

Investigation of pion–nucleon contributions to nucleon matrix elements

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LATTICE 2024



LIVERPOOL

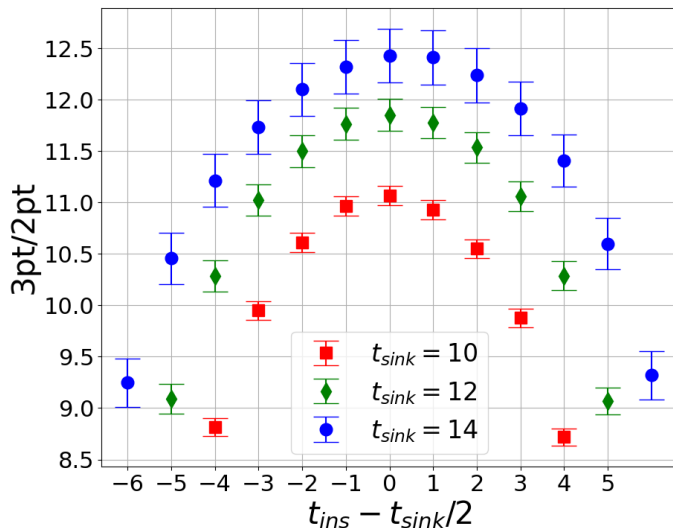
Background

- Nucleon structure: nucleon matrix elements

$$R_{\mathcal{O}} := \frac{\langle \mathcal{J}_N(t_{\text{sink}}) \mathcal{O}(t_{\text{ins}}) \mathcal{J}_N^\dagger(0) \rangle}{\langle \mathcal{J}_N(t_{\text{sink}}) \mathcal{J}_N^\dagger(0) \rangle} \xrightarrow{\text{all } t \text{ well-separated}} \langle N | \mathcal{O} | N \rangle$$

- Time-dependence indicates contamination from excited states
- Lowest excited state is a Nucleon-Pion state

$$\mathcal{O} = S^{u+d} = \bar{u}u + \bar{d}d \quad \vec{q} = (0, 0, 0)$$



Simulation details

Ensembles	Flavors	$N_L^3 \times N_T$	m_π (MeV)	L (fm)	$m_\pi L$	N_{cfg}
cA2.09.48	2	$48^3 \times 96$	131	4.50	2.98	1200

➤ Physical-point twisted-mass ensemble

An ensemble with $m_\pi = 346$ was also studied (upcoming preprint)

➤ Interpolating fields used:

$$J_p; \quad J_{N\pi}^{1/2} = \sqrt{2/3} J_{n\pi^+} - \sqrt{1/3} J_{p\pi^0}$$

➤ Generalized eigenvalue problem (GEVP)

➤ Do GEVP on 2pt functions

➤ Use the results to improve 3pt functions

2pt functions and GEVP

➤ 2pt functions:

$$\begin{bmatrix} \langle \mathcal{J}_N \mathcal{J}_N^\dagger \rangle & \langle \mathcal{J}_N \mathcal{J}_{N\pi}^\dagger \rangle \\ \langle \mathcal{J}_{N\pi} \mathcal{J}_N^\dagger \rangle & \langle \mathcal{J}_{N\pi} \mathcal{J}_{N\pi}^\dagger \rangle \end{bmatrix}$$

$$C_{ij}(t) = \langle \mathcal{J}_i(t) \mathcal{J}_j^\dagger(0) \rangle$$

➤ GEVP returns eigenvalues and eigenvectors:

$$C_{ij}(t) v_j^n = \lambda^n(t, t_0) C_{ij}(t_0) v_j^n$$

$$\lambda^n(t, t_0) = e^{-E_n(t-t_0)}, \quad v_j^n \mathcal{J}_j^\dagger(0) |0\rangle = |n\rangle$$

➤ We determine the optimal interpolating field:

$$\tilde{\mathcal{J}}_N |0\rangle = (\mathcal{J}_N + v_{N\pi}^N \mathcal{J}_{N\pi}) |0\rangle \propto |N\rangle$$

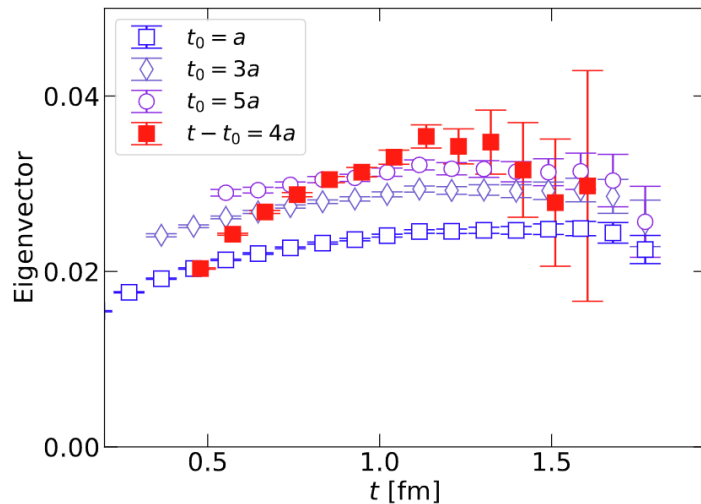
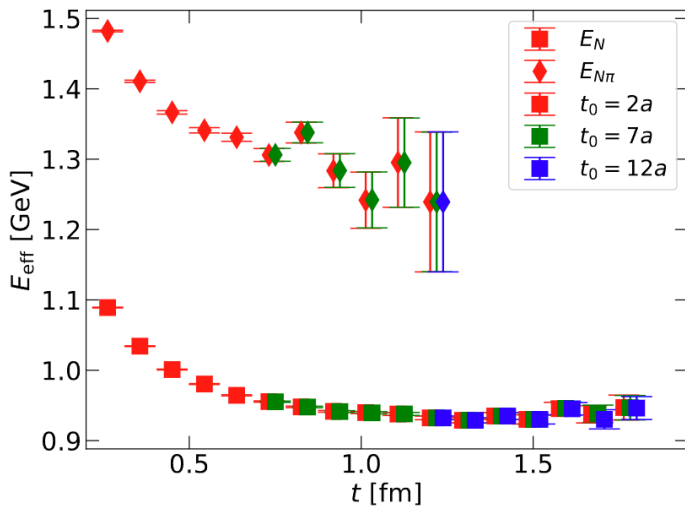
➤ We can use it to improve matrix elements:

$$\frac{\langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle}{\langle \mathcal{J}_N \mathcal{J}_N^\dagger \rangle} \xrightarrow{\text{GEVP improved}} \frac{\langle \tilde{\mathcal{J}}_N \mathcal{O} \tilde{\mathcal{J}}_N^\dagger \rangle}{\langle \tilde{\mathcal{J}}_N \tilde{\mathcal{J}}_N^\dagger \rangle}$$

t_0 dependence of GEVP

ALPHA Collaboration JHEP04(2009)094

$$C_{ij}(t)v_j^n = \lambda^n(t, t_0)C_{ij}(t_0)v_j^n$$



- E_{eff} are not sensitive to t_0
- Eigenvectors are sensitive to t_0
- With small t_0 , it converges to a wrong value
- In this work, we fix $t - t_0$, and do plateau fits to determine eigenvectors

3pt functions and GEVP improvement

$$\frac{\langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle}{\langle \mathcal{J}_N \mathcal{J}_N^\dagger \rangle} \xrightarrow{\text{GEVP improved}} \frac{\langle \tilde{\mathcal{J}}_N \mathcal{O} \tilde{\mathcal{J}}_N^\dagger \rangle}{\langle \tilde{\mathcal{J}}_N \tilde{\mathcal{J}}_N^\dagger \rangle} \rightarrow \begin{bmatrix} \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle & \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle \\ \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_N^\dagger \rangle & \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle \end{bmatrix}$$

$$\tilde{\mathcal{J}}_N = \mathcal{J}_N + v_{N\pi}^N \mathcal{J}_{N\pi}$$

➤ We compute everything except ~~$\langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle$~~

➤ RQCD: Last term is subleading by ChPT

Barca, Bali, Collins PRD 107, L051505 (2023)

Bar PRD 99, 054506 (2018) and 100, 054507 (2019)

➤ This work:

We found a new method that doesn't require such term

New method

➤ 3pt function without GEVP: $I_0 = \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle$

➤ Fully GEVP improved 3pt function:

$$I = \langle \tilde{\mathcal{J}}_N \mathcal{O} \tilde{\mathcal{J}}_N^\dagger \rangle = v_{N,N} v_{N,N}^* \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle + v_{N,N} v_{N,N\pi}^* \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle \\ + v_{N,N\pi} v_{N,N}^* \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_N^\dagger \rangle + v_{N,N\pi} v_{N,N\pi}^* \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle$$

➤ New method:

$$I_d = d_{N,N} v_{N,N} v_{N,N}^* \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_N^\dagger \rangle + d_{N,N\pi} v_{N,N} v_{N,N\pi}^* \langle \mathcal{J}_N \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle \\ + d_{N\pi,N} v_{N,N\pi} v_{N,N}^* \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_N^\dagger \rangle + 0 \times v_{N,N\pi} v_{N,N\pi}^* \langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle$$

- Coefficients d can be determined by GEVP, do not depend on the insertion operator \mathcal{O}
- I_d can remove the leading contamination terms

$$R_0 = \langle N | \mathcal{O} | N \rangle + a(e^{-\Delta E t_{\text{ins}}} + e^{-\Delta E (t_{\text{sink}} - t_{\text{ins}})}) + b e^{-\Delta E t_{\text{sink}}}$$

$$R = \langle N | \mathcal{O} | N \rangle$$

$$R_d = \langle N | \mathcal{O} | N \rangle - b e^{-\Delta E t_{\text{sink}}}$$

Overview of results

$$(\vec{p}_{sink}, \vec{p}_{src}) = (\vec{0}, \vec{0}), (\vec{0}, \vec{1}), (\vec{1}, \vec{1})$$



- We have investigated on **52 cases**
 $\{u+d \text{ \& } u-d\} \times \{S, V_\mu, P, A_\mu, \sigma_{\mu\nu}\} \times \{\text{various kinematics}\}$
- We include both connected and disconnected contributions for both isoscalar & isovector
(disc u-d is nonzero for tm1QCD)
- **46 cases**
 - No significant changes observed
 - Including **the sigma term** $\sigma_{\pi N}$
- **6 cases:** $(u-d) \times \{P, A_\mu\}$
 - Isovector pseudoscalar (**2**) and axial (**4**) currents
 - Significant changes observed
 - Insertion operator has same quantum number with the pion

Axial charge g_A^{u-d}

Open: no GEVP Filled: GEVP
 Grey: ETMC23, PRD 109 (2024) 3, 034503
 Dashed black: experiment

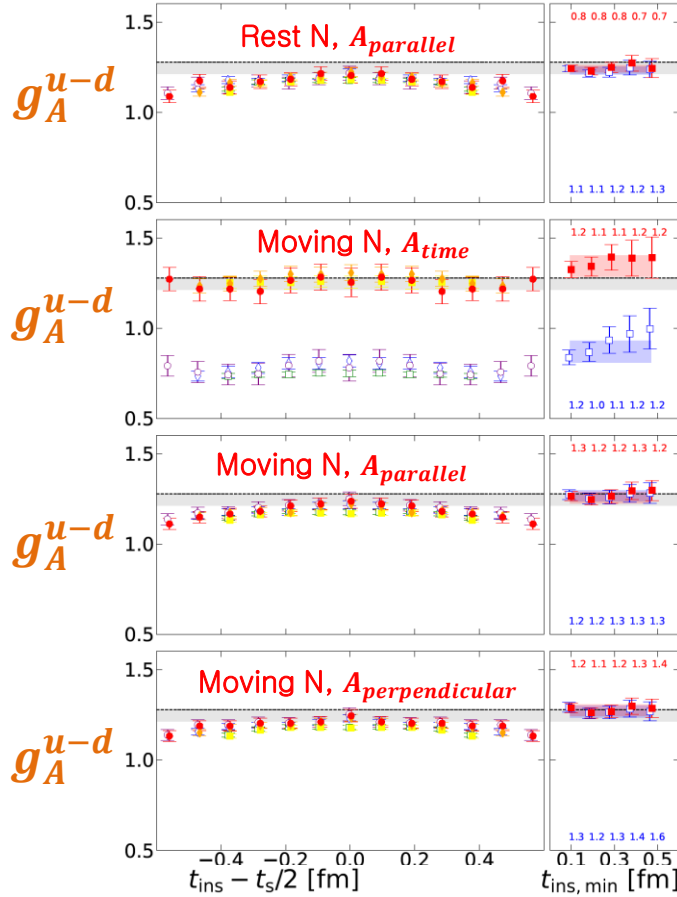
➤ Four cases for g_A^{u-d}
 Ratio (left) and two-state fits (right) =>

- Before GEVP: Open
 - After GEVP: Filled
 - ETMC23: Band
- (Excited state analysis & continuum limit)
 Talk by G. Koutsou, Monday @ 15.35

Small lattice artefact for g_A^{u-d}

- No changes for 3 cases
- GEVP brings 2nd case agreement with the other 3

This work: 1.258(18)
 Exp: 1.27641(56)



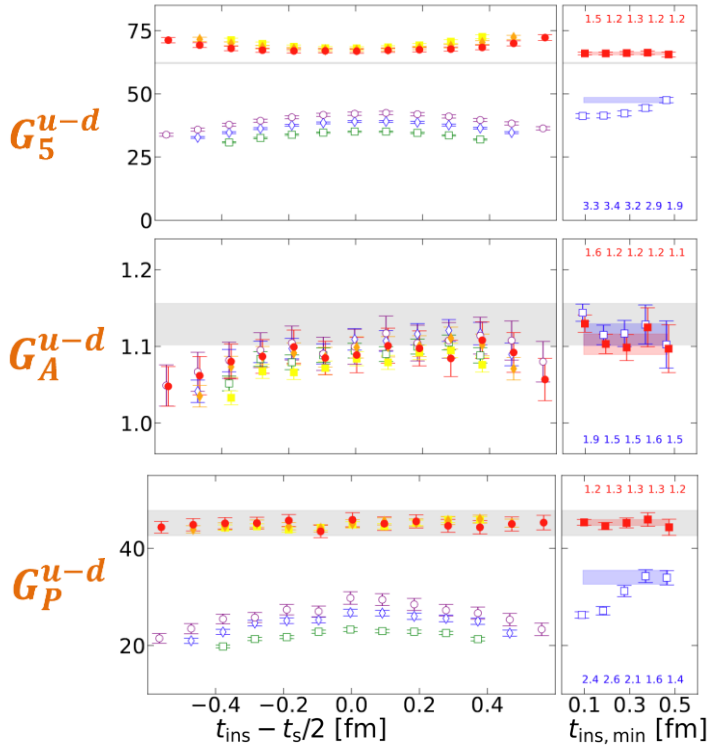
PCAC related quantities

@ 1-unit transfer

$$G_5^{u-d}, G_A^{u-d}, G_P^{u-d}$$

- Significant improvement observed for G_5^{u-d}, G_P^{u-d}
- Large lattice artefact expected for G_5^{u-d}
- Small lattice artefact for G_P^{u-d} after including the isovector insertion loop (nonzero for tm1QCD)

Open: no GEVP Filled: GEVP
 Grey: ETMC23, PRD 109 (2024) 3, 034503



Conclusions

- Strong t_0 dependence for GEVP eigenvectors
- New method without requiring $\langle \mathcal{J}_{N\pi} \mathcal{O} \mathcal{J}_{N\pi}^\dagger \rangle$
- We investigate on 52 cases = 46 (no) + 6 (yes)
 - 46 includes $\sigma_{\pi N}$
 - 6 are with isovector pseudoscalar & axial currents
- Reduced lattice artefacts with isovector insertion loop

THANKS



Με τη συγχρηματοδότηση
της Ευρωπαϊκής Ένωσης



Κυπριακή Δημοκρατία



ΙΔΡΥΜΑ
ΕΡΕΥΝΑΣ ΚΑΙ
ΚΑΙΝΟΤΟΜΙΑΣ

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