

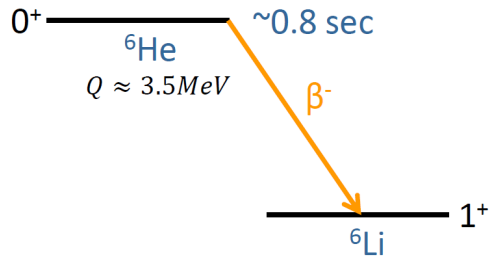


Exploring Nuclear Beta Decay Through Lattice Effective Field Theory

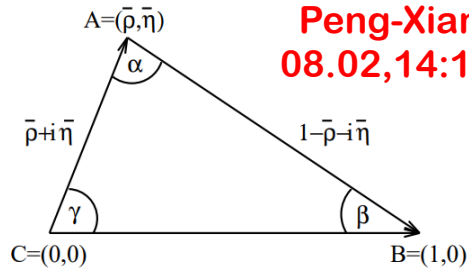
Teng Wang 2024.8.2
Collaborate with **Xu Feng** and **Bing-Nan Lu**

Nuclear β -Decay on the Lattice: Opportunities and Challenges

★ Nuclear β -decay is related to various fundamental problems in physics.

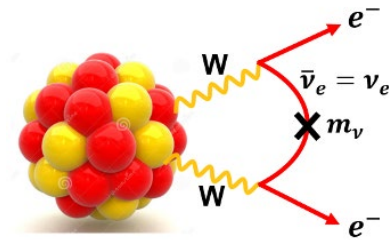


${}^6\text{He}$ β -decay spectrum and exotic currents



super allowed β -decay and CKM unitarity

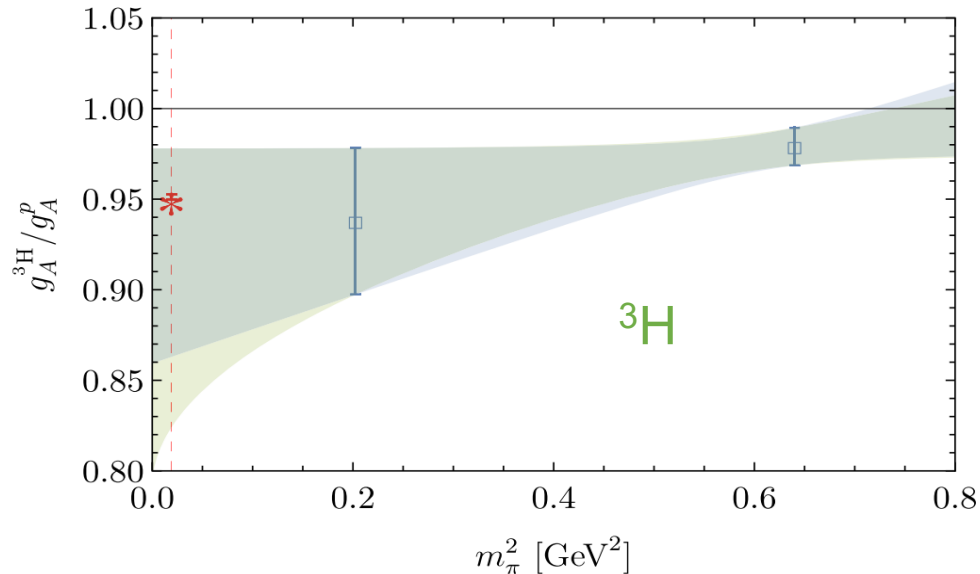
Peng-Xiang Ma's talk
08.02, 14:15-14:35, LT3



neutrinoless double β -decay and neutrino physics

Zi -Yu Wang's talk
07.31, LT1

★ Lattice QCD has made impressive progresses on the β -decay of few-body nuclei.



M. J. Savage et al, Phys. Rev. Lett. 119, 062002 (2017)

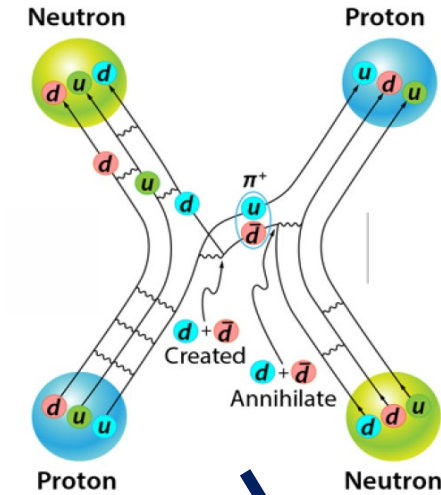
A. Parreño et al, Phys. Rev. D 103, 074511 (2021)

What about heavier nuclei?

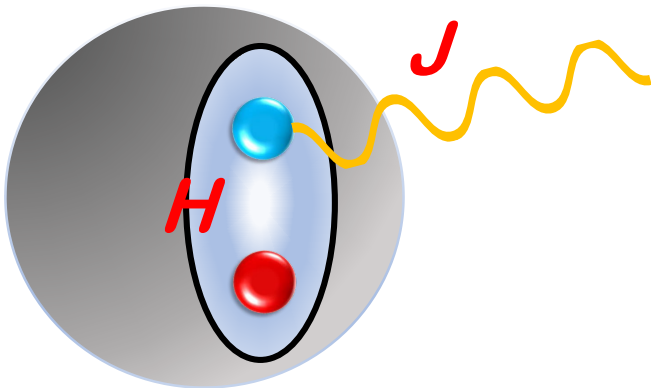
From QCD to Chiral Effective Field Theory

★ Chiral Effective field theory(EFT) provides a shortcut to the problem.

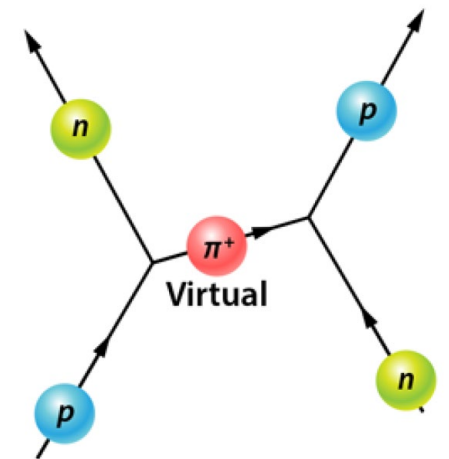
- Low energy effective field theory of QCD, works for $Q \ll \Lambda_\chi \approx 1\text{GeV}$
- Nucleons and pions as DOFs.
- Spontaneously broken chiral symmetry $SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$
- Order by order calculation through power counting.



★ Chiral EFT can provide the interaction of a nucleon with another nucleon inside the nucleus and with external fields.



- Chiral EFT can well describe nuclear β -decay.
- How to solve the nuclear many-body problem?



Lattice Effective Field Theory: an Introduction

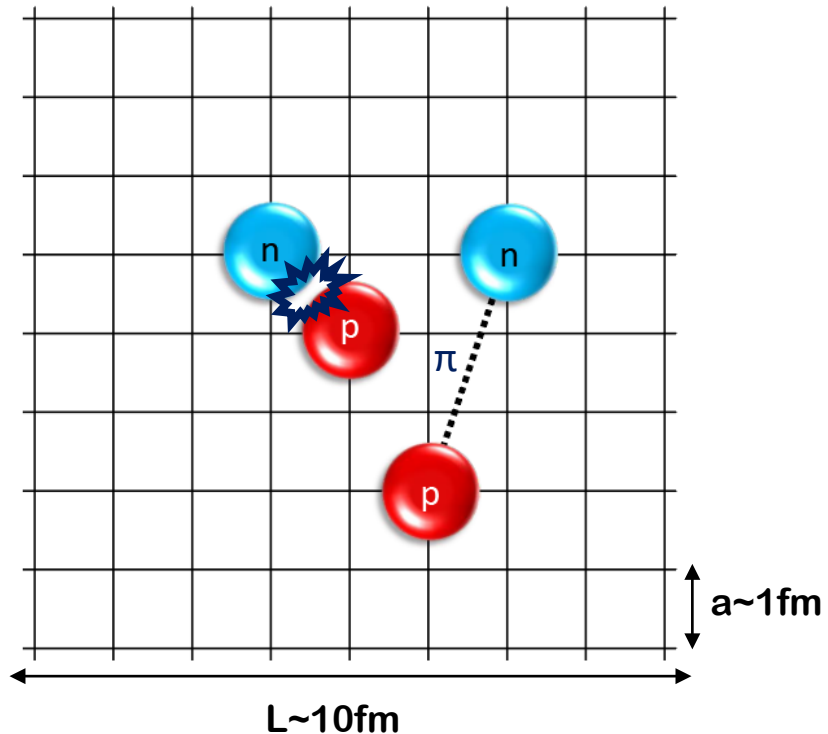
★ Lattice effective field theory(Lattice EFT) is a many body method to study nucleus on the lattice.

Lattice EFT = Chiral EFT + Lattice Method

Review: Dean Lee, Prog. Part. Nucl. Phys. 63, 117 (2009)

Lähde, Meißner, “Nuclear Lattice Effective Field Theory”, Springer (2019)

★ Main features of Lattice EFT.



- Discretize space-time into an $L_t \times L_x \times L_y \times L_z$ lattice.
- Nucleons are point-like particles on the site.
- Discretize the nuclear force, including the strong and Coulomb force on the lattice.
- lattice spacing $a \sim 1\text{fm}$ and lattice size $L \sim 10\text{fm}$.

An Analogy to Lattice QCD

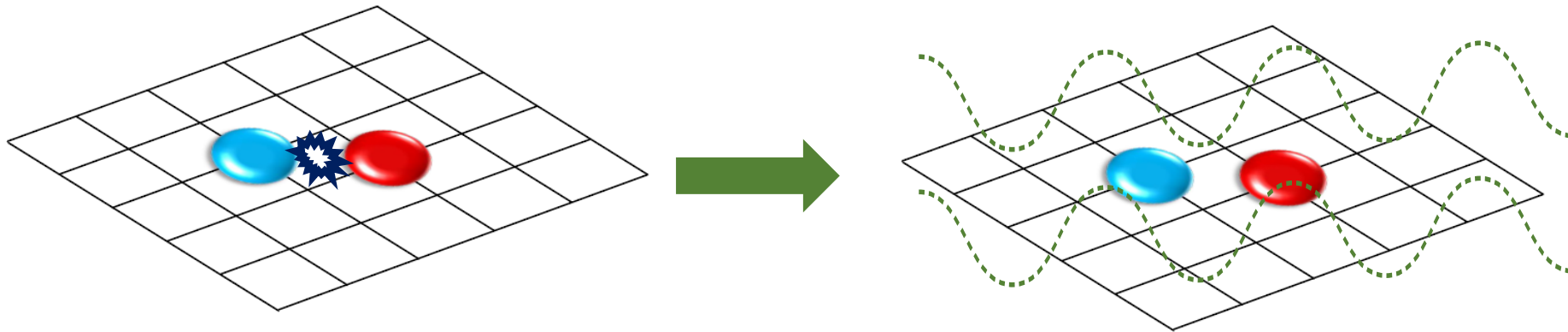
★ Calculations in Lattice EFT are quite similar to Lattice QCD.

- States are evolved via **Euclidean time projection**.

$$|\Phi_g\rangle = \lim_{T \rightarrow \infty} e^{-HT} |\Phi_i\rangle$$

- The NN interaction is decoupled into fermion-boson interaction via the **auxiliary field transformation**.

$$\exp \left[-\frac{C}{2} (N^\dagger N)^2 \right] = \sqrt{\frac{1}{2\pi}} \int ds \exp \left[-\frac{s^2}{2} + \sqrt{C} s (N^\dagger N) \right]$$



- The multi-dimensional integral over the auxiliary field is evaluated through **Monte Carlo**.

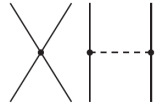
$$\langle \Phi | e^{-HT} | \Phi \rangle = \int \mathcal{D}s e^{-\frac{1}{2}s^2} Z[s]$$

Chiral Nuclear Force

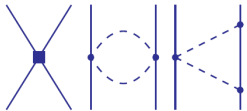
★ We use NNLO chiral nuclear force.

LO
(Q/Λ_χ)⁰

2N Force



NLO
(Q/Λ_χ)²



NNLO
(Q/Λ_χ)³



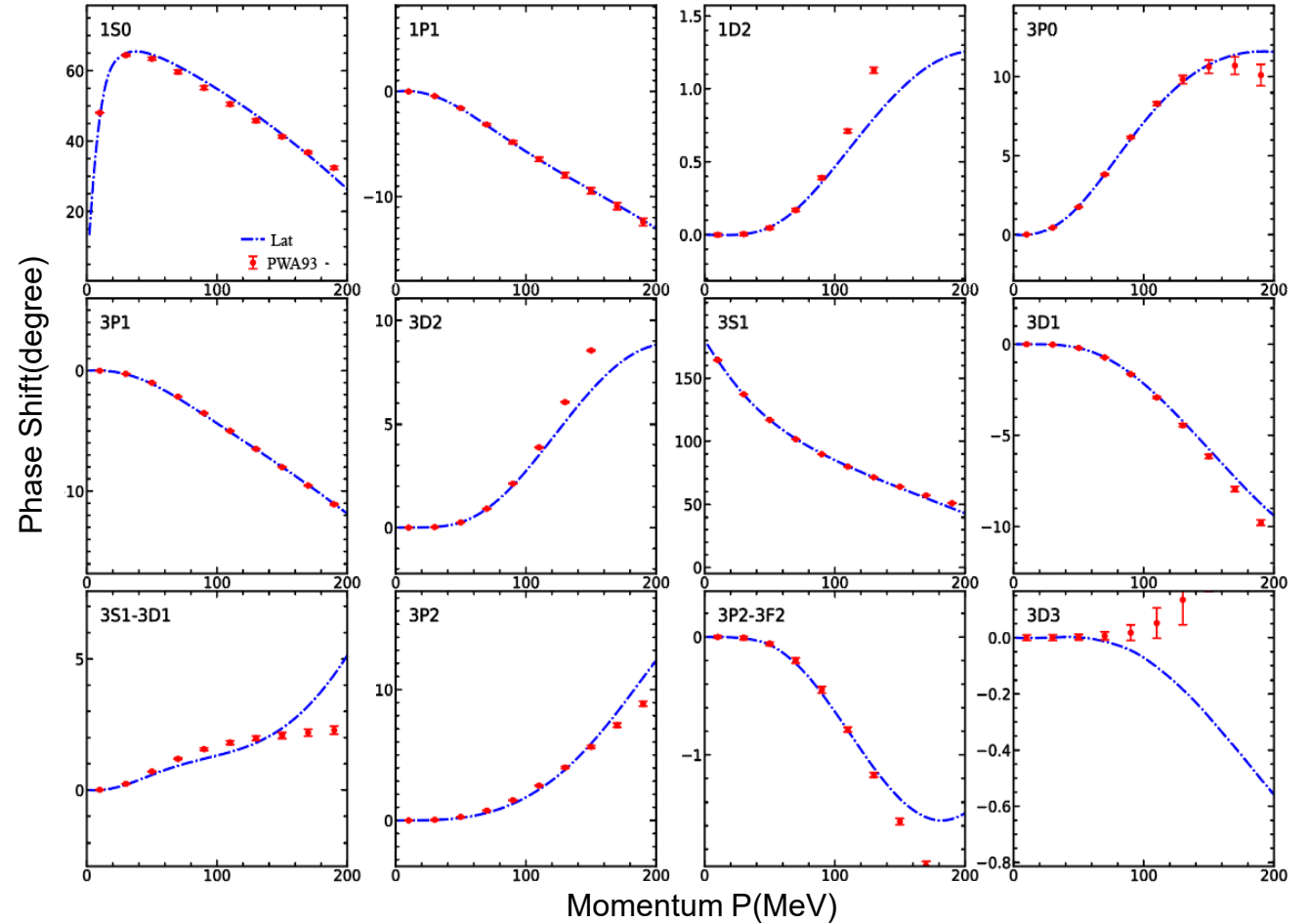
3N Force

R. Machleidt et al,
Phys.Rept. 503, 2011

$$\begin{aligned}
 V_{2N} = & \left[B_1 + B_2(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) + C_1 q^2 + C_2 q^2(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2) \right. \\
 & + C_3 q^2(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) + C_4 q^2(\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2)(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2) \\
 & + C_5 \frac{i}{2}(\mathbf{q} \times \mathbf{k}) \cdot (\boldsymbol{\sigma}_1 + \boldsymbol{\sigma}_2) + C_6(\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q}) \\
 & \left. + C_7(\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q})(\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2) \right] \\
 & - \frac{g_A^2 f_\pi(q^2)}{4F_\pi^2} \left[\frac{(\boldsymbol{\sigma}_1 \cdot \mathbf{q})(\boldsymbol{\sigma}_2 \cdot \mathbf{q})}{q^2 + M_\pi^2} + C'_\pi \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 \right] (\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2)
 \end{aligned}$$

9 LECs to fit, 10 data points for each channel

★ The 2N LECs are fitted to PWA93 n-p scattering phase shifts



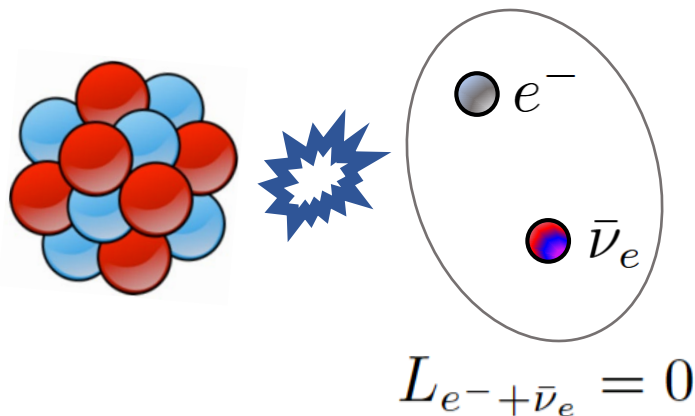
experiment data: <https://nn-online.org>

Fitting Method: B. N. Lu et al, Phys. Lett. B 760(2016)

Gamow Teller Matrix Element and Axial Current

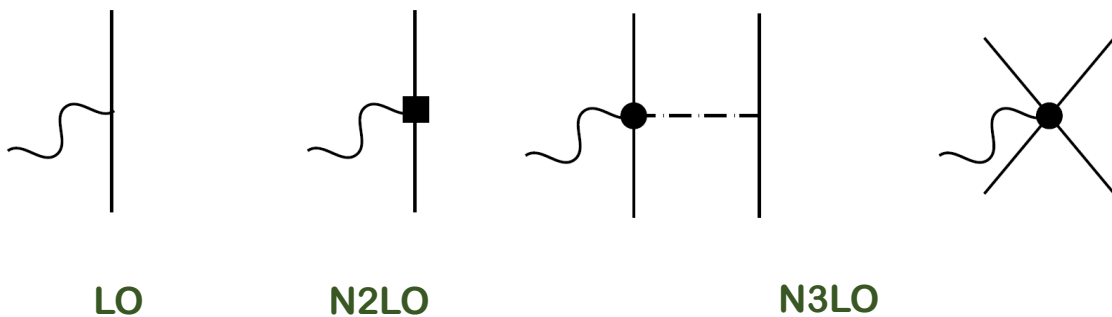
★ We focus on allowed nuclear β -decay.

★ There are 2 types of matrix elements, Fermi and Gamow-Teller type, with the later more interesting.



Fermi	$\langle \Psi_f J_V^0 \Psi_i \rangle$	constrained by CVC
GT	$\langle \Psi_f J_A^i \Psi_i \rangle$	sensitive to NS

★ We use nuclear axial currents up to N3LO consistently derived from chiral EFT

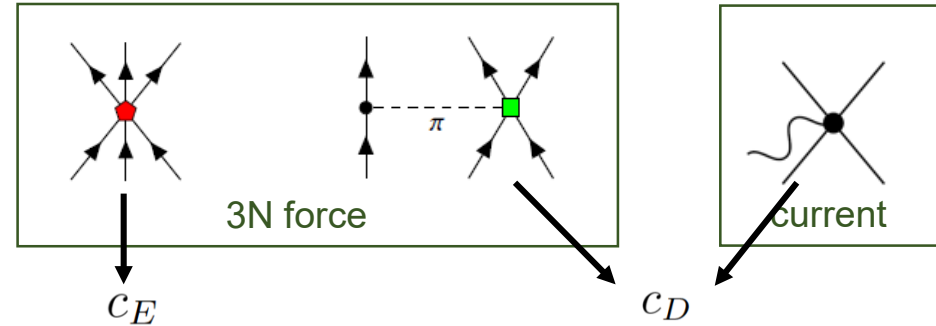


A. Baroni et al, Phys. Rev. C, 93 015501, 2016

H. Krebs et al, Annals of Physics 378, 2017

Determine 3N LECs: ${}^3\text{H}$ binding energy and ${}^3\text{H}$ β -decay

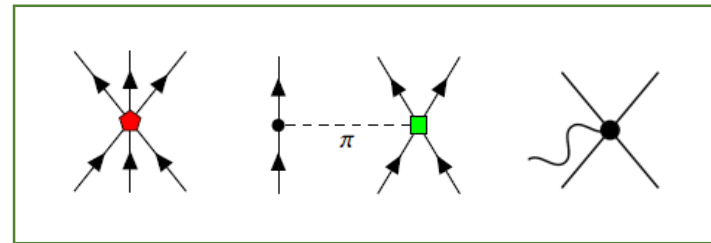
- ★ There are two unknown LECs, c_D and c_E , in the 3N force, the former also appears in the N3LO contact axial current.



- ★ ${}^3\text{H}$ is an ideal system to help determine the LECs.

$$E({}^3\text{H}) = -8.482\text{MeV}$$

$$\text{GT}({}^3\text{H} \rightarrow {}^3\text{He}) = 0.9511(13)$$



$$c_D = 1.01(12)$$

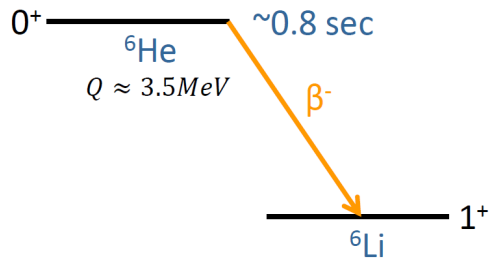
$$c_E = 0.81(1)$$

- ★ Contributions to the Gamow-Teller matrix element (GTME) from currents at different orders.

LO	N2LO	N3LO(OPE)	N3LO(CT)
0.9614(86)	-0.0041(11)	-0.0311(34)	0.0146(28)

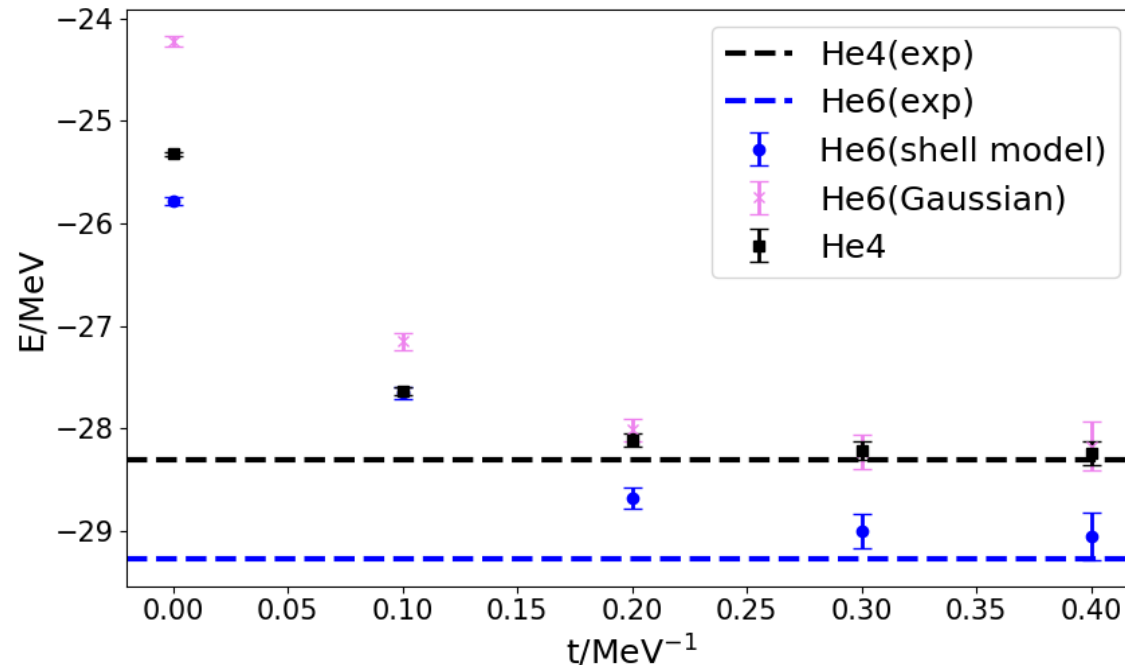
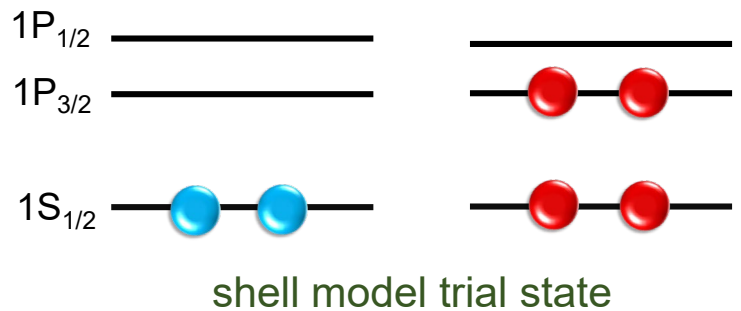
Predicting ${}^6\text{He}$ β -decay: Trial States' Preparation

- ★ We validate our method through calculating the GTME of ${}^6\text{He} \rightarrow {}^6\text{Li}$ β -decay process.



$$\text{GT}(t) = \frac{\sqrt{2J_f + 1}}{g_A} \frac{\langle \Phi_f | e^{-\frac{1}{2}Ht} \mathcal{J} e^{-\frac{1}{2}Ht} | \Phi_i \rangle}{\sqrt{\langle \Phi_i | e^{-Ht} | \Phi_i \rangle \langle \Phi_f | e^{-Ht} | \Phi_f \rangle}}$$

- ★ We find the convergence rate to ${}^6\text{He}$ ground depends strongly on the choice of the trial state.



Predicting ${}^6\text{He}$ β -decay: Overcoming Sign Problem

★ Among various competing uncertainties, statistical errors induced by the sign problem is the dominant.

● Monte Carlo Statistical error

● Model dependence

● Finite volume effect...

$$\int \mathcal{D}s e^{-\frac{1}{2}s^2} \underbrace{Z[s]}_{\text{complex}}$$

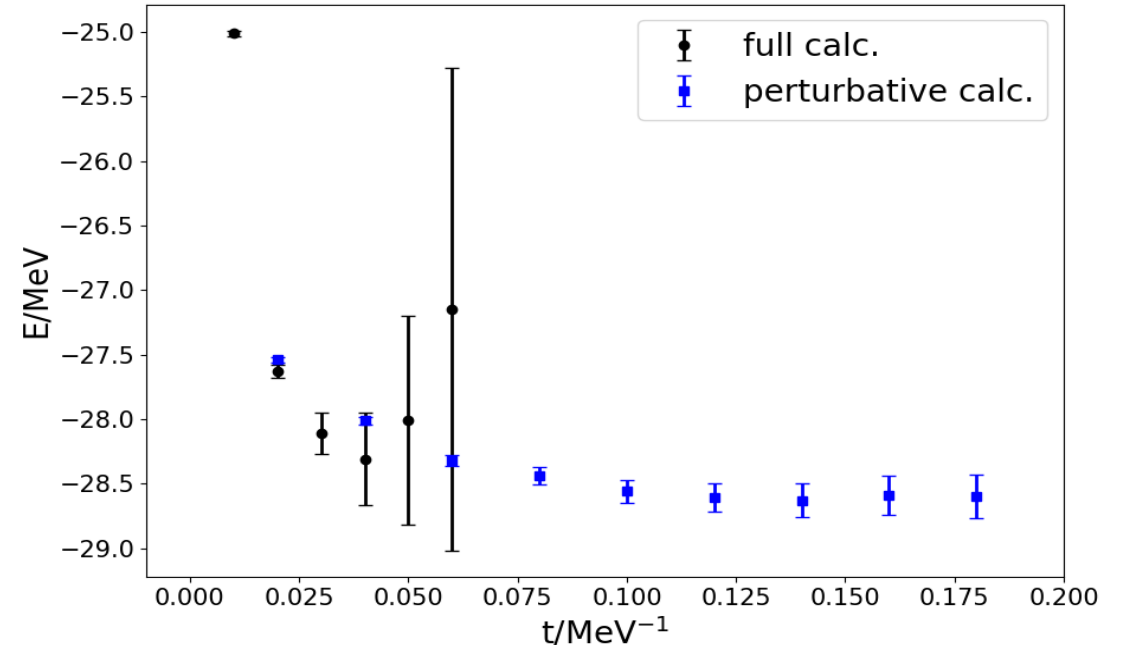
★ Wigner SU4 symmetry can help reduce sign problem

$$N = \begin{bmatrix} a_{p\uparrow} \\ a_{p\downarrow} \\ a_{n\uparrow} \\ a_{n\downarrow} \end{bmatrix}$$



E. Wigner, Phys. Rev. 51, 106 (1937)

J.-W. Chen et al, Phys. Rev. Lett. **93**, 242302 (2004)



★ Making use of the SU4 symmetry, we ease the sign problem using perturbative method

$$H = H_{\text{LO}} + \Delta V$$

$$\text{GT} = \text{GT}^{(0)} + \text{GT}^{(1)} + \mathcal{O}(\Delta V^2)$$

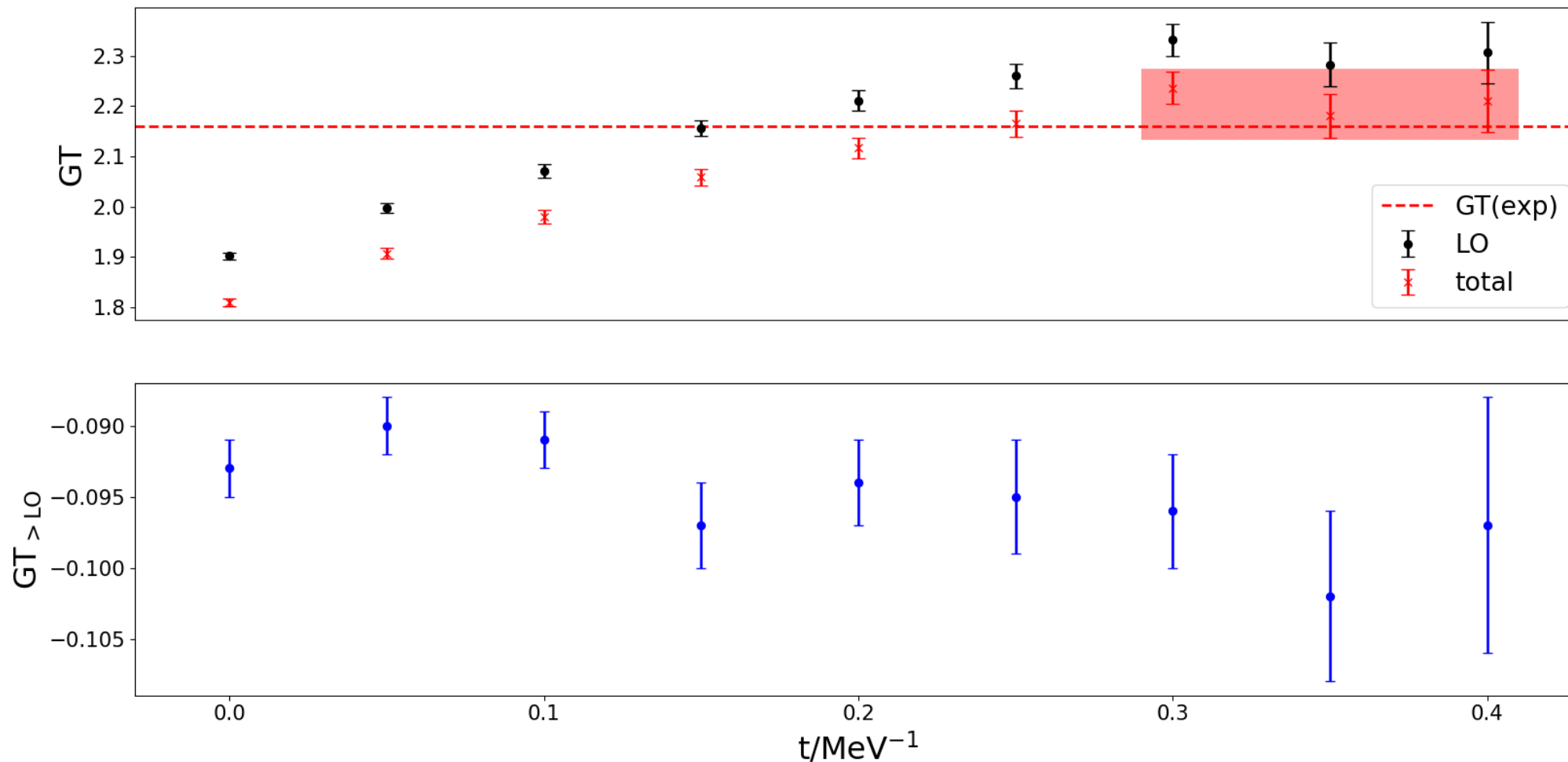
$$H_{\text{LO}} = K + \underbrace{\frac{C_2}{2!} \sum_n : \tilde{\rho}^2(\mathbf{n}) :}_{\text{SU4 symmetric}} + \frac{C_3}{3!} \sum_n : \tilde{\rho}^3(\mathbf{n}) : + \underbrace{\frac{C_I}{2!} \sum_{I,n} : \tilde{\rho}_I^2(\mathbf{n}) :}_{\text{SU4 breaking}} + V_{\text{OPE}}^{\Lambda'_\pi}$$

B. N. Lu et al, Phys. Rev. Lett. 128, 242501 (2022)

SU4 symmetric

SU4 breaking

Predicting ${}^6\text{He}$ β -decay: Result



	LO	N2LO	N3LO(OPE)	N3LO(CT)	Total-LO	Total	Exp
GT	2.304(64)	-0.024(9)	-0.105(1)	0.030(2)	-0.099(10)	2.205(65)	2.161(4)

Summary and Outlook

- ★ We calculated nuclear β -decay through lattice EFT **for the first time**.
- ★ **3N LECs** are determined from ${}^3\text{H}$ binding energy and β -decay.
- ★ **${}^6\text{He}$ β -decay** is predicted in good consistency with the experiment.
- ★ Tests for heavier nuclei are required.
- ★ A direct application on various EW processes in the nuclei is straightforward.

thank You!