

$X(3872)$ Relevant $D\bar{D}^*$ Scattering in $N_f = 2$ Lattice QCD ¹

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¹Based on: <https://arxiv.org/abs/2402.14541>

Outline

- 1 Motivation
- 2 Methodology: Lüscher's formulism
- 3 Results: bound state and $X(3872)$
- 4 Discussion: a possible resonance and $\chi_{c1}(2P)$
- 5 Summary and Outlook

Motivation

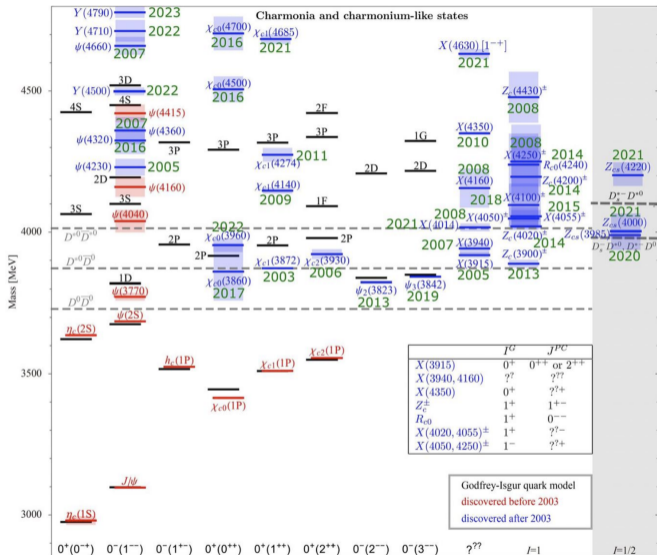
$X(3872)$ status

1 First observation from Belle. PDG:

- 1 $m_X = 3871.65 \pm 0.06 \text{ MeV}$
- 2 $\Gamma_X = 1.19 \pm 0.21 \text{ MeV}$
- 3 $I^G J^{PC} = 0^+ 1^{++}$

A lot of XYZ particles observed afterwards...

- 2 $X(3872)$ decays mainly to $D^0 \bar{D}^{0*}$, a small fraction to $J/\psi \omega$, and also isospin violating $J/\psi \rho$
- 3 Many intensive and extensive phenomenological studies
- 4 Interpreted as a $c\bar{c}$, a $D\bar{D}^*$ molecule or a tetraquark state
- 5 Main point of view: $D\bar{D}^* + c\bar{c}$



Finite volume scattering theory

$$\left\{ \begin{array}{l} \text{characteristic size: } R_a \ll L \\ \text{temporal size: } L_t \ll L \\ \text{scattering energy: } E_k \end{array} \right. \xrightarrow{\text{Lüscher's formulism}} \left\{ \begin{array}{l} \text{Scattering amplitude } \mathcal{M}, \text{ phase shift } \delta_l, \\ \text{ERE: } a_l, r_l, \\ \text{analytic pole position.} \end{array} \right.$$

Lüscher's formulism (QC)

single channel:^a

$$p \cot \delta_0(q^2) = \frac{2}{La_s \sqrt{\pi}} \mathcal{Z}_{00}(1, q^2). \quad (1)$$

multi-channel:^{b c}

$$\det \left[\delta_{ij} \delta_{JJ'} + i \rho_i t_{ij}^{(J)}(s) \left(\delta_{JJ'} + i \mathcal{M}_{JJ'}^{\bar{P}\Lambda}(p_i L) \right) \right] = 0. \quad (2)$$

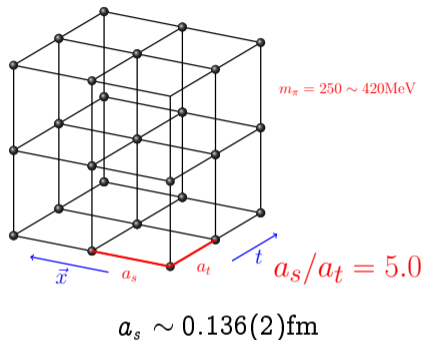
^a Lüscher *NPB*.354,2-3,(1991)531-578

^b L.Leskovec and S.Prelovsek *PRD*.(2015)85, 114507

^c J.J.Dudek et al.(Hadron Spectrum Collab.) *PRL*.(2014)113,182001

Lattice setup

- Tadpole improved Symanzik's **gauge action**
- Tadpole improved anisotropic **clover fermion action** for sea light and valence charm quarks.
(C. Morningstar, PRD60(1999)034509)
Scale setting: Jiang et al, Phys.Rev.D 107 (2023) 094510
- **Anisotropic** Lattice: $\xi = \frac{a_s}{a_t} \approx 5$.
- **Distillation method** (M. Peardon et al. (HSC), PRD80(2009)054506).
- Calculation of **disconnected contribution** of light dynamic quarks.

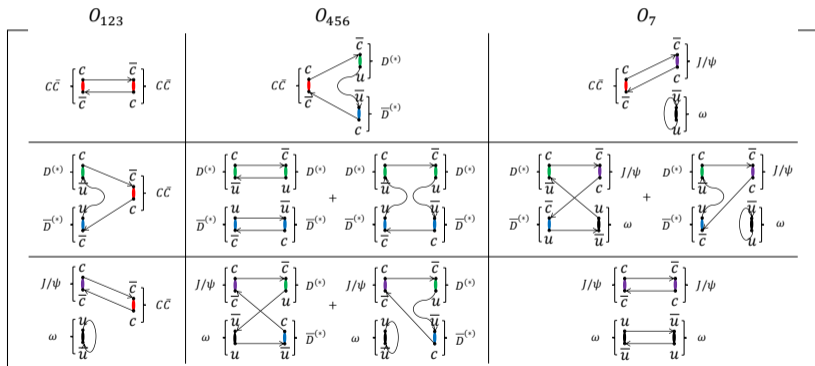


ens.	size	m_π (MeV)	$m_\pi L_s$	a_t^{-1} (GeV)	N_{cfg}	$N_V^{(1)}$	$N_V^{(c)}$	$m_{\chi_{c1}}$ (MeV)	m_D (MeV)	m_{D^*} (MeV)
M245	$16^3 \times 128$	250(3)	~ 2.7	7.276	401	70	120	3489(3)	1873(1)	1985(2)
M305	$16^3 \times 128$	307(2)	~ 3.4	7.187	401	70	120	3496(2)	1881(1)	1990(2)
M360	$16^3 \times 128$	362(1)	~ 4.0	7.187	401	70	120	3502(2)	1884(1)	2003(2)
M415	$16^3 \times 128$	417(1)	~ 4.7	7.219	401	70	120	3509(2)	1896(1)	2017(1)

Schematic calculation of $I^G J^{PC} = 0^+ 1^{++}$ system

flavor structure of $0^+ 1^{++} D \bar{D}^*$ operators

$$|D \bar{D}^*\rangle_{I=0}^{Q=0} = \frac{1}{2} (|D^+ \bar{D}^{*-}\rangle + |D^0 \bar{D}^{*0}\rangle - |\bar{D}^0 D^{*0}\rangle - |D^- D^{*+}\rangle) \quad (3)$$



$$\mathcal{S} = \{\mathcal{O}_\alpha | \alpha = 1, \dots, 7\} = \{\mathcal{O}_{c\bar{c}}^{r=0}, \mathcal{O}_{c\bar{c}}^{r=1}, \mathcal{O}_{c\bar{c}}^{r=2}, \mathcal{O}_{D\bar{D}}^{q=0, \gamma_5}, \mathcal{O}_{D\bar{D}}^{q=0, \gamma_4 \gamma_5}, \mathcal{O}_{D\bar{D}}^{q=1}, \mathcal{O}_{J/\psi \omega}^{q=0}\}.$$

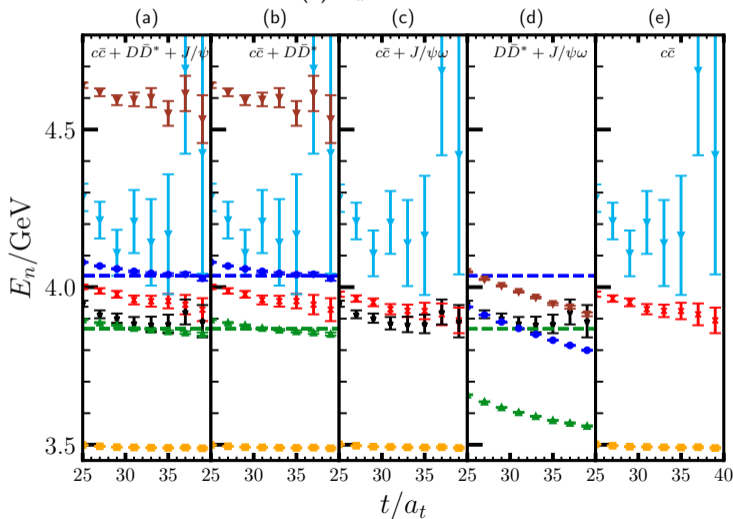
Energy levels of $I^G J^{PC} = 0^{+1^{++}}$ system

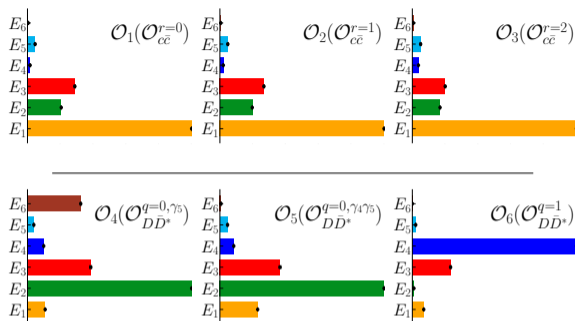
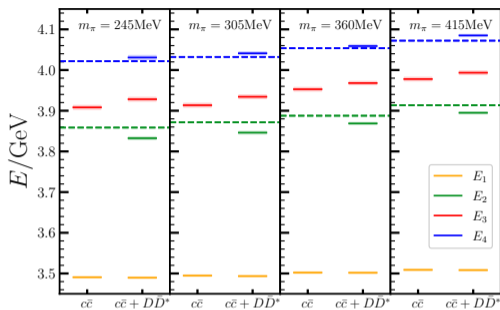
- $c\bar{c} + D\bar{D}^* + J/\psi\omega$ operators.
- $c\bar{c} + D\bar{D}^*$: black points (correspond to $J/\psi\omega$ state) disappear.
- $c\bar{c} + J/\psi\omega$: Energy levels close to non-interacting $D\bar{D}^*$ energies disappear.
- $D\bar{D}^* + J/\psi\omega$: Energy levels close to non-interacting $D\bar{D}^*$ and $J/\psi\omega$ states.
- $c\bar{c}$: Energy levels close to χ_{c1} states.

In all the cases, $J/\psi\omega$ energy has no sizable changes w/o $c\bar{c}$ and $D\bar{D}^*$ operators. So $J/\psi\omega$ almost decouples from other states and is **neglected** from the discussion in this work.

In agreement with [Prelovsek et. al., PRL\(2013\).111.192001](#).

(1) $m_\pi = 250$ MeV





Identify the energy levels:

- E_1 : ~ 3.5 GeV, should be χ_{c1} .
- E_2 : close but **below** the $D\bar{D}^*$ threshold.
- E_3 : far from and in middle of the non-interacting DD^* energies $E_{D\bar{D}^*}^{q=0}$ and $E_{D\bar{D}^*}^{q=1}$
- E_4 : close but **above** $E_{D\bar{D}^*}^{q=1}$

Operator couplings

- E_1 : coupled most by $c\bar{c}$ operators.
- E_2 : coupled most by $\mathcal{O}_{D\bar{D}^*}^{q=0}$ and substantially by $c\bar{c}$ operators.
- E_3 : coupled substantially by $\mathcal{O}_{D\bar{D}^*}^{q=0}$, $\mathcal{O}_{D\bar{D}^*}^{q=1}$ and $c\bar{c}$ operators.
- E_4 : coupled most by $\mathcal{O}_{D\bar{D}^*}^{q=1}$ and a little by $\mathcal{O}_{D\bar{D}^*}^{q=0}$ and $c\bar{c}$ operators.

Existence of a bound state below $D\bar{D}^*$ threshold

- **Lüscher's formalism** for S -wave single channel

$$p \cot \delta_0(q^2) = \frac{2}{La_s\sqrt{\pi}} \mathcal{Z}_{00}(1, q^2).$$

- $E_n(p_n) = \sqrt{m_D^2 + p_n^2} + \sqrt{m_{D^*}^2 + p_n^2}$,
 $q^2 \equiv \left(\frac{L}{2\pi}\right)^2 p^2$

Effective Eange Expansion(ERE)

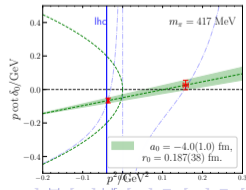
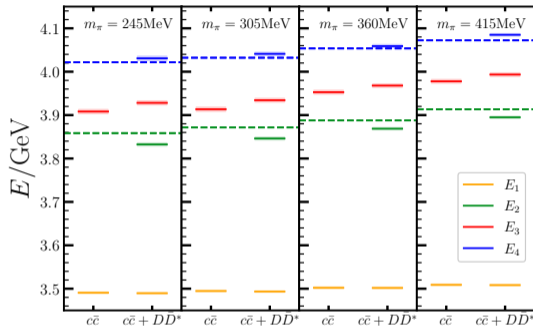
$$p \cot \delta_0(p) = \frac{1}{a_0} + \frac{1}{2}r_0 p^2 + \mathcal{O}(p^4).$$

Solving ERE with E_2 and E_3 , we can obtain the parameters (a_0, r_0)

pole singularity in \mathcal{T} matrix

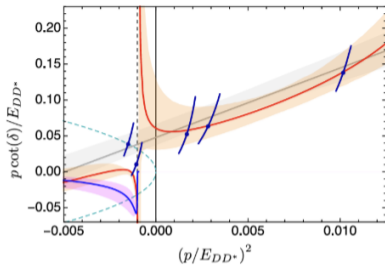
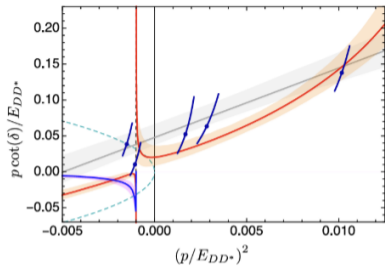
$$\mathcal{T} \propto (p \cot \delta_0(p) - ip)^{-1}.$$

- taking (a_0, r_0) as the approximation in the $V \rightarrow \infty$ limit, the pole equation $p_B \cot \delta_0(p_B) - ip_B = 0$ gives the banding energy $E_B = E_{D\bar{D}^*}(p_B) - (m_D + m_{D^*})$.

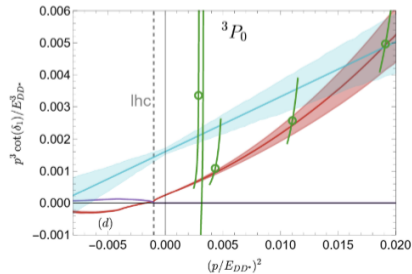
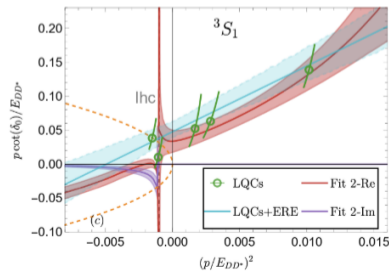


Left-hand cut issue: example of $T_{cc}^+(3875)$

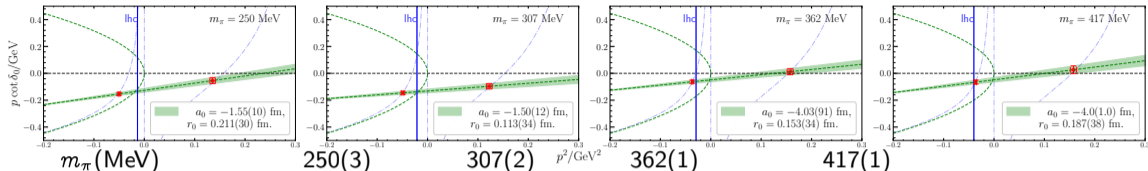
■ M.-L. Du et al., PRL131(2023)131901
 Case studies on $T_{cc}^+(3875)$ relevant $D\bar{D}^*$ scattering. The data points are from lattice QCD calculation (M. Padmanath et al. PRL129(2022)032002)



■ L. Meng et al., PRD109(2024)L071506
 The lattice finite volume energy levels are used to fix the parameters in the EFT involved. The prediction of the EFT (red curves) are compared with ERE with out OPE.



Existence of a bound state below $D\bar{D}^*$ threshold



m_π (MeV)	250(3)	307(2)	362(1)	417(1)
$(p_{\text{lhc}}^{1\pi})^2$ (GeV^2)	-0.0135(4)	-0.0210(4)	-0.0292(3)	-0.0400(3) (★)
a_0 (fm)	-1.55(10)	-1.50(12)	-4.03(91)	-4.0(1.0) (★)
r_0 (fm)	0.211(30)	0.113(34)	0.153(34)	0.187(38) (★)
E_B (MeV)	$-9.7_{-2.2}^{+2.1}$ (*)	$-9.7_{-2.0}^{+1.9}$ (*)	$-1.3_{-0.8}^{+0.6}$ (*)	$-1.3_{-1.0}^{+0.8}$ (★)

■ For $m_\pi \lesssim 360$ MeV, our results suffering from the LHC issue. (M.-L. Du et al., PRL131(2023)131901,

$$(p_{\text{lhc}}^{1\pi})^2 \approx \frac{1}{4} [(\Delta M)^2 - m_\pi^2]$$

■ E_2 : The lattice energy is lower than the $D\bar{D}^*$ threshold by 20MeV or even more.

■ a_0 : Large negative: existence of a bound state.

■ r_0 : Small positive: the compositeness $X \sim 1$ up to a $\mathcal{O}(p^2)$ correction (Y. Li et al., PRD105(2022)L071502) .

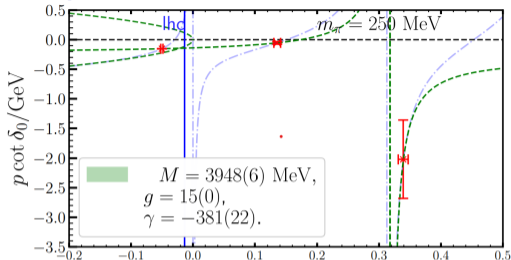
■ The bound state is predominantly a $D\bar{D}^*$ molecule.

Preliminary: Pole position of K -matrix parameterization

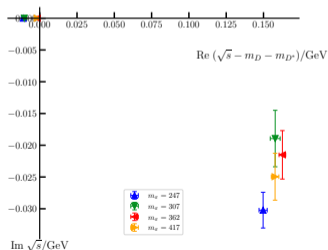
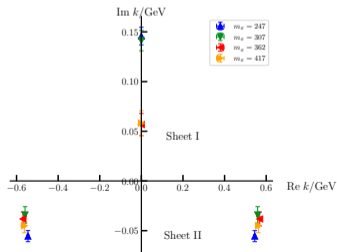
$$\mathcal{T}(s) = \frac{K(s)}{1 - K(s)i\rho(s)} = \frac{8\pi E}{p \cot \delta - ip}.$$

$$\text{Im}\mathcal{T}^{-1}(s) = -\rho(s) = -\frac{p}{8\pi E} \Theta(E - E_{\text{thr.}}),$$

$$\text{Re}\mathcal{T}^{-1}(s) = K^{-1}(s) = \frac{p}{8\pi E} \cot \delta, .$$



poles in k plane

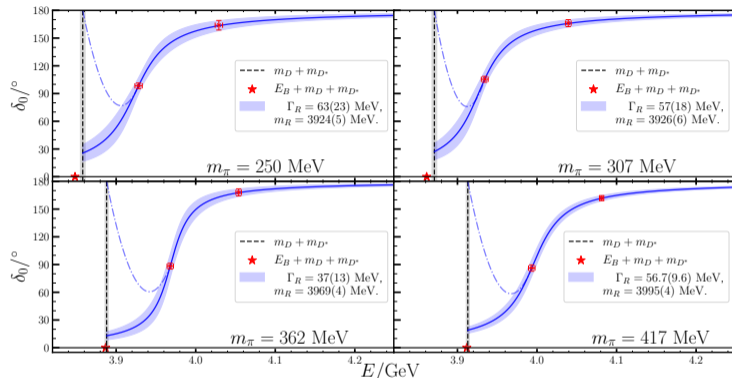


poles in s plane

K -matrix parameterization,

$$K = \frac{g^2}{s - M^2} + \gamma.$$

A possible resonance below 4.0 GeV



- E_3 : a scattering phase around $\delta(E_3) \sim 90^\circ$
- E_4 : a scattering phase close to $\delta(E_4) \sim 180^\circ$
- Exactly as the expectation of the **Generalized Levinson's theorem**
- **Hint at the existence of a resonance**

Generalized Levinson's theorem

$$\delta_l(0^+) - \delta_l(p_{\max}) = (n_l - n_b)\pi$$

- n_l : # of bound states,
- n_b : # of bare states below the energy corresponding to p_{\max} . (F. Vidal et al., PRC45(1992)418, Y. Li et al., PRD105(2022)116024)

A possible resonance below 4.0 GeV

- **Breit-Wigner ansatz** for a resonance :

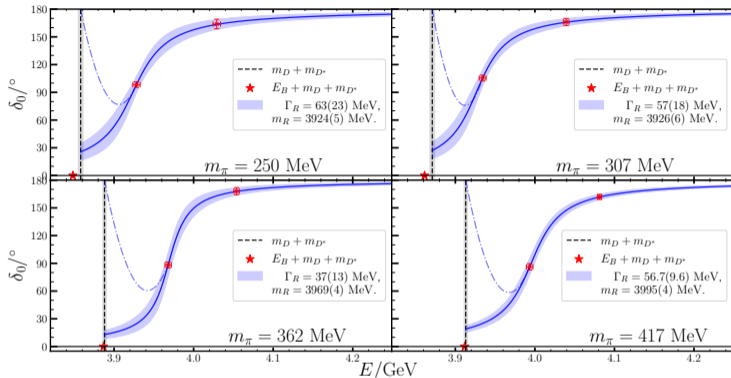
$$\mathcal{T} \approx \frac{1}{\cot \delta_0 - i} \sim \frac{1}{(m_R - E) - \frac{i\Gamma_R}{2}}$$

- Resonance parameters through

$$\delta_0(E) = \arctan\left(\frac{\Gamma_R}{2(m_R - E)}\right)$$

by using E_3 and E_4 .

- **Caution:** The parameters (m_R, Γ_R) may change , since they are only two energy levels.



potential experimental observation

X(3940): $m_X = 3942(9)\text{MeV}$

$\Gamma_X = 37_{-17}^{+27} \text{ MeV}$

(Belle, PRL98(2007)082001; PRL100(2008)20200)

m_π (MeV)	250(3)	307(2)	362(1)	417(1)
m_R (MeV)	3924(5)	3926(6)	3969(4)	3995(4)
Γ_R (MeV)	63(23)	57(18)	37(13)	57(10)

A possible resonance below 4.0 GeV and $\chi_{c1}(2P)$

Related phenomenological studies

Ref.	E. Cincioglu et al., EPJC76(2016)576	F. Giacosa et al., IJMPA34(2019)1950173	Q. Deng et al., 2312.10296	G.J. Wang et al., 2306.12406
m_R (MeV)	3910 – 3925	3995	3990	3958
Γ_R (MeV)	5 – 70	72	~ 60	~ 17

potential experimental observation (Belle)

$$X(3940): m_X = 3942(9)\text{MeV}$$

$$\Gamma_X = 37_{-17}^{+27} \text{ MeV}$$

(Belle, PRL98(2007)082001; PRL100(2008)20200)

New resonant structures in

$B^+ \rightarrow D^{*+} D^- K^+$, $B^+ \rightarrow D^* - D^+ K^+$ decays from LHCb :

	$\eta_c(3945)$	$h_c(4000)$	$\chi_{c1}(4010)$	and $h_c(4300)$
Consistent with X(3940) seen by Belle	10σ 0^{-+}	9.1σ 1^{+-}	16σ 1^{++}	6.4σ 1^{+-}

See Plenary talk in Mon. 10:00 AM. and arXiv:[2406.03156]

LHCb (arXiv: 2406.03156):

$$\chi_{c1}(4010) \quad J^{PC} = 1^{++}$$

$$m_0 = 4012.5_{-3.9}^{+3.6+4.1} \quad \Gamma_0 = 62.7_{-6.4}^{+7.0+6.4}$$

m_π (MeV)	250(3)	307(2)	362(1)	417(1)
m_R (MeV)	3924(5)	3926(6)	3969(4)	3995(4)
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Summary & Outlook

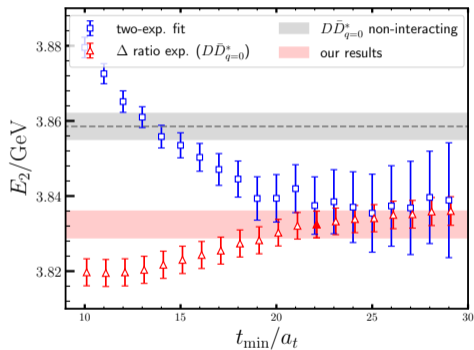
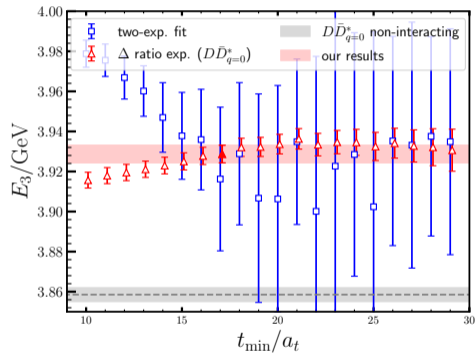
Summary

- 1 The pion mass m_π dependence (from 250 MeV to 417 MeV) of the S -wave $D\bar{D}^*(I=0)$ scattering process was studied in $N_f = 2$ lattice QCD.
- 2 A state below the $D\bar{D}^*$ threshold was observed at each pion mass.
- 3 Particularly for $m_\pi = 417$ MeV, this state is considered as a definite **bound state** in the infinite volume limit, located about 1 MeV below the $D\bar{D}^*$ threshold. This bound state is primarily a $D\bar{D}^*$ **molecular component**, possibly corresponding to the physical $X(3872)$ state.
- 4 The possibility of a **resonance** existing below 4.0 GeV was proposed for the first time from lattice QCD perspective.

Thank You

Backup slides

Backup for FV energy level determination



Backup for GEVP mass plateau

