

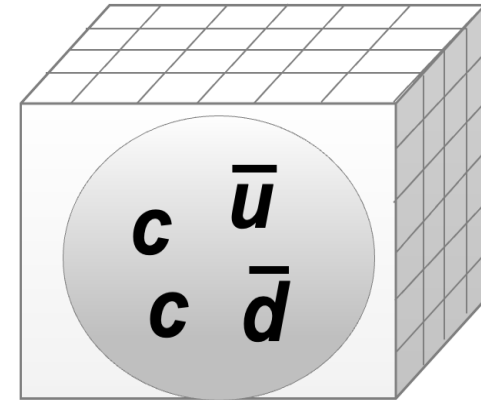
# Towards quark mass dependence of $T_{cc}$

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Lattice 2024, July 2024, Liverpool



Sara Collins, Alexey Nefediev, M. Padmanath, SP [2402.14715, PRD]



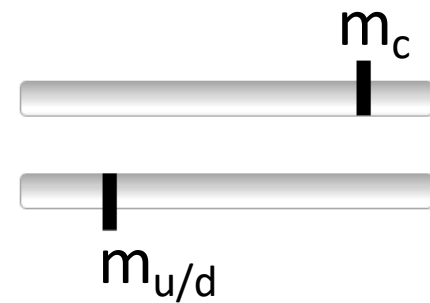
Regensburg



Ljubljana



Chennai



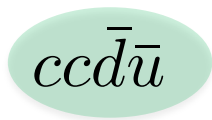
follow-up of M. Padmanath, SP [2202.10110, PRL]

# $T_{cc}$ from LHCb experiment

$$D^* \rightarrow D\pi$$

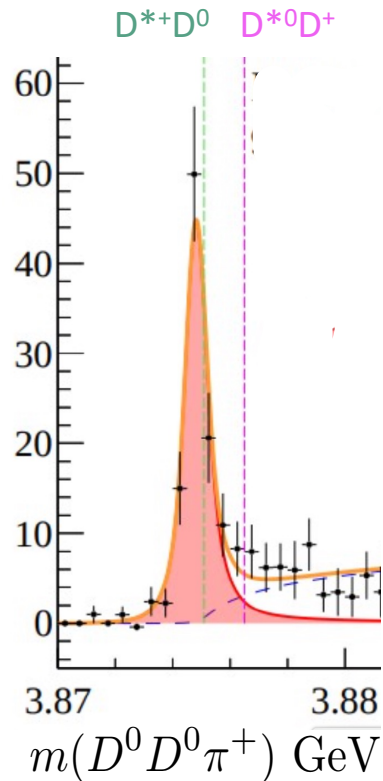
$$m_{\pi^0} \simeq 135 \text{ MeV}$$

$$m_{D^{*+}} - m_{D^+} \simeq 140 \text{ MeV}$$



$I=0, J^P=1^+$  (most likely)

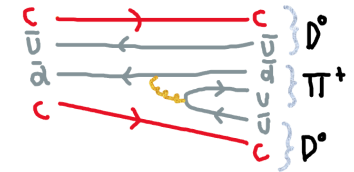
The longest lived exotic hadron ever discovered



$$\delta m = m - (m_{D^{*+}} + m_{D^0})$$

$$\delta m_{pole} = -0.36 \pm 0.04 \text{ MeV}$$

LHCb 2109.01038, 2109.01056, Nature Physics



Omitting  $D^* \rightarrow D\pi$ ,  $T_{cc} \rightarrow DD\pi$   
 $T_{cc}$  would be a bound state

# T<sub>cc</sub> from lattice

all analyzed in 2402.14715, PRD

Collins, Nefediev, Padmanath , SP

all simulations :

$$m_u = m_d > m_{u,d}^{ph}$$

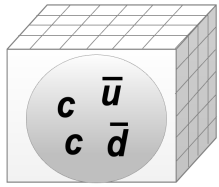
$$m_\pi > m_{D^*} - m_D$$

$$D^* \not\rightarrow D\pi$$

single lattice spacing

(J. Green et al are exploring

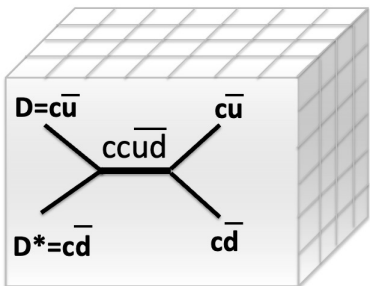
several lattice spacings, lat 2023, unpublished)



mc	mpi	L	ensembles	ref.
five values m <sub>D</sub> =1.7–2.4 GeV	280 MeV	~ 2.1, 2.8 fm	CLS Nf=2+1	our, 2402.14715, PRD eigenenergies



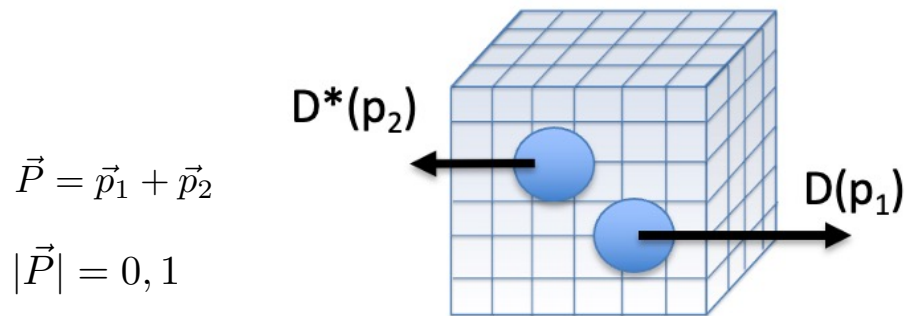
mc	mpi	L	ensembles	ref.
~ physical	146 MeV	~ 8 fm	Nf=2+1	HALQCD, 2302.04505, PRL HALQCD potentials
~ physical	280 MeV	~ 2.1, 2.8 fm	Nf=2+1, CLS	our, 2402.14715, PRD eigenenergies
~ physical	348 MeV	~ 2.4 fm	Nf=2	CLQCD, 2206.06186, PLB eigenenergies



# Interpolators and $E_n$ [our simulation, CLQCD]

$$I=0, J^P=1^+$$

$cc\bar{u}\bar{d}$



$$\mathcal{O} = \begin{matrix} D(p_1) & D^*(p_2) \\ (\bar{u}\gamma_5 c)_{\vec{p}_1} & (\bar{d}\gamma_i c)_{\vec{p}_2} - (\vec{p}_1 \leftrightarrow \vec{p}_2) \\ (\bar{u}\gamma_5 \gamma_t c)_{\vec{p}_1} & (\bar{d}\gamma_i \gamma_t c)_{\vec{p}_2} \end{matrix} \quad \vec{p}_{1,2} = \vec{n}_{1,2} \frac{2\pi}{L}$$

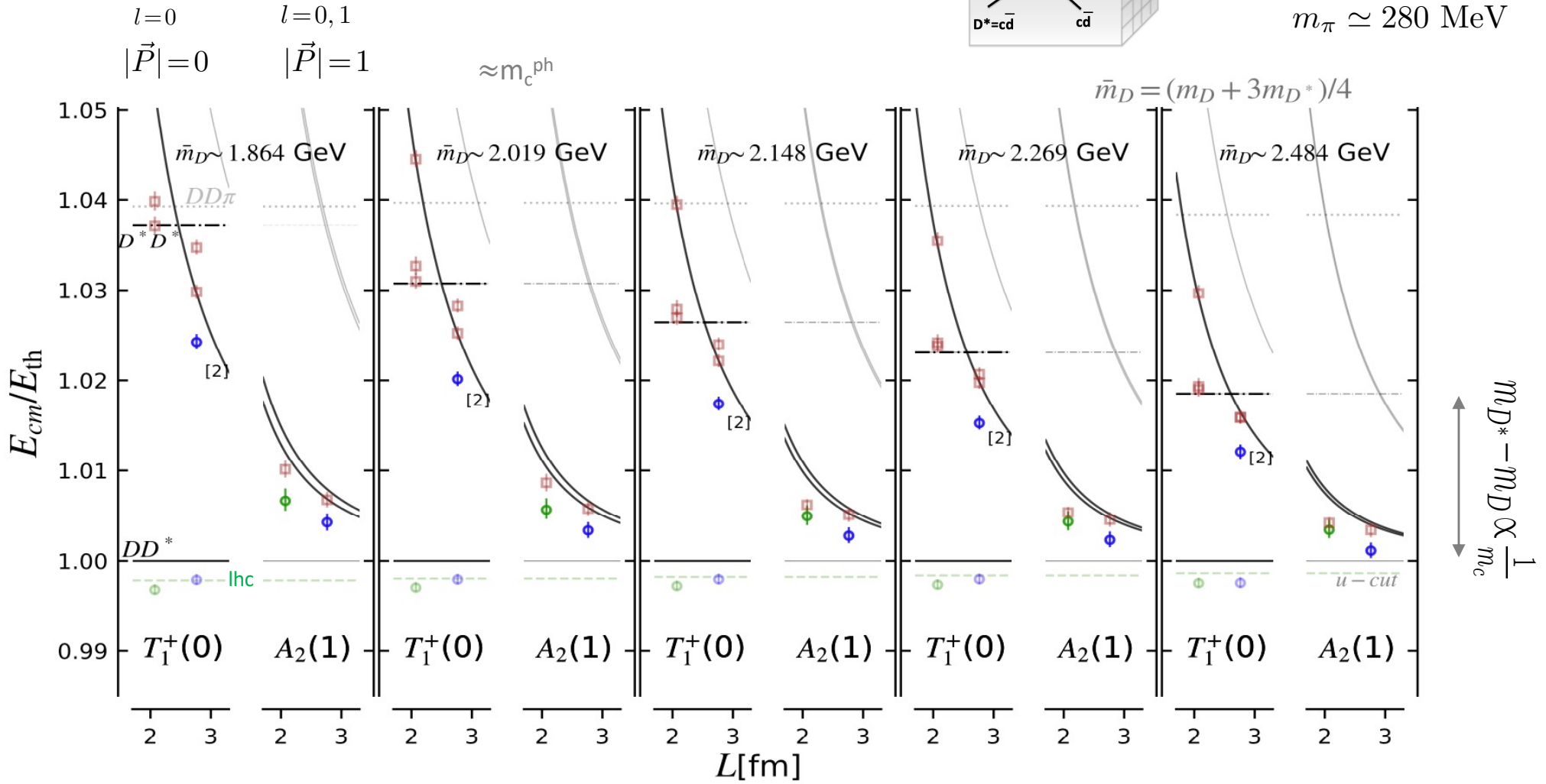
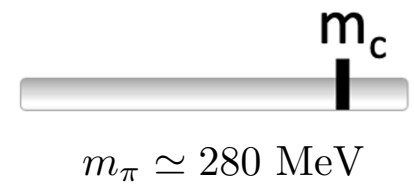
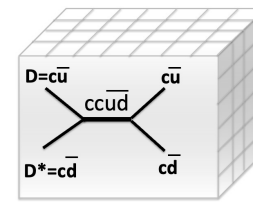
[cc][ud] interpolators not employed

[their effect discussed in the next talk by Ivan Vujmilovic]

$$C_{ij}^{2pt}(t) = \langle 0 | \mathcal{Q}_i(t) \mathcal{Q}_j^+(0) | 0 \rangle = \sum_n \langle 0 | \mathcal{Q}_i | n \rangle e^{-E_n t} \langle n | \mathcal{Q}_j^+ | 0 \rangle$$

- our simulation: distillation
- GeVP

# $T_{cc}$ : finite-volume eigen-energies



lines

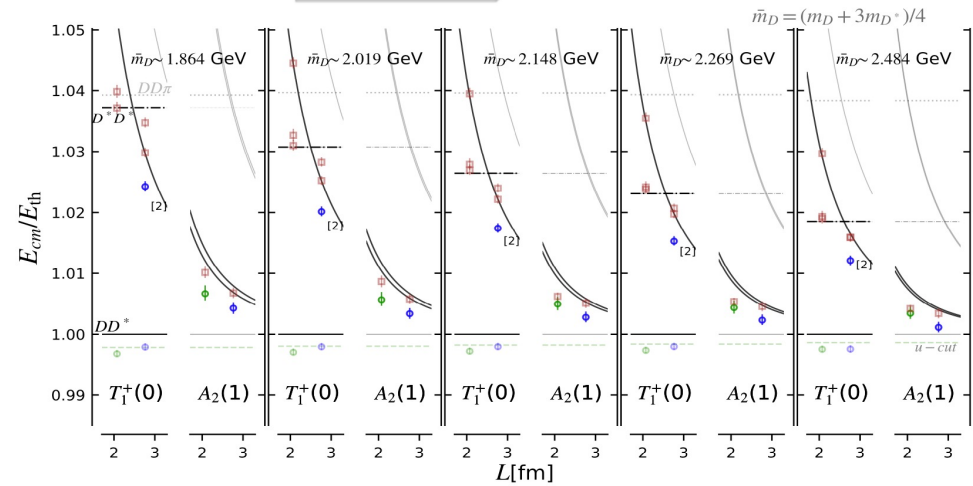
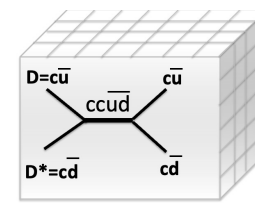
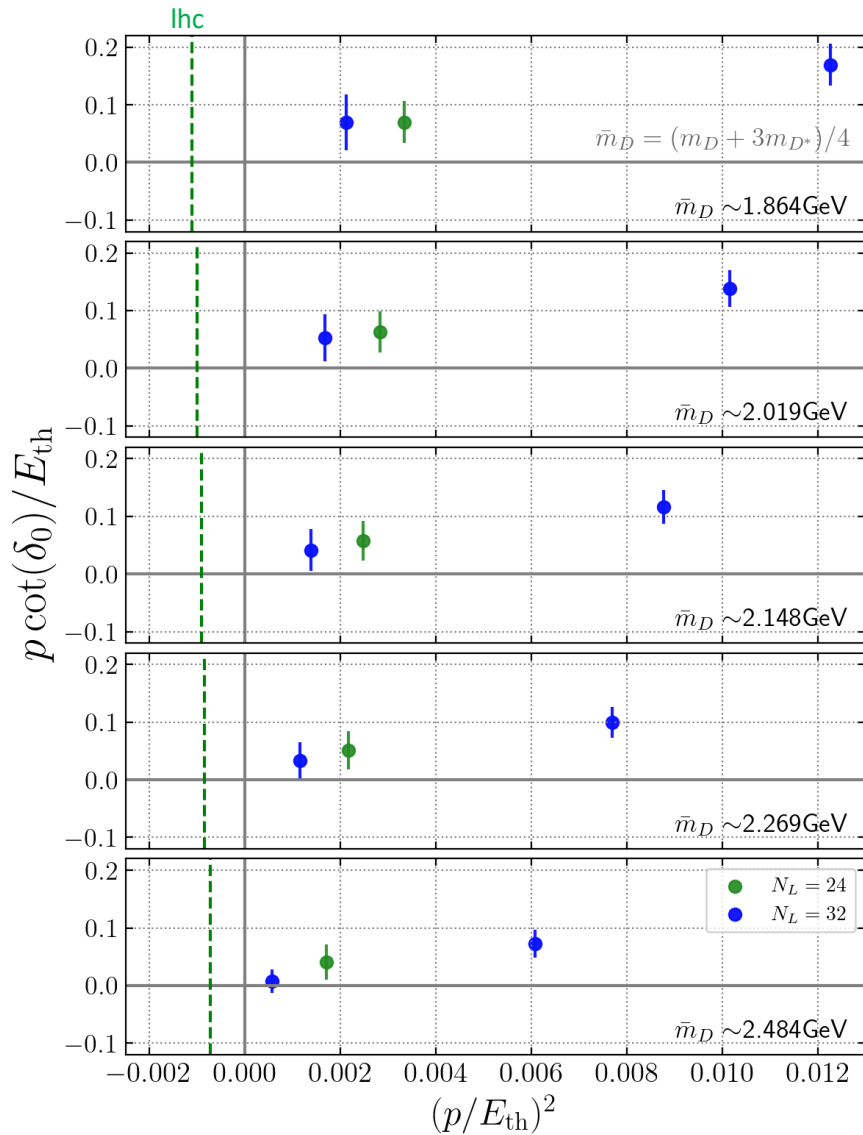
$$E^{n,i} = \sqrt{m_D^2 + \vec{p}_1^2} + \sqrt{m_{D^*}^2 + \vec{p}_2^2}$$

$$\vec{p}_i = \vec{n}_i \frac{2\pi}{L}$$

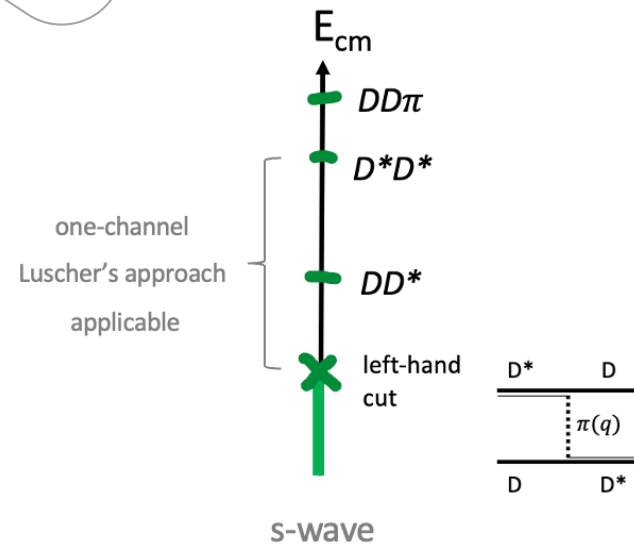
- all levels below threshold are omitted from the scattering analysis (throughout this talk)

# $T_{CC}$ : scattering amplitude

- all levels below threshold are omitted from the scattering analysis
- result on s-wave are consistent if p-wave included in analysis of  $A_2(P=1)$  or not
- p-wave constrained also from  $A_1^-(P=0)$  which is not shown

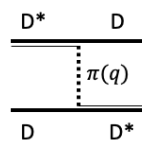
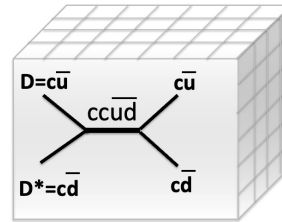
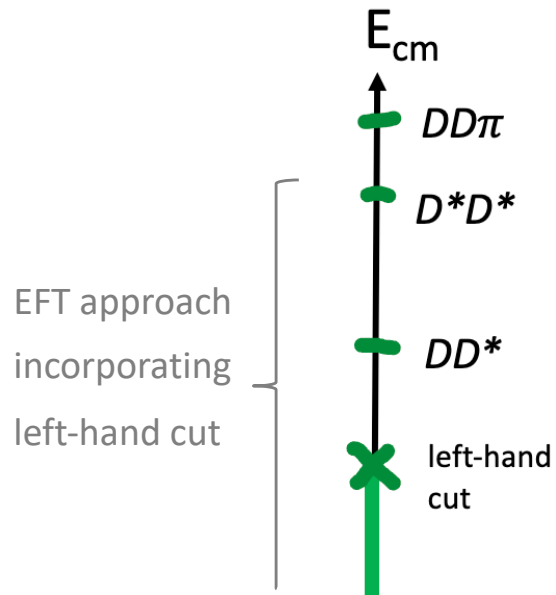


$\delta(E)$  E



# Plan: interpolating/extrapolating energy dependence of DD\* scattering amplitude

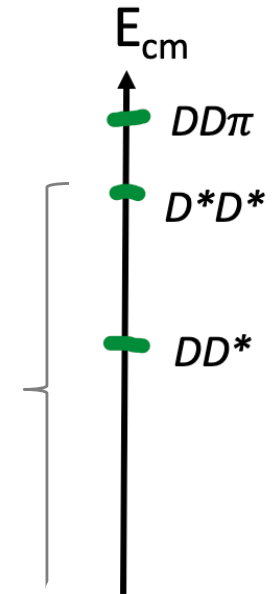
Part 1 (main result)



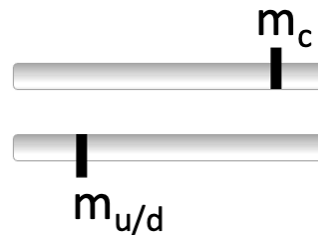
Part 2 (briefly)

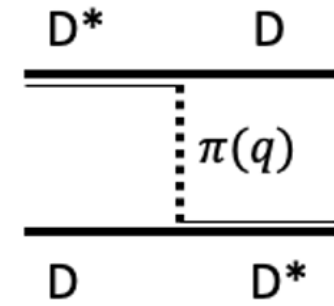
assuming eff. range  

$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$
 ignoring left-hand cut



In both cases:





## Analysis of $T_{cc}$ lattice results based on Effective Field Theory

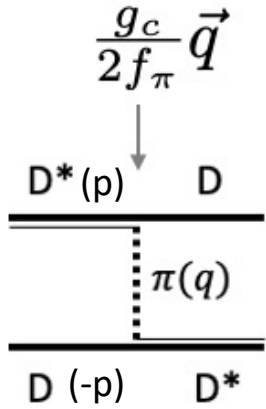
taking into account effect from left-hand cut



# Pion exchange, left-hand cut etc

Heavy meson ChPT

$g_c(m_\pi), f_\pi(m_\pi)$



$$q^2 = q_0^2 - \vec{q}^2 \simeq (m_{D^*} - m_D)^2 - \vec{q}^2$$

$$V_\pi^{cent}(\vec{q}) = \frac{g_c^2}{4f_\pi^2} \frac{\vec{q}^2}{q^2 - m_\pi^2} = \frac{g_c^2}{4f_\pi^2} \left( -1 + \frac{\mu_\pi^2}{\vec{q}^2 + \mu_\pi^2} \right)$$

$$\mu_\pi^2 = m_\pi^2 - (m_{D^*} - m_D)^2$$

lat :  $\mu_\pi^2 > 0$

ph :  $\mu_\pi^2 < 0$

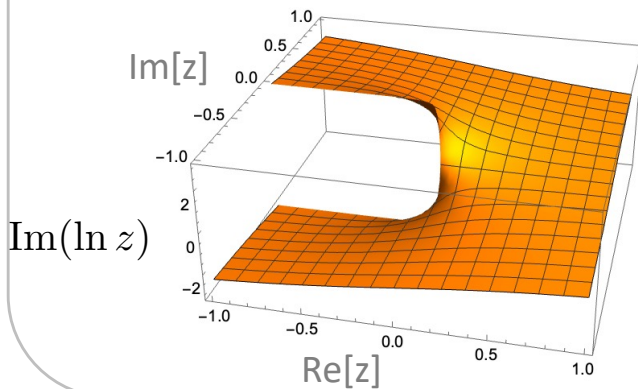
attraction at short distance      slight repulsion at long distance

$$-\delta^{(3)}(\vec{r}) \quad \frac{\mu_\pi^2}{r} e^{-\mu_\pi r}$$

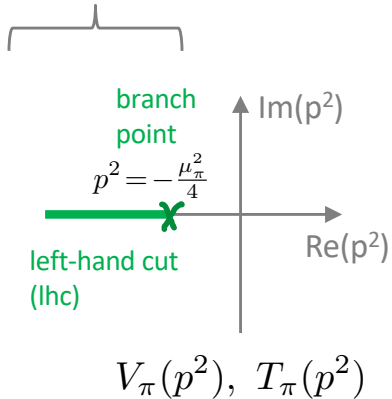
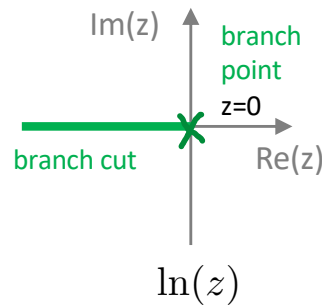
s-wave projection

$$V_\pi^S(p, p) \propto \int V_\pi(\vec{q}) d\cos\theta, \quad \vec{q}^2 = 2p^2(1 - \cos\theta)$$

$$V_\pi^S(p, p) \propto \ln\left(1 + \frac{4p^2}{\mu_\pi^2}\right)$$



complex  $p \cot \delta$  (Luscher's eq would render it real)



lhc slightly below DD\*, BB\*, NN ... th.

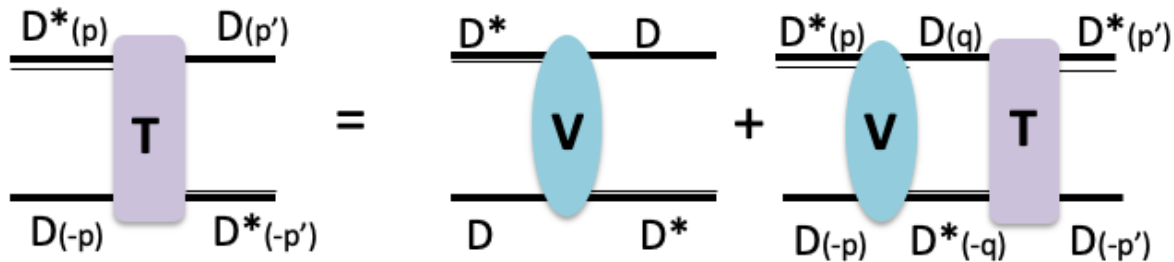
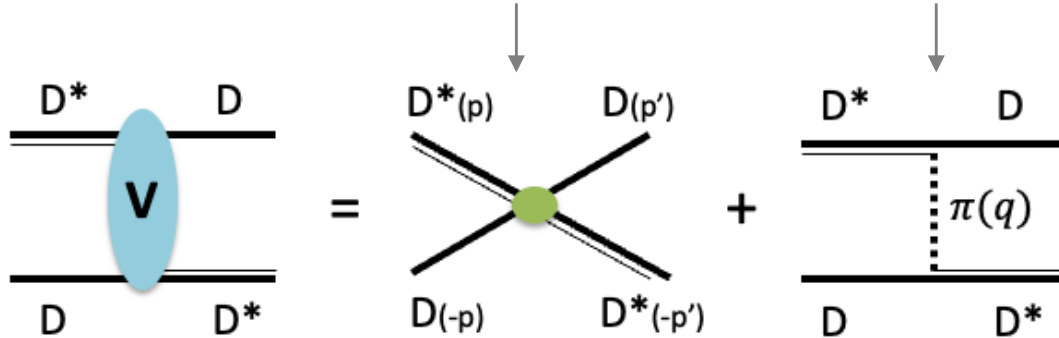
# $T_{cc}$ analysis based on EFT

$c_{0,2}$  fitted from lat. data

significant short-distance attraction

$$V_{CT} = 2c_0 + 2c_2(p^2 + p'^2)$$

$$\frac{g_c}{2f_\pi} \vec{q}$$



$$T(\mathbf{p}, \mathbf{p}'; E) = V(\mathbf{p}, \mathbf{p}') - \int \frac{d^3q}{(2\pi)^3} V(\mathbf{p}, \mathbf{q}) G(\mathbf{q}; E) T(\mathbf{q}, \mathbf{p}'; E)$$

Lipmann-Schwinger eq.

$$T = V - VGT$$

$$T = \frac{1}{V^{-1} + G}$$

$$|\vec{q}| < \Lambda = 0.5 \text{ GeV}$$

cut-off on q:

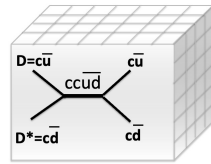
$c_{0,2}$  depend on  $\Lambda$ ,

poles dont

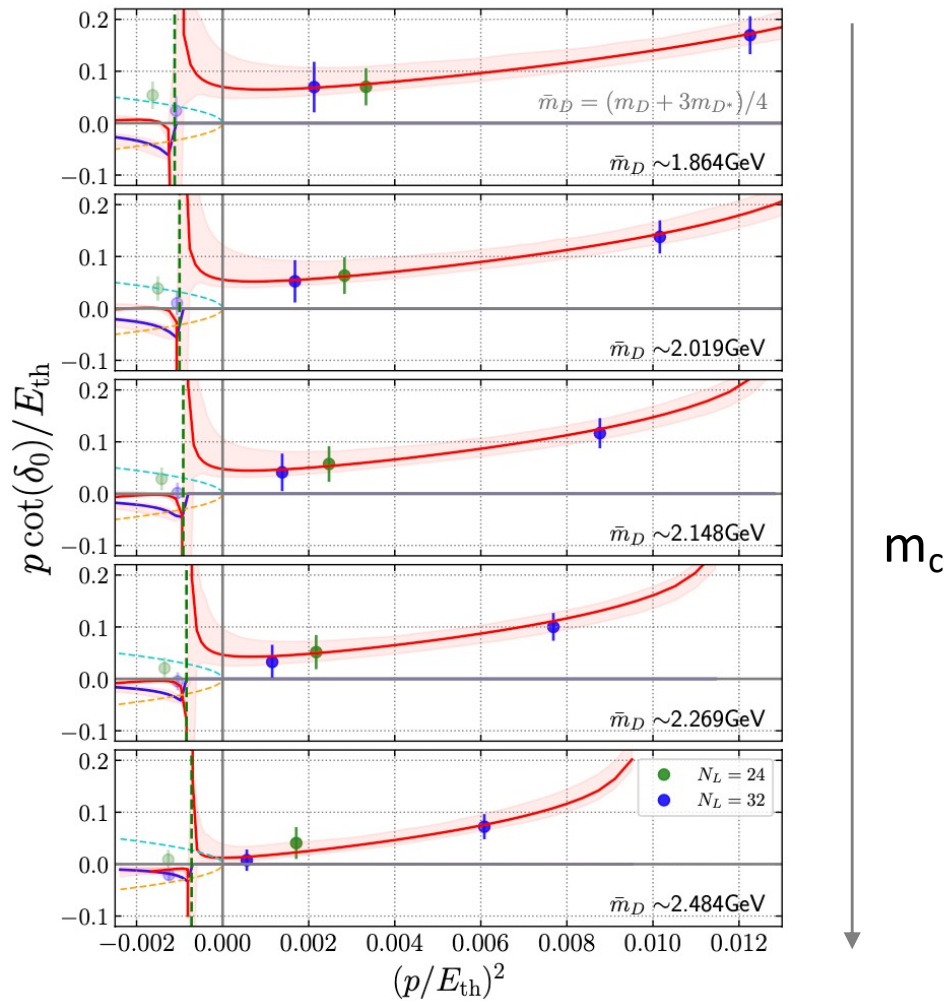
inspired by

Du, Hanhart, Guo, Nefediev, Filin, et al, PRL 2023, 2303.09441

# $T_{cc}$ : pole trajectory



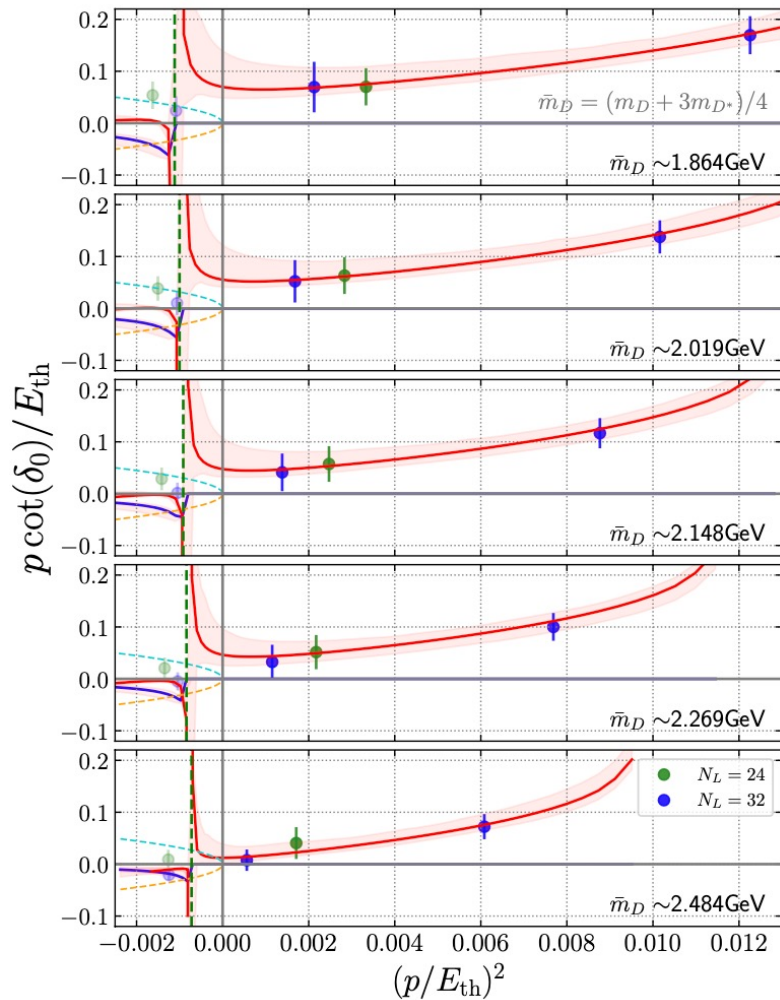
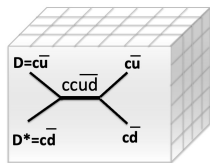
$m_\pi \simeq 280$  MeV



levels below lhc omitted from the fit

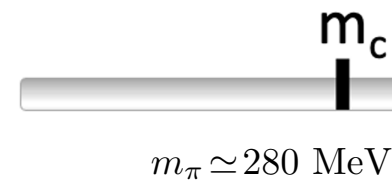
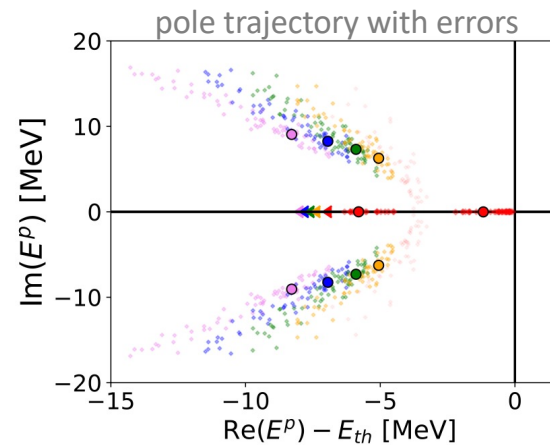
it is reassuring that consistent results are obtained via plane-wave method, which can incorporate also levels on lhc, [Meng, Baru, Epelbaum et al., 2312.01930, PRD](#); [Vujmilovic \(next talk\)](#)

# $T_{cc}$ : pole trajectory

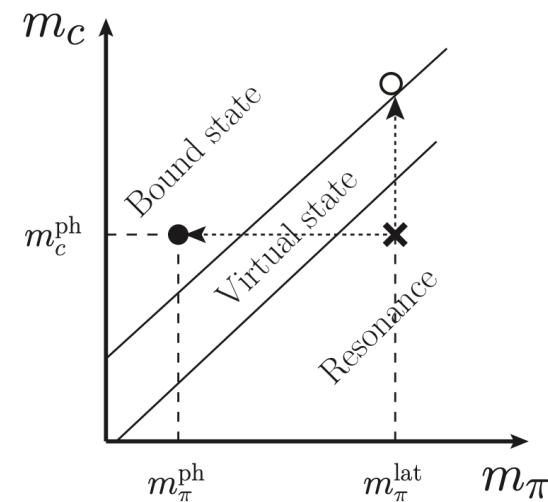
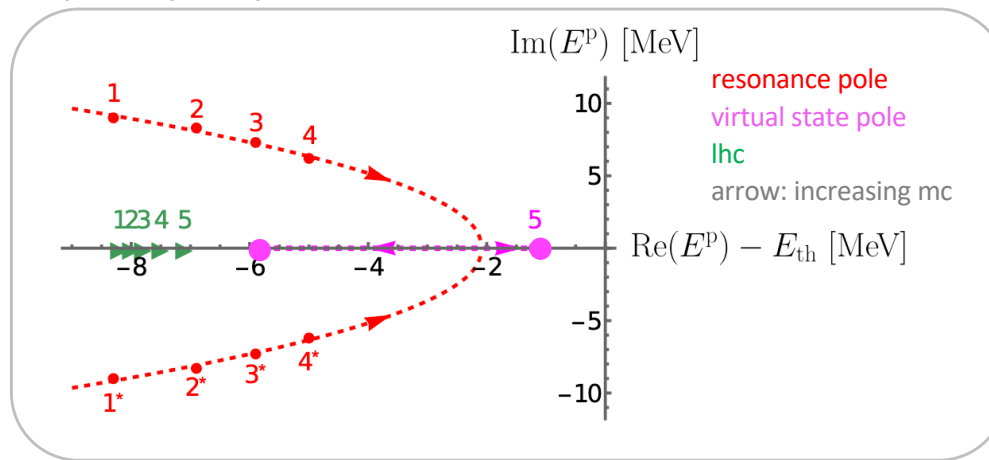


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it is reassuring that consistent results are obtained via plane-wave method, which can incorporate also levels on lhc, Meng, Baru, Epelbaum et al., 2312.01930, PRD; Vujmilovic (next talk)



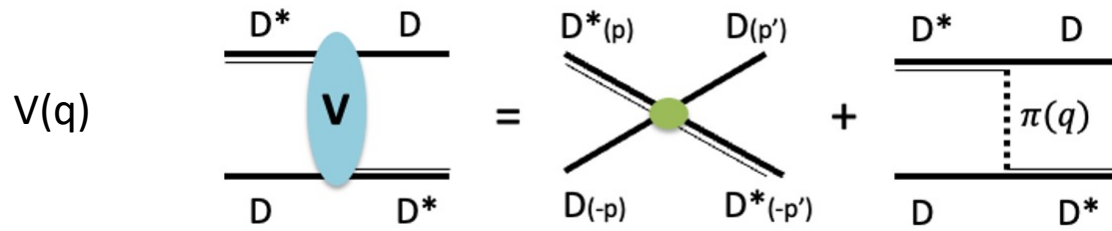
pole trajectory: central values



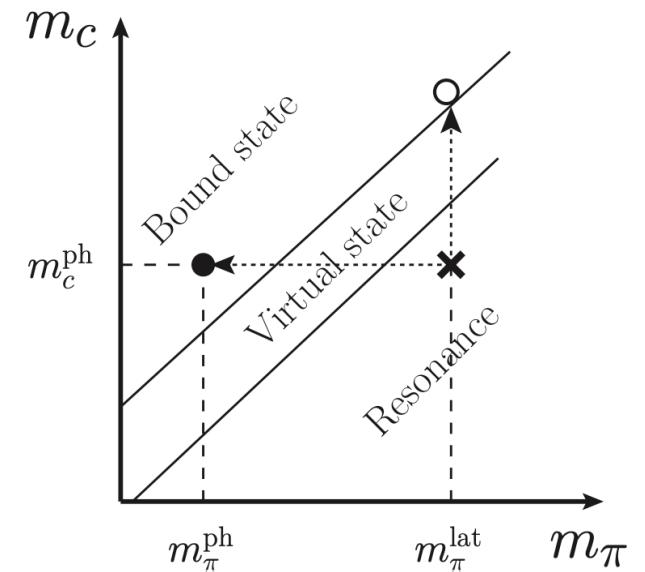
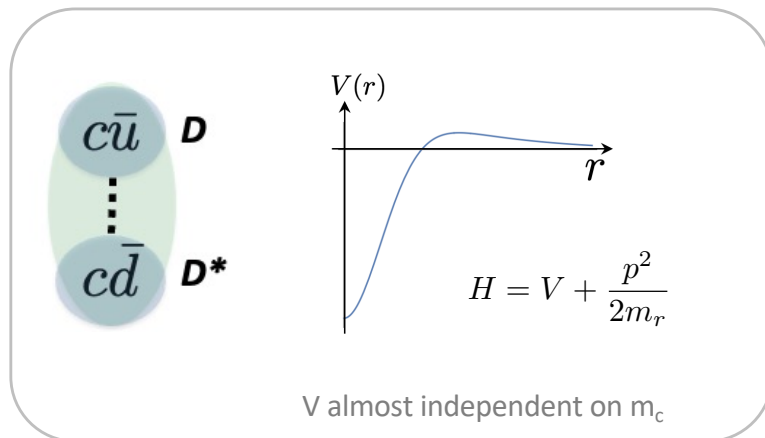
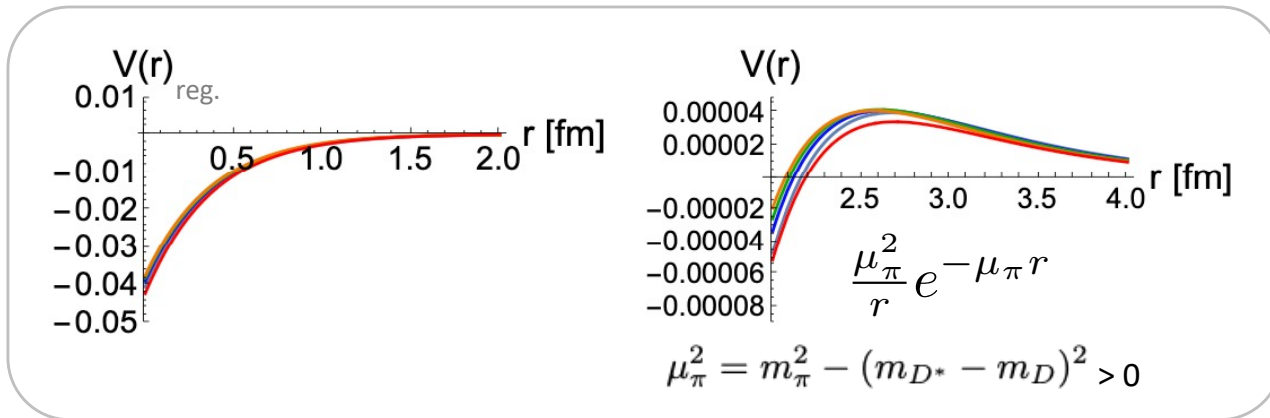
# $T_{cc}$ : interpretation



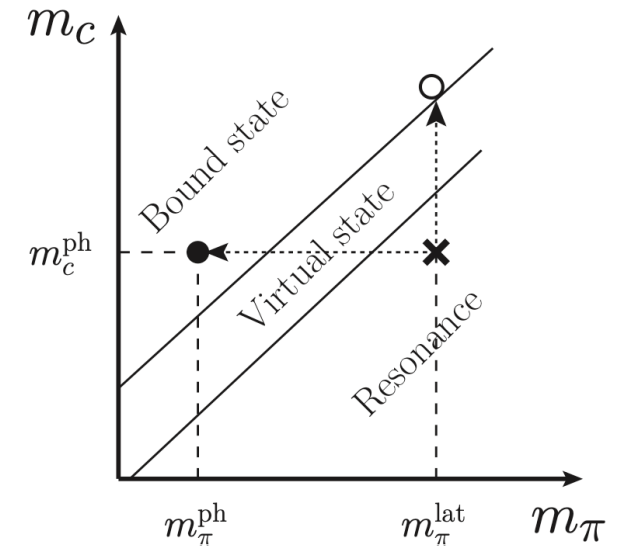
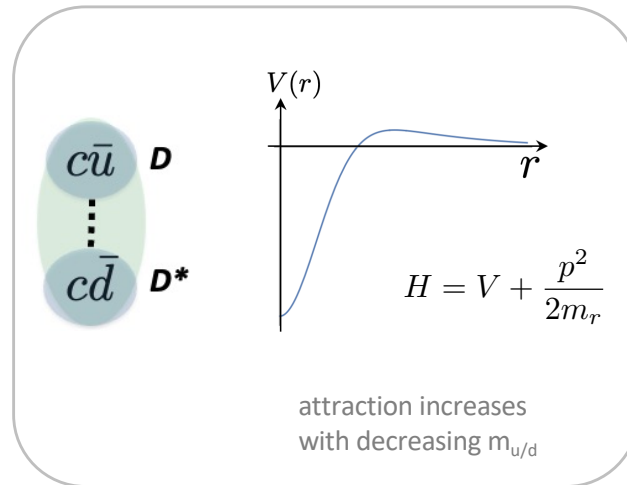
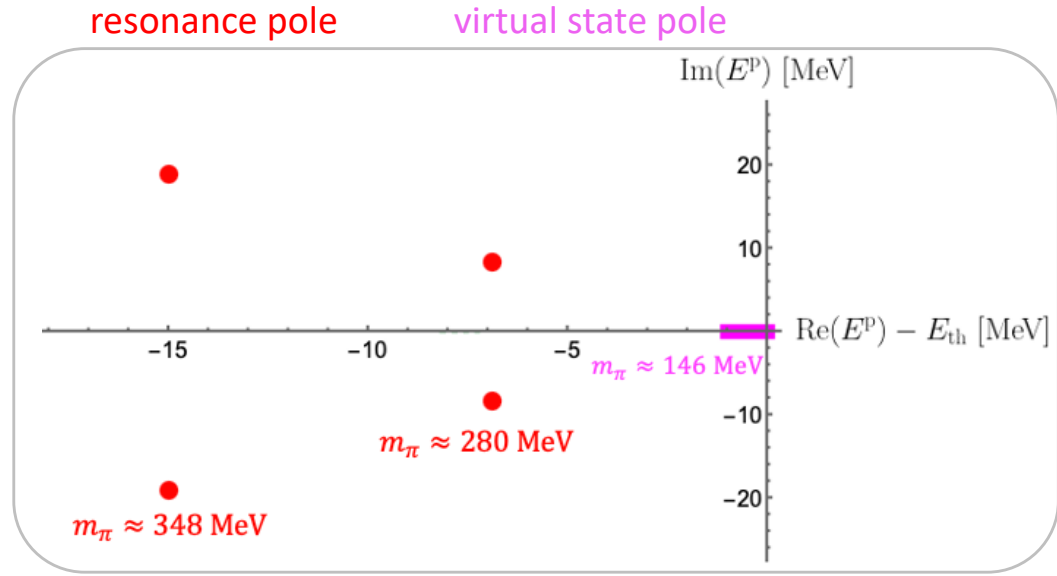
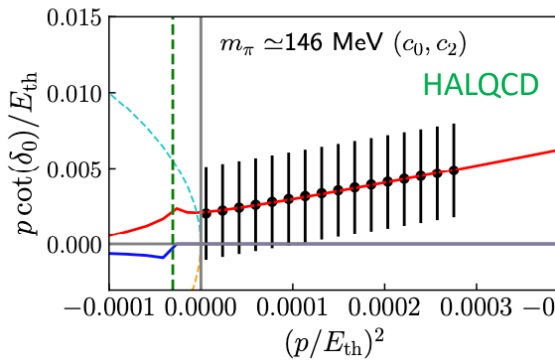
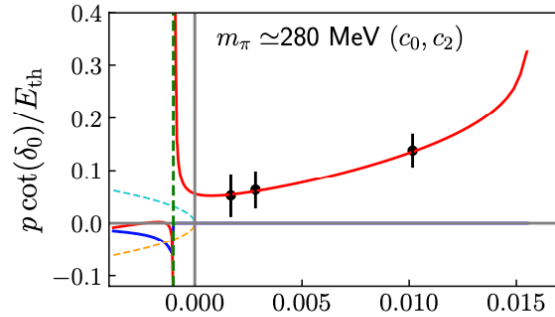
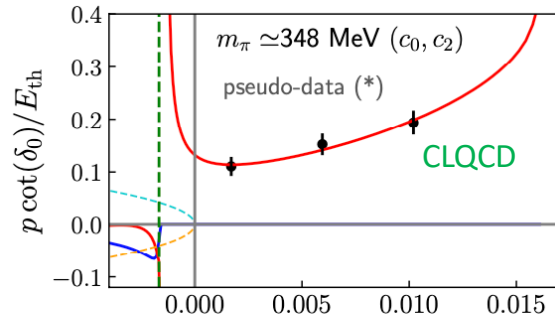
$$m_\pi \simeq 280 \text{ MeV}$$



↓ FT



# $T_{cc}$ analysis based on EFT



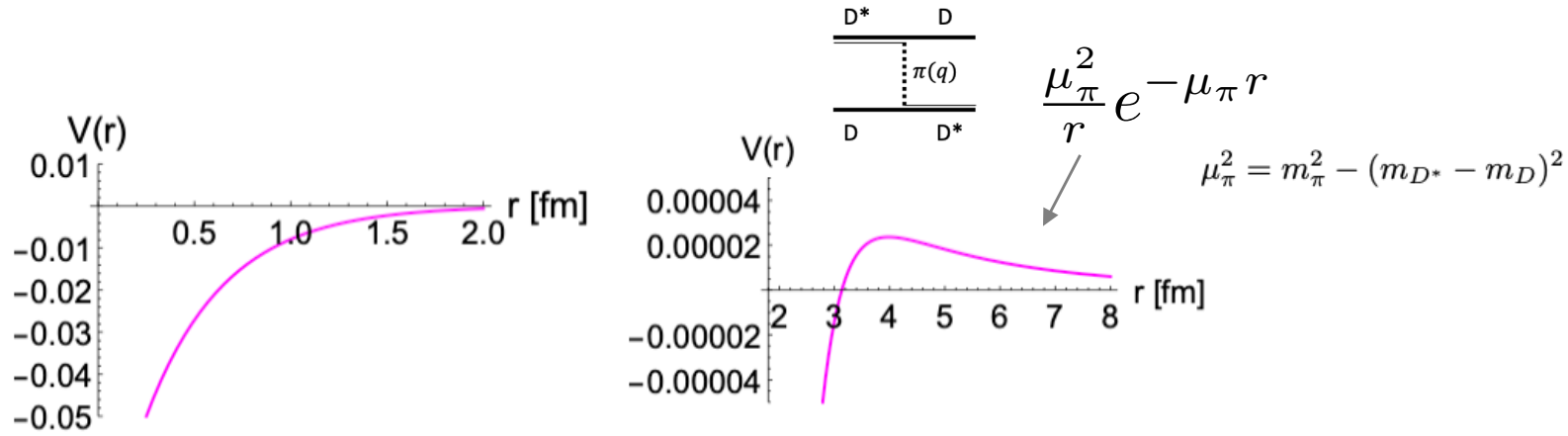
errors on  $m_{u/d}$  dependence not reliably determined  
(see disclaimers in our paper)

see also: [2407.04649](#),

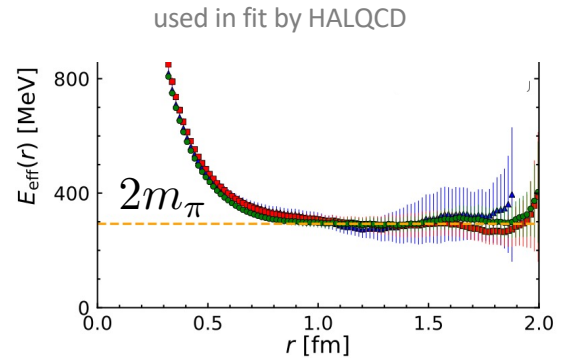
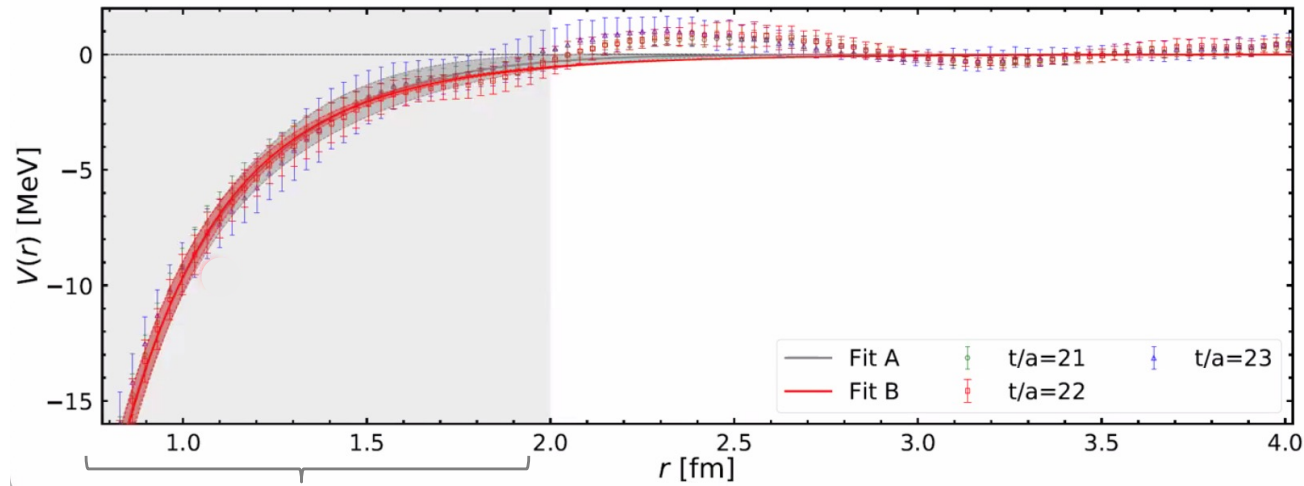
Abolnikov, Baru, Epelbaum, Filin, Hanhart, Meng

# Signature of one-pion exchange in HALQCD potential ?

$$m_\pi \simeq 146 \text{ MeV}$$



HALQCD, 2302.04505, PRL, talk by Yan Lyu



## Analysis of $T_{cc}$ lattice results assuming effective range expansion

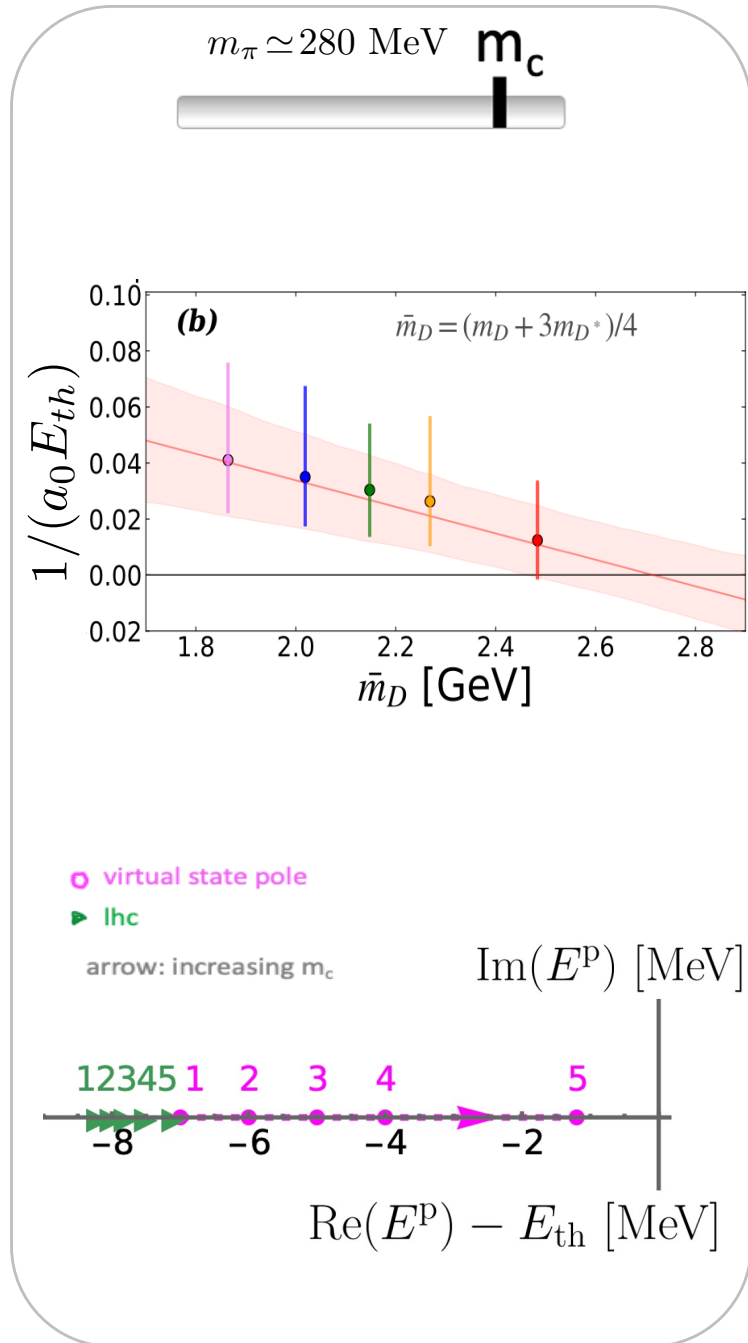
omitting effects from left-hand cut

$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2 + \dots$$



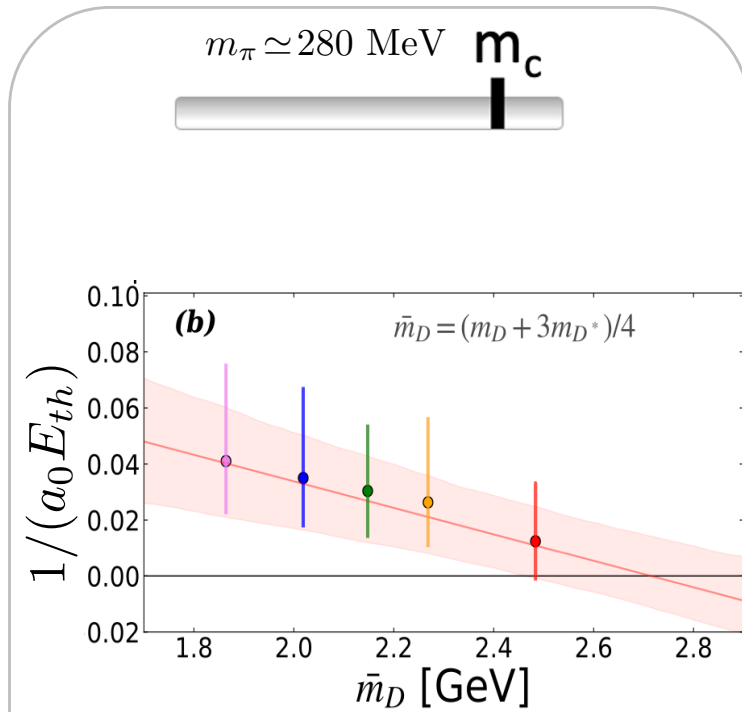
$T_{CC}$  assuming eff. range expansion

$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$



$T_{CC}$  assuming eff. range expansion

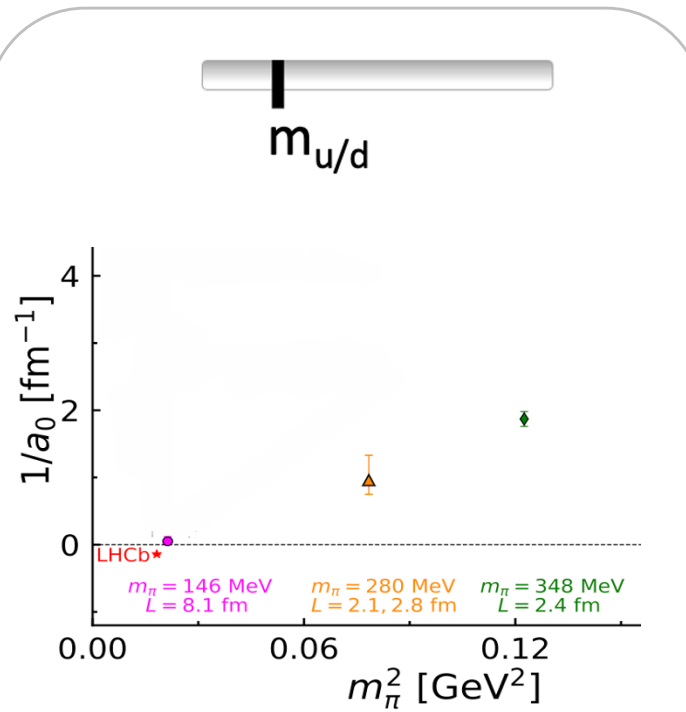
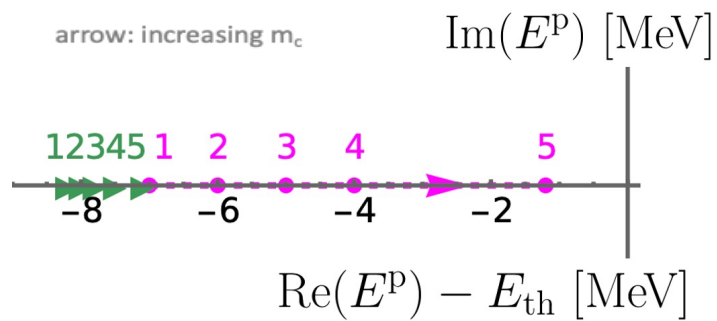
$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$



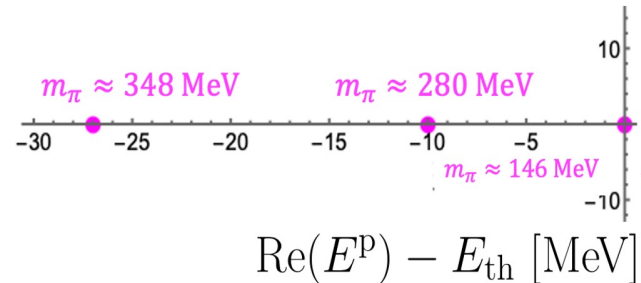
○ virtual state pole

▶ lhc

arrow: increasing  $m_c$

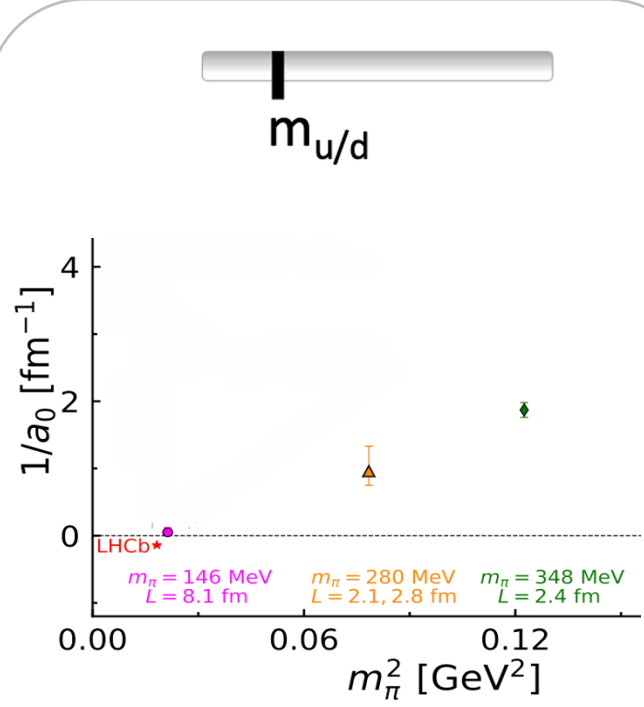
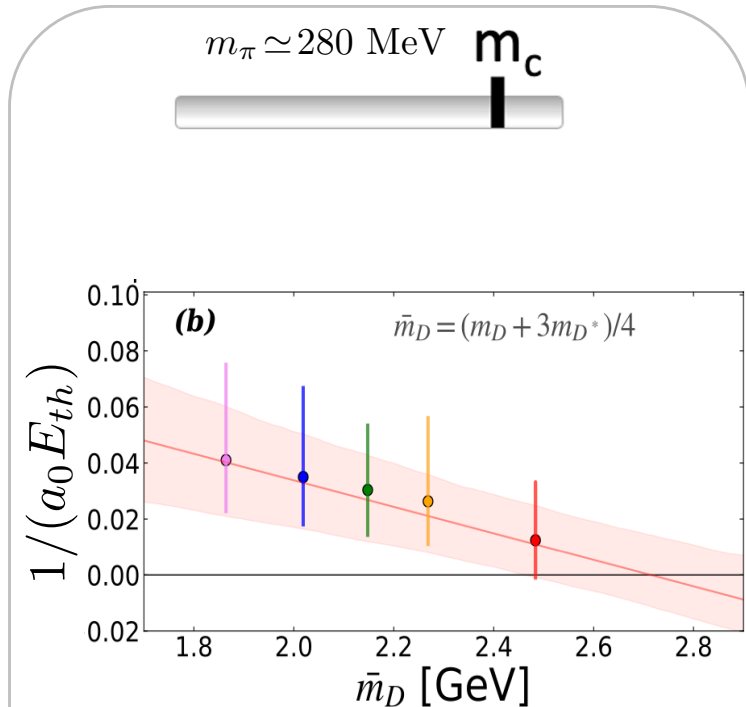


$\text{Im}(E^p)$  [MeV]

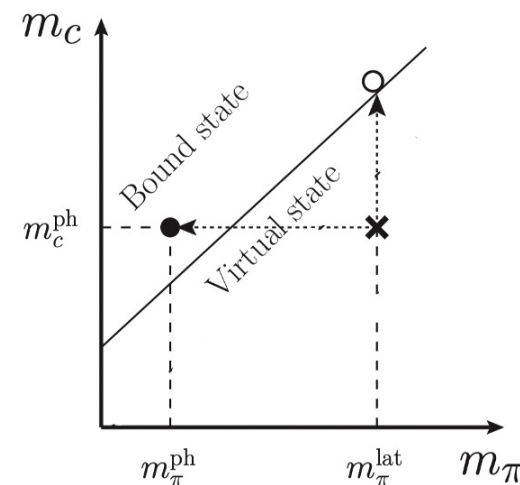
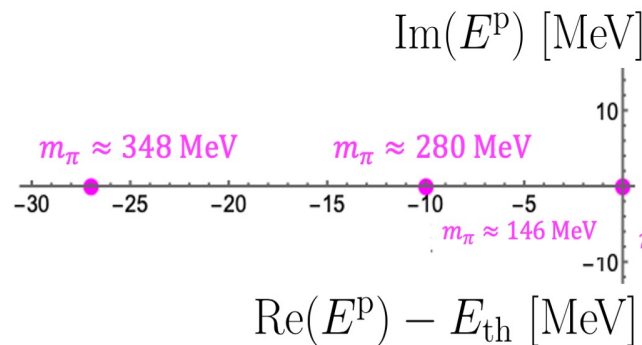
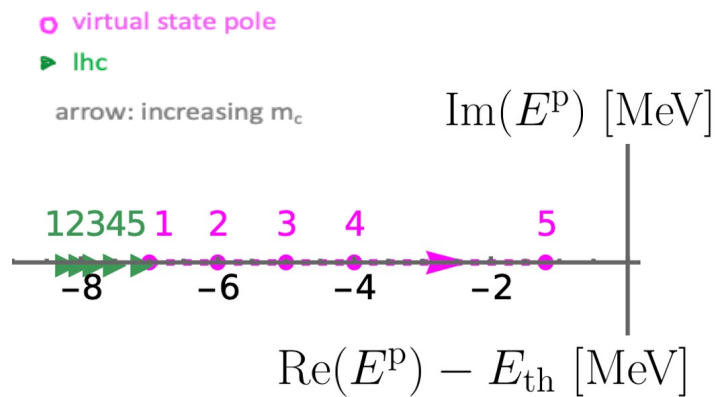
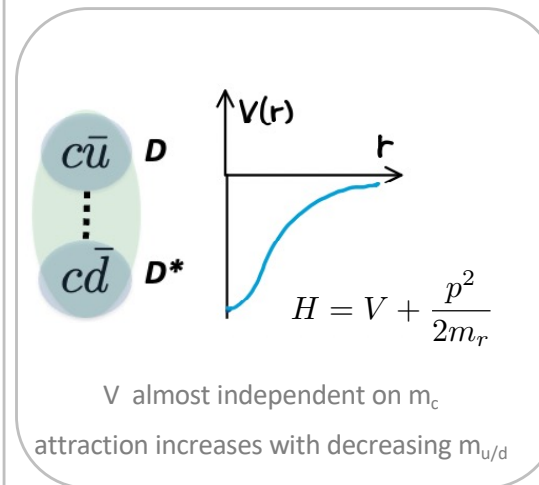


$T_{CC}$  assuming eff. range expansion

$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$

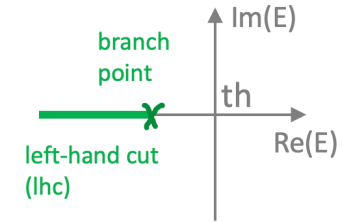


interpretation



# Scattering on the lattice and the left-hand cut

$$E_n \leftrightarrow T(E_n)$$



- **two-body formalism**

- EFT approach here inspired by: [Du, Hanhart, Guo, Nefediev, Filin, et al, PRL 2023, 2303.09441](#)
- plane-wave approach : [Meng, Epelbaum, 2108.02709](#); [Meng, Baru, Epelbaum et al., 2312.01930, PRD](#)  
Vujmilovic: next talk
- generalization of Luscher's equation on left-hand cut [[Raposo, Hansen, 2311.18793](#)]
- Luscher eq. with long-range forces [[Bubna, Hammer, Muller, Pand, Rusetsky, Wu, 2402.12985](#)], comparison with existing approaches

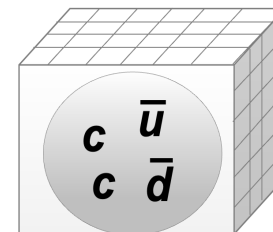
- **three-body formalism**

- $T_{cc}$  channel,  $D^*$  as  $D\pi$  bound state, [Romero-Lopez, Sharpe, Hansen, Draper \[2401.06609](#)
- [Islam, Dawid, Briceno, Lattice 2023, 2303.04394](#) (signal for break-down of two-body Luscher's equation)

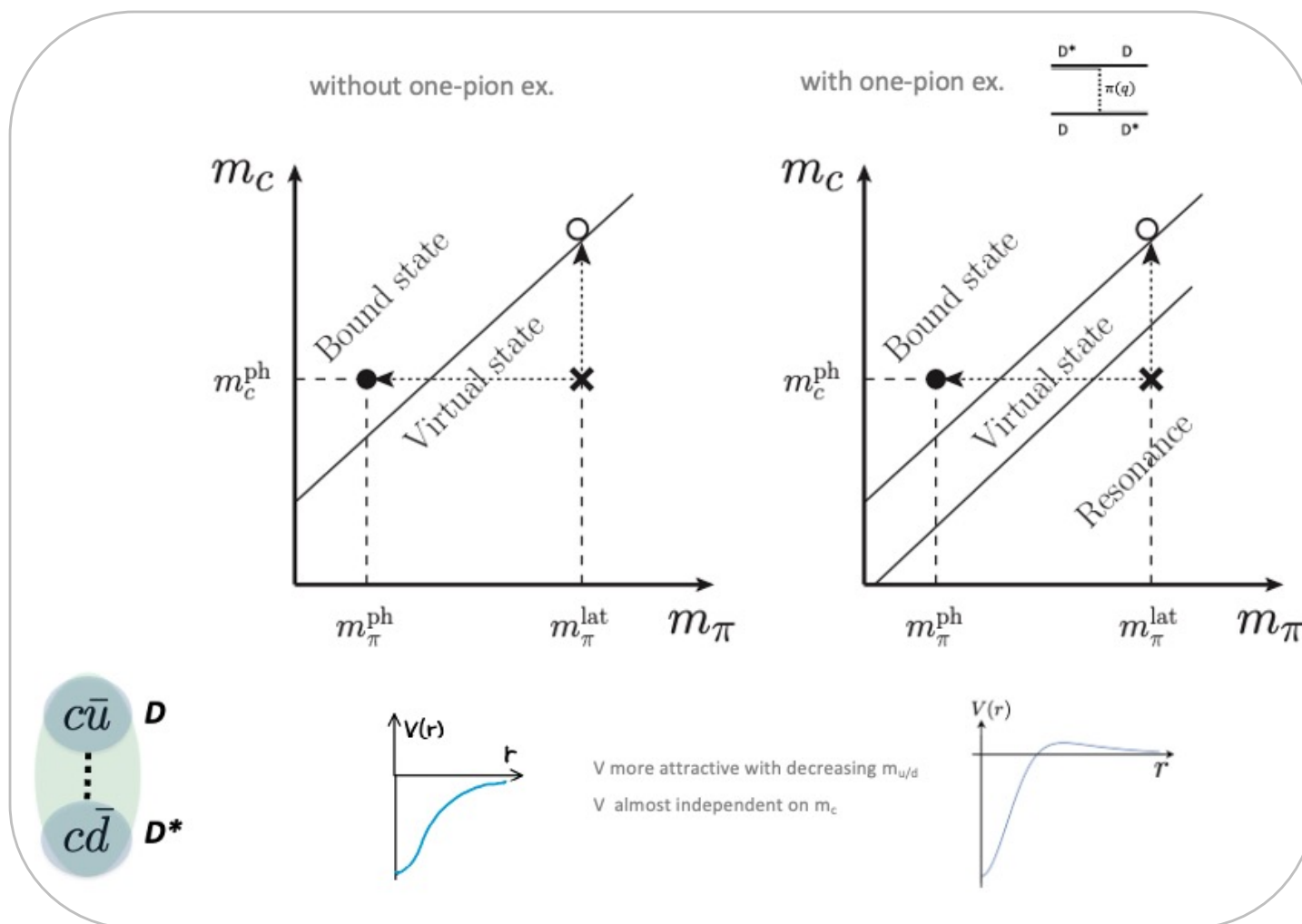
# Conclusions

$T_{cc}$

- $T_{cc}$  is the longest-lived exotic hadron discovered in experiment
- lies near threshold  $\rightarrow$  has to be extracted from  $DD^*$  scattering amplitude
- lattice studies find attraction



quark mass dependence of  $T_{cc}$ . for stable  $D^*$ ,  $m_\pi > m_{D^*} - m_D$

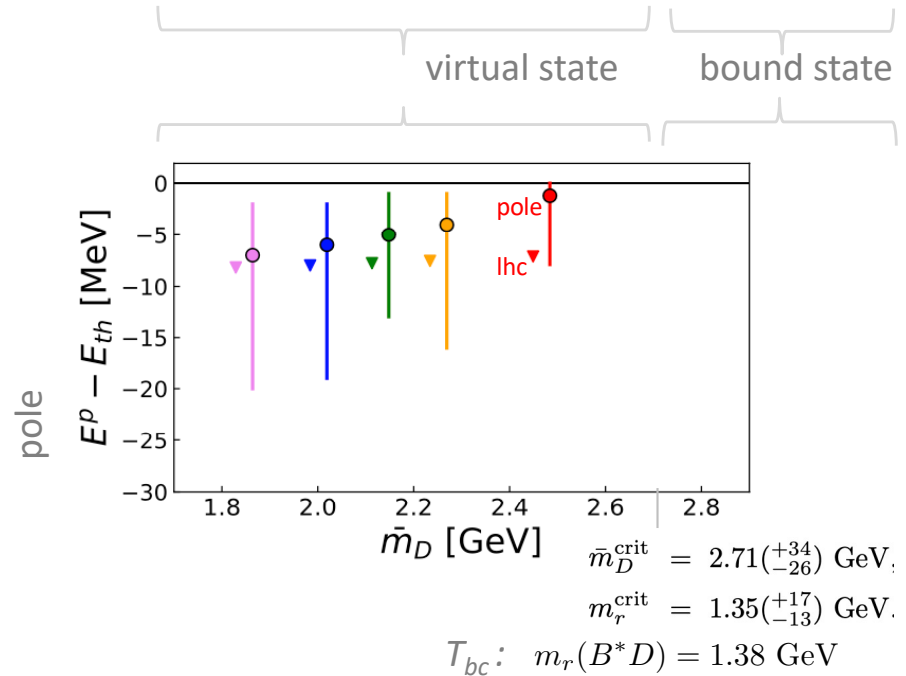
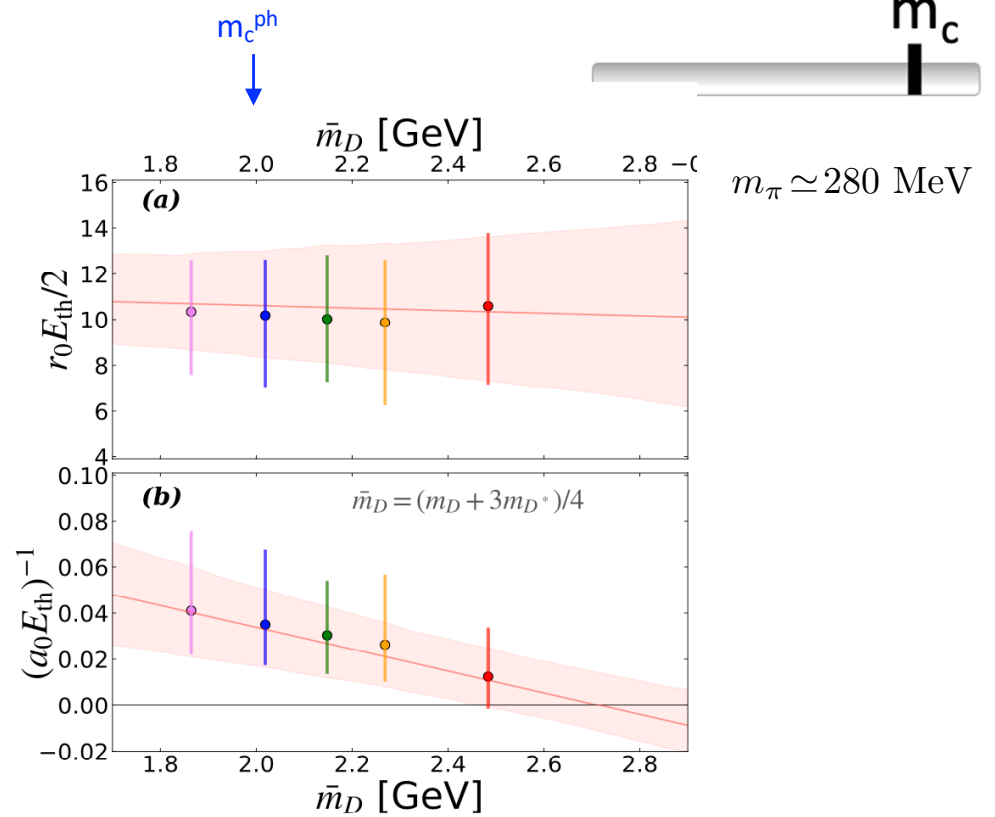
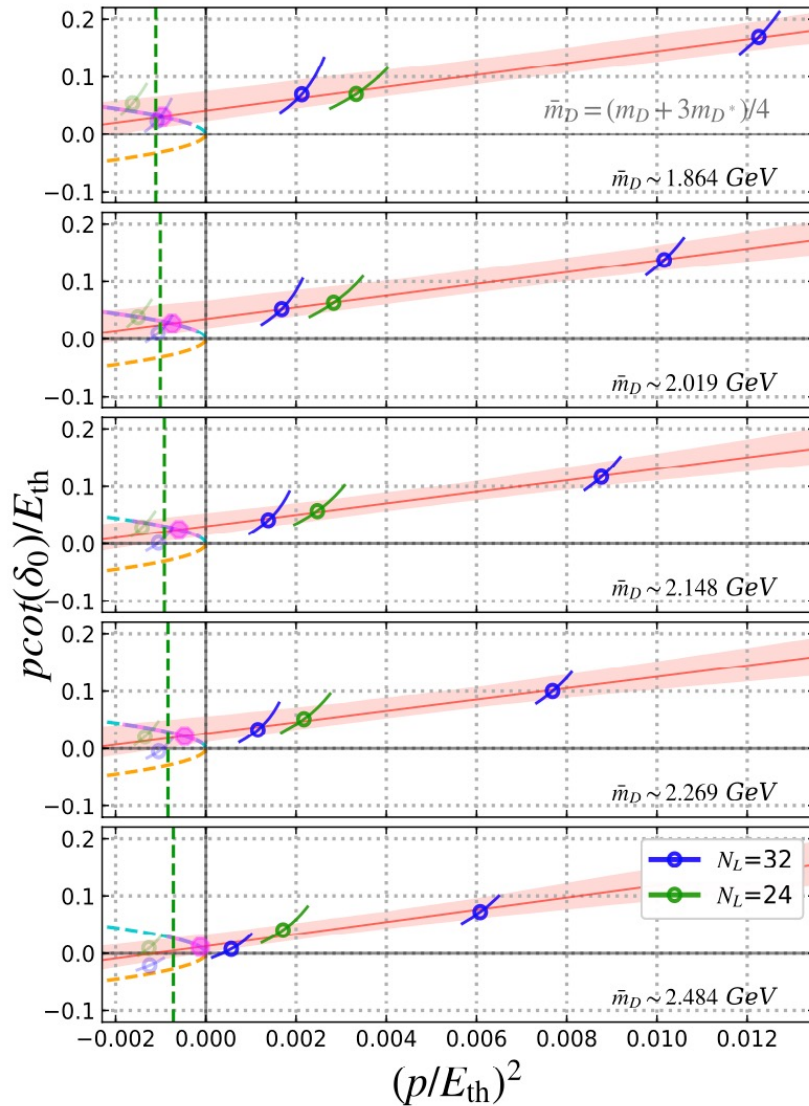
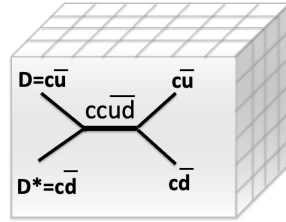


# Backup

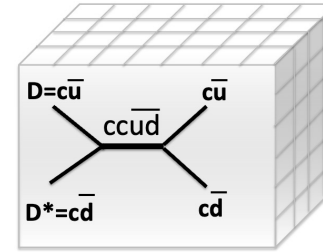
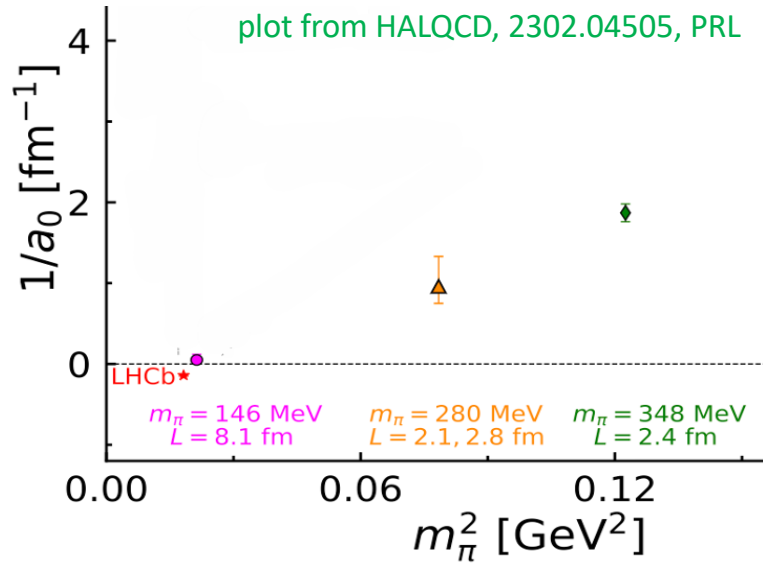
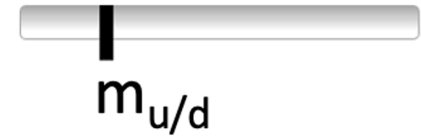
# $T_{CC}$ assuming eff. range expansion

$$T \propto \frac{1}{p \cot \delta - ip} \quad \text{s-wave}$$

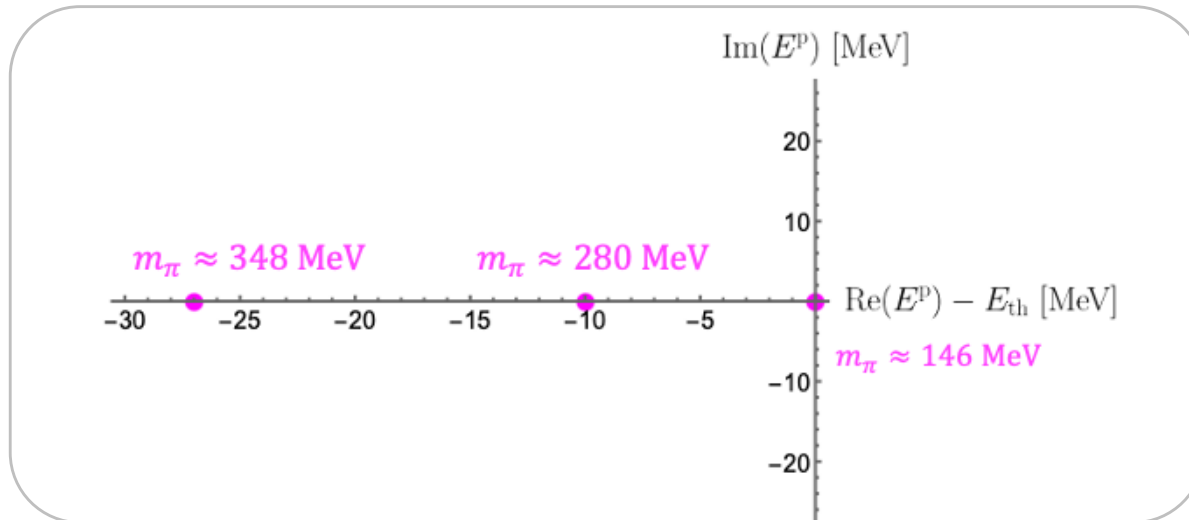
$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$



# $T_{cc}$ assuming effective range expansion: pole trajectory



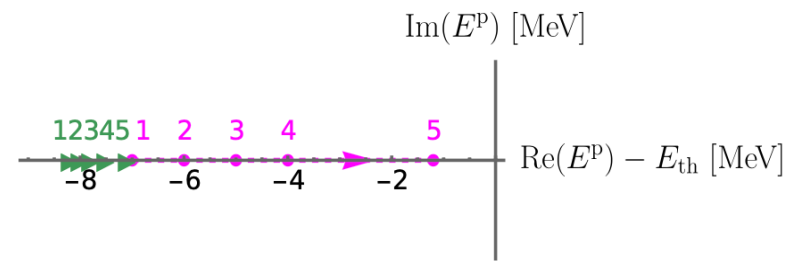
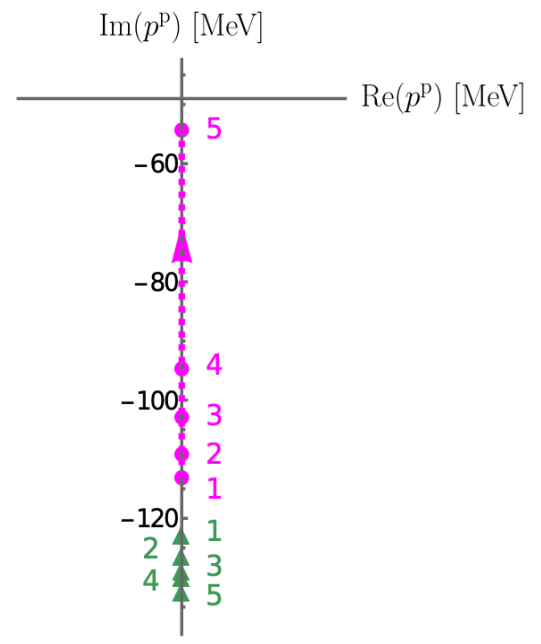
virtual state pole



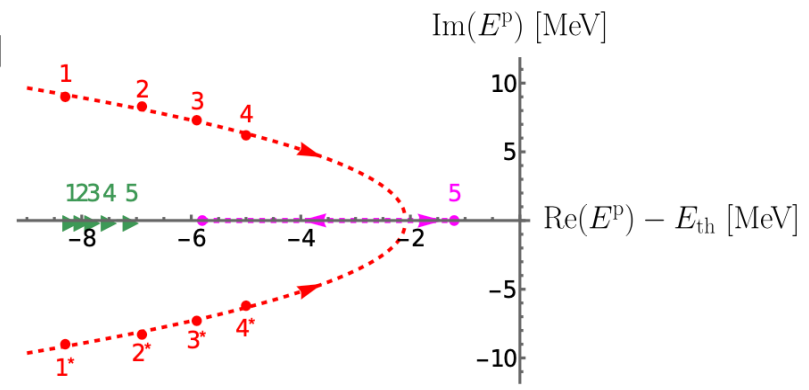
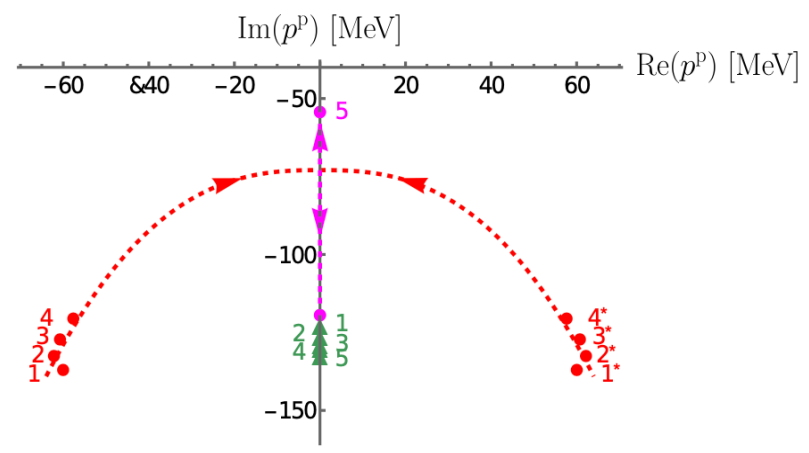




effective  
range  
expansion



EFT



# $T_{cc}$ : plane wave approach

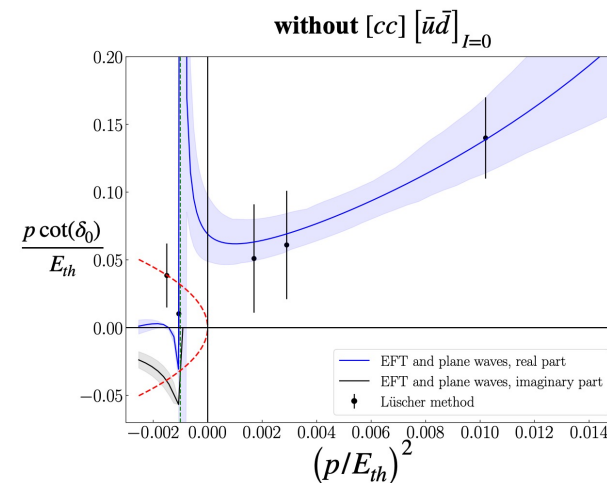
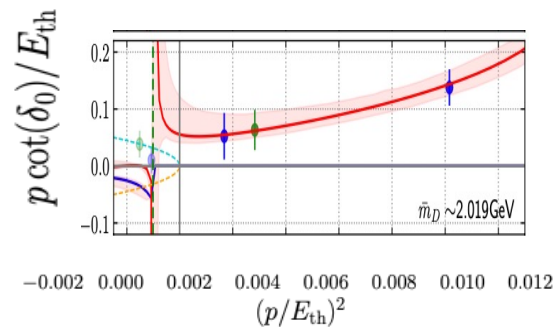
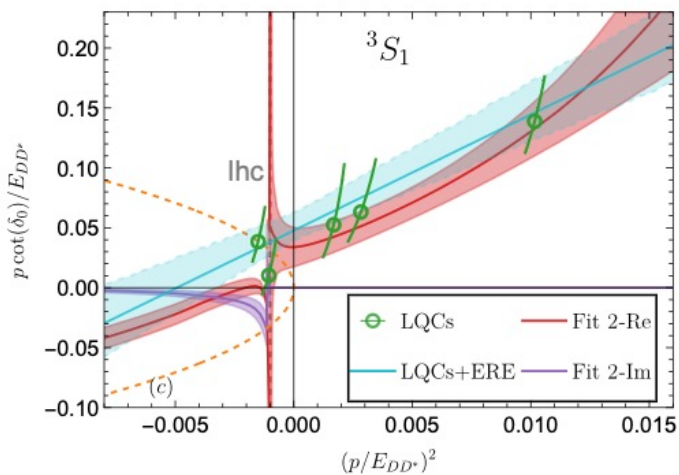
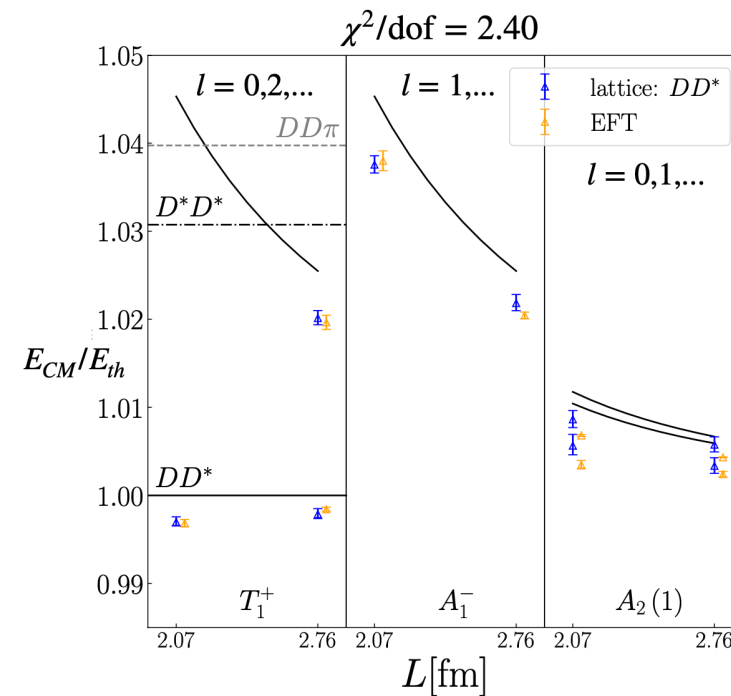
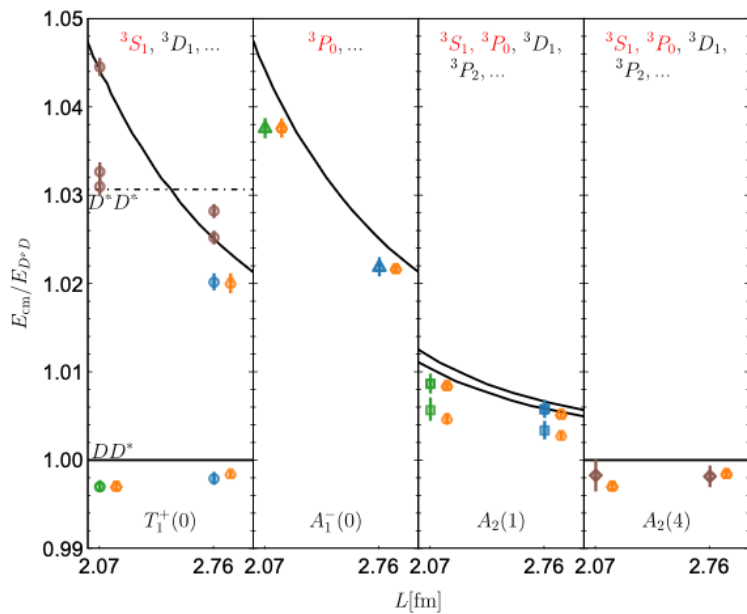
$m_{\pi} \sim 280 \text{ MeV}$ ,  $m_c^{\text{ph}}$

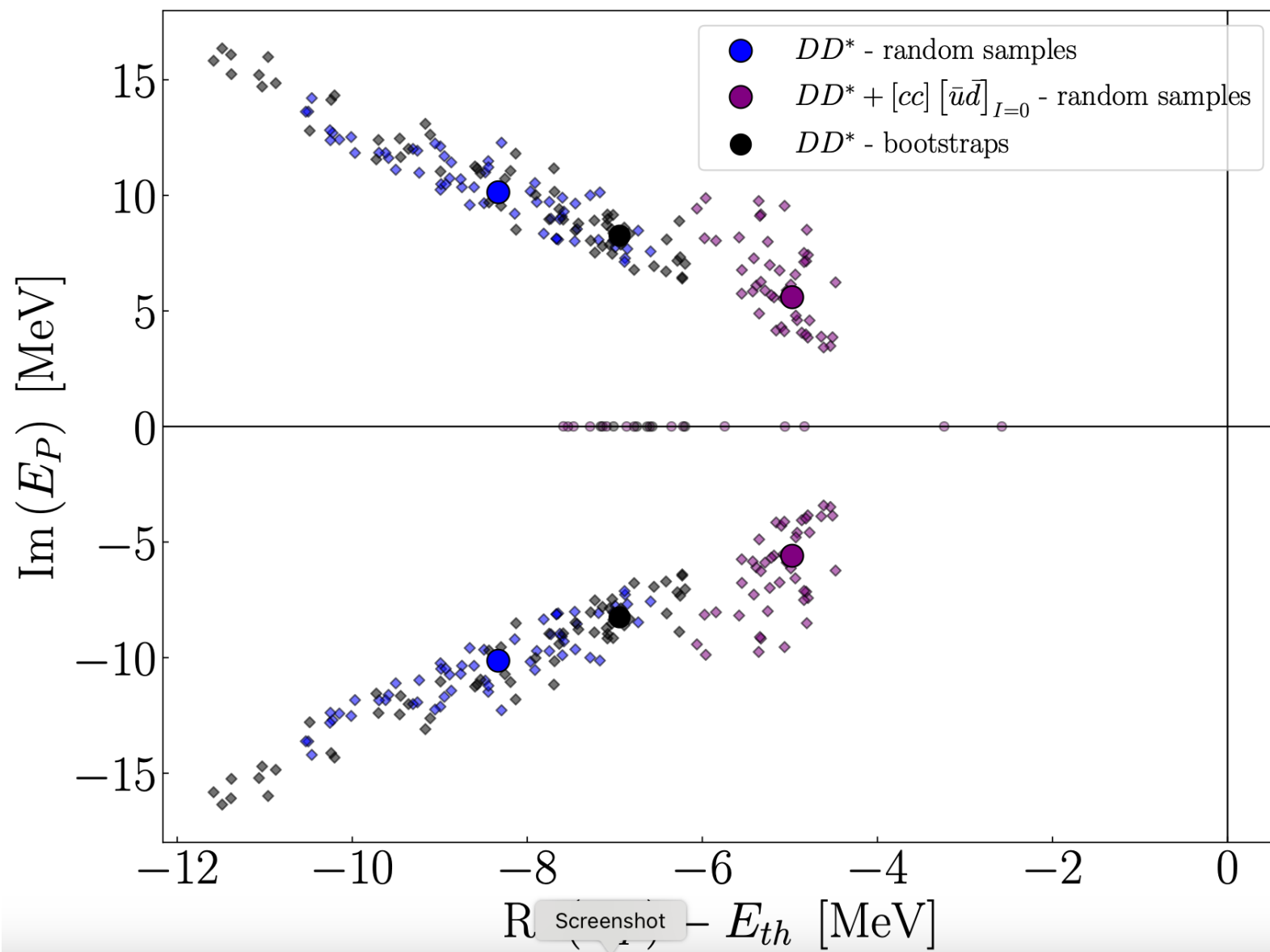
Meng, Baru, Epelbaum et al., 2312.01930,

this talk, 2402.14715

talk by Ivan Vujmilovic @ lat24

orange: reconstructed levels, blue, green: lattice levels

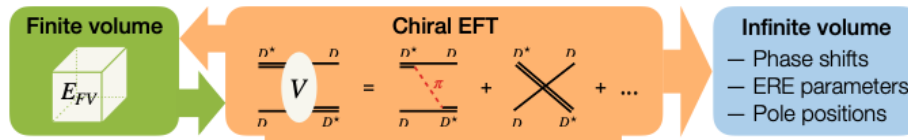




this talk  
plane waves  
plane waves

# $T_{cc}$ : plane wave approach

Meng, Baru, Epelbaum et al., 2312.01930, PRD  
 formalism: Meng, Epelbaum, 2108.02709;  
 lattice data: Padmanath, SP. 2202.10110, PRL



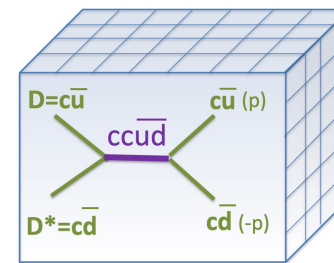
$$H = \frac{p^2}{2m_r} + V_\pi + V_{CT}$$

$$Hv = Ev$$

$$H_\Gamma v = E_\Gamma v$$

# $T_{cc}$ : dependence on $m_{u,d}$

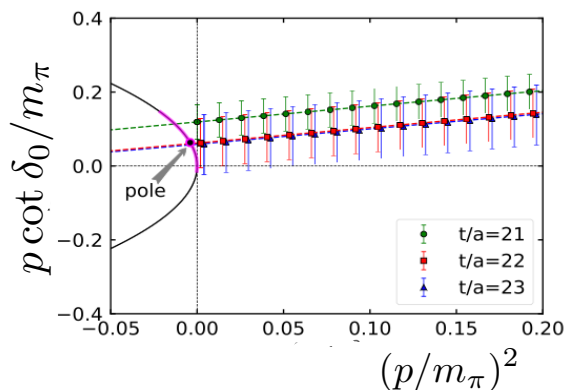
$$p \cot \delta_0 = \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$



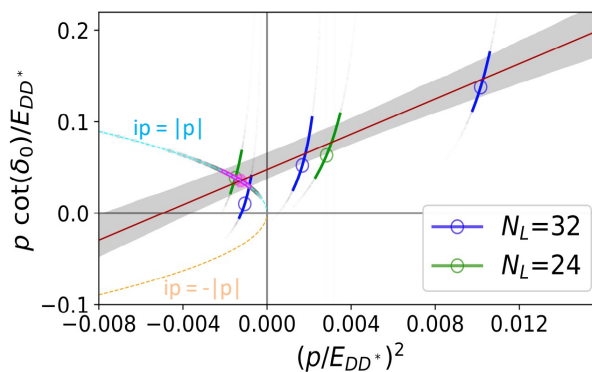
$m_{u,d}$

LHCb

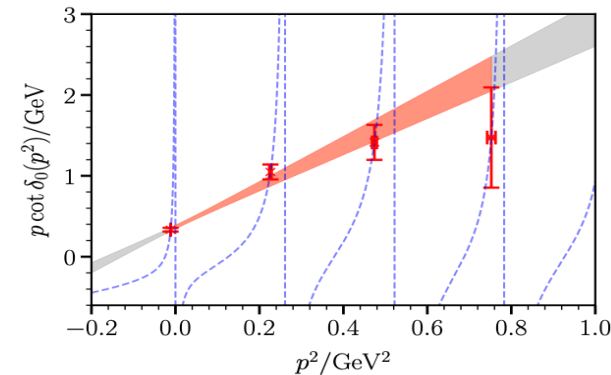
HALQCD method, 2302.04505,  $m_\pi \approx 146$  MeV



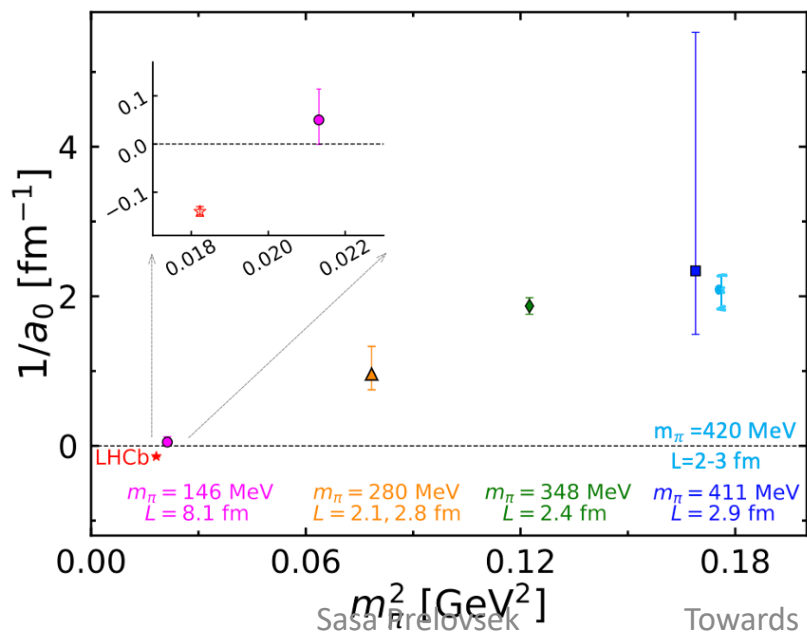
Padmanath, SP: 2202.10110, PRL,  $m_\pi \approx 280$  MeV



CLQCD 2206.06185, PLB,  $m_\pi \approx 348$  MeV

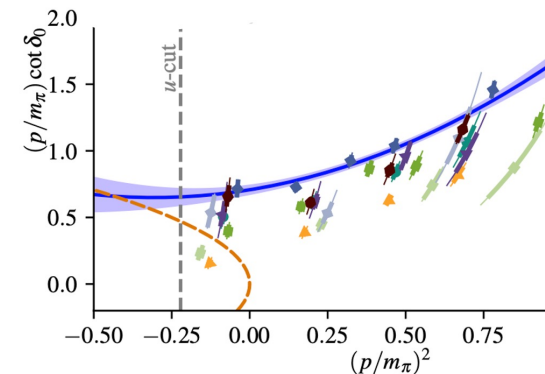


adapted from 2302.04505

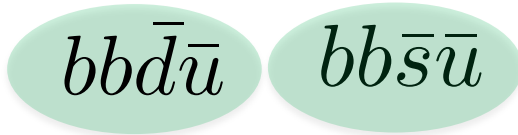


several a

Green et al, Mainz, Lat 2023,  $m_\pi \approx 420$  MeV

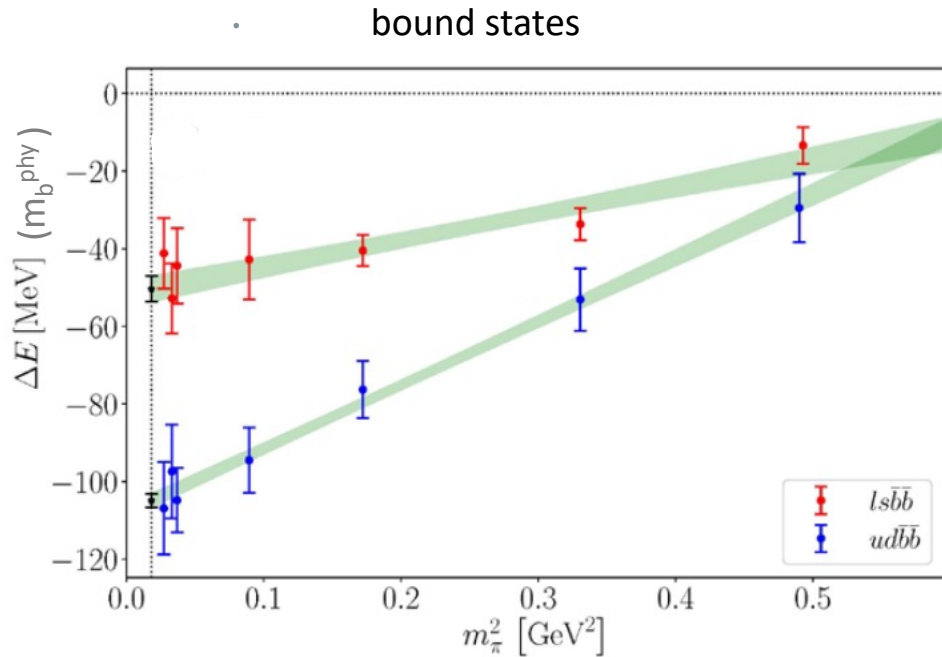


# Doubly bottom tetraquarks

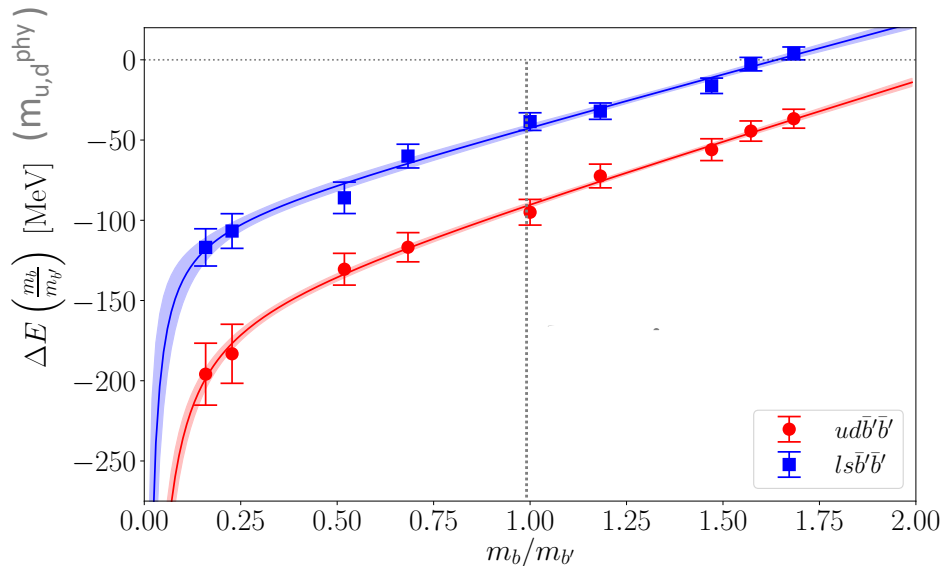


$I=0, J^P=1^+$

Frances, Colquhoun, Lewis, Maltman, Hudspith  
PoS LATTICE2021 (2022) 144



$m_{u,d}$  increases →



$m_{b'}$  decreases →

