

# The Equation of State of QCD up to the Electro-Weak scale

## part 2

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# Overview

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1. Thermal QCD in a moving frame
2. Strategy
3. Numerical results
4. Conclusions & outlook

# Thermal QCD in a moving frame

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- QCD in a box  $L_0 \times L^3$  in a moving frame:

$$Z(L_0, L, \boldsymbol{\xi}) = \text{Tr} \left\{ e^{-L_0(\hat{H} - i\boldsymbol{\xi} \cdot \hat{\mathbf{P}})} \right\}$$

$\boldsymbol{\xi} \rightarrow$  euclidean boost parameter, 3d vector

[Giusti, Meyer, JHEP 01 (2013), 140]

- Shifted boundary conditions

$$A_\mu(x_0 + L_0, \mathbf{x}) = A_\mu(x_0, \mathbf{x} - L_0\boldsymbol{\xi})$$

$$\psi(x_0 + L_0, \mathbf{x}) = -\psi(x_0, \mathbf{x} - L_0\boldsymbol{\xi})$$

$$\bar{\psi}(x_0 + L_0, \mathbf{x}) = -\bar{\psi}(x_0, \mathbf{x} - L_0\boldsymbol{\xi})$$

- Temperature range 1 – 100 GeV: Scale setting through a NP-renormalized coupling  $\bar{g}(\mu)$

[Leonardo Giusti's talk]

# Strategy

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Definition of free-energy density:

$$f(L_0, L, \boldsymbol{\xi}) = -\frac{1}{L_0 L^3} \ln Z(L_0, L, \boldsymbol{\xi})$$

Computation of the entropy density:

$$\frac{s}{T^3} = \frac{(1 + \boldsymbol{\xi}^2)}{\xi_k} \frac{1}{T^4} \frac{\partial f_{\boldsymbol{\xi}}}{\partial \xi_k}, \quad k = 1, 2, 3$$

[Dalla Brida, Giusti, Pepe, JHEP **04** (2020), 043]

Split in two contributions:

$$\frac{s}{T^3} = \frac{(1 + \boldsymbol{\xi}^2)}{\xi_k} \frac{1}{T^4} \left\{ \frac{\partial}{\partial \xi_k} (f_{\boldsymbol{\xi}} - f_{\boldsymbol{\xi}}^{\infty}) + \frac{\partial}{\partial \xi_k} f_{\boldsymbol{\xi}}^{\infty} \right\}$$

where  $f_{\boldsymbol{\xi}}^{\infty}$  free-energy of shifted QCD with quarks at infinite mass

# Mass integral

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$$\frac{s}{T^3} = \frac{(1 + \boldsymbol{\xi}^2)}{\xi_k} \frac{1}{T^4} \left\{ \frac{\partial}{\partial \xi_k} (f_{\boldsymbol{\xi}} - f_{\boldsymbol{\xi}^\infty}) + \frac{\partial}{\partial \xi_k} f_{\boldsymbol{\xi}^\infty} \right\}$$

Write the difference as an integral in subtracted mass  $m_q = m_0 - m_{cr}$ :

$$\frac{\partial}{\partial \xi_k} (f_{\boldsymbol{\xi}} - f_{\boldsymbol{\xi}^\infty}) = -\frac{\partial}{\partial \xi_k} \int_0^\infty dm_q \frac{\partial f_{\boldsymbol{\xi}}}{\partial m_q} = -\int_0^\infty dm_q \frac{\partial}{\partial \xi_k} \langle \bar{\psi} \psi \rangle_{\boldsymbol{\xi}}$$

On the lattice at given  $L_0/a$  and  $g_0$ :

$$\frac{\Delta}{\Delta \xi_k} (f_{\boldsymbol{\xi}} - f_{\boldsymbol{\xi}^\infty}) = -\int_0^\infty dm_q \frac{1}{\xi_+ - \xi_-} \left[ \langle \bar{\psi} \psi \rangle_{\boldsymbol{\xi}_+} - \langle \bar{\psi} \psi \rangle_{\boldsymbol{\xi}_-} \right]$$

We choose  $\boldsymbol{\xi} = (1, 0, 0)$ , and  $\boldsymbol{\xi}_\pm = (1 \pm \frac{a}{L_0}, 0, 0)$

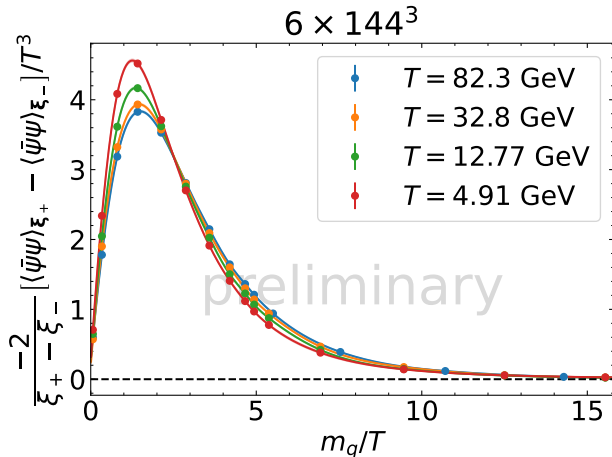
[Giusti, Pepe, Phys. Rev. D **91** (2015), 114504]

# Mass integral

- $N_f = 3$   
 $\mathcal{O}(a)$ -improved Wilson fermions
- Gauss quadratures:

interval	points
$0 \leq m_q/T \leq 5$	10
$5 \leq m_q/T \leq 20$	7
$m_q/T \geq 20$	3 (in $\kappa$ )

- Relative error on integral  
 $\sim 0.5\%$



## $g_0^2$ integral

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$$\frac{s}{T^3} = \frac{(1 + \xi^2)}{\xi_k} \frac{1}{T^4} \left\{ \frac{\partial}{\partial \xi_k} (f_\xi - f_\xi^\infty) + \frac{\partial}{\partial \xi_k} f_\xi^\infty \right\}$$

where  $f_\xi^\infty = f_\xi^{\text{YM}}$ . We sample at many bare couplings

$$\frac{d}{dg_0^2} \left( \frac{\partial}{\partial \xi_k} f_\xi^{\text{YM}} \right) = -\frac{1}{g_0^2} \frac{\partial}{\partial \xi_k} \langle S^G \rangle_\xi$$

and then integrate up to the desired value of  $g_0^2$ :

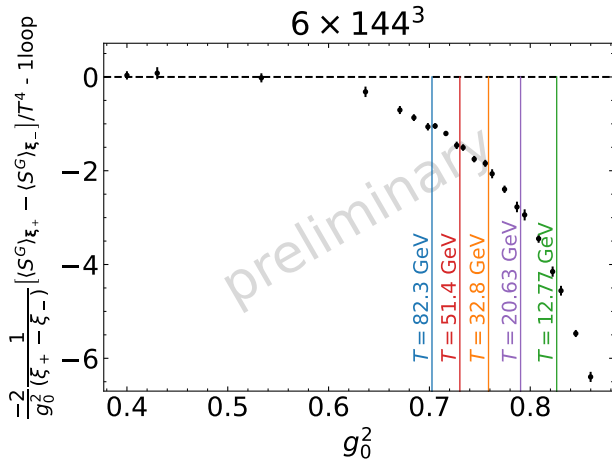
$$\frac{\partial}{\partial \xi_k} f_\xi^\infty = \left( \frac{\partial}{\partial \xi_k} f_\xi^{\text{YM}} \right)_{\text{free}} - \int_0^{g_0^2} du \frac{1}{u} \frac{\partial \langle S^G \rangle_\xi}{\partial \xi_k} \Big|_u$$

# $g_0^2$ integral

- Integration rules:

interval	points
$0 \leq g_0^2 \leq 0.4$	1 (trapz)
$0.4 \leq g_0^2 \leq 0.67$	3 (Gauss)
$T_i \rightarrow T_{i+1}$	3 (Gauss)

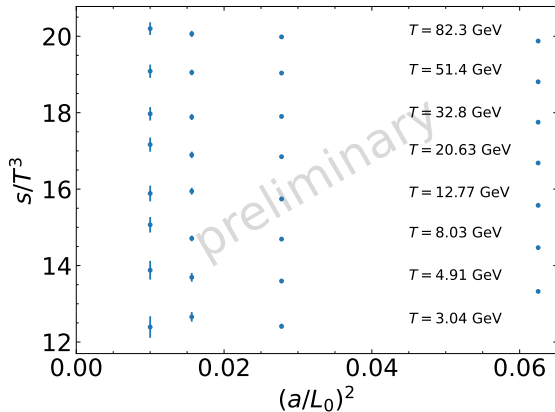
- Relative error on integral  
 $\sim 0.5\%$





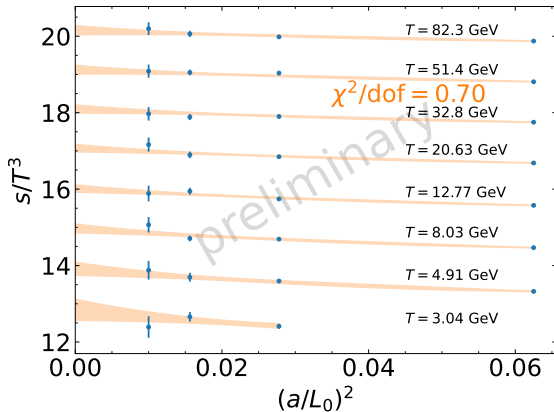
# Continuum limit

- $L_0/a = 4, 6, 8, 10$  and  $L/a = 144$
- 1loop improvement



# Continuum limit

- $L_0/a = 4, 6, 8, 10$  and  $L/a = 144$
- 1loop improvement
- Global fit of cutoff effects

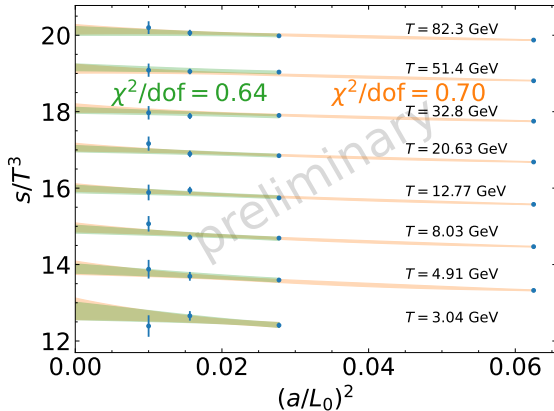


$$s(g_i, a/L_0)/T^3 = s_i + \left(\frac{a}{L_0}\right)^2 (s_{23} g_i^3 + s_{24} g_i^4) + \left(\frac{a}{L_0}\right)^3 (s_{33} g_i^3 + s_{34} g_i^4)$$

where  $g_i = g(T_i)$

# Continuum limit

- $L_0/a = 4, 6, 8, 10$  and  $L/a = 144$
- 1loop improvement
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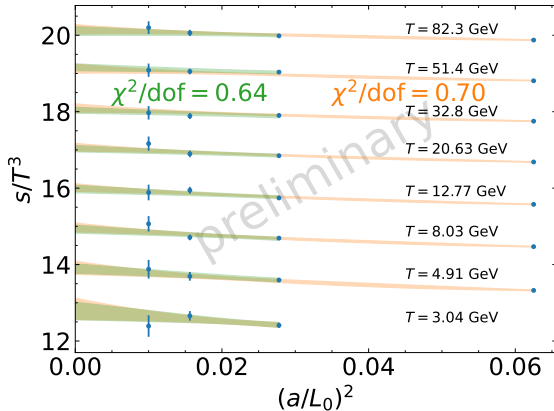


$$s(g_i, a/L_0)/T^3 = s_i + \left(\frac{a}{L_0}\right)^2 (s_{23} g_i^3 + s_{24} g_i^4 + s_{25} g_i^5)$$

where  $g_i = g(T_i)$

# Continuum limit

- $L_0/a = 4, 6, 8, 10$  and  $L/a = 144$
- 1loop improvement
- Global fit of cutoff effects
- No impact from  $\gamma \in [-1, 1]$
- Final accuracy  $\lesssim 1\%$



$$s(g_i, a/L_0)/T^3 = s_i + [\alpha(a^{-1})]^\gamma \left(\frac{a}{L_0}\right)^2 (s_{23} g_i^3 + s_{24} g_i^4 + s_{25} g_i^5)$$

where  $g_i = g(T_i)$

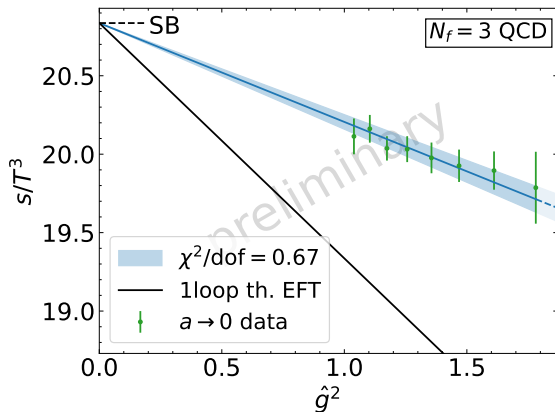
# $N_f = 3$ QCD entropy density at high temperature

- $\hat{g}(T) \equiv g_{\overline{MS}}^{5\text{loop}}(2\pi T)$

- Fit of continuum values

$$s/T^3 = s_0 + s_1 \hat{g}^2$$

- $s_0$  compatible with SB
- Fit with  $s_0 = s_{SB}/T^3$



# Conclusions & outlook

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- QGP thermodynamics at high temperature made accessible by thermal QCD in a moving frame
- First computation of the entropy density of  $N_f = 3$  QCD in the temperature range 1 – 100 GeV
- Continuum extrapolations with 4 lattice spacings, final accuracy  $\sim 1\%$

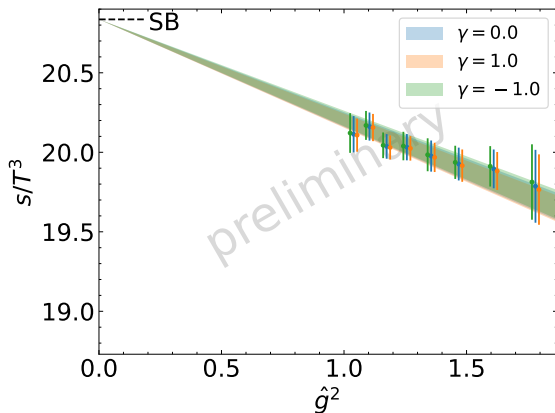
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Non-perturbative renormalization of lattice QCD Energy-Momentum tensor

$$s/T^3 = -\frac{1 + \xi^2}{\xi_k} \langle T_{0k}^{R,\{6\}} \rangle_{\xi} / T^4$$

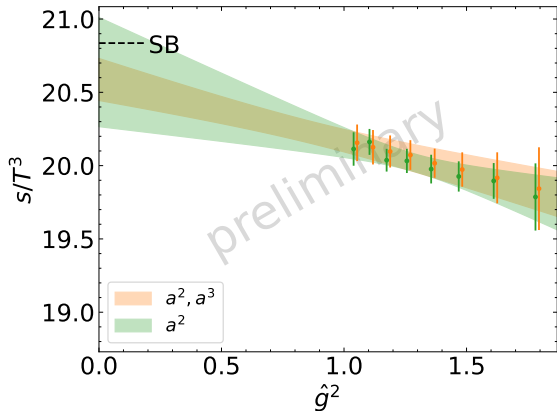
$$\langle T_{\mu\nu}^{R,\{i\}} \rangle_{\xi} = Z_G^{\{i\}} \langle T_{\mu\nu}^{G,\{i\}} \rangle_{\xi} + Z_F^{\{i\}} \langle T_{\mu\nu}^{F,\{i\}} \rangle_{\xi} \quad i = 3, 6$$

# Log corrections to continuum limit



$$s(g_i, a/L_0)/T^3 = s_i + [\alpha(a^{-1})]^\gamma \left(\frac{a}{L_0}\right)^2 (s_{23} g_i^3 + s_{24} g_i^4 + s_{25} g_i^5)$$

# Entropy vs $\hat{g}^2$



Fit with  $s_0, s_1$  free parameters

$$s(T)/T^3 = s_0 + s_1 \hat{g}^2$$