Simulating an SO(3) Quantum Link Model with Dynamical Fermions in 2+1 Dimensions



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Quantum Technologies for Fundamental Physics

• Quantum Link Models: What and Why?

• Exact gauge symmetry with a finite-dimensional Hilbert space

• SO(3) QLM and Nuclear Physics

- Imposing gauge-invariance on a per-site basis: 1+1d and 2+1d
- Confinement and chiral symmetry-breaking in 1+1d
- Phase diagram information from 1+1d

• Exact Diagonalization for a Single Plaquette

- The plaquette observable as an order parameter
- Explicit and spontaneous chiral symmetry-breaking
- Exploiting discrete symmetries for variational quantum simulation

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Quantum Link Models (QLMs)

- Generalisations of Wilsonian LGTs, with link operators on finite-dimensional Hilbert spaces.
- Exact gauge symmetry

 Choice of embedding algebra
- Well-posed continuum limits
 - Dim. reduction (D-theory)¹
 - Large spin representations²



1] Wiese, 2022, "From QLMs to D-theory…"

[2] Zache et al., 2022, "Toward the continuum limit..."

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Constructing an SO(3) QLM

- so(3) algebra realized on link operators O^{ab}, L^a, R^b
- Embed into so(6) bilinears:
 - $\circ \quad O^{ab} = \quad \sigma^a \otimes \sigma^b$
 - $\circ L^a = \sigma^a \otimes Id$
 - $\circ \quad \mathsf{R}^{\mathsf{b}} = \quad \mathsf{Id} \otimes \sigma^{\mathsf{b}}$
- Adjoint fermions ψ^a

a) _{SU(3)} Barvon	b)	[adapted from Rico et al., 2018]			
$\langle \rangle$		3-d QCD	1-d SO(3)	2-d SO(3)	
	gauge	SU(3)	SO(3)	SO(3)	
	symmetry				
	chiral	$SU(2)_L \times SU(2)_R$	\mathbb{Z}_2	$\mathbb{Z}_2 \times \mathbb{Z}_2$	
	symmetry				
	flavor	$SU(2)_{L=R}$	I	\mathbb{Z}_2	
SO(3) "Barvon"	symmetry				
	baryon	U(1)	U(1)	U(1)	
	symmetry				
	charge	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_2	
	conjugation				
$\langle \bigcirc \rangle$	parity	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_2	

• H includes hopping, staggered mass, plaquette, and four-Fermi couplings

Hamiltonian:
$$H = -t \sum_{x,k} \left[\eta_{x,k} B_{x,\hat{k}}^{\dagger} B_{x+k,-\hat{k}} + h.c. \right] + m \sum_{x} \eta_{x} M_{x}$$

 $- \frac{1}{4g^{2}} \sum_{x} \Phi_{x,34} \Phi_{x,23} \Phi_{x,12} \Phi_{x,41} + G \sum_{x} M_{x}^{2} + V \sum_{x,k} M_{x} M_{x+k}$

Gauge-Invariant Basis

Four spins per site: 2 singlets, 9 triplets
 Fermions: singlets 0,3 ; triplets 1,2



[Rico et al., 2018]

- Combining gives 10 gauge-singlets per site
 Solo
 <
- Gauge-invariant operators

$$B_{x,k} = \psi_x^a \sigma_{x,k}^a, \, M_x = \psi_x^{a,\dagger} \psi_x^a, \, \Phi_{x,k,l} = \sigma_{x,k}^a \sigma_{x,l}^a.$$



SO(3) "Nuclear Physics" in (1+1)d

- For V>0 & G<V, chiral symmetry breaking restored at finite B density
- Conformal phase for G >> V > 0, few-body bound states for G<0 or G << V



(1+1)d Gauge-Invariant Basis

- 4 gauge-singlet states
- Operators M, B⁺, B⁻
- Exact correspondence to spin-3/2 chain

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Magnetic Observable (Plaquette)

• Three limiting phases – sharp massless transition to high-field, at $g_{M} \approx 1.28$.



Explicit Chiral Symmetry Breaking

• m>0 explicitly breaks chiral symmetry – weakly within strong-field phase.



Spontaneous Chiral Symmetry Breaking

- Closer look reveals SSB at $g_x \approx 0.28$ for m = 0.
- Coincides with degenerate ground state.



Summary

- We have identified the gauge-invariant state space of an SO(3) QLM w/ fermions in (2+1)d.
- We have performed exact diagonalisation for a single plaquette, finding:
 - Distinct magnetic phases, discontinuous transition at $g_M \approx 1.28$,
 - Explicit chiral symmetry breaking with two phases,
 - Spontaneous chiral symmetry breaking at strong-coupling, $g_v \approx 0.28$.
- *Caution*: we should take single-plaquette results as broad approximations.
- Next: study larger (2+1)d lattices, and simulate via variational QAs.

THANK YOU FOR LISTENING!

Works Cited

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Thank you to Lattice 2024, and our hosts here in Liverpool!

Extra: Variational QAs and Discrete Symmetries

- QLMs generically have numerous discrete symmetries (C, P, translation)
- Eliminating redundant basis states = smaller variational ansätze.

