THE QED CONTRIBUTIONS TO THE SHORT AND INTERMEDIATE WINDOWS OF THE HADRONIC VACUUM POLARIZATION CONTRIBUTION TO THE MUON  $g-2$ . Alexei Bazavov<sup>1</sup>, Christine Davies<sup>1</sup>, David Clarke<sup>1</sup>, Carleton DeTar<sup>1</sup>, Aida El-Khadra<sup>1</sup>, Anthony Grebe<sup>1</sup>, Steven Gottlieb<sup>1</sup>, Andrew Lytle<sup>1</sup>, William Jay<sup>1</sup>, Hwancheol Jeong<sup>1</sup>, Andreas Kronfeld<sup>1</sup>, Shaun Lahert<sup>1</sup>, Peter Lepage<sup>1</sup>, Andrew Lytle<sup>1</sup>, Michael Lynch<sup>1</sup>, **Craig McNeile**<sup>1,2</sup>, Ethan Neil<sup>1</sup>, Curtis Peterson<sup>1</sup>, Gaurav Ray<sup>1</sup>, James Simone<sup>1</sup>, Jake Sitison<sup>1</sup>, Ruth Van de Water<sup>1</sup>, and Alejandro Vaquero<sup>1</sup> <sup>1</sup>Fermilab Lattice, HPQCD and MILC collaborations 2 craig.mcneile@plymouth.ac.uk

• The exciting recent results from the Fermilab Muon g-2 experiment for the Muon Anomalous Magnetic Moment (2104.03281) motivates reducing the error on lattice calculations

- of the hadronic contribution to  $a_{\mu}^{LO}$  $\frac{LO}{\mu}$  .
- The lattice QCD results from the BMW collaboration (2002.12347, 2407.10913) for  $a_{\mu}^{LO}$  $\mu$ are in tension with the data driven estimates.
- We have reported lattice-QCD calculations of the light-quark connected contribution to window observables associated with the leading-order hadronic vacuum polarization contribution to the anomalous magnetic moment of the muon (2207.04765,2301.08274). Here we present the connected QED contributions to window observables to further compare to data driven estimates.

BMW estimate that the connected QED contributions to  $a_{\mu}^{LO}$  $_{\mu}^{LO}$  are  $\sim 0.2$  %.

#### Formalism for computig  $a$  $HVP(LO)$  $\overline{\mu}$

### Introduction

 $K_G^w(t) \equiv K_G(t)W(t, t_i, \Delta t).$ The short distance (SD) window is (T. Blum et al, arXiv:1801.07224)

 $W^{SD}(t;t_1) \equiv 1 - \Theta(t;t_1,\Delta)$ 

with  $t_1 = 0.4$  fm. This regime may be described by perturbation theory. The intermediate window (W) is defined by

 $W(t; t_1, t_2) \equiv \Theta(t; t_1, \Delta) - \Theta(t; t_2, \Delta)$ 

with  $t_1 = 0.4$  fm and  $t_2 = 1.0$  fm. This is a standard benchmark number for comparison of lattice and phenomenological numbers.

The Windows are defined (with  $\Delta = 0.15$  fm) using

 $\Theta(t;t',\delta) \equiv$ 1 2  $+$ 1 2  $\tanh[\frac{t-t'}{2}]$  $\Delta$ ]

We estimate the QED contribution to the strange or light quarks via

 $\delta a_{\mu} = a_{\mu} [QCD + QED] - a_{\mu} [QCD]$ 

We have **not** yet retuned the quark masses to include the QED contribution.

- We use the electro-quenched approximation and only compute connected correlators.
- The calculation used quenched QED fields fixed to the Feynman gauge with zero modes dealt with using the  $QED<sub>L</sub>$  prescription.

Fig. 5: Light short distance Fig. 6: Light intermediate window For the short distance window we use the continuum limit model  $\delta a^l_{\iota}$  $\mu_\mu^l = a_0(1 + a_1 \alpha_s (\mu =$ 2  $\overline{a}$  $)(\Lambda a)^2 + a_2(\Lambda a)^4$ 

Ensembles at physical pion masses, but because of noise increasing we use  $3/5/7m_l$  valence quark masses (following BMW) and extrapolate to  $m_l$  for vector mesons. Lower statistics at physical point for pseudo-scalar mesons to test scheme dependence.





Lattice-QCD calculations of the HVP are based on the Euclidean time vector-vector correlation

$$
G_{ff'}(t) = Q_f Q_{f'} \sum_{\vec{x}} Z_V^2 \langle j_f^i(\vec{x}, t) j_{f'}^i(0) \rangle.
$$
 (1)

where f and  $f'$  are flavour indices,  $Q_f$  is the electric charge for that flavour in units of  $e$ ,  $Z_V$  is the renormalisation factor for the lattice vector (electromagnetic) current. The contribution to  $a_{\mu}$  from a window (that isolates a region in time) is

> with fit parameters  $a_0$ ,  $a_1$  and  $a_2$ .  $\alpha_s$  in the V scheme, and  $\Lambda = 1.2$  GeV. For the W window we use

$$
a^w_\mu(t_1,\Delta t) = \left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dt \, G_{ff'}(t) \; K_G^w(t) \,,
$$

with a modified kernel,

 $\delta a_\mu^s$  $\frac{s}{\mu}=a_0(1+a_1\alpha_s(\mu=$ 2  $\overline{a}$  $)(\Lambda a)^2 + a_2(\Lambda a)^4$ with fit parameters  $a_0$ ,  $a_1$  and  $a_2$ .  $\alpha_s$  in the V scheme and  $\Lambda = 0.5$  GeV.

- We next plan to compute the QED contributions to the disconnected diagrams.
- We are investigating using the perturbative approach (1706.05293) to estimate the QED corrections.
- We will then compute the remaining QED corrections.

## Connected quenched QED corrections

The basis of the calculation are MILC HISQ  $nf = 2 + 1 + 1$  QCD gauge-field ensembles.

Connected strange quenched QED corrections

We have computed the SD and W contribution to the QED contribution to  $a_{\mu}^{LO}$  $\frac{LO}{\mu}$  .



# Connected light quenched QED corrections

We extrapolate the light quark data from quark masses  $7m_l$ ,  $5m_l$ , and  $3m_l$  to the physical light quark mass  $m_l$ .



$$
\delta a_{\mu}^{l} = a_0 + a_1 \alpha_s (\mu = \frac{2}{a})(\Lambda a)^2 + a_2(\Lambda a)
$$

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because  $a_0$  is small.



• We use the truncated solver method with 16 sloppy inversions and 1 precise inversion. Ensemble  $\mid L^3 \times T \mid a[\text{fm}] \mid \text{no. meas} \mid$  masses very coarse  $32^3 \times 48$  0.15 1844  $m_u m_d$  3/5/7 $m_l m_s$ coarse  $|48^3 \times 64|$  0.12 967  $|3/5/7m_l m_s|$ 

fine  $|64^3 \times 96|0.09|$  596  $|3/5/7m_l m_s|$ 

## Future work