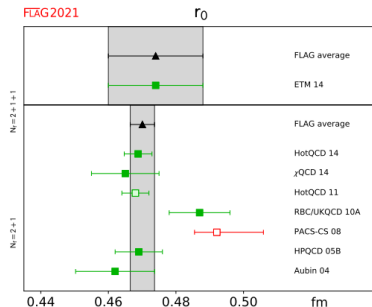


The scales r_0 & r_1 in $N_f = 2 + 1$ QCD

Tom Matty Bo Asmussen
with Roman Höllwieser, Francesco Knechtli, and Tomasz
Korzec

Motivation

- ▶ r_0 as reference scale
 - Weak quark mass dependence
 - Reasonable statistical precision
 - Mild lattice artifacts
- ▶ **Goal:** To add to existing measurements with state of the art simulated data



[FLAG Review 2021, 2111.09849]

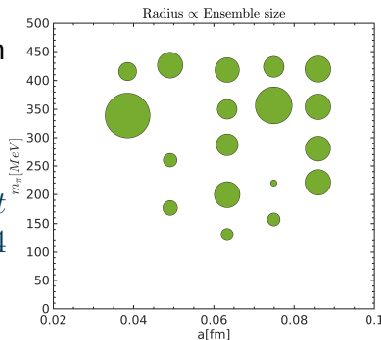


Coordinated Lattice Simulations ($N_f = 2 + 1$)

“Simulation of QCD with $N_f = 2 + 1$ flavors of non-perturbatively improved Wilson fermions”

[M.Bruno et al, (2014) 1411.3982]

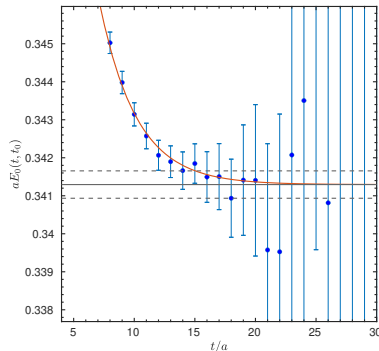
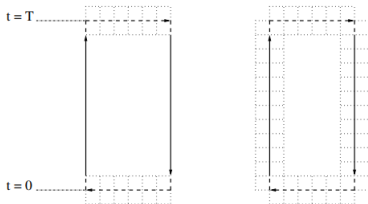
- ▶ 2+1 flavors of $O(a)$ improved Wilson fermions using Lüscher-Weisz gauge action with open boundary conditions in time.
- ▶ Chiral trajectory with $tr(M) = m_u + m_d + m_s = const$ while making sure that $m_\pi L > 4$
- ▶ 5 lattice spacings from 0.085 fm to 0.037 fm
- ▶ m_π from 430 MeV to 134 MeV



Wilson loops

“Determination of the Static Potential with Dynamical Fermions” [M.Donnellan, F.Knechtli, B.Leder, R.Sommer, 1012.3037]

- ▶ Operator basis with different levels of HYP smearing.



- ▶ HYP2 static quark action



Finding r_0 and r_1

We are using an improved

definition of the force:

$$F(r_I) = [V(r) - V(r - a)]/a$$

where $r_I = r - a/2 + \mathcal{O}(a^2)$

[S. Necco and R. Sommer, [hep-lat/0108008](#)]

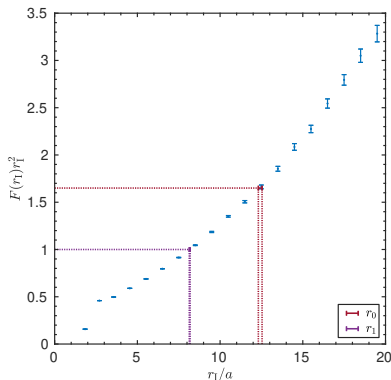
Using that, we find the following scales:

$$r_I^2 F(r_I)|_{r_I=r_0} = 1.65$$

[R. Sommer, [hep-lat/9310022](#)]

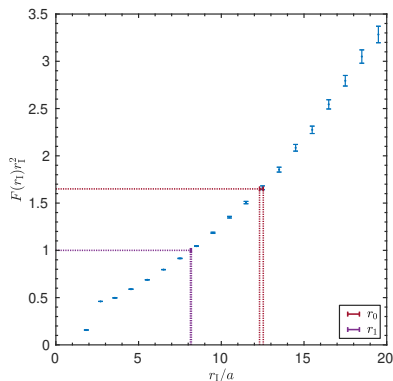
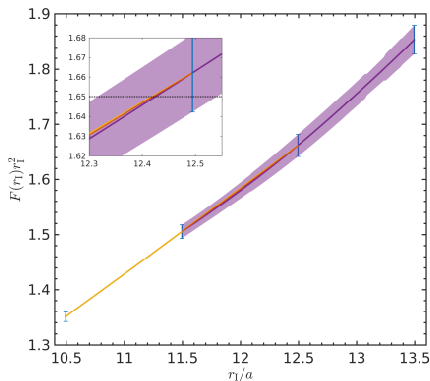
$$r_I^2 F(r_I)|_{r_I=r_1} = 1$$

[C. Bernard et al, (2000)]



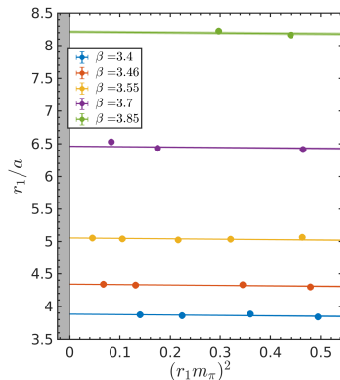
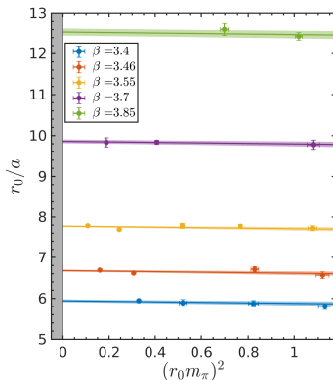
Finding r_0 and r_1

Zooming in at the r_0 interpolation: 2-points (red), 3-points early (yellow), and 3-point late (purple)



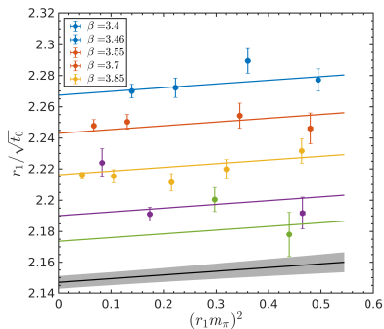
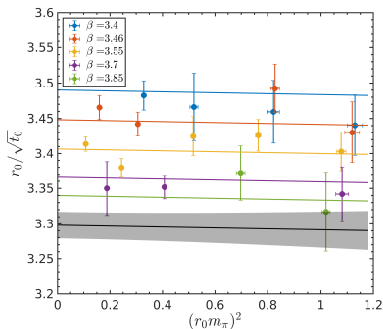
Chiral extrapolations in lattice units

Global fit using: $\frac{r_i}{a} = c_1|\beta + c_2(r_i m_\pi)^2$



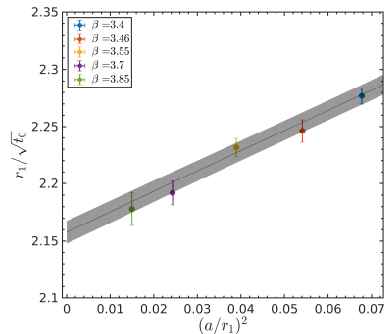
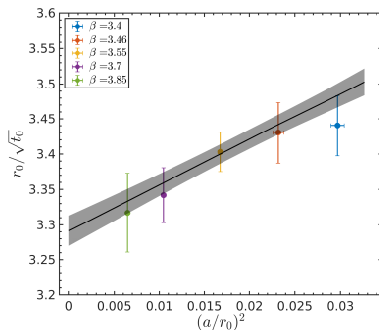
Continuum & chiral extrapolation using a global fit

Global fit using: $\frac{r_i}{\sqrt{t_0}} = c_1 + c_2 \left(\frac{a}{r_{0,sym}} \right)^2 + c_3 (r_0 m_\pi)^2$



At the symmetric point

Using the parameters from previous fit results



Different Fits

$$\frac{r_0}{\sqrt{t_0}} = F_{fit}(a, m_u, m_d, m_s)$$

$$F_1 : c_1 + c_2 \left(\frac{a}{r_{0,sym}} \right)^2 + c_3 (r_0 m_\pi)^2$$

$$F_2 : c_1 + c_2 \left(\frac{a}{r_{0,sym}} \right)^2 + c_3 (r_0 m_\pi)^2 + c_4 (m_\pi a)^2$$

$$F_3 : c_1 + c_2 \left(\frac{a}{r_{0,sym}} \right)^2 + c_3 \phi_2 + c_4 (1.098 - \phi_4)$$

$$F_4 : c_1 + c_2 \left(\frac{a}{r_{0,sym}} \right)^2 + c_3 (t_0 m_\pi^2) + c_4 (1.098 - \phi_4),$$

$(1.098 - \phi_4)$ takes mistuning into account

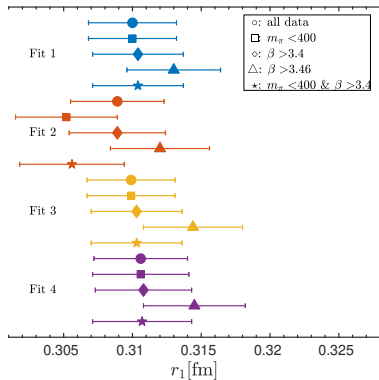
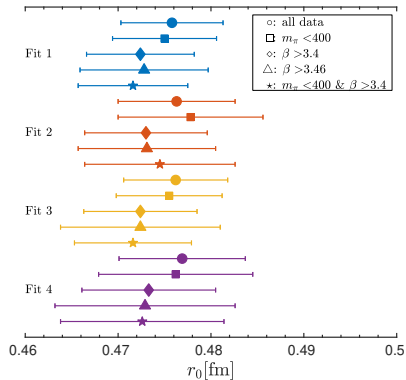
[B. Straßberger et al, (2021) [hep-lat/2112.06696](#)]

$$\phi_2 \sim m_u + m_d, \quad \phi_4 \sim m_u + m_d + m_s$$

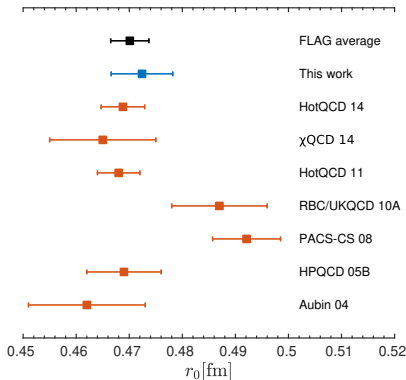


Comparison between the fits with datacuts

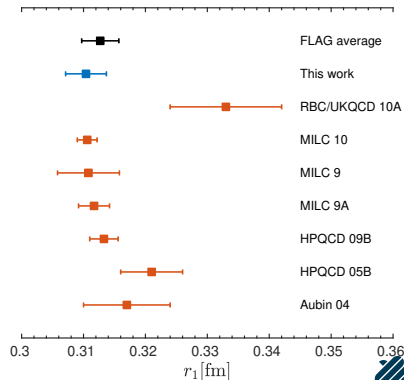
A cut in the symmetrical mass (\square), 1-2 coarse lattice spacings (\diamond, \triangle), and a combination (\star)



The physical value

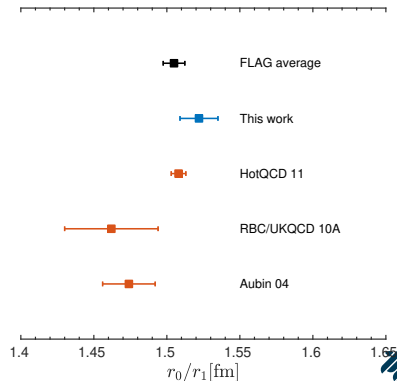
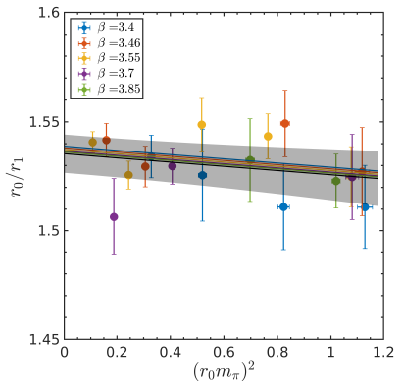


$$r_0^{\text{Phys}} = 0.4724(58) [\text{fm}]$$



$$r_1^{\text{Phys}} = 0.3104(33) [\text{fm}]$$



r_0/r_1


$$\frac{r_0}{r_1}^{\text{Phys}} = 1.522(13)[\text{fm}]$$



Structure of the static potential

Through the motivation of the Cornell potential:

$$V = \sigma r - \frac{K}{r}$$

$$K = 0.52 \quad [\text{E. Eichten, K. Gottfried, T.}$$

Kinoshita, K.D. Lane and T.M. Yan,

Phys. Rev. D21 (1980) 203]

Building following quantity

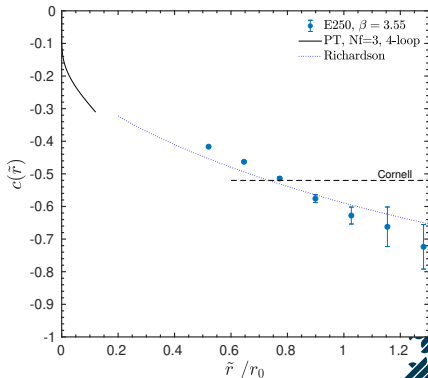
$$c(\tilde{r}) = \frac{1}{2} \tilde{r}^3 F'(r)$$

with $\tilde{r} = r + \mathcal{O}(a^2)$ such that

$$c_{tree}(\tilde{r}) = -C_F \frac{g_0^2}{4\pi}$$

[M. Lüscher and P. Weisz, hep-lat/0207003]

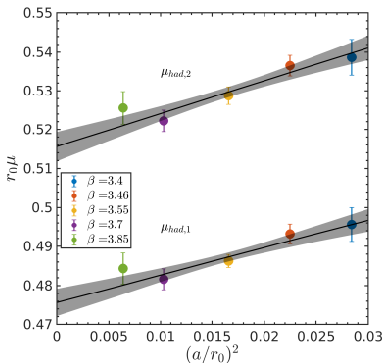
[J.L. Richardson, Phys. Lett. B82 (1979) 272]



Lambda parameter

Nonperturbative determination of $\beta(\bar{g})$ between μ_{had} and μ_{PT}

[M.Bruno et al, (2017) 1706.03821]



- ▶ $\mu_{had,1}$: scale where $\bar{g}_{GF}^2 = 11.31$
- ▶ $\mu_{had,2}$: scale where $\bar{g}_{GF}^2 = 10.20$

- ▶ $\frac{\Lambda_{\overline{MS}}^{(3)}}{\mu_{had,1}} = 1.729(57)$

- ▶ $\frac{\Lambda_{\overline{MS}}^{(3)}}{\mu_{had,2}} = 1.593(53)$

We combine our $\frac{r_0}{a}$ with $a\mu_{had}$

- ▶ $r_0 \Lambda_{\overline{MS}}^{(3)} = 0.820(28)$

- ▶ Flag: $r_0 \Lambda_{\overline{MS}}^{(3)} = 0.808(29)$



Conclusion

- The results were shown to be stable under alterations of different cuts, fits, as well as the determination of the effective masses
- Detailed analysis with masses down to the pion mass and lattice spacing down to 0.037 fm giving:
 - ▶ $r_0 = 0.4724(58)\text{fm}$
 - ▶ $r_0 = 0.3104(33)\text{fm}$
 - ▶ $\frac{r_0}{r_1} = 1.537(5)$
 - ▶ $r_0 \Lambda_{\overline{\text{MS}}}^{(3)} = 0.820(28)$



That's it, questions?



The difficulties of E300

