

# Multilevel samplings for glueball calculations

## Collaboration

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Based on

*[arXiv: 2406.12656]*

SU(3) YM

*[Work in progress...]*

Quenched QCD



Lattice 2024, Liverpool 28/07-03/08

# Multilevel samplings for pure gauge glueballs

*[G. Parisi (1983)]*

If action and observables are local, the correlation function

*[M. Lüscher, P. Weisz (2001)]*

*[H. Meyer (2002), (2003)]*

$$\langle O(t_1)O(t_0) \rangle = \frac{1}{\mathcal{Z}} \int [dU] e^{-S_g[U]} O(U, t_1)O(U, t_0)$$

$O(t)$  = glueball operator

can be factorised into a product of integrals

$$\langle O(t_1)O(t_0) \rangle = \frac{1}{\mathcal{Z}} \int [dU_B] e^{-S_B[U_B]} [O^{(2)}(U_B, t_1)] [O^{(1)}(U_B, t_0)]$$

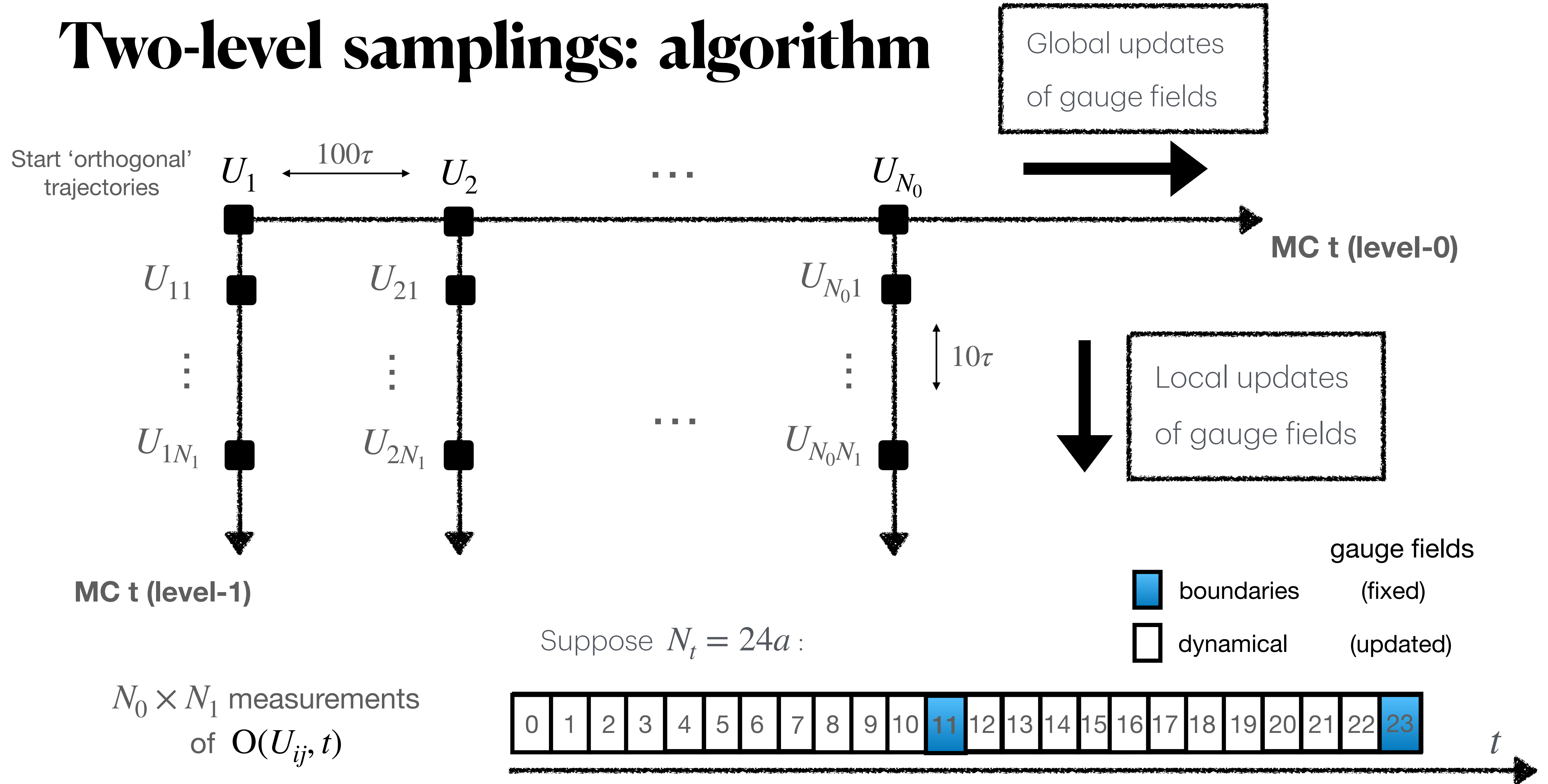
$$[O^{(r)}(U_B, t)] = \int [dU^{(r)}] e^{-S_r[U^{(r)}|U_B]} O(U^{(r)}, t)$$

# Two-level samplings: algorithm



In the same way as with the standard algorithm  
we generate  $N_0$  gauge configurations (HMC)

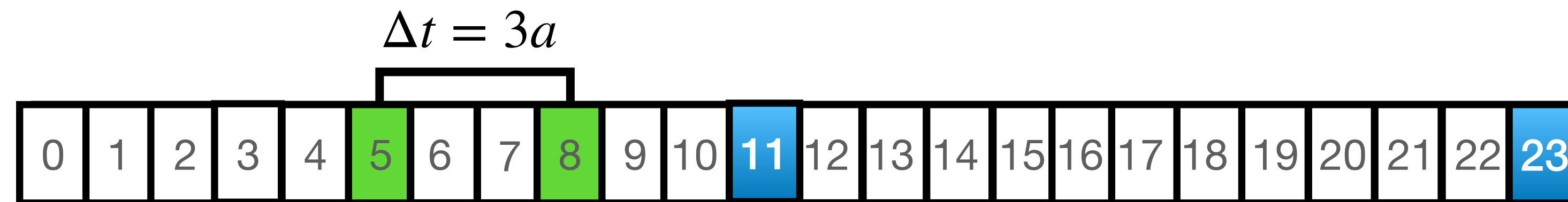
# Two-level samplings: algorithm



# Two-level samplings: algorithm

if  $t_0$  and  $t_1$  belong to the same region

e.g.



Correlate over  $N_0 \times N_1$  samples

$$C(t_1 - t_0) = \langle O(U_{ij}, t_1) O(U_{ij}, t_0) \rangle = \frac{1}{N_0 N_1} \sum_{i=1}^{N_0} \sum_{j=1}^{N_1} O(U_{ij}, t_1) O(U_{ij}, t_0)$$

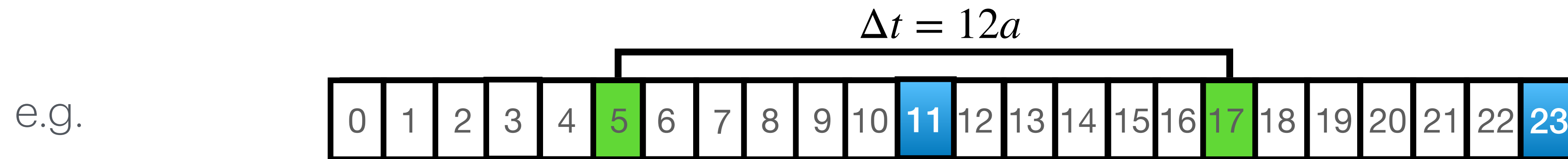
Error scaling

$$\sigma^2(t_1 - t_0) \propto \frac{1}{N_0 N_1} + \dots$$

**Standard**

# Two-level samplings: algorithm

if  $t_0$  and  $t_1$  belong to different regions



Take average over level-1

$$[O(U_i, t)] = \frac{1}{N_1} \sum_{j=1}^{N_1} O(U_{ij}, t)$$

Correlate

$$C(t_1 - t_0) = \langle [O(U_i, t_1)] [O(U_i, t_0)] \rangle = \frac{1}{N_0} \sum_{i=1}^{N_0} [O(U_i, t_1)] [O(U_i, t_0)]$$

Error scaling

$$\sigma^2(t_1 - t_0) \propto \frac{1}{N_0 N_1^2} + \dots \quad \textbf{Two-level}$$

# Numerical simulations

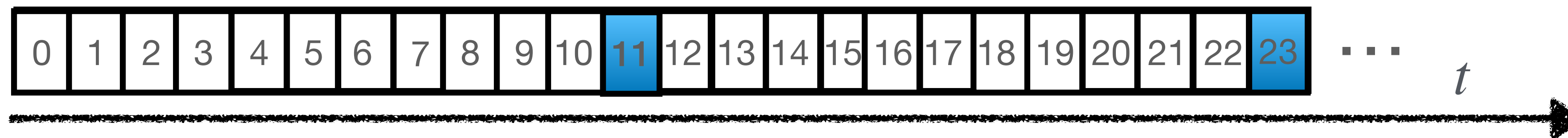
[arXiv: 2406.12656]

*L.B., F. Knechtli, S. Martins, M. Peardon, S. Schaefer, J.A. Urrea-Niño*

In the following:

4D SU(3) theory,  $\beta = 6.2$ ,  $V = 24^3 \times 48$   $N_0 = 101$ ,  $N_1 = 1, \dots, 1000$

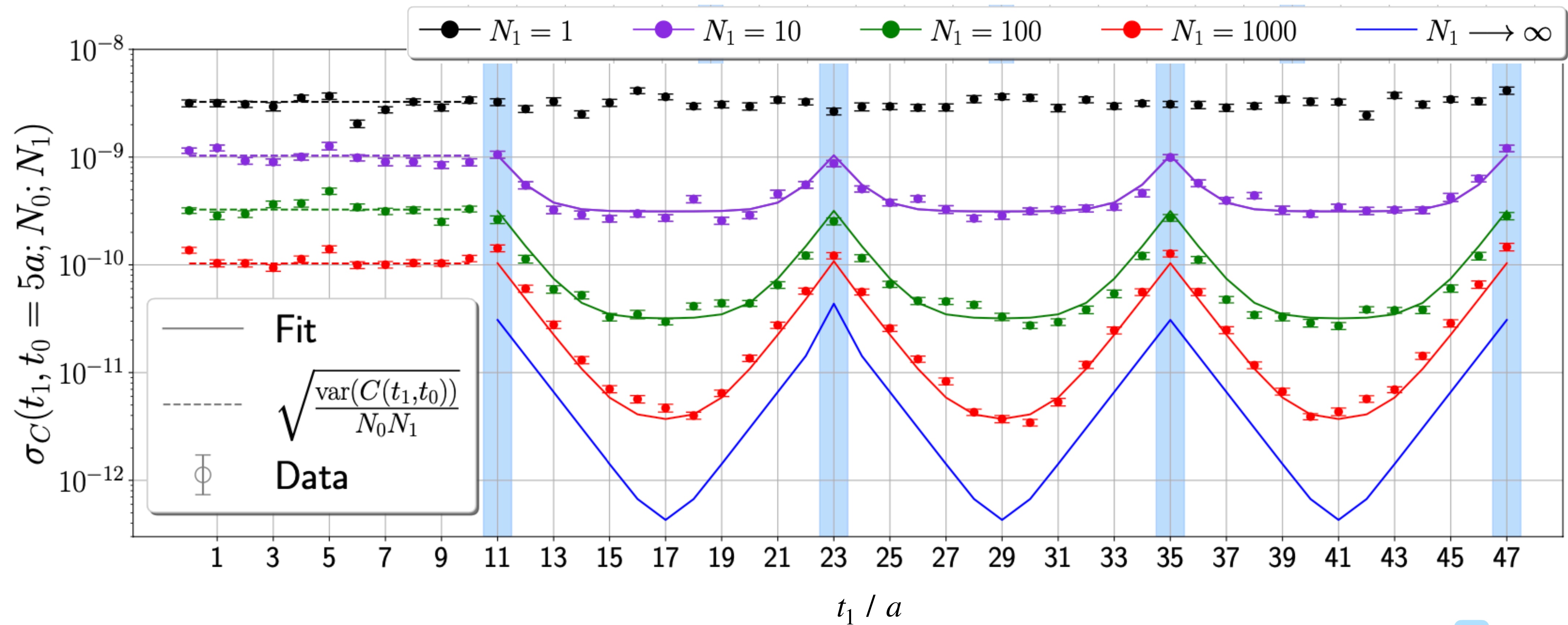
Sub-lattice decomposition on level-1:



$$C(t_1 - t_0) = \langle O(U_{ij}, t_1) O(U_{ij}, t_0) \rangle$$

# Results for the statistical errors

4D SU(3) theory,  $\beta = 6.2$ ,  $V/a^4 = 24^3 \times 48$



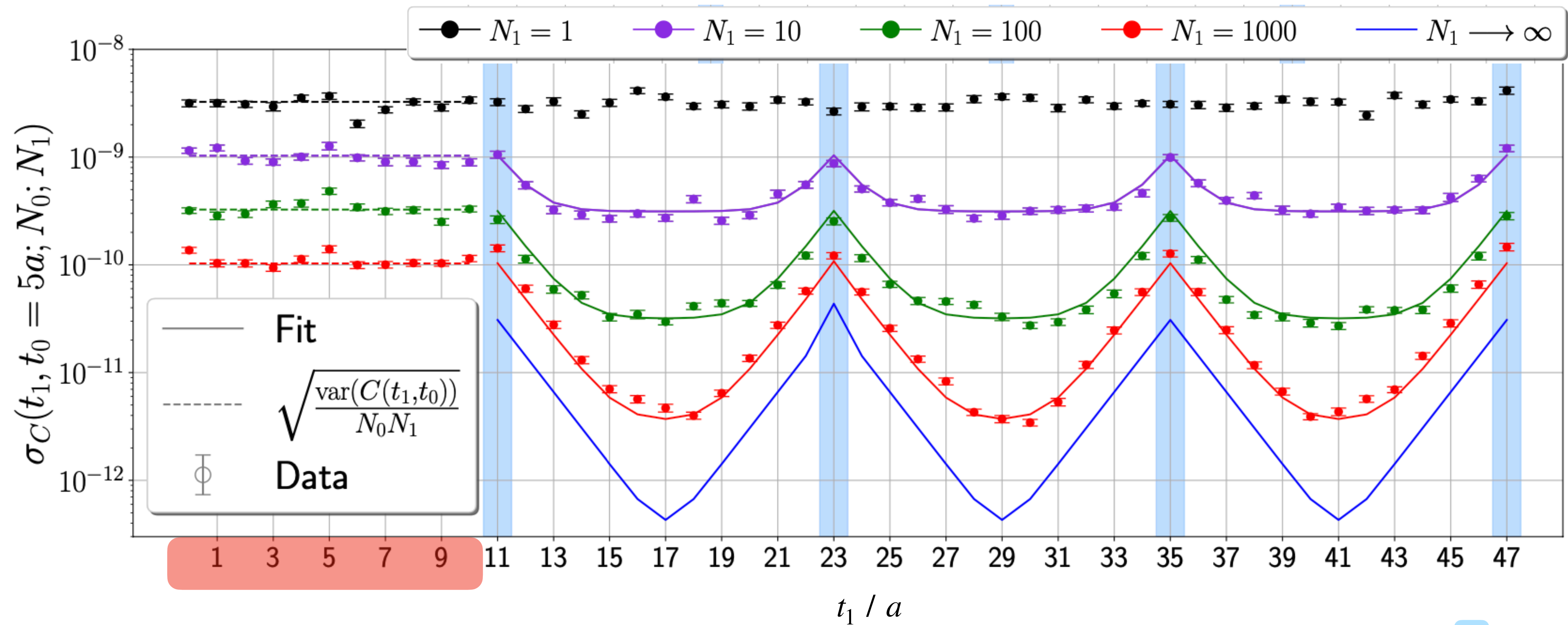
boundaries

$t_B/a = 11, 23, 35, 47$ .



# Results for the statistical errors

4D SU(3) theory,  $\beta = 6.2$ ,  $V/a^4 = 24^3 \times 48$



Error scales like  $1/\sqrt{N_0 N_1}$

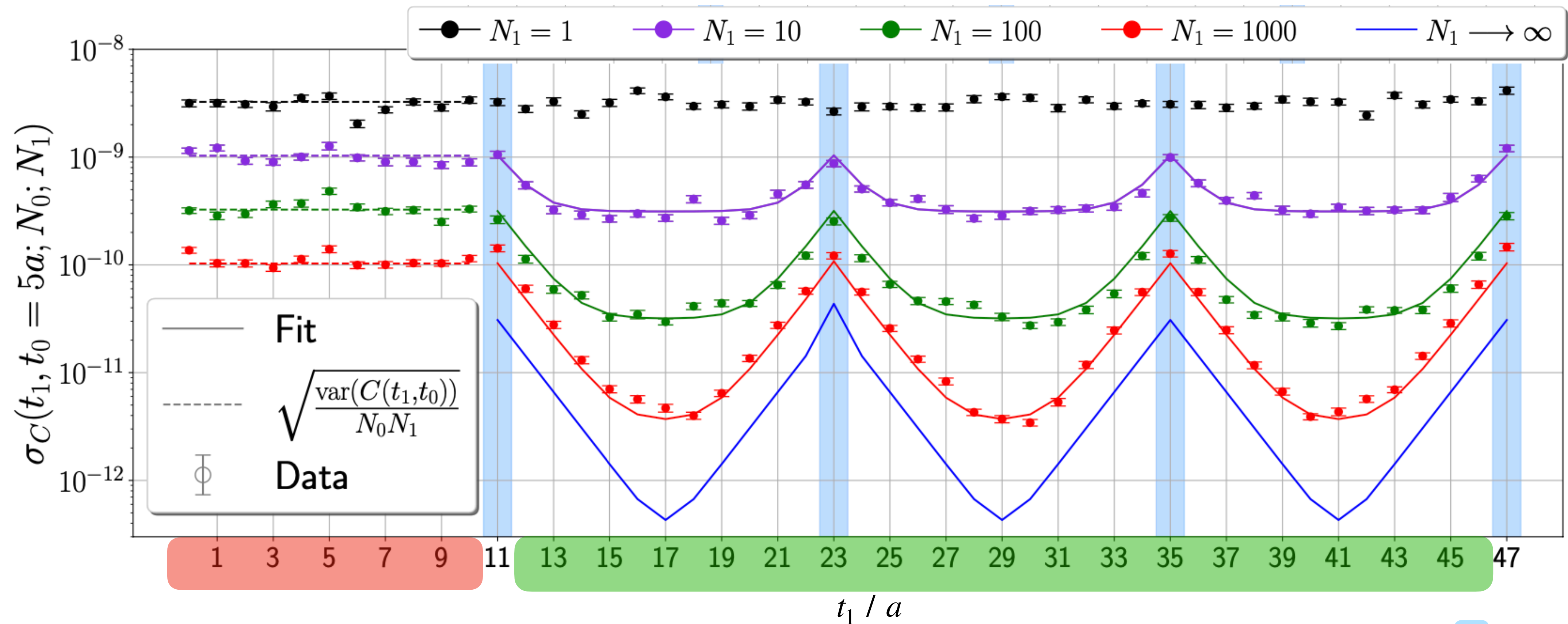
when operators in same region

boundaries

$t_B/a = 11, 23, 35, 47$ .

# Results for the statistical errors

4D SU(3) theory,  $\beta = 6.2$ ,  $V/a^4 = 24^3 \times 48$



Error scales like  $1/\sqrt{N_0 N_1}$

when operators in same region

Error scales exponentially with distance from boundaries

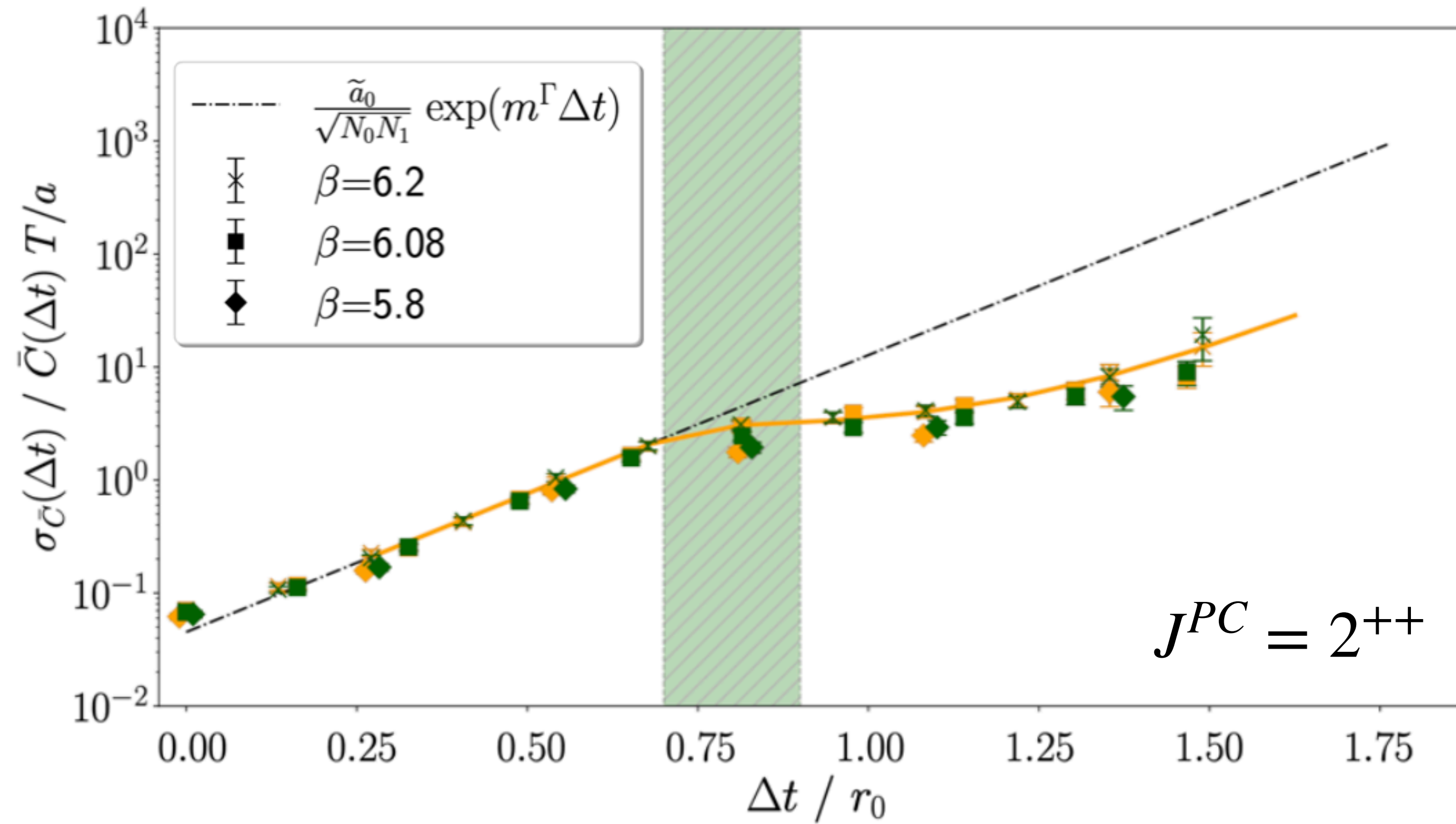
when operators in different regions

$t_B/a = 11, 23, 35, 47$ .

$$\sigma_C^2(t_1, t_0) \approx \frac{c_0^2}{N_0 N_1^2} + \frac{c_1^2}{N_0 N_1} \left[ e^{-m\Delta_1} + e^{-m\Delta_0} \right] + \frac{c_2^2}{N_0} e^{-m\Delta_1} e^{-m\Delta_0}$$

$$\Delta_{1,0} = |t_{1,0} - t_B|$$

# Noise/Signal towards continuum limit



## Weighted average

$$\bar{C}(\Delta t) = \frac{\sum_{t_0} w(t_0) C(\Delta t = t_1 - t_0)}{\sum_{t_0} w(t_0)}$$

$$w(t_0) = 1 / \sigma_C^2(t_0)$$

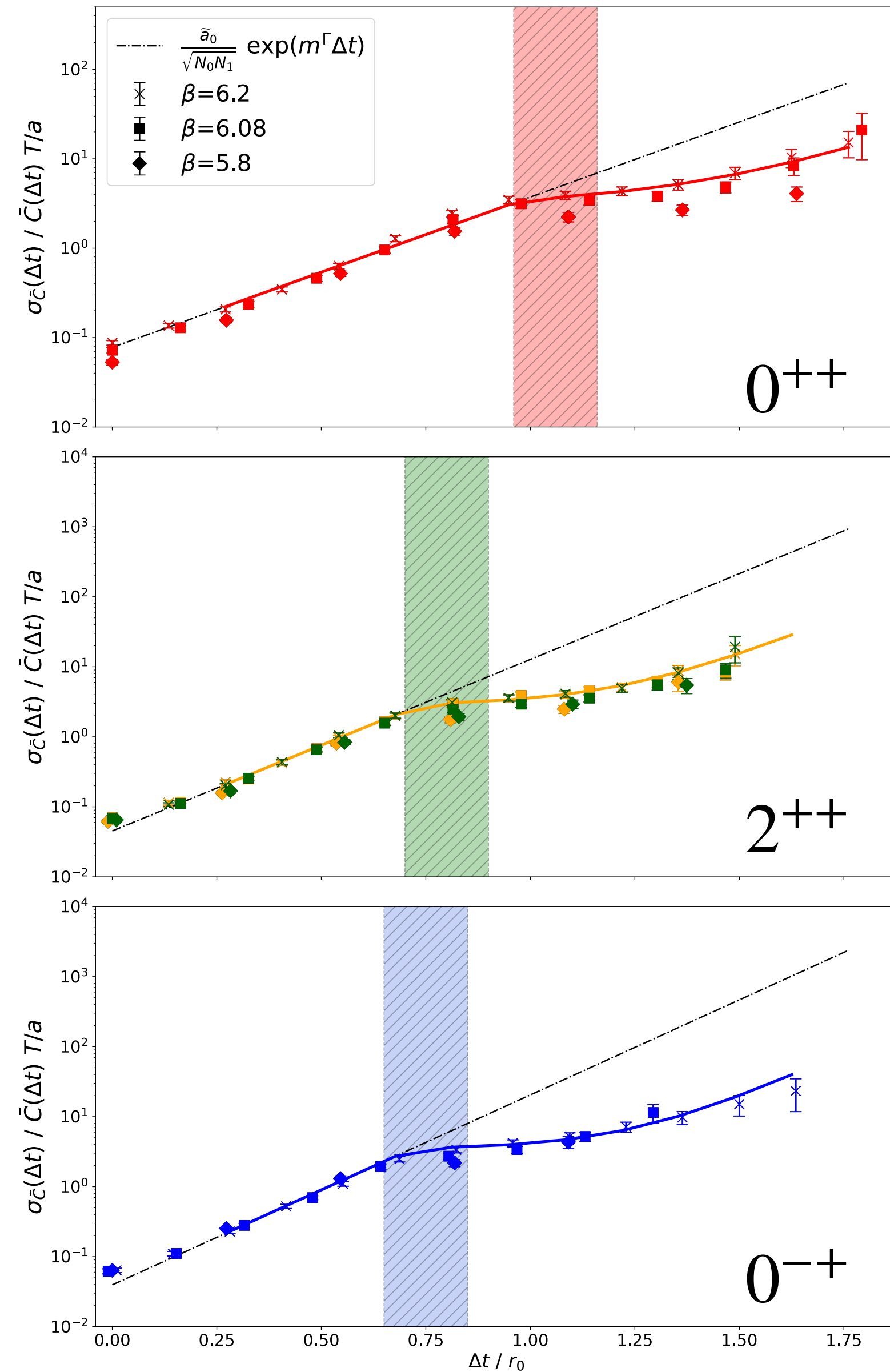
Coloured band highlights transition

between different scalings

◆ Short distance scaling:  $\frac{\sigma_C^2(\Delta t)}{\bar{C}^2(\Delta t)} \approx \tilde{a}_0^2 \frac{e^{2m^\Gamma \Delta t}}{N_0 N_1}$

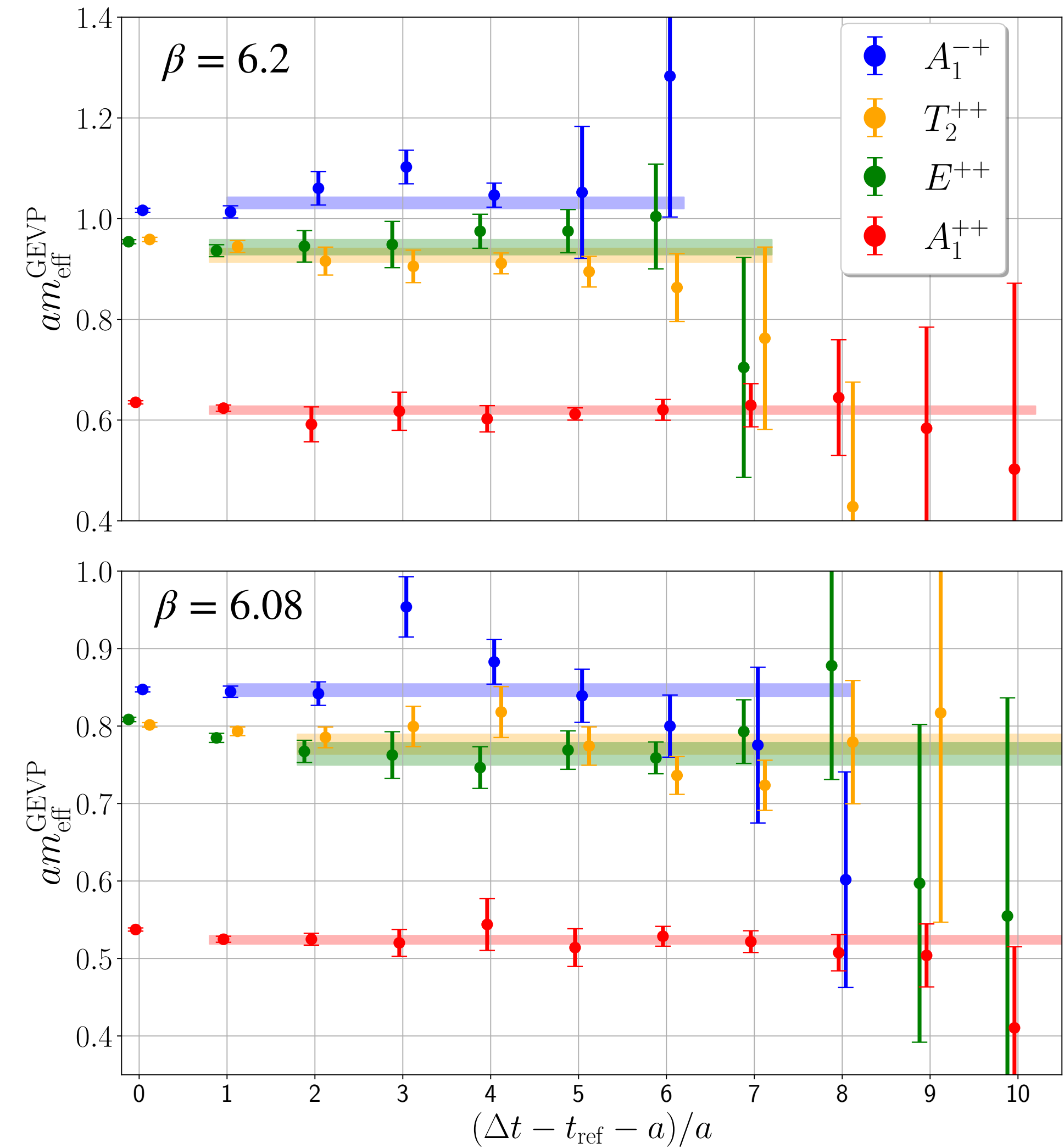
◆ Long distance scaling:  $\frac{\sigma_C^2(\Delta t)}{\bar{C}^2(\Delta t)} \approx \frac{\tilde{c}_0^2}{N_0 N_1^2} e^{2m^\Gamma \Delta t} + \frac{2\tilde{c}_1^2}{N_0 N_1} e^{2m^\Gamma \Delta t/2} + \frac{\tilde{c}_2^2}{N_0}$

# Noise/Signal towards continuum limit



Consistency across  
all glueball channels

# GEVP effective masses



Results agree with state-of-the-art calculations that use  $\mathcal{O}(10^5)$  configs

[A. Athenodoros, M. Teper, 2020]

# Glueballs in Quenched QCD

$$\langle O(t_1)O(t_0) \rangle = \frac{1}{\mathcal{Z}} \int [dU][d\bar{q}][dq] e^{-S_g[U]} O(U, t_1) O^\dagger(U, t_0)$$

$O(U, t)$  are now purely gluonic and **fermionic** operators like  $O_{\pi\pi}(U, t)$  (Action is still SU(3))

Scalar glueballs candidates e.g.  $f_0(1500)$ ,  $f_0(1730)$  decay into multi-meson states like  $f_0(1500) \rightarrow \pi\pi$

Need to compute

$$\langle O_{\pi\pi}(t_1) \bar{O}_{\pi\pi}(t_0) \rangle = \frac{1}{\mathcal{Z}} \int [dU] e^{-S[U]} \langle O_{\pi\pi}(U, t_1) O_{\pi\pi}^\dagger(U, t_0) \rangle_{\text{Wick}}$$

$$\langle O_{\pi\pi}(U, t_1) O_{\pi\pi}^\dagger(U, t_0) \rangle_{\text{Wick}} = \text{Tr} [D^{-1}(t_1, t_1) D^{-1}(t_1, t_1)^\dagger] \text{Tr} [D^{-1}(t_0, t_0) D^{-1}(t_0, t_0)^\dagger] + \dots$$

Disconnected diagram



[J. A. Urrea-Niño talk, Tuesday]

# Two-level Samplings for Glueballs in Quenched QCD

## Challenge:

$$\text{Tr} \left[ D^{-1}(t, t) D^{-1}(t, t)^\dagger \right] \quad \text{Propagators depend on gauge fields over all space-time}$$

## Solution:

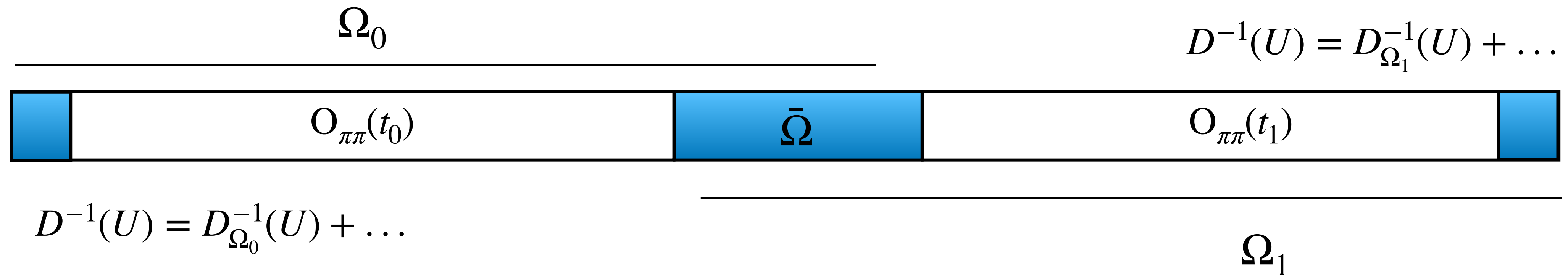
- Full factorisation of propagators in different regions allow multilevel integration

*[arXiv:1601.04587]*

L. Giusti, M. Cè, S. Schaefer, 2016

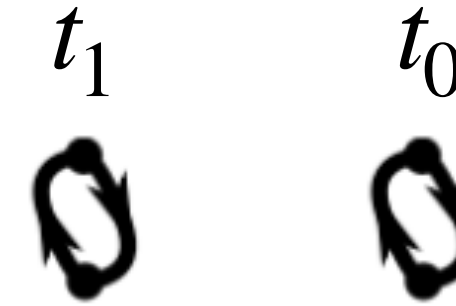
## New approach:

- Propagators are numerically approximated on sub-domains of the full temporal extent



# Two-level Samplings for Glueballs in Quenched QCD

This approximation allows to factorise the disconnected diagram



$$\langle \mathbf{O}_{\pi\pi}(t_1) \bar{\mathbf{O}}_{\pi\pi}(t_0) \rangle_{disc.} = \frac{1}{\mathcal{Z}} \int [dU] e^{-S[U]} \text{Tr} \left[ \mathbf{D}_{\Omega_1}^{-1}(t_1, t_1) \mathbf{D}_{\Omega_1}^{-1}(t_1, t_1)^\dagger \right] \text{Tr} \left[ \mathbf{D}_{\Omega_0}^{-1}(t_0, t_0) \mathbf{D}_{\Omega_0}^{-1}(t_0, t_0)^\dagger \right] + \text{correction}$$

Now amenable for multilevel integration:

$$= \frac{1}{\mathcal{Z}} \int [dU_B] e^{-S_B[U_B]} \left[ \mathbf{O}_{\pi\pi}^{(1)}(U_B, t_1) \right]_{\Omega_1} \left[ \mathbf{O}_{\pi\pi}^{(0)}(U_B, t_0) \right]_{\Omega_0} + \text{correction}$$

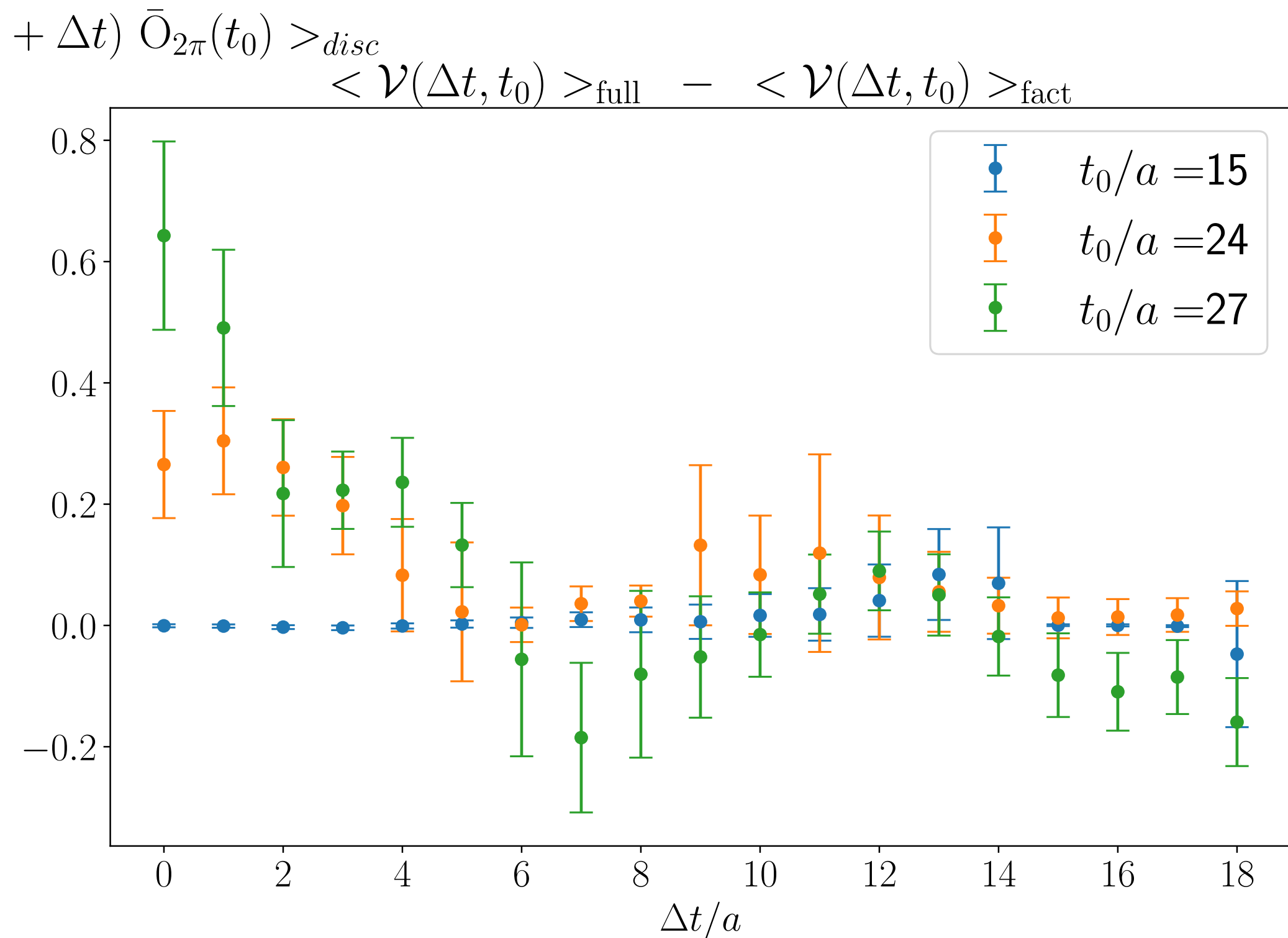
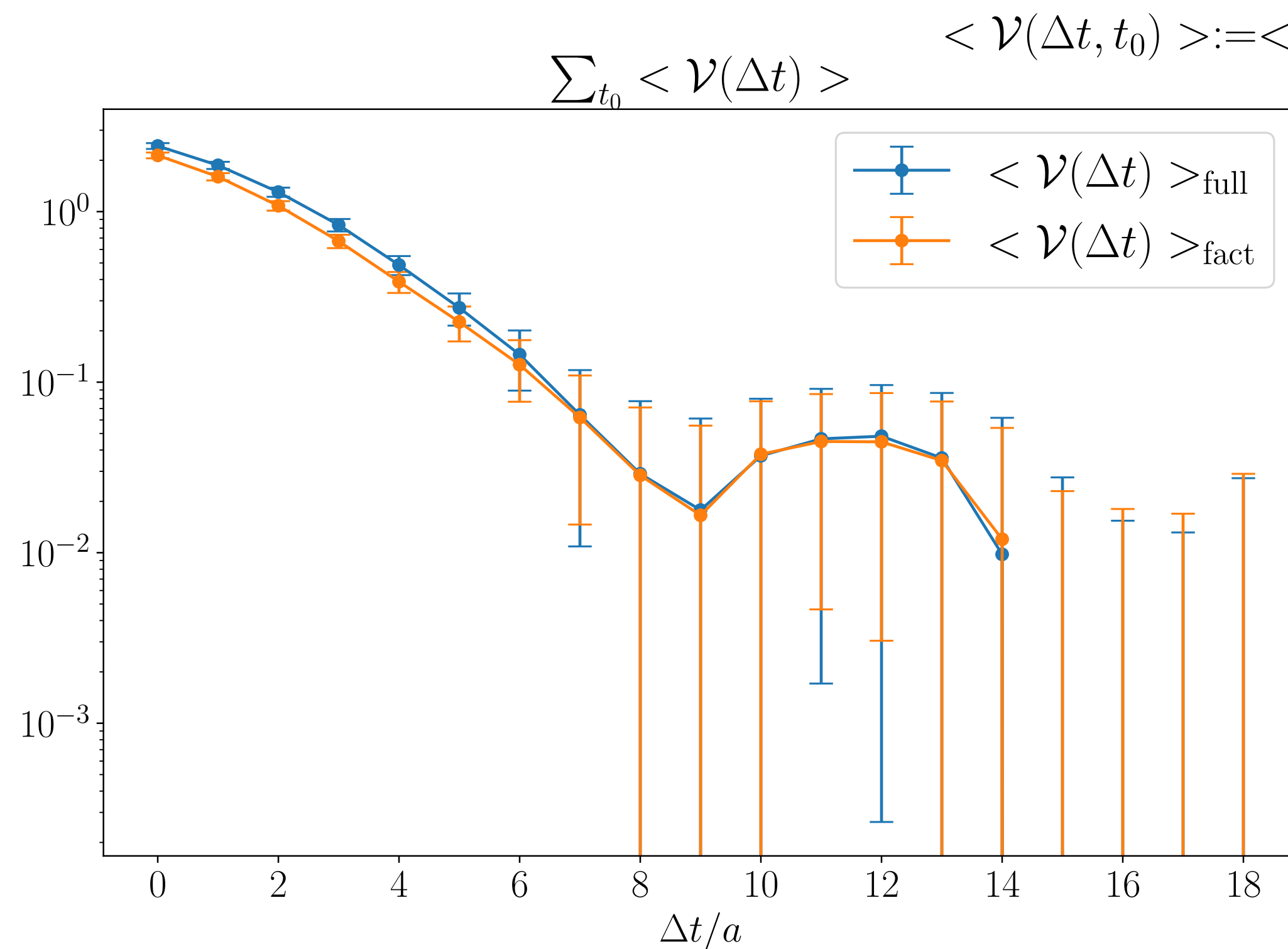
$$\left[ \mathbf{O}_{\pi\pi}^{(r)}(U_B, t) \right] = \int [dU^{(r)}] e^{-S_r[U^{(r)}, U_B]} \text{Tr} \left[ \mathbf{D}_{\Omega_r}^{-1}(t, t) \mathbf{D}_{\Omega_r}^{-1}(t, t)^\dagger \right] \quad (\text{in practice we use distillation})$$

# How good is the approximation?

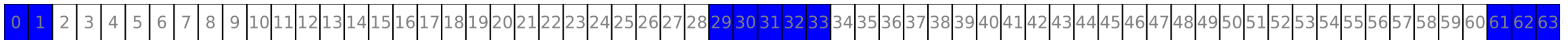
Longer the distance from the other region, better the approximation

Matches well our use case

(multilevel for longer distances)



current sub-lattice decomposition





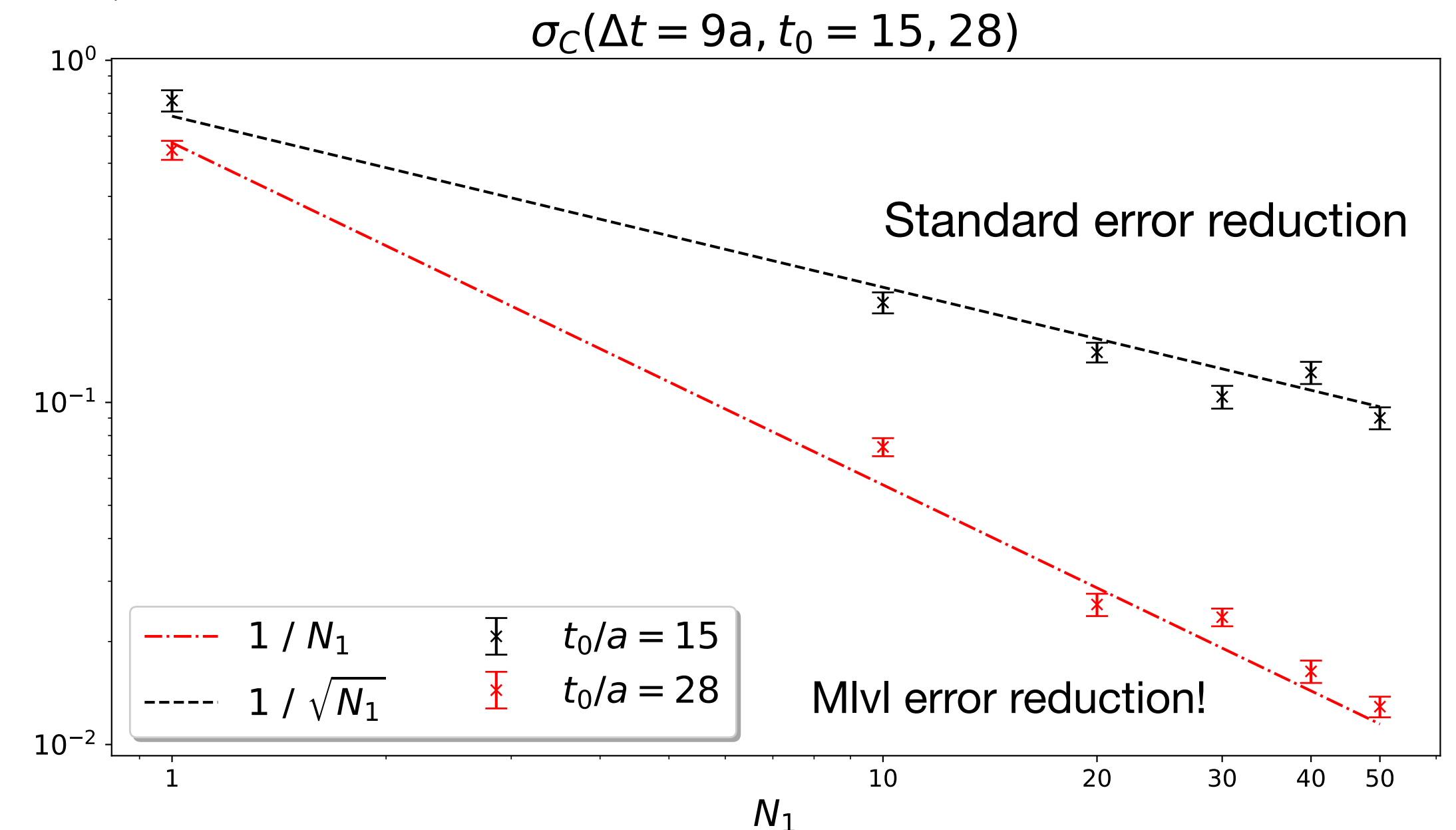
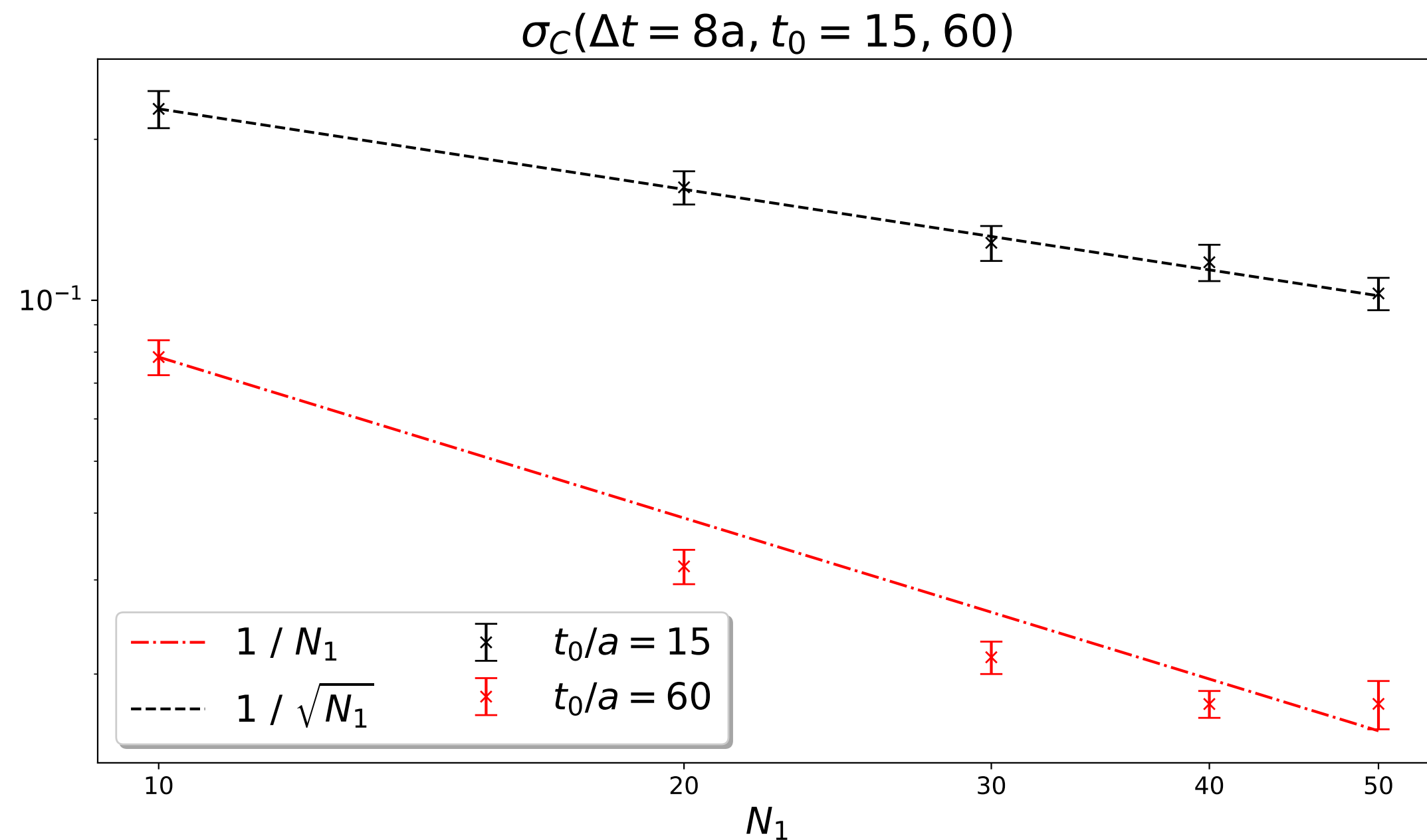
# Multilevel Error Reduction

$$m_\pi \approx 800 \text{ MeV} \quad \text{s.t.} \quad 2m_\pi \approx m_{f_0(1500)}$$

$$V/a^4 = 16^3 \times 64 \quad N_0 = 20, \quad N_1 = 50$$

$$\sigma_C(\Delta t, t_0) = \text{Error of } \langle O_{\pi\pi}(t_0 + \Delta t) \bar{O}_{\pi\pi}(t_0) \rangle$$

## Preliminary



Error  $\propto 1 / N_1$  when the operators are located in different regions!

# Conclusions

[arXiv: 2406.12656]

L.B., F. Knechtli, S. Martins, M. Peardon, S. Schaefer, J.A. Urrea-Niño

- ✓ Two-level sampling for pure gauge glueballs very effective and well understood
  - ✓ Fit ansatz describes short and long distance scalings, as well as transition point
  - ✓ GEVP effective masses agree with literature

- 📌 Two-level sampling for glueballs in quenched QCD
  - ★ Approximating Dirac inversions to local domains enables multilevel measurements
  - ✓ Preliminary results show multilevel error reduction for  $\langle O_{2\pi}(t)\bar{O}_{2\pi}(t_0) \rangle$ ,  $\langle O_{2\pi}(t)W(t_0) \rangle$
  - ✓ Approximation improves with distance between operators



📌 Next:

L.B., F. Knechtli, J. Finkenrath, M. Peardon, S. Schaefer, J.A. Urrea-Niño

- Study dependence of mlvl error reduction & approximation with sub-lattice decomposition,  $N_1$ ,  $m_\pi$
- Solve GEVP using both fermionic and gluonic operators (First quenched results ever afaik)