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# STOCHASTIC NORMALIZING FLOWS FOR EFFECTIVE STRING THEORY

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31/07/2024

Lattice 2024

Liverpool

Based on:

M. Caselle, E.C., A. Nada, M. Panero

• JHEP 07 (2022) 015, arxiv:2201.08862

M. Caselle, E.C., A. Nada

• JHEP 02 (2024) 048, arxiv:2307.01107

• In prep., arxiv:2408.XXXX



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# OUTLINE

1. **EFFECTIVE STRING THEORY**
2. **STOCHASTIC NORMALIZING FLOWS**
3. **NUMERICAL RESULTS**
4. **OUTLOOKS**

# EFFECTIVE STRING THEORY

# EFFECTIVE STRING THEORY

Correlators of Polyakov loops modelled in terms of string partition functions:

$$\langle P(0)P^\dagger(R) \rangle \sim \int DX e^{-S_{EST}[X]} \equiv Z_{EST}$$

The main choice for  $S_{EST}$  is the Nambu-Goto (NG) action:

$$S_{NG} = \sigma \int d\xi^2 \sqrt{g}$$

- Anomalous at quantum level  $\rightarrow$  effective, large-distance description of Yang-Mills theories (low-energy universality theorem).
- Works only up to order  $1/R^5$   $\rightarrow$  first order approximation of a more general theory  $\rightarrow$  Beyond Nambu-Goto (BNG)  
[Aharony and Komargodski; 1302.6257], [Brandt and Meineri; 1603.06969],[Caselle; 2104.10486]

# NAMBU-GOTO STRING

Main method: zeta-function regularization

Main observables:

- Partition function → directly associated with the interquark potential. Well known at all the order.
- Correlation functions (e.g. width  $\sigma w^2$ ) → measure of the density of the chromoelectric flux tube

Analytical limits:

- Correlation functions
- Higher order corrections (Beyond NG)

See talks by:

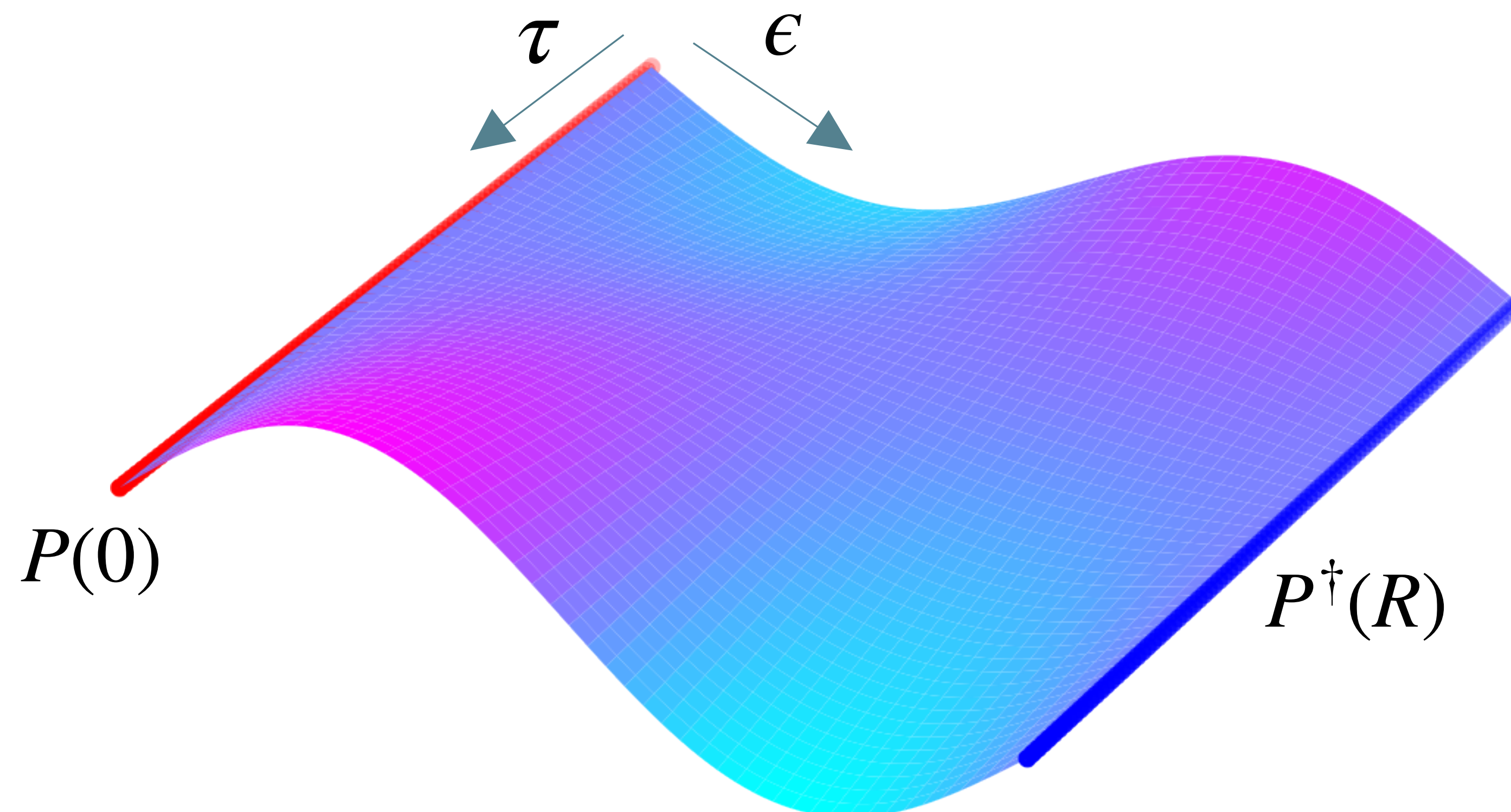
- D. Panfalone 29/07, 11:35
- A. Athenodorou 31/07, 11:35
- L. Verzichelli 31/07, 11:55
- M. Caselle 31/07, 12:35



# LATTICE NAMBU-GOTO STRING

$$S_{NG}(\phi) = \sigma \sum_{x \in \Lambda} \left[ \sqrt{1 + (\partial_\mu \phi(x))^2 / \sigma} - 1 \right]$$

- $d = 2 + 1$  target Yang-Mills
- $\sigma$  string tension
- $\Lambda$  : square lattice of size  $L \times R$ ,  $a = 1$
- $\phi(x) = \phi(\tau, \epsilon) \in \mathbb{R}$
- $\phi(\tau + L, \epsilon) = \phi(\tau, \epsilon)$
- $\phi(\tau, 0) = \phi(\tau, R) = 0$
- $\sigma w^2 = \langle \phi(\tau, R/2)^2 \rangle_\tau$

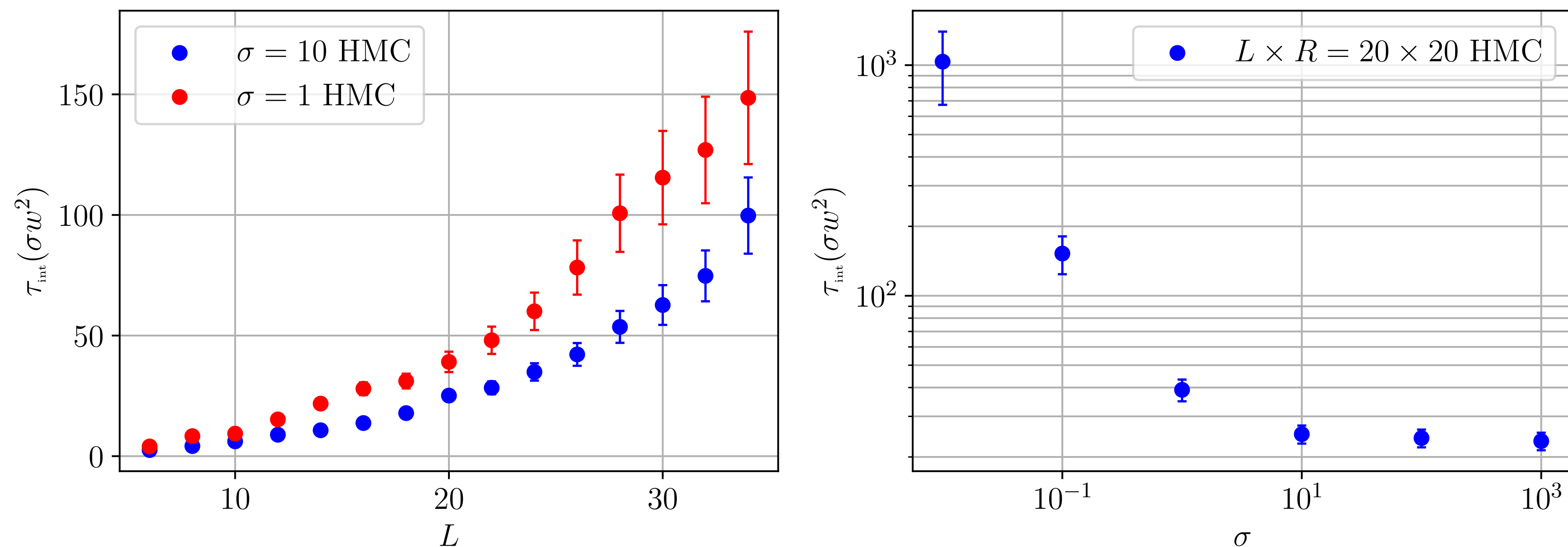


[Caselle, [EC](#), Nada; 2307.01107]

# LACKS OF NUMERICAL METHODS

Numerical problems:

- Strong non-linearity → critical theory (Critical Slowing Down)
- Estimation of partition functions



We did a proof-of-concept with continuous normalizing flows; however, we found scaling issues.

[Caselle, [EC](#), Nada; 2307.01107]

# STOCHASTIC NORMALIZING FLOWS



# NON-EQUILIBRIUM MCMC (NE-MCMC)

$$q_0 \simeq e^{-S_0} \xrightarrow{P_1} e^{-S_1} \xrightarrow{P_2} \dots \xrightarrow{P_N} e^{-S_N} \simeq p$$

1. **Thermalized**  $q_0$  "**prior**"
2.  $P_i \propto \exp(-S_i)$  **change along the processes** and satisfy detailed balance.
3.  $p = \exp(-S_N)/Z_N \rightarrow$  "target" distribution

**Forward** and **reverse** probability density:

$$q_0(\phi_0) \prod_{n=0}^{N-1} P_{i+1}[\phi_i \rightarrow \phi_{i+1}] = q_0(\phi_0) P_f[\phi_0, \dots, \phi_N]$$

$$p(\phi_N) \prod_{n=0}^{N-1} P_{i+1}[\phi_{i+1} \rightarrow \phi_i] = p(\phi_N) P_r[\phi_N, \dots, \phi_0]$$

**Remark:** no thermalization during the processes.

# CROOKS' THEOREM

Observe that:

$$\ln \frac{q_0(\phi_0)P_f[\phi_0, \dots, \phi_N]}{p(\phi_N)P_r[\phi_N, \dots, \phi_0]} = S_N(\phi_N) - S_0(\phi_0) - Q - \Delta F = W(\phi_0, \dots, \phi_N) - \Delta F = W_d(\phi_0, \dots, \phi_N)$$

(dimensionless) Work  $W$

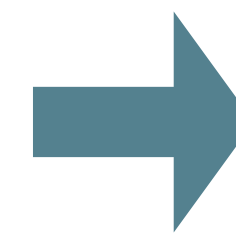
Where:

$$Q = \ln \frac{P_r[\phi_N, \dots, \phi_0]}{P_f[\phi_0, \dots, \phi_N]} = \sum_{n=0}^{N-1} \ln \frac{q_{n+1}(\phi_{n+1})}{q_{n+1}(\phi_n)} = \sum_{n=0}^{N-1} \left( S_{n+1}(\phi_{n+1}) - S_{n+1}(\phi_n) \right)$$

Detailed Balance

Thus:

$$\frac{q_0(\phi_0)P_f[\phi_0, \dots, \phi_N]}{p(\phi_N)P_r[\phi_N, \dots, \phi_0]} = \frac{\mathcal{P}_f(W_d)}{\mathcal{P}_r(-W_d)} = e^{W_d}$$



**Crooks Theorem**

[Crooks; cond-mat/9901352]

# JARZYNSKI'S EQUALITY

Observe also:

$$1 = \int \prod_{i=0}^N d\phi_i q_0(\phi_0) P_f[\phi_0, \dots, \phi_N] \left( \frac{p(\phi_N) P_r[\phi_N, \dots, \phi_0]}{q_0(\phi_0) P_f[\phi_0, \dots, \phi_N]} \right) = \langle e^{-W_d} \rangle_f$$

**Jarzynski's equality**

[Jarzynski; cond-mat/9610209]

$$\langle e^{-W} \rangle_f = e^{-\Delta F} \quad \langle \mathcal{O} \rangle_{\phi \sim p} = \langle \mathcal{O} e^{-W_d} \rangle_f$$

**Non-Equilibrium Ensemble**

**Equilibrium Quantity**

Equivalent to **Annealed Importance Sampling** and widely applied to LFT

[Neal; physics/9803008]

See:

- Talk by A. Bulgarelli 29/07, 14.35
- Poster by D. VDACCHINO 30/07

# STOCHASTIC NORMALIZING FLOWS

Stochastic Normalizing Flows (SNFs) combine NE-MCMC update and Normalizing Flows layers:

$$\phi_0 \longrightarrow g_{\theta}^1(\phi_0) \xrightarrow{P_1} \phi_1 \longrightarrow g_{\theta}^2(\phi_1) \xrightarrow{P_2} \dots \xrightarrow{P_N} \phi_N = \phi$$

Where  $P_i$  are MCMC update and  $g_{\theta}^i$  are Normalizing Flows (NFs) layers:

$$g_{\theta}^{n+1} : q^n \rightarrow q_{\theta}^{n+1} \quad \phi_{n+1} = g_{\theta}(\phi_n) \quad \ln(q_n(\phi_n)/q_{n+1}(\phi_{n+1})) = \ln |\det J_{g^{n+1}}(\phi_n)|$$

[Rezende et al.; 1505.05770],[Wu et al.; 2002.06707],[Caselle, EC, Nada, Panero; 2201.08862]

# DISSIPATED WORK $W_d^\theta$

We have now:

$$W_d^\theta = W_\theta(\phi_0, \dots, \phi_N) - \Delta F = S_N(\phi_N) - S_0(\phi_0) - Q_\theta - \Delta F$$

Where:

$$Q_\theta = \sum_{n=0}^{N-1} \left( S_{n+1}(\phi_{n+1}) - S_{n+1}(\phi_n) + \ln |\det J_{g_\theta^n}| \right)$$

We can now train a SNF by minimizing:

$$\mathcal{L}(\theta) = \langle W_d^\theta \rangle_f = D_{KL}(q_0 P_f || p P_r) \geq 0 \quad \longrightarrow \quad \langle W_\theta \rangle_f \geq \Delta F \quad \longrightarrow \quad \underline{\text{Second Law!}}$$

Measure how reversible the process is.

Lower  $W_d \rightarrow$  numerically stabler exponential averages

[Wu+; 2002.06707],[Caselle, EC, Nada, Panero; 2201.08862]

See talk by A. Nada

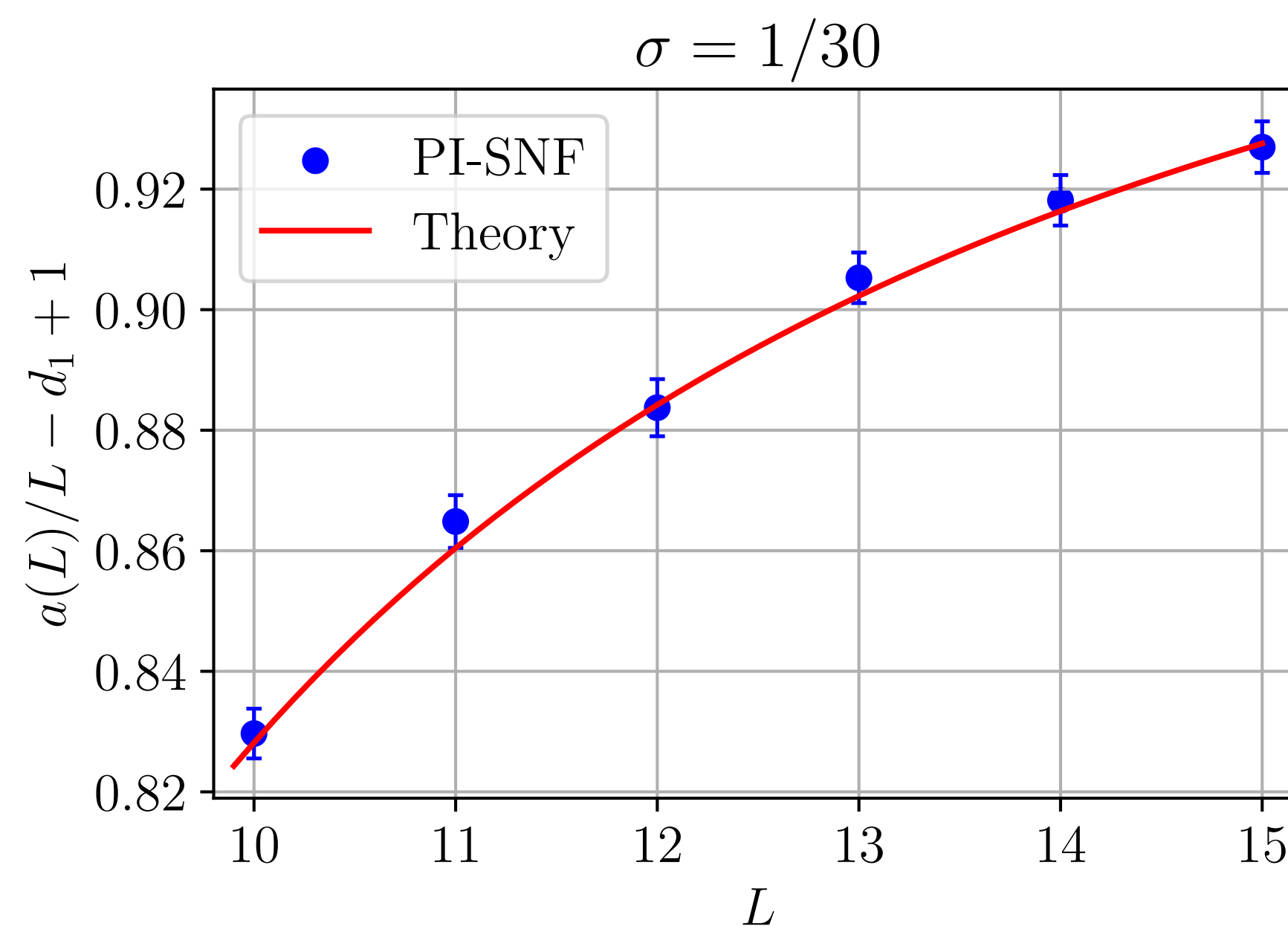
# NUMERICAL RESULTS



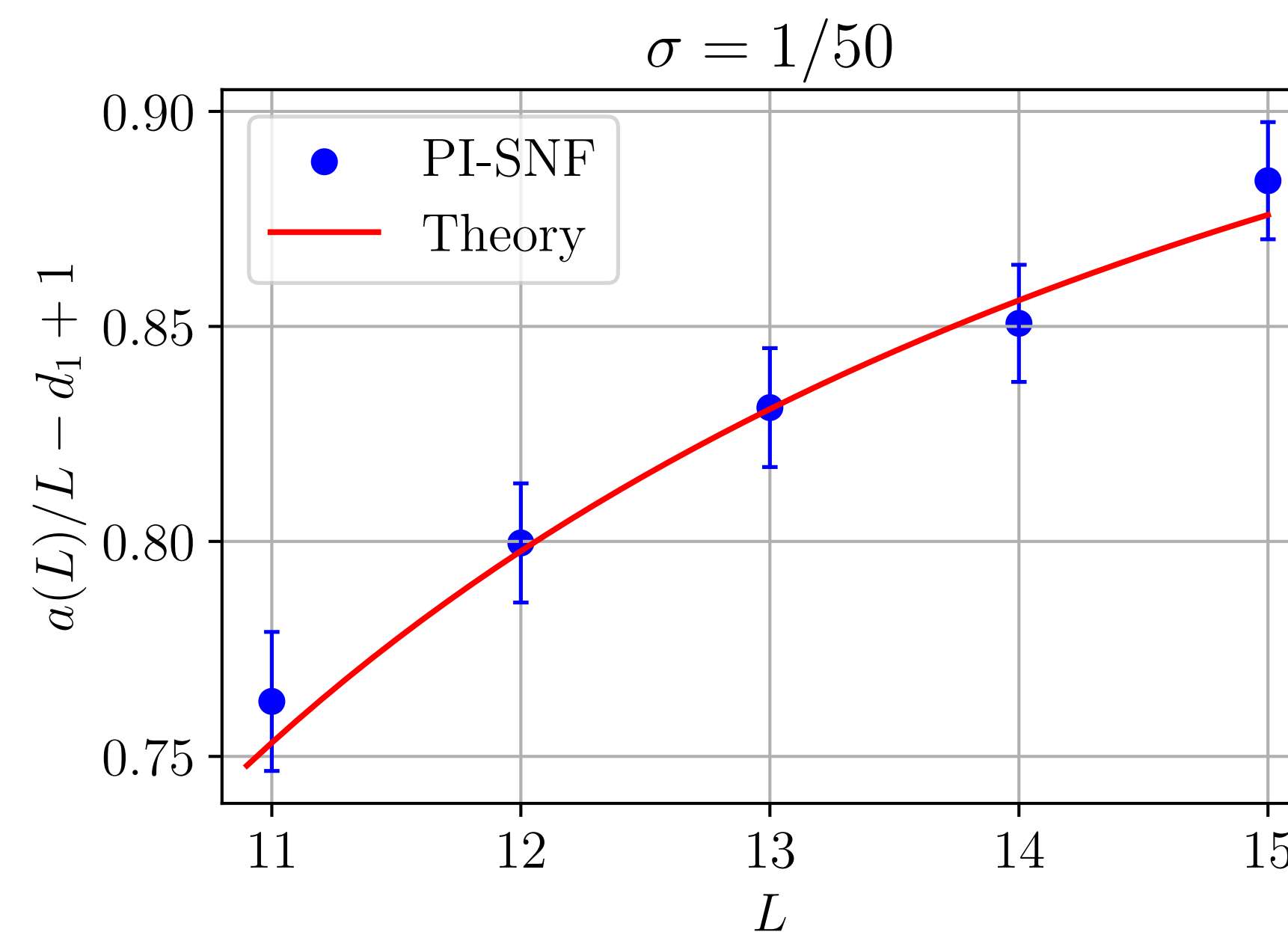
# NG FREE ENERGY $R \gg L$

$$-\log Z = \sigma RL \sqrt{1 - \frac{\pi}{3\sigma L^2}} + \dots$$

Fitted:  $-1.03(2)$ ,  
Target:  $-1.047\dots$



Fitted:  $-1.04(7)$ ,  
Target:  $-1.047\dots$



# NG WIDTH $R \gg L$

$\sigma = 1/10$

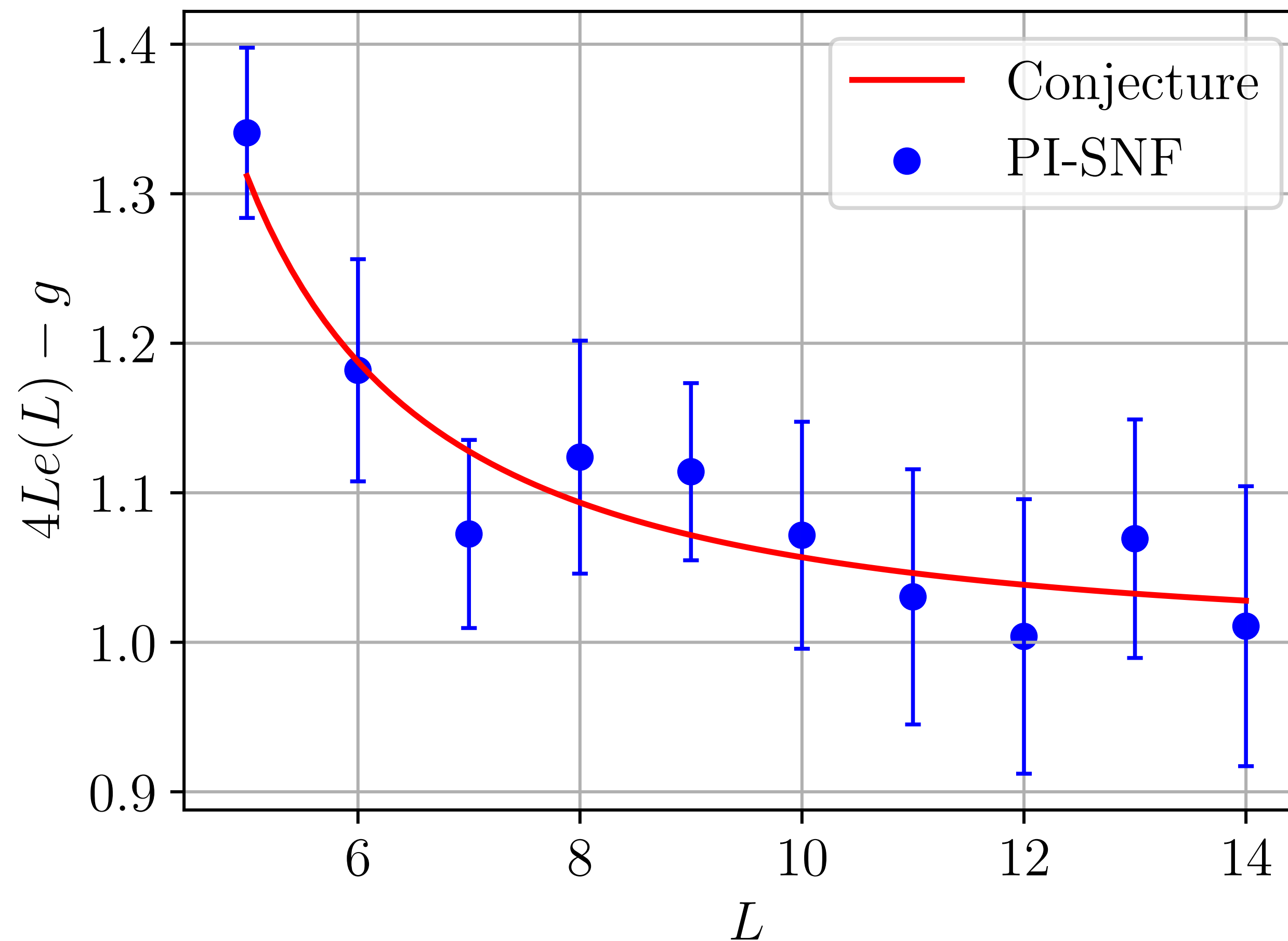
Conjecture:

$$\sigma w^2 = \frac{1}{\sqrt{1 - \frac{\pi}{3\sigma L^2}}} \frac{R}{4L} + \dots$$

NG:  $\sigma(L)/\sigma$

Gaussian solution

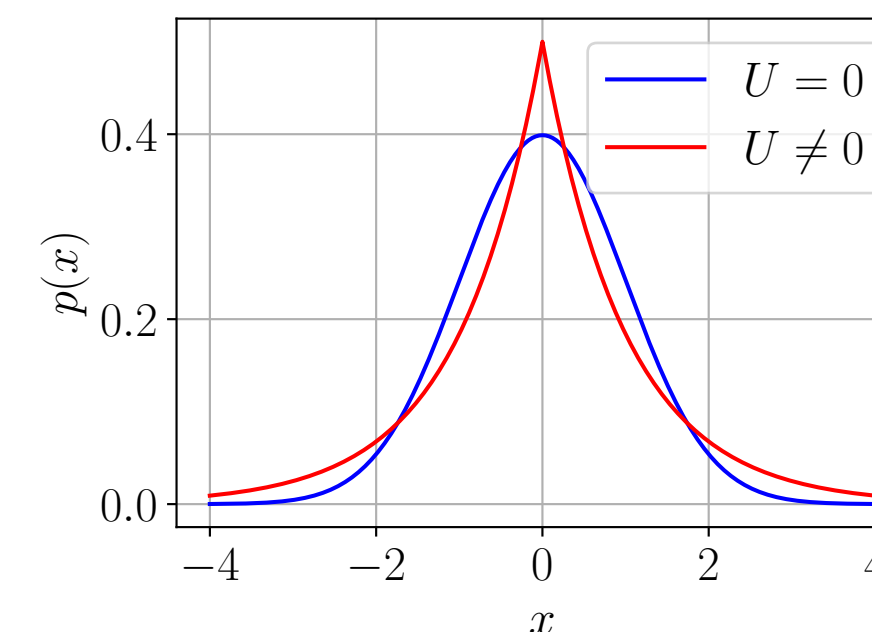
Fitted:  $-1.09(8)$ , target:  $-1.047\dots$



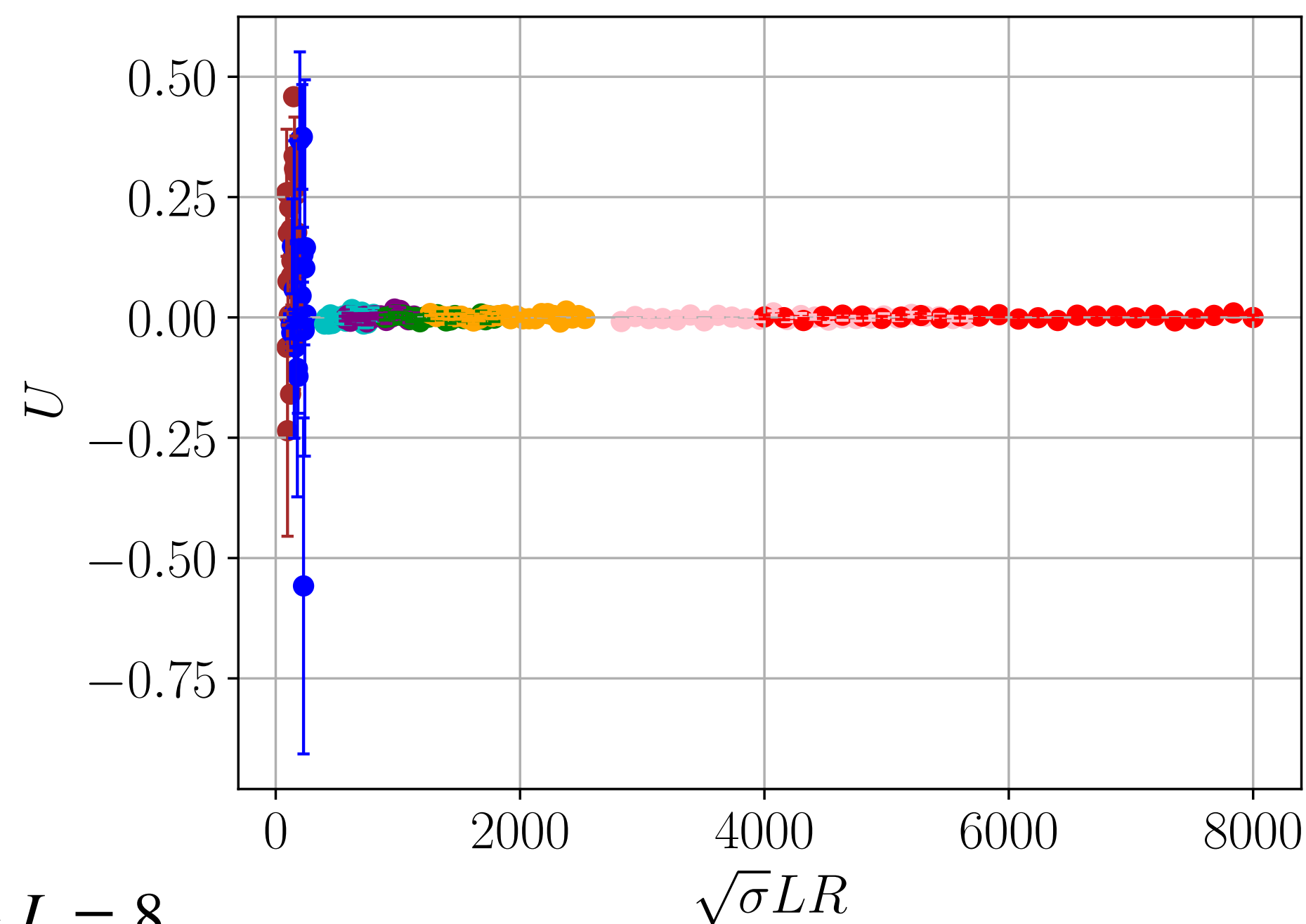
[Caselle, [EC](#), Nada; 2309.14983][Caselle;1004.3875]

# IS THE NG FLUX TUBE SHAPE GAUSSIAN?

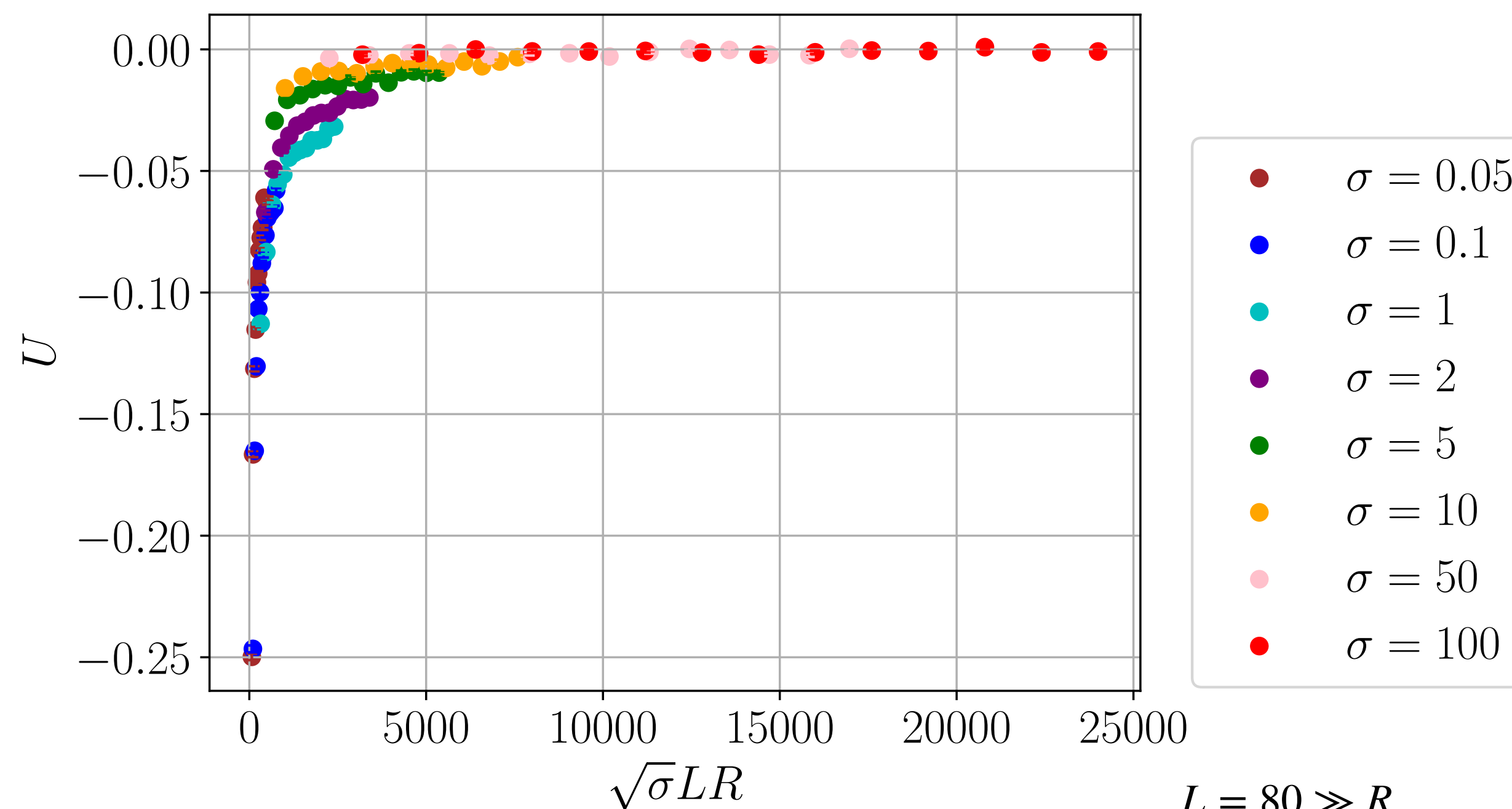
$$U = 1 - \frac{\langle \phi^4(\tau, R/2) \rangle_\tau}{3 \langle \phi^2(\tau, R/2) \rangle_\tau^2}$$



$\sigma = 100$



$R \gg L = 8$

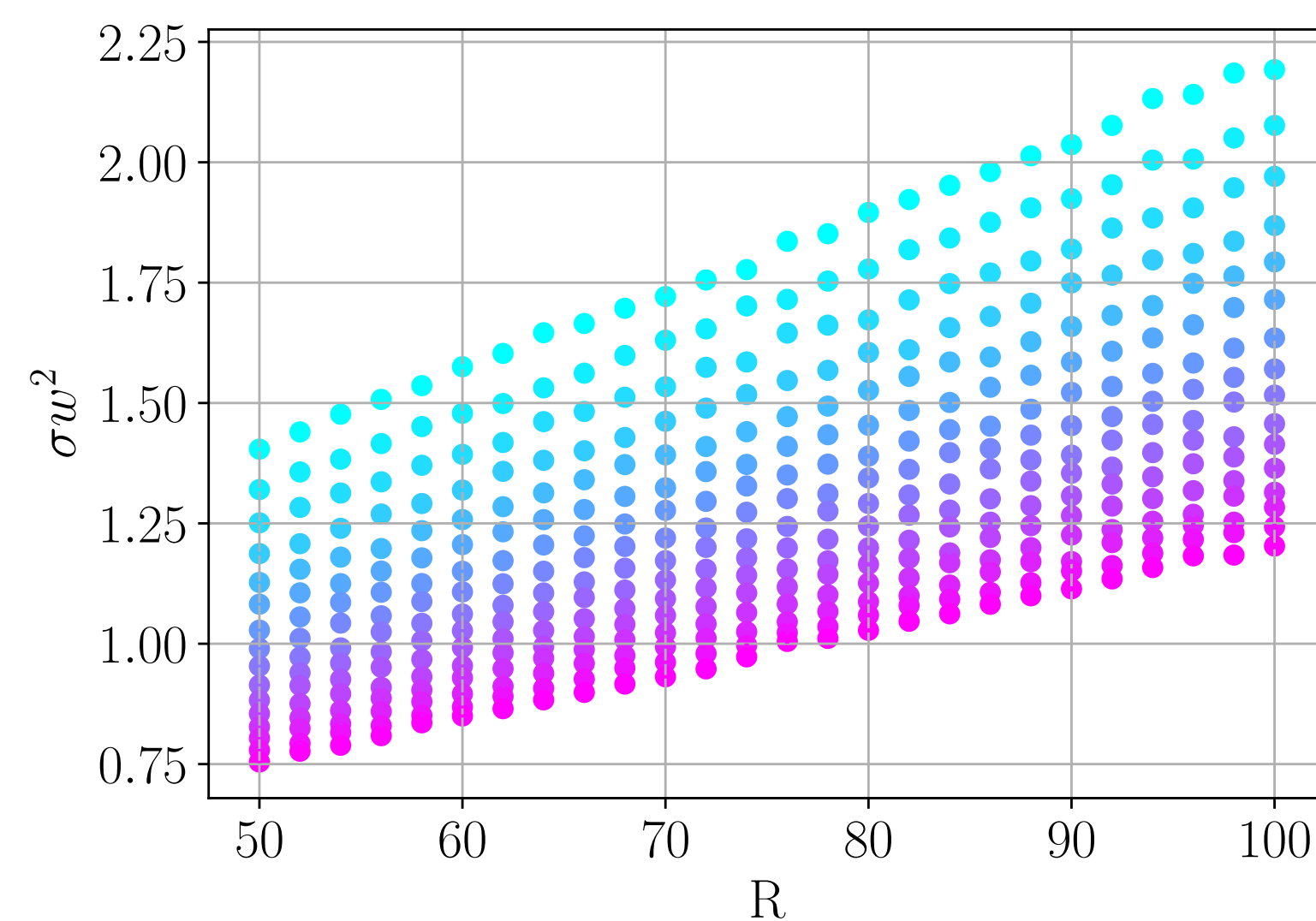


$L = 80 \gg R$

# BEYOND NG: WIDTH

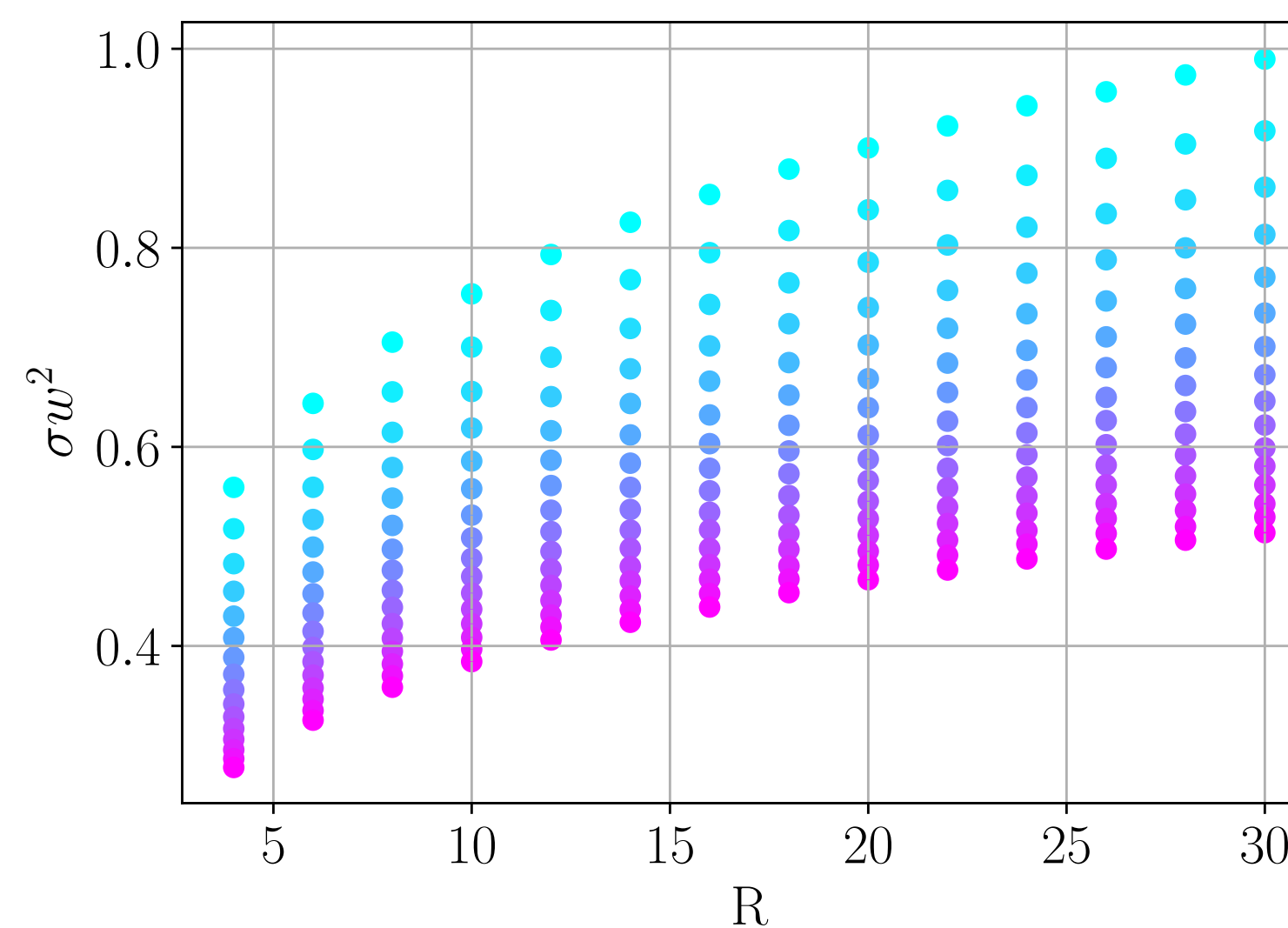
$$S_{EST}(\phi) = S_{NG}(\phi) + \kappa_2 \mathcal{K}^2(\phi)$$

$$\mathcal{K}^2(\phi) = \sum_{(\tau, \epsilon) \in \Lambda} \mathcal{L}^2(\phi(\tau, \epsilon)) = \sum_{(\tau, \epsilon) \in \Lambda} (\partial_\tau \partial_\tau \phi(\tau, \epsilon))^2 + (\partial_\epsilon \partial_\epsilon \phi(\tau, \epsilon))^2 + 2(\partial_\tau \partial_\epsilon \phi(\tau, \epsilon))^2$$

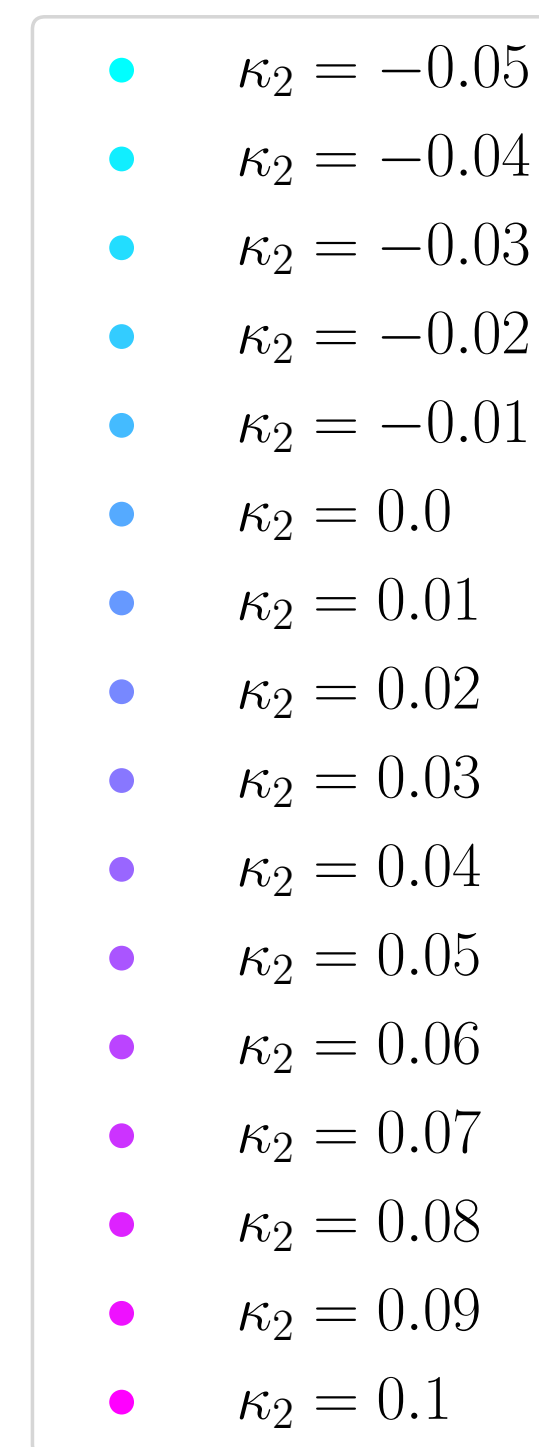


$\sigma = 100$

$R \gg L = 20$



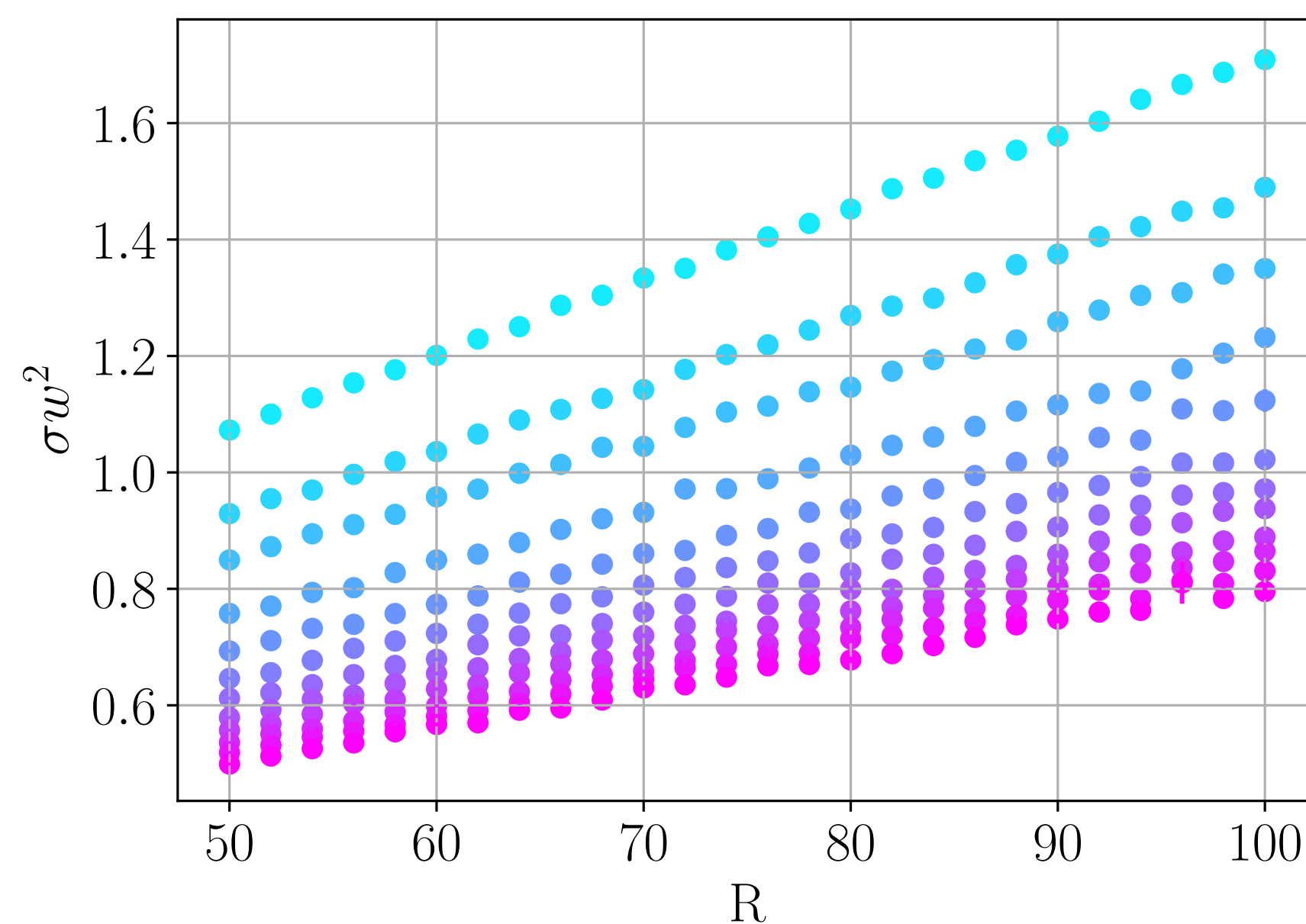
$L = 80 \gg R$



# BEYOND NG: WIDTH

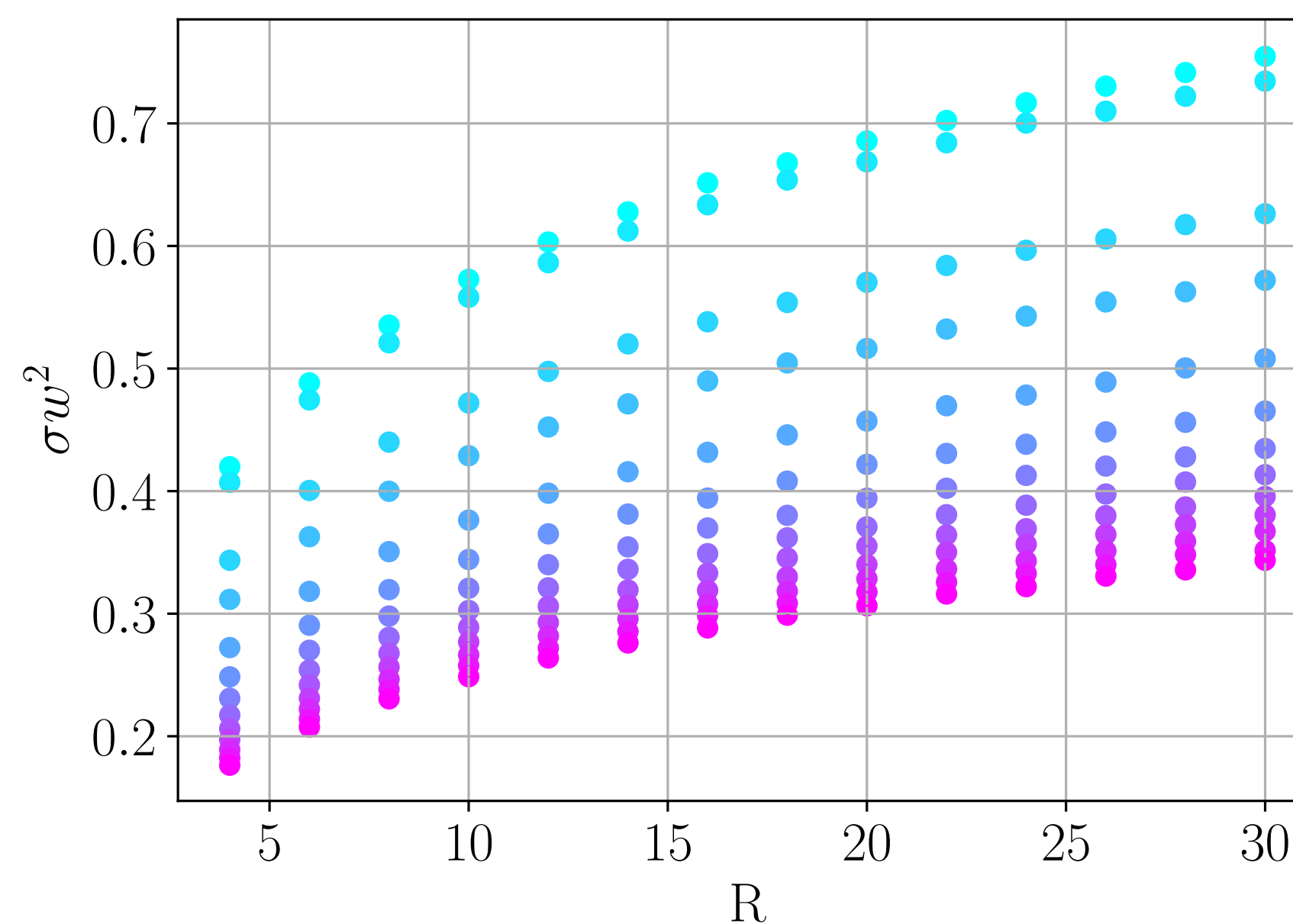
$$S_{EST}(\phi) = S_{NG}(\phi) + \kappa_4 \mathcal{K}^4(\phi)$$

$$\mathcal{K}^4(\phi) = \sum_{x \in \Lambda} (\mathcal{L}^2(\phi(x)))^2$$

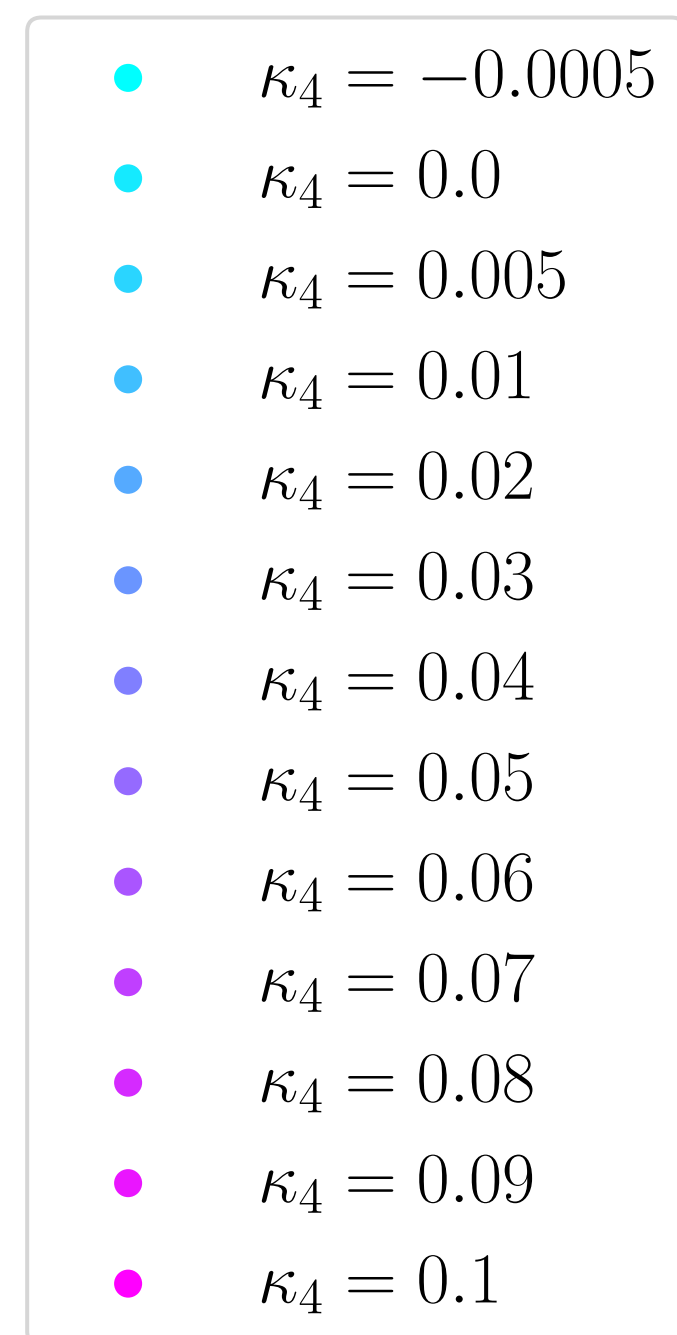


$\sigma = 100$

$R \gg L = 20$

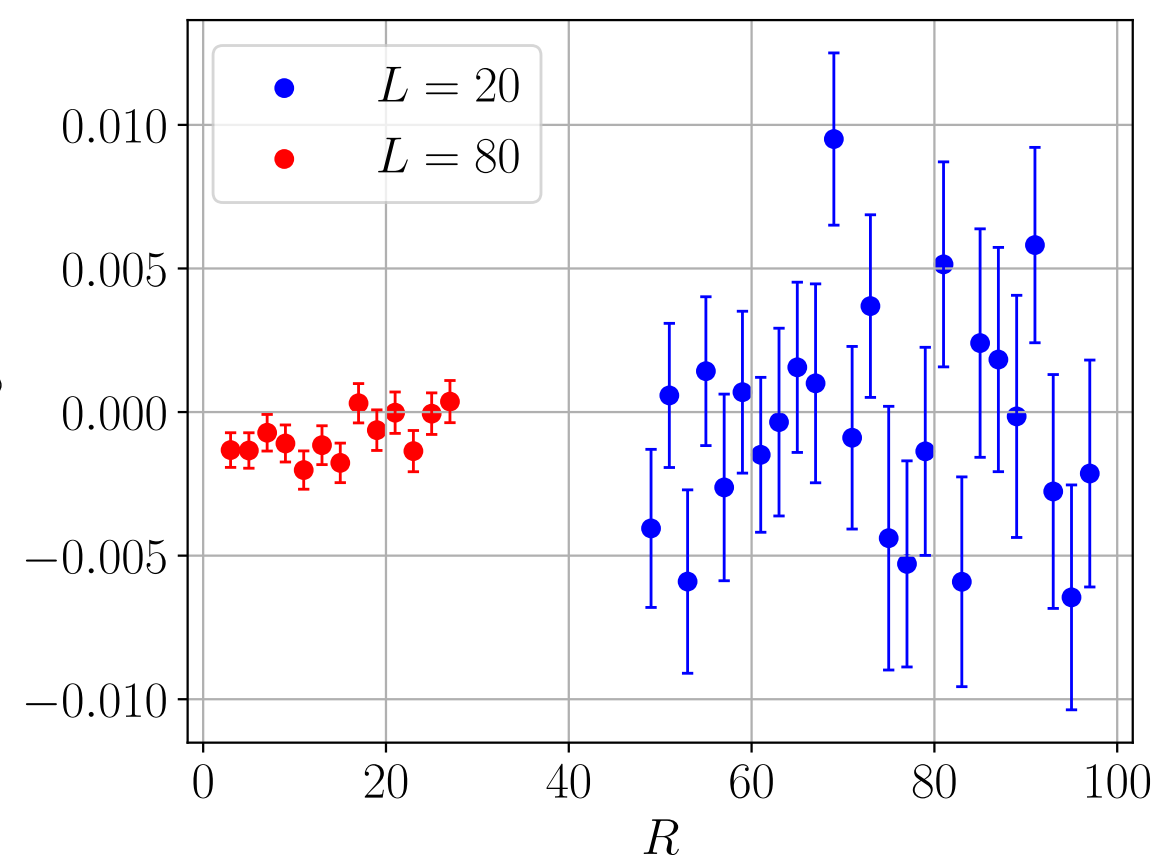
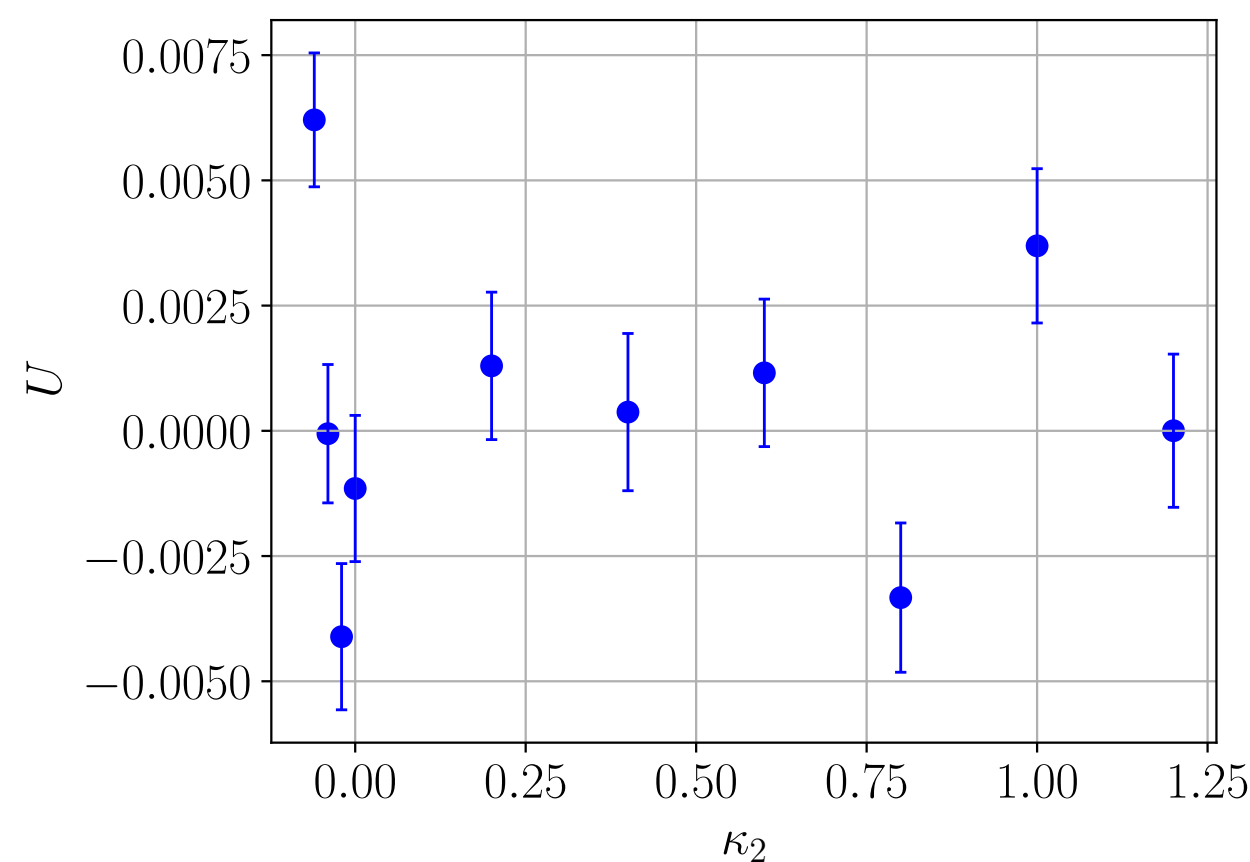


$L = 80 \gg R$

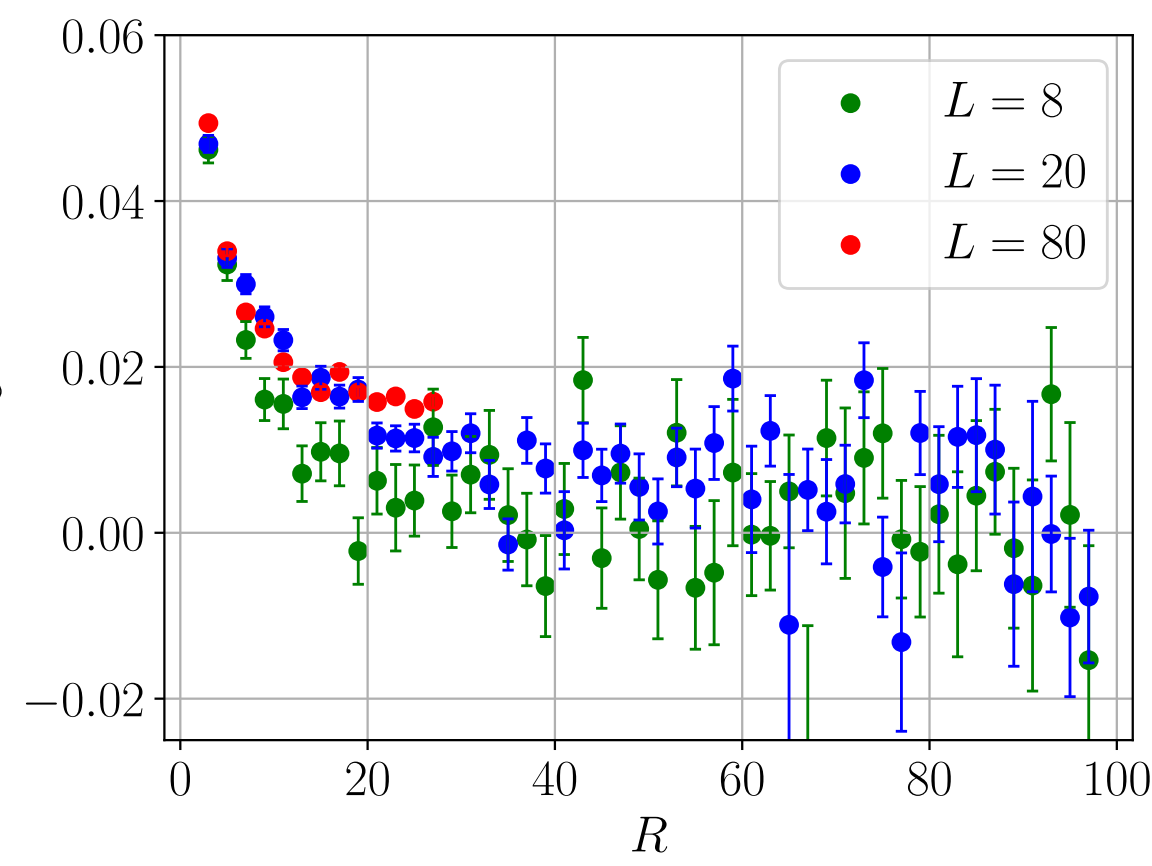
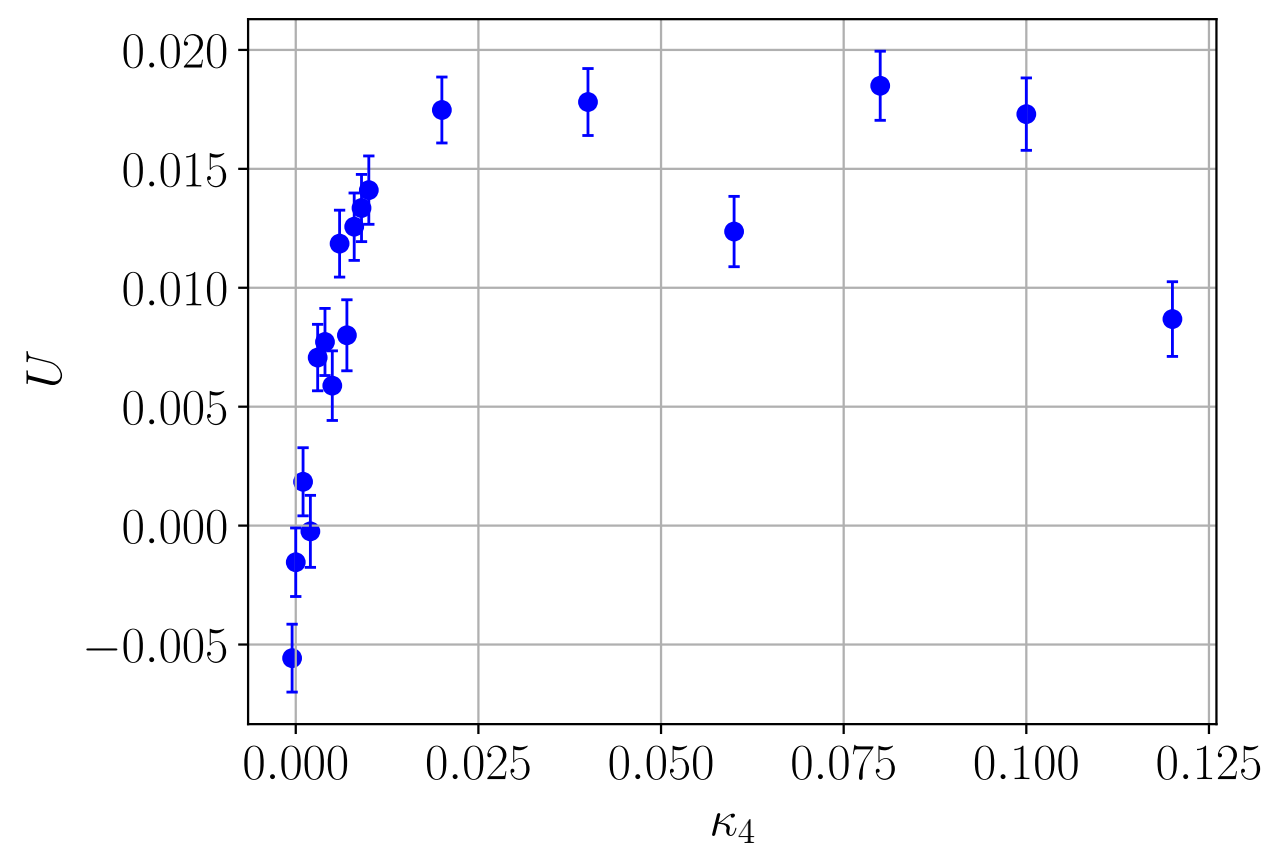


# BEYOND NG: GAUSSIANTY

$\sigma = 100$



$$S_{NG} + \kappa_2 \mathcal{K}^2$$



$$S_{NG} + \kappa_4 \mathcal{K}^4$$



# OUTLOOKS

- Flow-based sampler can be successfully applied to sample Lattice EST:
  1. Numerical solution of the high temperature NG width
  2. Numerical studies of the Beyond NG EST and new interesting observable
  
- EST provided a challenging laboratory for SNFs:
  1. Toward new applications in lattice field theory (see talk by A. Nada on SNFs for  $SU(3)$ )

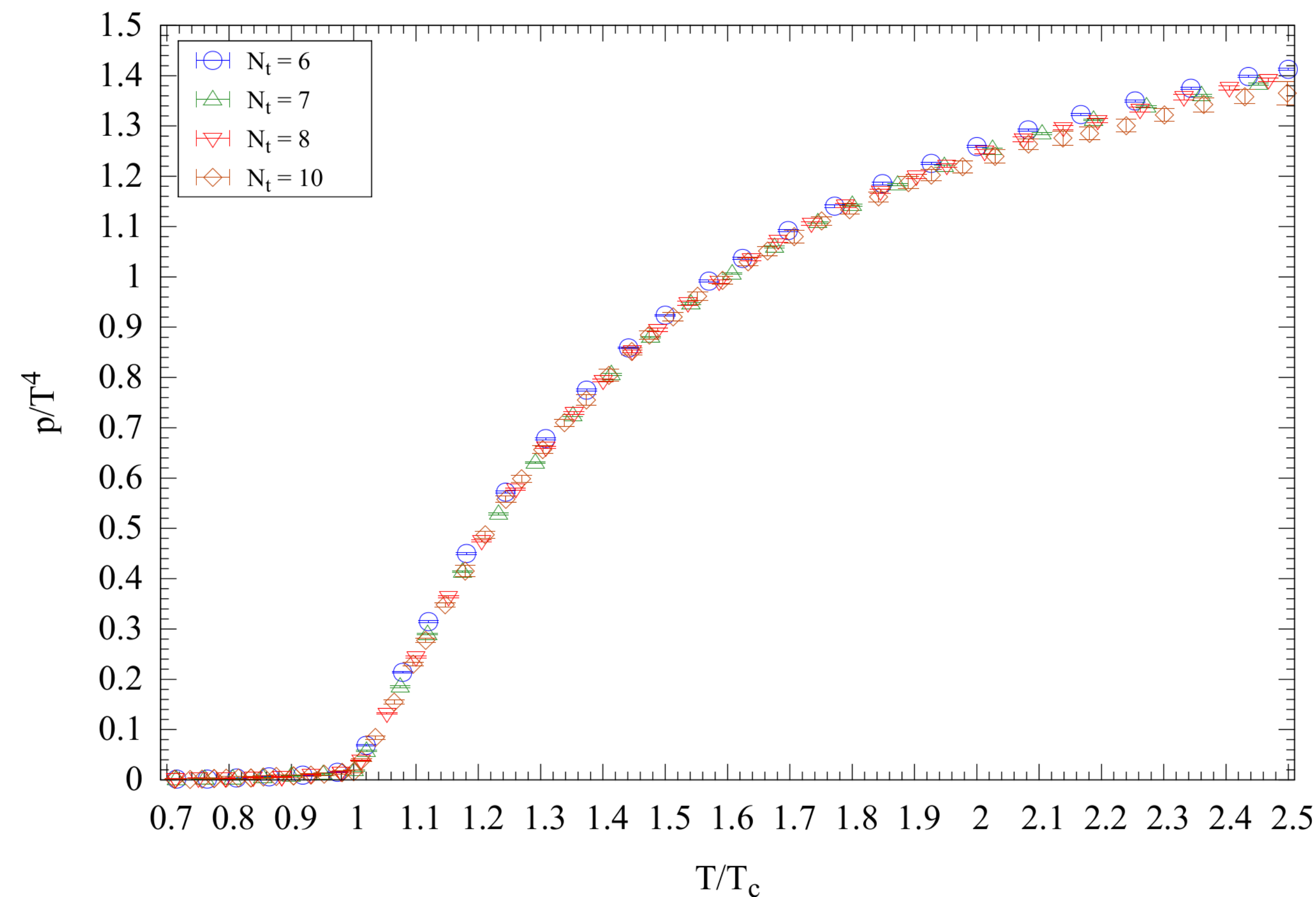
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**THANK YOU FOR YOUR  
ATTENTION!**

# NE-MCMC FOR LFT

Jarzynski's equality has been exploited to obtain state-of-the-arts results in LFT:

- Interface free energy.  
[Caselle et al.; 1604.05544]
- $SU(3)$  e.o.s.  
[Caselle et al.; 1801.03110]
- Running coupling  
[Francesconi et al.; 2003.13734]
- Entanglement entropy  
[Bulgarelli and Panero; 2304.03311, 2404.01987]
- Topological freezing  
[Bonanno et al.; 2402.06561]



# STOCHASTIC NORMALIZING FLOWS: $W_d$

Forward and Reverse transition probabilities of NF layers can be written as:

$$P[\phi_n \rightarrow \phi_{n+1}] = \delta(\phi_{n+1} - g_\theta^n(\phi_n)) \quad P[\phi_{n+1} \rightarrow \phi_n] = \delta(\phi_n - (g_\theta^n)^{-1}(\phi_{n+1}))$$

And satisfies:

$$\ln(P[\phi_{n+1} \rightarrow \phi_n]/P[\phi_n \rightarrow \phi_{n+1}]) = \ln(q_n(\phi_n)/q_{n+1}(\phi_{n+1})) = \ln |\det J_{g_\theta^n}(\phi_n)|$$

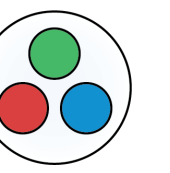
We have now:

$$W_d^\theta = W_\theta(\phi_0, \dots, \phi_N) - \Delta F = S_N(\phi_N) - S_0(\phi_0) - Q_\theta - \Delta F$$

Where:

$$Q_\theta = \sum_{n=0}^{N-1} \left( S_{n+1}(\phi_{n+1}) - S_{n+1}(\phi_n) + \ln |\det J_{g_\theta^n}| \right)$$

[Wu+; 2002.06707],[Caselle, E.C., Nada, Panero; 2201.08862]



# SNFS: RELATED WORKS

- Annealed Importance Sampling: Equivalent to Jarzynski's equality. Used in the original SNF paper  
[Neal; physics/9803008]
- Sequential Monte Carlo: Generalization of AIS.  
[Dai+; 2007.11936]
- SNF idea reworked in CRAFT  
[Matthews+; 2201.13117]
- An hybrid (deterministic/stochastic) approach with no neural networks has been proposed also by Jarzynski in 2011  
[Vaikuntanathan and Jazynski; 1101.2612]
- FAB: combination of NFs and AIS.  
[Midgley+; 2208.01893]
- Exact work for discretized Langevin dynamics.  
[Sivak+; 1107.2967]

# TECHNICAL DETAILS: NG

- Prior massless free boson and linear protocol in  $t = 1/\sigma \rightarrow$  Inspired by Irrelevant Perturbations
- HMC for stochastic updates
- Affine coupling layers, 3 convolutional layers with  $3 \times 3 \times 16$  kernels and a two channels output layer. Each blocks (even-odd) share the same network



# NG BINDER 2

