

The Four Gluon Vertex from Lattice QCD

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Phys. Rev. D109 (2024) 7, 074502 [arXiv:2401.12008]



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Motivation

QCD Dynamics at a fundamental level

importance of higher order corrections (two loop contributions to gluon DSE, etc.)

Test our understanding of QCD Green functions (ghost dominance at IR)

Effective charge $\frac{g^2}{4\pi} \Gamma^{(4g)}(p^2) p^2 D(p^2)$

Landau gauge

(pure gauge) Lattice point of view

Lattice Approach

Importance Sampling **allows to access the gluon field** $A_{\mu}^a(x)$

$$\mathcal{G}^{(n)}(x_1, \dots, x_n) = \langle 0 | T \left(A_{\mu_1}^{a_1}(x_1) \cdots A_{\mu_n}^{a_n}(x_n) \right) | 0 \rangle$$

Larger n:



Noisy

Disconnected parts



Larger sets of configurations

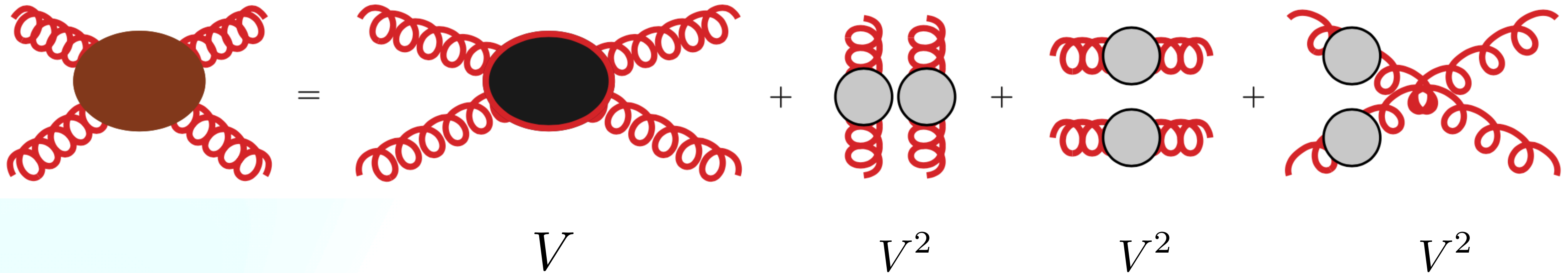


kinematical configurations
or
Combinations of components

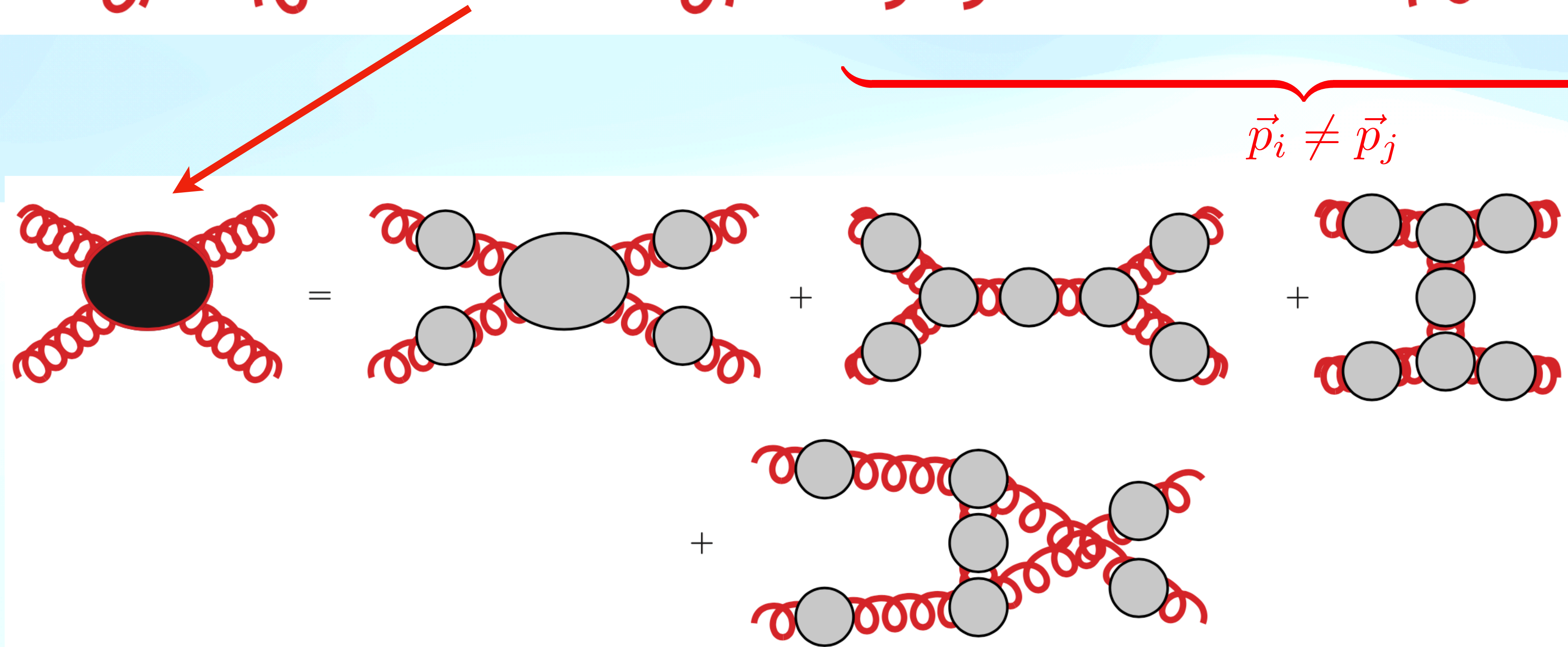
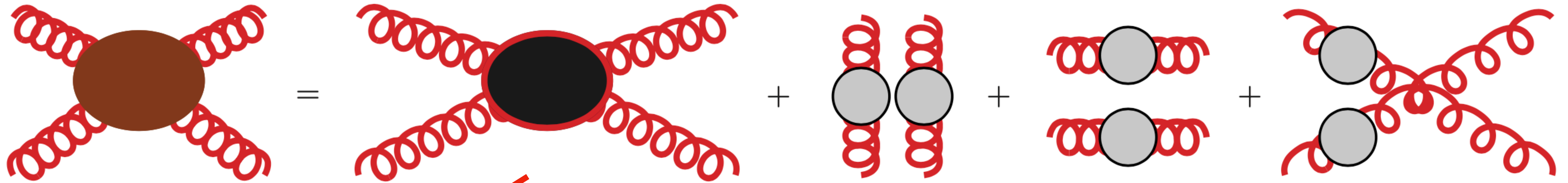
4-gluon correlation function

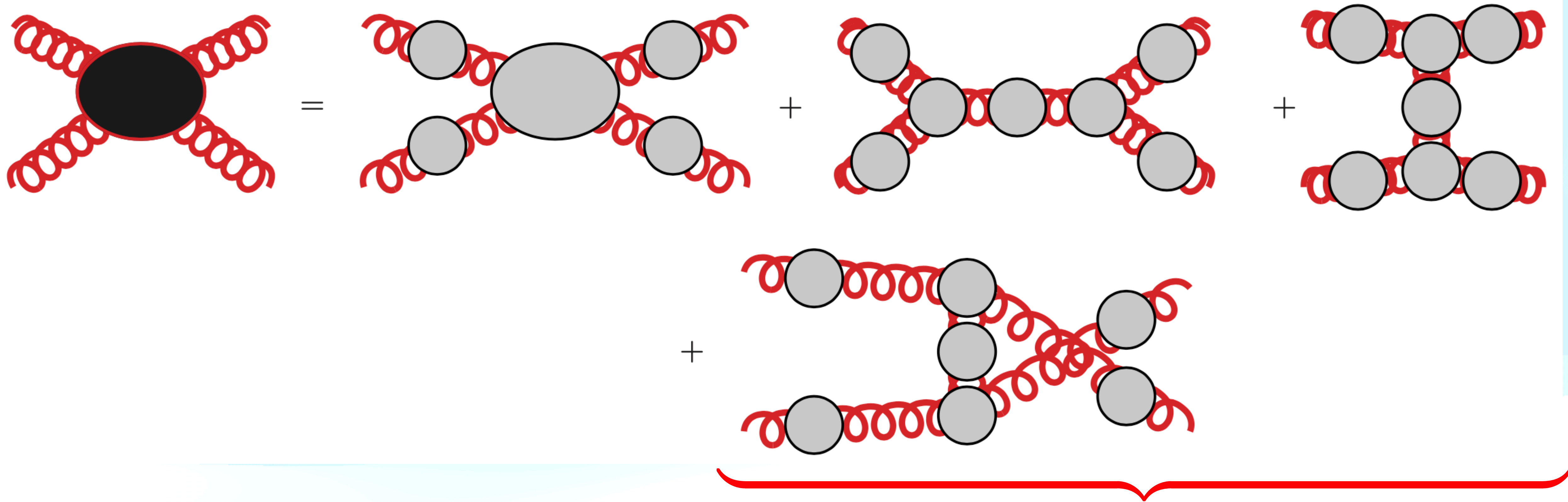
$$\langle A_{\mu_1}(p_1) A_{\mu_2}(p_2) A_{\mu_3}(p_3) A_{\mu_4}(p_4) \rangle =$$

$$= \int \mathcal{D}A A_{\mu_1}(p_1) A_{\mu_2}(p_2) A_{\mu_3}(p_3) A_{\mu_4}(p_4) e^{-S}$$



4-gluon correlation function





Single momentum scale

$$\vec{p}_i \propto \vec{p}$$

In **Landau gauge** simplifies the tensor analysis

Contributions only from the tensors proportional to $\delta_{\mu\nu}$

Recent Continuum Calculations

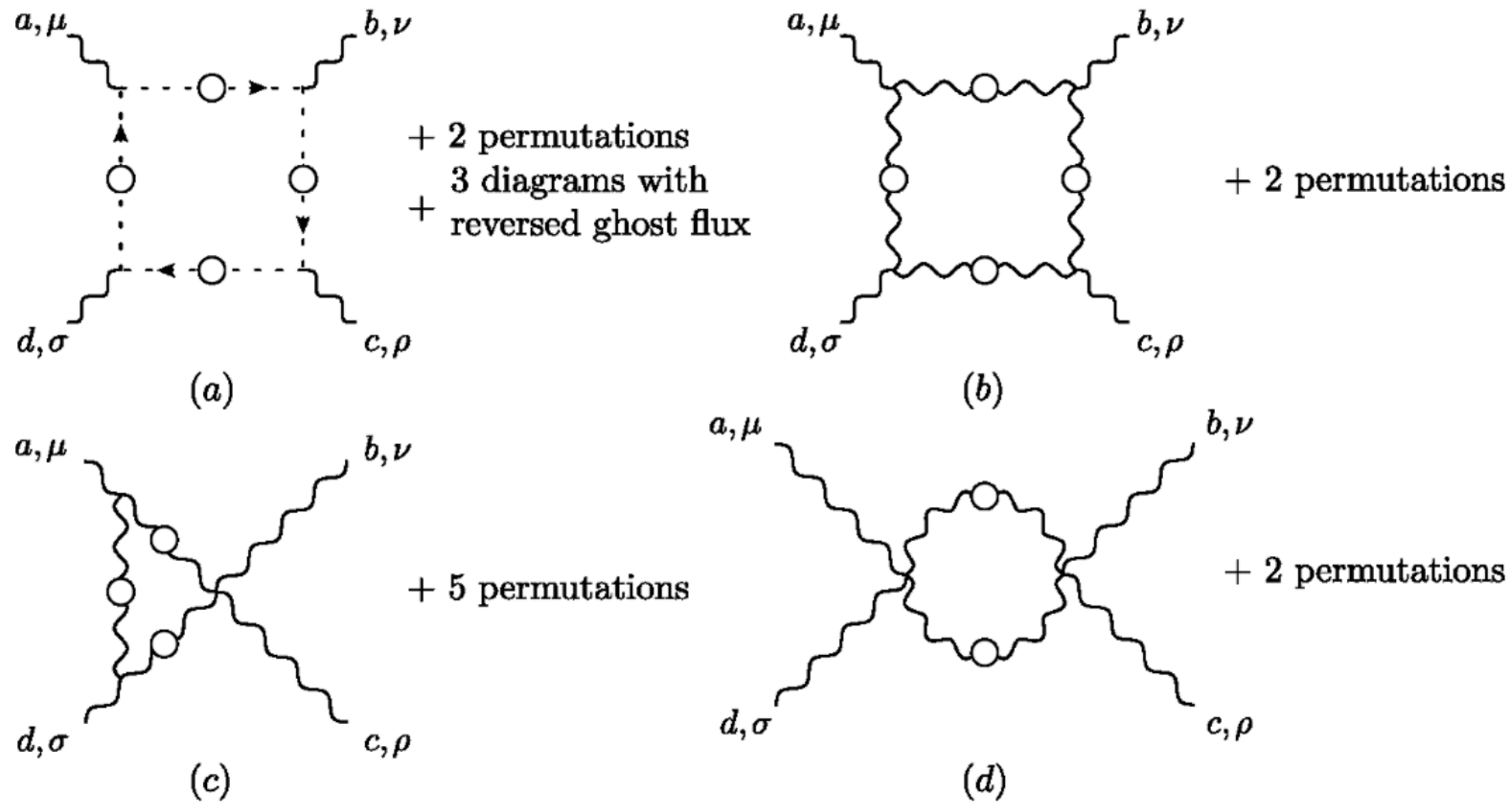
C Kellermann, C S Fischer, Phys. Rev. D **78** (2008) 025015

D Binosi, D Ibañez, J Papavassiliou, JHEP **1409** (2014) 059

A K Cyrol, M Q Huber, L von Smekal, Eur Phys J C **75** (2015) 102

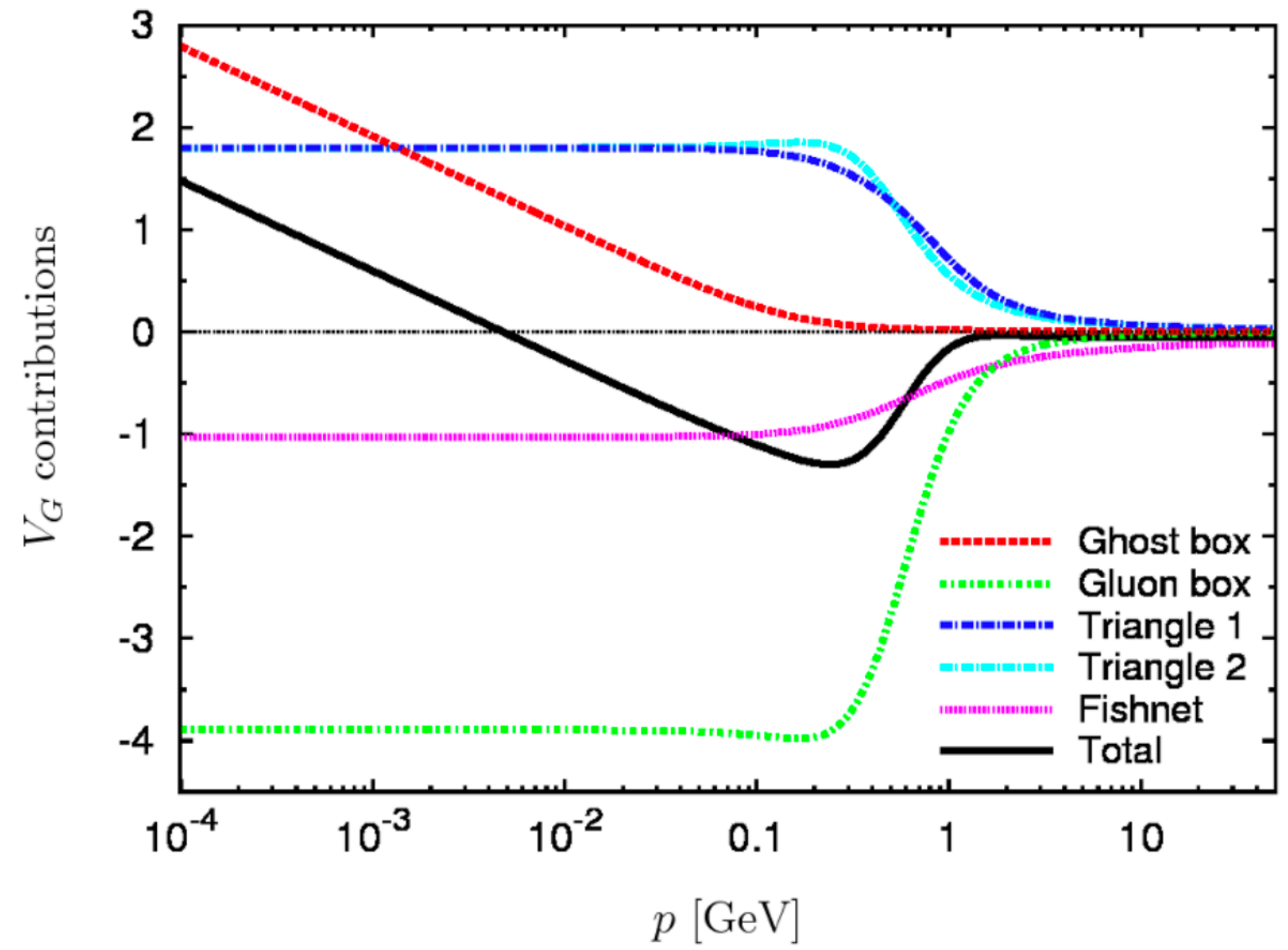
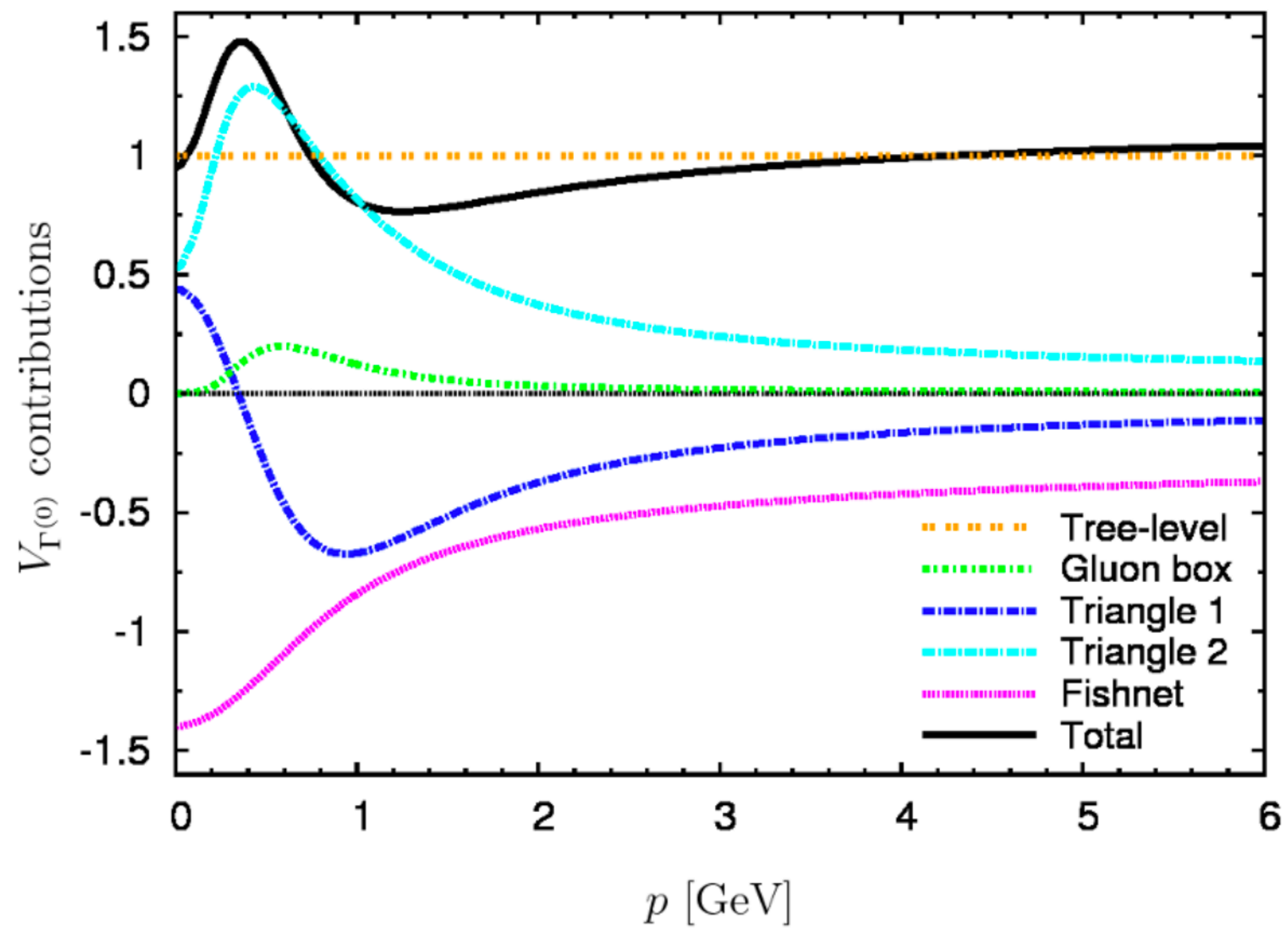
A C Aguilar, M N Ferreira, J Papavassiliou, L R Santos, Eur Phys J C **84** (2024), 676

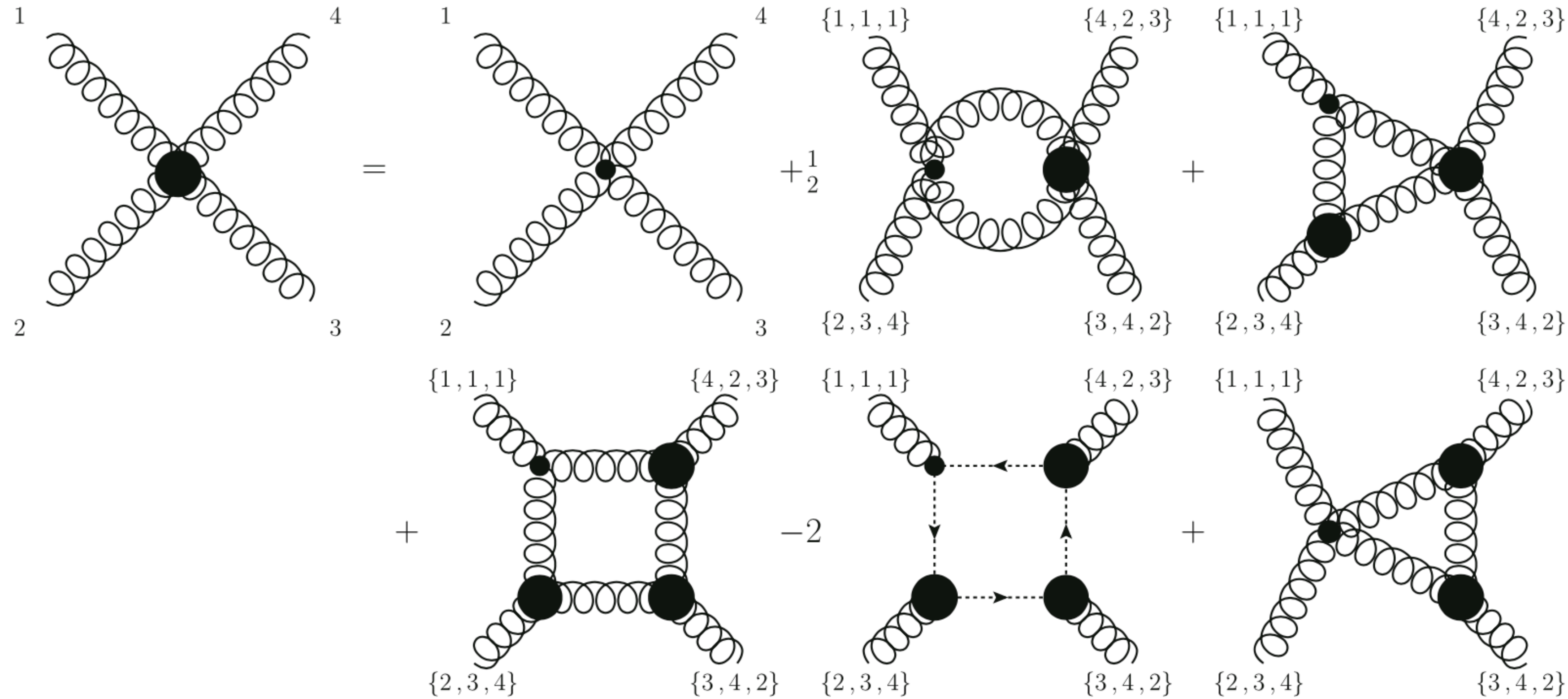
N Barrios, P De Fabritiis, M Peláez, Phys Rev **109** (2024) L091502



$$\Gamma_{\mu\nu\rho\sigma}^{abcd}(p, p, p, -3p) \Big|_{gg} = V_{\Gamma(0)}(p^2) \Gamma_{\mu\nu\rho\sigma}^{abcd(0)} + V_G(p^2) G_{\mu\nu\rho\sigma}^{abcd},$$

$$G_{\mu\nu\rho\sigma}^{abcd} = (\delta^{ab} \delta^{cd} + \delta^{ac} \delta^{bd} + \delta^{ad} \delta^{bc}) \underbrace{(g_{\mu\nu} g_{\rho\sigma} + g_{\mu\rho} g_{\nu\sigma} + g_{\mu\sigma} g_{\nu\rho})}_{R_{\mu\nu\rho\sigma}}$$

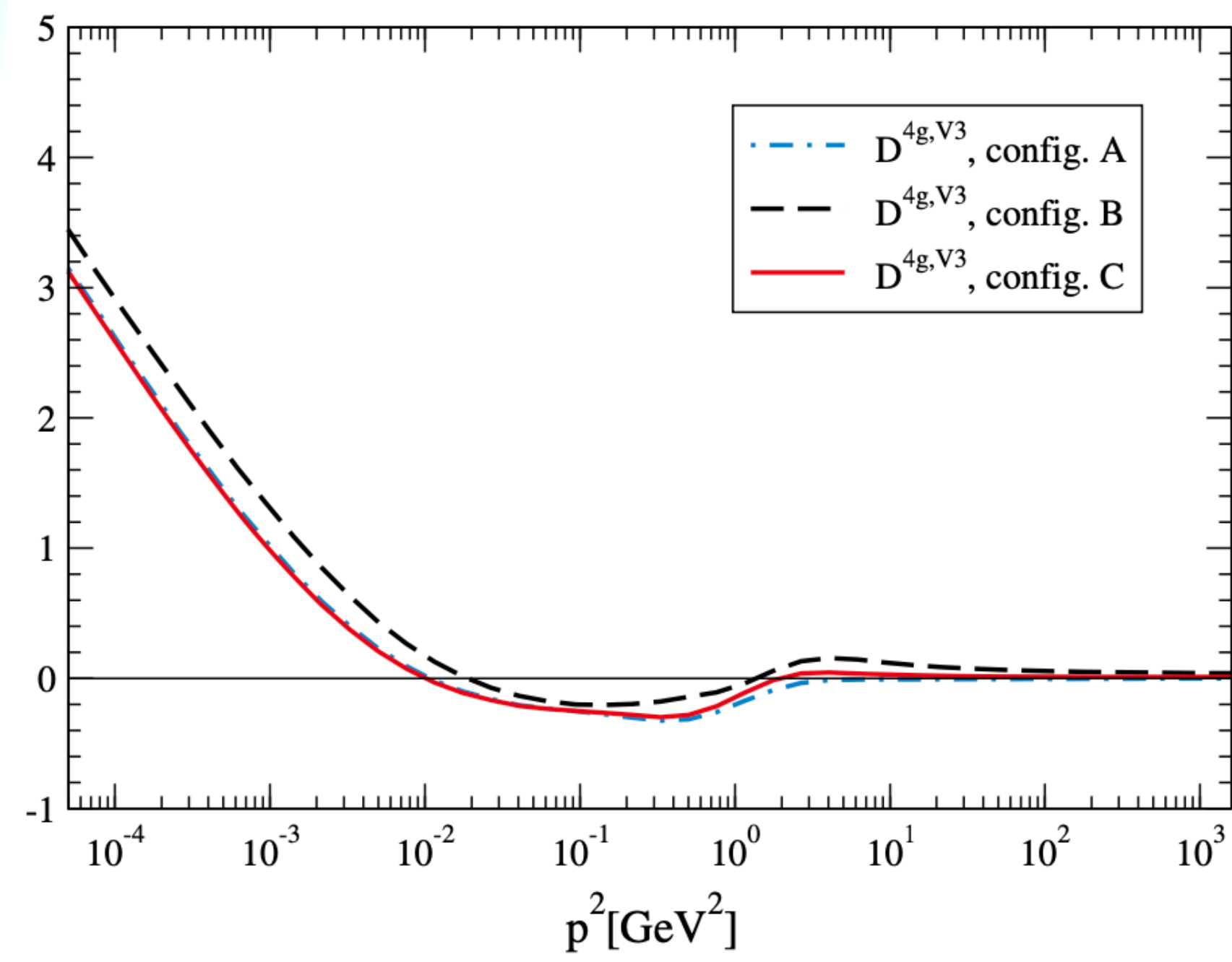
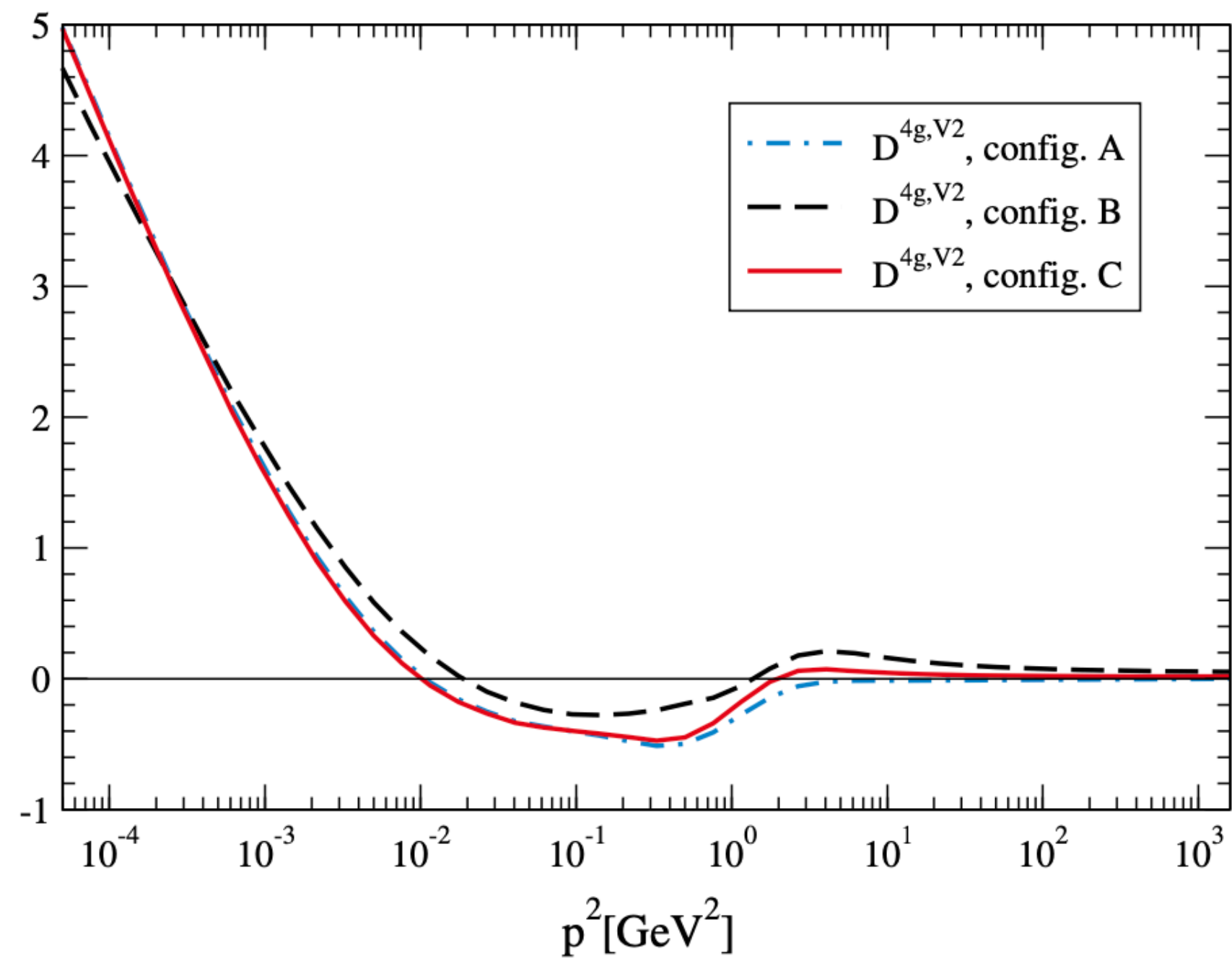
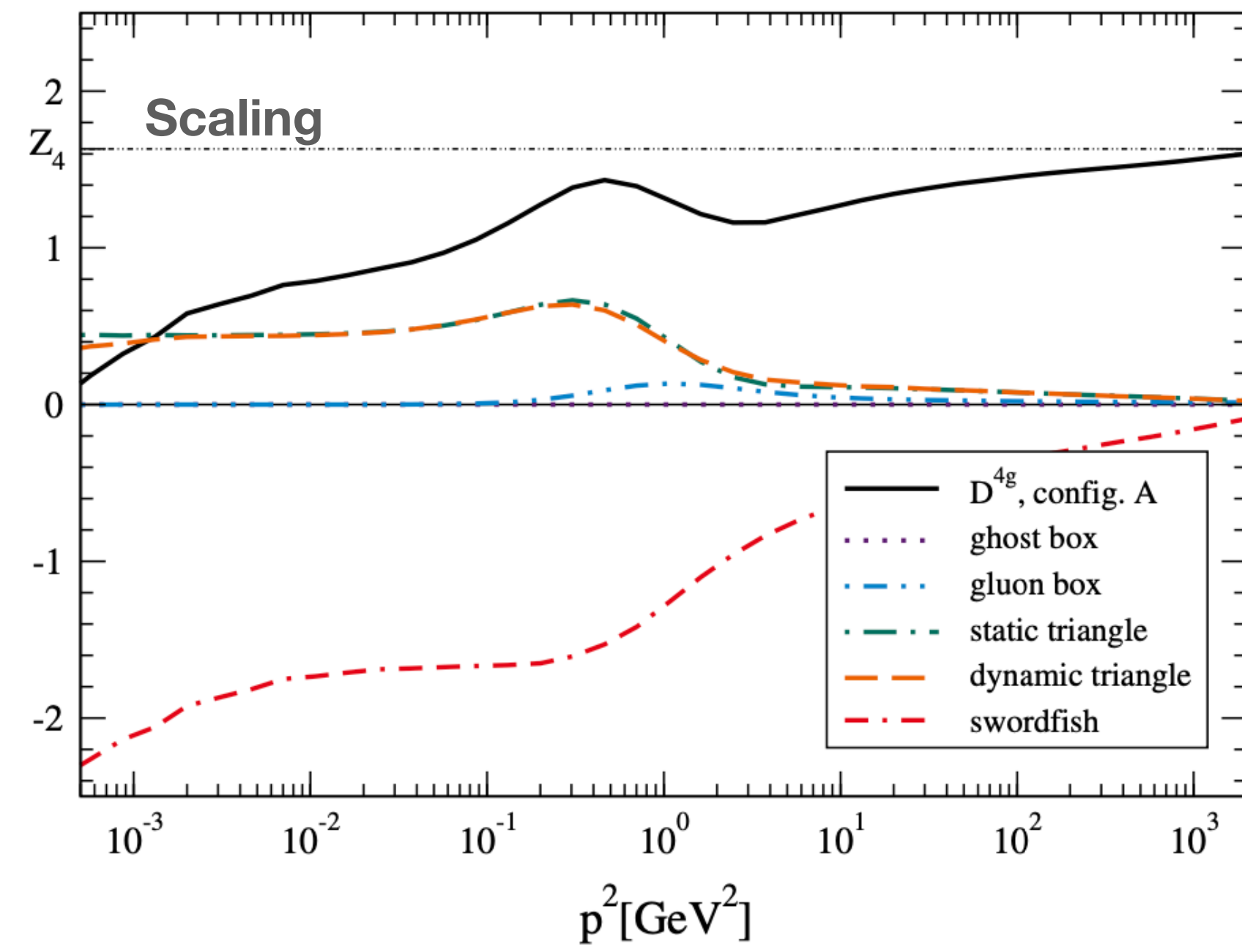
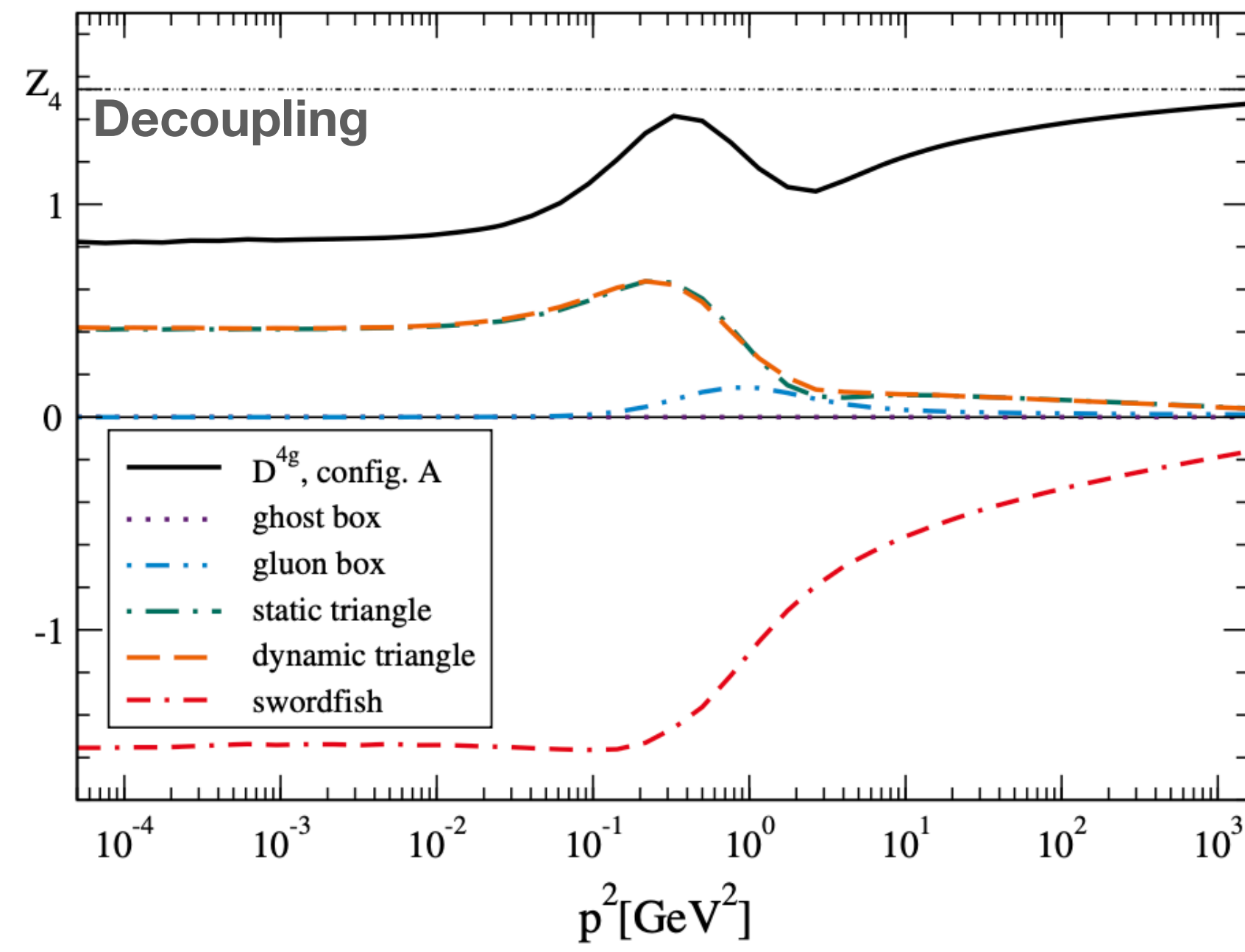




IR dominant (primitively divergent diagrams)

One-loop contributions (disregard contributions with five-point diagrams and ghost-gluon five-point functions)

Take tree level tensor structure for three-gluon and four-gluon diagrams



For **tensor analysis** see:

J A Gracey, Phys Rev D90, 025011 (2014) arXiv: 1406.1618

G Eichmann, C S Fischer, W Heupel, Phys Rev D92, 056006 (2015) arXiv: 1505.06336

$$\tilde{\Gamma}^{(0)}{}_{\mu\nu\eta\zeta}{}^{abcd} = f_{abr}f_{cdr}(\delta_{\mu\eta}\delta_{\nu\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{acr}f_{bdr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\zeta}\delta_{\nu\eta}) + f_{adr}f_{bcr}(\delta_{\mu\nu}\delta_{\eta\zeta} - \delta_{\mu\eta}\delta_{\nu\zeta})$$

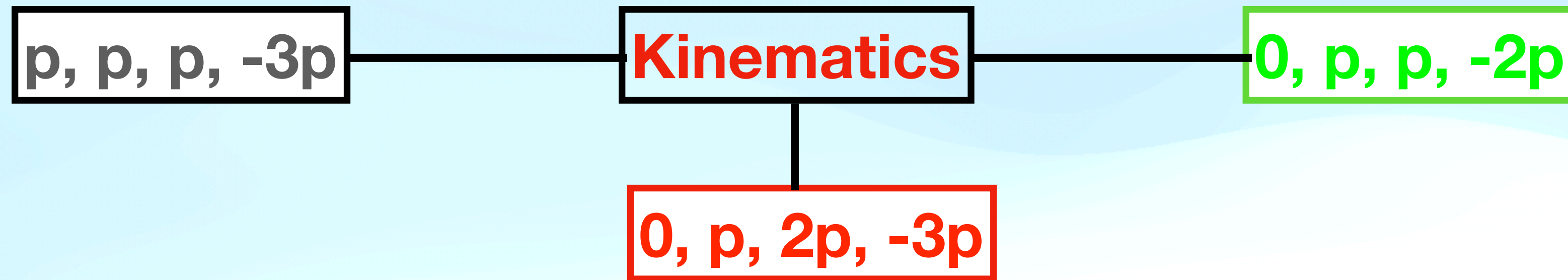
$$\tilde{\Gamma}^{(1)}{}_{\mu\nu\eta\zeta}{}^{abcd} = d_{abr}d_{cdr}(g_{\mu\eta}g_{\nu\zeta} + g_{\mu\zeta}g_{\nu\eta}) + d_{acr}d_{bdr}(g_{\mu\zeta}g_{\nu\eta} + g_{\mu\nu}g_{\eta\zeta}) + d_{adr}d_{bcr}(g_{\mu\nu}g_{\eta\zeta} + g_{\mu\eta}g_{\nu\zeta})$$

$$\tilde{\Gamma}^{(2)}{}_{\mu\nu\eta\zeta}{}^{abcd} = (\delta^{ab}\delta^{cd} + \delta^{ac}\delta^{bd} + \delta^{ad}\delta^{bc})(\delta_{\mu\nu}\delta_{\eta\zeta} + \delta_{\mu\eta}\delta_{\nu\zeta} + \delta_{\mu\zeta}\delta_{\nu\eta})$$

$$\tilde{\Gamma} = F(p^2)\tilde{\Gamma}^{(0)} + G(p^2)\tilde{\Gamma}^{(1)} + H(p^2)\tilde{\Gamma}^{(2)}$$

Not an orthogonal basis

$$\mathcal{G}^{(4)} = \tilde{\Gamma} \left(P^\perp(p) D(p^2) \right)^3 \left(P^\perp(3p) D(9p^2) \right)$$



Measure the three form factors

$$F^{(0)}(p^2)$$

$$F^{(1)}(p^2)$$

$$F^{(2)}(p^2)$$

Tree Level Tensor

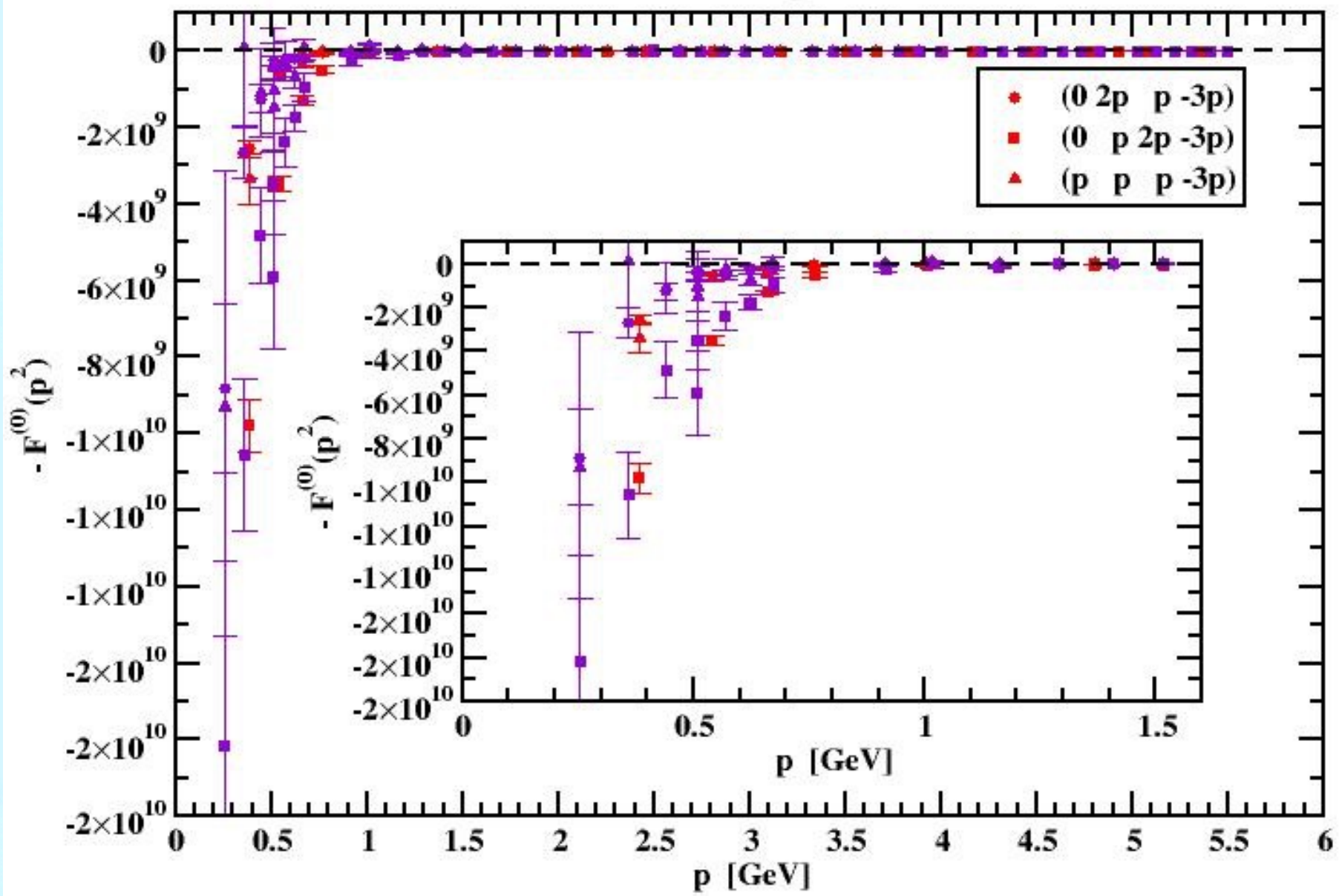
$$\beta = 6.0 \quad a = 0.102 \text{ fm} \quad a^{-1} = 1.943 \text{ GeV}$$

Lattice	Configs	p_{\min}
32^4	9038	381 MeV
48^4	9035	254 MeV

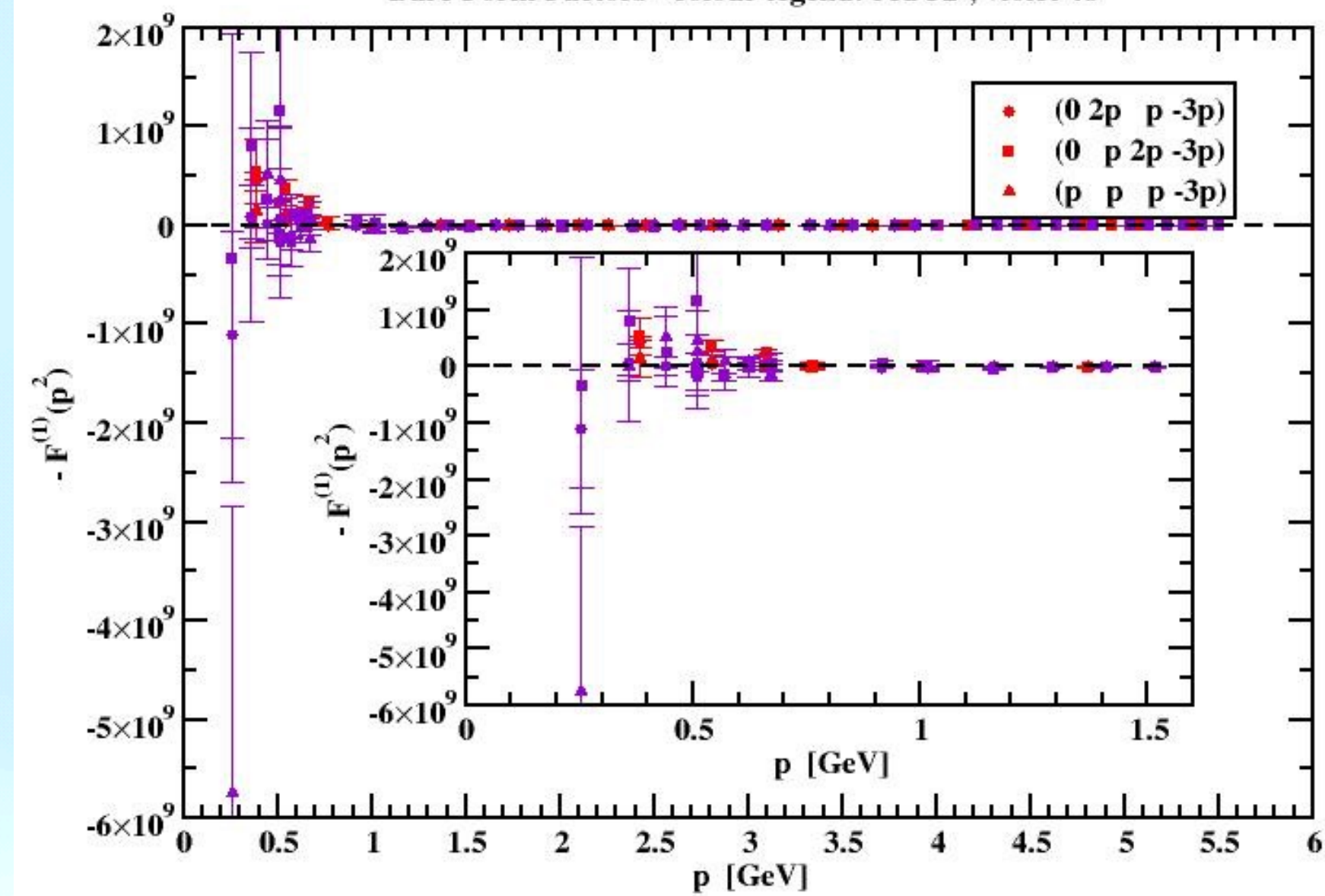
Averaged over equivalent momenta, including the negative momenta !

(1 1 1 1) (-1 1 1 1) (1 -1 1 1) (1 1 -1 1) (1 1 1 -1)
 (-1 -1 1 1) (-1 1 -1 1) (-1 1 1 -1) (1 -1 -1 1) (1 -1 1 -1) (1 1 -1 -1)
 (1 -1 -1 -1) (-1 1 -1 -1) (-1 -1 1 -1) (-1 -1 -1 1)
 (-1 -1 -1 -1)

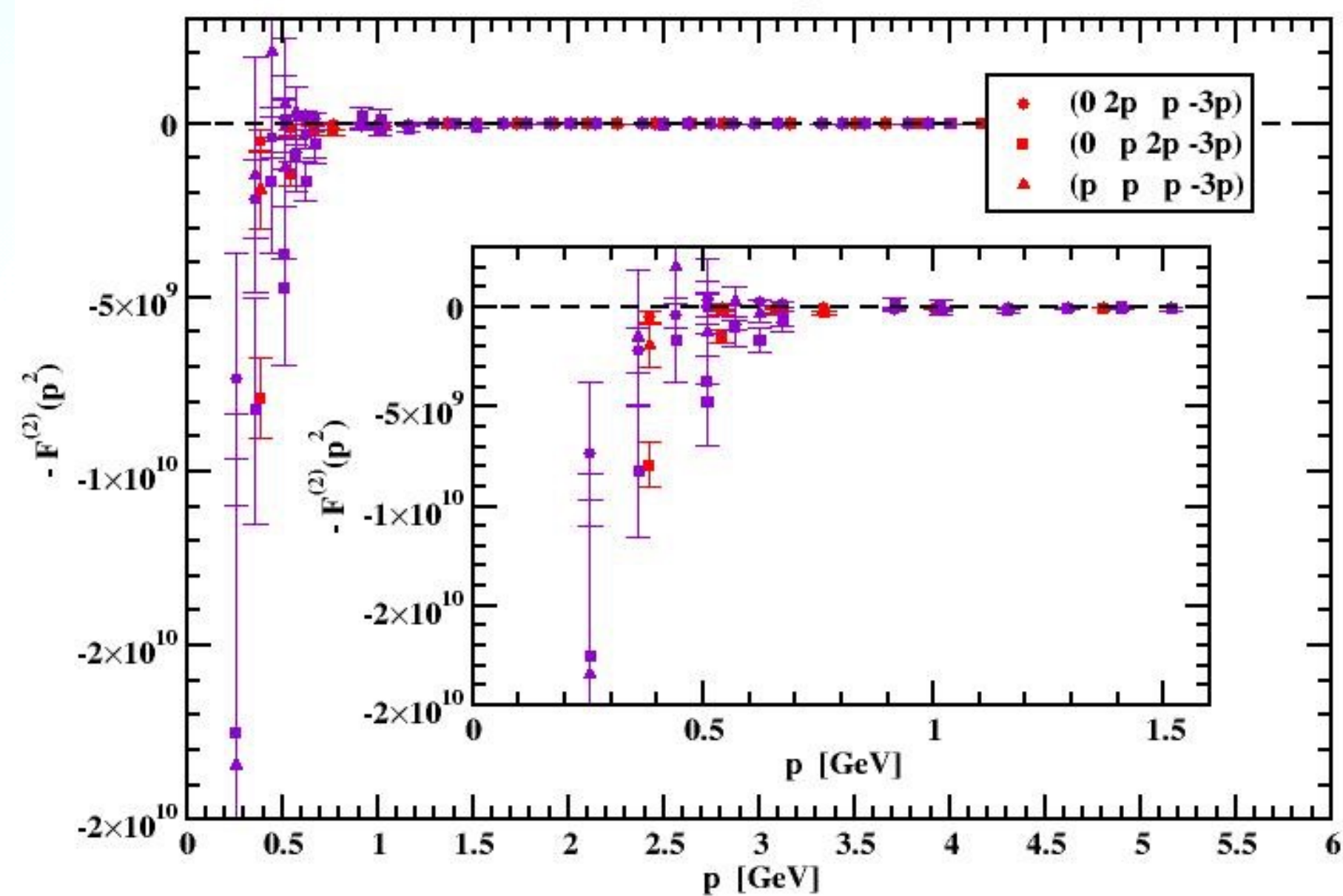
Bare Form Factors - colour legend: red 32^4 , violet 48^4



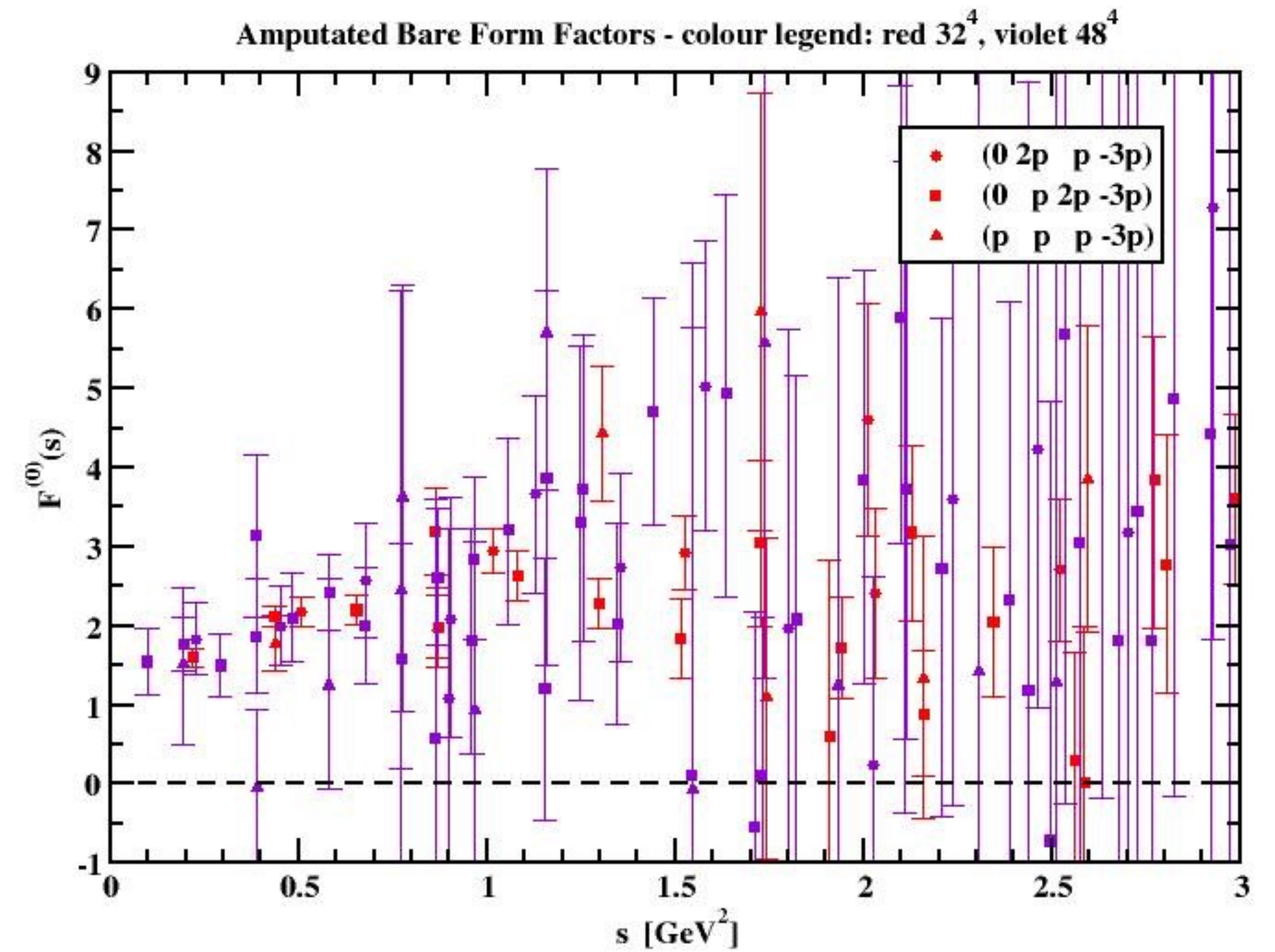
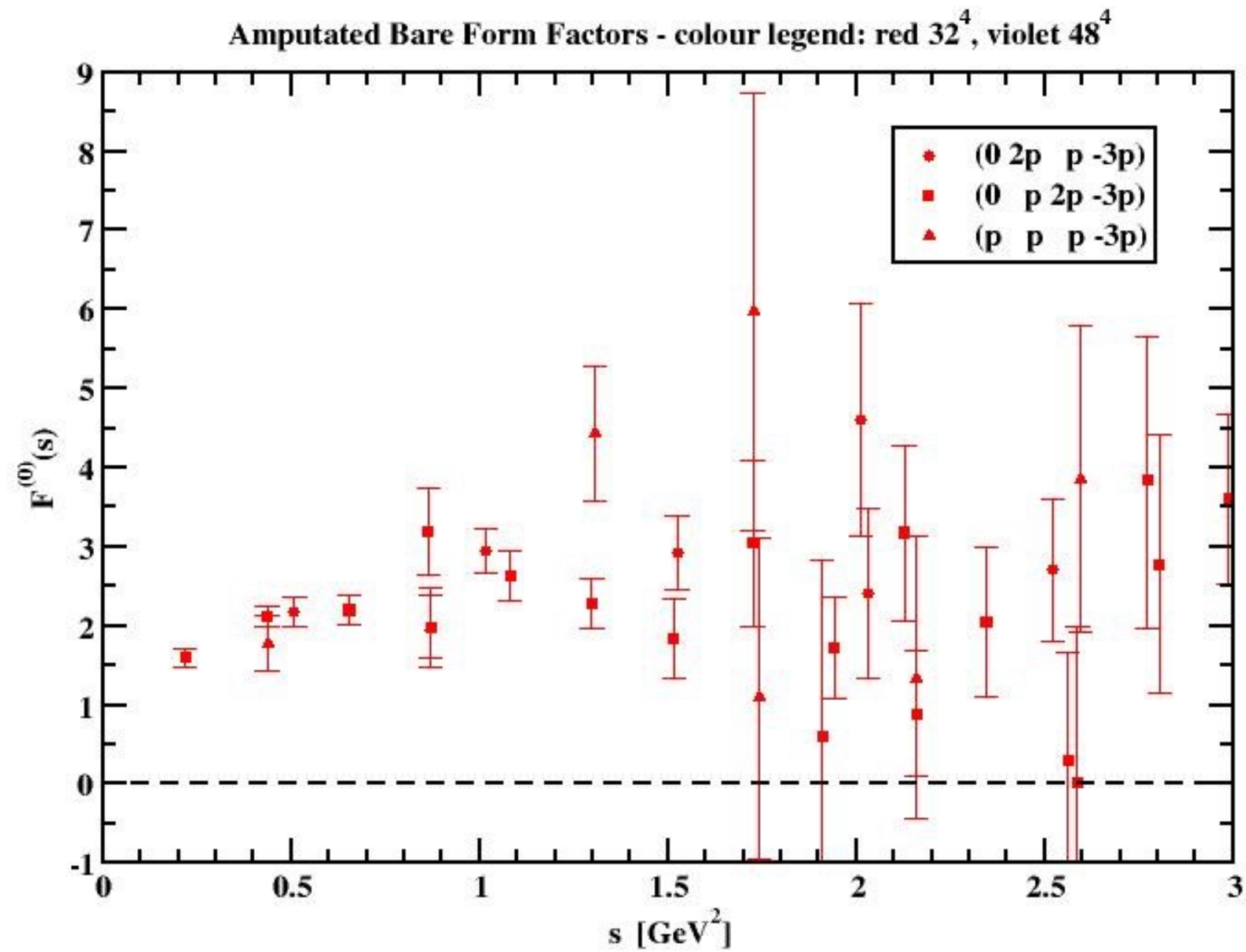
Bare Form Factors - colour legend: red 32^4 , violet 48^4



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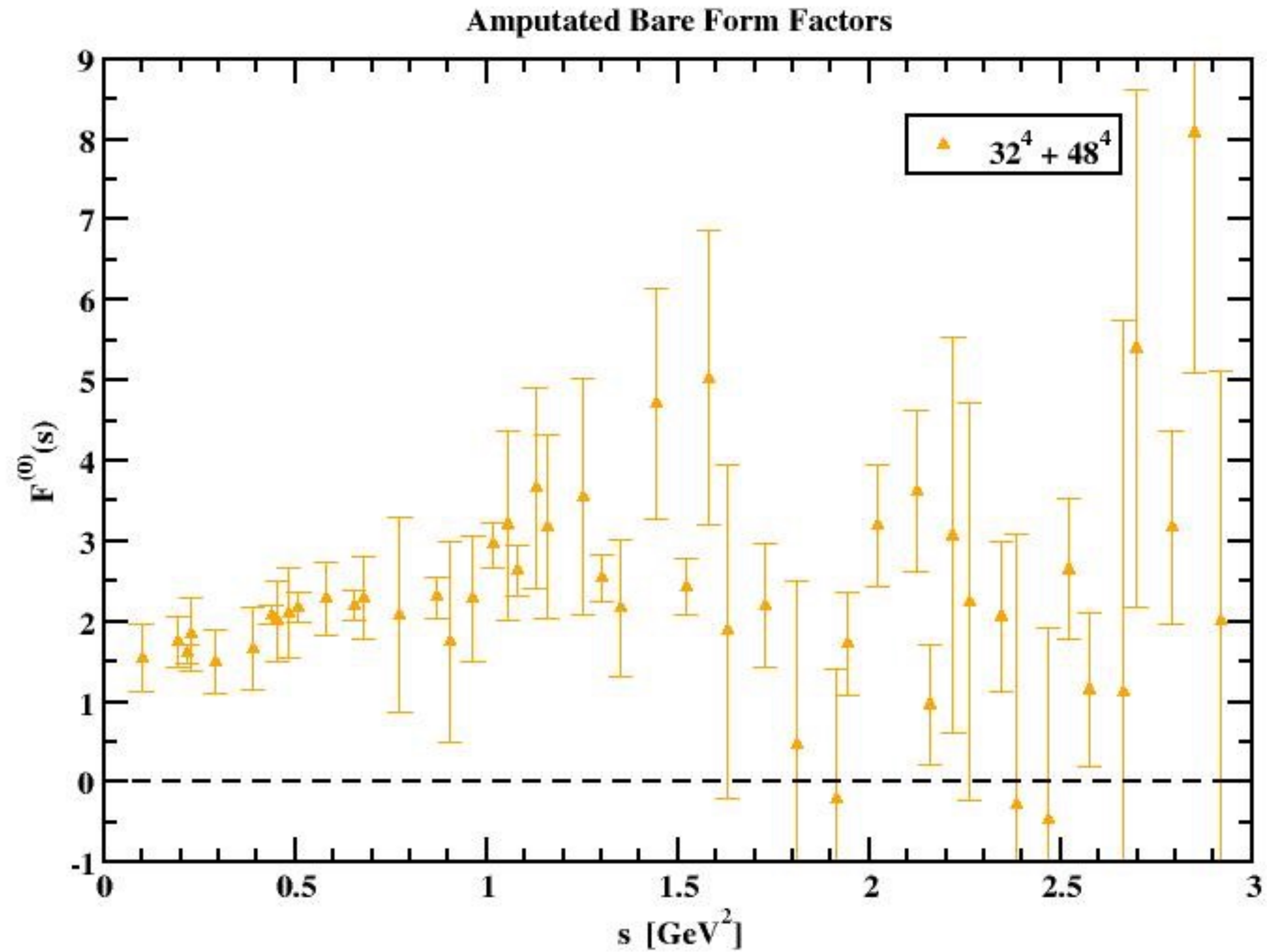


Bare Amputated Green Function

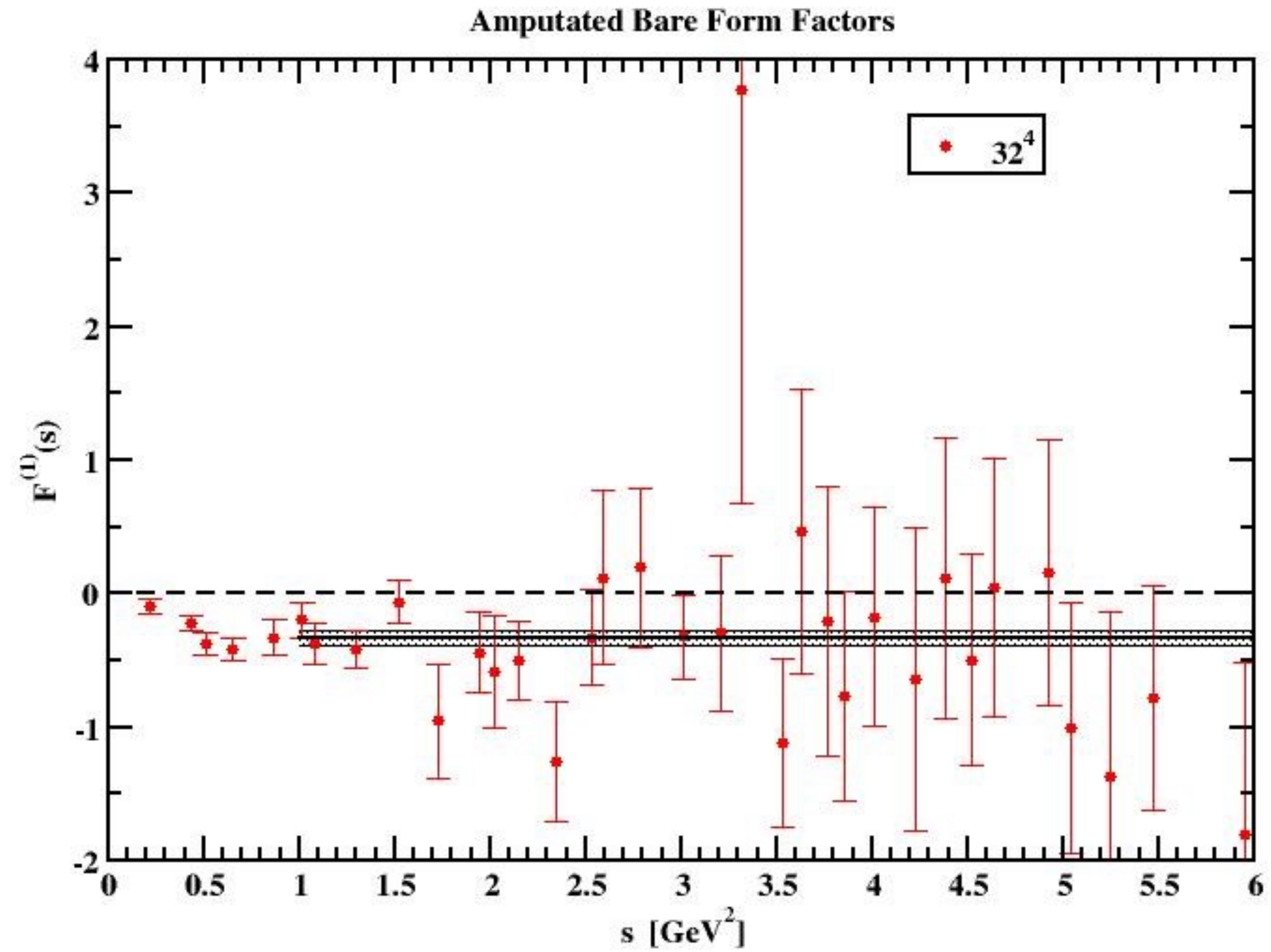


$$s = \frac{1}{4} \sum_i p_i^2$$

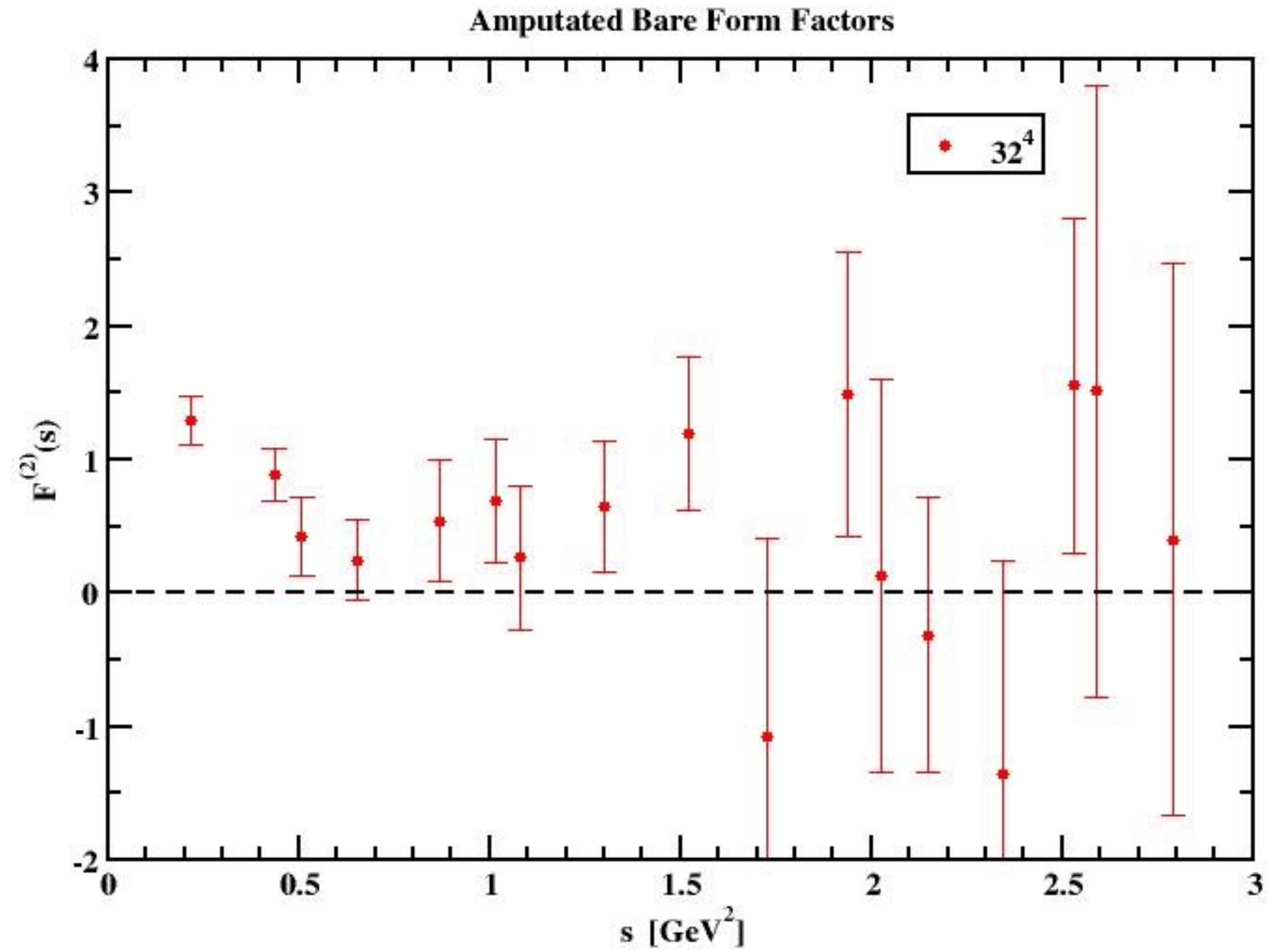
Bare Amputated Green Function



Bare Amputated Green Function



Bare Amputated Green Function



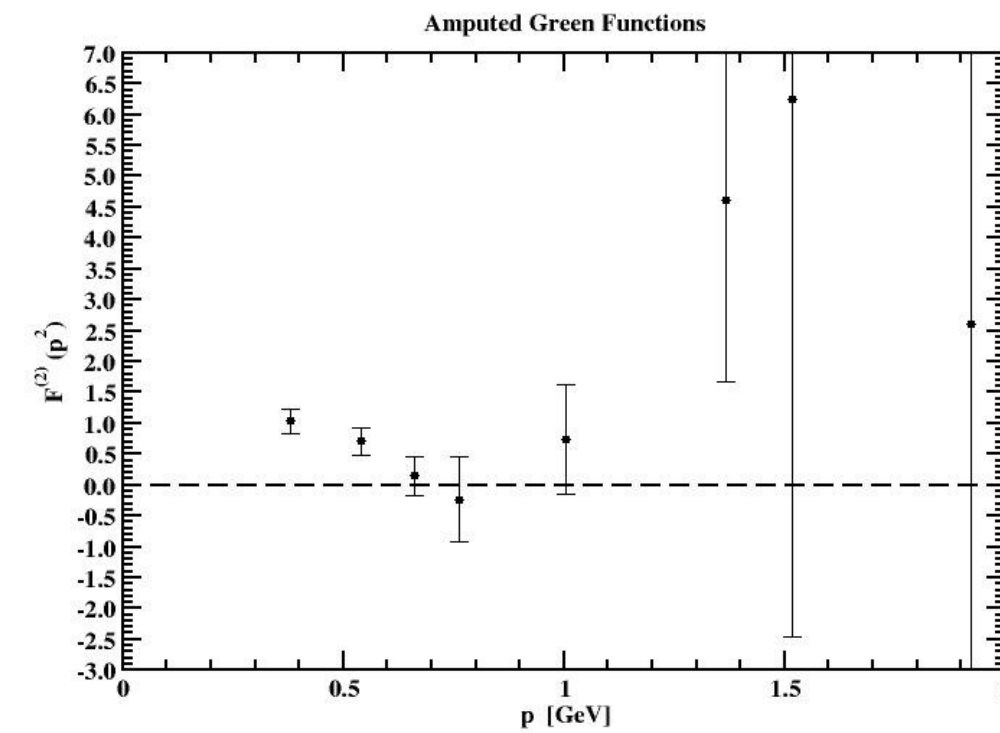
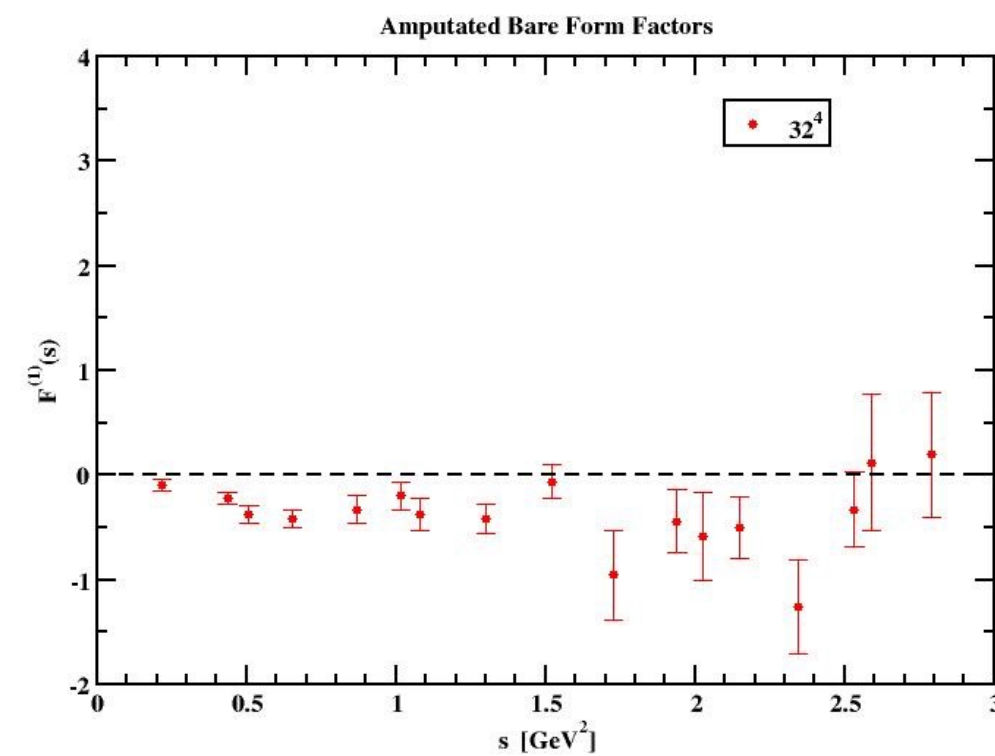
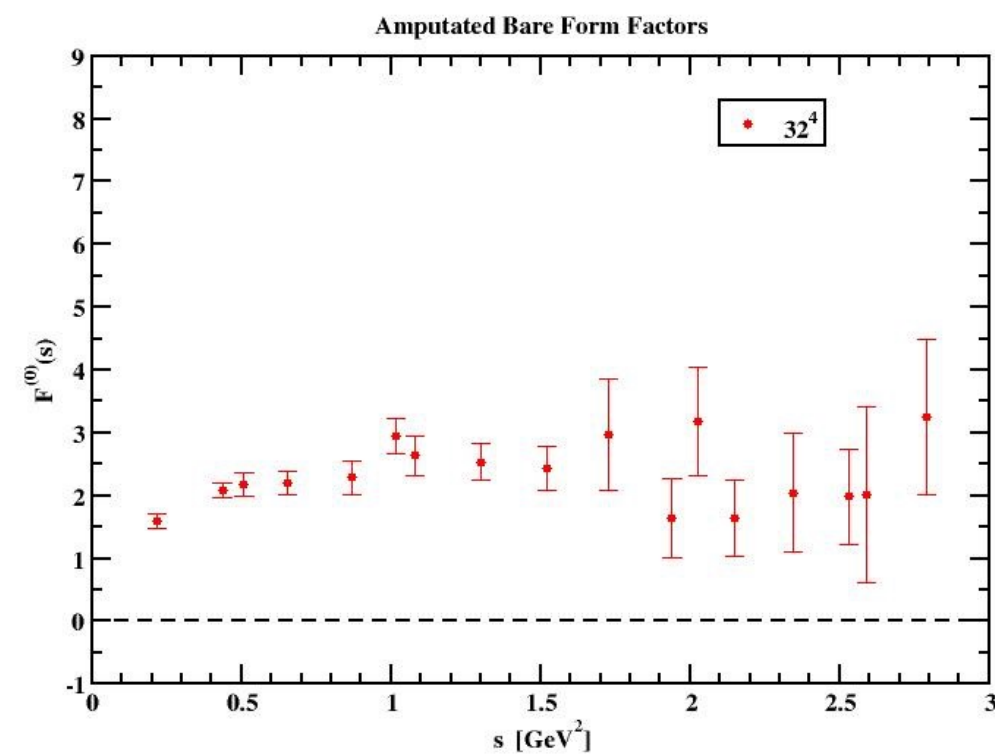
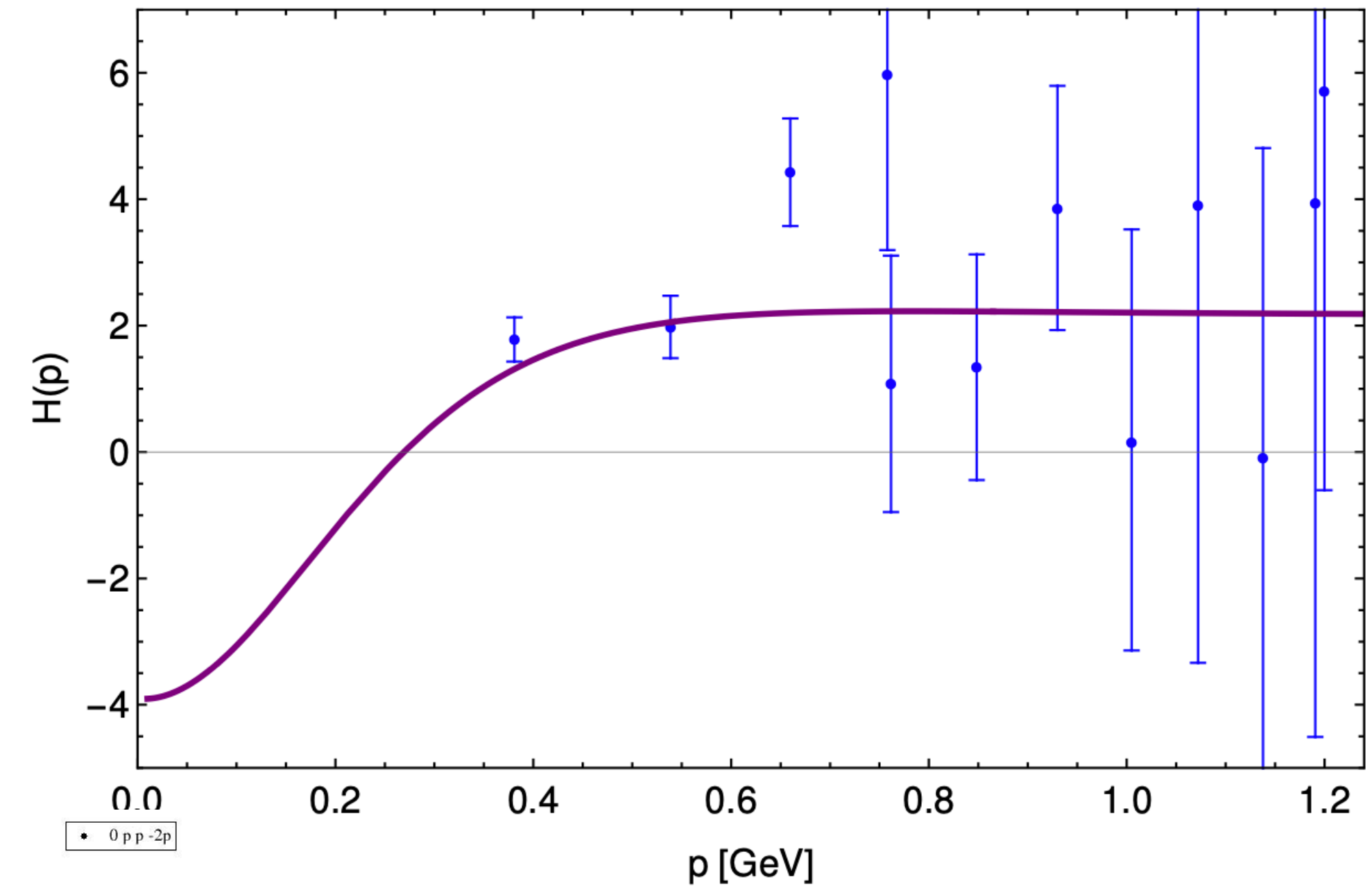
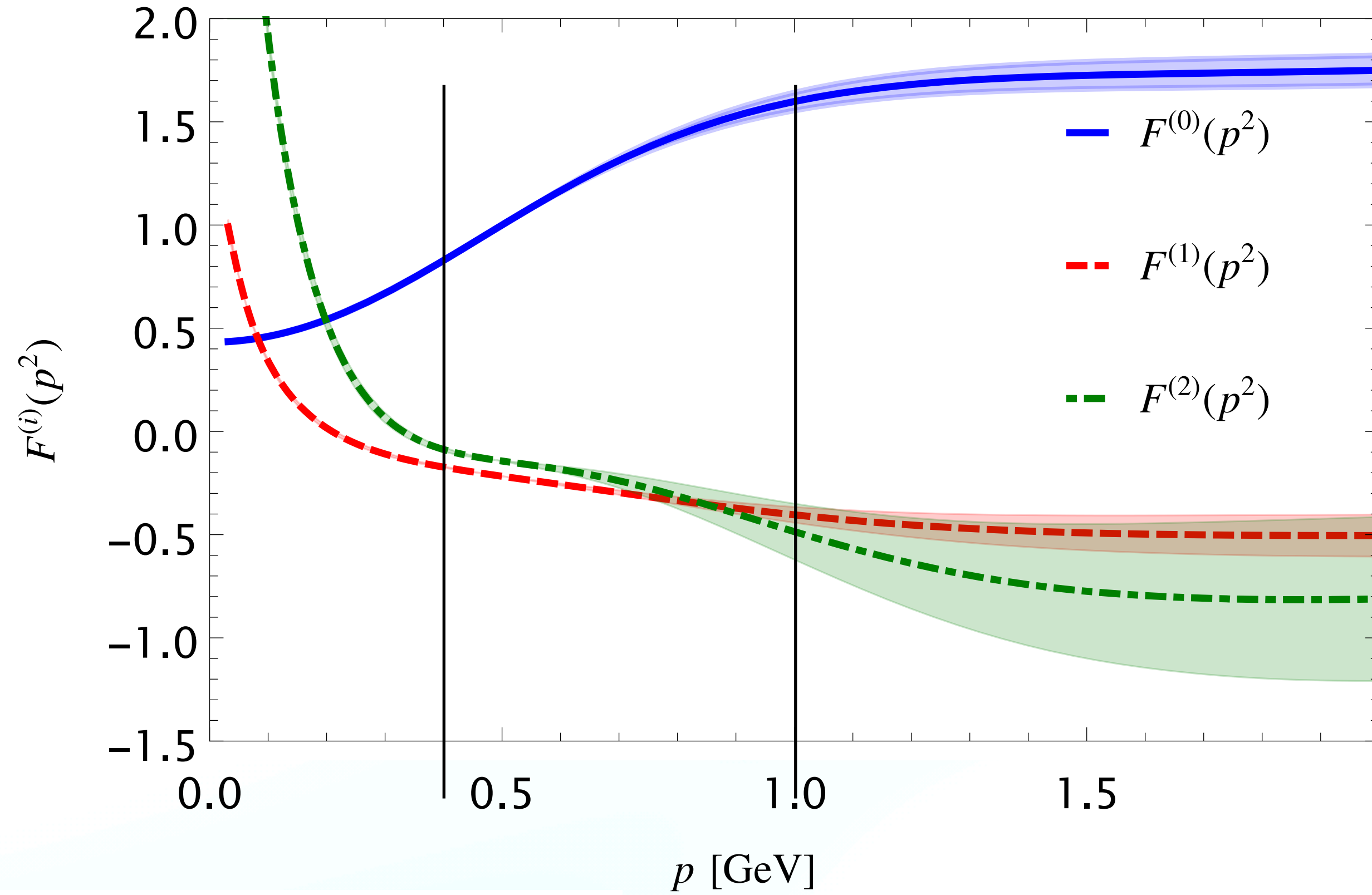
Continuum Calculations

One loop truncated Dyson-Schwinger equation by

A C Aguilar, M N Ferreira, J Papavassiliou, L R Santos, Eur Phys J C 84 (2024) 676

Curci-Ferrari model (massive QCD)

N Barrios, P De Fabritiis, M Peláez, Phys Rev D 109 (2024) L091502



Summary and Conclusions

- gluon correlation functions up to 4-external legs are possible with standard lattice methods (call for large statistical ensembles)
- Dominated by tree level tensor structure but not only
- Good (at least) qualitative agreement between Lattice and continuum approaches
- Look at other kinematics and extend collaboration to get a better picture
- Need large statistics or methods to reduce variance

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