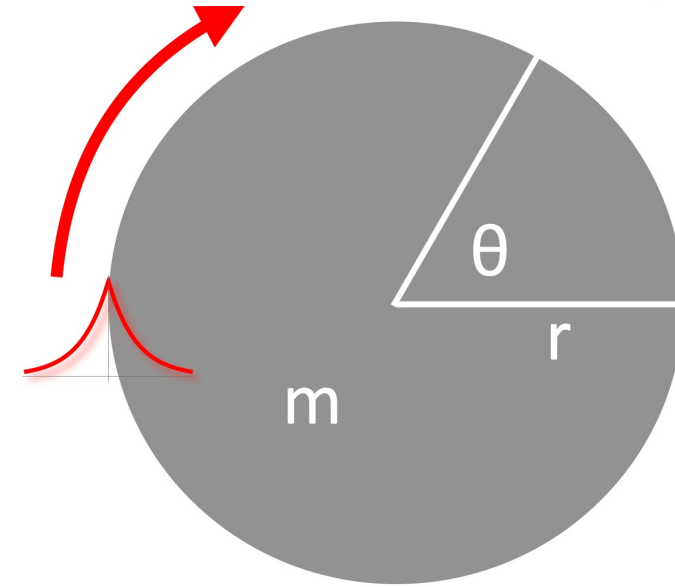
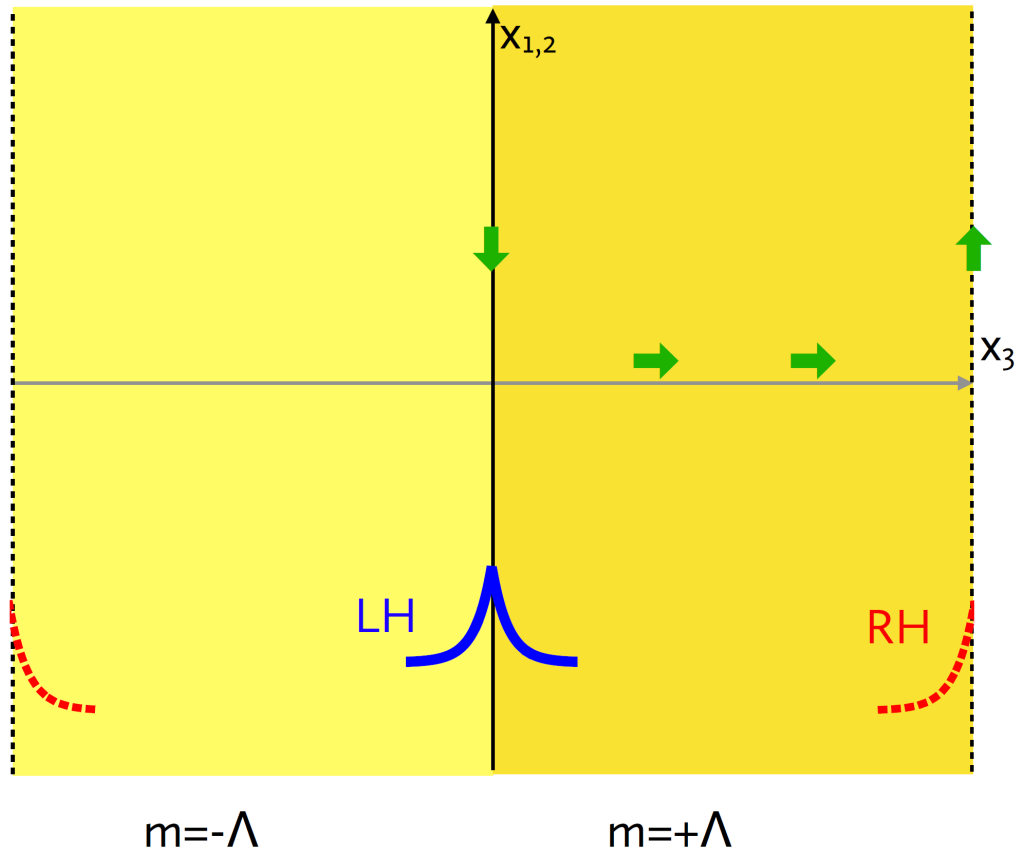


Weyl fermion on a lattice: A path to lattice chiral gauge theory



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Lattice 2024

Based on *Phys.Rev.Lett.* **132** (2024) **14**, 141604

with David Kaplan, University of Washington

Plan of the talk

- What are chiral gauge theories?
- Why is it hard to put them on the lattice?
- A few of the past attempts, that are yet to work or don't work.
- What is new and why is it exciting?

Chiral gauge theories

Even dimensional world with massless fermions and gauge field.

Chiral symmetry is gauged.

Fermion mass terms transform under gauge transformation.

So, a simple mass term is disallowed.

Example: The standard model.

Good chiral symmetry is essential

Why is chiral symmetry hard?

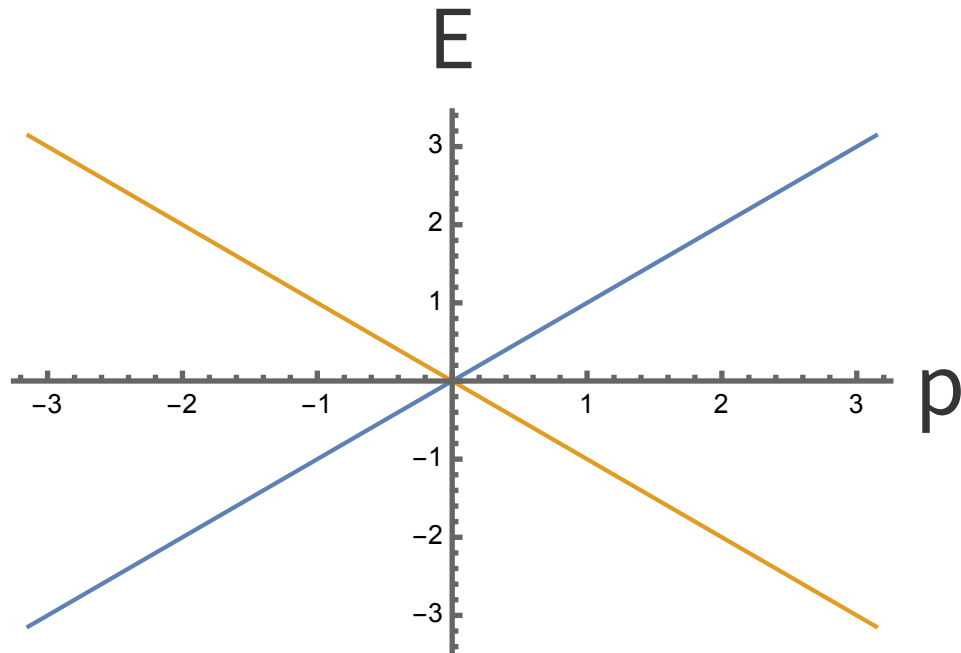
Nielsen-Ninomiya theorem is one of the major reasons.

Nielsen Ninomiya (1981): Cannot formulate Dirac fermion with exact chiral symmetry without an unwanted doubling of all fermion species.

Any number of massless Dirac fermions needed for global chiral symmetry

Odd number of Weyl fermions needed for gauging chiral symmetry

Why chiral symmetry is hard: Dispersion (1 spatial dimension)



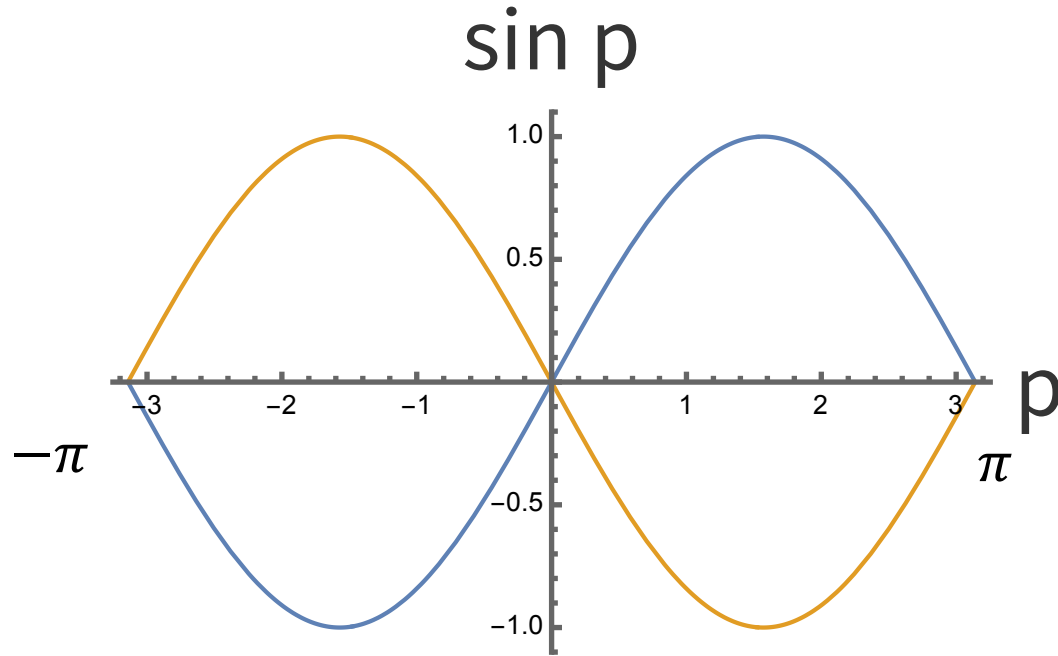
Continuum dispersion for a
Dirac Hamiltonian

The no-go is better visualized
using dispersion relation in
Minkowski space-time (time
continuous).

Hamiltonian formulation.

$$E = \pm p$$

Brilluoin zones (Dirac)



Two Dirac fermions

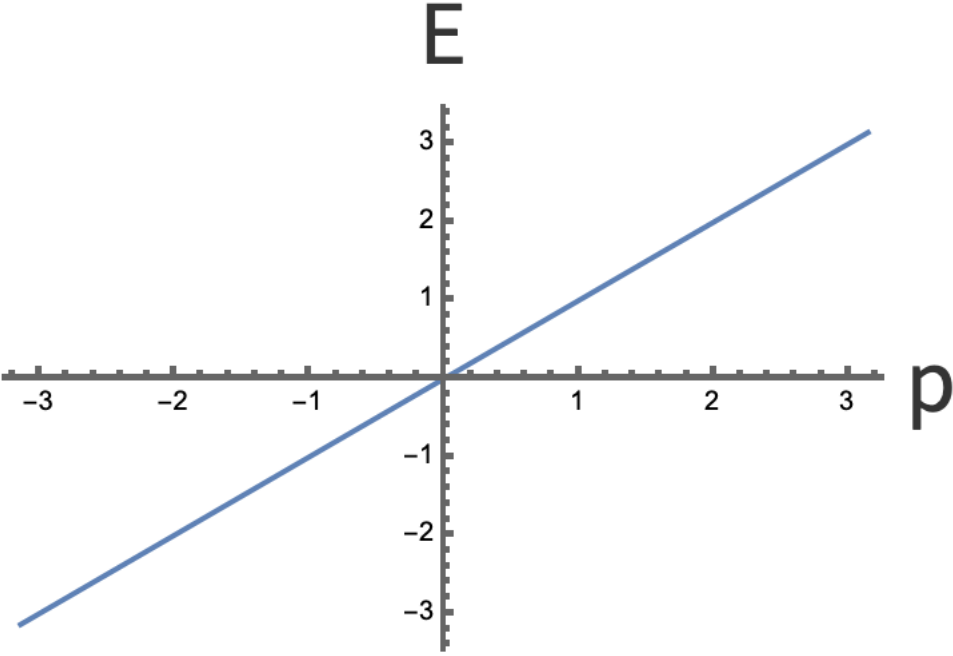
Lattice in space.

Time not discretized.

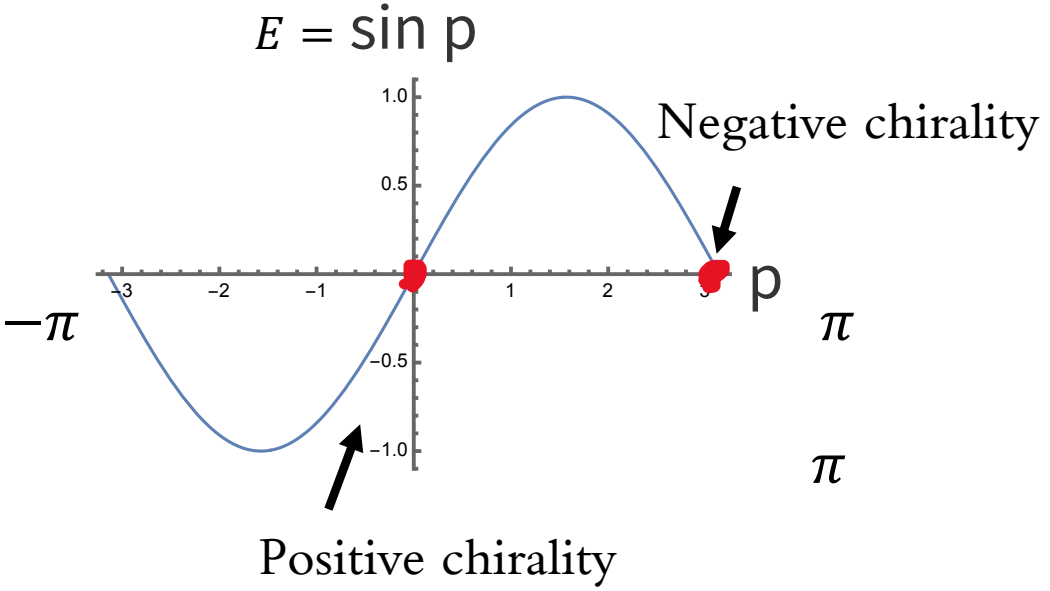
Solving the naively discretized Dirac Hamiltonian with eigenvalues $\pm \sin p$

$$E = \pm \sin p$$

Brilliuoin zones(Weyl)



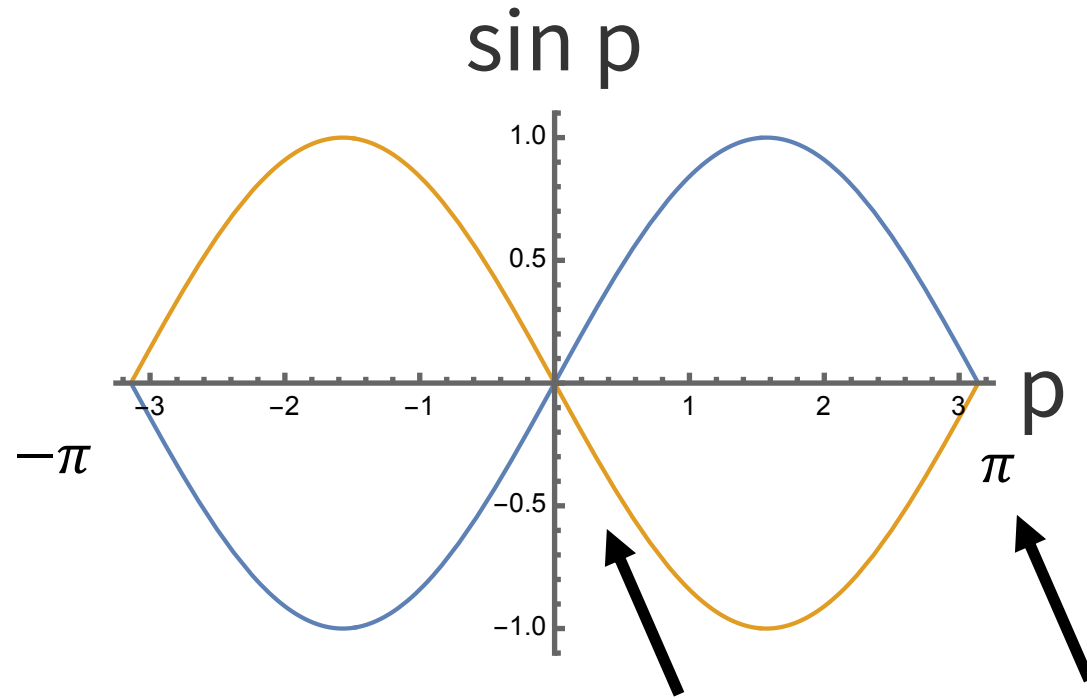
Continuum



Lattice

Even number of zero crossing of periodic functions

Wilson term for Dirac



Gaplessness is not protected

Wilson term removes this.
But kills chiral symmetry

Lattice in space.

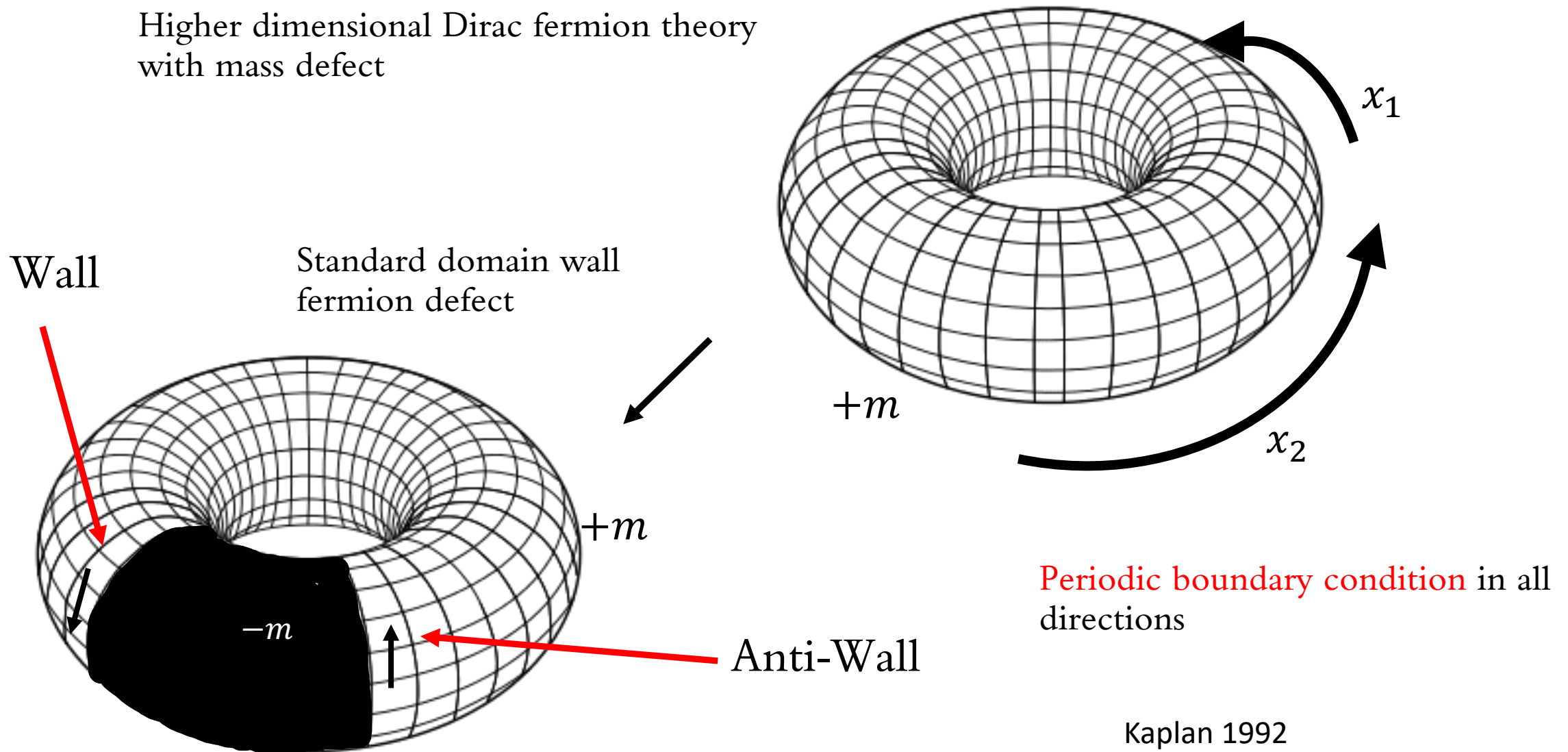
Time not discretized.

Solving the naively discretized
Dirac Hamiltonian with
eigenvalues $\pm \sin p$

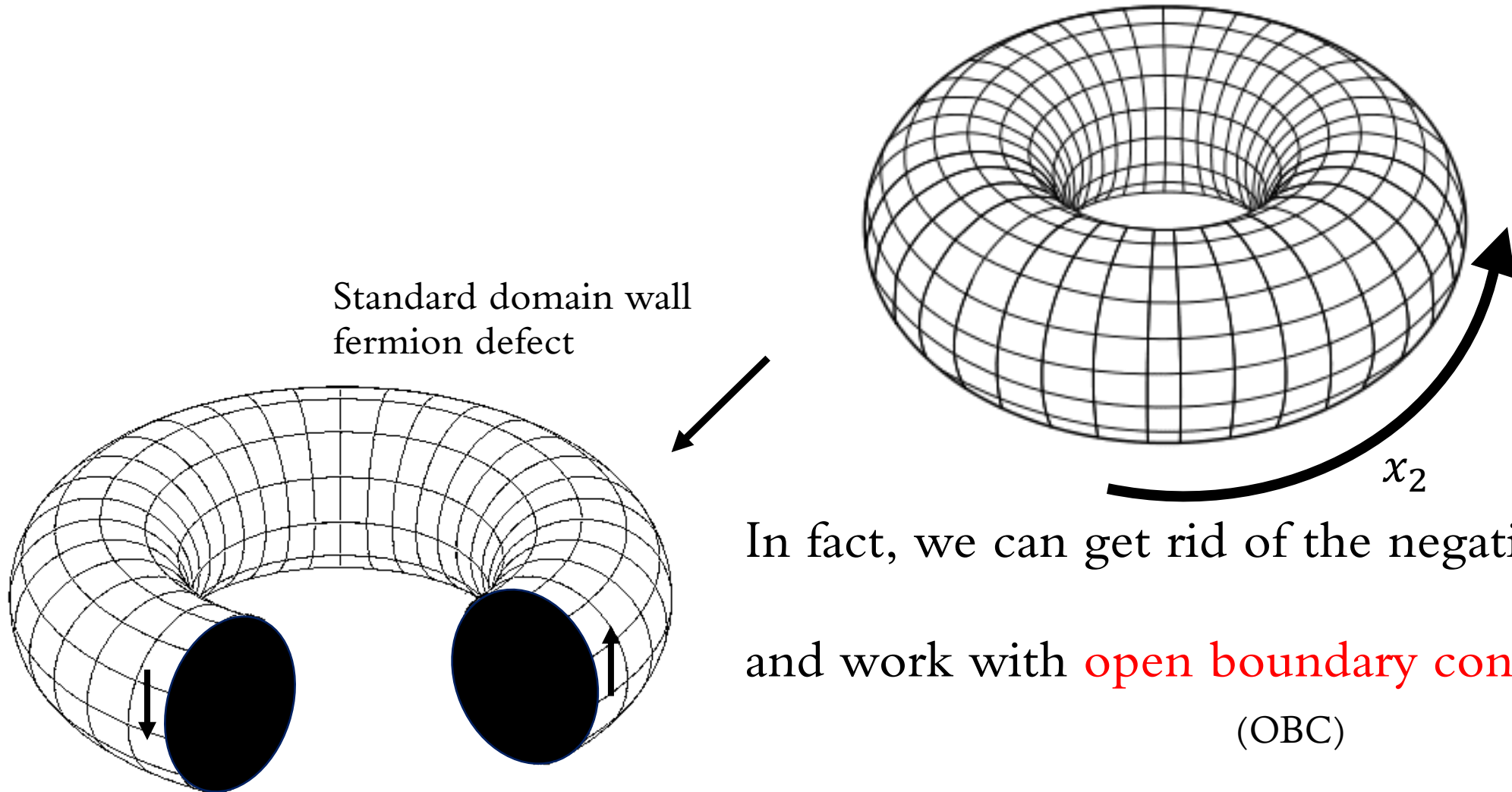
$$E = \pm \sin p$$

Domain wall fermion for global chiral symmetry

Higher dimensional Dirac fermion theory with mass defect



Domain wall fermion



Standard domain wall fermion defect

x_2

In fact, we can get rid of the negative mass region and work with **open boundary condition** in x_2 (OBC)

Towards the spectrum: the DW Hamiltonian

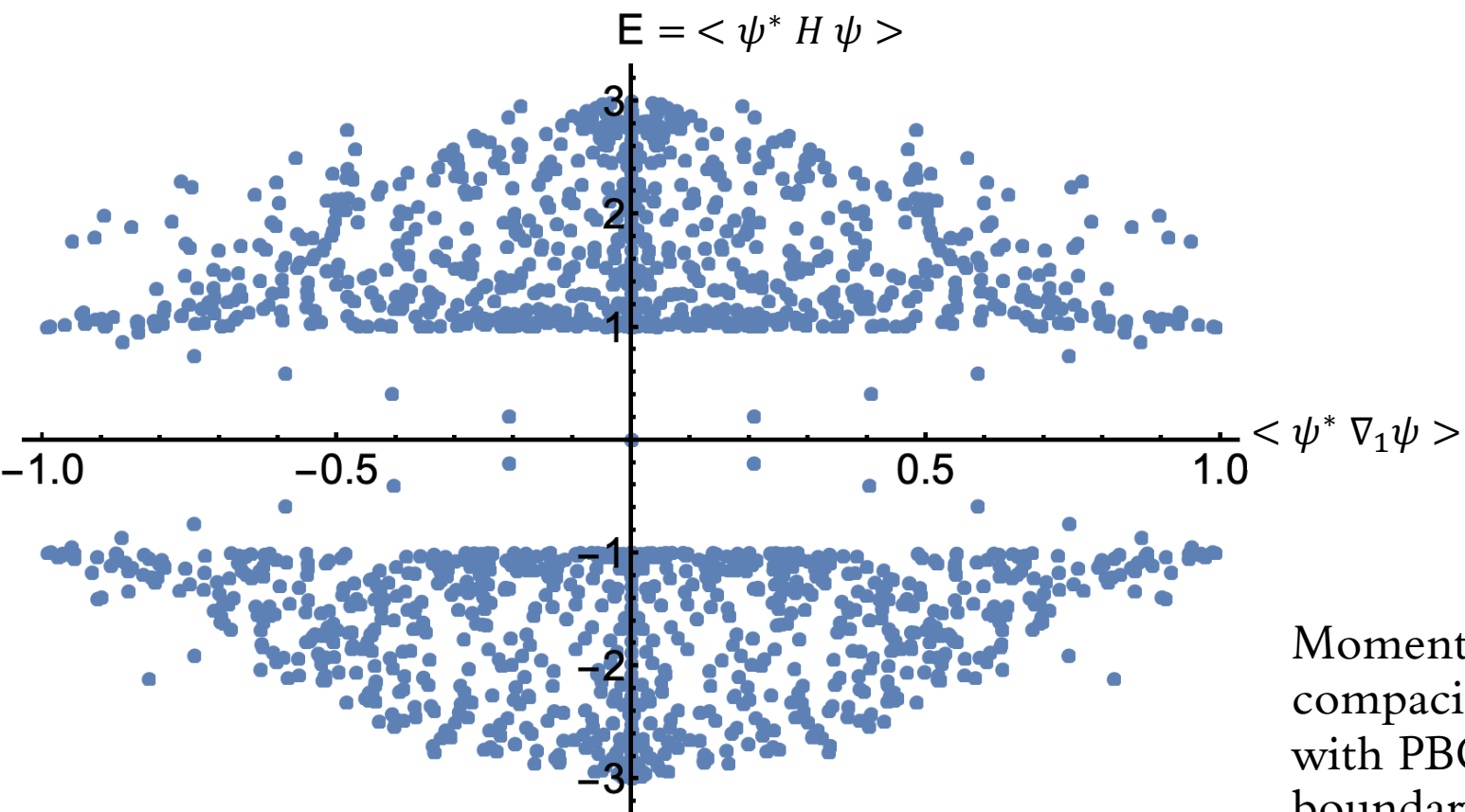
It's the Wilson fermion with no discretization in time.

Single particle Hamiltonian: $H = -i\gamma^i \nabla_i + m + \frac{R}{2} \nabla$

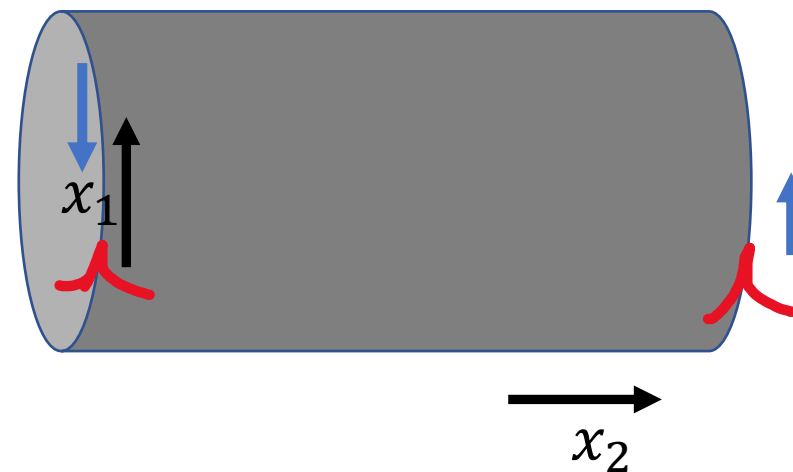
$\nabla_i =$ Symmetric finite difference in space

$\nabla =$ symmetric discrete spatial Laplacian

Spectrum



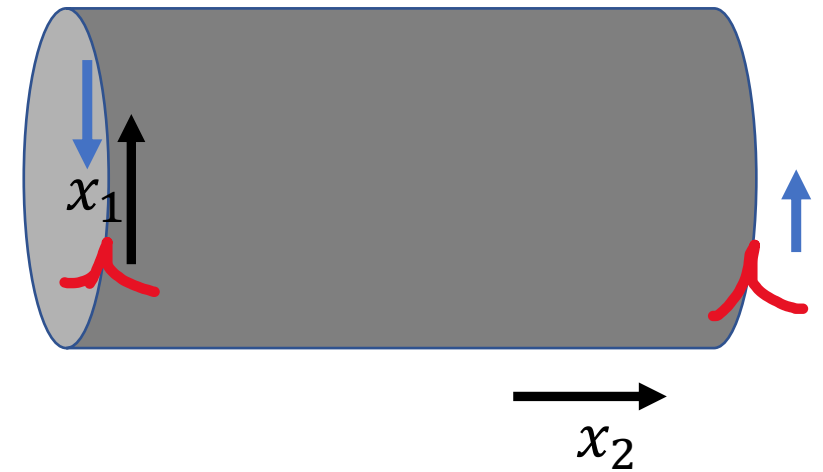
Opposite chiralities



Momentum along the compactified dimension with PBC (periodic boundary condition)

Solved using domain wall fermions

- Right and left moving modes separated in space. So, any quantum correction to mass exponentially suppressed.
- Allow gauge fields to talk to both walls in the same way producing a vector gauge theory.
- Very useful in QCD simulations.



Doesn't work for chiral gauge theories

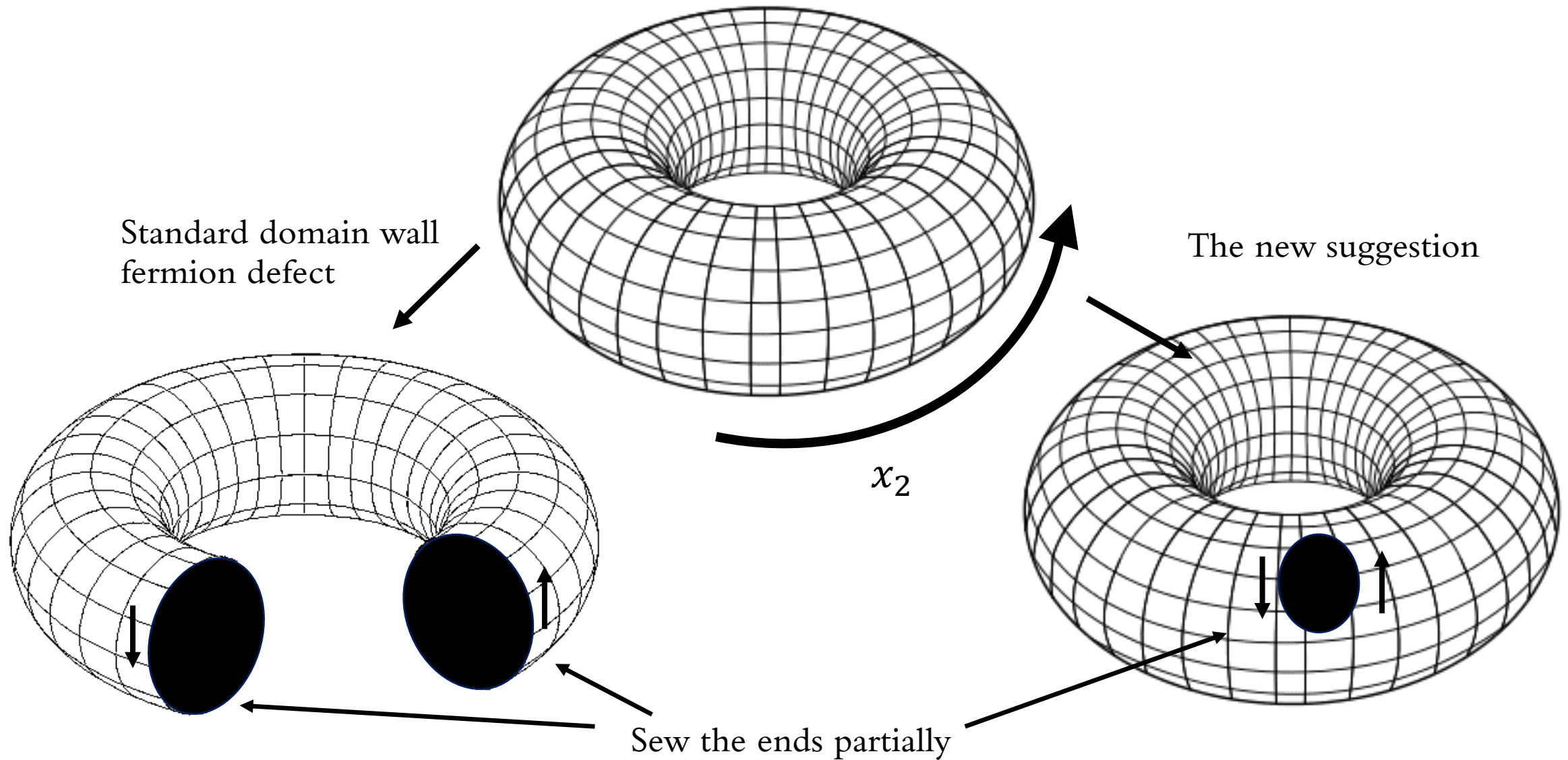
The idea does not work for **chiral gauge theories** though.

The construction in finite volume necessarily has two defects.

Two defects lead to opposite chiralities producing vector theory.

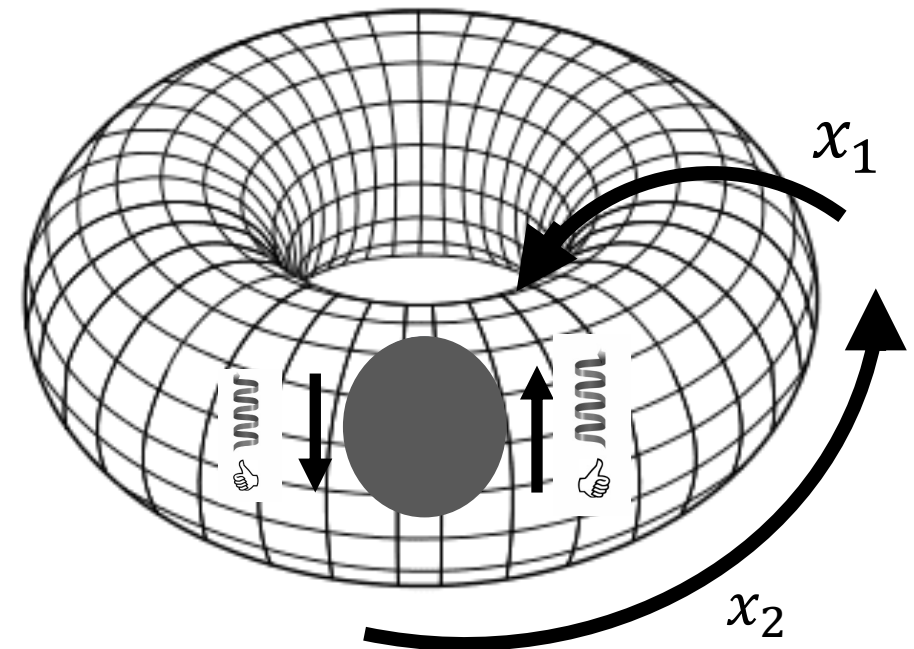
We need to isolate Weyl fermions of a particular chirality --- impossible with the standard domain wall setup.

How about a single disk-like defect?



Opposite chirality on the two sides..

Maybe the problem is that we are keeping the definition of chirality position independent.



Define chirality in a position dependent manner

Define chirality as clockwise travel vs anticlockwise travel:

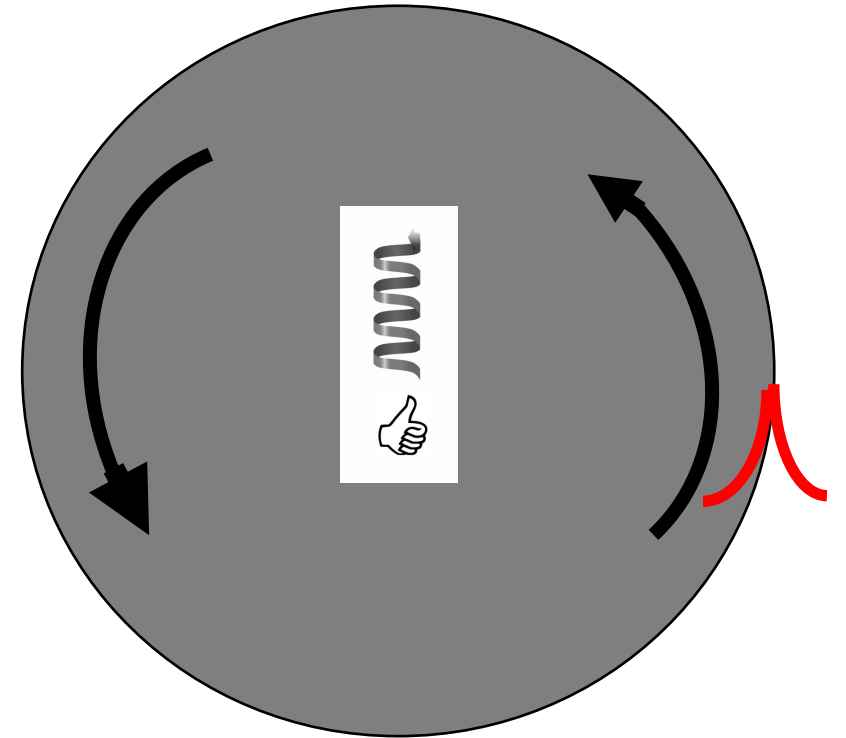
counter-clockwise



clockwise

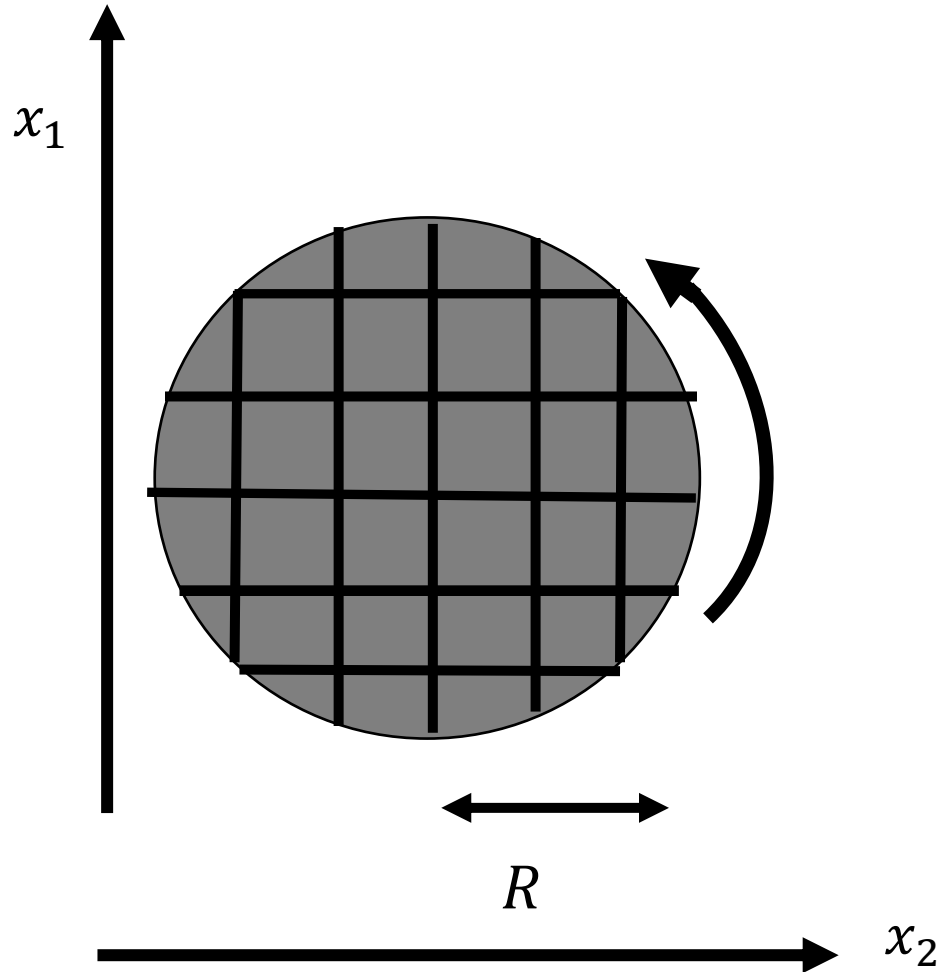


Phys.Rev.Lett. 132 (2024) 14, 141603 (Kaplan)



Single chirality: Weyl mode

Disc



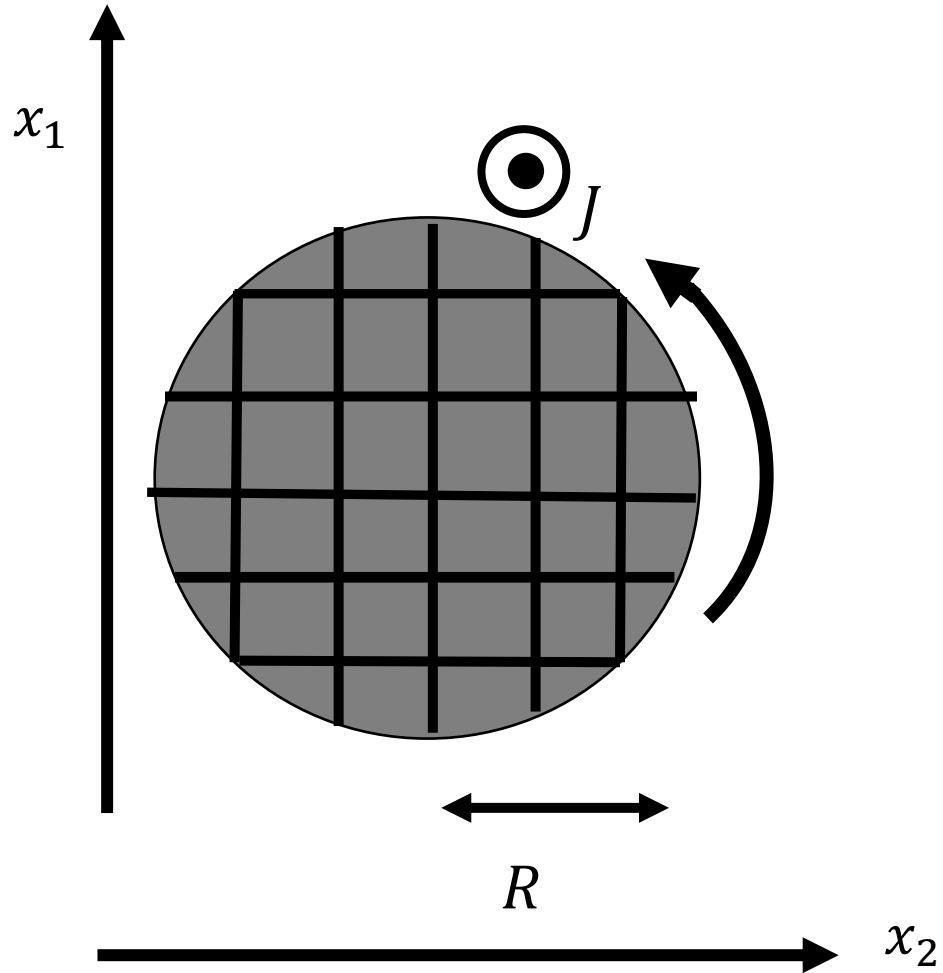
Check the dispersion.

How?

Broken translation invariance
along both x_1 and x_2

Does not make sense to plot
 E vs p_1

Disc



We have rotational invariance (approx).

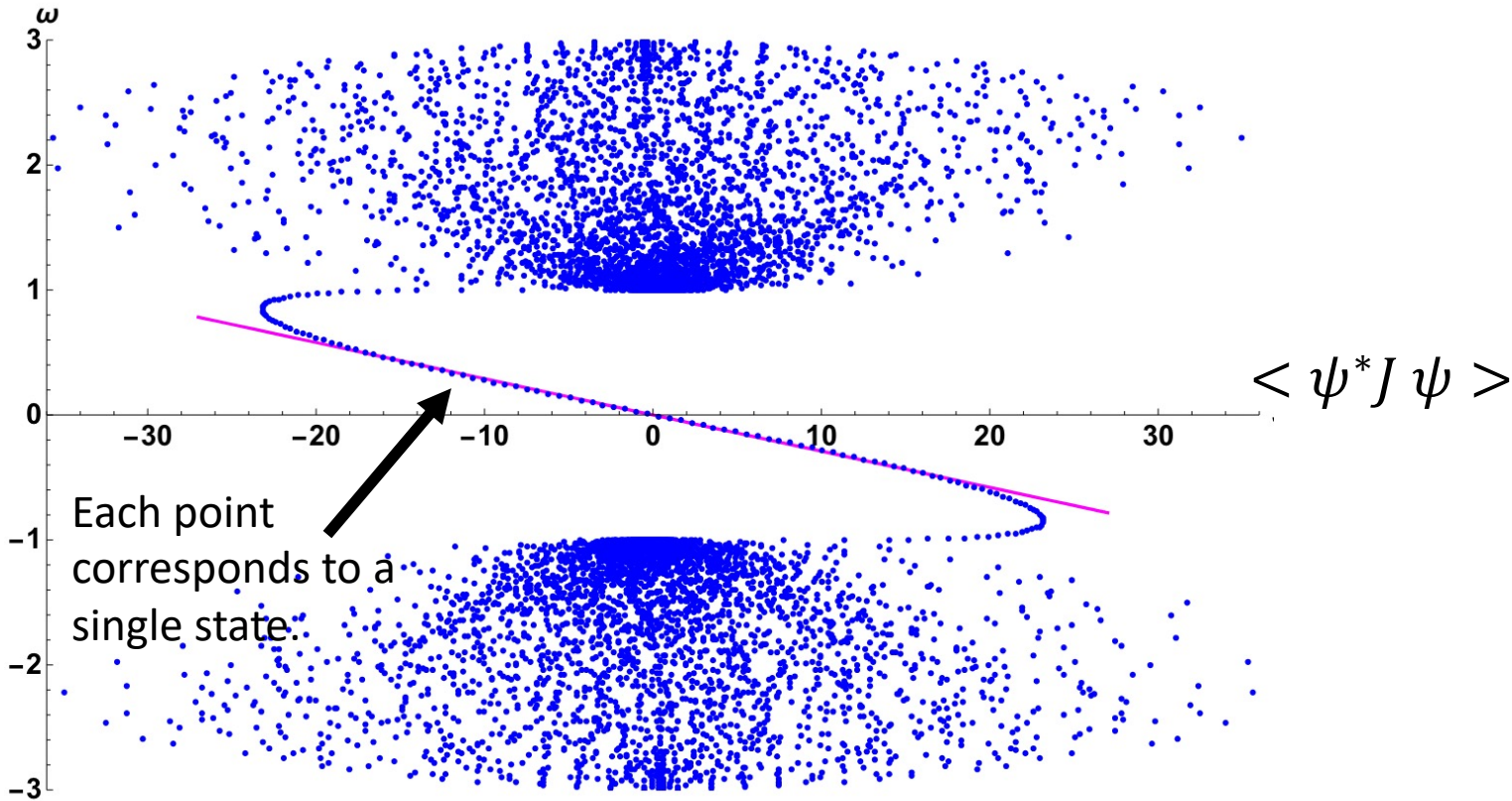
Diagonalize the lattice Hamiltonian.

Compute expectation values of angular momentum J

Plot E vs J

Dispersion for the disk

$$E = \langle \psi^* H \psi \rangle$$

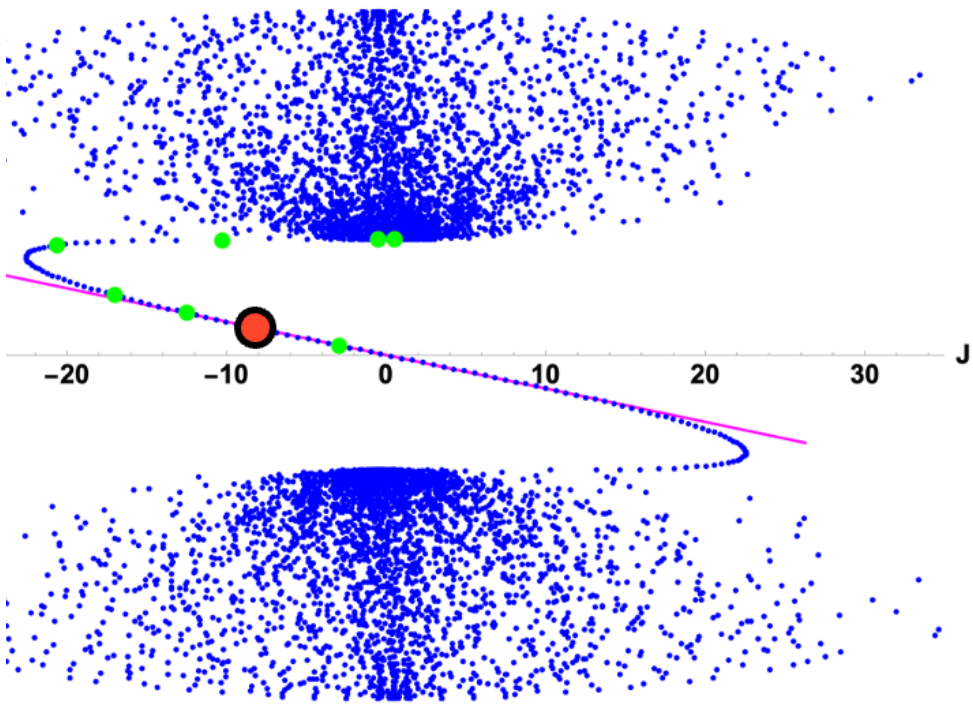


Exactly as expected
from the continuum

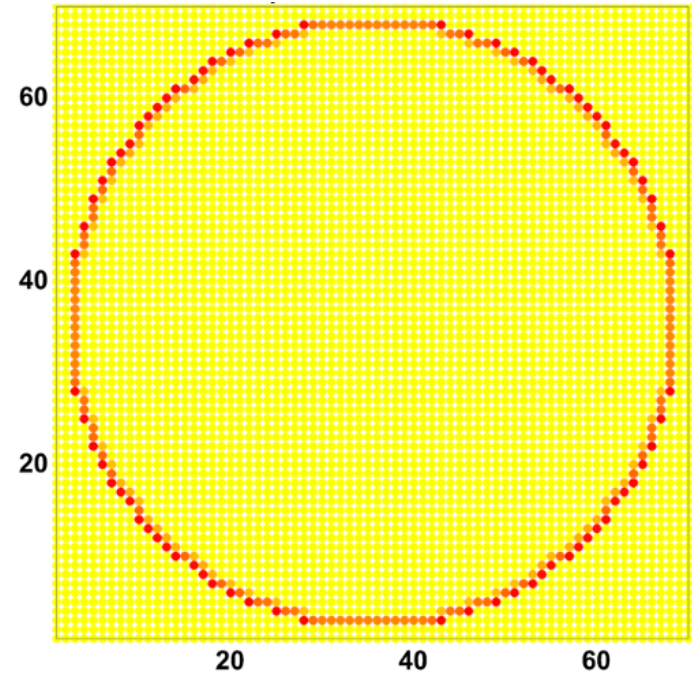
Disk of radius $R = 34$ in
lattice units.

Linear dispersion:

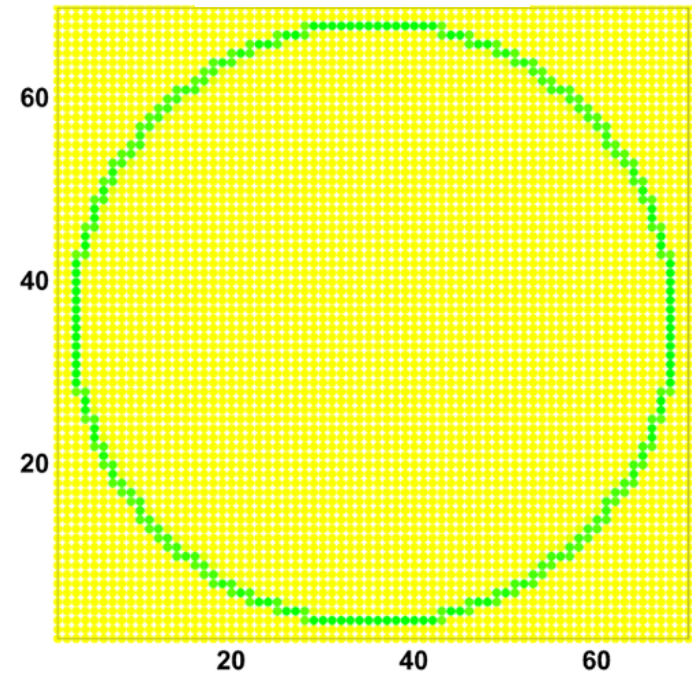
$$E = -J/R$$

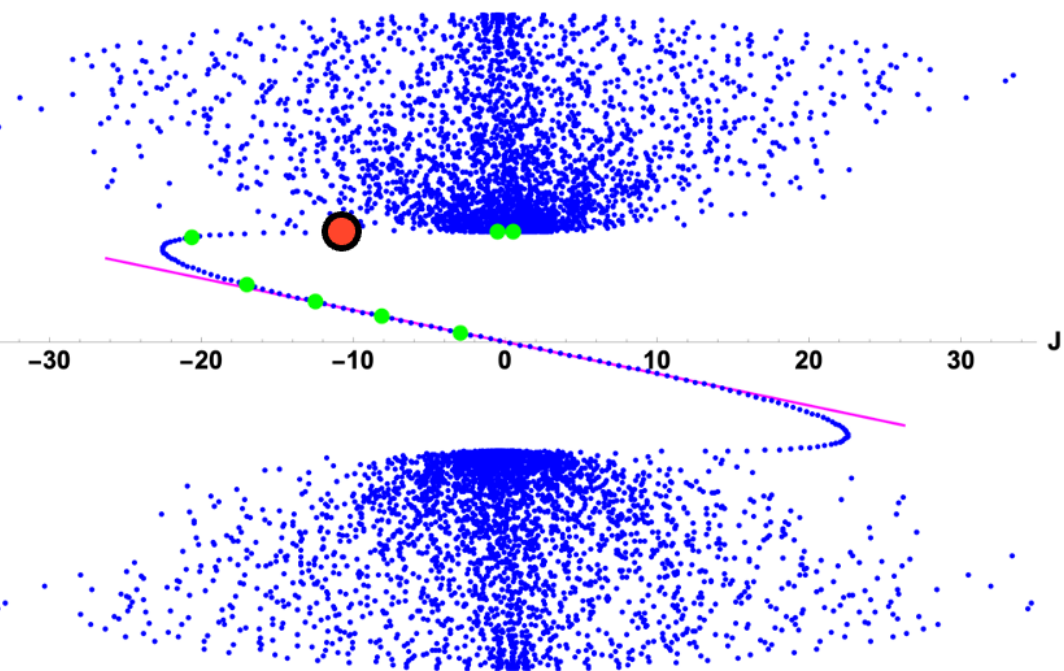


charge density ρ

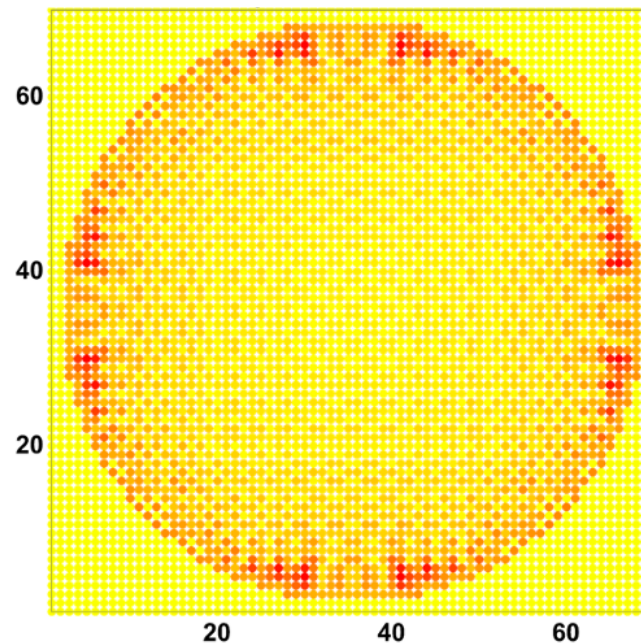


current density j_θ

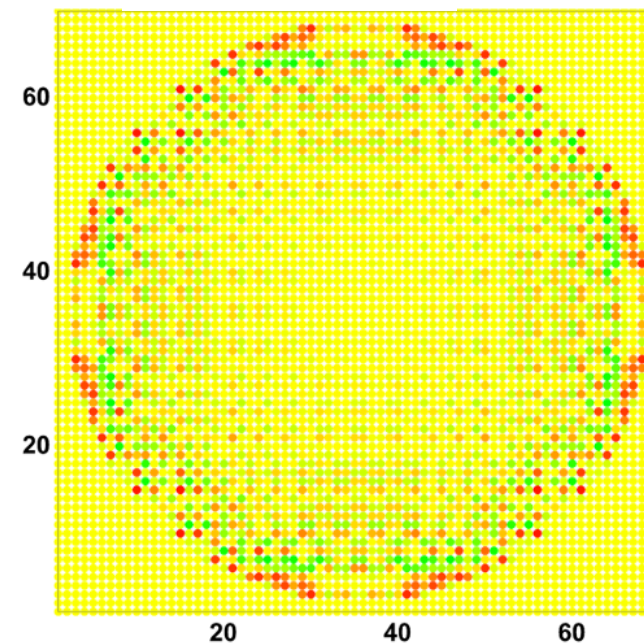




charge density ρ



current density j_θ



Summary

We have a sensible microscopic theory of a Dirac fermion which at low energy produces a single Weyl fermion on the lattice.

Nielsen Ninomiya is not an obstacle. We were fixated on the wrong kind of defect.

Removes one of the most significant obstacles of realizing a chiral gauge theory.

There is more to do though!

Future work

What's the overlap operator for this setup?

How does the latticized version of the overlap operator (lattice boundary theory) realize a Weyl fermion?

Gauge this theory on a small lattice and compute the path integral exactly.

What's the ideal way to simulate this theory? (gauging the full theory or the overlap operator?)