# **Smeared** *R***-ratio in isoQCD with Low Mode Averaging**

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Low Mode Averaging (LMA) is a technique to increase the quality of the signal-to-noise ratio in the long time separation of Euclidean correlators. It is highly beneficial in computing the vector-vector light connected two-point correlators and derived physical quantities, e.g.  $a_{\mu}^{\rm HVP}$  $_{\mu}^{\rm HVP}, \,\, {\rm in \,\, the \,\, ETMC}$ mixed action lattice setup [1]. Here, we focus on preliminary results for the smeared *R*-ratio [2] using the Hansen-Lupo-Tantalo (HLT) spectral-density reconstruction method [3], now with Gaussian kernels of width down to  $\sigma \sim 250$  MeV. This is enough to appreciate the  $\rho$  resonance around 770 MeV.

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### **Introduction**

#### *R***-ratio via HLT method (mini review)**  $\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{L}^{\text{max}}(\mathcal{$

$$
\sigma(e^+e^- \to \mu^+\mu^-) \qquad \qquad e^- \nearrow \qquad \qquad \qquad
$$



#### γ **hard (a) (b)**

#### $R(\omega) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{(1 + \omega^2 + \omega^2)^2}$  $\mu(\omega) = \frac{\partial (e^+ e^- \to \text{hadrons})}{\partial (e^+ e^- \to \mu^+ \mu^-)}$   $\lambda \sim \sqrt{2}$  hadrons

 $\frac{1}{2}$  $\frac{1}{2}$  st-t<sup>1</sup> and equal leads to  $\frac{1}{2}$  **Fig. 5.9** and the measurement of the  $\frac{1}{2}$  the  $\frac{1}{2}$  cm section by KLOE at the φ–factory  $\frac{1}{2}$  and  $\frac{1}{2}$  at the  $\frac{1}{2}$  cm section by KLOE at the  $\frac{1}{2}$  and  $\frac{1}{2}$  and  $\frac{1}{2}$  are the  $R(\omega)$  is a distribution and has to be probed by using smearing kernels. We extract the smeared *R*-ratio

 $\int_{-\infty}^{\infty}$  $R_{\sigma}(E) = \int d\omega G_{\sigma}(E-\omega)R(\omega)$ , with normalized Gaussian kernels  $\mathcal{L}_{\mathbf{0}}$  $R_{\sigma}(E) = \int$ ∞ 0  $d\omega G_{\sigma}(E-\omega)R(\omega)$ , with normalized Gaussian kernels, √

at  $\alpha$   $\alpha$   $\beta$   $\beta$   $\beta$   $\beta$   $\beta$   $\beta$   $\beta$  $G_{\sigma}(E - \omega) = \exp(-(E - \omega)^2/2\sigma^2)/2$  $2\pi\sigma^2$  .

for the R-rational by and the gradieric current is connected to the R-ratio by Is connected to the It-ratio by is connected to the *R*-ratio by

 $\frac{1}{2}$  with an accuracy measured by with an accuracy measured by

> $\overline{\Omega}$ *s*−2 weig √*s* de a<br>
> de space. The g coefficients result from minimizi spond √*s* that, for weight-functions  $w_n > 0$ , correspond to a class of weighted *L*2-norms in functional space. The **g** coefficients result from minimizing

 $\Delta$  is a parameter varied/optimized in stability analysis, see [2]  $\alpha$ upp. *λ* is a parameter varied/optimized in stability analysis, see [2] Supp. Mat. **Figure 2:** The frequency histogram of the eigenvalues. As the frequency of  $\lambda_{\min} \propto L^3 T$ , we scale *K* roughly in this way while varying *a*: typically *K* is such that  $\lambda_{\text{max}}/\mu_{ud} \sim 8$ .

where  $r = \pm 1$  and  $\mu$  are Wilson and twisted mass parameters. The two-point functions read



al<br>**Rigure 1:** The first principle lattice results of 2023 [2] compared to the smoares **Figure 1:** The first-principle lattice results of 2023 [2] compared to the smeared experimental *R*-ratio from KNT19 compilation [4]. Errors increase as  $\sigma$  decreases.

level [59, 60] at the end. The first dedicated radiative return experiment has been per-Two-point Euclidean correlator of the quark electromagnetic current *C*(*t*)

$$
C(t) = -\frac{1}{3} \sum_{i=1}^{3} \int d^{3}x T \langle 0 | J_{i}(x)J_{i}(0) | 0 \rangle = \frac{1}{12\pi^{2}} \int_{2m_{\pi}}^{\infty} d\omega \, e^{-\omega t} \, \omega^{2} R(\omega) .
$$
  
In [3] the smearing kernels are approximated as  $K(\omega; \mathbf{g}) = \sum_{\tau=1}^{T/a} g_{\tau} e^{-a\omega \tau}$ 

**Figure 3:** Left: vector-vector light connected two-point correlator on B64 ensemble at  $L = 5.1$  fm using old [9] and new LMA setup. Right: the gain in the signal-to-noise ratio.

$$
A_{\rm n}[g] = \int_{E_0}^{\infty} d\omega w_{\rm n}(\omega) \left| K(\omega; g) - \frac{12\pi^2 G_{\sigma}(E - \omega)}{\omega^2} \right|^2,
$$

$$
W[\lambda, \mathbf{g}] = (1 - \lambda) A_{n}[\mathbf{g}] + \lambda B[\mathbf{g}], \text{ where } B[\mathbf{g}] = \sum_{\tau_{1,2}=1}^{T/a} g_{\tau_1} g_{\tau_2} \text{Cov}_{\tau_1 \tau_2},
$$

*τ*=1



## **Low Mode Averaging**

$$
D_r = D_W + ir\mu\gamma_5 \ , \quad Q_r \equiv \gamma_5 D_r = Q_W + ir\mu \ , \quad S_r = (D_r)^{-1} = Q_r^{-1}\gamma_5 \ ,
$$

$$
C_{rr'}^{AB}(t) = \frac{1}{L^3} \text{Tr} \left[ \sum_{\vec{x}, \vec{y}} S_r (\vec{x}, t_0; \vec{y}, t_0 + t) \Gamma_A S_{r'} (\vec{y}, t_0 + t; \vec{x}, t_0) \Gamma_B \right].
$$

Following  $[5-7]$ , we split the propagator as:

$$
S_r(x,y) = |P_{\text{IR}}Q_r^{-1}P_{\text{IR}}\eta(x)\rangle\langle\eta(y)|\gamma_5 + |P_{\text{UV}}Q_r^{-1}P_{\text{UV}}\eta(x)\rangle\langle\eta(y)|\gamma_5
$$
  
= 
$$
\sum_{j=1}^K \frac{|v_j(x)\rangle\langle v_j(y)|\gamma_5}{\lambda_j + ir\mu} + \frac{1}{N}\sum_{\eta} |\widetilde{\phi}_r^{\eta}(x)\rangle\langle\eta(y)|\gamma_5|_{N\gg 1},
$$

*K* and *N* count the lowest modes and the stochastic sources respectively.



**Table 1:** ETMC ensembles [8] used for two-point vector correlators in isoQCD below.



### *R***-ratio via HLT** + **LMA Preliminary**



**Figure 4:** Smeared *R*-ratio with Gaussian kernels of width  $\sigma = 250$  MeV and central energy up to 1.6 GeV.

# **References**

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