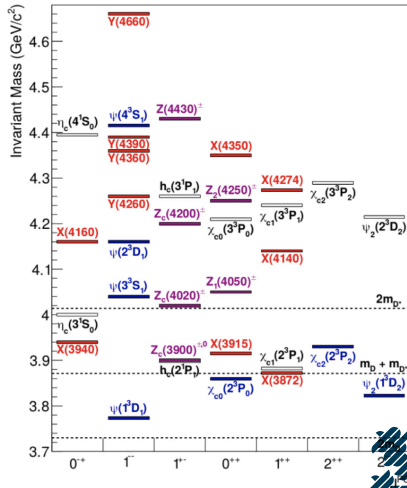
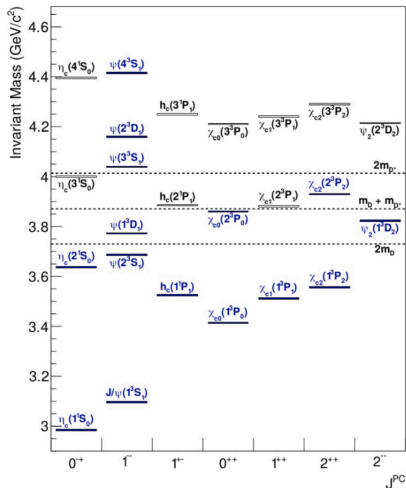


# Hybrid static potentials and gluelumps on $N_f = 3 + 1$ ensembles

R. Höllwieser\*, T. Korzec, F. Knechtli,  
M. Peardon, L. Struckmeier, J.-A. Urrea-Niño

Bergische Universität Wuppertal, Trinity College Dublin

# Motivation: charmonium spectrum



2003 and nowadays (above open charm threshold)

# Hybrid static potentials

- ▶ *Hybrid mesons* are states consisting of a quark-antiquark pair bound by an excited gluon field.
- ▶  $\Lambda_\eta^\epsilon$ , cylindrical symmetry group  $D_{\infty h} = D_\infty \times Z_2$  ( $D_{4h}$ )
- ▶ angular momentum about molecular axis:  
 $\Lambda = 0, 1, 2, 3, \dots \equiv \Sigma, \Pi, \Delta, \Phi, \dots$
- ▶  $\eta = +, - \equiv g, u$  under parity and charge conjugation  $\mathcal{P} \circ \mathcal{C}$
- ▶ for  $\Lambda = 0$  additional parity index  $\epsilon = +, -$
- ▶  $R \rightarrow \infty$ , multiplets associated with excitations of a *relativistic string*
- ▶  $R \rightarrow 0$ , multiplets associated with *gluelumps*, which are energy levels of QCD in the presence of a color-octet source ( $\mathbf{3} \otimes \mathbf{3}^* = \mathbf{1} \oplus \mathbf{8}$ )



# Gluelumps

- ▶ A system containing a valence gluon connected by an adjoint string to an adjoint source is called *gluelump*.
  - ▶ adjoint source  $\rightarrow$  infinitely heavy gluon  $\rightarrow$  *gluino*?
  - ▶ no physical object, not measured experimentally
- ▶ Why consider gluelumps?
  - ▶ gluelump masses define decay of field strength correlators in QCD vacuum
$$G_{gluelump}(x, y) = \langle F_{\mu\nu}^a(x) F_{\lambda\sigma}^b(y) \rangle$$
  - ▶ inverse mass of lowest gluelump equals the gluon correlation length  $T_g$
  - ▶ testing ground for models of low energy QCD (MIT bag model, flux tube model, pNRQCD  $\rightarrow$  BOEFT)
  - ▶ needed for getting insights in non-perturbative features of QCD (heavy quarkonia)

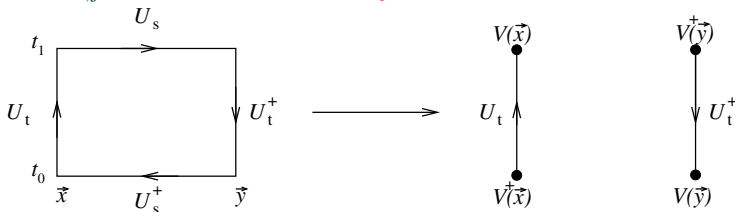


# Laplace trial states

- ▶  $N_v$  eigenvectors  $v(\vec{x})$  of the 3D covariant Laplace operator
- ▶ spatial Wilson line:  $U'_s(\vec{x}, \vec{y}) = G(\vec{x})U_s(\vec{x}, \vec{y})G^\dagger(\vec{y})$   
 $v'(\vec{x})v'^\dagger(\vec{y}) = G(\vec{x})v(\vec{x})v^\dagger(\vec{y})G^\dagger(\vec{y})$
- ▶ Wilson loop of size ( $R = |\vec{r}| = |\vec{x} - \vec{y}|$ )  $\times$  ( $T = |t_1 - t_0|$ )

$$\Sigma_g^+(R, T) = \langle \text{tr} [U_t(\vec{x}; t_0, t_1) U_s(\vec{x}, \vec{y}; t_1) U_t^\dagger(\vec{y}; t_0, t_1) U_s^\dagger(\vec{x}, \vec{y}; t_0)] \rangle$$

$$\rightarrow \langle \sum_{i,j}^{N_v} \text{tr} [U_t(\vec{x}) \rho_j(t_1) v_j(\vec{x}, t_1) v_j^\dagger(\vec{y}, t_1) U_t^\dagger(\vec{y}) \rho_i(t_0) v_i(\vec{y}, t_0) v_i^\dagger(\vec{x}, t_0)] \rangle$$



# Laplacian static-hybrid operators

- ▶ Gaussian profile functions:  $\rho_i^{(k)}(\lambda_i) = \exp(-\lambda_i^2/2\sigma_k^2)$
- ▶ we can realize gluonic excitations via covariant derivatives
 
$$\nabla_{\vec{k}} v(\vec{x}) = \frac{1}{2} [U_k(\vec{x})v(\vec{x} + \hat{k}) - U_k^\dagger(\vec{x} - \hat{k})v(\vec{x} - \hat{k})]$$
- ▶  $\Sigma_{u/g}^\mp(R, T) = \sum_{\vec{x}, t_0, i, j, \vec{k} | \vec{r} = \vec{y} - \vec{x}}$ 

$$\langle \text{tr} [U_t(\vec{x}; t_0, t_1) \rho(\lambda_j) \{ [\nabla_{\vec{k}} v_j](\vec{x}, t_1) v_j^\dagger(\vec{y}, t_1) \pm v_j(\vec{x}, t_1) [\nabla_{\vec{k}} v_j]^\dagger(\vec{y}, t_1) \}$$

$$U_t^\dagger(\vec{y}; t_0, t_1) \rho(\lambda_i) \{ [\nabla_{\vec{k}} v_i](\vec{y}, t_0) v_i^\dagger(\vec{x}, t_0) \pm v_i(\vec{y}, t_0) [\nabla_{\vec{k}} v_i]^\dagger(\vec{x}, t_0) \}] \rangle$$
- ▶  $\Pi_{u/g}(R, T) = \Pi_{\mp}(R, T) = \sum_{\vec{x}, t_0, i, j, \vec{k} \perp \vec{r} = \vec{y} - \vec{x}}$ 

$$\langle \text{tr} [U_t(\vec{x}; t_0, t_1) \rho(\lambda_j) \{ [\nabla_{\vec{k}} v_j](\vec{x}, t_1) v_j^\dagger(\vec{y}, t_1) \pm v_j(\vec{x}, t_1) [\nabla_{\vec{k}} v_j]^\dagger(\vec{y}, t_1) \}$$

$$U_t^\dagger(\vec{y}; t_0, t_1) \rho(\lambda_i) \{ [\nabla_{\vec{k}} v_i](\vec{y}, t_0) v_i^\dagger(\vec{x}, t_0) \pm v_i(\vec{y}, t_0) [\nabla_{\vec{k}} v_i]^\dagger(\vec{x}, t_0) \}] \rangle$$
- ▶ ...  $\nabla_i \nabla_j$



# Gluelump operators

- ▶ 35 different 3D loop shapes  $W_s(\vec{x}, t)$
- ▶ fix parity by considering the sum ( $P=+1$ ) or difference ( $P=-1$ ) of each loop with its parity twin
- ▶ fix charge conjugation by using  $W_s^{P,C=\pm} = W_s^P \pm W_s^{P\dagger}$
- ▶ elements of the cubic group can act on  $W_s^{PC}(\vec{x}, t)$ , resulting in 24 loops which form a basis which generates the regular (or permutation) representation of the group
- ▶ linear combinations of the 24 loops project to the five irreps  $\Lambda \in \{A1, A2, E, T1, T2\}$
- ▶ projection coefficients differ from those for glueballs, since symmetries have to hold locally



# Gluelump correlator

$$C_{\Lambda PC}(t_1 - t_0) =$$

$$\sum_{\vec{x}, t_0} \left\{ \text{tr} [W_{\Lambda PC}(\vec{x}, t_0) U_t(\vec{x}, t_0, t_1) W_{\Lambda PC}^\dagger(\vec{x}, t_1) U_t^\dagger(\vec{x}, t_0, t_1)] \right. \\ \left. - \text{tr} [W_{\Lambda PC}(\vec{x}, t_0)] \text{tr} [W_{\Lambda PC}^\dagger(\vec{x}, t_1)] / 3 \right\}$$

effective masses contain a self-energy contribution  $m_{self}(a)$  from adjoint color sources, which differs from the self-energy in the (hybrid) static potentials  $V_{self}(a)$ , in leading order PT

$$m_{self}(a) = \frac{C_A}{C_F} \frac{V_{self}(a)}{2} = \frac{N^2}{N^2 - 1} V_{self}(a) > V_{self}(a),$$

with  $C_A = N$  and  $C_F = (N^2 - 1)/2N$  the quadratic Casimir invariants of the adjoint and fund. representations of  $SU(N)$



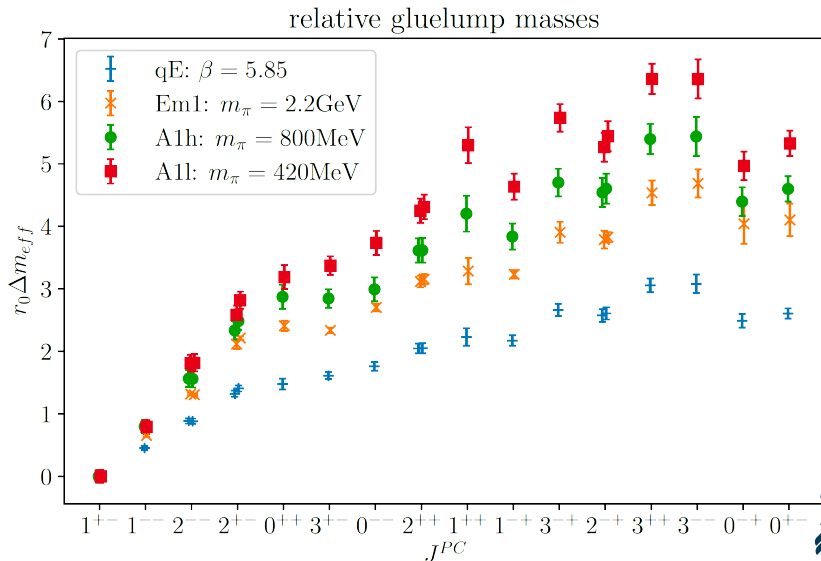


# Lattice ensembles

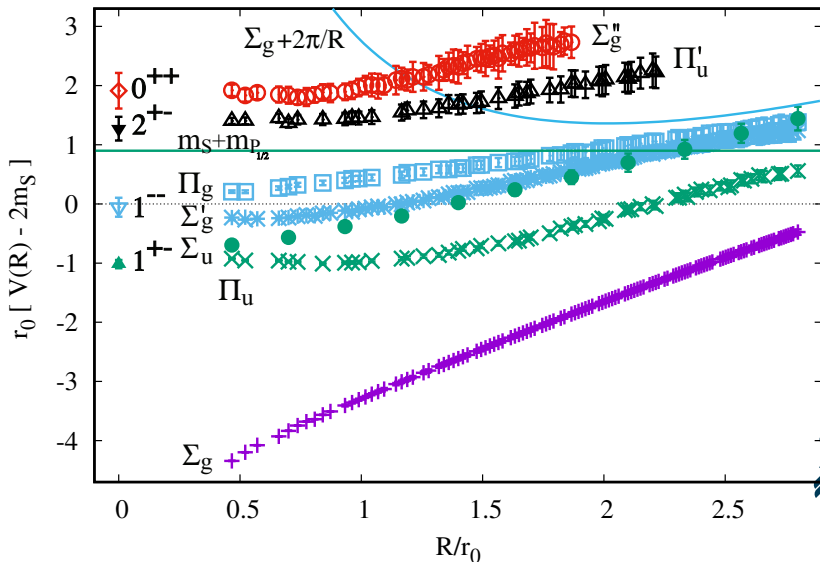
ens.	qE	Em1	A1h	A1	A2
volume	$48 \times 24^3$	$48 \times 24^3$	$96 \times 32^3$	$96 \times 32^3$	$128 \times 48^3$
$\beta$	5.85	5.3	3.24	3.24	3.24
$N_f$	-	0+2	3+1	3+1	3+1
$\kappa_l$	-	-	0.13392	0.134407	0.134396
$\kappa_c$	-	0.13270	0.12834	0.12784	0.12798
$m_\pi$	-	2.2(1) GeV	788(6) MeV	406(3) MeV	409(1) MeV
$r_0/a$	4.2612(42)	4.2866(24)	7.278(27)	9.023(63)	8.998(54)
$t_0/a^2$	-	1.8477(3)	5.078(14)	7.438(32)	7.434(20)
$a$ [fm]	0.0662(12)	0.0658(10)	0.0690(11)	0.05359(15)	0.05355(13)
$N_{cfg}$	6480	7098	4000	4000	2000
$N_{vec}^l$	-	-	200	100	100
$N_{per}^l$	-	-	2000	4000	2000
$N_{vec}^c$	-	200	200	200	-
$N_{per}^c$	-	1000	2000	4000	-

for  $N_f = 3 + 1$  ensembles we use open b.c., Lüscher-Weisz gauge action  
 Wilson quarks with a non-perturbatively determined clover coefficient in a  
 massive O(a) improvement scheme → *Fritzsch, 2018*

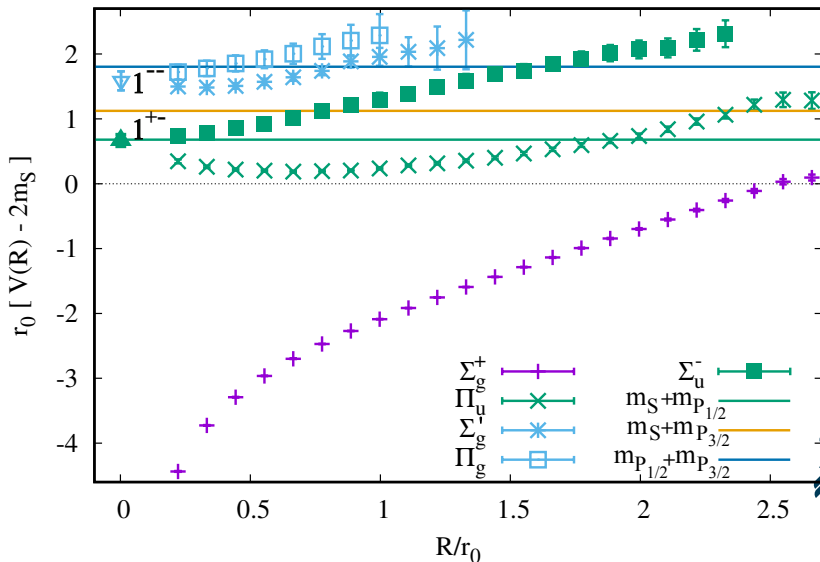




# Static-hybrid states, $N_f = 2$ , $m_\pi = 2.2\text{GeV}$



# Preliminary!!! $N_f = 3 + 1$ , $m_\pi = 420$ MeV



# Conclusions & Outlook

- ✓ hybrid static potentials (hybrid meson masses) from Laplacian trial states with covariant derivatives of eigenvectors as gluonic excitations
- ✓ effective gluelump masses from extended Wilson loop basis correlators ( $R \rightarrow 0$  limit)
- ✓ inputs for Born-Oppenheimer Effective Field Theories
- 🔧 BOEFT allows systematically improvable approximations to various properties of hadrons (XYZ)
- 🔧 more (hybrid) static(-light) and multi-quark potentials
- 🔧 hybrid string breaking study non-trivial (various avoided and broad level crossings)
- ? questions, discussion, many possible applications...

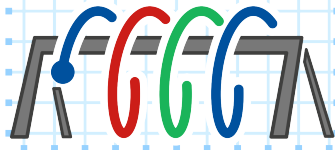


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## THANK YOU



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