

Towards glueball scattering in lattice Yang-Mills theory

Maxwell T. Hansen

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with Mattia Bruno and Antonio Rago

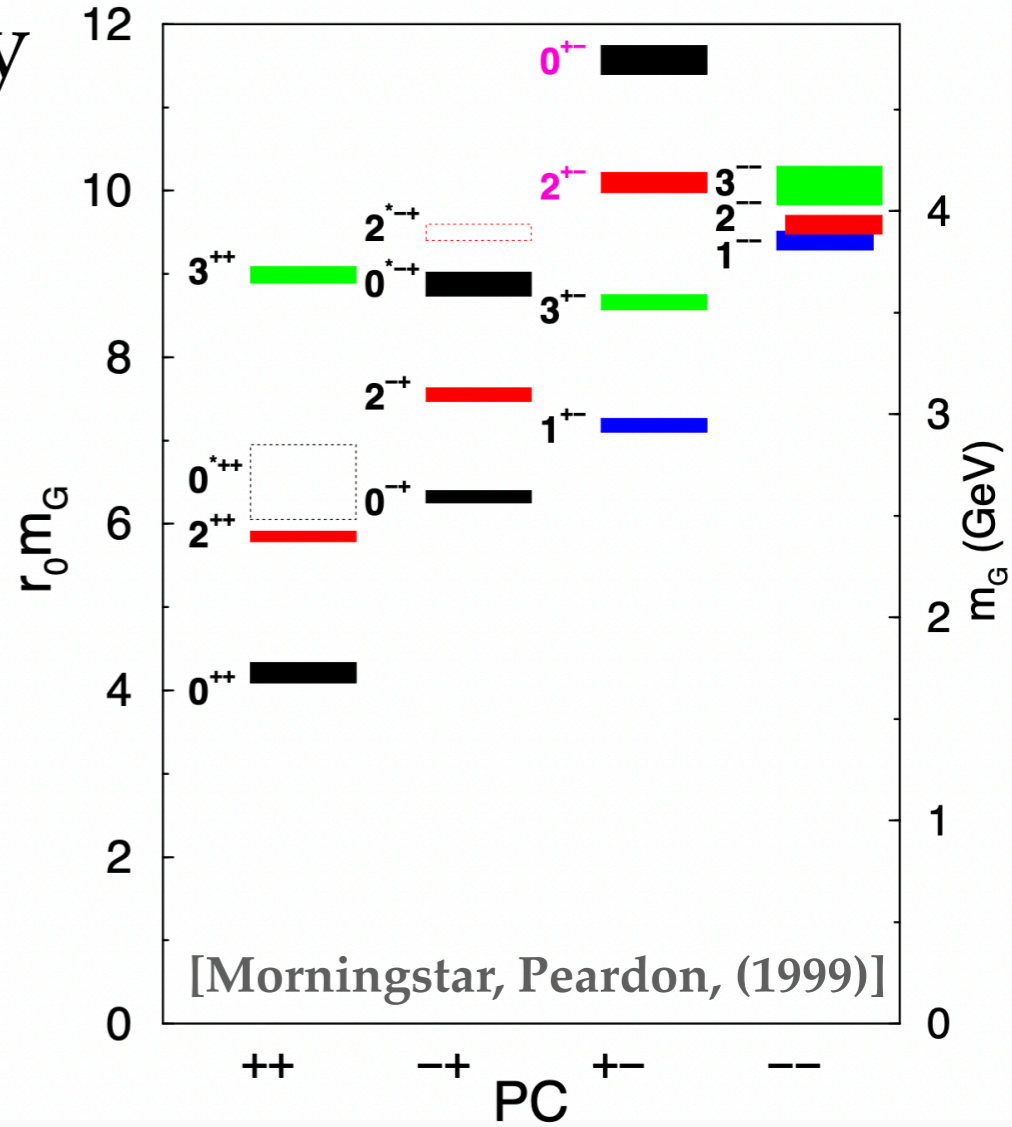


Scattering glue balls in Yang-Mills theory

➤ In Yang-Mills theory, low-lying spectrum consists of stable glue balls

$$\mathcal{L}_{\text{YM}} = -\frac{1}{4} F^{a\mu\nu} F_{\mu\nu}^a$$

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Can we use standard finite-volume methods to calculate glue-ball scattering amplitudes?



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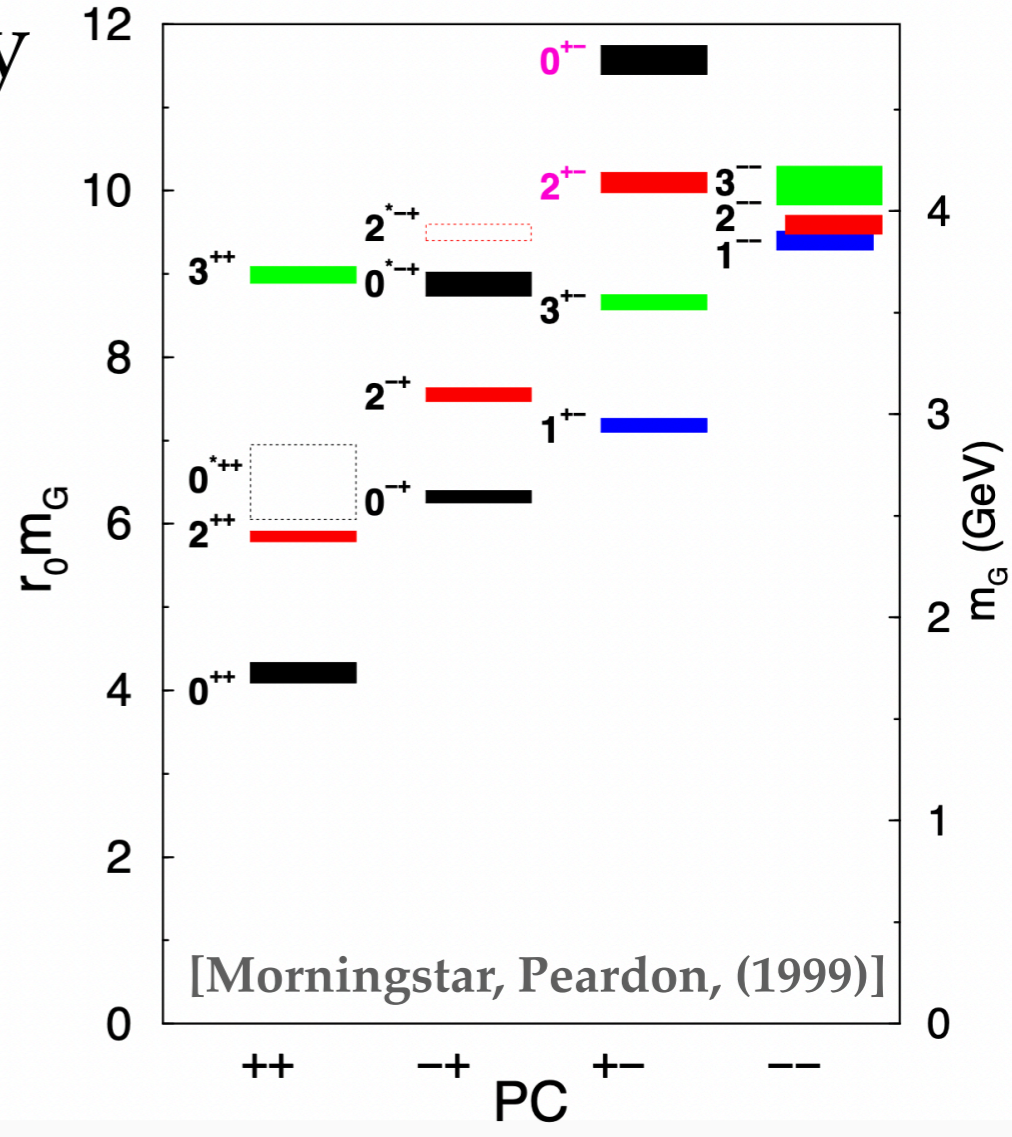
➤ Motivations

constrain strongly-interacting dark-matter models

test S-matrix-bootstrap predictions

better understand QCD glue balls in a *theoretically clean* environment

experimental evidence that X(2370) has quantum numbers 0^{-+}
 (from $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta'$)



[BES III (2024)]

Challenges for scattering glue balls

- ▶▶ Glue-ball correlators suffer from a signal-to-noise problem
- ▶▶ Require a large and varied set of operators to reliably extract excited states
- ▶▶ Window problem: finite-volume must be...
 - large enough to ignore unwanted effects
 - small enough to resolve effects of interest

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related to fact that YM = *a single-scale theory*
- ▶▶ Missing the analogous quark model → less obvious operator construction
- ▶▶ Finite-volume spectrum is more complicated than, say pions
 - one-, two-, three-particle energies have the same quantum numbers
 - di-torelon states could contaminate spectrum (more later)

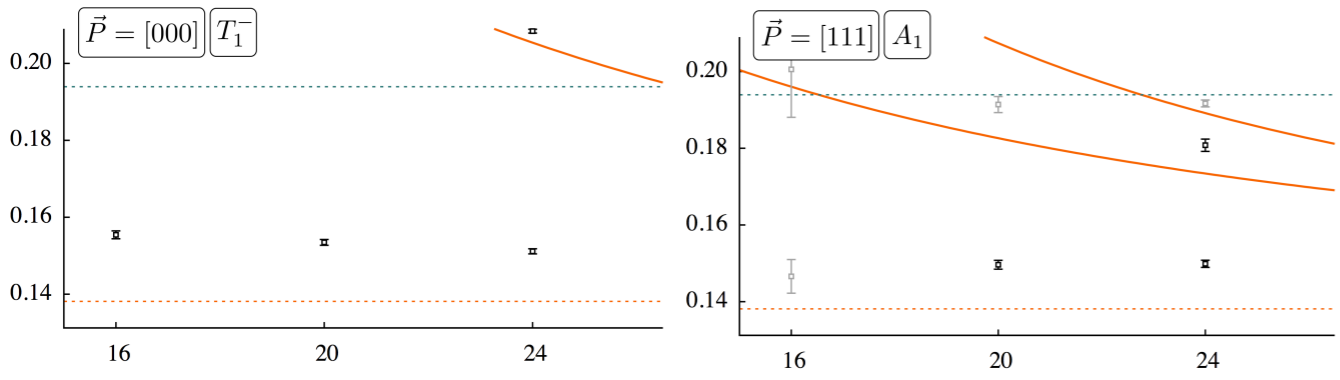
Finite-volume scattering (1 / 2)

➤ Standard in lattice QCD to use finite-volume energies to extract scattering amplitudes

scattering geometric function

$$\rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

equal at finite-volume energies



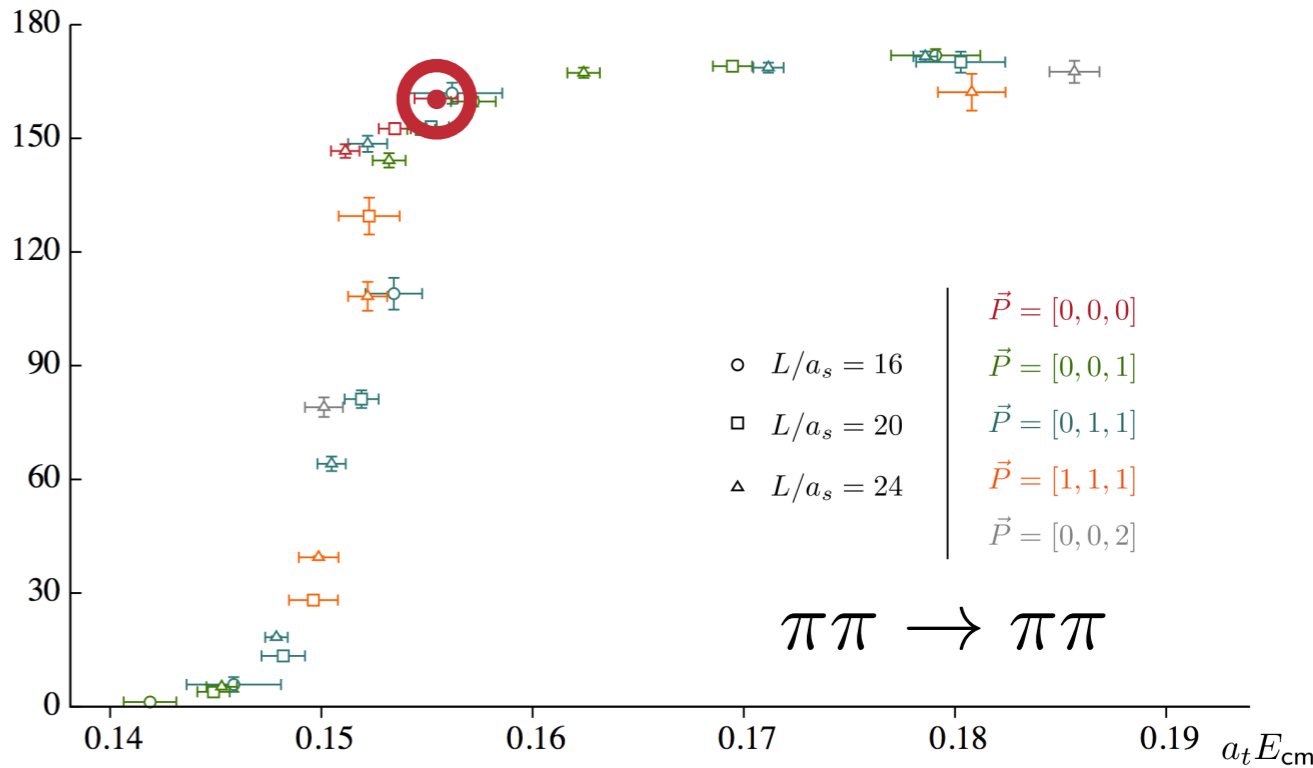
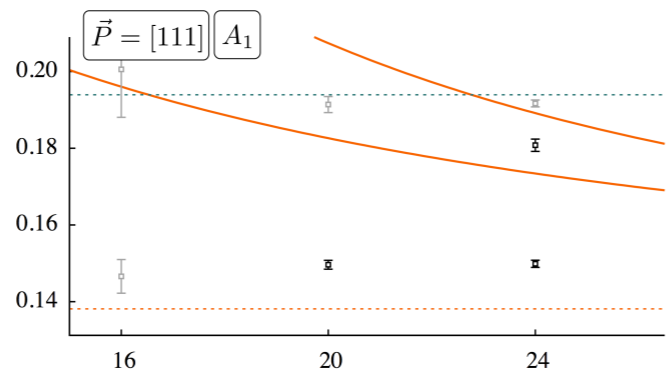
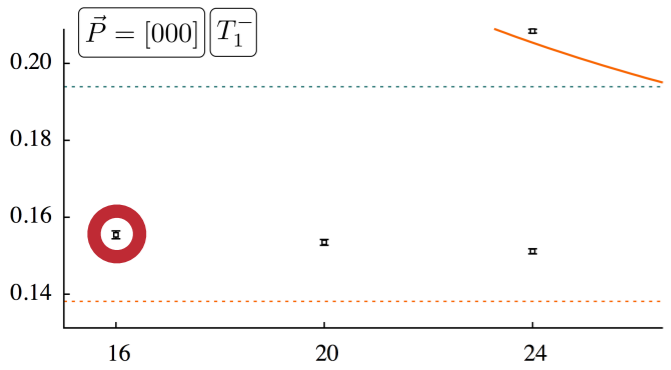
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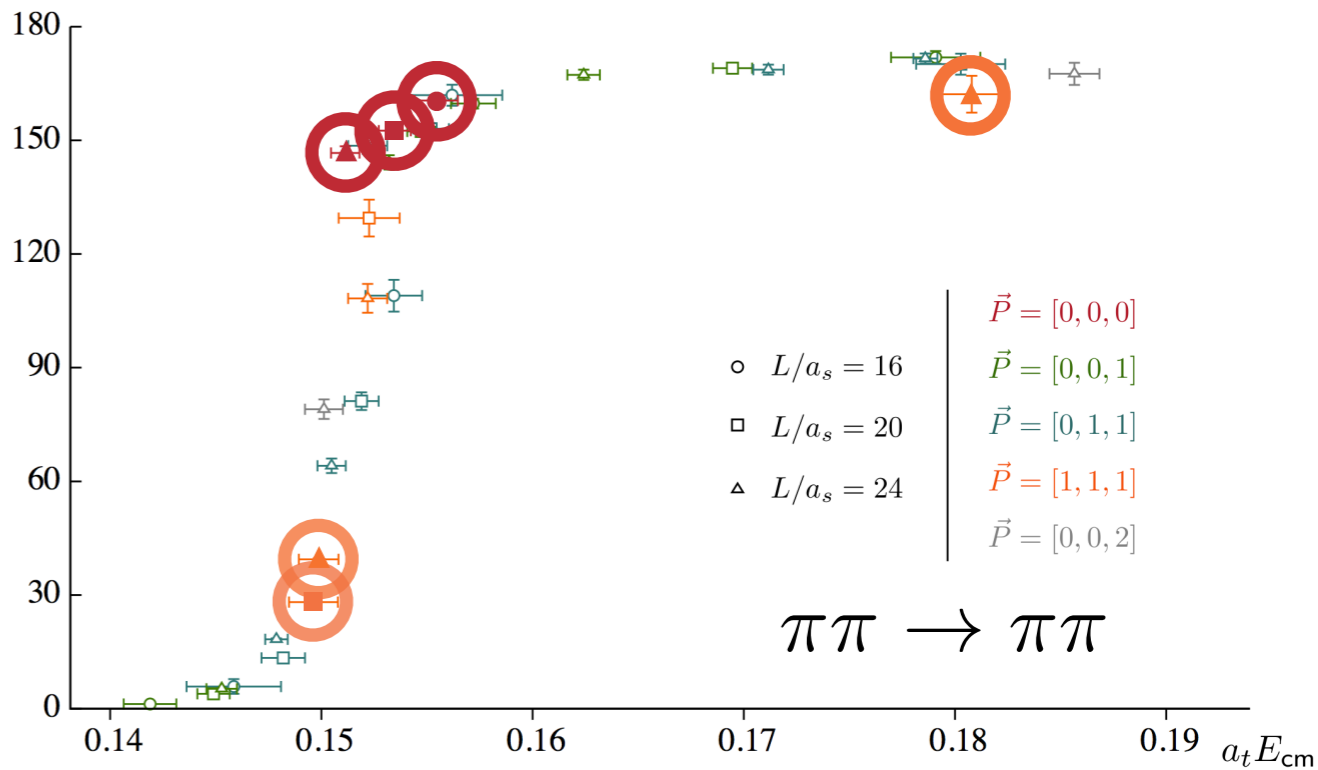
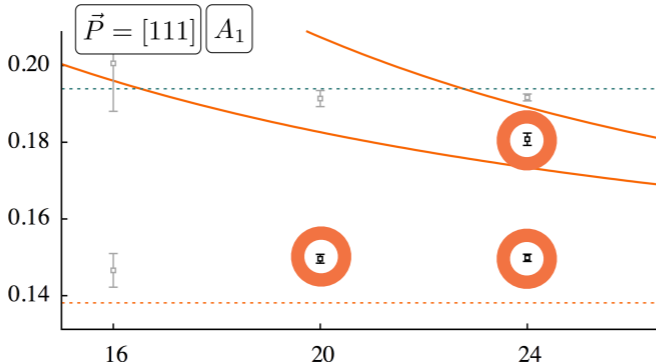
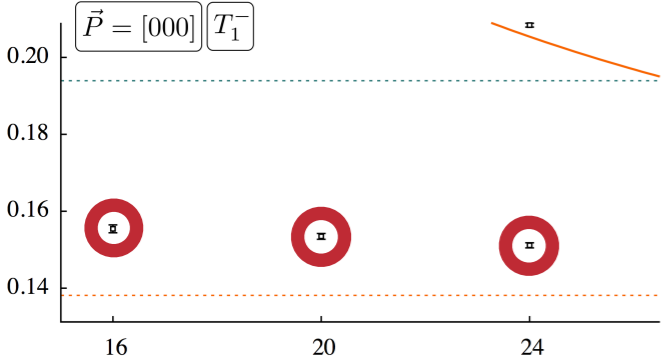


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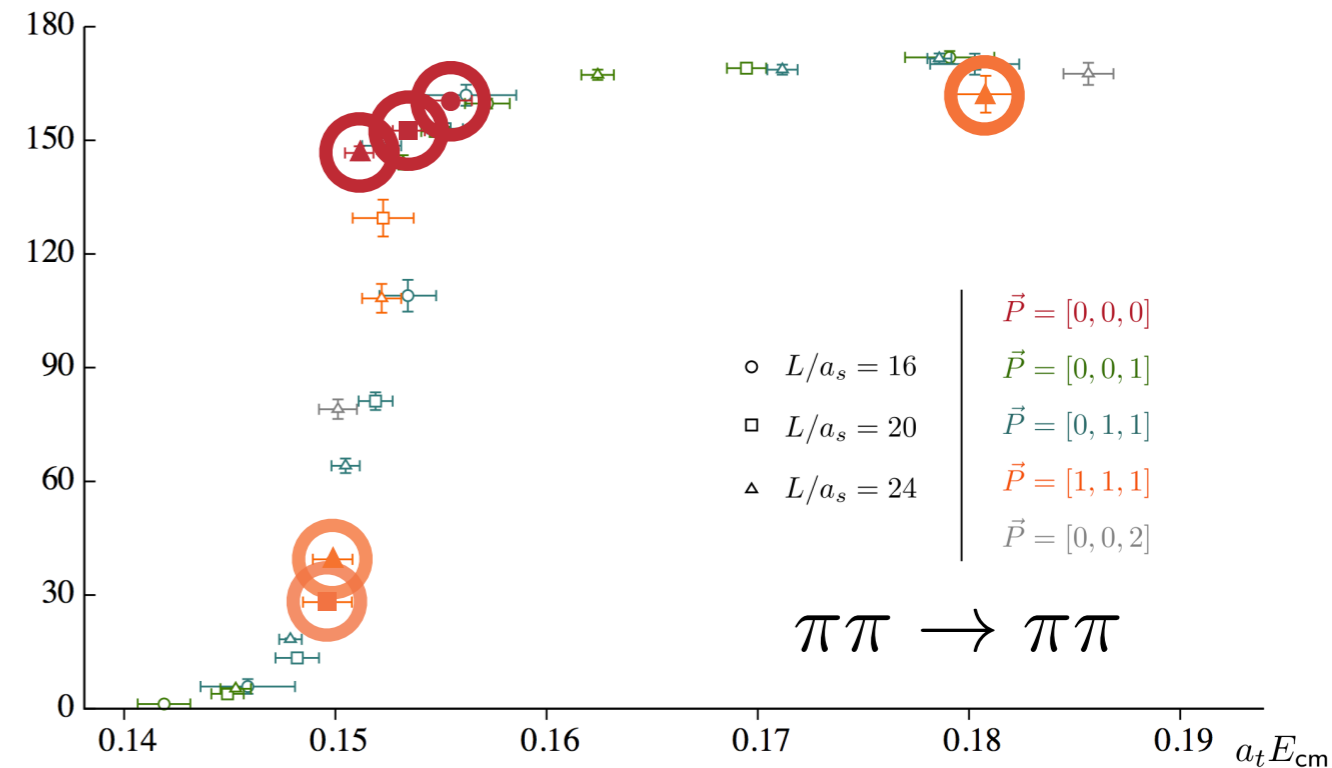
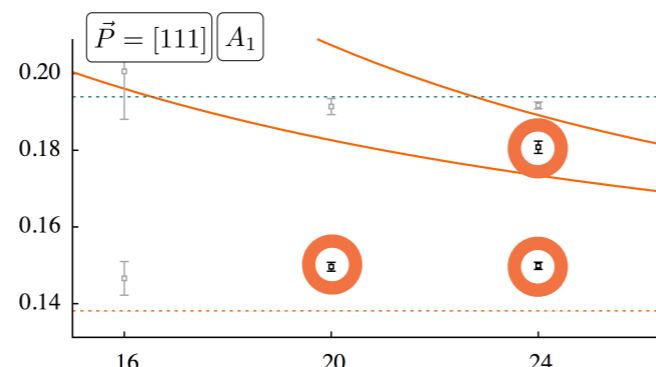
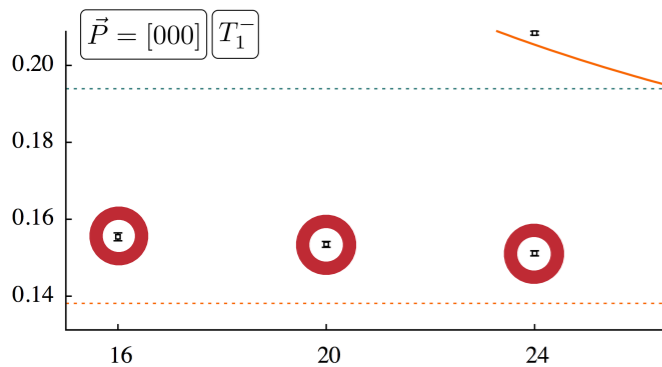
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➤ Goal: Calculate finite-volume energies in Yang-Mills to extract

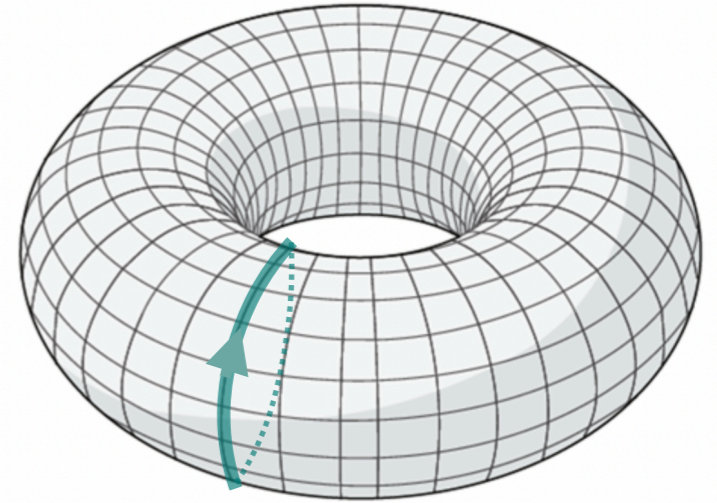
$GG \rightarrow GG$ amplitudes

($G = 0^{++}$ ground state)

Expectations for the finite-volume energies

➤ *torelon* = closed flux tube, wrapping spatial direction

$$E_{\text{tor}}(L) = \sigma L + \mu + \mathcal{O}(1/L)$$



Expectations for the finite-volume energies

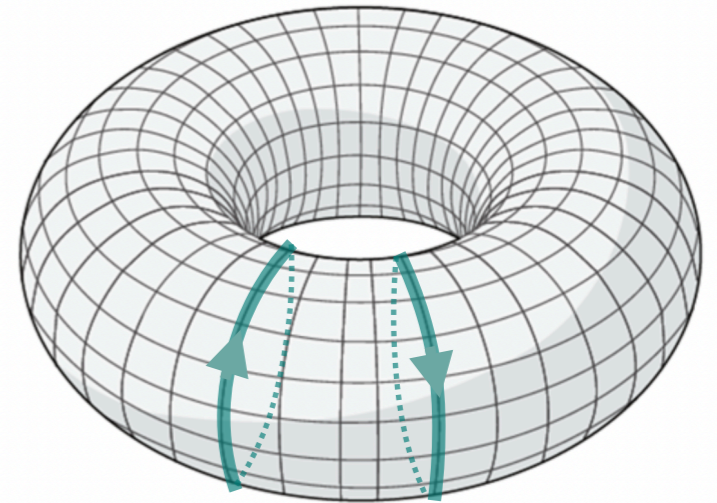
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same quantum numbers as local 0^{++} states



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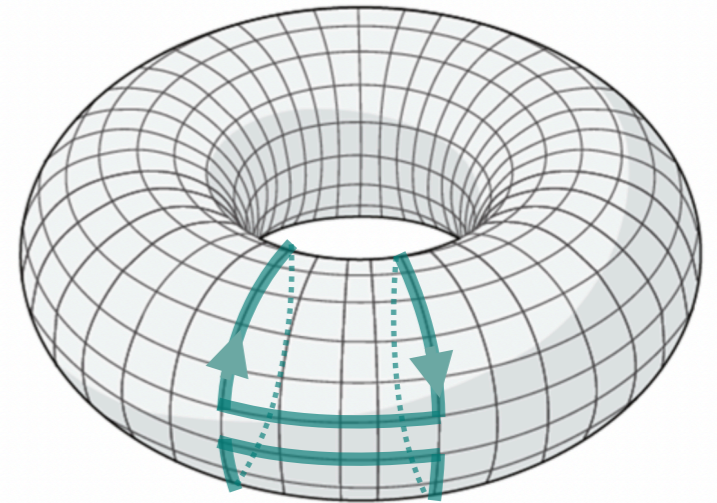
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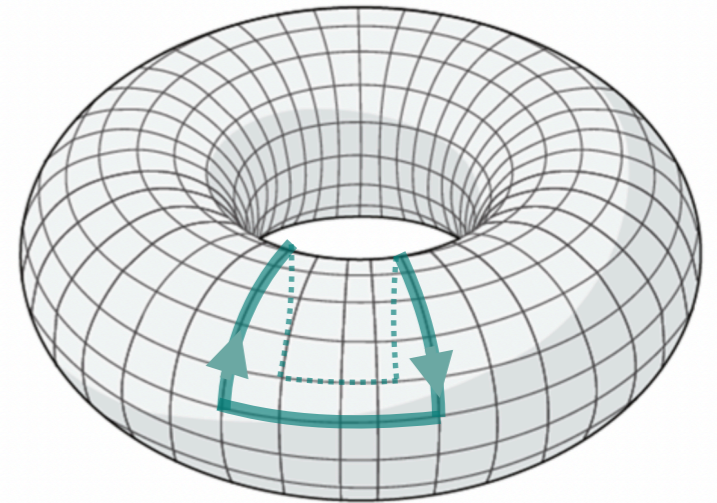
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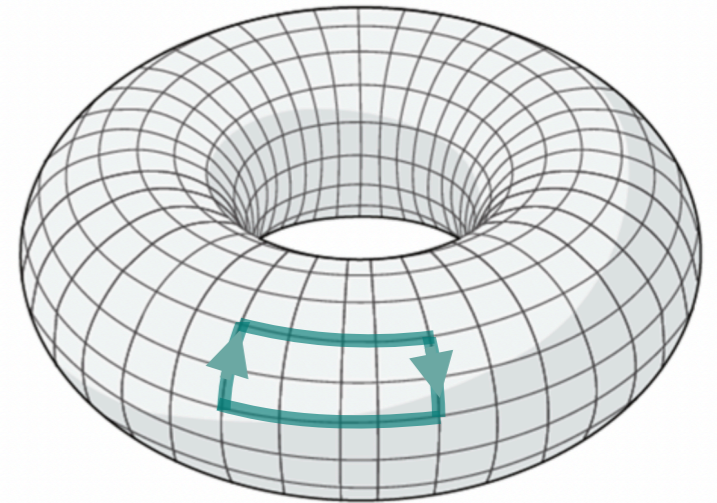
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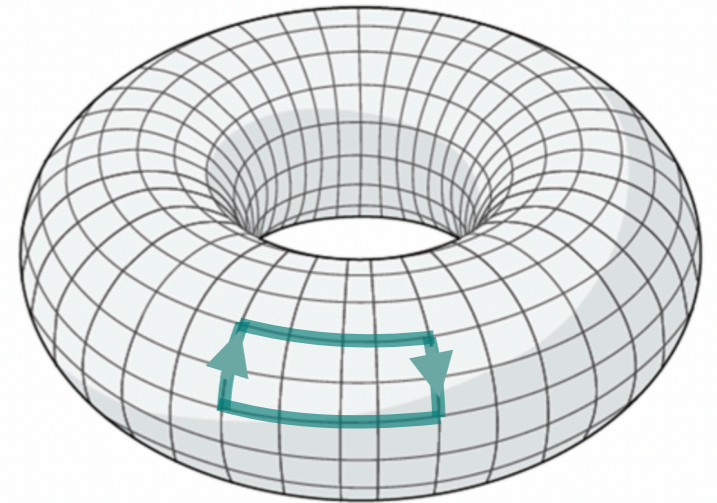
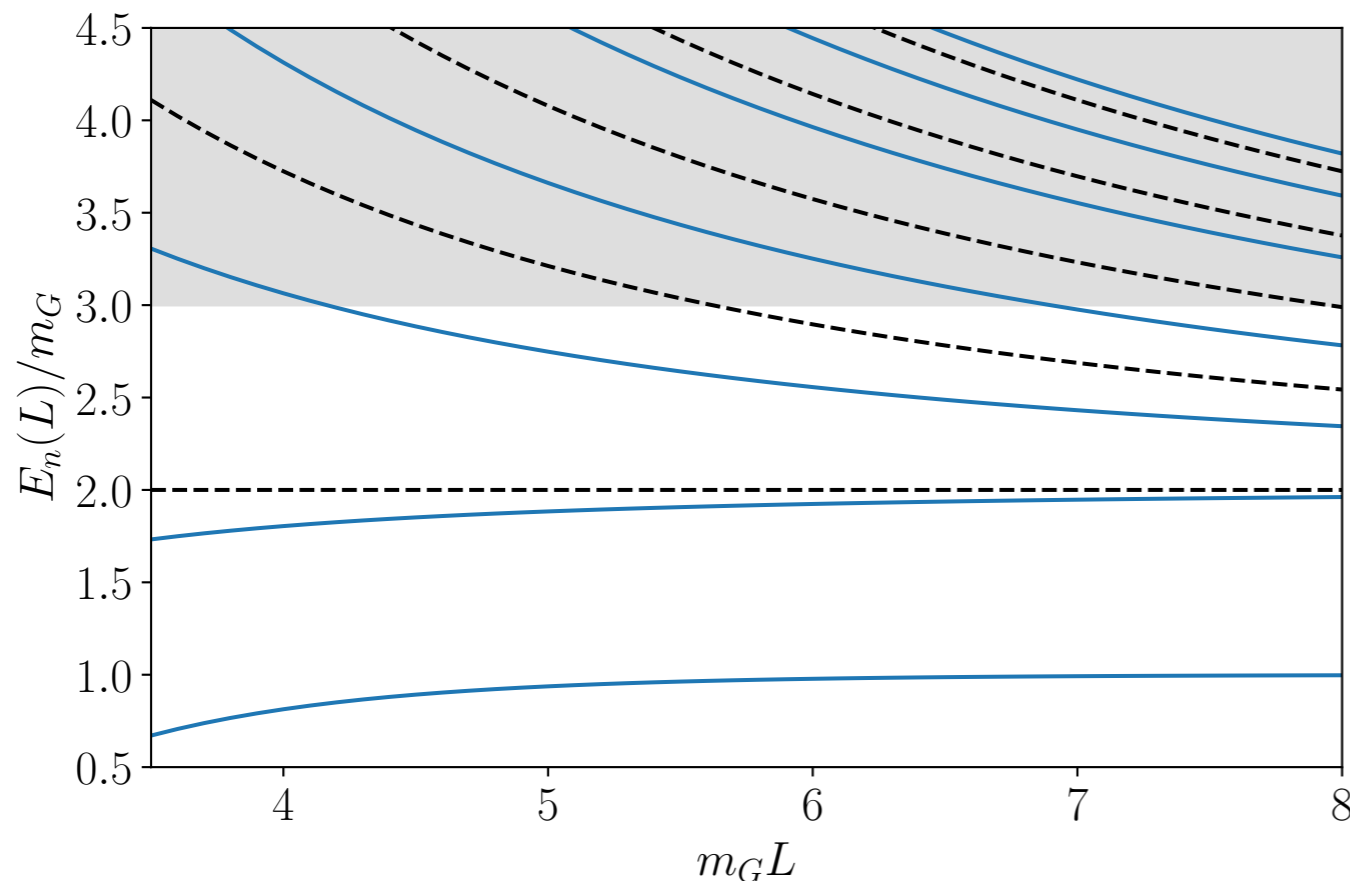
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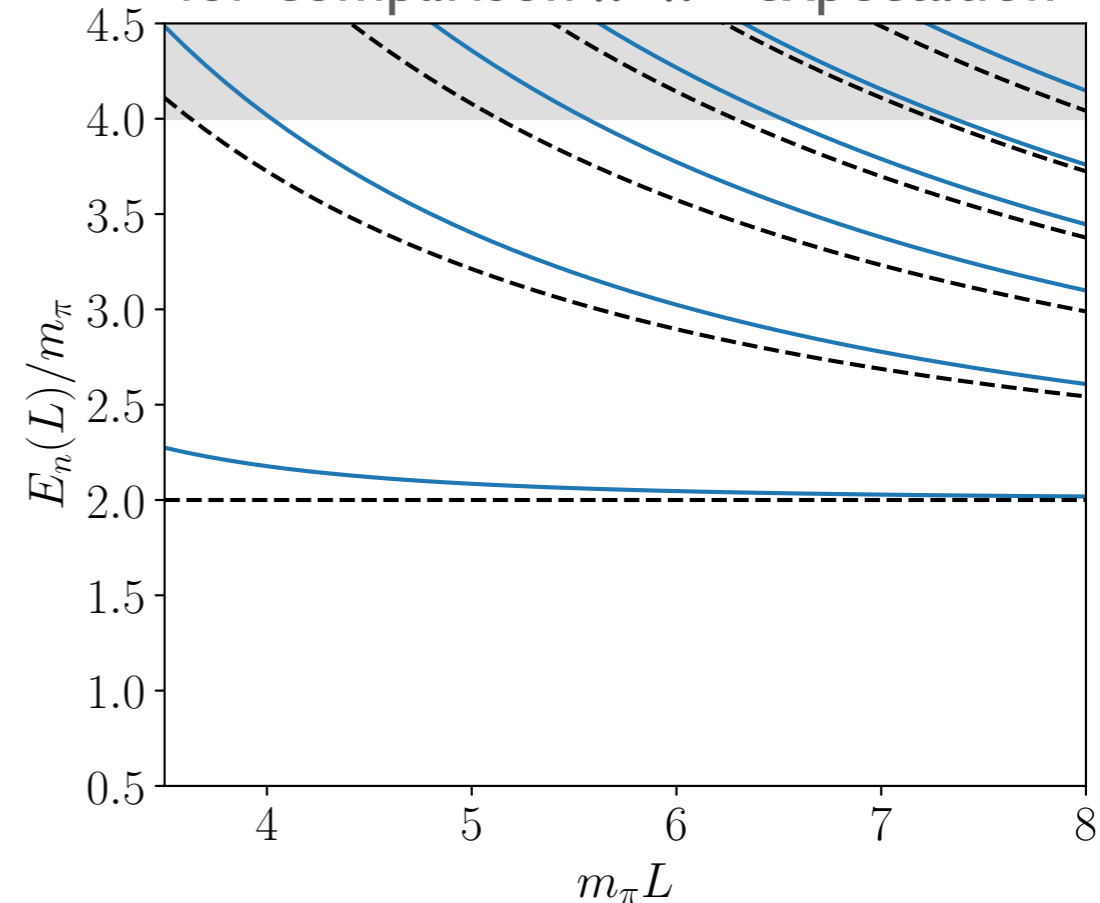
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- What do we expect for the spectrum?...

0^{++} qualitative expectation



for comparison $\pi^+\pi^+$ expectation



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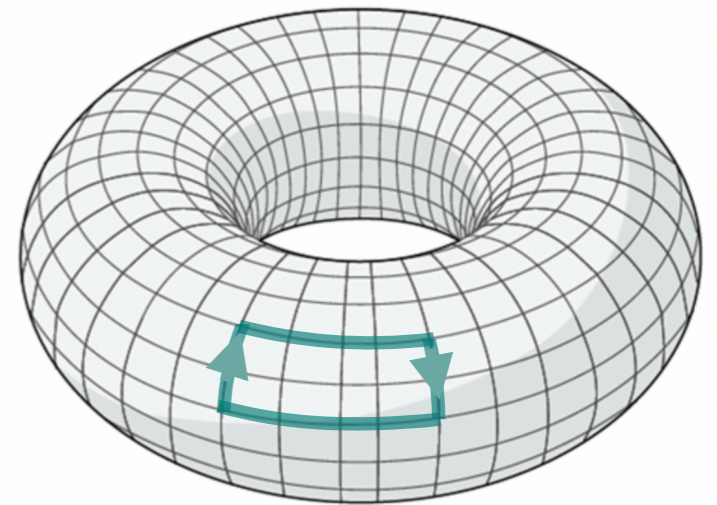
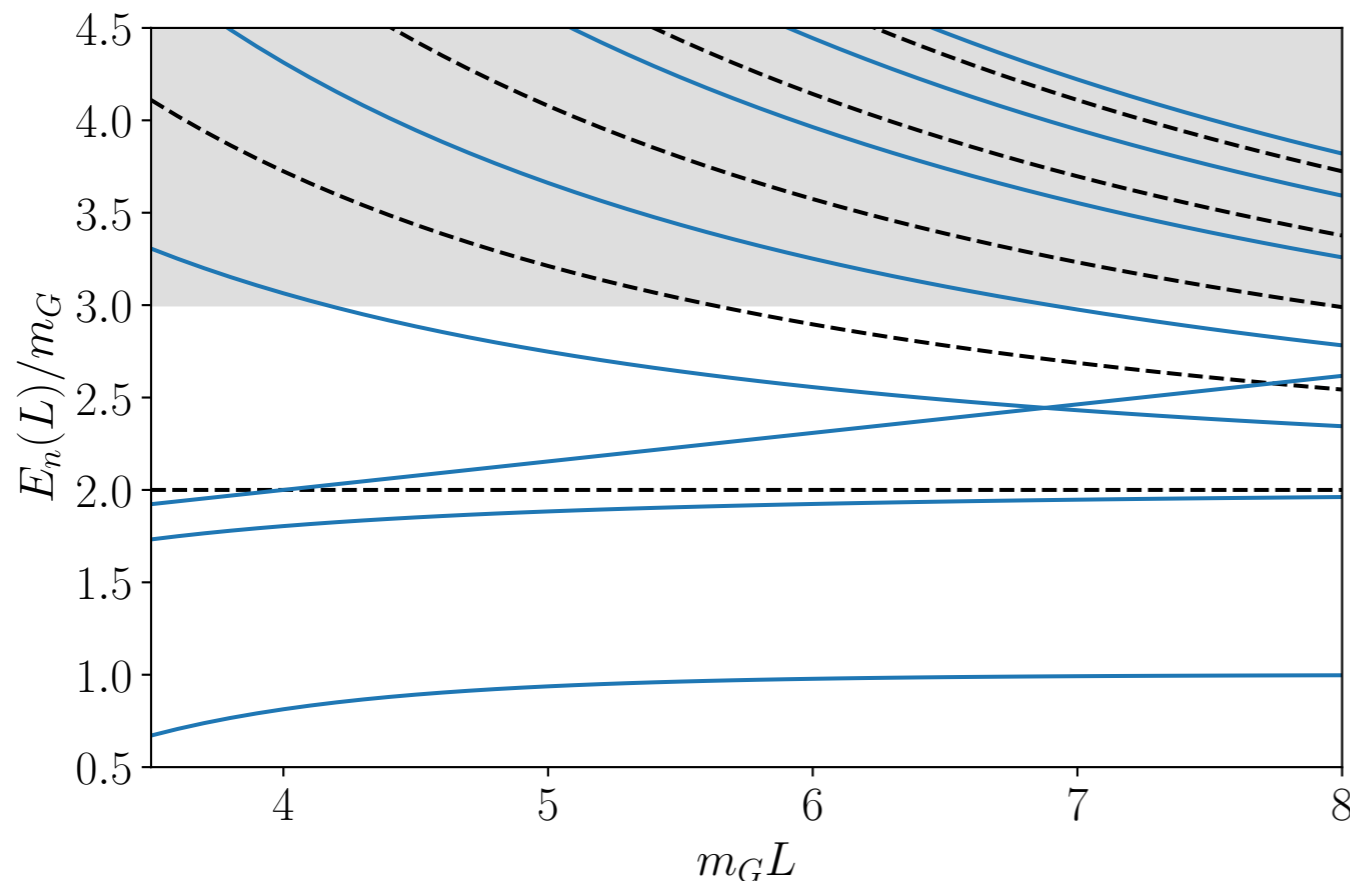
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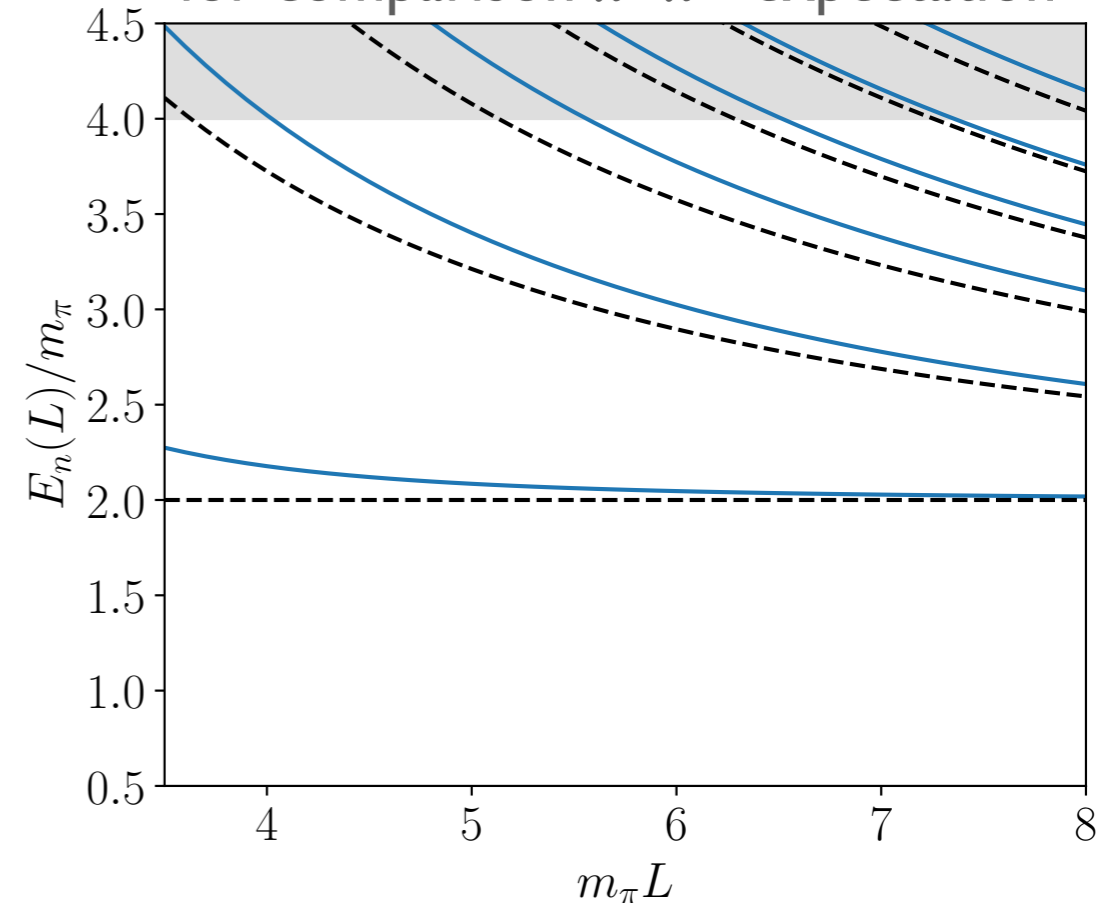
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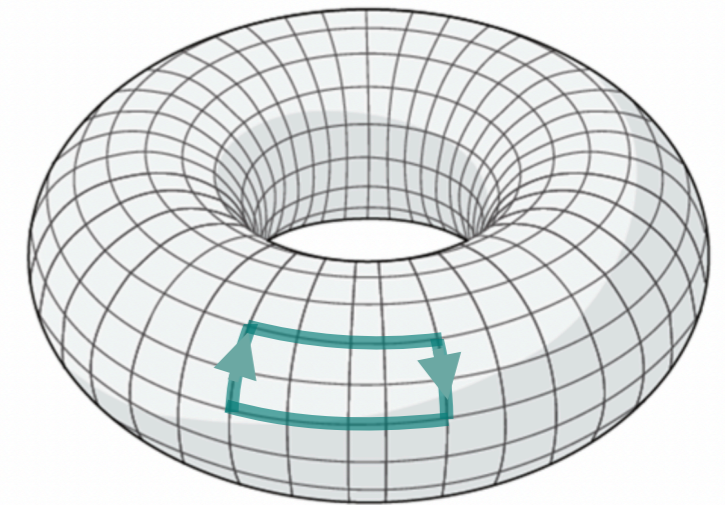
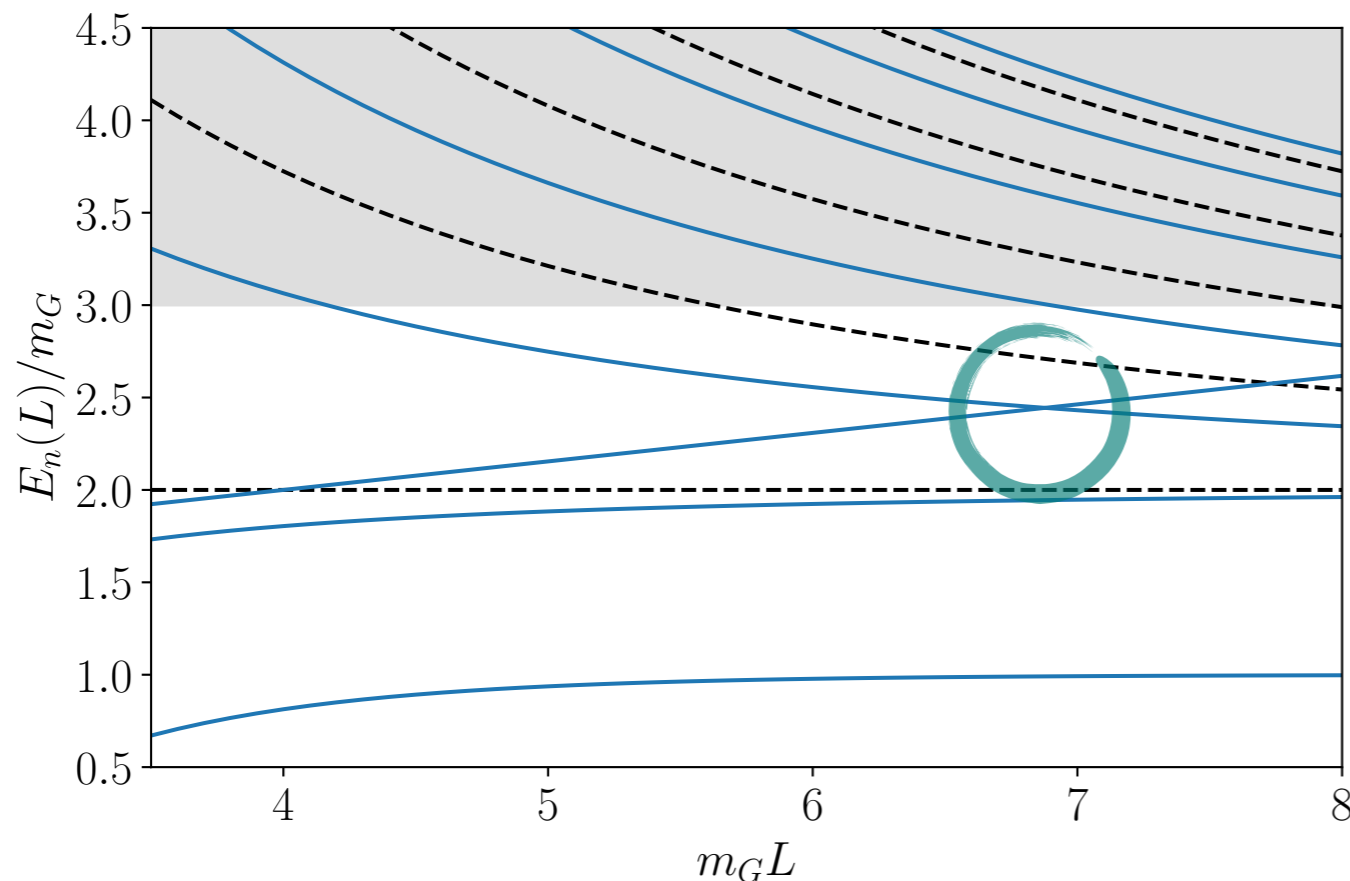
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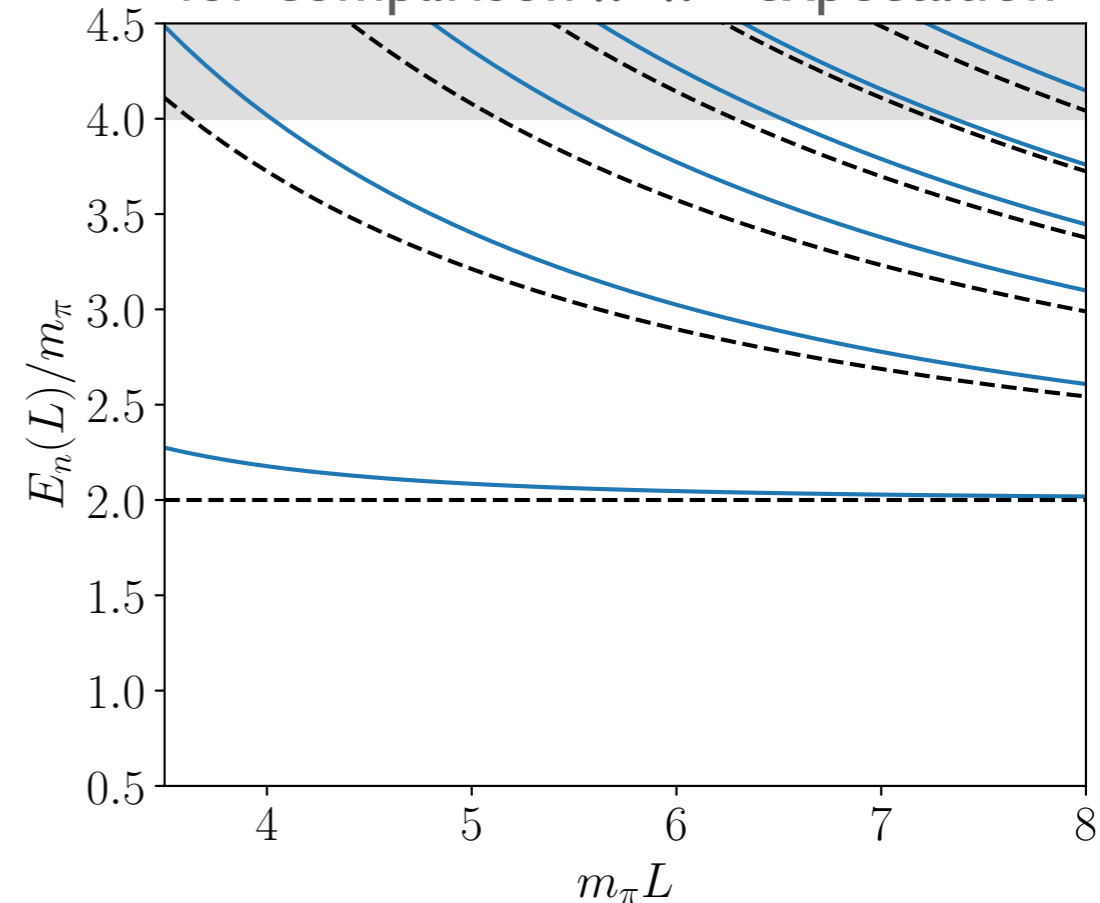
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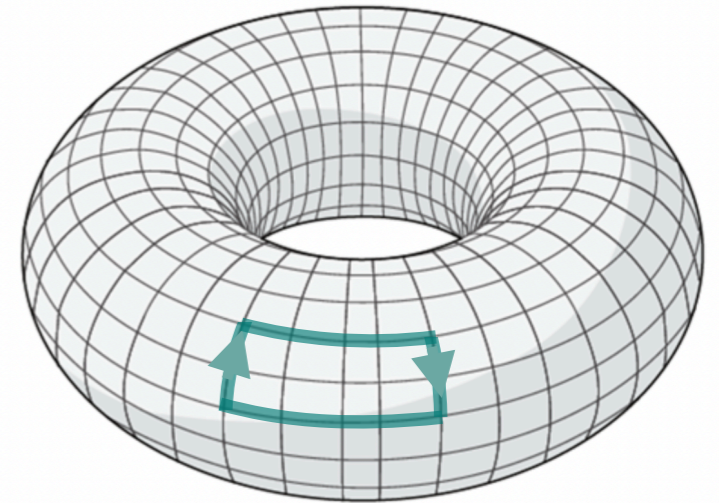
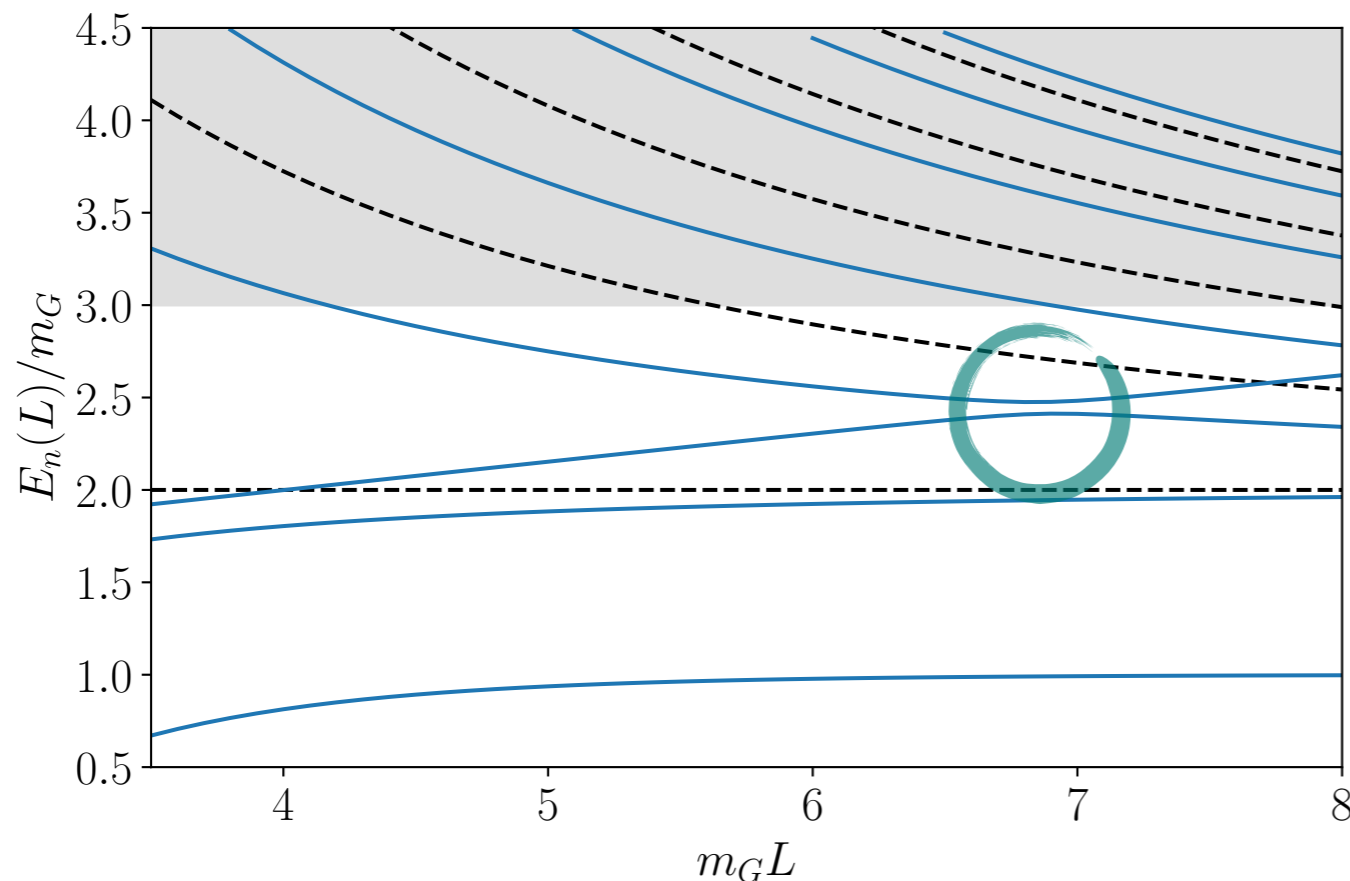
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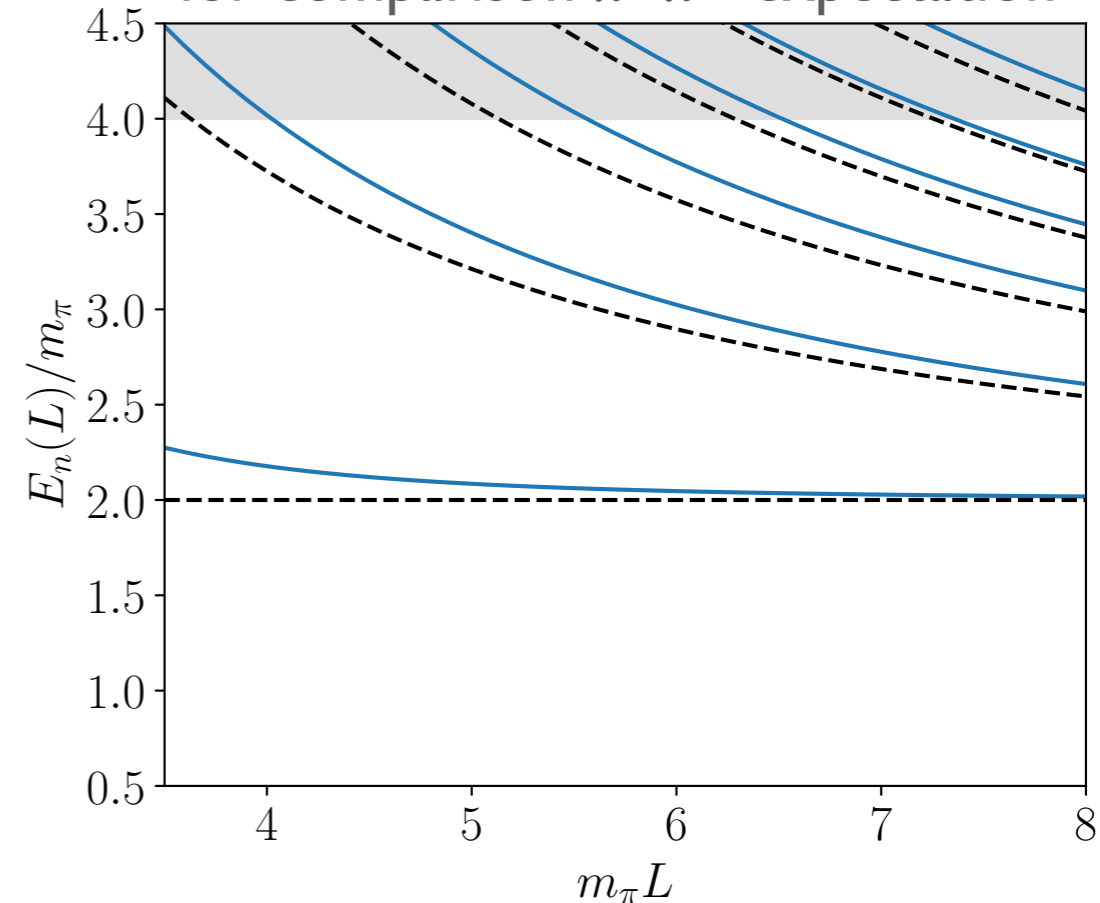
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Set-up

- Anisotropic Wilson plaquette action ($a_t \neq a_s$)
- Renormalised anisotropy tuned via *spatial-spatial and spatial-temporal energy density*

$$\xi_R(\beta, \xi) = \frac{a_s(\beta, \xi)}{a_t(\beta, \xi)} = 4 \quad \text{[BMW]}$$

- Multi-level algorithm (*described more on next slide*)

Set-up

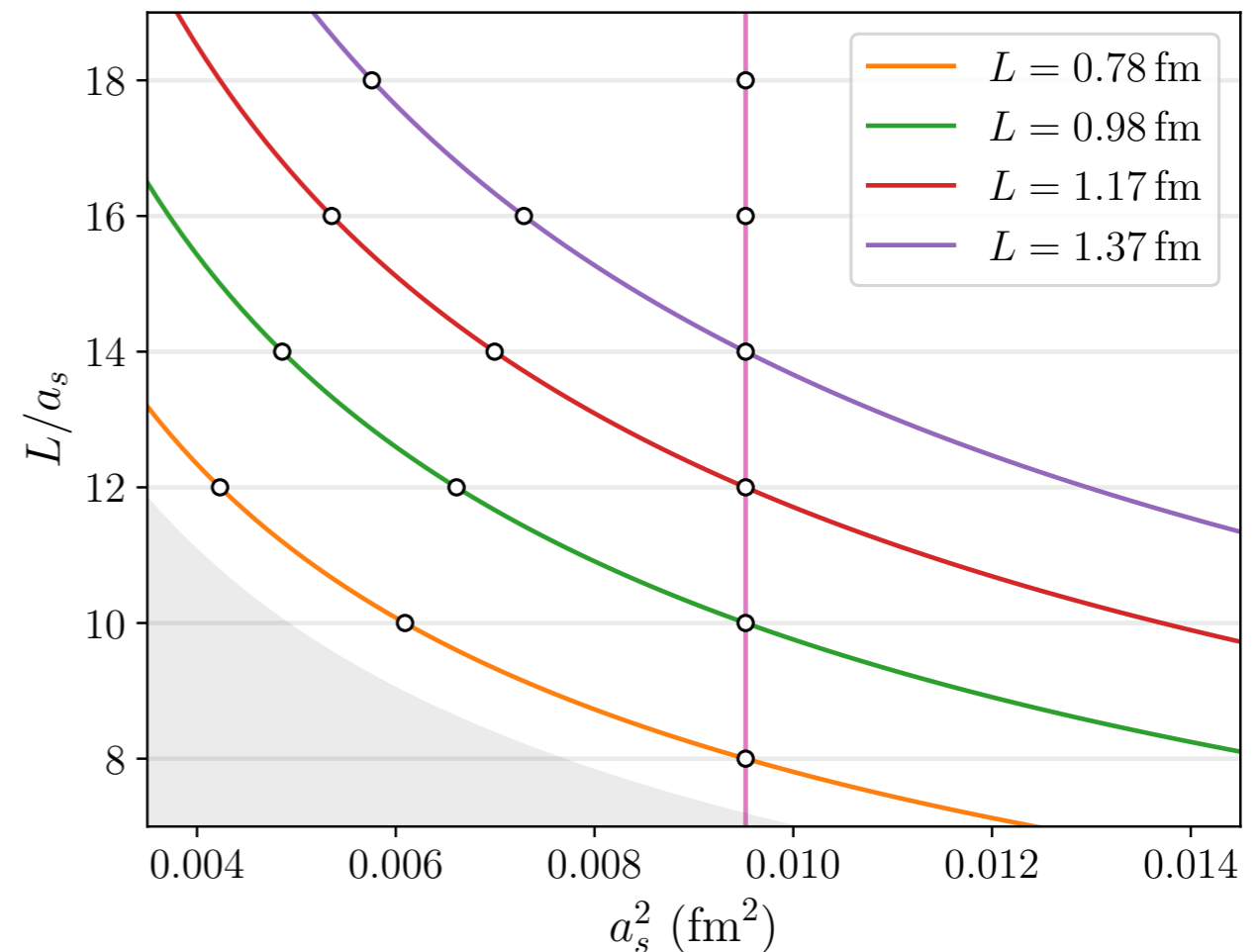
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- Multi-level algorithm (*described more on next slide*)
- All ensembles tuned... so far only **vertical pink line** has been generated

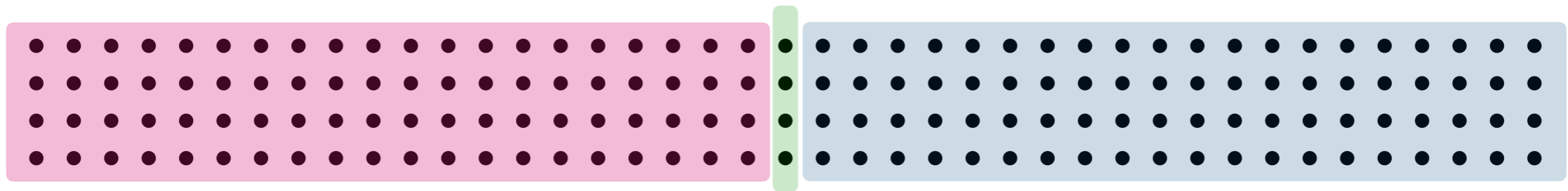
- Scale-setting via t_0

$$t^2 \langle E(t) \rangle \Big|_{t=t_0} = 0.3$$



Multi-level algorithm

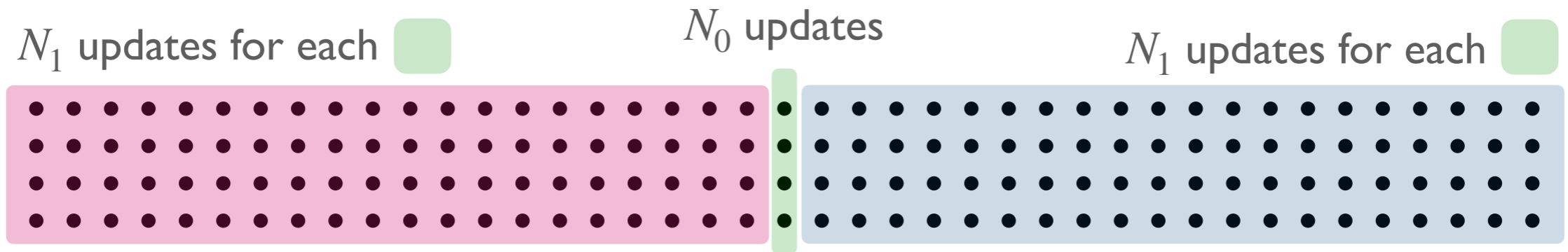
➤ Update parts of lattice to improve signal at fixed computational cost



[Lüscher and Weisz, (2001)] - [Meyer (2003,2004)] - [Ce, Giusti, Schaefer (2017)] - [Barca et al. (2024)]

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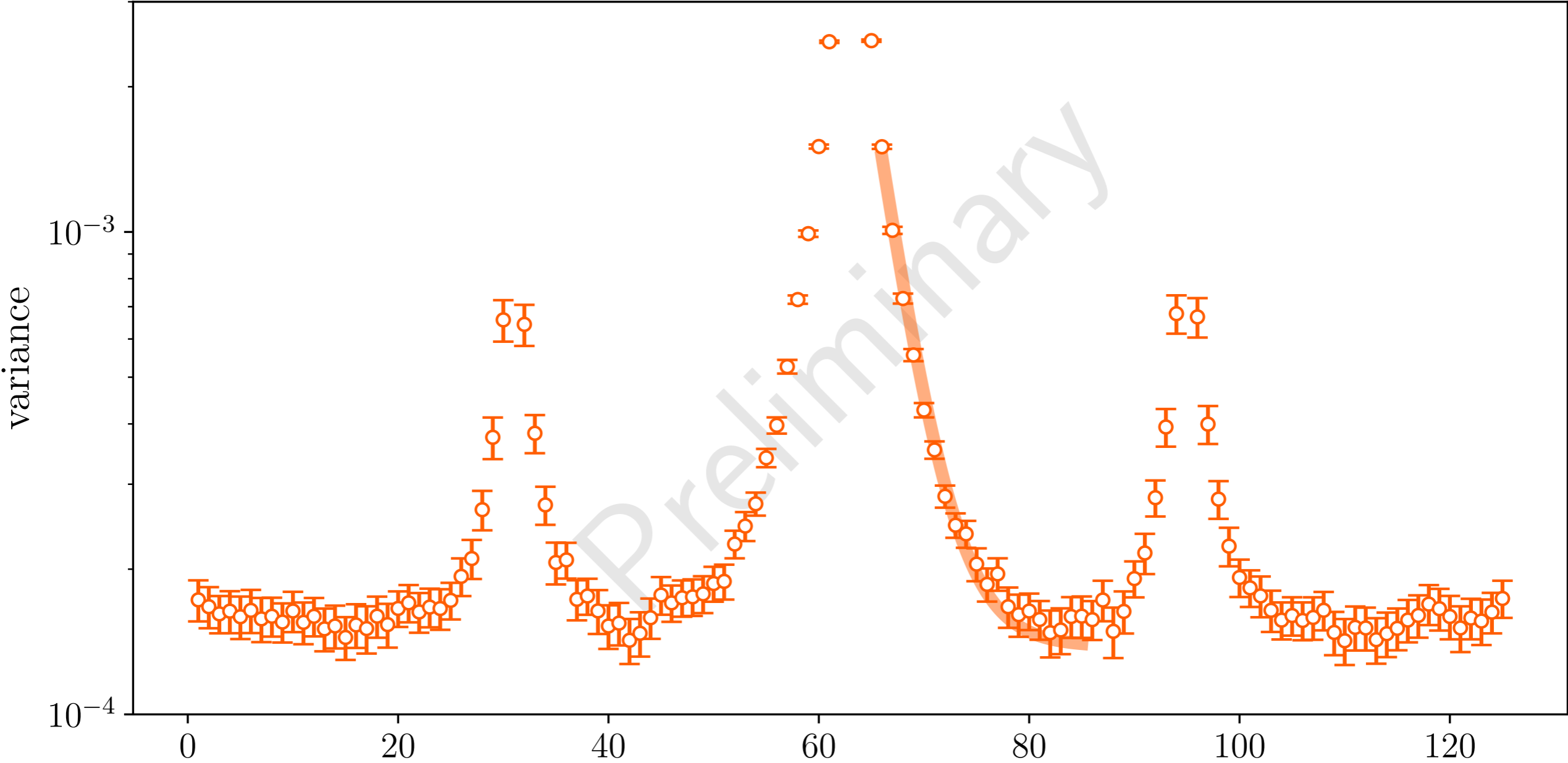
$$\begin{aligned}
 \langle \mathcal{O}_x^{(1)} \mathcal{O}_y^{(2)} \rangle &= \sum_{\mathcal{A} \in \Omega} p(\mathcal{A}) \left(\sum_{X \in \Omega_{\text{left}}} p_{\mathcal{A}}(X) \mathcal{O}_x^{(1)}[\mathcal{A}, X] \right) \left(\sum_{Y \in \Omega_{\text{right}}} p_{\mathcal{A}}(Y) \mathcal{O}_y^{(2)}[\mathcal{C}, Y] \right) \\
 &= \sum_{\mathcal{A} \in \Omega} p(\mathcal{A}) \langle \mathcal{O}_x^{(1)} \rangle_{\text{left}, \mathcal{A}} \langle \mathcal{O}_x^{(2)} \rangle_{\text{right}, \mathcal{A}}
 \end{aligned}$$

➤ Ideally variance should scale as $N_1^2 N_0$ (and cost as $N_1 N_0$)

[Lüscher and Weisz, (2001)] - [Meyer (2003,2004)] - [Ce, Giusti, Schaefer (2017)] - [Barca et al. (2024)]

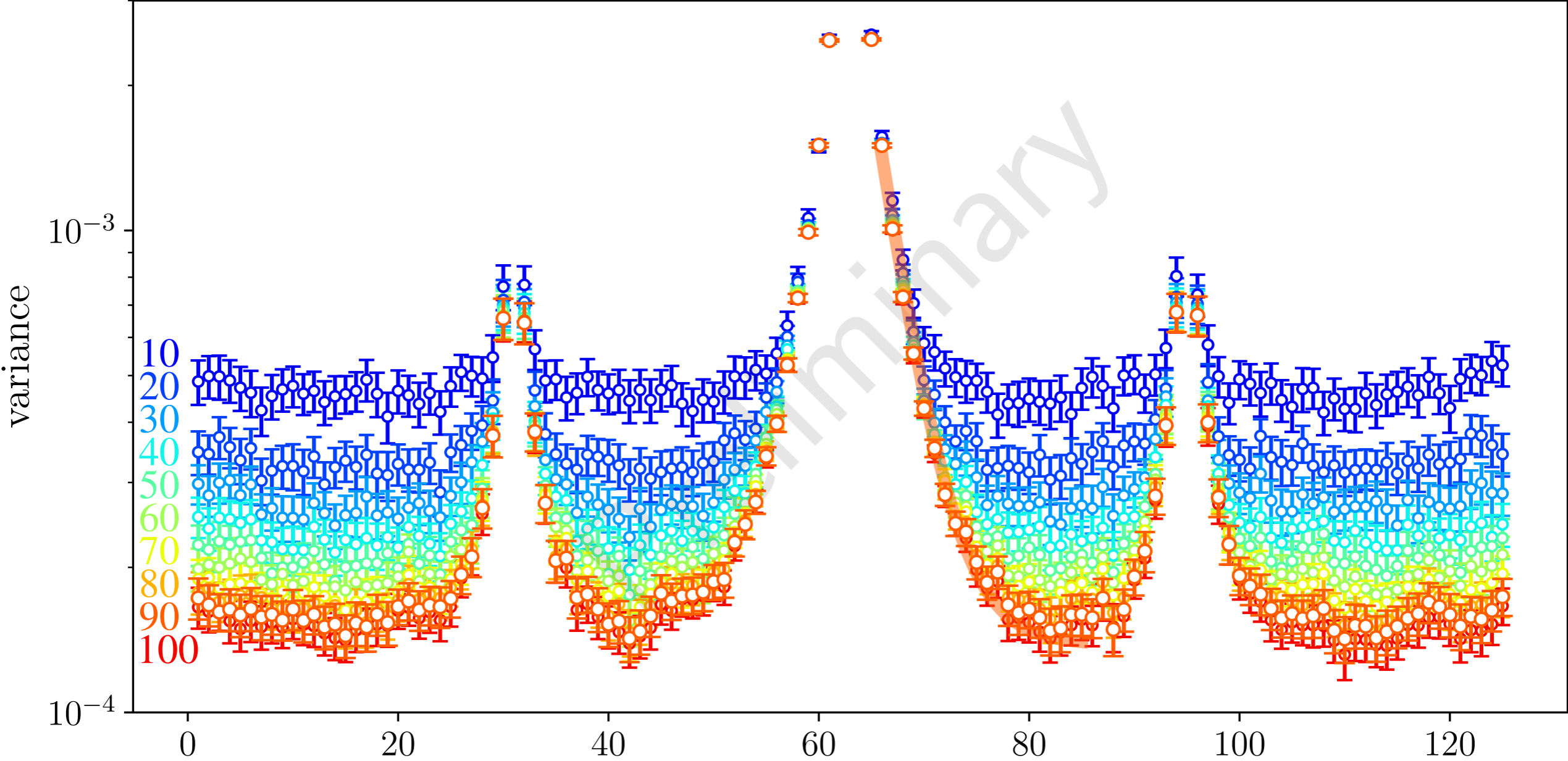
Theoretical understanding of scaling

➤ Variance (and uncertainty of variance) of the one point off the boundary



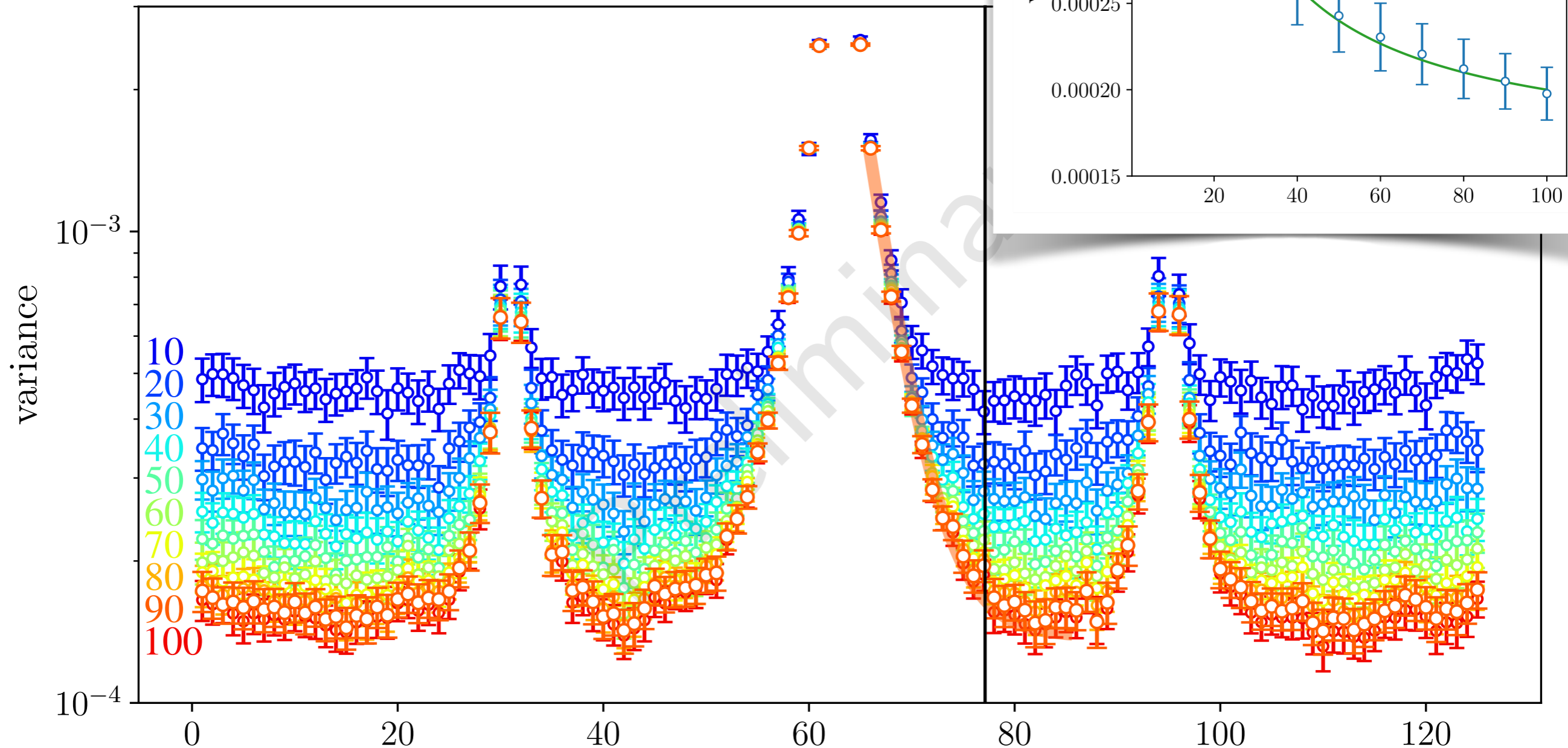
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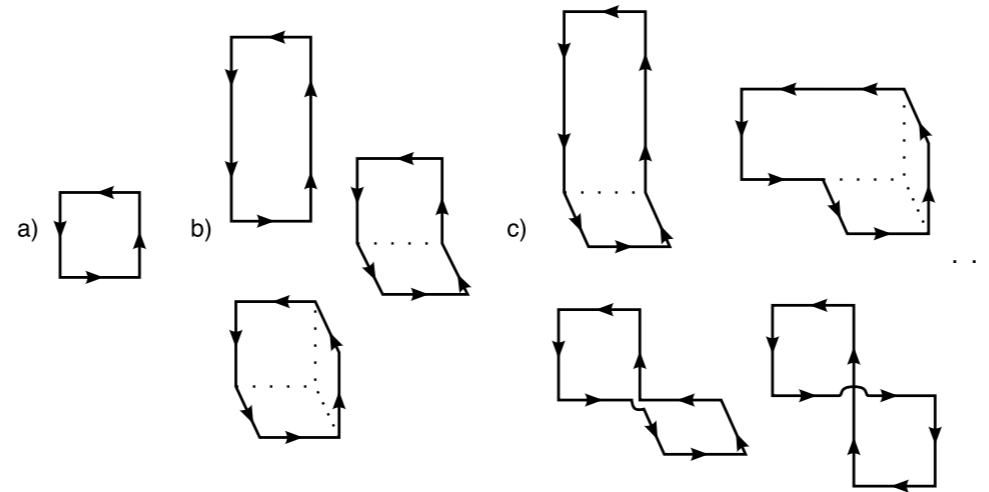
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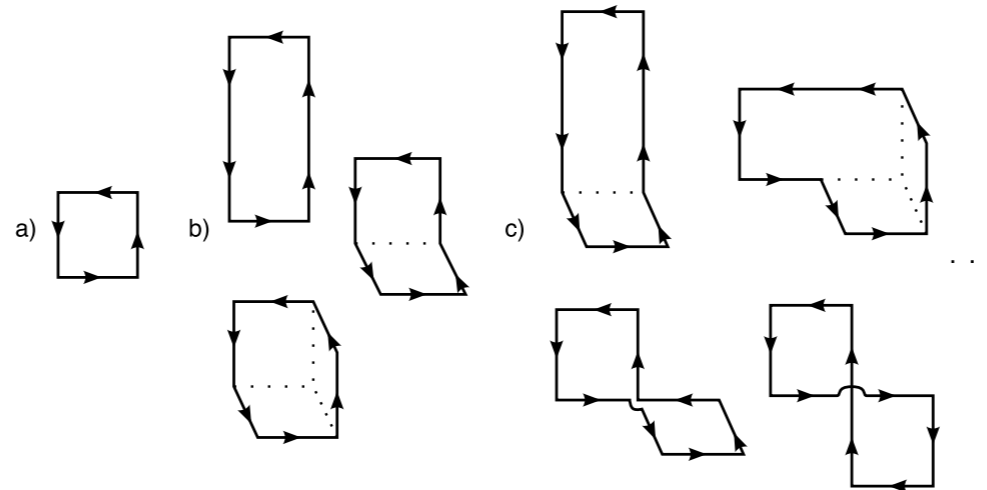
Operator construction

➤ General approach = many shapes of Wilson loops + smearing + blocking



Operator construction

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➤ Project to all good finite-volume quantum numbers as usual

$$\left(\text{Image of a blue and white airplane} \right)_{A_1^+} = 1 \times \left(\text{Image of a white airplane} \right) + 1 \times \left(\text{Image of a white airplane from a different angle} \right) + 1 \times \left(\text{Image of a blue and white airplane from a different angle} \right) + 1 \times \left(\text{Image of a dark airplane} \right) + \dots$$

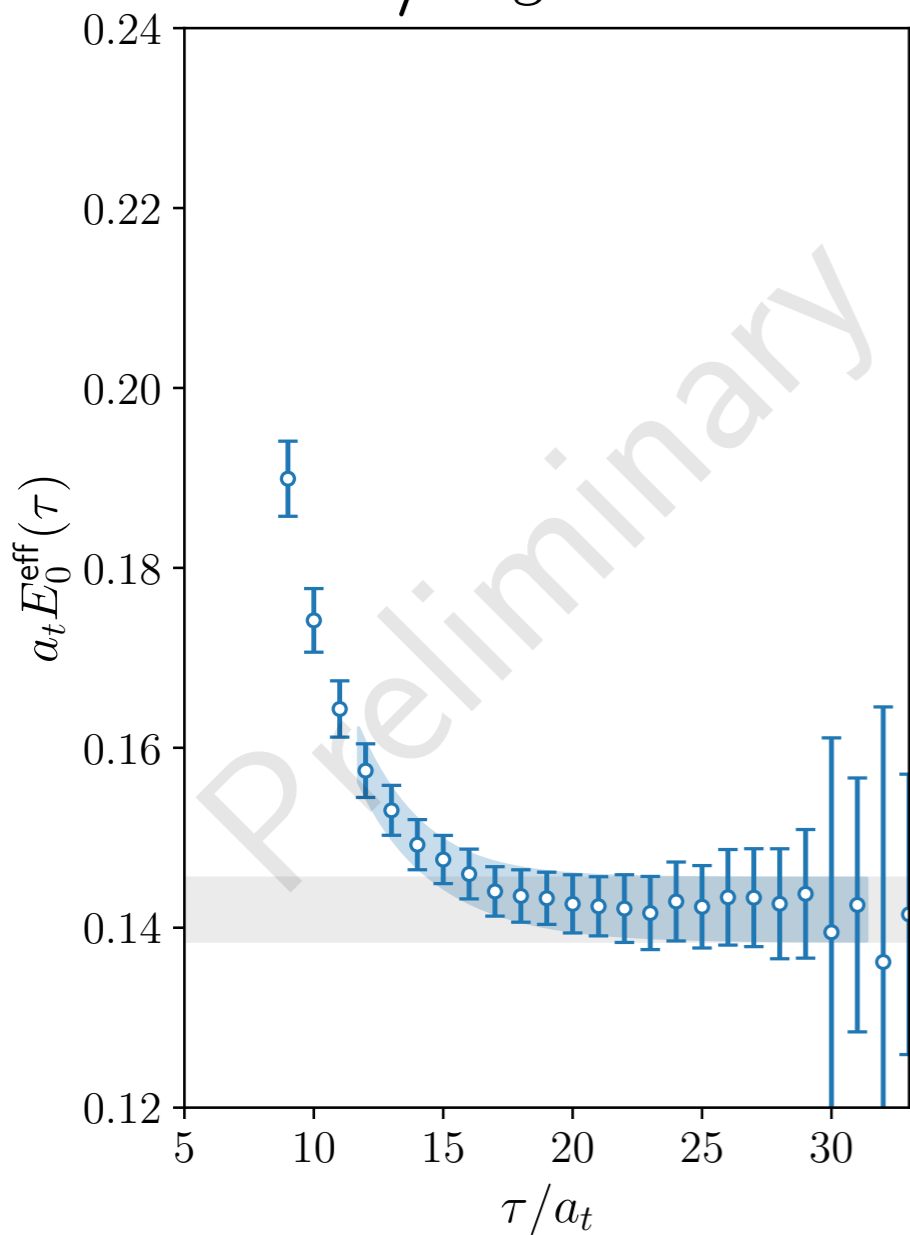
➤ Also have two-trace operators to extract “two-gluon-ball-like” excited states

$$\sum_{\mathbf{x}, \mathbf{y}} e^{-i\mathbf{p} \cdot \mathbf{x} - i\mathbf{k} \cdot \mathbf{y}} \text{tr}[W^{(1)}(\tau, \mathbf{x})] \text{tr}[W^{(2)}(\tau, \mathbf{y})]$$

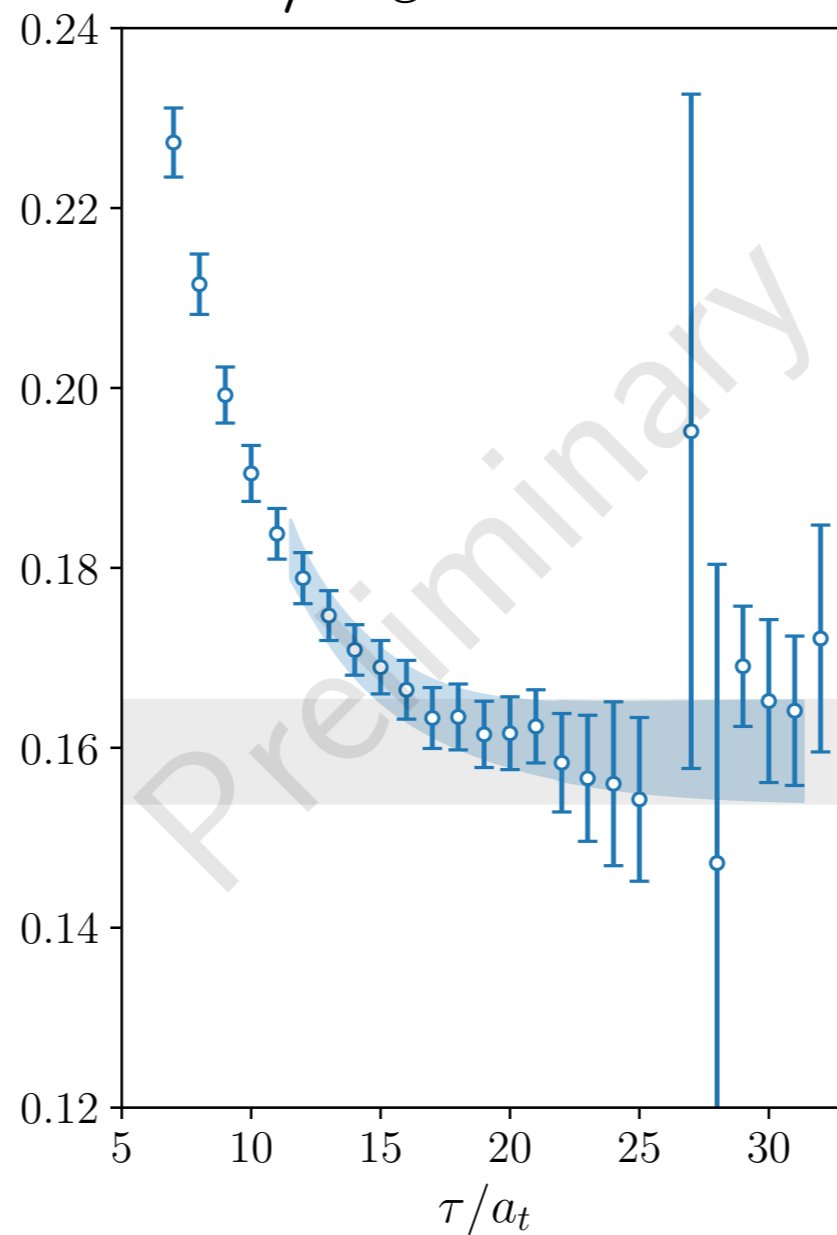
Extracting energies

➤ Use variational method to extract the ground state energy

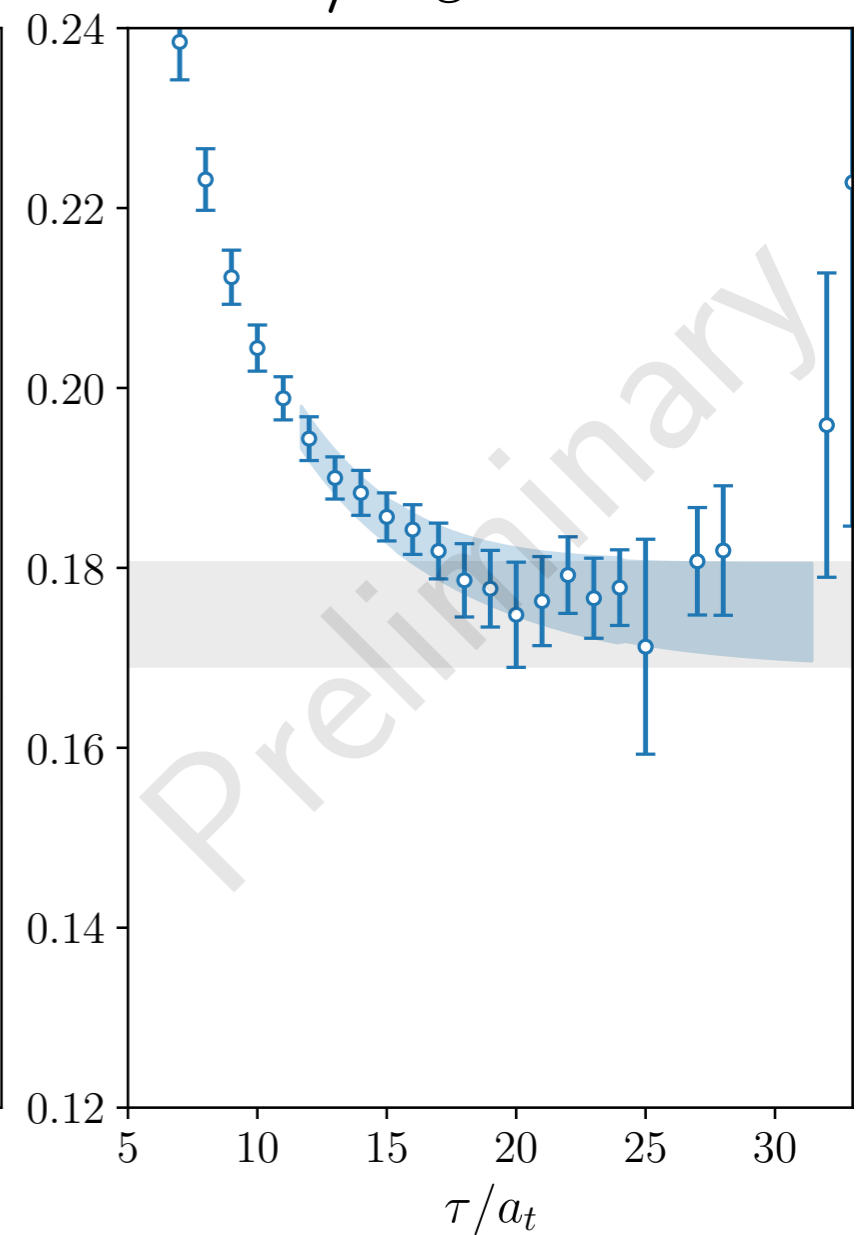
$L/a_s = 8$



$L/a_s = 10$

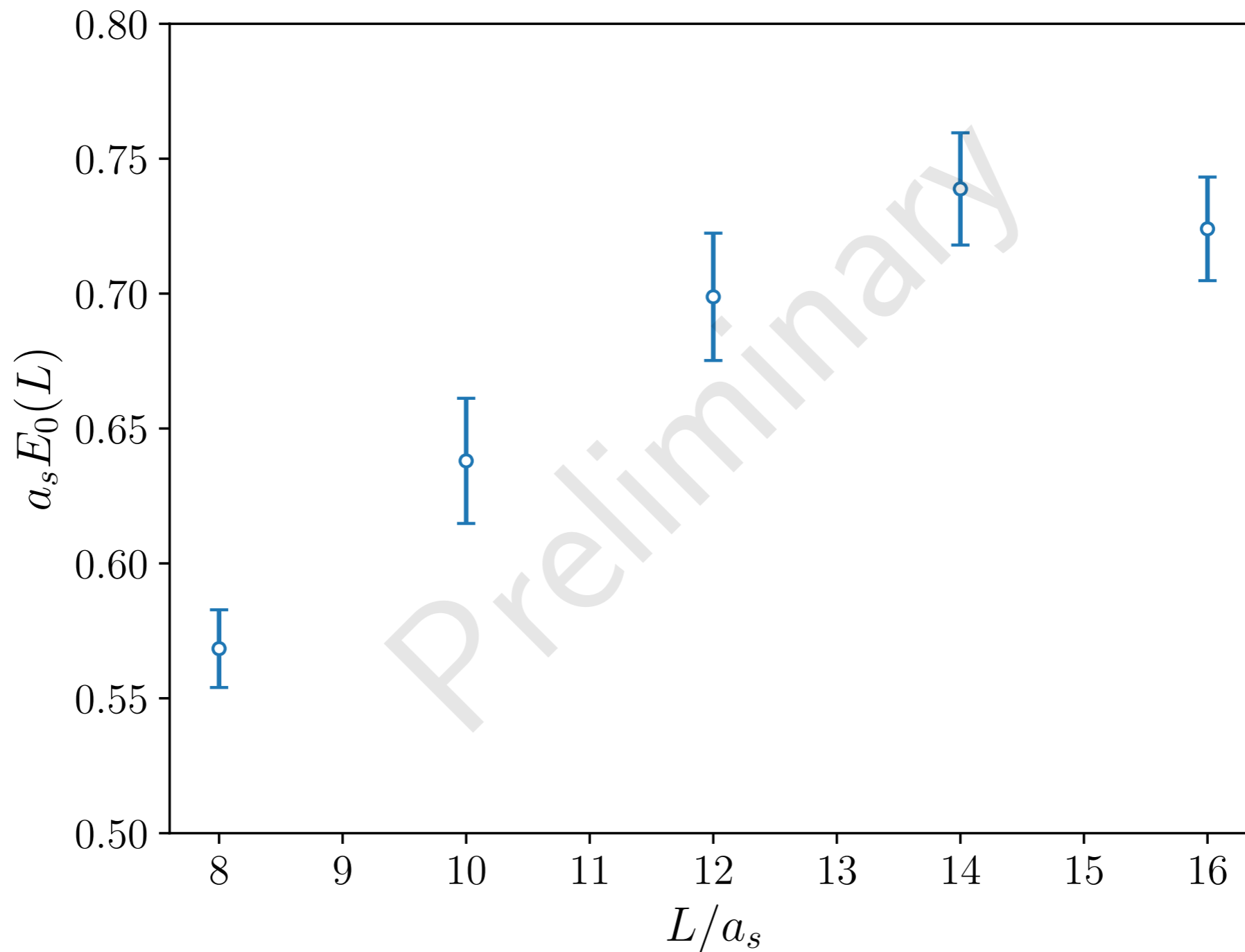


$L/a_s = 12$



$$E_0^{[A_1^{++}]}(L) \text{ vs } L$$

➤ Volume dependence of the ground state



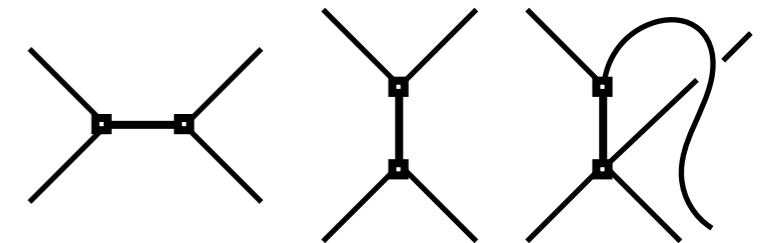
Interpreting the ground state L dependence

➤ Lüscher stable-states paper relates L dependence to trilinear coupling λ

$$\Delta m = -\frac{3}{16\pi m^2 L} \left\{ \lambda^2 e^{-\frac{\sqrt{3}}{2}mL} + \frac{m}{\pi} \int_{-\infty}^{\infty} dy e^{-\sqrt{m^2+y^2}L} F(iy) + O(e^{-\bar{m}L}) \right\}$$

[Lüscher (1986)]

➤ **WARNING:** If λ is too large, subleading exponentials could be important (also F contains terms proportional to λ^2)



Interpreting the ground state L dependence

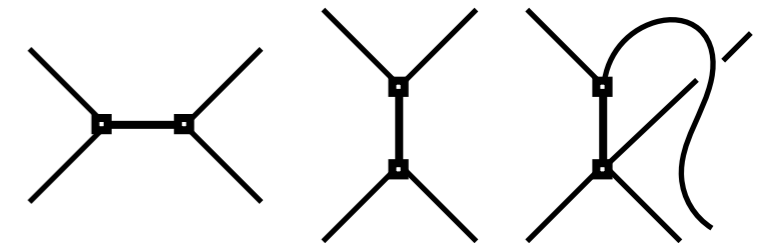
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➤ Method was already used to estimate the coupling in 1984



$$\frac{\lambda}{m} = 50 \pm 8$$

[de Forcrand et al., (1984)]

Interpreting the ground state L dependence

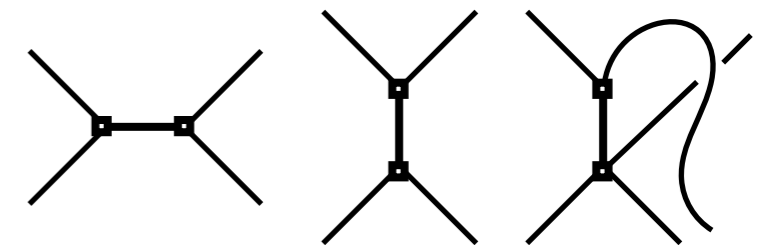
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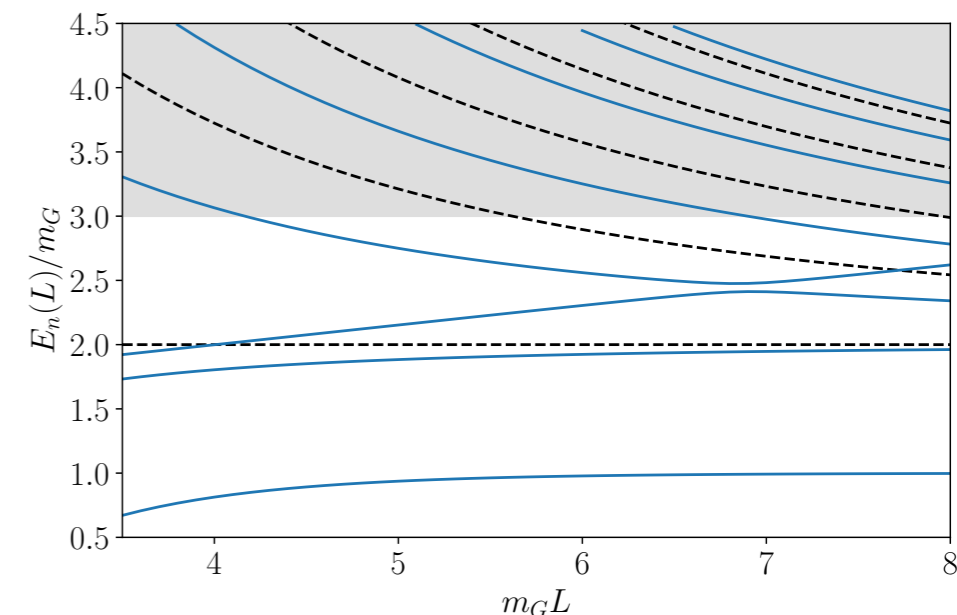
[de Forcrand et al., (1984)]

- Can also sum enhanced exponentials by continuing scattering-states formula below threshold

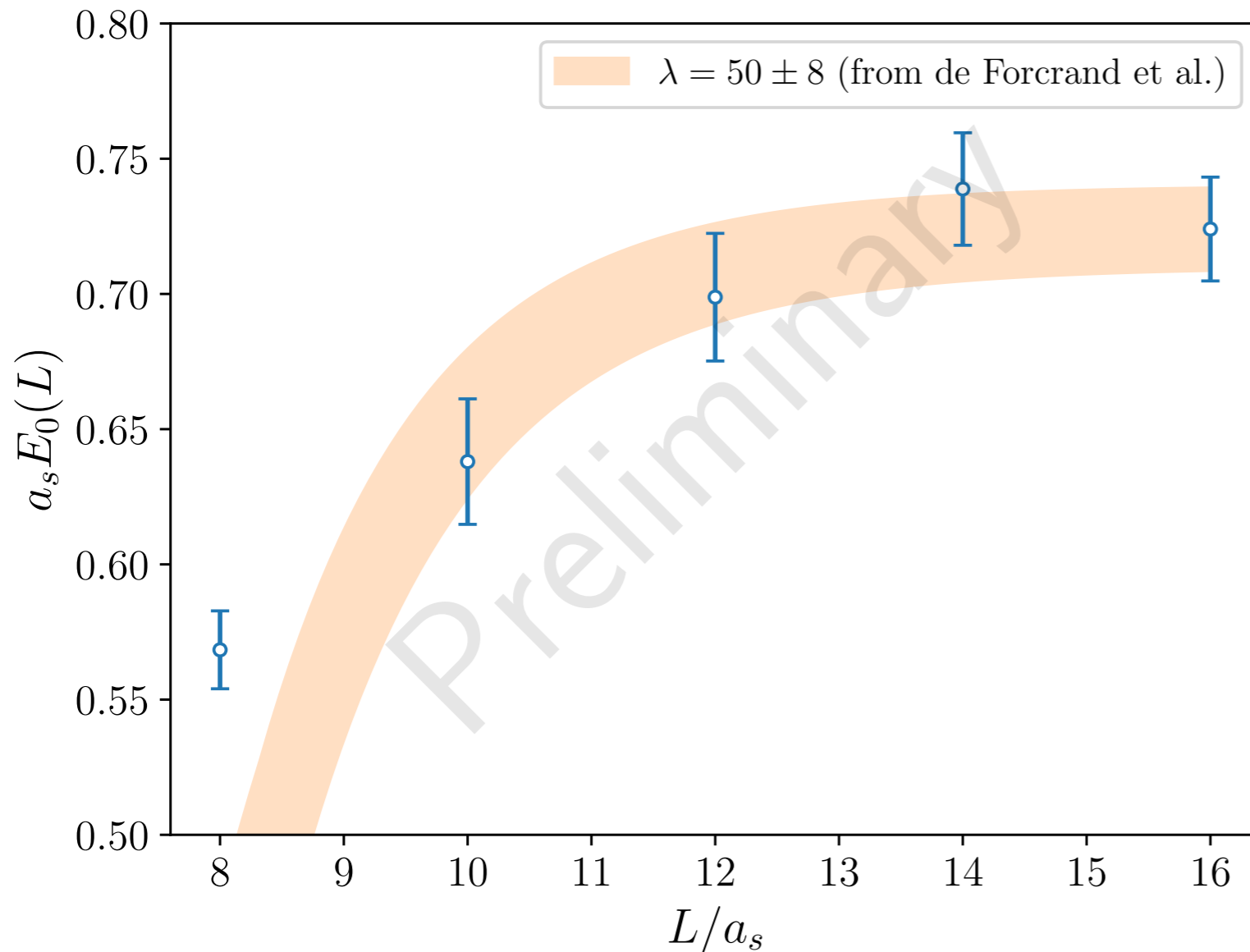
[Lüscher, (1986)]

- Requires left-hand cut formalism

[MTH and Raposo, (2024)]



$E_0^{[A_1^{++}]}(L)$ vs L



$$\Delta m = -\frac{3\lambda^2}{16\pi m^2 L} e^{-\sqrt{3}mL/2}$$

- Our interpretation = matching the *model and coupling* to de Forcrand et al. gives a good description
- For these parameters *sub-leading exponentials* could significantly change the story

Conclusions and next steps

➤ Presented results for 0^{++} ground state vs L :

$$E_0^{[A_1^{++}]}(L)$$

➤ L -dependence gives a first probe of scattering: *tri-linear coupling*

➤ S-matrix bootstrap also constrains $\lambda/m < 200$

[Guerrieride Forcrand et al., (2023)]

➤ Next steps (most ambitious scenario)

extract excited-state energies... in moving frames... for various quantum numbers

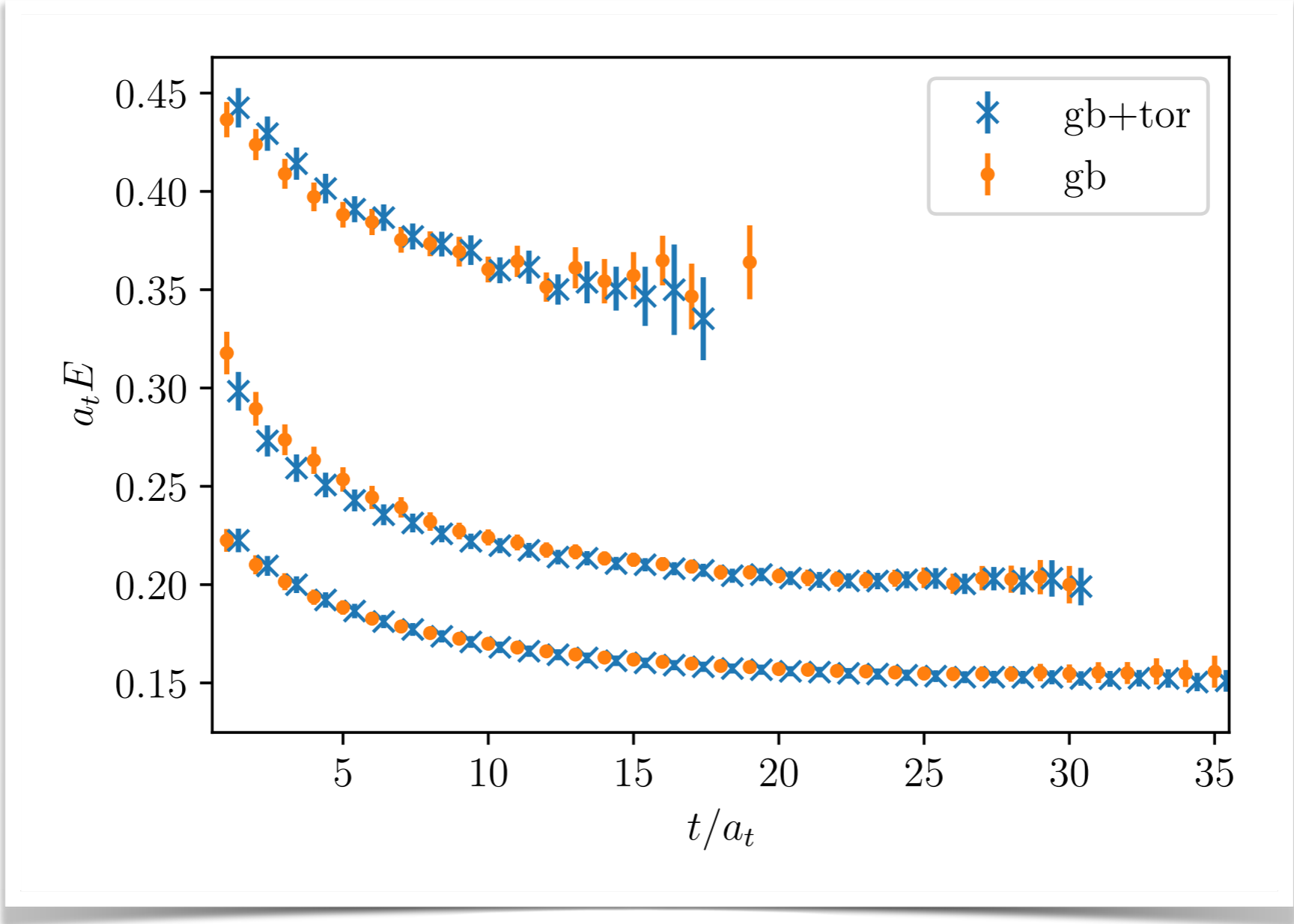
determine and publish fully controlled, continuum-limit, finite-volume energies

analyse two- (and three-particle energies) and predict scattering amplitudes

Backup slides

Di-torelon stability

➤ Stability with respect to inclusion/exclusion of di-torelon operator



GEVP improvement

