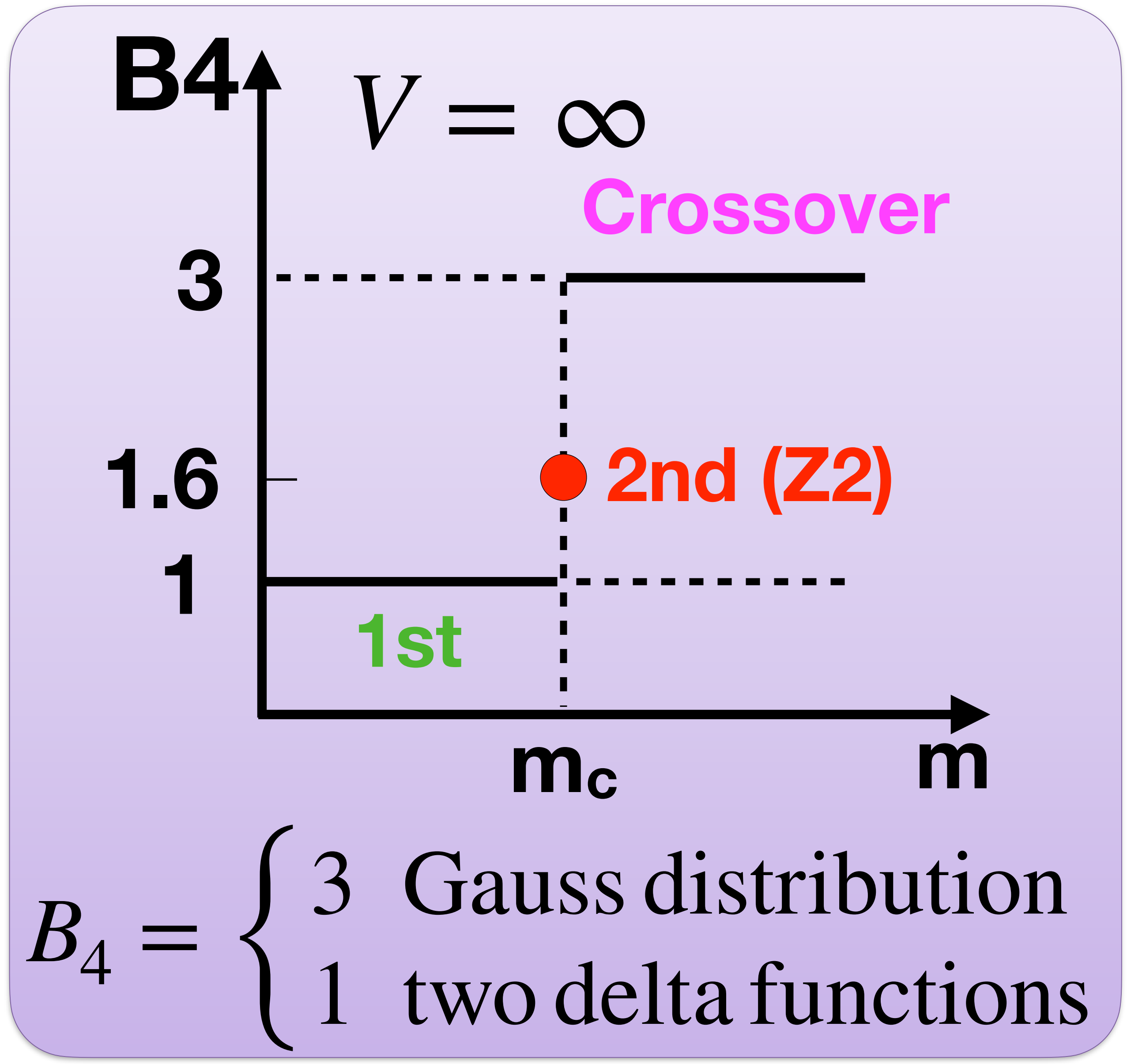
➡ **Consistent with the crossover transition**

- **The transition mass point is around 3.6 MeV**

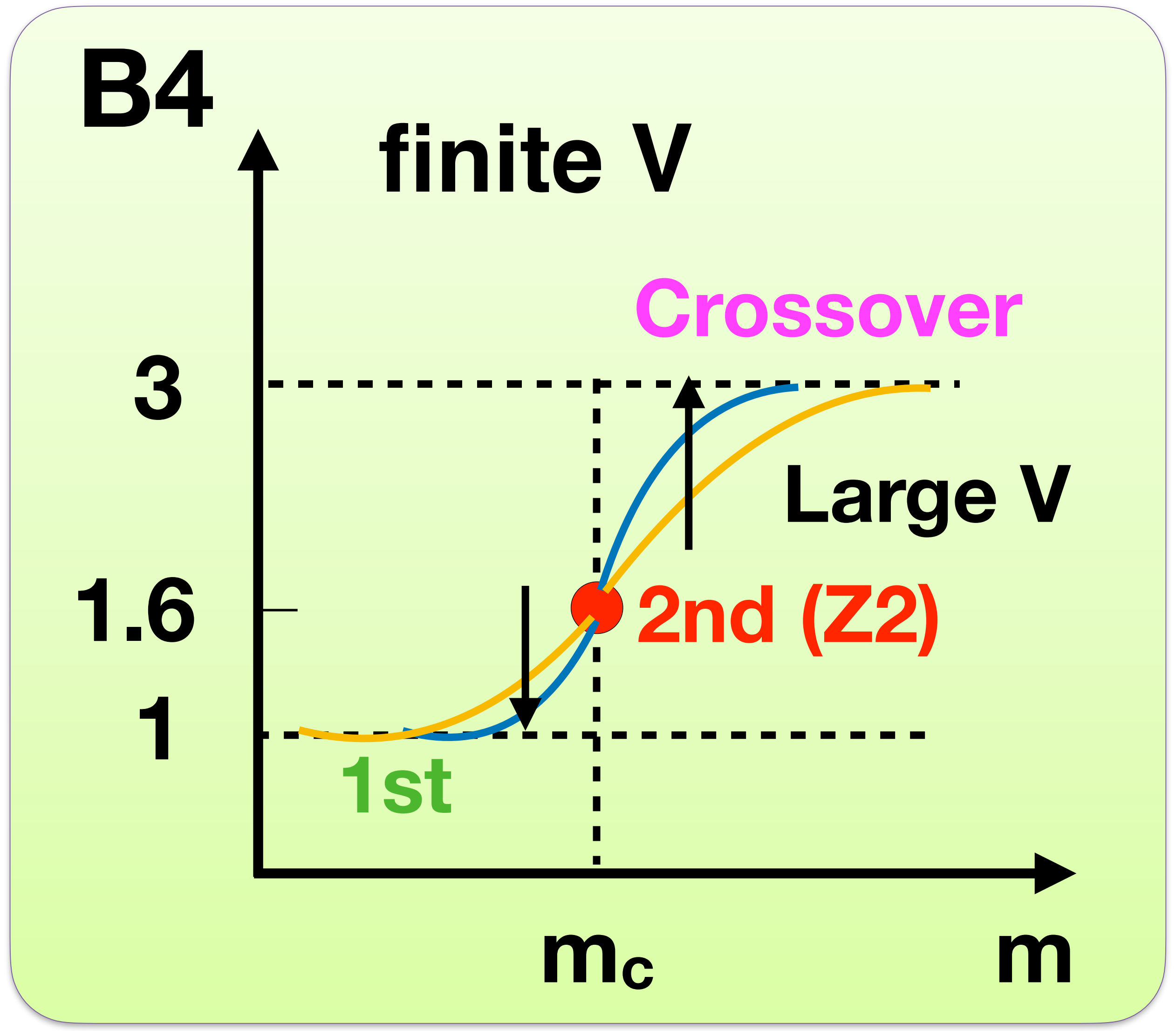
(FLAG Review '21: $m_{ud}^{\overline{\rm MS}}(2\,{\rm GeV}) = 3.381(40)\,{\rm MeV}$)

- χ_{disc} **seems to be function of total quark mass**

Methodology to determine the order of transition

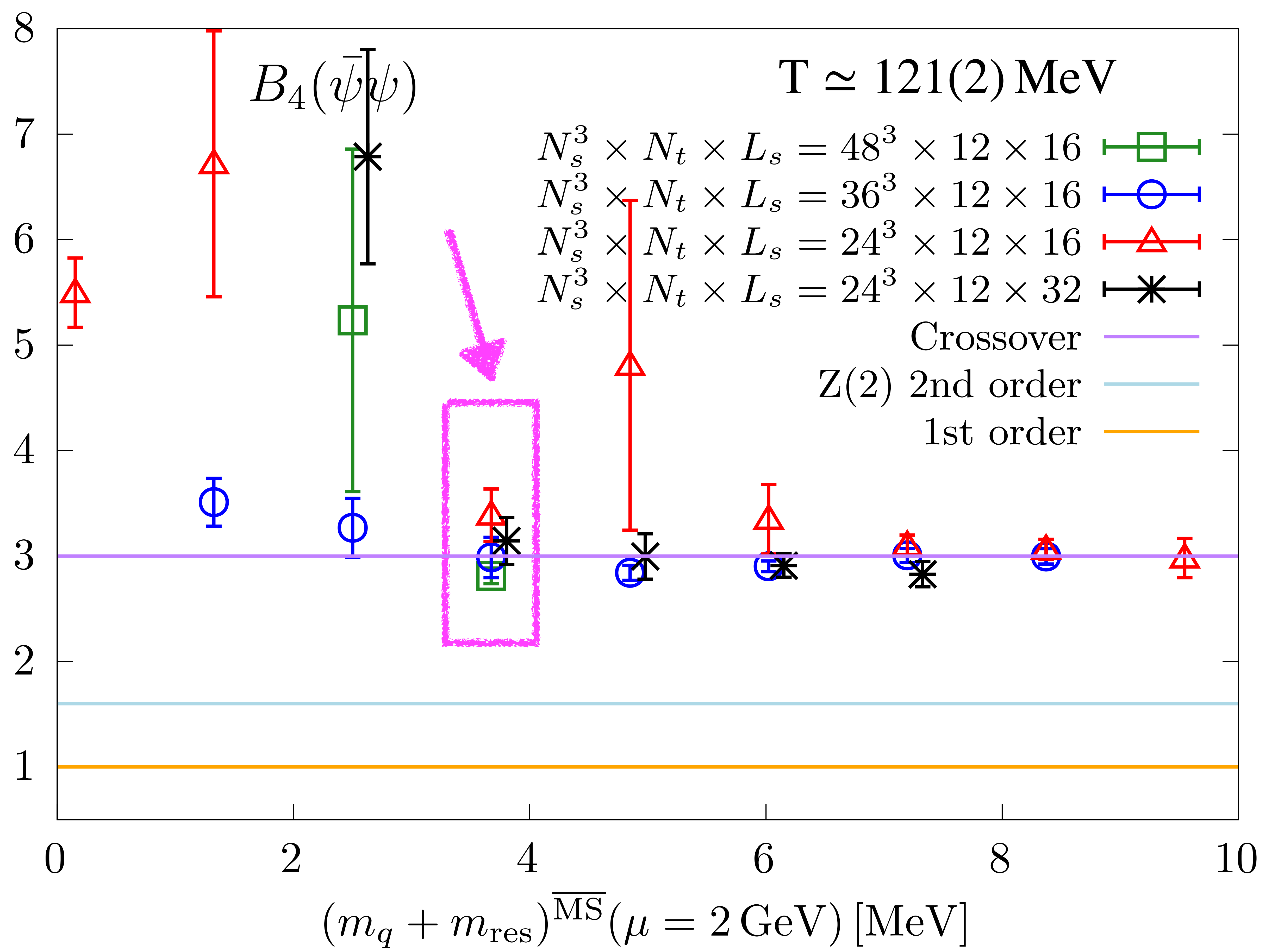


A.Kiyohara et al., PRD 104, 114509 (2021)



Binder Cumulant: $B_4(\bar{\psi}\psi) = \dfrac{\langle(\delta\bar{\psi}\psi)^4\rangle}{\langle(\delta\bar{\psi}\psi)^2\rangle^2}, \quad \delta\bar{\psi}\psi = \bar{\psi}\psi - \langle\bar{\psi}\psi\rangle$

Suggests a crossover transition at $(m_q + m_{\text{res}})^{\overline{\text{MS}}}(2\,\text{GeV}) \simeq$ 3.6 MeV



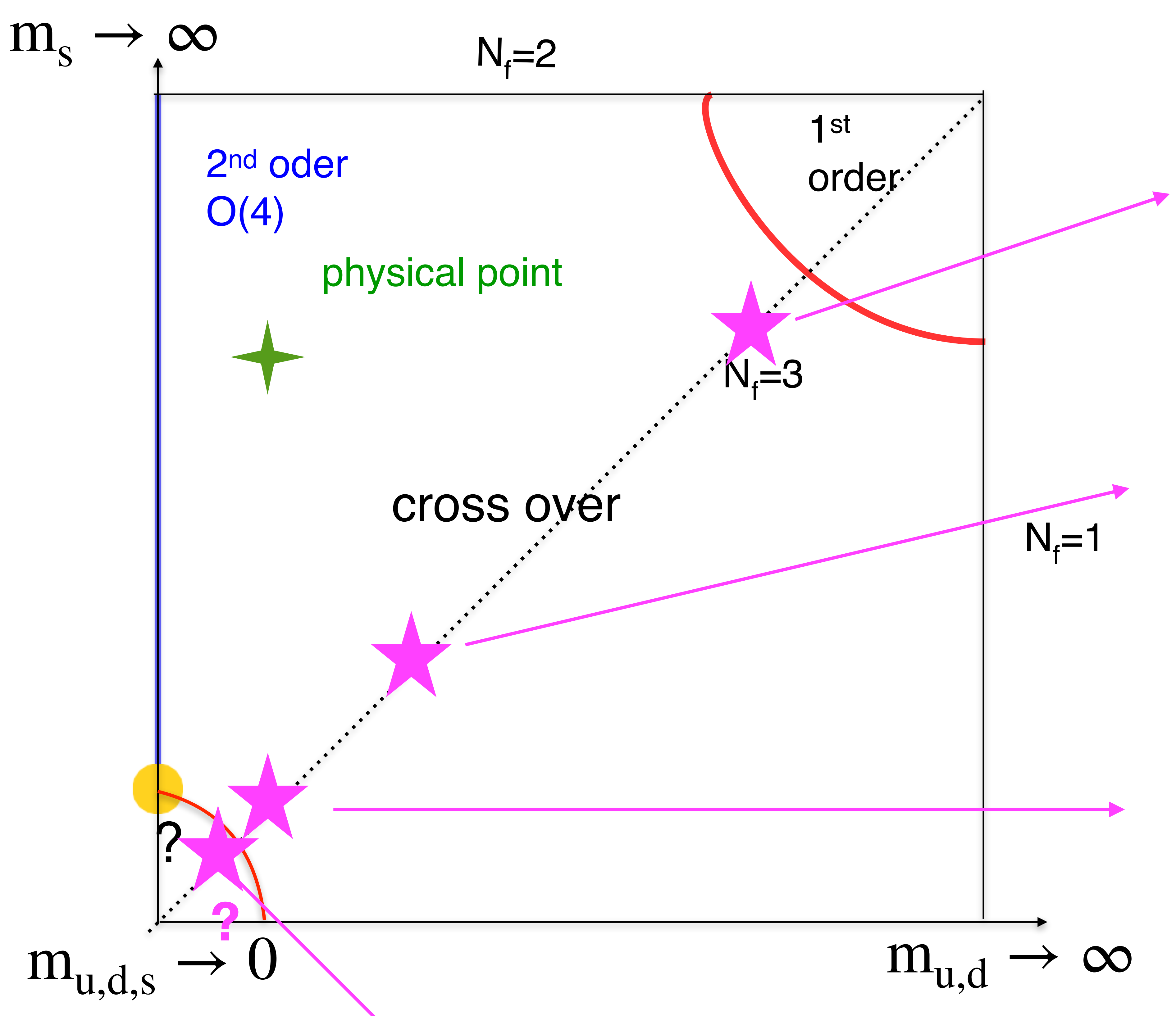
Histogram of chiral condensate near transition point

Behaves like a Gaussian distribution, no evidence of a double peak structure would appear as $V \uparrow$

Summary and outlook

No evidence of a 1st order PT in our explored quark mass range

If a 1st order region exist, the critical mass should be less than 3.6 MeV



- $T \sim 242\,\mathrm{MeV}(N_t = 6), (m_q + m_{\mathrm{res}})^{\overline{\mathrm{MS}}}_{pc} \sim 183\,\mathrm{MeV}$
  $\Leftrightarrow m_\pi^{pc} \sim 997\,\mathrm{MeV}$, crossover transition

  Y. Nakamura, Y. Zhang et al., *PoS* LATTICE2021

- $T \sim 181\,\mathrm{MeV}(N_t = 8), (m_q + m_{\mathrm{res}})^{\overline{\mathrm{MS}}}_{pc} \sim 42\,\mathrm{MeV}$
  $\Leftrightarrow m_\pi^{pc} \sim 476\,\mathrm{MeV}$, crossover transition

  Y. Zhang et al., *PoS* LATTICE2022

- $T \sim 121\,\mathrm{MeV}(N_t = 12), (m_q + m_{\mathrm{res}})^{\overline{\mathrm{MS}}}_{pc} \sim 3.6\,\mathrm{MeV}$
  $\Leftrightarrow m_\pi^{pc} \sim 141\,\mathrm{MeV}$, crossover transition

  Y. Zhang et al., *PoS* LATTICE2023

- $T \sim 104\,\mathrm{MeV}\,(N_t = 14)$, lighter quark mass simulation is underway

Acknowledgements

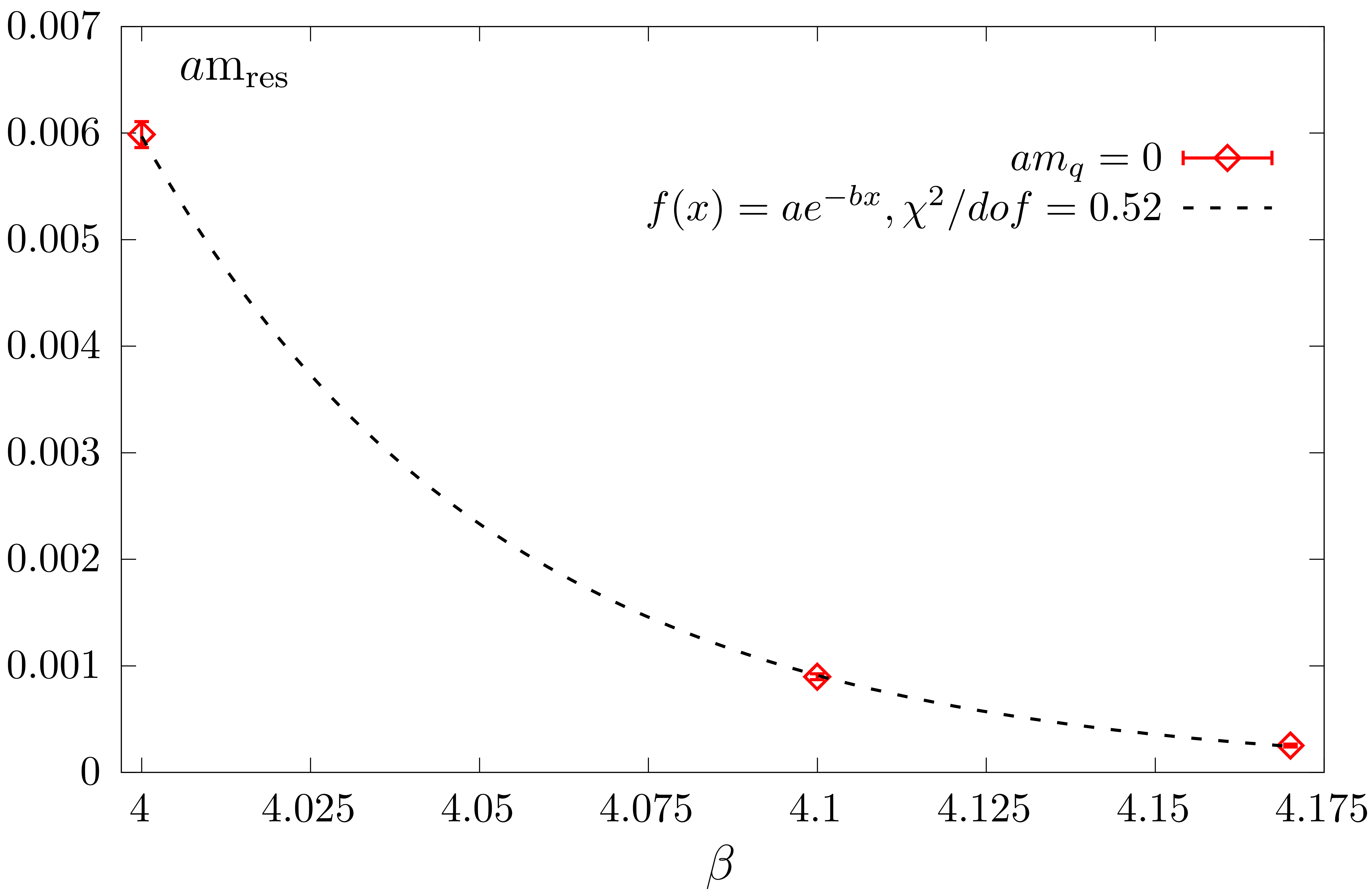
- Codes
  - HMC
    - Grid (implementation for A64FX: thanks to the Regensburg group)
  - Measurements
    - Bridge++
    - Hadrons / Grid

- Computers
  - Supercomputer Fugaku provided by the RIKEN Center for Computational Science through HPCI project #hp210032 and Usability Research ra000001.
  - Wisteria/BDEC-01 Oddysey at Univ. Tokyo/JCAHPC through HPCI project #hp220108 Ito supercomputer at Kyushu University through HPCI project #hp190124 and hp200050
  - Hokusai BigWaterfall at RIKEN

- Grants
  - JSPS Kakenhi (20H01907)

Backup slide

Residual chiral symmetry breaking



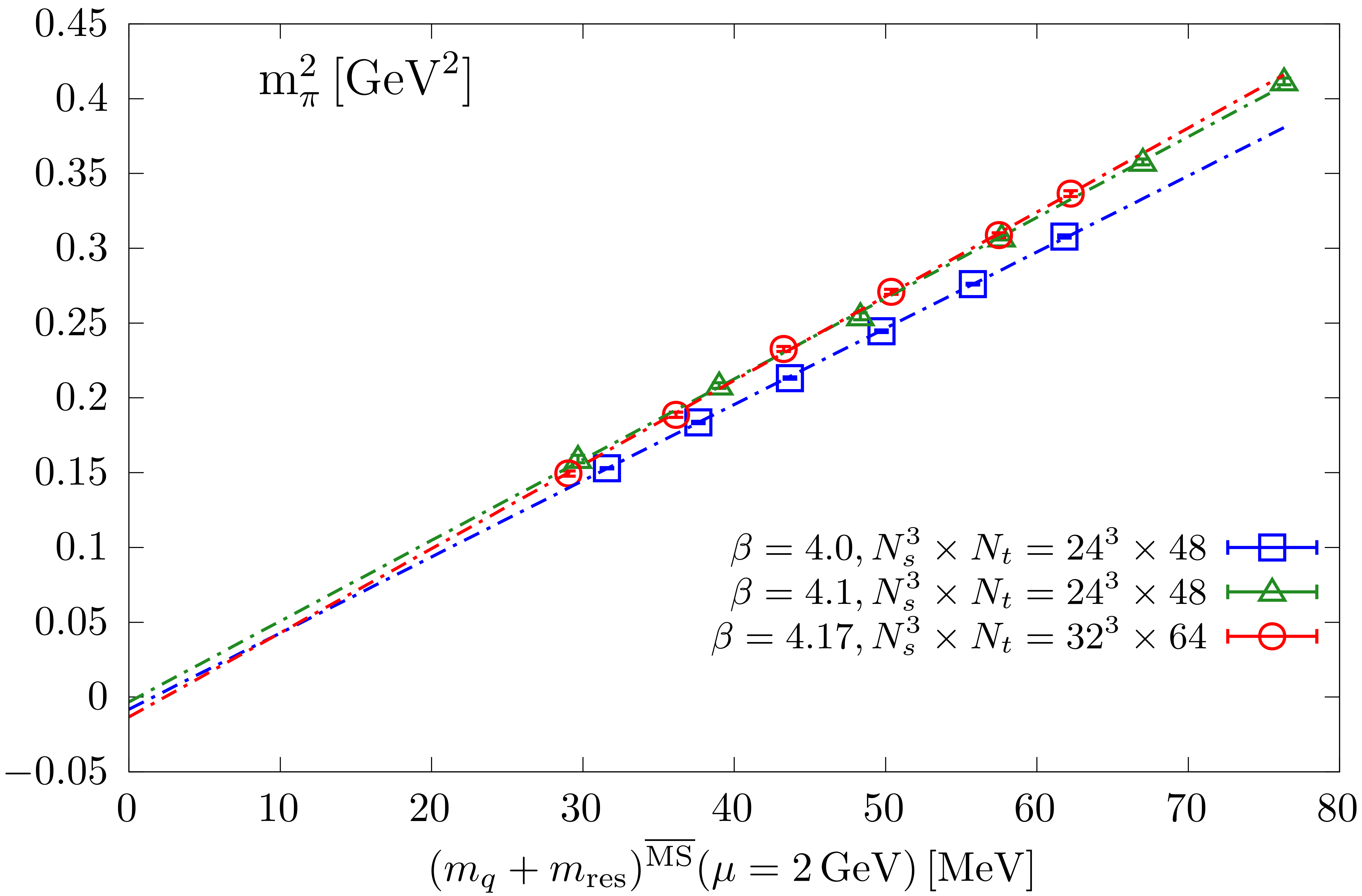
Coupling dependence: m_{res} decreases exponentially with increasing β

Pion mass

- **At leading order in chiral perturbation theory**
  $m_\pi^2 \propto m_q + m_{res}$

- **Evidence of good linearity**

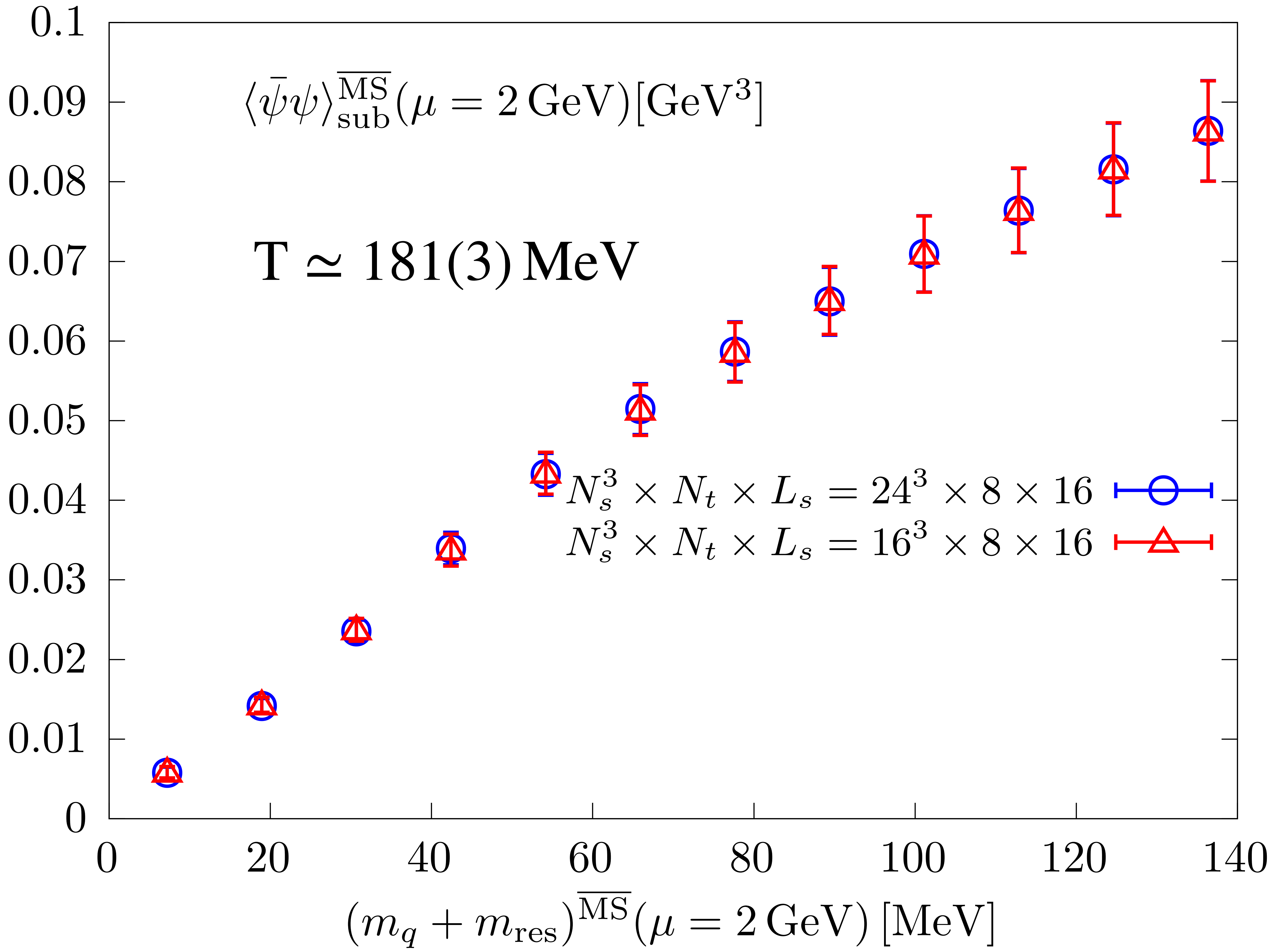
- m_π **close to zero, but not exact**
zero at chiral limit:
  due to not considering chiral logarithm term & finite volume effects?



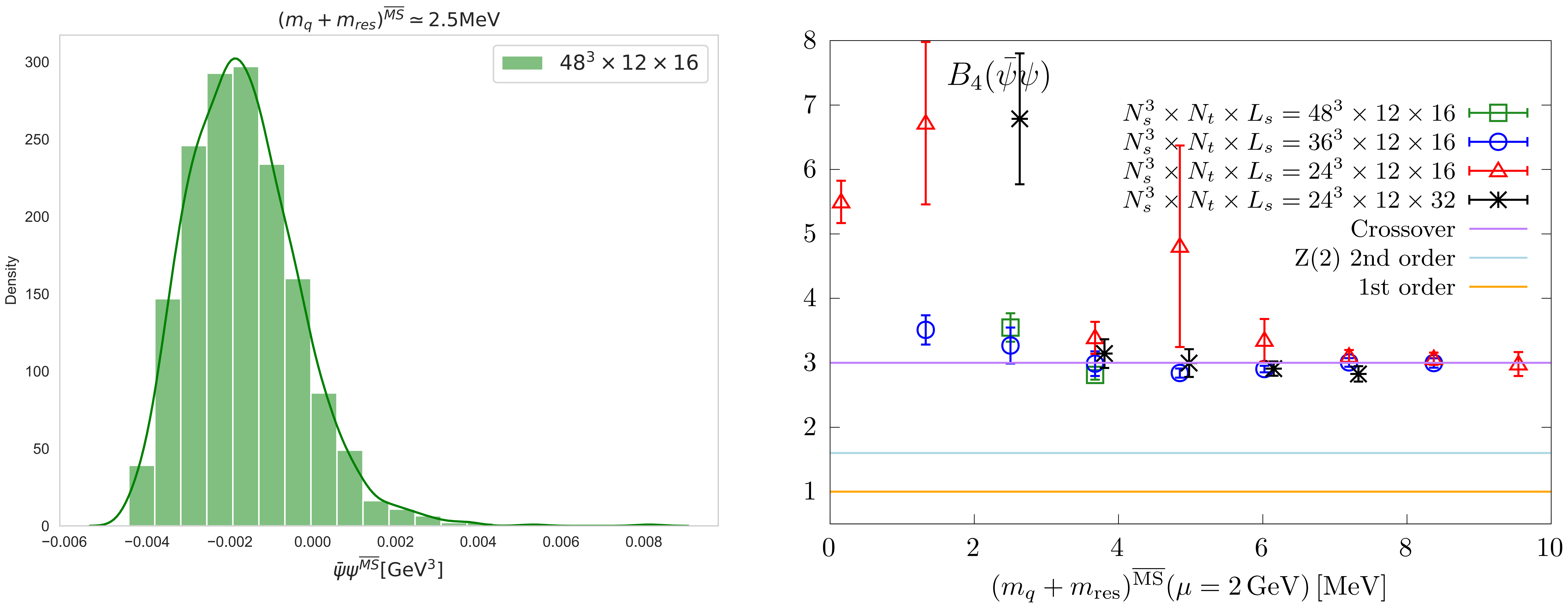
Good test of chirality

Renormalized chiral condensate

$$\langle\bar{\psi}\psi\rangle_{\rm sub}^{\overline{\rm MS}}(2\,{\rm GeV}) = \left(\langle\bar{\psi}\psi\rangle - C^D\frac{m_q + x m_{res}}{a^2}\right)/Z_m^{\overline{\rm MS}}(2\,{\rm GeV})$$



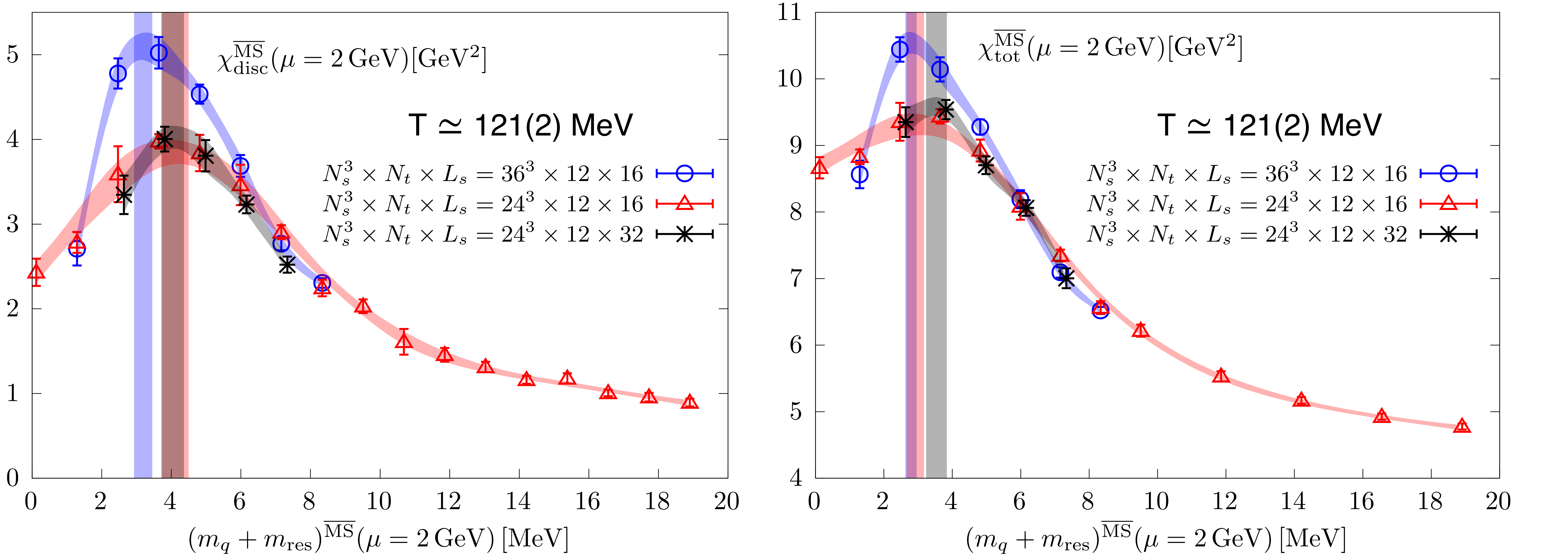
- **Subtracted chiral condensate vanishes in the chiral limit**
- **No volume dependence**



$(m_q + m_{res})^{\overline{MS}} \simeq 2.5$MeV
$48^3 \times 12 \times 16$
Density
$\bar{\psi}\psi^{\overline{MS}}$[GeV^3]
$B_4(\bar{\psi}\psi)$
$N_s^3 \times N_t \times L_s = 48^3 \times 12 \times 16$
$N_s^3 \times N_t \times L_s = 36^3 \times 12 \times 16$
$N_s^3 \times N_t \times L_s = 24^3 \times 12 \times 16$
$N_s^3 \times N_t \times L_s = 24^3 \times 12 \times 32$
Crossover
Z(2) 2nd order
1st order
$(m_q + m_{\rm res})^{\overline{\rm MS}}(\mu = 2\,{\rm GeV})$ [MeV]

Chiral susceptibility at T~121 MeV

$$\chi_{\rm tot} = \frac{T}{V}\frac{\partial^2 \ln Z}{\partial m^2} = \chi_{\rm disc} + \chi_{\rm con},\ \ \chi_{\rm disc} = \frac{N_f^2}{N_s^3 N_t}\left[\left\langle \left(\mathrm{Tr}\, M^{-1}\right)^2\right\rangle - \left\langle \mathrm{Tr}\, M^{-1}\right\rangle^2\right],\ \ \chi_{\rm con} = -\frac{N_f}{N_s^3 N_t}\langle \mathrm{Tr}\, M^{-2}\rangle$$



Finite size scaling: susceptibility at T~121 MeV

- Crossover: $\chi_\sigma^{\max}(N_s, N_t)$ independent of V
- 1st order PT: $\chi_\sigma^{\max}(N_s, N_t) \propto V$
- Z(2) 2nd order PT: a singular behavior should be observed in $\chi_\sigma^{\max}(N_s, N_t)$ with V $((N_s^3 \times N_t)^\alpha,\ \alpha = 1.966$ is the critical exponent)

The peak height of χ_σ does not scale like a 1st or Z(2) second order PT

finite-size scaled renormalized chiral susceptibility