Bringing near-physical QCD+QED calculations beyond the electro-quenched approximation

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Motivation

 $FLAG¹$ reports averages for observables calculable from $K \to \ell \overline{\nu}$, $\pi \to \ell \overline{\nu}$ at a sub-percent level.

\n- $$
N_f = 2 + 1
$$
 $f_{\pi^{\pm}} = 130.2(0.8)$ MeV (0.61%)
\n- $N_f = 2 + 1 + 1$ $f_{K^{\pm}} = 155.7(0.3)$ MeV (0.19%)
\n- $N_f = 2 + 1$ $f_{K^{\pm}} = 155.7(0.7)$ MeV (0.45%)
\n

• These inform $|V_{\mu s}|/|V_{\mu d}|$.

- **o** Lattice results based on partial evaluation of first-order isospin-breaking corrections (or χ PT).
- \bullet < 1% errors without a full ab-initio correction?

Plot from FLAG Review 2021 (February 2024 Revision). Full citation list at end of talk.

1 FLAG Review 2021 (February 2024 Revision),<http://flag.unibe.ch/2021/>

Motivation

- Similar situation for $D \to \ell \overline{\nu}$, $D_s \to \ell \overline{\nu}$.
	- $N_f = 2 + 1 + 1$ $f_D = 212.0(0.7)$ MeV (0.33%)
	- $N_f = 2 + 1$ $f_D = 209.0(2.4)$ MeV (1.15%)
	- $N_f = 2 + 1 + 1$ $f_{D_s} = 249.9(0.5)$ MeV (0.2%)
	- $N_f = 2 + 1$ $f_D = 248.0(1.6)$ MeV (0.65%)

o Important to include a complete ab-initio calculation of first-order isospin-breaking corrections.

Plot from FLAG Review 2021 (February 2024 Revision). Full citation list at end of talk.

Several published results addressing disconnected QED IB:

Electromagnetic Splittings and Light Quark Masses in Lattice QCD

Duncan et al. [1996 Phys. Rev. Lett. 76, 3894](https://doi.org/10.1103/PhysRevLett.76.3894)

- Computing electromagnetic effects in fully unquenched QCD Duncan et al. [2005 Phys. Rev. D 71, 094509](https://doi.org/10.1103/PhysRevD.71.094509)
- Full QED+QCD Low-Energy Constants through Reweighting Ishikawa et al. [2012 Phys. Rev. Lett. 109, 072002](https://doi.org/10.1103/PhysRevLett.109.072002)
- \bullet 1+1+1 flavor QCD+QED simulation at the physical point Aoki et al. [2012 Phys. Rev. D 86, 034507](https://doi.org/10.1103/PhysRevD.86.034507)
- Isospin splittings of meson and baryon masses from three-flavor lattice $QCD + QED$ Horsley et al. [2016 J. Phys. G: Nucl. Part. Phys. 43 10LT02](https://doi.org/10.1088/0954-3899/43/10/10LT02)

Quark-disconnected $O(\alpha)$ diagrams have also been calculated for g_{μ} – 2 HVP:

- [Leading hadronic contribution to the muon magnetic moment](https://doi.org/10.1038/s41586-021-03418-1) [from lattice QCD](https://doi.org/10.1038/s41586-021-03418-1) Borsanyi et al. [2021 Nature 593, 51–55](https://doi.org/10.1038/s41586-021-03418-1)
- [High precision calculation of the hadronic vacuum polarisation](https://doi.org/10.48550/arXiv.2407.10913) [contribution to the muon anomaly](https://doi.org/10.48550/arXiv.2407.10913) Boccaletti et al. [2024 \[arXiv:2407.10913\]](https://doi.org/10.48550/arXiv.2407.10913)

Lattice Strategy: RM123 method

• Working at $\mathcal{O}(\alpha)$: $m_{\mu} = m_{d}$.

- Introduce IB effects using the RM123 method 12 :
	- IB corrections via perturbative expansion in $\alpha = \frac{e^2}{4\pi}$ $rac{e}{4\pi}$, m .

$$
\langle O \rangle = \langle O \rangle \Big|_{e=0} + \underbrace{\frac{1}{2} (e^{\phi})^2 \left[\frac{\partial}{\partial e} \frac{\partial}{\partial e} \langle O \rangle \right]_{e=0}}_{QED IB} + \underbrace{(m^{\phi} - m^{(0)}) \left[\frac{\partial}{\partial m} \langle O \rangle \right]_{e=0}}_{Strong IB} + \dots (1)
$$

- IB corrections take the form of additional diagrams evaluated in the isospin-symmetric limit.
- $m^{\phi} =$ physical mass, $m^{(0)} =$ simulation-point mass.

 $^{\text{1}}$ [de](https://doi.org/10.1007/JHEP04%282012%29124) Divitiis *et al.* JHEP 04 (2012) 124 [\[arXiv:1110.6294\]](https://arxiv.org/abs/1110.6294)

² de Divitiis et al. [PRD 87 \(2013\) 114505](https://doi.org/10.1103/PhysRevD.87.114505) [\[arXiv:1303.4896\]](https://arxiv.org/abs/1303.4896)

Lattice Strategy: Isospin-Breaking Corrections to $P \to \ell \overline{\nu}$

Diagrams from Boyle et al. [JHEP 02 \(2023\) 242.](https://doi.org/10.1007/JHEP02(2023)242) Red diamonds: scalar insertions.

- Quark-connected contributions (left) to the isospin-breaking correction have been calculated for $P\to\ell\overline{\nu}$ in lattice QCD¹⁻².
- Quark-disconnected contributions (right) omitted.
- Referred to as the "electro-quenched" approximation.
- Uncontrolled systematic.

¹ Di Carlo et al. [PRD 100 \(2019\) 034514](https://doi.org/10.1103/PhysRevD.100.034514) [\[arXiv:1904.08731\]](https://arxiv.org/abs/1904.08731)

² EP 02 (2023) 242 [\[arXiv:2211.12865\]](https://arxiv.org/abs/2211.12865)

Lattice Strategy: Propagator Loops in Lattice QCD

- Quark-disconnected diagrams are difficult to estimate—loops given by factors like $D^{-1} (x,x).$
- This requires one propagator solve per lattice site. \rightarrow Computationally infeasible.
- Instead stochastically estimate Dirac operator inverse using noise vectors η obeying

$$
\langle \eta(y)\eta^{\dagger}(x)\rangle_{\eta} = \delta_{xy}, |\eta(x)|^2 = 1, \langle \eta(x)\rangle_{\eta} = 0, \qquad (2)
$$

where $\langle \cdot \rangle_n$ is an average over η . This gives

$$
D^{-1}(x, x) = \sum_{y} D^{-1}(x, y) \delta_{xy}
$$
\n
$$
\approx \frac{1}{N_{\eta}} \sum_{\eta} \left(\sum_{y} D^{-1}(x, y) \eta(y) \right) \eta^{\dagger}(x). \tag{3}
$$

Lattice Strategy: EM Currents in the Isospin Limit

 \bullet $\mathcal{O}(\alpha)$ correlation function for operator O:

$$
\sum_{x}\sum_{y}\langle J_{\mu}(x)A_{\mu}(x)J_{\nu}(y)A_{\nu}(y)O\rangle \tag{5}
$$

- EM current insertions: $J_{\mu}(x) = \sum_{f} Q_{f} \psi_{f}(x) \gamma_{\mu} \psi_{f}(x)$.
	- 2 + 1f: Consider sum over quark flavours $f \in \{u, d, s\}$.
	- Q_f : EM charge (*i.e.* $Q_u = 2/3$, $Q_d = -1/3$, $Q_s = -1/3$).
- \bullet u and d terms sum in single-propagator loops.
	- Light and strange quarks equally-weighted; relative minus sign.
	- \Rightarrow $J_\mu(x) = 1/3 \left(\overline{\psi}_I(x) \gamma_\mu \psi_I(x) \overline{\psi}_s(x) \gamma_\mu \psi_s(x)\right) A_\mu.$
	- This leads to differences of single-propagator traces in several disconnected diagrams.

Lattice Strategy: Split-Even Estimator

Giusti et $al¹$ have demonstrated a successful variance-reduction strategy for differences of single-propagator loops: "split-even" estimators.

For e.g. Wilson, DWF Dirac Operators differing only by mass,

$$
D_1^{-1} - D_2^{-1} = D_1^{-1} (D_2 - D_1) D_2^{-1}, \tag{6}
$$

$$
= (m_2 - m_1)D_1^{-1}D_2^{-1}.
$$
 (7)

Choice in how to stochastically estimate propagator traces:

"Standard"
$$
(m_2 - m_1) \text{Tr} \left\{ \gamma^{\mu} \left\{ D_1^{-1} D_2^{-1} \eta \right\} (\mathbf{x}) \eta^{\dagger} (\mathbf{x}) \right\},
$$
 (8)

"Split-Even"
$$
(m_2 - m_1) \text{Tr} \left\{ \gamma^{\mu} \left\{ D_1^{-1} \eta \right\} (x) \{ \eta^{\dagger} D_2^{-1} \}(x) \right\},
$$
 (9)

 \rightarrow c.f. Raoul Hodgson 11:35 2nd August – Use in rare K decays

1 [EPJC 79, 586 \(2019\)](https://doi.org/10.1140/epjc/s10052-019-7049-0) [\[arXiv:1903.10447\]](https://arxiv.org/abs/1903.10447)

Lattice Strategy: Ensemble Parameters

Current run performed on the RBC-UKQCD 'C0' ensemble.

- $2+1$ flavour, $L^{3} \times T = 48^{3} \times 96$, $a^{-1} = 1.73$ GeV.
- Physical-scale light-, strange-quark masses.
- zMöbius Domain-Wall action.
	- \rightarrow Cheaper than Möbius DWF; requires bias correction step.
	- \rightarrow Accumulate statistics on cheaper zMöbius estimator.
- Light quarks deflated with 2000 low modes.
- Following techniques developed on non-physical mass 'C1' ensemble for quark-disconnected diagrams (Harris *et al.* $^{1})$

Runs also planned on the 'M0' ensemble.

- $2{+}1$ flavour, $L^3\times\mathcal{T}=64^3\times128$, $a^{-1}=2.36$ GeV.
- Same physical volume as C0.
- Also at physical-scale light-, strange-quark masses.

1 TICE2022 (2023) 013 [\[arXiv:2301.03995\]](https://arxiv.org/abs/2301.03995)

Lattice Strategy: Photon Action

- \bullet Finite volume $+$ periodic boundary conditions:
	- \rightarrow Charged states forbidden by Gauss' Law.
- Need to choose a QED prescription.
	- QED_L : Remove spatial zero-mode¹.
		- \rightarrow Can express as a special case of QED $_L^{\rm IR}$ 2
		- \rightarrow Large finite-volume effects at $\mathcal{O}(1/L^3)$ for $\mathcal{K} \rightarrow \ell \overline{\nu} ?^3$
	- QED_r : Redistribute zero-mode to neighbouring modes⁴⁵.
		- \rightarrow Investigated to remove $\mathcal{O}(1/L^3)$ finite-volume effects
		- \rightarrow Also a particular case of QED $_L^{\rm IR}$
		- \rightarrow Used for this project.

¹ [Hayakawa and Uno, PTP 120 \(2008\) 413](https://doi.org/10.1143/PTP.120.413) [\[arXiv:0804.2044\]](https://arxiv.org/abs/0804.2044)

² Davoudi et al. [PRD 99 \(2019\) 034510](https://doi.org/10.1103/PhysRevD.99.034510) [\[arXiv:1810.05923\]](https://arxiv.org/abs/1810.05923)

³
Boyle *et al.* [JHEP02\(2023\)242](https://doi.org/10.1007/JHEP02%282023%29242) [arXiv: \[2211.12865\]](https://arxiv.org/abs/2211.12865)

⁴ [Di Carlo, PoS LATTICE2023 \(2024\) 120](https://doi.org/10.22323/1.453.0120) [\[arXiv:2401.07666\]](https://arxiv.org/abs/2401.07666)

⁵ Hermansson-Truedsson et al.[, PoS LATTICE2023 \(2024\) 265](https://doi.org/10.22323/1.453.0265) [\[arXiv:2310.13358\]](https://arxiv.org/abs/2310.13358)

Lattice Strategy: Software

- Calculation performed with Grid^1 , and the Grid-based workflow management software $\sf Hadrons^2.$
- Split-even and quark-disconnected diagram contractions implemented as Hadrons modules (code review TBC).
- Thanks to Antonin Portelli, Raoul Hodgson, and Tim Harris for assisting with code development.

1 <https://github.com/paboyle/Grid>

2 github.com/aportelli/Hadrons

$\mathcal{O}(\alpha)$ Quark-Disconnected Diagrams for $P \to P$

[Schematics courtesy of Matteo Di Carlo.]

Specs error analysis

- Specs subdiagram contains two $l - s$ loops.
- Can be computed as a difference of propagators or with the split-even estimator.

• Left: error scaling of specs subdiagram.

- Significant reduction in error with the split-even estimator.
- Reaching the gauge noise with ~32 72 noise hits.

Tadpole error analysis

- Tadpole diagrams also feature an $l - s$ loop.
- **.** Left: Relative error of the kaon tadpole diagrams.
- Once again, the split-even estimator significantly reduces the error.

Burger error analysis

- \bullet 'Burger' diagram \rightarrow same flavour in both propagators.
- $e_q^2 \rightarrow$ no relative cancellation between diagrams.

- 'Burger' diagram falls off exponentially with propagator separation \Rightarrow short-distance dominated.
- Prior study on non-physical mass 'C1' ensemble (Harris et $al.$ ¹): concentrate computational effort on short-distance behaviour.
	- Volume-averaged stochastic estimation of all-to-all propagators within a radius $|x - y| < R$.
	- Random point sources for $|x y| \ge R$.

¹ ICF2022 (2023) 013 [\[arXiv:2301.03995\]](https://arxiv.org/abs/2301.03995)

- Right: Error scaling of the strange burger subdiagram $(R=4)$.
- Volume-averaging strategy provides an efficient method for obtaining a small error.
- Not yet at gauge noise with 64 hits.

- Quark-disconnected QED IB corrections are challenging, but important to quantify.
- Diagrams at $\mathcal{O}(\alpha)$ have exploitable characteristics:
	- Precision of 'Specs' and Tadpole diagrams can be greatly improved with the split-even estimator.
	- 'Burger' diagram is short-distance dominated.
- Findings at non-physical masses reproduced on a physical-point ensemble.
- Building towards a physical-point electro-unquenched calculation of $P \to \ell \nu$.

Backup

Dowdall et al. PRD 88 (2013) 074504 Carrasco et al. PRD 91 (2015) 054507 Bazavov et al. PRD 98 (2018) 074512 Miller et al. PRD 102 (2020) 034507 Alexandrou et al. PRD 104 (2021) 074520 Follana et al. PRL 100 (2008) 062002 Bazavov et al. PoS LATTICE2010 (2010) 074 Durr et al. PRD 81 (2010) 054507 Blum et al. PRD 93 (2016) 074505 Durr et al. PRD 95 (2017) 054513 Bornyakov et al. PLB 767 (2017) 366–373 Blossier et al. JHEP 07 (2009) 043

Balasubramamian, Blossier EPJC 80 (2020) 5, 412 Carrasco et al. JHEP 03 (2014) 016 Davies et al. PRD 82 (2010) 114504 Bazavov et al. PRD 85 (2012) 114506 Boyle et al. JHEP 12 (2017) 008 Yang et al. PRD 92 (2015) 034517 Na et al. PRD 86 (2012) 054510 Bazavov et al. PRD 98 (2018) 074512 Carrasco et al. PRD (2015) 054507