

# Reconstruction of the vector meson propagator using a generalized eigenvalue problem

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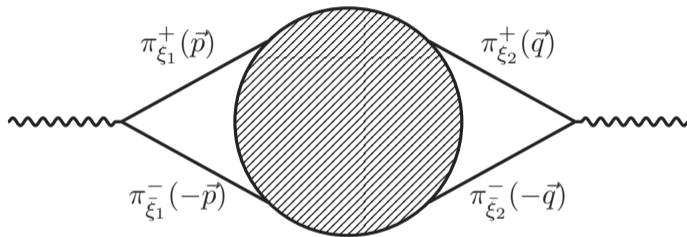
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# Introduction

- HVP contribution to  $g_\mu - 2$  given by  $\rho$



⇒ Use GEVP to extract  $\pi\pi$ -contributions and reconstruct the long-distance behavior **[1910.01083, 2406.19193, 2406.19194]**

# The staggered symmetry group

- Clebsch-Gordan coefficients of the staggered symmetry group

$$\mathbb{Z}_N^3 \rtimes [\Gamma^{4,1} \rtimes W_3]$$

- Two-pion state with the same quantum numbers as the  $\rho$ -meson

$$|s, m, \alpha\rangle_\rho = \sum_{(\vec{p}^a, \xi_\mu^a, \vec{p}^b, \xi_\mu^b) \in s} |\vec{p}^a, \vec{\xi}^a, \xi_4^a; \vec{p}^b, \vec{\xi}^b, \xi_4^b\rangle \langle \vec{p}^a, \vec{\xi}^a, \xi_4^a; \vec{p}^b, \vec{\xi}^b, \xi_4^b | s, m, \alpha\rangle_\rho$$

Derivation in **[2401.00514, 2112.11647]**

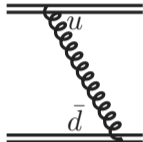
# Propagators

$$\pi^\dagger(\vec{x}, 0) \quad \bar{d} \quad \pi(\vec{z}, t)$$



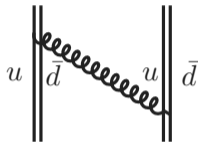
$$\pi(\vec{y}, 0) \quad \bar{d} \quad \pi^\dagger(\vec{w}, t)$$

$$\pi^\dagger(\vec{x}, 0) \quad \bar{d} \quad \pi(\vec{z}, t)$$

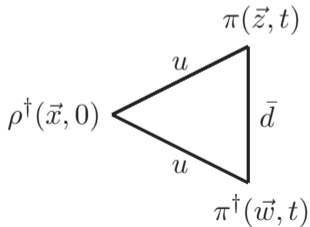


$$\pi(\vec{y}, 0) \quad u \quad \pi^\dagger(\vec{w}, t)$$

$$\pi^\dagger(\vec{x}, 0) \quad \pi(\vec{z}, t)$$



$$\pi(\vec{y}, 0) \quad \pi^\dagger(\vec{w}, t)$$



$$\rho^\dagger(\vec{x}, 0) \quad \begin{array}{c} u \\ \hline \hline \\ \bar{d} \end{array} \quad \rho(\vec{z}, t)$$

# Implementation

- Random wall sources  $\xi$  at  $t_0$  and additional low-mode averaging for  $\rho$
- Sequential inversion for the connected diagramm (square)

$$\langle M^{-1}O_t^B M^{-1}O_t^B | M^{-1}O_0^A M^{-1}O_0^A \rangle$$

- Number of inversion:

$$T \times \#O \times \#\xi \quad \text{times}$$

# Implementation (LMA for connected diagramm)

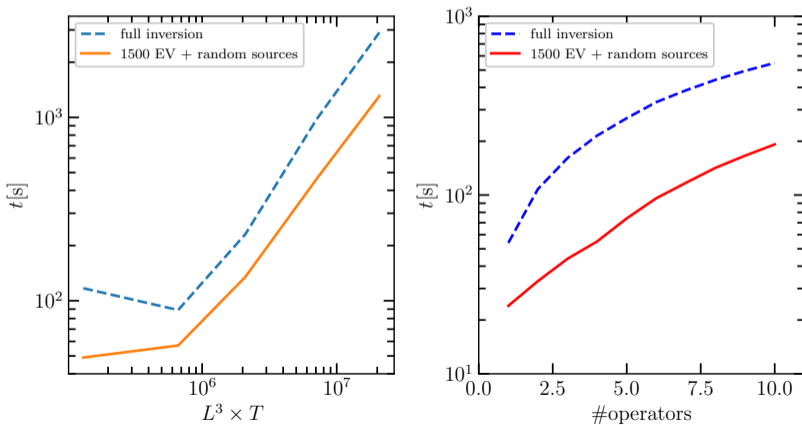
$$M^{-1} = \sum_{i=1}^{N_{\text{eig}}} \frac{1}{\lambda_i} |v_i\rangle \langle v_i| + M_{\text{rest}}^{-1}$$

Get rid of  $T$ -factor in the total number of inversions

$$\begin{aligned} \langle M^{-1} O_t^B M^{-1} O_t^B | M^{-1} O_0^A M^{-1} O_0^A \rangle &\rightarrow \sum_i \frac{1}{\lambda_i} \langle M^{-1} O_t^B | v_i \rangle \langle v_i | O_t^B | M^{-1} O_0^A M^{-1} O_0^A \rangle \\ &+ \sum_{\sigma} \langle M^{-1} O_t^B | \sigma \rangle \langle \sigma | M_r^{-1} O_t^B | M^{-1} O_0^A M^{-1} O_0^A \rangle \end{aligned}$$

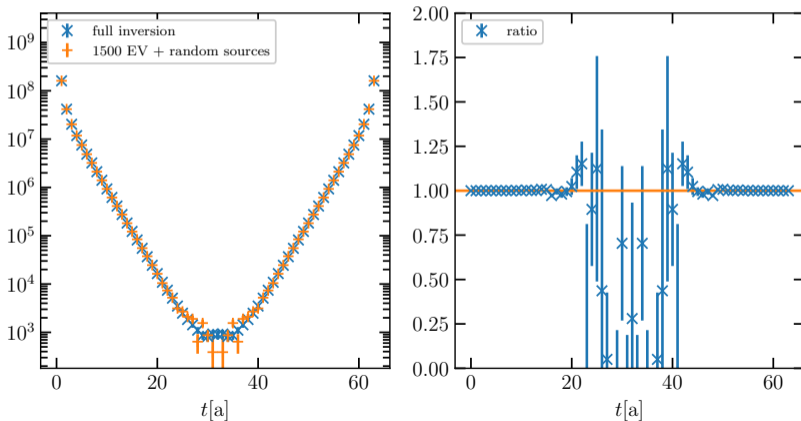
# Computational time of LMA and full inversion

- 1500 eigenvectors, 18 wall sources,  $T/2$  sources for the rest part



# Comparison LMA vs full inversion

- $48^3 \times 64$  box, 1500 eigenvectors, 18 wall sources,  $T/2$  sources for the rest part





# Simulation setup

- $L^3 \times T = 48^3 \times 64$
- $a = 0.1315$  fm
- Staggered fermions with four steps of stout smearing
- 48 lattice configurations

- Vector meson with LMA  
**[2407.10913]**
- $\pi\pi$ -orbits (with  $E < m_\rho$ ):
  - $|\vec{p}| = 1$ : pseudoscalar, axial vector and tensor taste
  - $|\vec{p}| = 2$ : pseudoscalar and axial vector taste
  - $|\vec{p}| = 3$ : pseudoscalar taste

# Analysis strategy

- Solve GEVP to extract contributing modes  
 $a_i \exp(-m_i t)$

$$\mathbf{C}(t)v_i(t_0) = \lambda_i(t, t_0)\mathbf{C}(t_0)v_i(t_0)$$

$$\lambda_i(t, t_0) \propto \exp(-m_i(t - t_0))$$

$$\mathbf{C}(t) = \left( \begin{array}{ccc|ccc} C_{t-4}^{\rho \rightarrow \rho} & C_{t-3}^{\rho \rightarrow \rho} & C_{t-2}^{\rho \rightarrow \rho} & C_{t-2}^{\rho \rightarrow \pi \pi^1} & C_{t-2}^{\rho \rightarrow \pi \pi^2} & \dots \\ C_{t-3}^{\rho \rightarrow \rho} & C_{t-2}^{\rho \rightarrow \rho} & C_{t-1}^{\rho \rightarrow \rho} & C_{t-1}^{\rho \rightarrow \pi \pi^1} & C_{t-1}^{\rho \rightarrow \pi \pi^2} & \dots \\ C_{t-2}^{\rho \rightarrow \rho} & C_{t-1}^{\rho \rightarrow \rho} & C_t^{\rho \rightarrow \rho} & C_t^{\rho \rightarrow \pi \pi^1} & C_t^{\rho \rightarrow \pi \pi^2} & \dots \\ \hline C_{t-2}^{\pi \pi^1 \rightarrow \rho} & C_{t-1}^{\pi \pi^1 \rightarrow \rho} & C_t^{\pi \pi^1 \rightarrow \rho} & C_t^{\pi \pi^1 \rightarrow \pi \pi^1} & C_t^{\pi \pi^1 \rightarrow \pi \pi^2} & \dots \\ C_{t-2}^{\pi \pi^2 \rightarrow \rho} & C_{t-1}^{\pi \pi^2 \rightarrow \rho} & C_t^{\pi \pi^2 \rightarrow \rho} & C_t^{\pi \pi^2 \rightarrow \pi \pi^1} & C_t^{\pi \pi^2 \rightarrow \pi \pi^2} & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots \end{array} \right)$$

# Analysis strategy

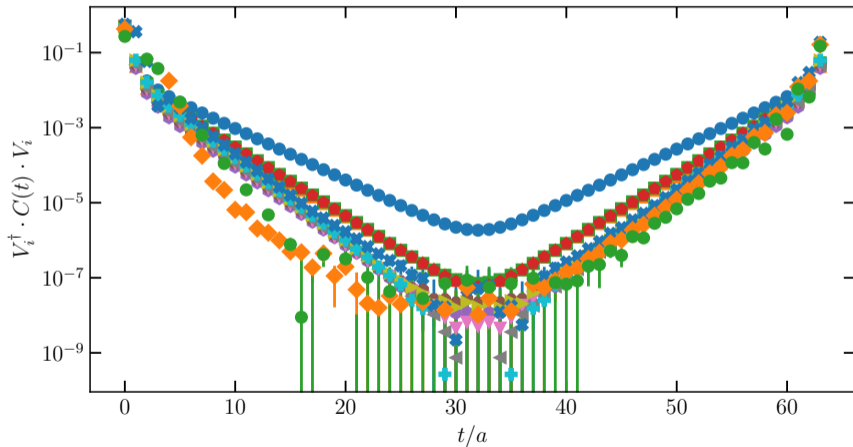
- Fit plateaus to effective masses
- Reconstruct the vector meson propagator

$$\tilde{\lambda}_i(t, t_0) = v_i^T(t_0) \mathbf{C}(t) v_i(t_0)$$

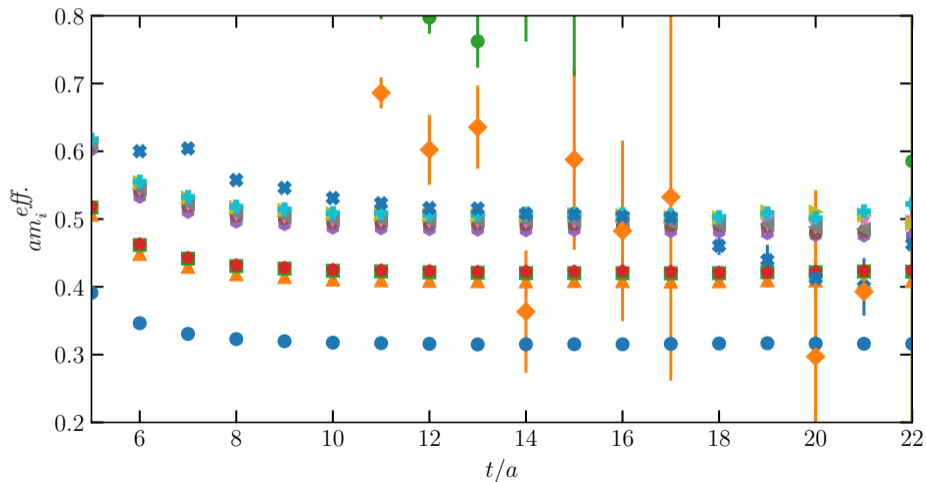
$$\rho_{\text{rec.}}(t) = \sum_i a_i \exp(-m_i t)$$

$$a_i^{\text{eff.}}(t) = \frac{(v_i^T \cdot \mathbf{C}_{\rho \bullet}(t))^2}{v_i^T \cdot \mathbf{C}(t) \cdot v_i} \exp(m_i^{\text{eff.}}(t) \cdot t)$$

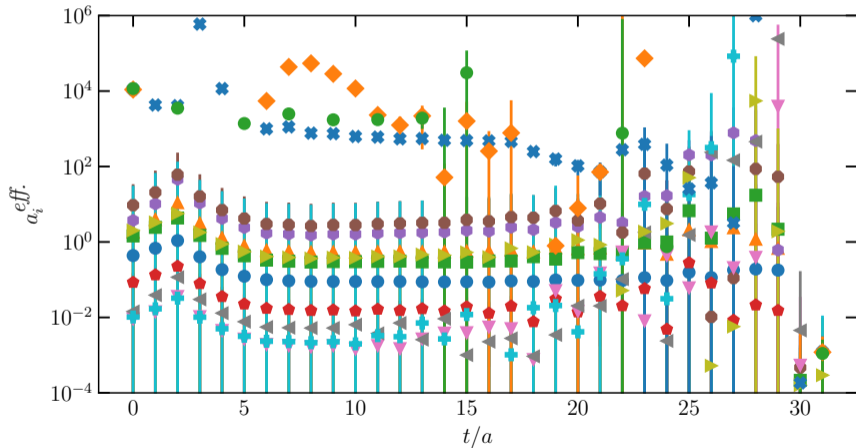
# Reconstructed states



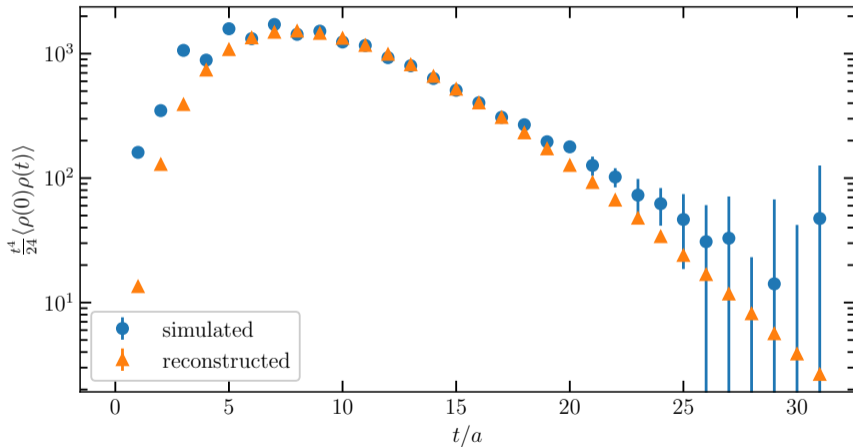
# Effective energies of reconstructed states



# Plateau fits of the contributions



# Reconstruction of the vector meson propagator



# Conclusion

Application of the formalism  
from **[2401.00514]**

Efficient method to compute  
the connected diagram for two  
pions

GEVP and reconstruction have to be in-  
vestigated further





# APPENDIX

# Multiplicities of the $\pi\pi$ -states

	$\ \vec{p}\ ^2 = 0$	$\ \vec{p}\ ^2 = 1$	$\ \vec{p}\ ^2 = 2$	$\ \vec{p}\ ^2 = 3$	$\ \vec{p}\ ^2 = 4$
$\ \vec{\xi}\ ^2 = 0$	0	1	1	1	1
$\ \vec{\xi}\ ^2 = 1$	0	2	3	2	2
$\ \vec{\xi}\ ^2 = 2$	0	2	3	2	2
$\ \vec{\xi}\ ^2 = 3$	0	1	1	1	1